

Space-time evolution of hadronization probed in e+A collisions

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31 January 2007

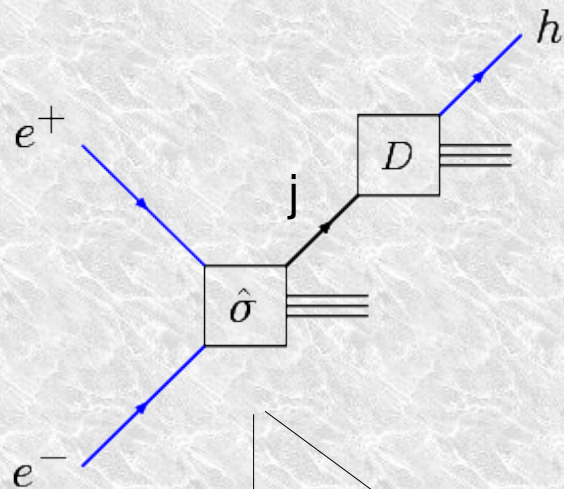
Based on: A.A., PLB in press [nucl-th/0604041]
A.A., EPJC in press [nucl-th/0609010]

Outline

- **Hadronization in nuclear matter**
 - space-time evolution, non-perturbative – confinement
 - calibration of jet-quenching signal in A+A
- **Short review of experimental data**
- **Hadron quenching in nDIS**
 - formation times
 - prehadron vs. hadron
 - energy loss vs. absorption
- **Can we distinguish energy loss from hadron absorption?**
 - failure of the " $A^{2/3}$ power law"
 - formation time scaling
 - p_T -broadening
- **Perspectives and conclusions**

Hadronization in nuclear matter

Hadronization in elementary collisions



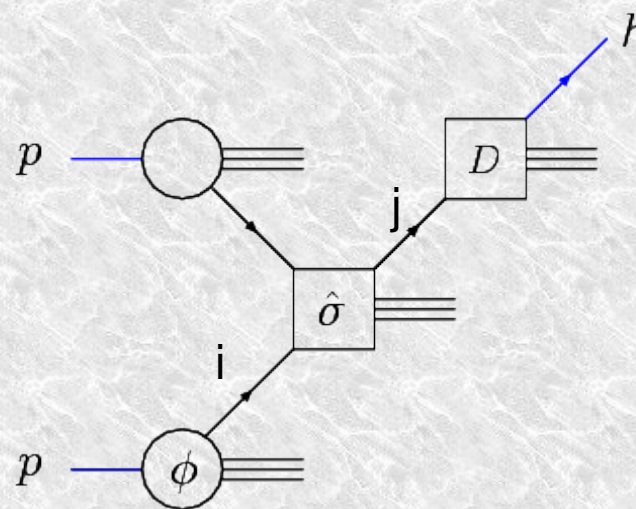
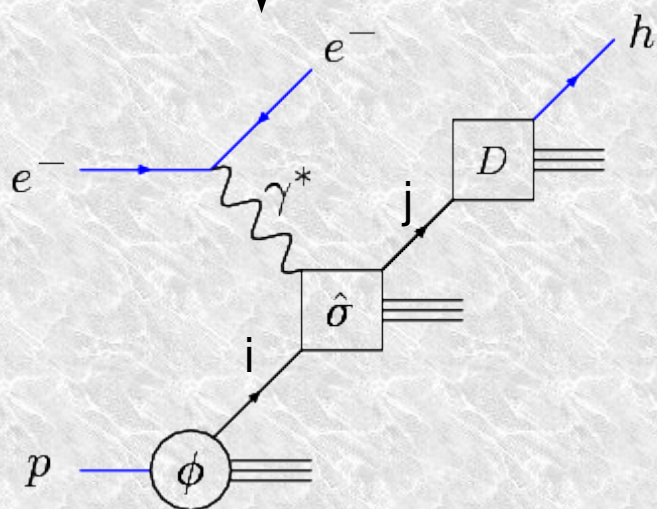
- ◆ perturbative QCD factorization of short and long distance physics

$$d\sigma_{\text{hadronic}} = \prod_i \phi_i \otimes \hat{\sigma}_{\text{partonic}}^{ij} \otimes D_{j|h}$$

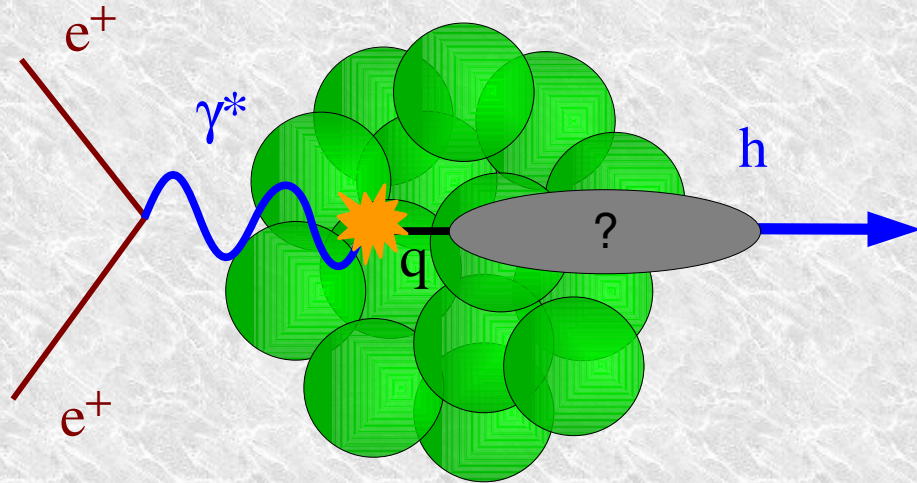
Parton Distribution Fns
(from inclusive DIS)

Fragmentation Fns
(from $e^+e^- \rightarrow h+X$)

- ◆ **Universality:** Fragm. Fns. from $e^+e^- \rightarrow h+X$ describe hadronization in DIS and $p+p \rightarrow h+X$

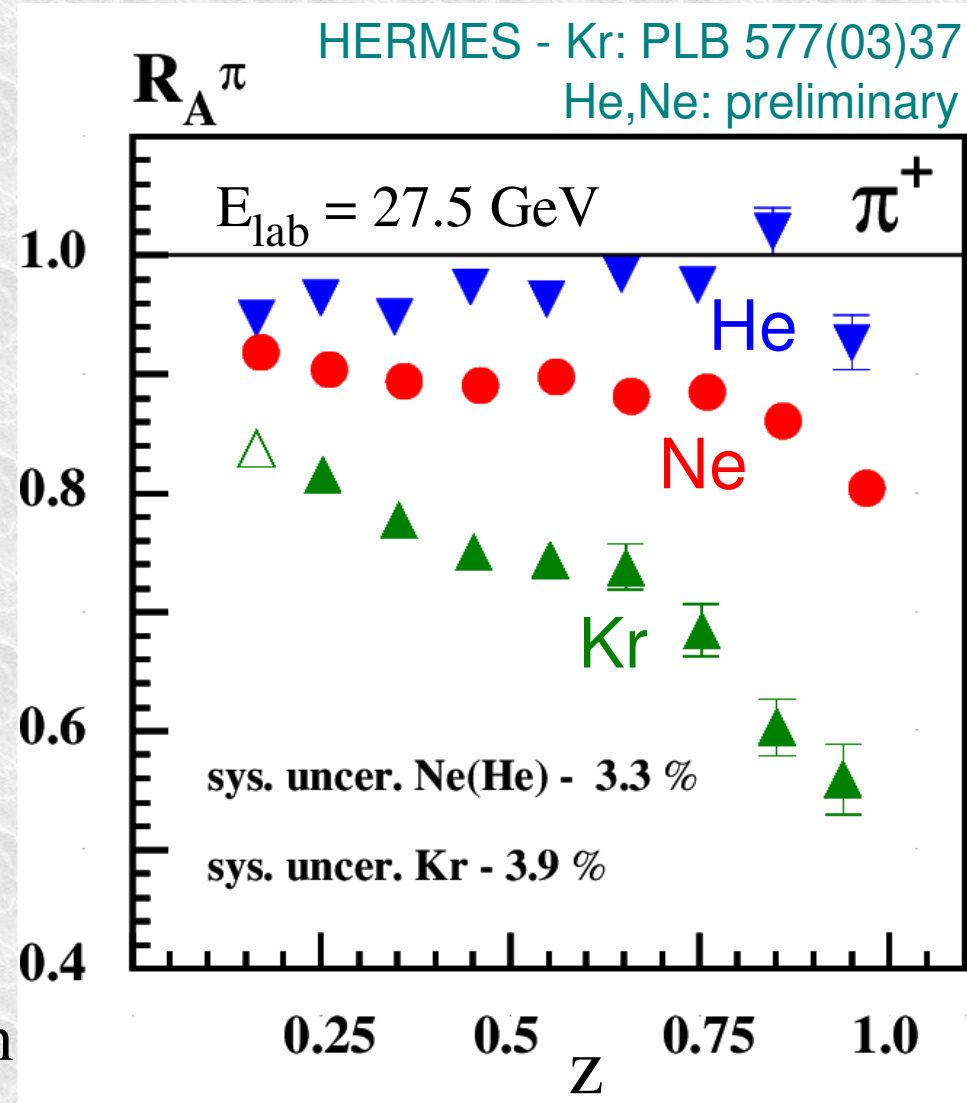


Nuclear collisions 1 - nDIS



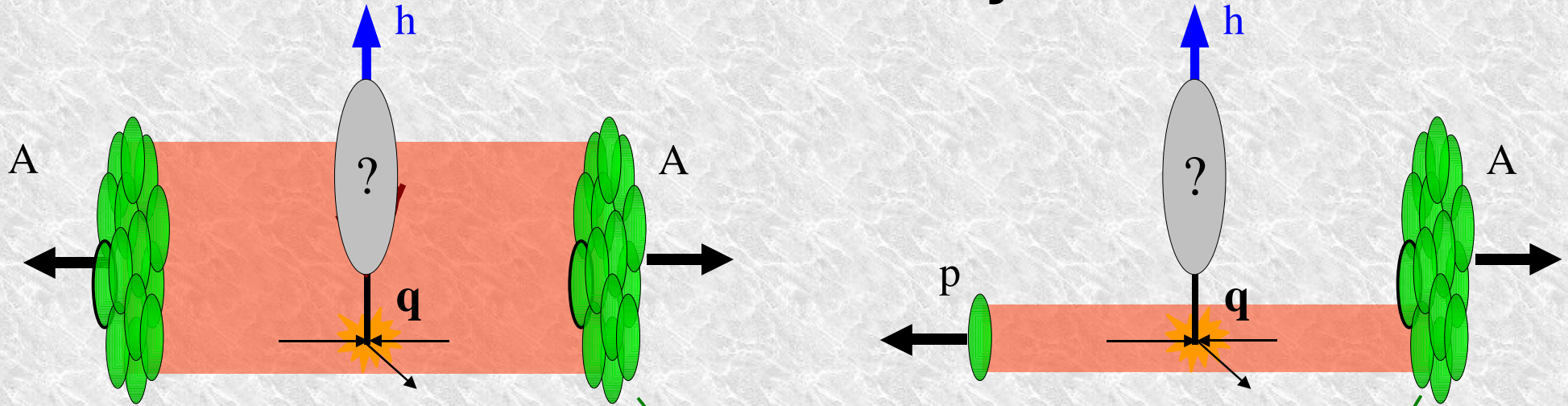
$$R_M^h(z) = \frac{\frac{1}{N_A^{DIS}} \frac{dN_A^h(z)}{dz}}{\frac{1}{N_D^{DIS}} \frac{dN_D^h(z)}{dz}}$$

- Nuclear effects on PDF cancel in ratios
- Exposes modifications of hadronization



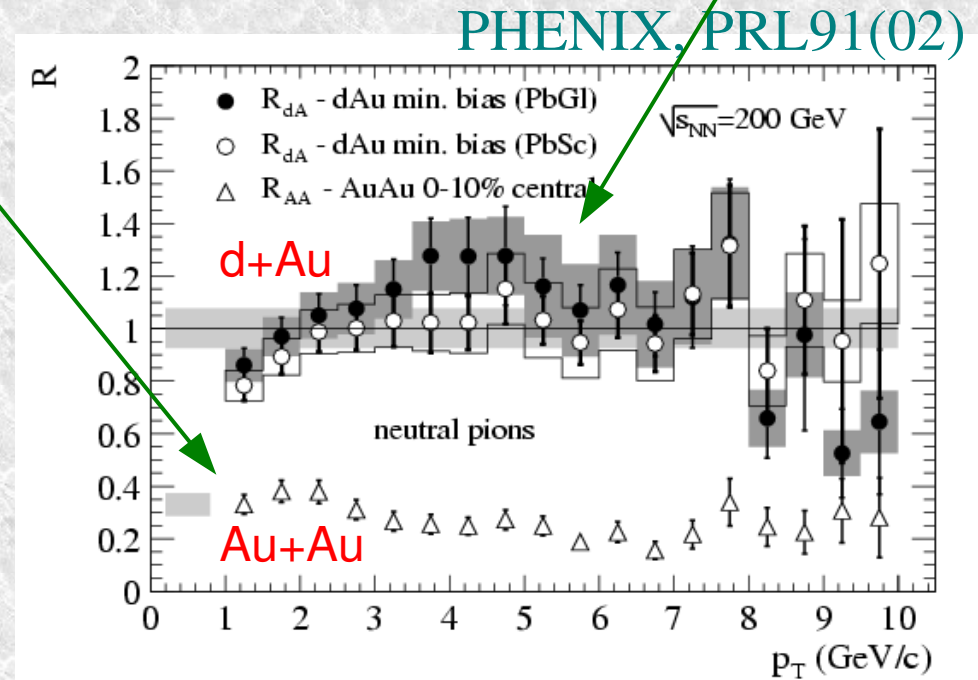
$R_M < 1 \Rightarrow$ hadron attenuation in cold nuclear matter

Nuclear collisions 2 – Heavy ion collisions



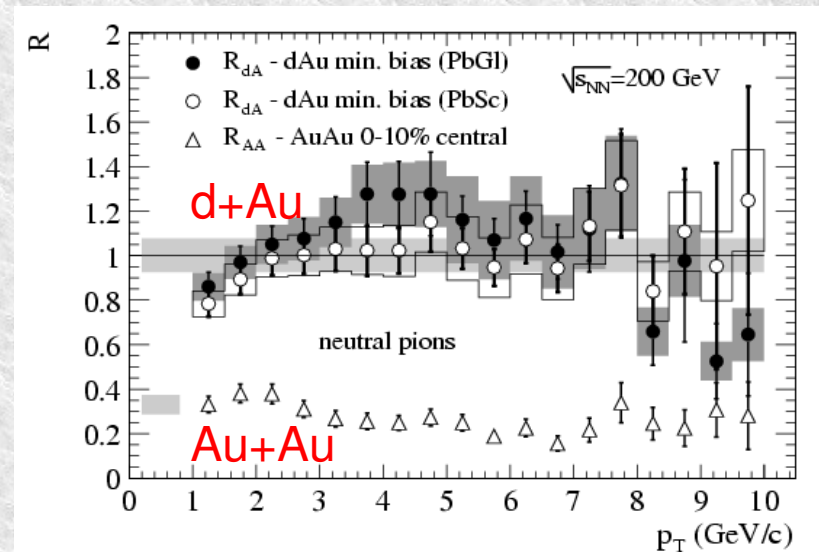
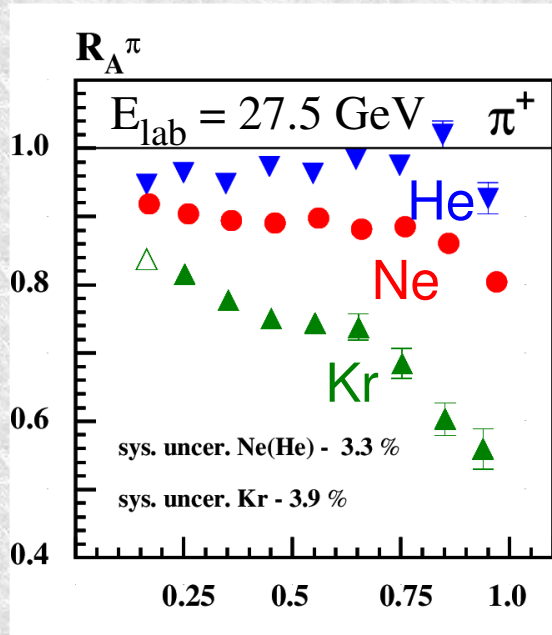
$$R_{AB} = \frac{(dN^h/d^2p_T)_{A+B}}{T_{BA}(b) (d\sigma^h/d^2p_T)_{p+p}}$$

➤ Medium modifications of hadronization isolated by comparison of h+A and A+A



$R_{AuAu} < 1$ & $R_{dAu} > 1 \Rightarrow$ hadron attenuation in hot nuclear matter

Breakdown of universality in nuclei



➔ Hadronization is no more process-independent

➔ Among possible causes:

➔ struck quark interactions with the medium

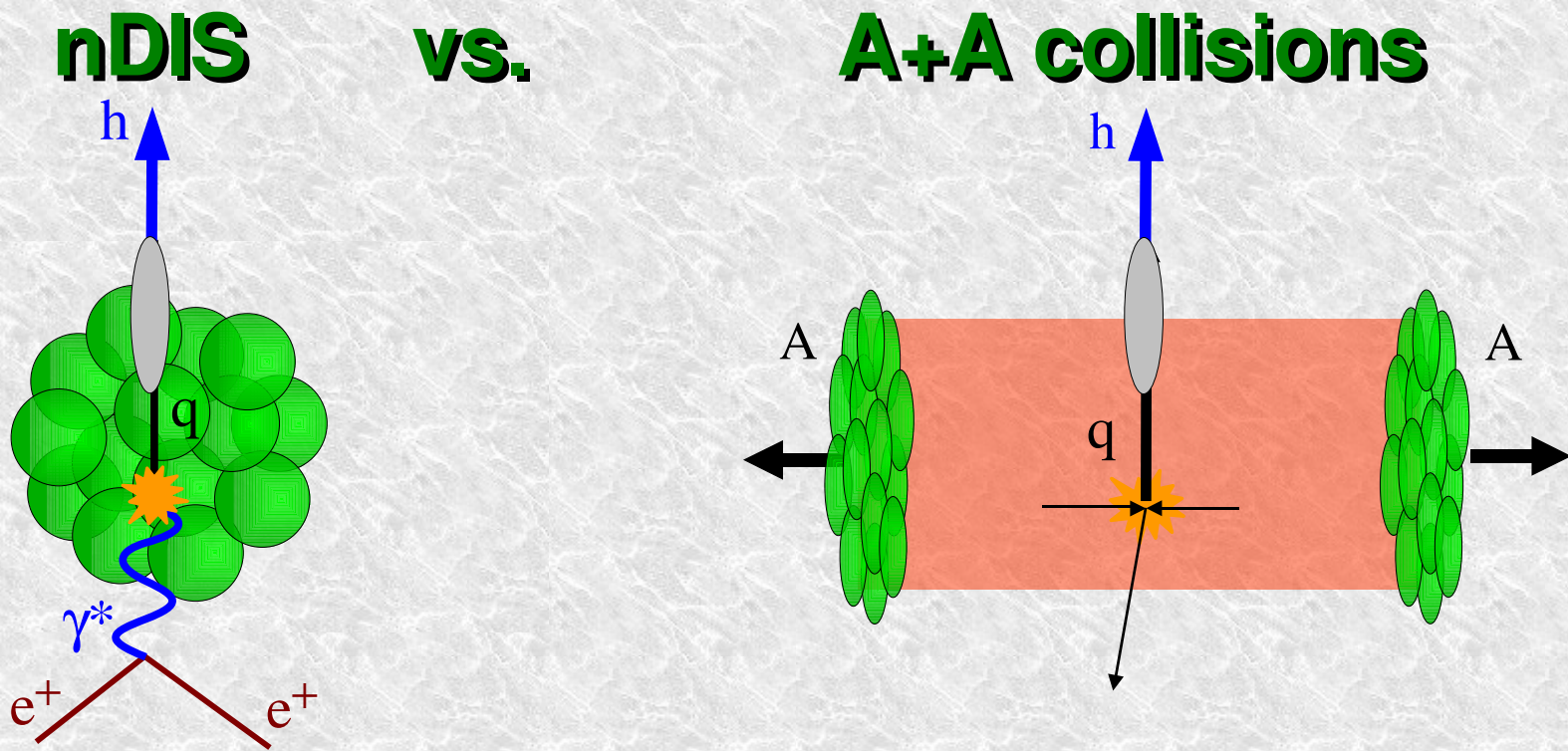
➔ (pre)hadron interactions with the medium

➔ other medium nuclear, e.g., partial deconfinement [Dias de Deus '87]

➔ in-medium modifications of parton showers [Borghini, Wiedemann '05]

➔ breakdown of factorization [for nuclear PDF, see Qiu, Sterman '02]

This talk:
 space-time evolution
 of hadronization



- nDIS is a clean environment for
 - (1) **space-time evolution of hadronization**
 - nucleons as micro-detectors
 - medium rather well known
 - (2) **Cold nuclear matter effects**
 - quark energy loss
 - nuclear modifications of FF

Jet-quenching in A+A

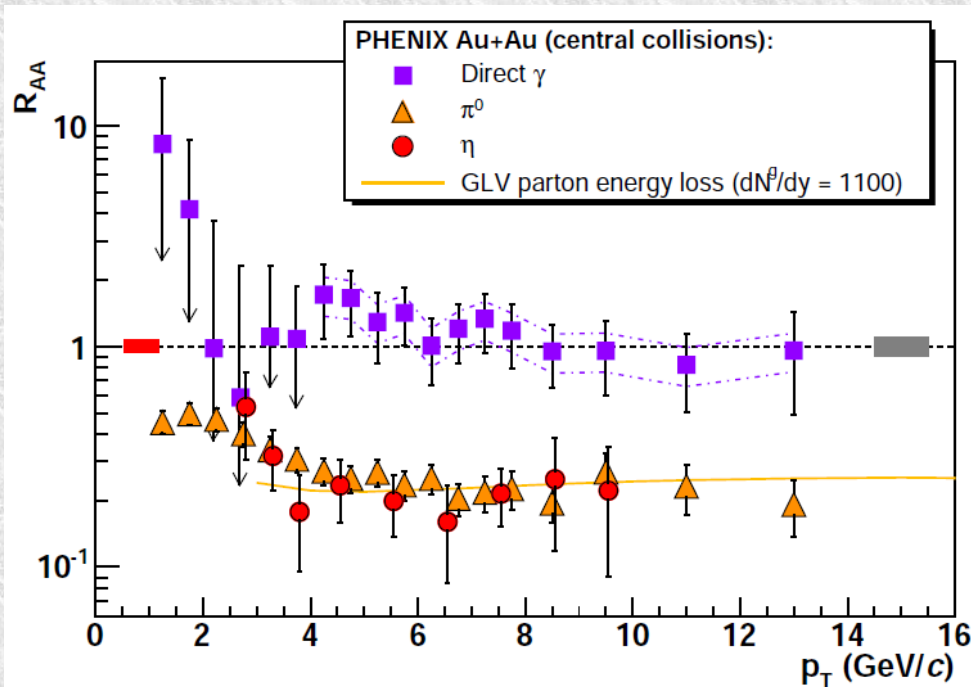
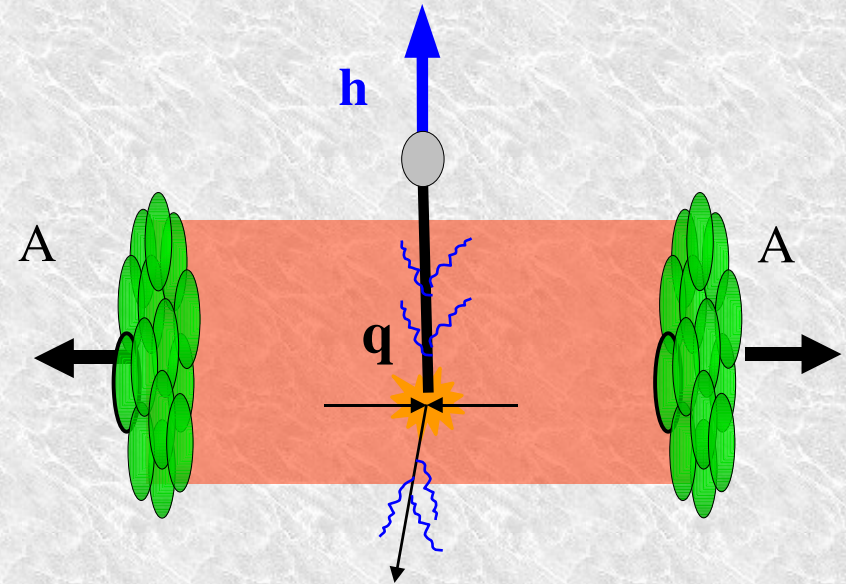


properties of hot nuclear matter

Energy loss paradigm: success!

- Paradigm: long lived partons
hadron quenching due to energy lost
by gluon brehmstrahlung

$$R_{AA} = \frac{(dN^h/d^2p_T)_{A+A}}{N_{coll}(b) (dN^h/d^2p_T)_{p+p}}$$



- pQCD energy-loss theory used to extract initial gluon density
 $dN^g/dy \approx 1000$

- Can be converted to initial state energy density:
 $\epsilon \sim 14 - 20 \text{ GeV}/\text{fm}^3 > \epsilon_c$
allowing QGP phase transition

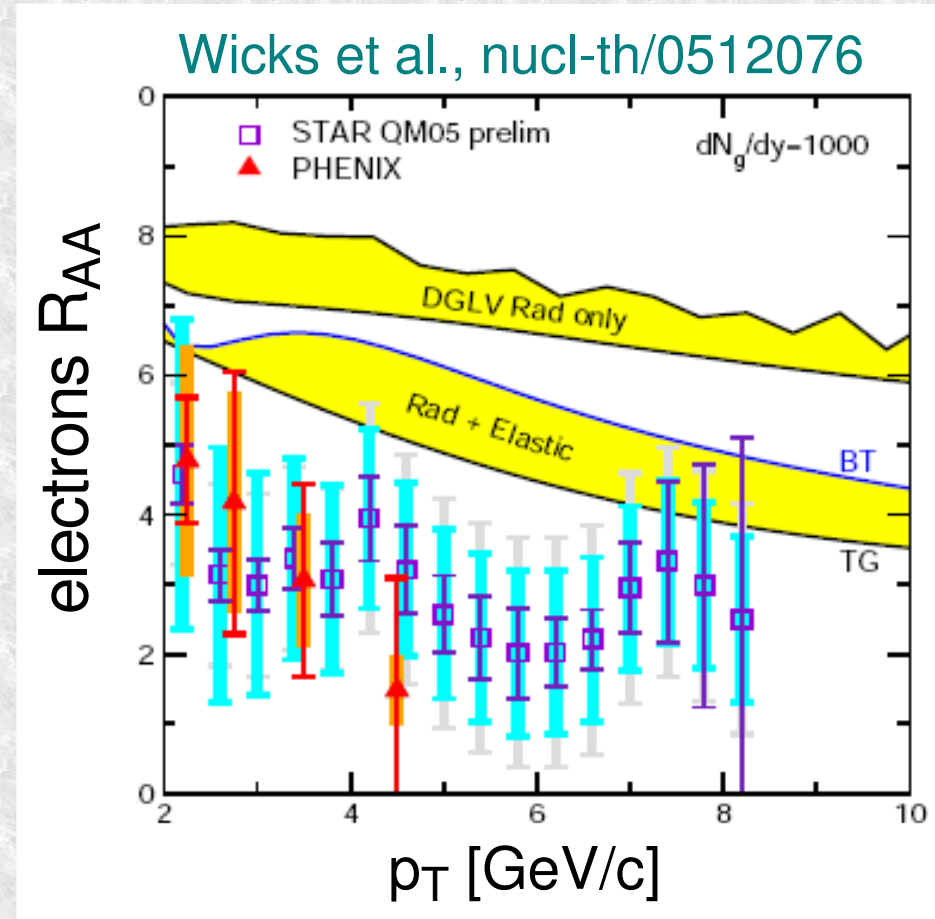
Challenges: heavy quarks

- ◆ Heavy flavour puzzle at RHIC
[QM05: Djordjevic, Armesto, STAR, PHENIX]
 - single non-photonic e^- as much suppressed as π
 - *theory gives half of the observed suppression!*

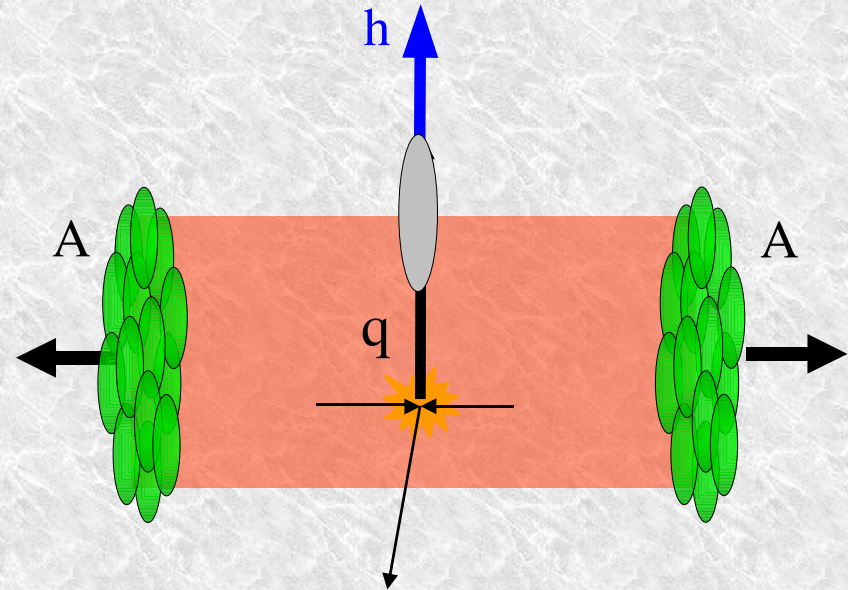
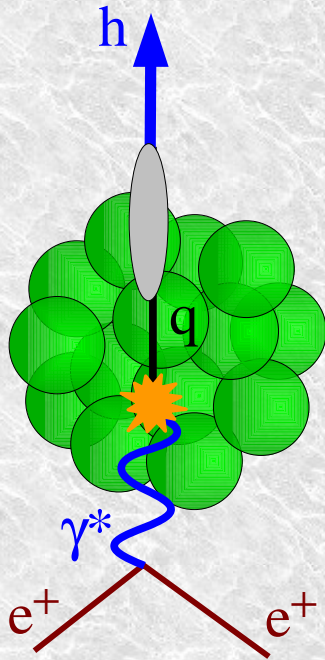
- ◆ Additional mechanisms:
 - elastic energy loss?
[Mustafa, Thoma]
with running α_s ? [Peshier '06]
 - Uncertainties in NLO cross sect.
(B/D ratio vs. p_T)?

This talk:

- **do hadrons form inside QGP?**



Similarities and differences



$$E_q = \nu = E_e - E_{e'} \approx 2-25 \text{ GeV}$$

at HERMES/Jlab

$$E_h = z \nu \approx \mathbf{2 - 20 \text{ GeV}}$$

$$E_q = p_T / z$$

$$E_h = p_T \approx \mathbf{2 - 20 \text{ GeV}}$$

★ HERMES/JLAB kinematics is relevant to RHIC mid-rapidity

...but beware the virtuality...

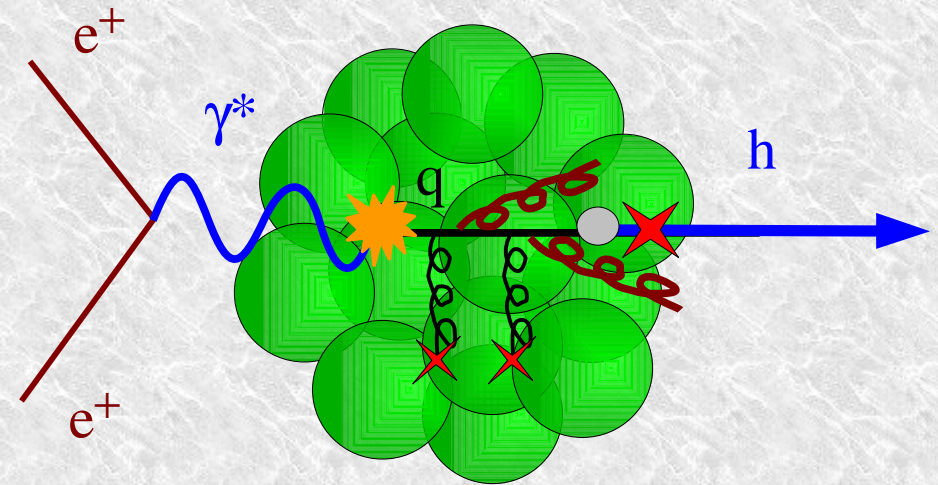
$$Q^2 = -q^2 \text{ is measured}$$

$$Q^2 \equiv E_q^2 \propto (p_T/z)^2$$

Summary of motivations

➤ Nuclei as space-time analyzers

- nucleons as femto-detectors
- medium rather well known
- low final-state multiplicity



➤ Non perturbative aspects of hadronization

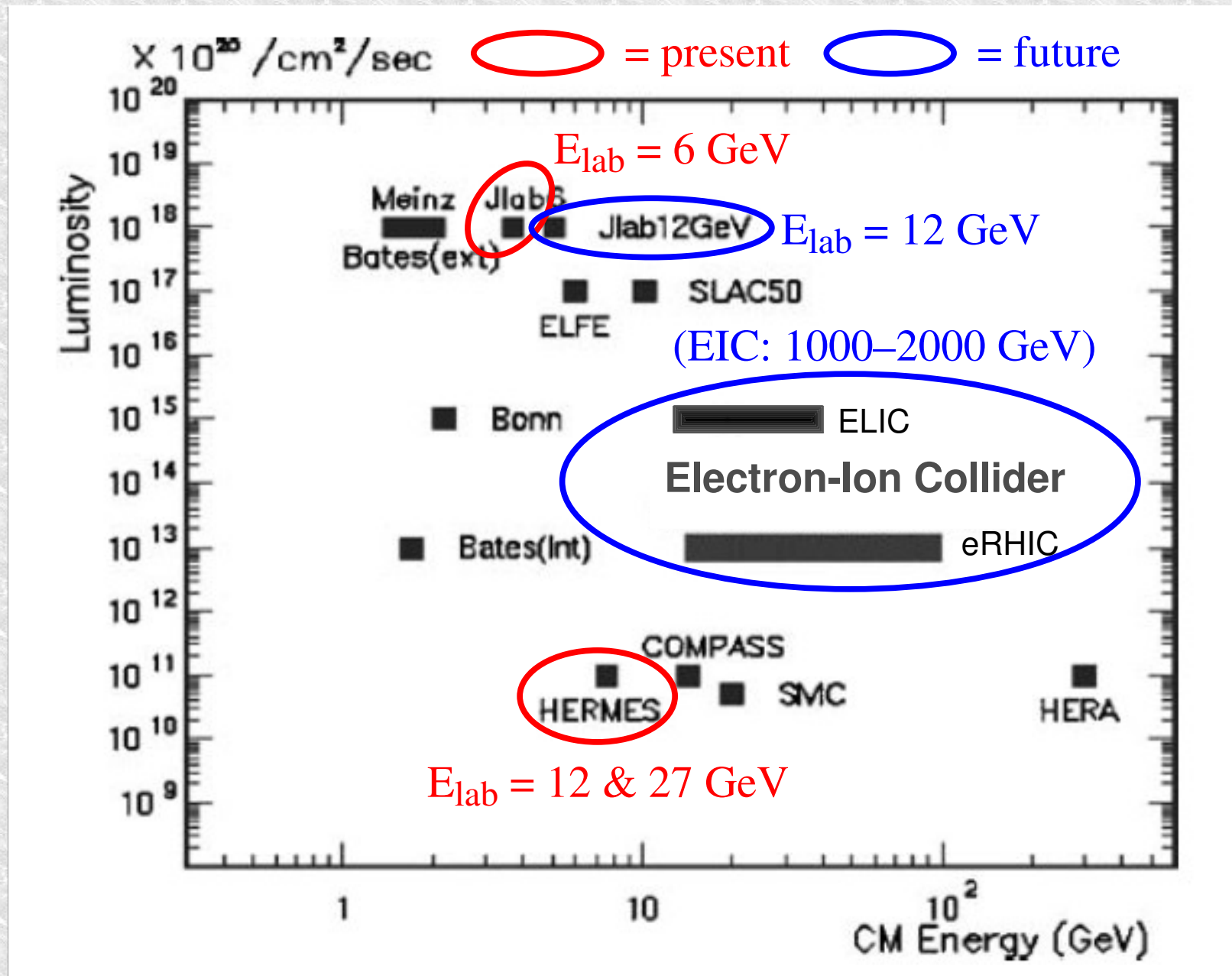
- approaching microscopic understanding of Fragmentation Functions
- how do partons dress up?
- understanding of color confinement

➤ Calibration of jet-quenching signal in A+A collisions

- properties of Quark-Gluon Plasma

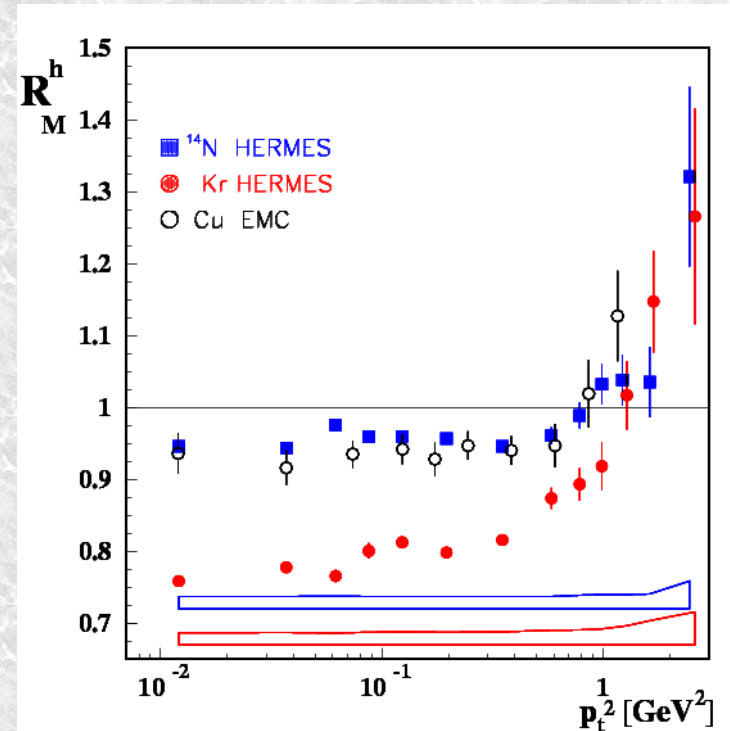
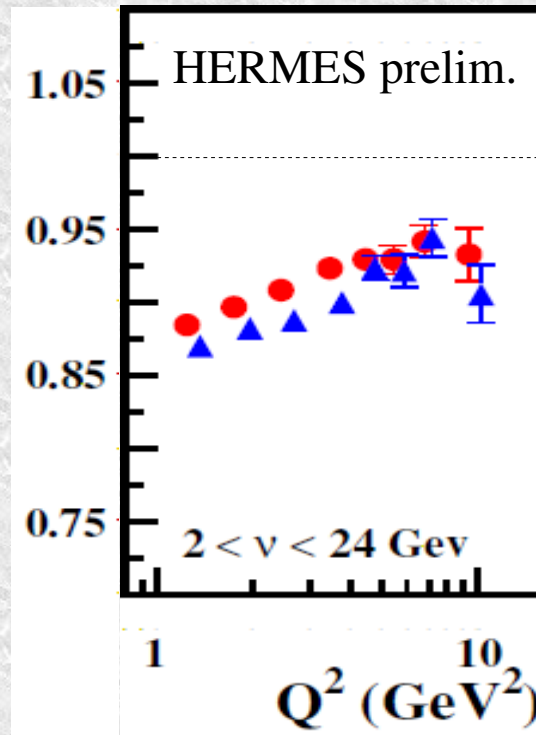
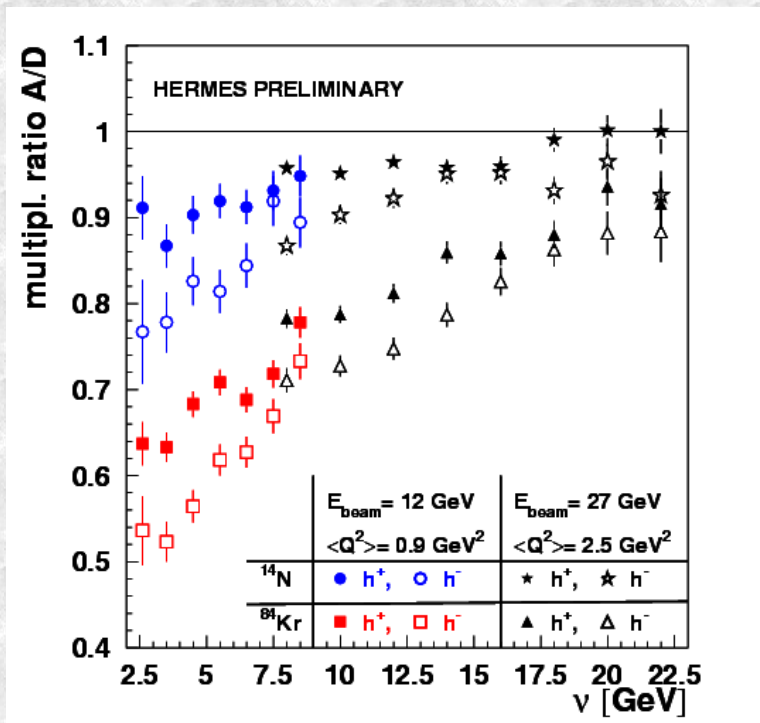
Short review of e+A data

Present and future facilities



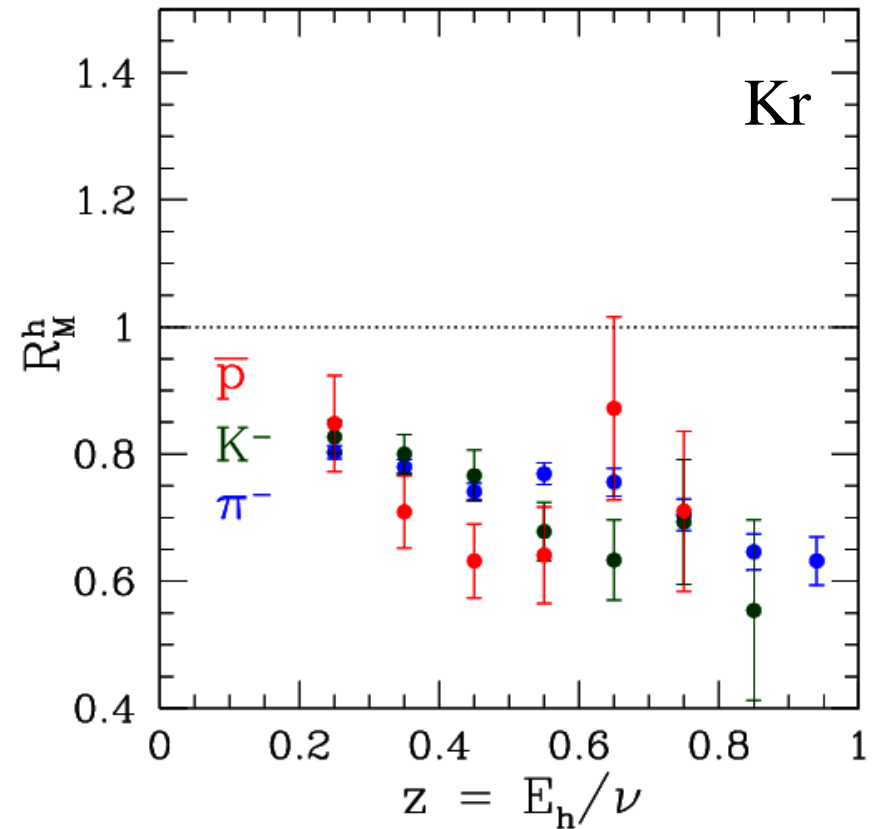
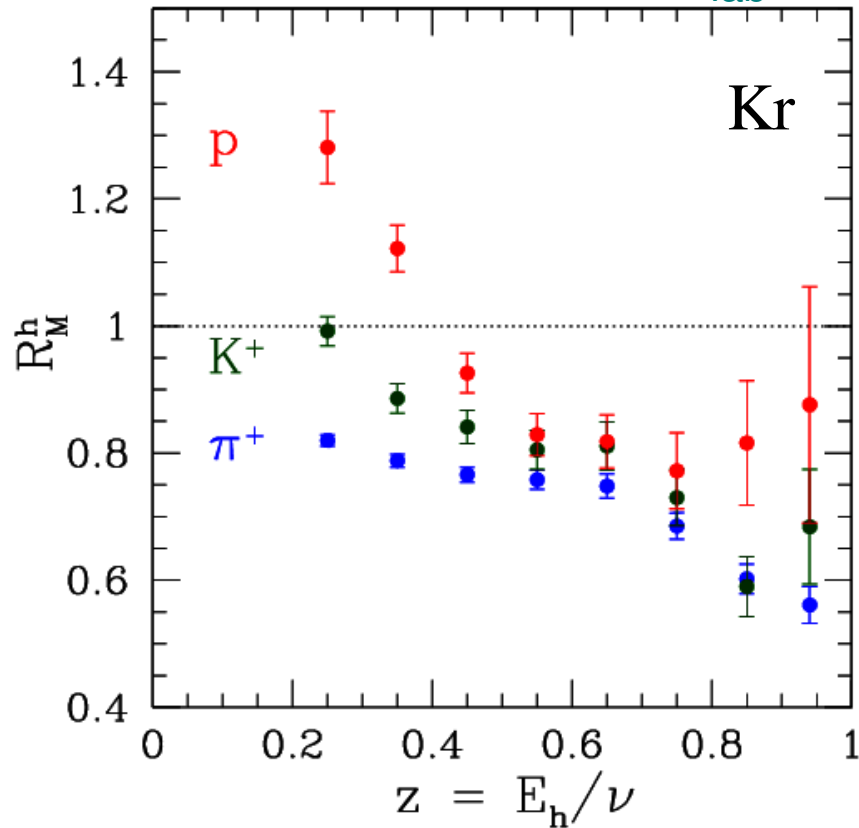
Measurements at HERMES

- ◆ HERMES: fixed target, $E_{\text{lab}} = 27.5 \text{ GeV}$ and 12 GeV
- ◆ Hadron attenuation as a function of (fixed target kinematics)
 - ➔ $\nu =$ virtual γ energy
 - ➔ $z = E_h/\nu =$ hadron's fractional energy
 - ➔ $Q^2 =$ photon virtuality
 - ➔ $p_T =$ hadron transv. momentum
 - ➔ hadron flavour = $\pi^\pm, K^\pm, p, \bar{p}$
 - ➔ $A =$ target mass number



Measurements at HERMES

HERMES $E_{\text{lab}}=27 \text{ GeV}$ - Phys.Lett.B577(03)37



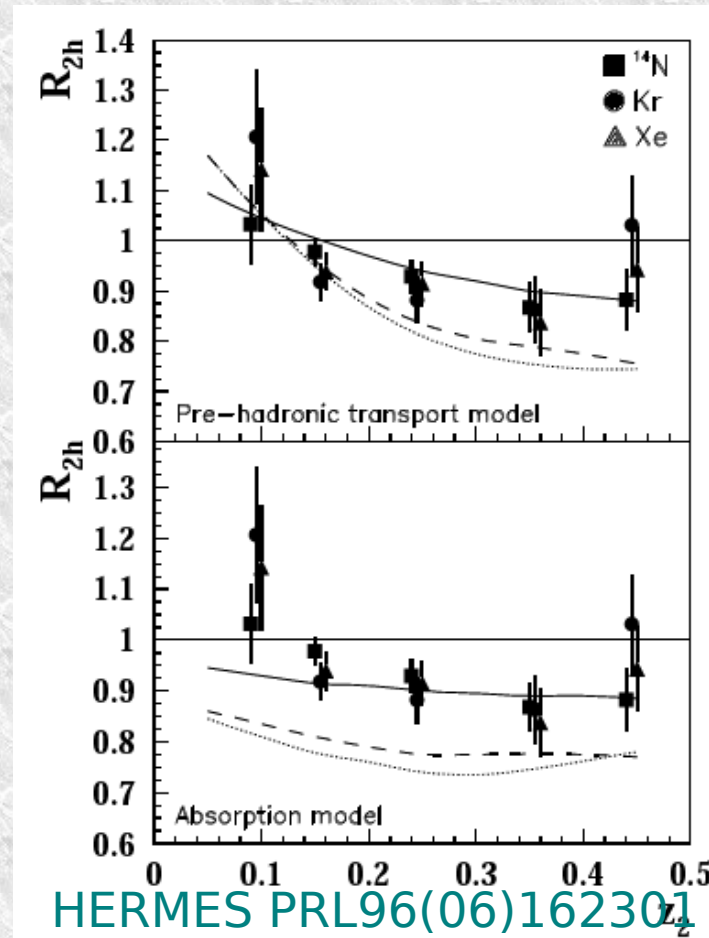
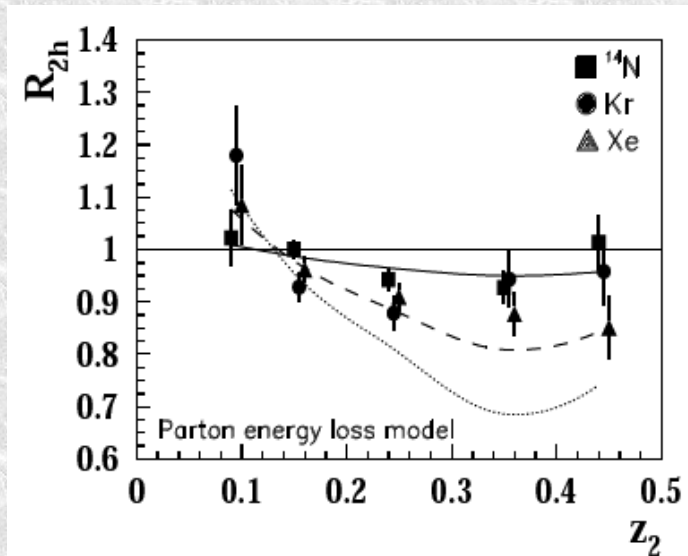
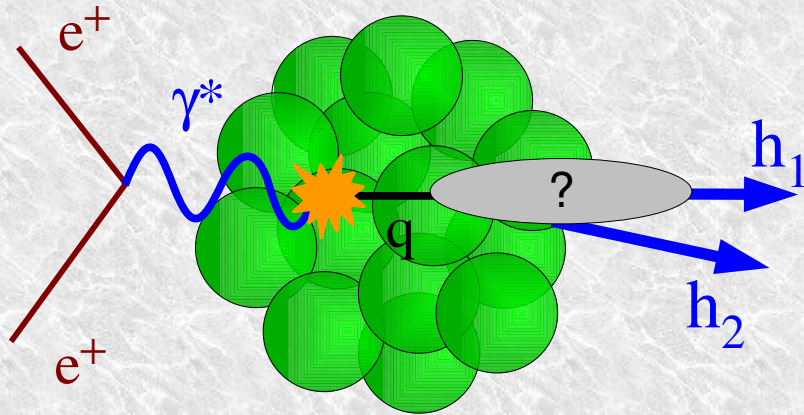
- ◆ **proton anomaly!**
- ◆ analogous to “baryon/meson anomaly” in $p+p$, $p+A$ and $A+A$
- ➡ what do they have in common?

Measurements at HERMES

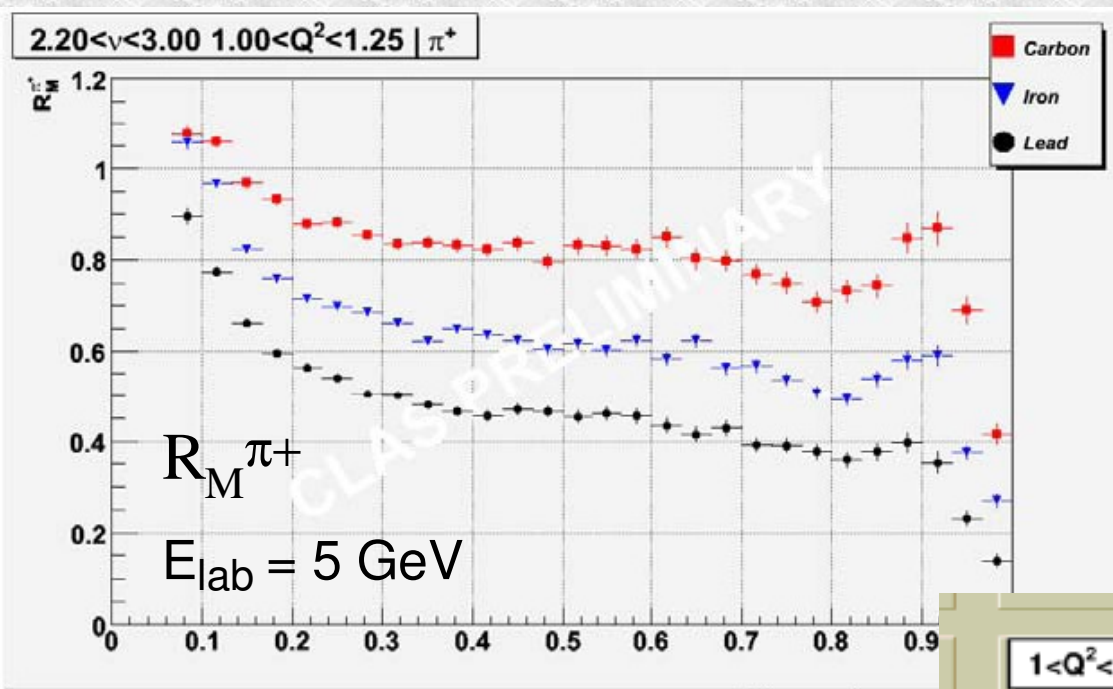
- ◆ Double hadron attenuation R_2
- ◆ in A+A: “same-side correlations”

$$R_2(z_2) = \frac{\left. \frac{N_2(z_2)}{N_1} \right|_A}{\left. \frac{N_2(z_2)}{N_1} \right|_D}$$

$$z_2 < z_1 \quad z_1 > 0.5$$



Preliminary results at CLAS

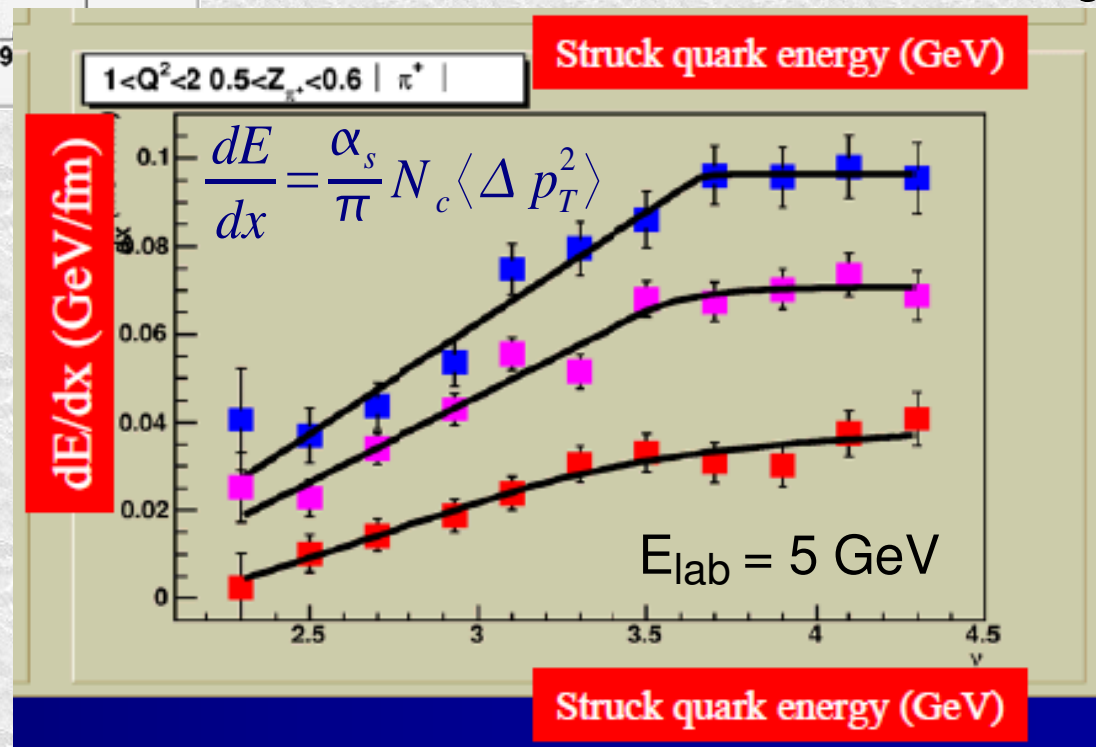
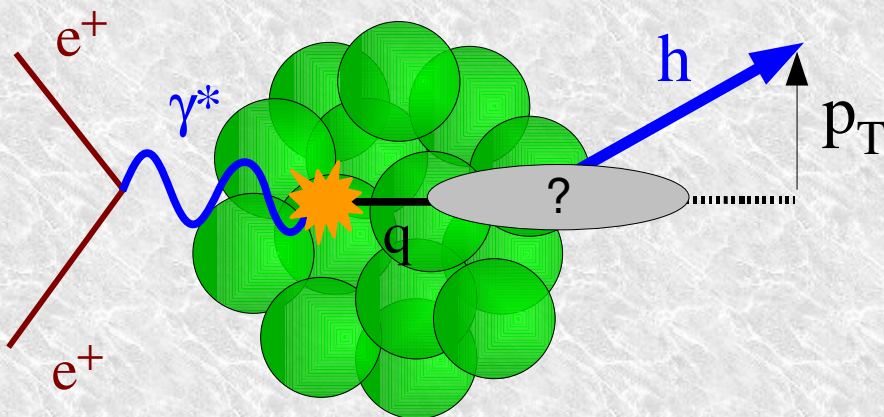


CLAS: fixed target
 $E_{lab} = 5 \text{ GeV}$

Will Brooks, 13 June 2006
JLAB user's group workshop

Transverse momentum broadening

K.Hafidi, nucl-ex/0609005

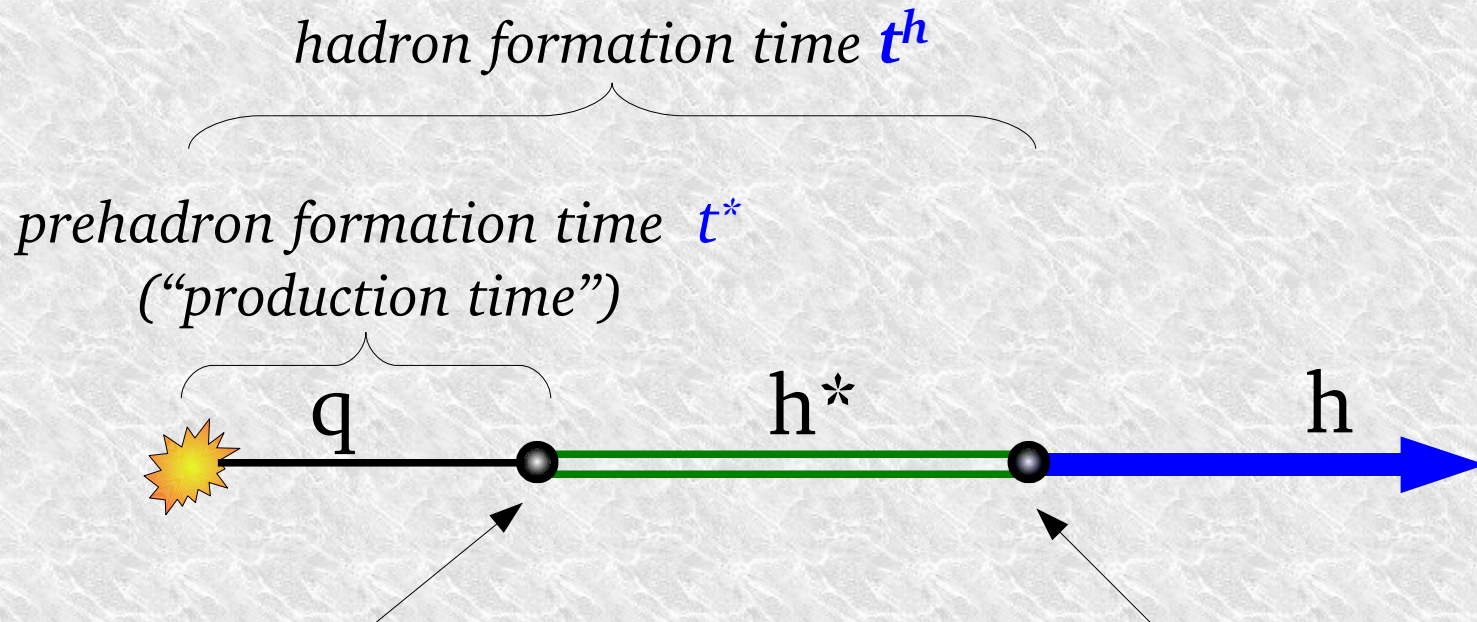


Hadron formation times

Prehadron vs. hadron

➡ Hadronization is non perturbative \Rightarrow (many) models

➡ General features:



Color neutralization \Rightarrow “prehadron” h^*
- gluon radiation stops
- large inelastic cross-section for h^*

prehadron collapses on
hadron's h wavefunction

➡ NOTE:

➡ It's tricky to rigorously define t^* , t^h : consider them as working tools

Hadron attenuation in nDIS

$$R_M^h(z) = \frac{\frac{1}{N_A^{DIS}} \frac{dN_A^h(z)}{dz}}{\frac{1}{N_D^{DIS}} \frac{dN_D^h(z)}{dz}}$$

Energy loss (gluon bremsstrahlung)

[Arleo; Wang *et al.*]

- hadronization outside the medium
- gluon radiation off struck quark

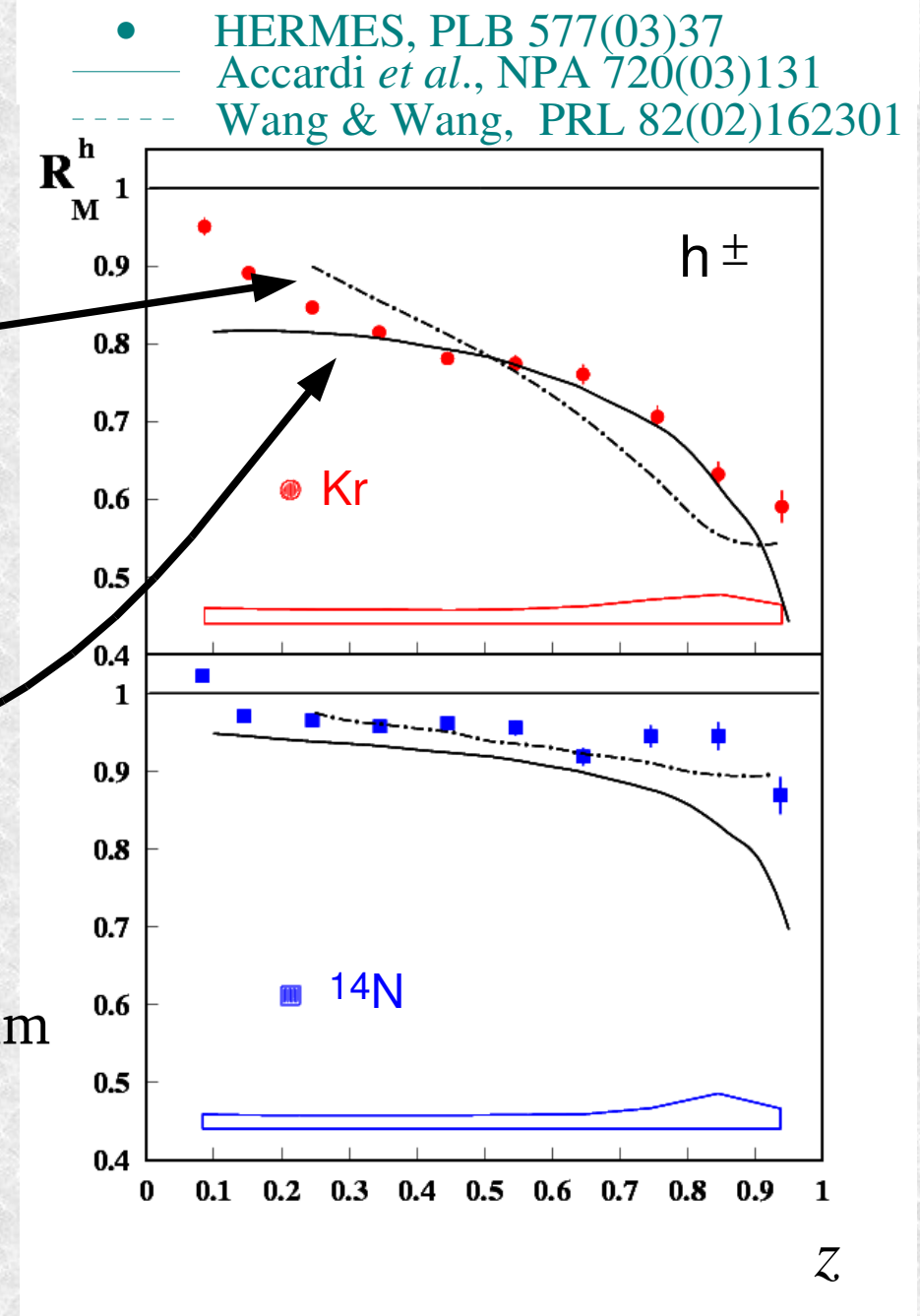
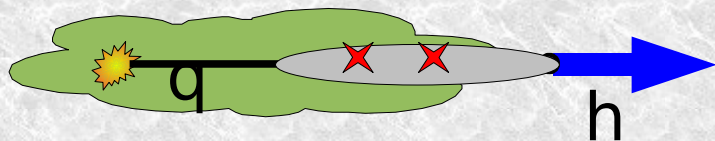


Prehadron absorption

[Accardi *et al.*;

Falter *et al.*; Kopeliovich, *et al.*]

- color neutralization inside the medium
- prehadron-nucleon scatterings



Formation time estimates 1

◆ Hadron formation time = time for partons to build up a color field and develop hadron wave function

◆ hadron's rest frame: $\tau^h \sim R_h$
 ◆ lab frame: $t^h \sim R_h \frac{E_h}{m_h}$ (in DIS $t^h \sim R_h \frac{zV}{m_h}$))

boost factor

◆ For a 7 GeV pion at HERMES $t^\pi \sim 35 \text{ fm} \gg R_A$

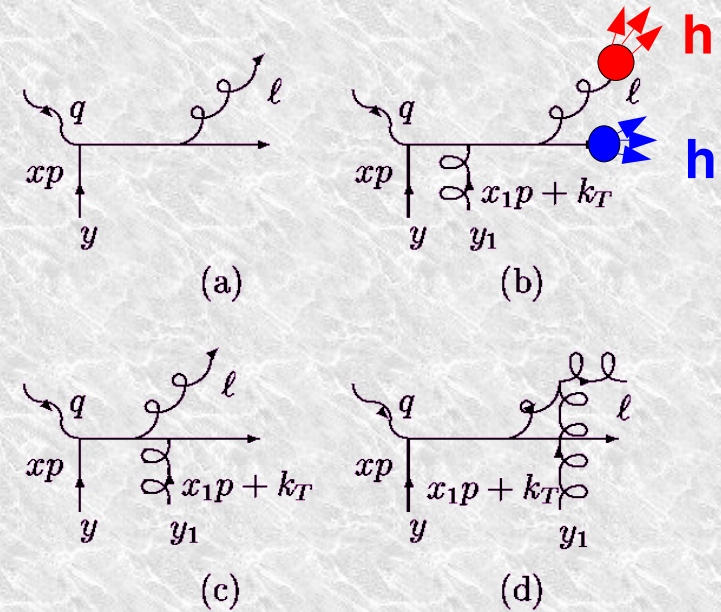
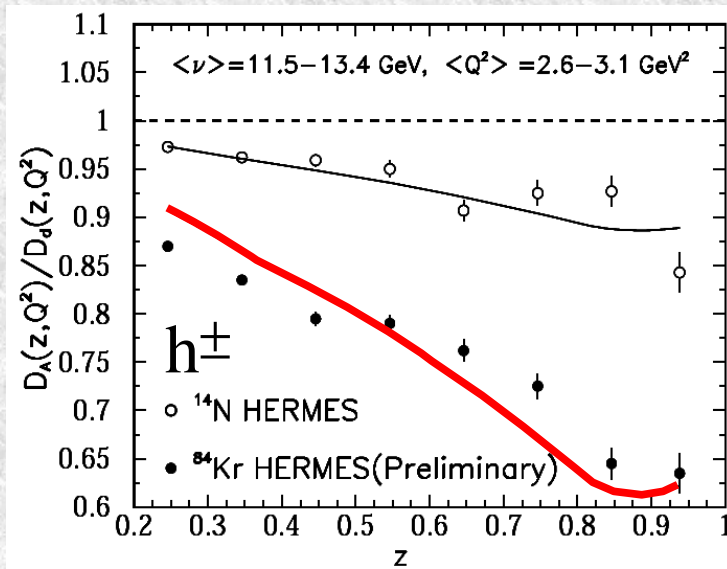
◆ pions are formed outside A, quark energy loss only

◆ Note, however: $t^K \sim t^P \sim 8 \text{ fm} !!$

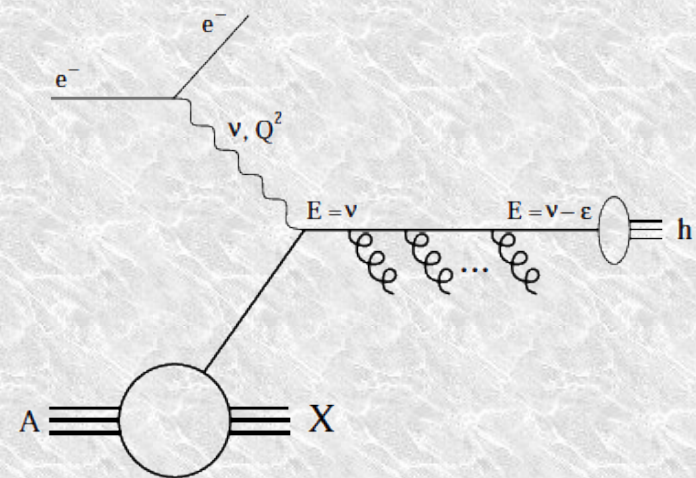
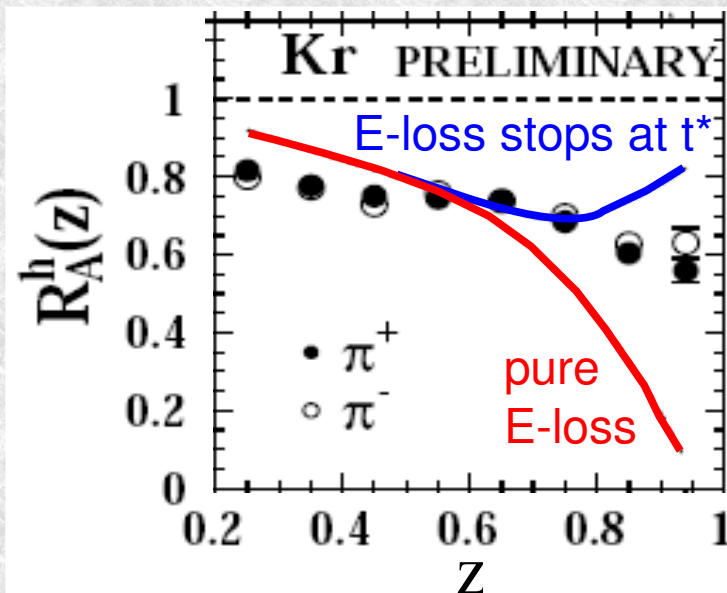
◆ This is used in energy loss models to justify the assumptions, but **neglects interactions of forming color field with the medium**

Formation time estimates 1 – energy loss models

➔ Twist-4 modified Fragmentation Fns. [Wang&Guo '00, Wang & Wang '02]



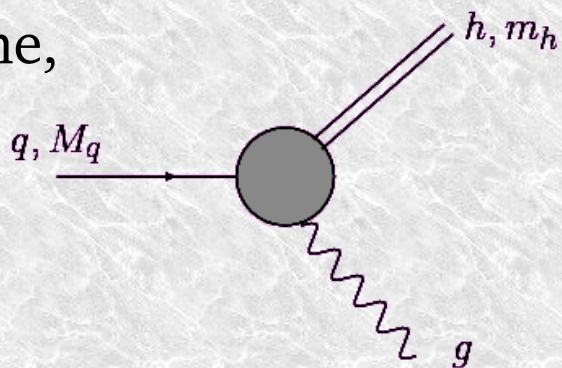
➔ Quark energy loss à la BDMPS [Arleo '02]



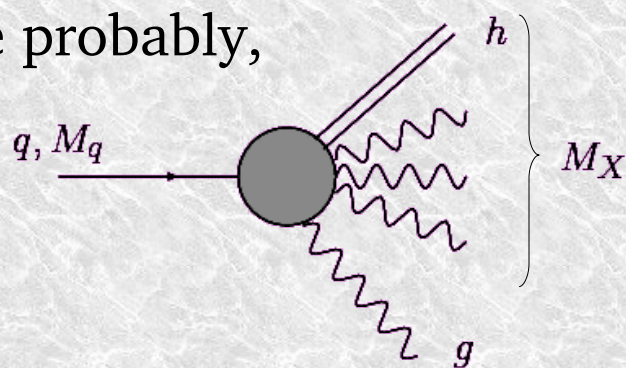
Formation time estimates 2 – pQCD estimate

➔ pQCD estimate [see Vitev, QM'05]

assume,



more probably,



$$\left[p^+, \frac{M_q^2}{2p^+}, 0 \right] \rightarrow \left[zp^+, \frac{\mathbf{k}^2 + m_h^2}{2zp^+}, \mathbf{k} \right] + \left[(1-z)p^+, \frac{\mathbf{k}^2}{2(1-z)p^+}, -\mathbf{k} \right]$$

$$\Delta y^+ \simeq \frac{1}{\Delta p^-} = \frac{2z(1-z)p^+}{\mathbf{k}^2 + (1-z)m_h^2 - z(1-z)M_q^2}$$

	π	K	p	D	B
HERMES ($v \sim 13$ GeV, $z \sim 0.5$)	37 fm	11 fm	4 fm	1.2 fm	0.1 fm
RHIC ($p_T^h \sim 7$ GeV, $z \sim 0.7$)	26 fm	6 fm	4 fm	1.2 fm	0.1 fm

~ Inside the medium !!

Formation time estimates 3 – Lund model

★ Prehadrons and hadrons are identified as follows [Bialas-Gyulassy '87]

➔ Prehadron formed at $q\bar{q}$ creation (string breaking) – C_i

➔ Hadron h_i formed when q and \bar{q} meet – P_i

★ Average formation times are analytically computable

➔ At large $z \rightarrow 1$

$E_h \rightarrow v \Rightarrow$ string breaks early to leave

all energy to the hadron: $\langle t^* \rangle \rightarrow 0$

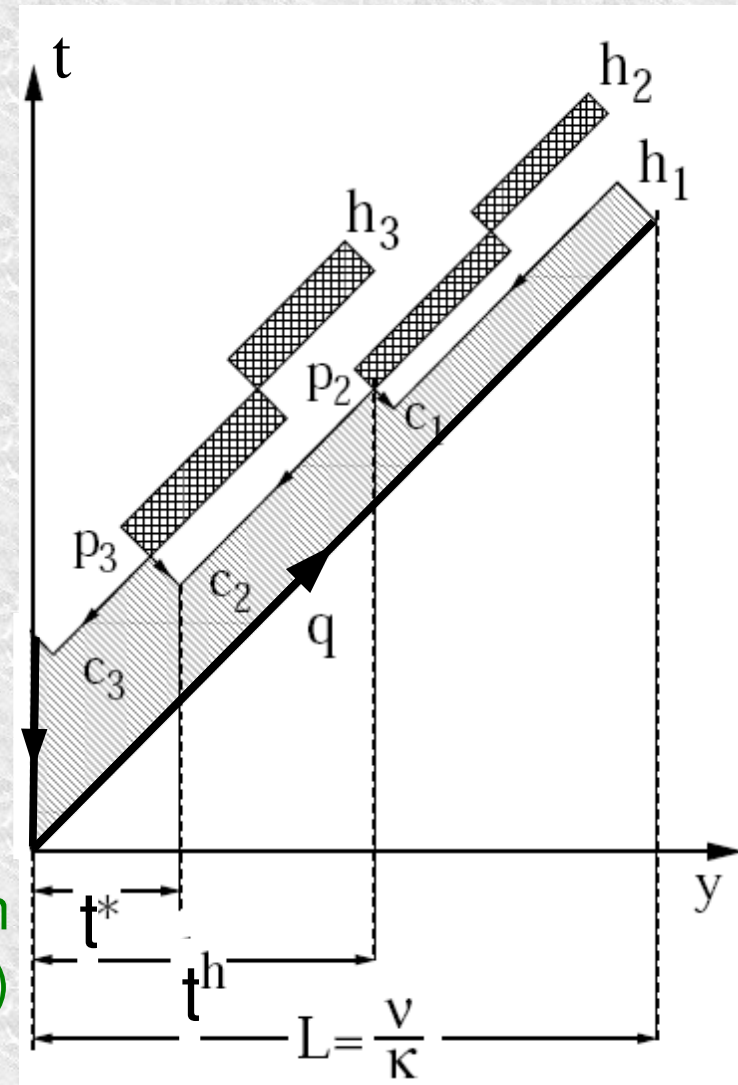
➔ At small $z \rightarrow 0$

hadron created at high rank after many string breakings: $\langle t^* \rangle \rightarrow 0$

nucleus remnant X

$$\left\{ \begin{array}{l} \langle t^* \rangle = f(z) (1-z) \frac{zV}{K} \\ \langle t^h \rangle = t^* + \frac{zV}{K} \end{array} \right.$$

← string-tension (non perturbative scale)
← energy conservation



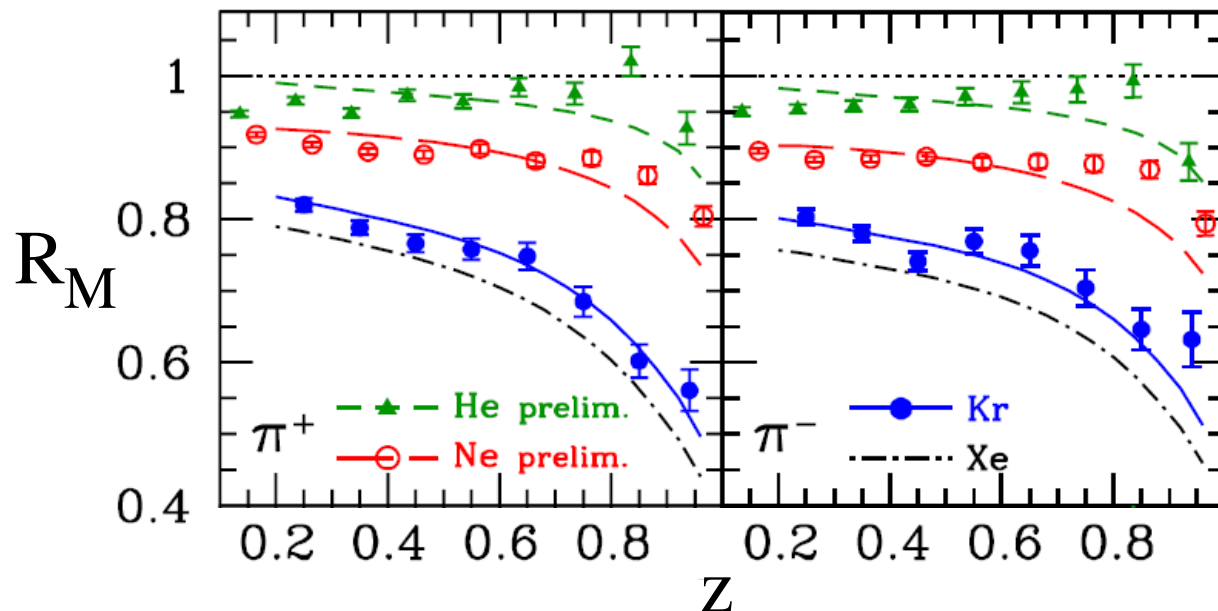
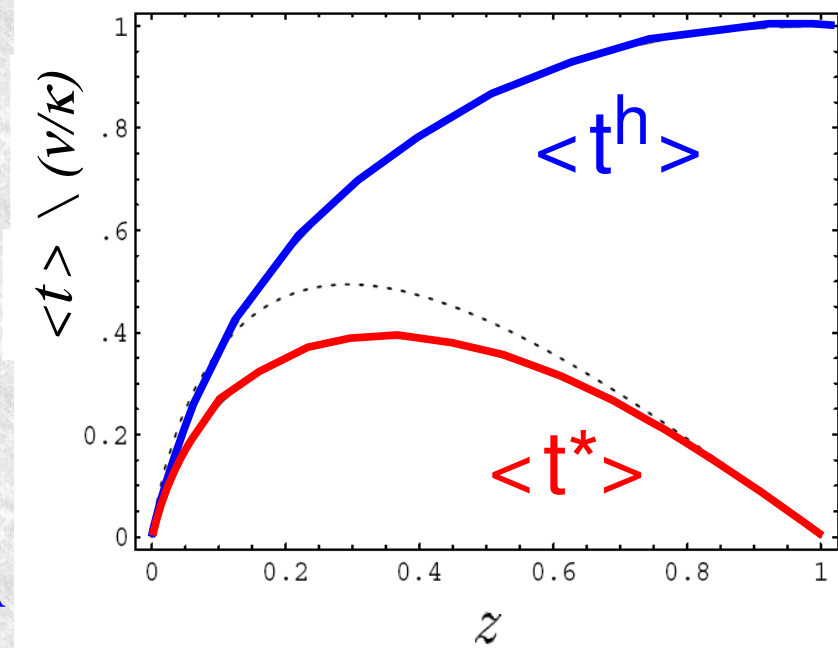
Formation time estimates 3 – Lund model

$$\begin{cases} \langle t^* \rangle = f(z) (1-z) \frac{zV}{K} \\ \langle t^h \rangle = t^* + \frac{zV}{K} \end{cases}$$

★ For a 7 GeV pion at Hermes

$$\langle t^* \rangle < 5 \text{ fm} \sim O(R_A) \quad \langle t^h \rangle \sim 10 \text{ fm} > R_A$$

★ Prehadron absorption with this estimate [A.A. et al., NPA 761(05)67]

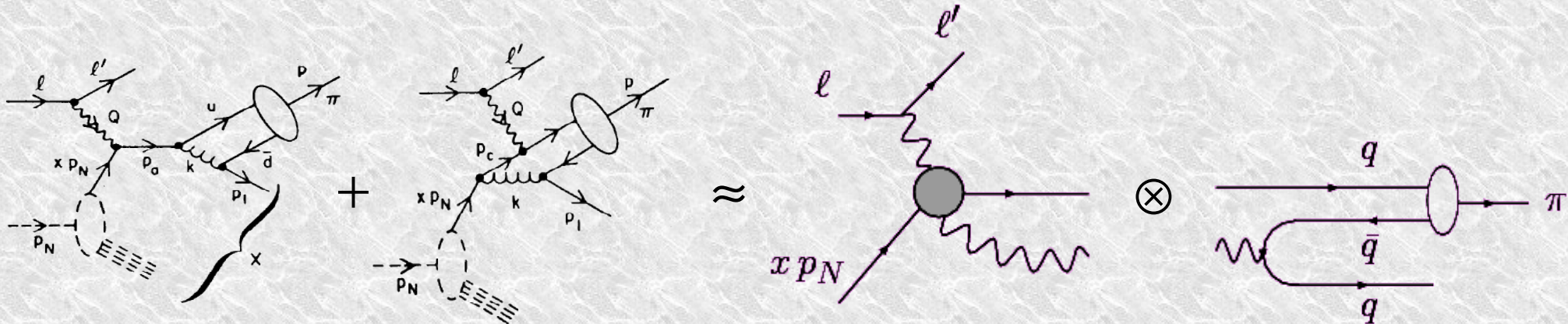


see also:

Falter, Gallmeister, nucl-th/0512104
for similar ideas in a transport model

Formation time estimates 4 – Dipole model

★ Leading hadron formation ($z > 0.5$) [Kopeliovich et al., NPA 740(04)2111]



★ Prehadron formation time t^*

= time at which gluon becomes decoherent with parent quark

➔ At large $z \rightarrow 1$

$E_h \rightarrow v \Rightarrow$ quark must be short-lived

(otherwise radiates too much energy)

$$\langle t^* \rangle \propto (1-z) \frac{z v}{Q^2}$$

energy conservation \rightarrow $(1-z)$ \leftarrow boost \leftarrow $z v$
 virtuality (perturbative scale) \leftarrow Q^2

★ Evolution into hadron by path-integral formalism

➔ usually $\langle t^h \rangle \gg R_A$

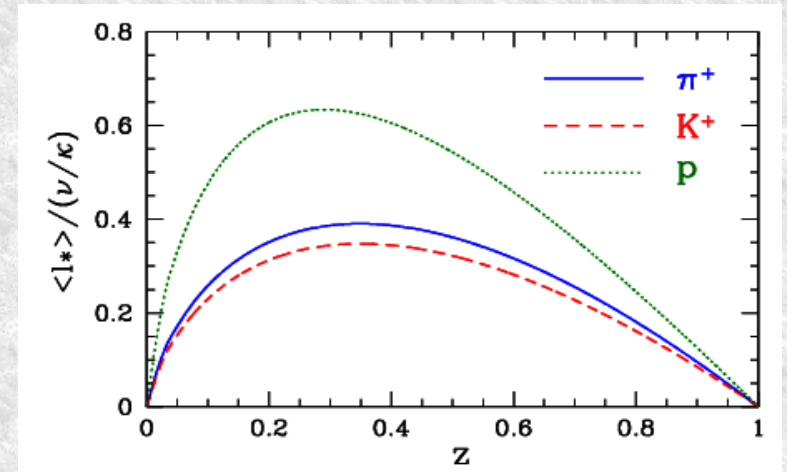
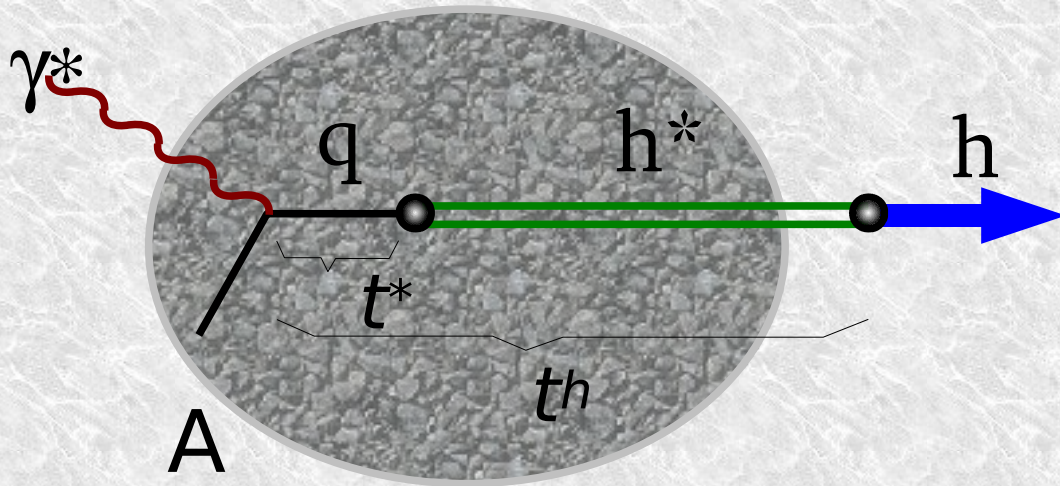
**Can we distinguish energy loss
from prehadron absorption?**

I - Hadron absorption model

A.A. et al., NPA 761(05)67

Hadron absorption model – 1

A.A. et al., NPA 761(05)67



- Two-step hadronization inside the nucleus:
 - quark q neutralizes color \Rightarrow prehadron h^*
 - hadron h 's wavefunction fully develops
- Average formation lengths $\langle t^* \rangle (z, \nu)$, $\langle t^h \rangle (z, \nu)$ taken from theLund model
 - for a 10 GeV π : $\langle t^* \rangle \sim 4$ fm $\langle t^h \rangle \sim 10$ fm

Hadron absorption model – 2

A.A. et al., NPA 761(05)67

- ◆ (Pre)hadron suppression factor S^A by transport diff. eqns.
- ◆ (Pre)hadron-nucleon cross sections:
 - $\sigma^* = 0.35 \sigma^h$ - fitted to $e^+ + \text{Kr} \rightarrow \pi^+ + X$
 - σ^h - from Particle Data Group
- ◆ Full integration over γ^*q interaction point (b,y)

$$S_{f,h}^A = \int d^2b dy \rho_A(b,y) \int_y^\infty dx' \int_y^{x'} dx \frac{e^{-\frac{x-y}{\langle l^* \rangle}}}{\langle l^* \rangle} e^{-\sigma^* \int_x^{x'} ds A \rho_A(b,s)} \frac{e^{-\frac{x'-x}{\langle \Delta l \rangle}}}{\langle \Delta l \rangle} e^{-\sigma_h \int_{x'}^\infty ds A \rho_A(b,s)}$$

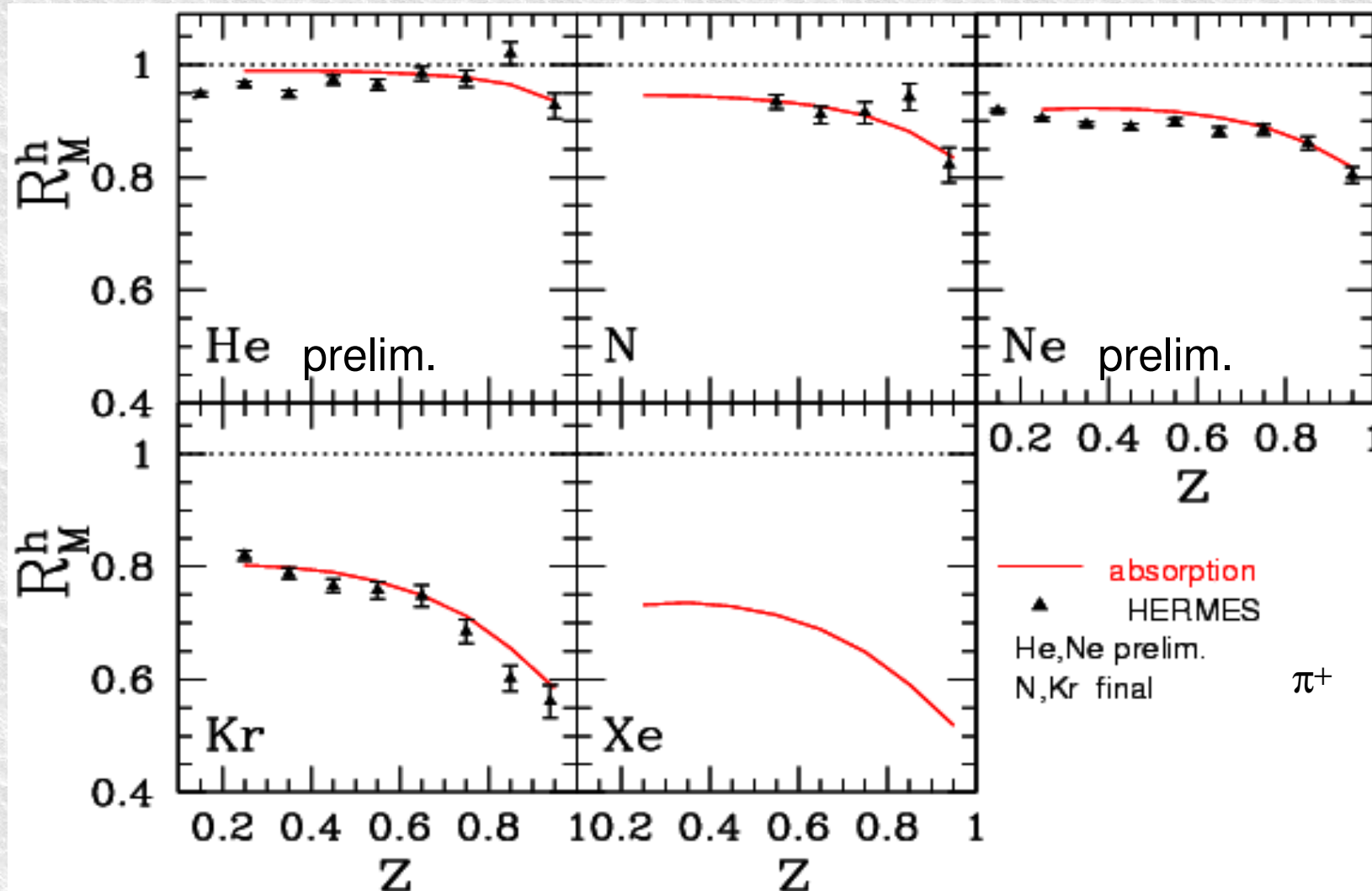
prob. that h^* is formed at x
absorption of h^* up to x'
prob. that h is formed at x'
absorption of h from x' to ∞

$$\frac{1}{N_A^{DIS}} \frac{dN_A^h(z)}{dz} = \frac{1}{\sigma^{lA}} \int dx d\nu \sum_f e_f^2 q_f(x, Q^2) \frac{d\sigma^{lq}}{dx d\nu} S_{f,h}^A(z, \nu) D_f^h(z, Q^2)$$

exp. cuts

Hadron absorption model - results

A.A. et al., NPA 761(05)67



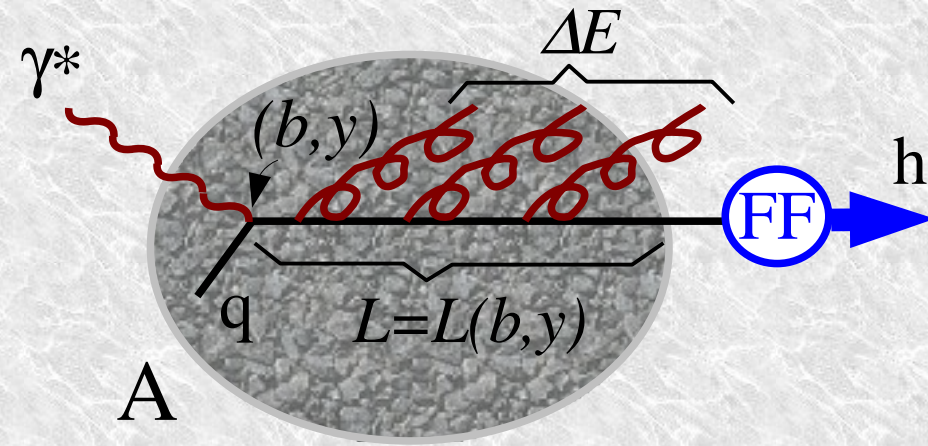
Note: in ref. above, curves differ by inclusion of Q^2 -rescaling due to additional hypothesis of partial deconfinement in nuclei

II - Energy loss model

A.A. (QM'05) nucl-th/0510090

Energy loss model

A.A. (QM'05) nucl-th/0510090



- Quark hadronizes outside the nucleus
- Gluon bremsstrahlung $\Rightarrow \Delta E = v \Delta z$
 \Rightarrow modified Fragmentation Function

[see Arleo EPJ C30(02)213]

$$D_q^h(z, Q^2) \longrightarrow \frac{1}{1 - \Delta z} D_q^h\left(\frac{z}{1 - \Delta z}, Q^2\right)$$

- New: realistic geometry + quenching weights $\mathcal{P}(\Delta z, L)$ corrected for finite in-medium path $L=L(b, y)$

[Salgado-Wiedemann '03]

$$\tilde{D}_f^h(z, Q^2; L) = \int_0^{(1-z)} d\Delta z \mathcal{P}(\Delta z; \hat{q}, L) \frac{1}{1 - \Delta z} D_f^h\left(\frac{z}{1 - \Delta z}, Q^2\right) + p_0(\hat{q}, L) D_f^h(z, Q^2)$$

Quenching weight:

prob. of radiating Δz

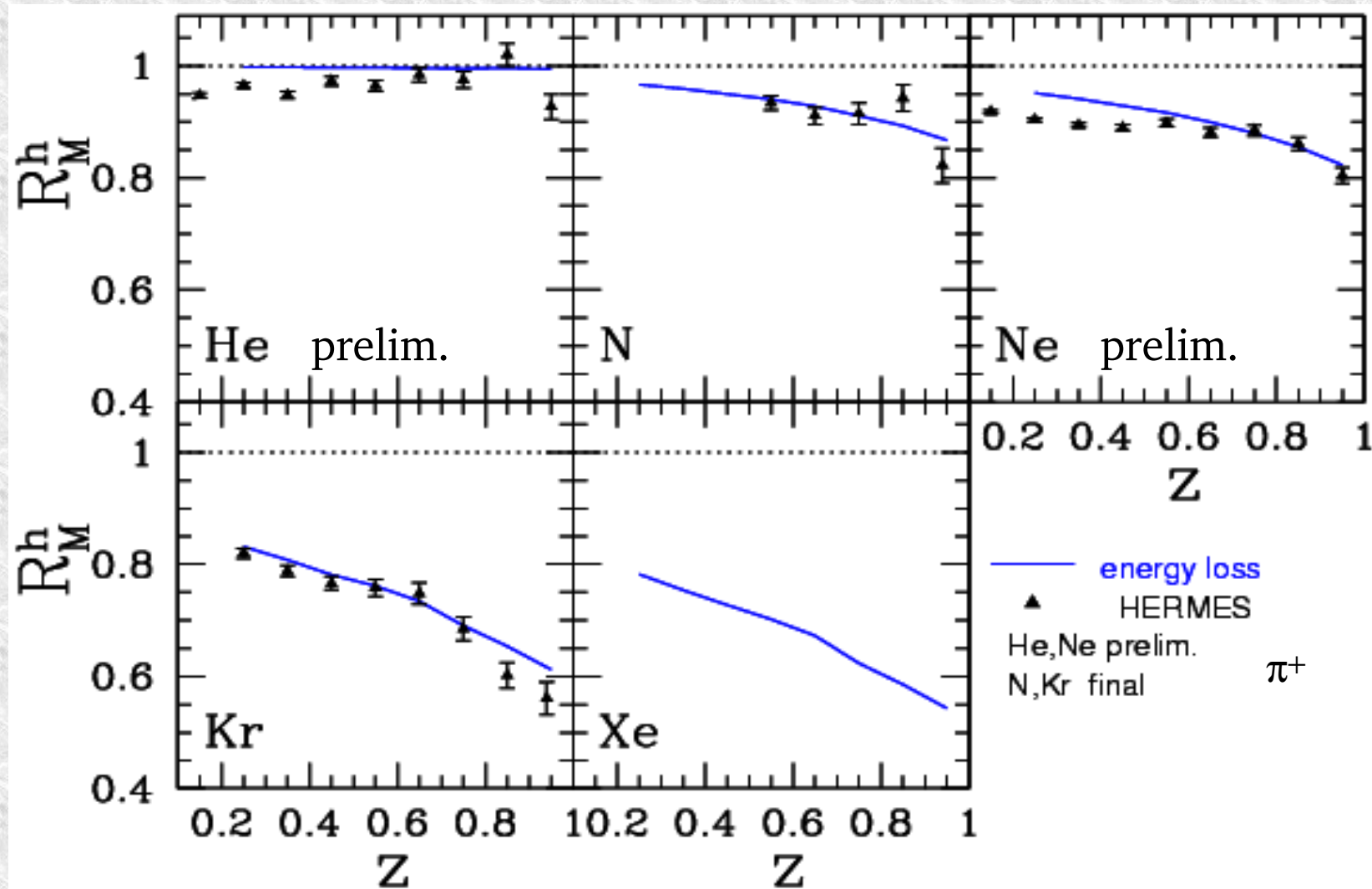
prob. of radiating no gluons

- Transport coefficient $\langle \hat{q} \rangle(b, y) = \hat{q}_0 \frac{\langle \rho \rangle(b, y)}{\rho(0, 0)}$ where $\hat{q}_0 = \langle \hat{q} \rangle(0, 0)$

with $\hat{q}_0 = 0.5 \text{ GeV}^2/\text{fm}$ - fitted to $e^+ + \text{Kr} \rightarrow \pi^+ + X$

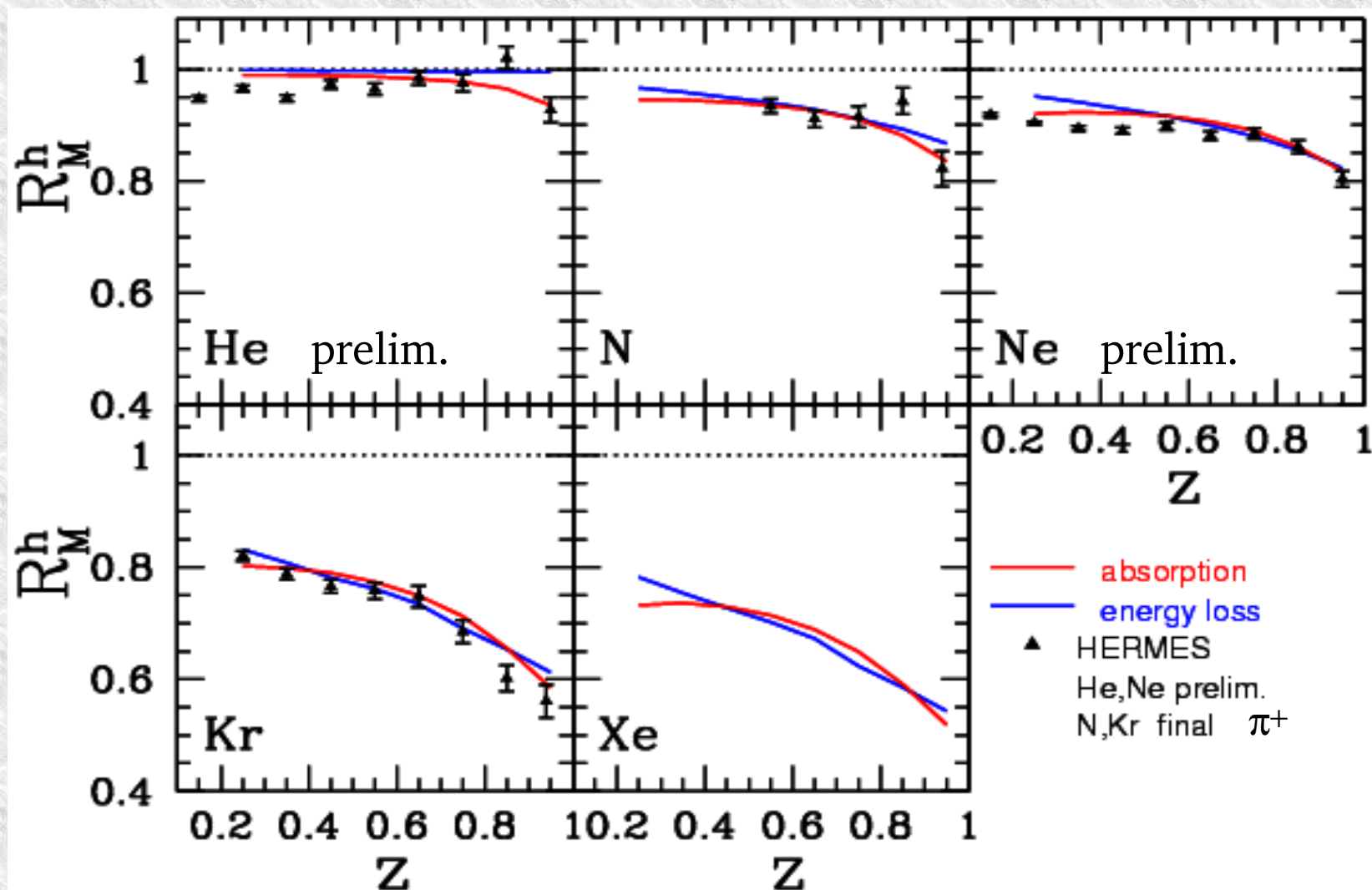
Energy loss model - results

A.A. (QM'05) nucl-th/0510090



Energy loss vs. absorption

A.A. (QM'05) nucl-th/0510090



- Both models account well for HERMES R_M data
- Surprisingly similar up to heavy nuclei

III – How to tell energy loss
from absorption?

1) The “ $A^{2/3}$ power law”

A.A., Gruenewald, Muccifora, Pirner, NPA 761(2005)67

◆ Conventional thinking: the $A^{2/3}$ law

◆ Energy loss (LPM effect in QCD): $1 - R_M \sim \langle \Delta z \rangle \sim L^2 \sim A^{2/3}$

◆ Hadron absorption: $1 - R_M \sim \langle \text{no. of nucleons seen} \rangle \sim L \sim A^{1/3}$

WRONG!

◆ $A^{2/3}$ also for absorption models!

◆ extra dimensionful scale: prehadron formation length $\langle l^* \rangle$

◆ neutralize it \Rightarrow extra power of A $(R_A / \langle l^* \rangle)^n \sim A^{n/3}$

◆ typically $n=1$

1) The “ $A^{2/3}$ power law”

A.A., Gruenewald, Muccifora, Pirner, NPA 761(2005)67

➤ Don't believe it?

➤ try fitting $1-R_M = cA^\alpha$

[Accardi et al., NPA 761 (05);
Accardi, Acta Phys. Hung. in press]

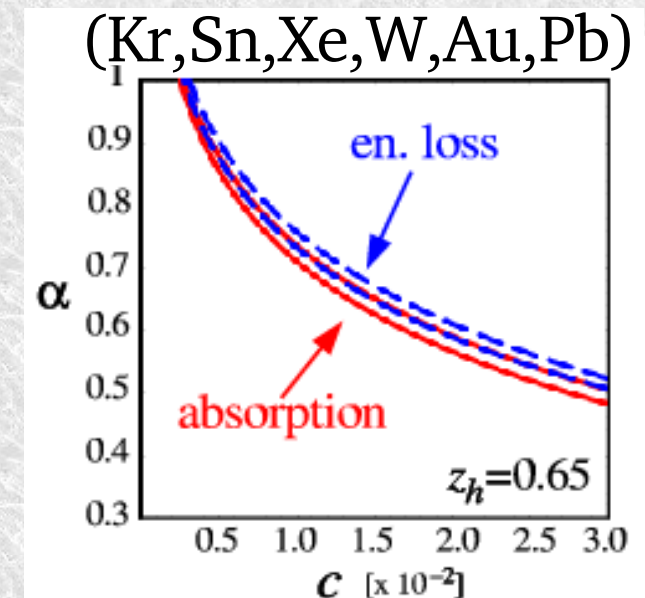
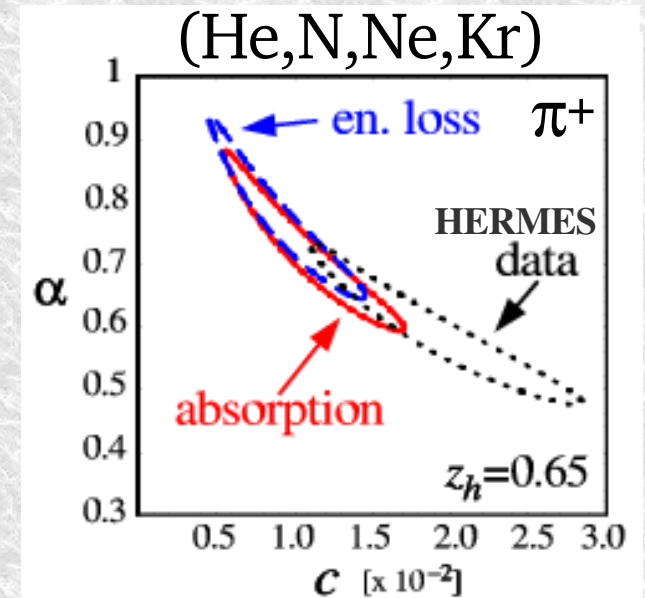
➤ $A^{2/3}$ broken on large nuclei

➤ Absorption & en. loss mimick each other

➤ not possible to distinguish (separate)
the 2 mechanisms

➤ not even using heavy nuclei up to Pb

**A-dependence of R_M does not test
dominance of partonic or prehadronic physics**



2) Scaling of R_M – basic idea

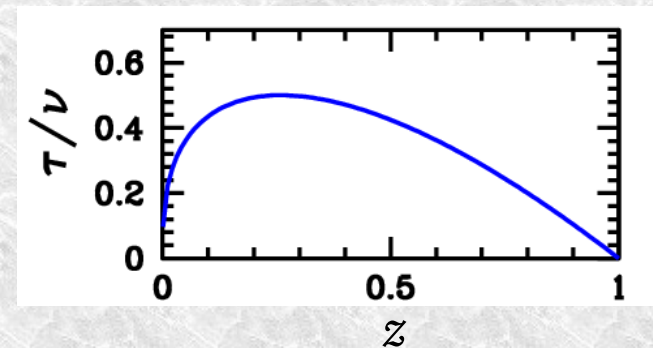
A.A., nucl-th/0604041, PLB in press

- R_M should scale with $\tau = \tau(z, \nu)$ not with z and ν separately

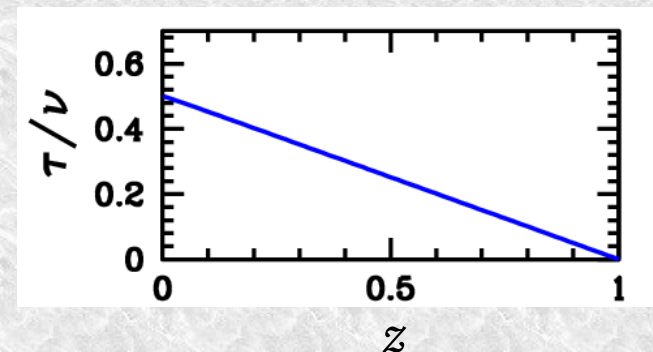
$$R_M = R_M[\tau(z, \nu)] \quad \text{with} \quad \tau = C z^\lambda (1-z) \nu$$

- “Scaling exponent” λ can distinguish absorption and energy-loss

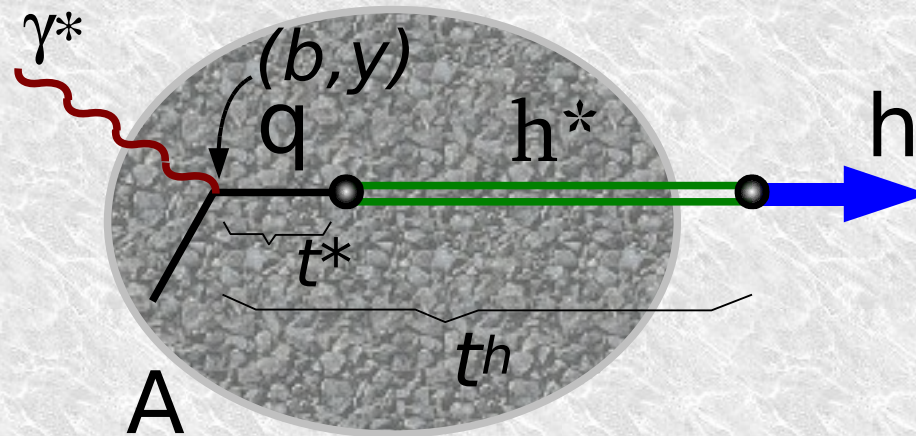
- **absorption models:** $\lambda > 0$
[finite formation time]



- **energy loss models:** $\lambda \leq 0$
[from energy conservation:
radiated energy $\rightarrow \epsilon < (1-z_h) \nu$]



2) Scaling of R_M – Absorption models



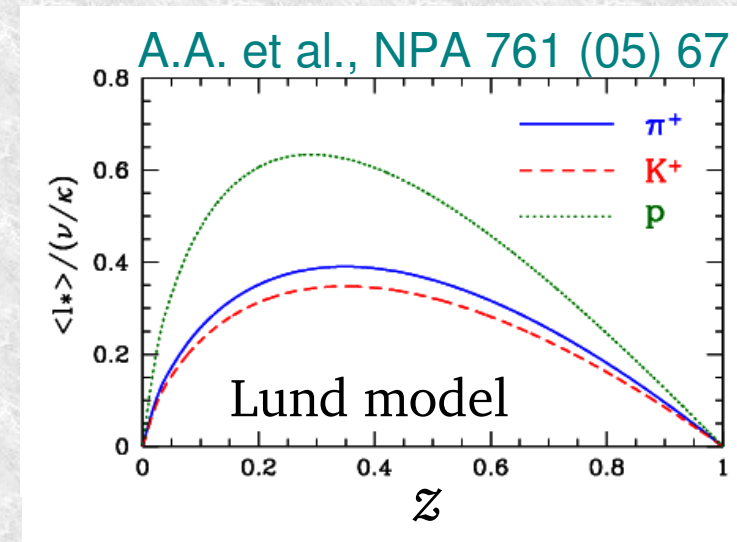
◆ Extreme situation: if quark has no interactions,
 $\Rightarrow R_M$ depends only on the prehadron's in-medium path, i.e. $\langle t^* \rangle$

◆ General form of $\langle t^* \rangle$:

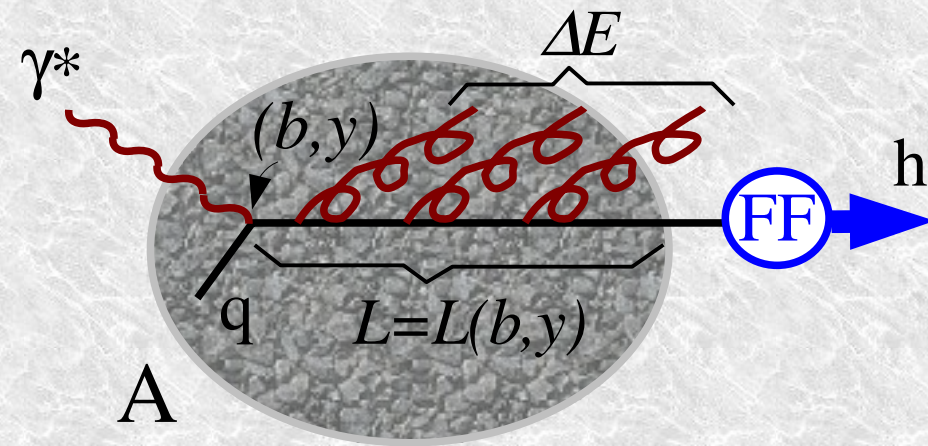
$$\langle t^* \rangle = f(z) (1-z) \frac{zV}{\kappa} \sim \tau(z, V)$$

← Lorentz boost
← energy conservation
← model dep. scale (e.g. string tension)

◆ obtain τ -scaling with $\lambda > 0$



2) Scaling of R_M – Energy loss models



define average energy loss:

energy conservation: $\epsilon < \nu - E_h$

$$\langle \epsilon \rangle = \int_0^{(1-z_h)\nu} d\epsilon \epsilon \mathcal{P}(\epsilon) = f[(1-z_h)\nu]$$

Approximate: $R_M \approx \frac{1}{1 - \langle \epsilon \rangle / \nu} D\left(\frac{z_h}{1 - \langle \epsilon \rangle / \nu}\right) / D(z_h)$

Use KKP at $Q^2 = 2 \text{ GeV}^2$ $D(z_h) = C z_h^\alpha (1 - z_h)^\beta$

Obtain **approximate τ scaling with $\lambda \approx 0$** :

$$R_M \approx \frac{1}{\left(1 - \frac{1}{\nu} f[(1-z_h)\nu]\right)^{\alpha+\beta+1}} \left(1 - \frac{f[(1-z_h)\nu]}{(1-z_h)\nu}\right)^\beta \sim \tau(\lambda=0)$$

in practice, $\lambda \leq 0$

2) Scaling analysis

A.A., nucl-th/0604041, PLB in press

◆ HERMES data presented as:

➡ z -distributions: $R_M(z)$ with $\langle v \rangle = \langle v \rangle(z)$

➡ v -distributions: $R_M(v)$ with $\langle z \rangle = \langle z \rangle(v)$

◆ Scaling analysis of R_M – remember $\tau = C z^\lambda (1-z) v$

1) Fix λ

2) For each z compute $\tau = \tau(z, \langle v \rangle)$ and $R_M(\tau) = R_M(z)$

For each v compute $\tau = \tau(\langle z \rangle, v)$ and $R_M(\tau) = R_M(v)$

3) Fit $\phi(\tau)$ to $\{\tau, R_M\}$ obtained above – $\chi^2 = \chi^2(\lambda)$

$$\phi(\tau) = a + b\tau + c\tau^2 + d\tau^3 + e\tau^4$$

$$\left(\text{cross-checked by } \phi(\tau) = \frac{a + b\tau + c\tau^2}{a + d\tau + e\tau^2} \right)$$

4) Best-fit λ_{best} by χ^2 minimization

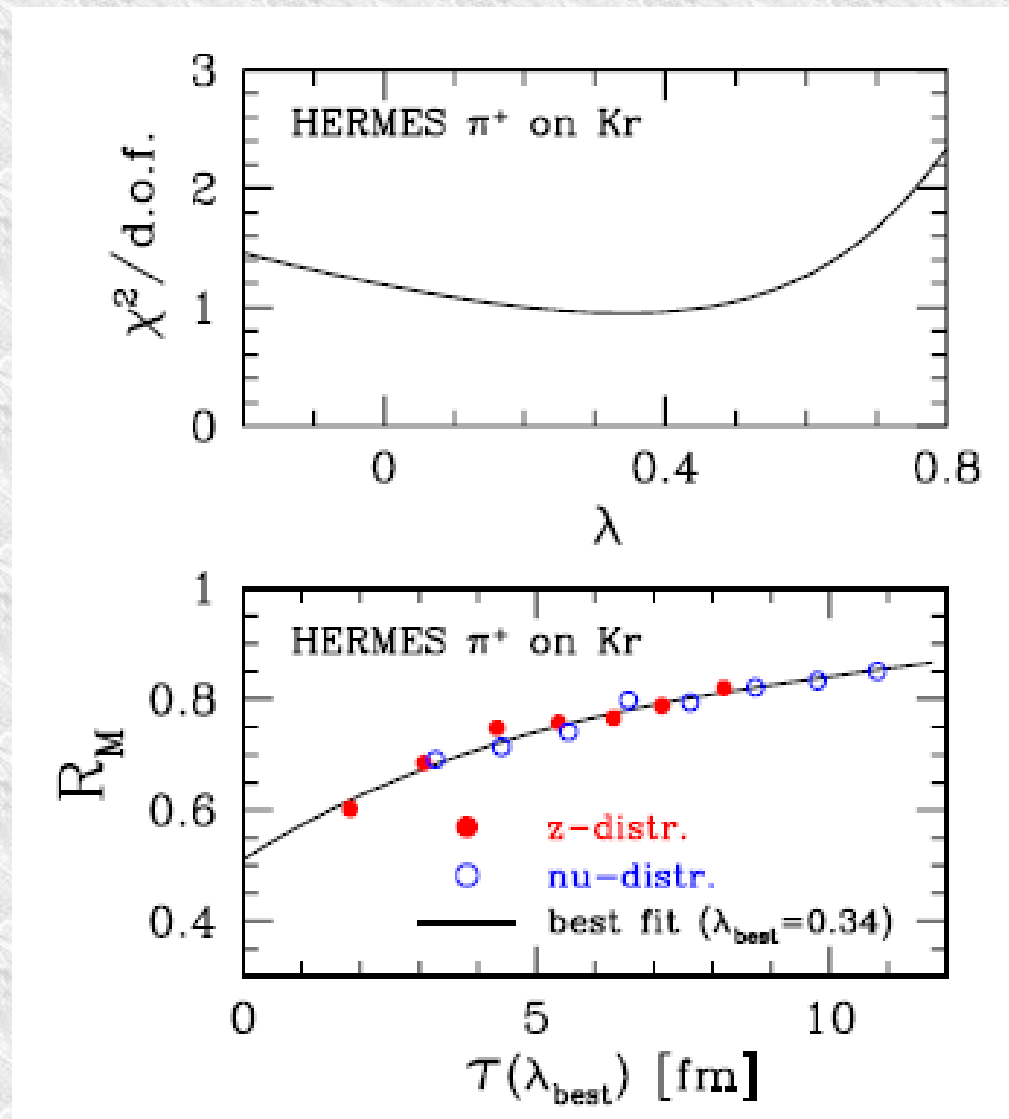
5) If $\chi^2(\lambda_{\text{best}})/\text{d.o.f.} \lesssim 1 \Rightarrow R_M$ scales with τ , and is characterized by λ_{best}

◆ NOTE: the overall scale of τ cannot be extracted

(R_M scales, or doesn't scale, independently of the value of the overall constant C)

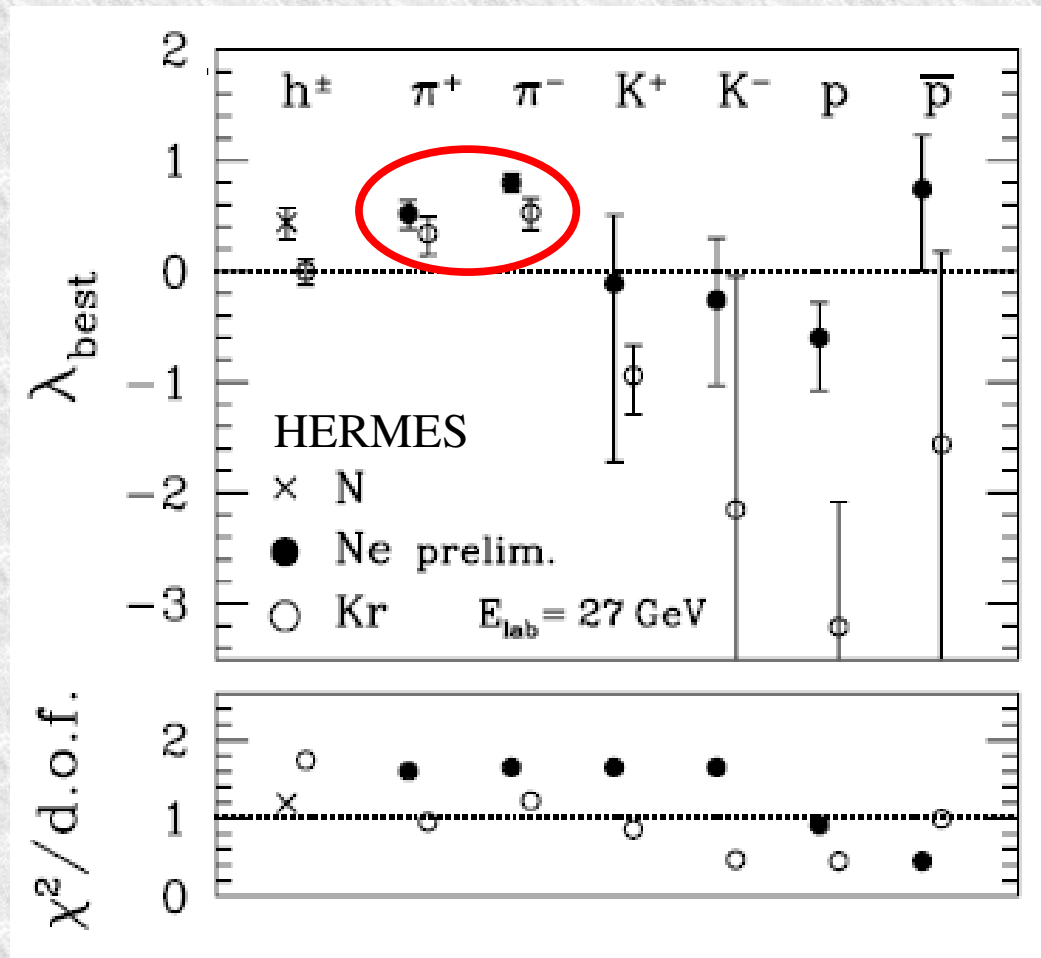
2) Scaling analysis - example

A.A., nucl-th/0604041, PLB in press



2) Results – $E_{\text{lab}} = 27 \text{ GeV}$

A.A., nucl-th/0604041, PLB in press

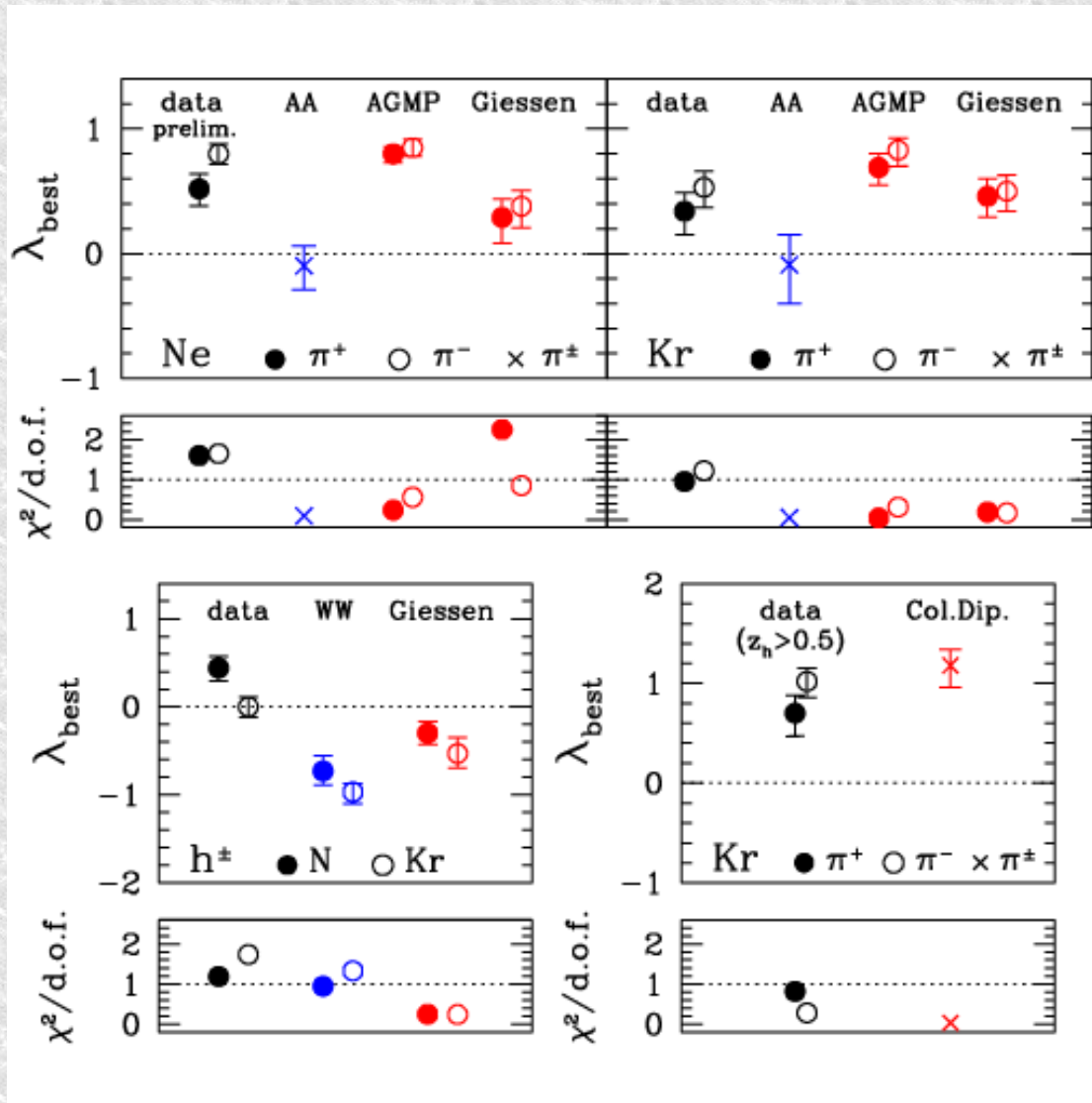


◆ $\lambda(\pi) > 0$: Formation-time scaling for pions!

Hadronization starts inside the nucleus!

2) Data vs. theory

A.A., nucl-th/0604041, PLB in press



➤ **Experimental data:**

➤ HERMES

➤ **Energy loss models:**

➤ AA = Arleo + realistic geometry (A.A.)

➤ WW = Wang & Wang

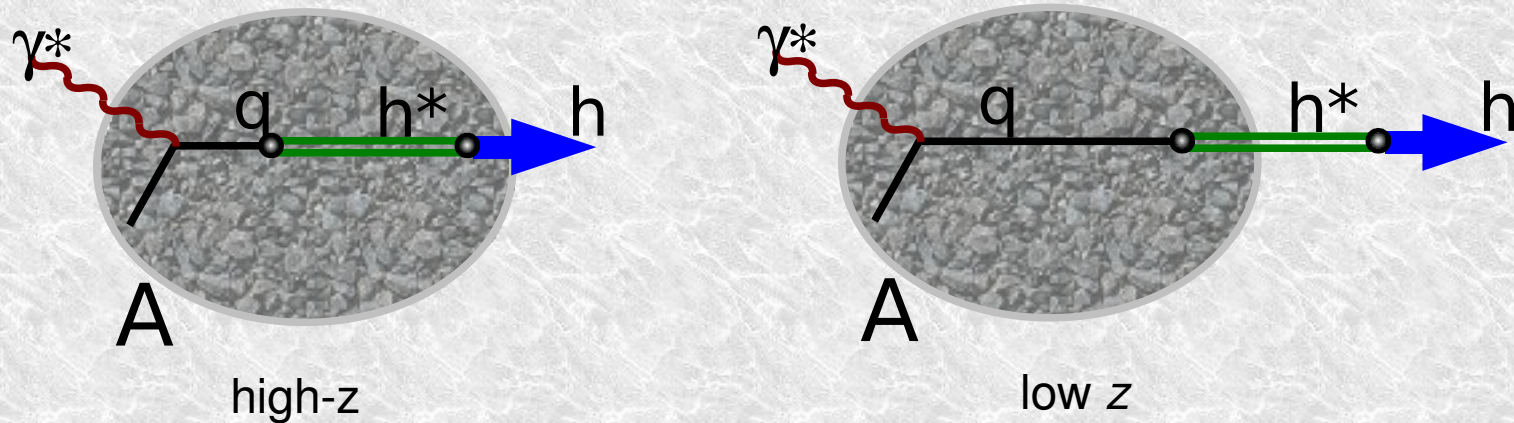
➤ **Absorption models:**

➤ AGMP = Accardi, Gruenewald, Muccifora, Pirner

➤ Giessen = Falter et al.

➤ Col.Dip. = Kopeliovich et al.

3) p_T – broadening



◆ $\langle p_T^2 \rangle$ broadening [Kopeliovich et al., NPA 740(04)211]

- 1) Directly proportional to quark's in-medium path
- 2) Can measure prehadron formation time t^*
- 3) Detect hadronization inside or outside the nucleus

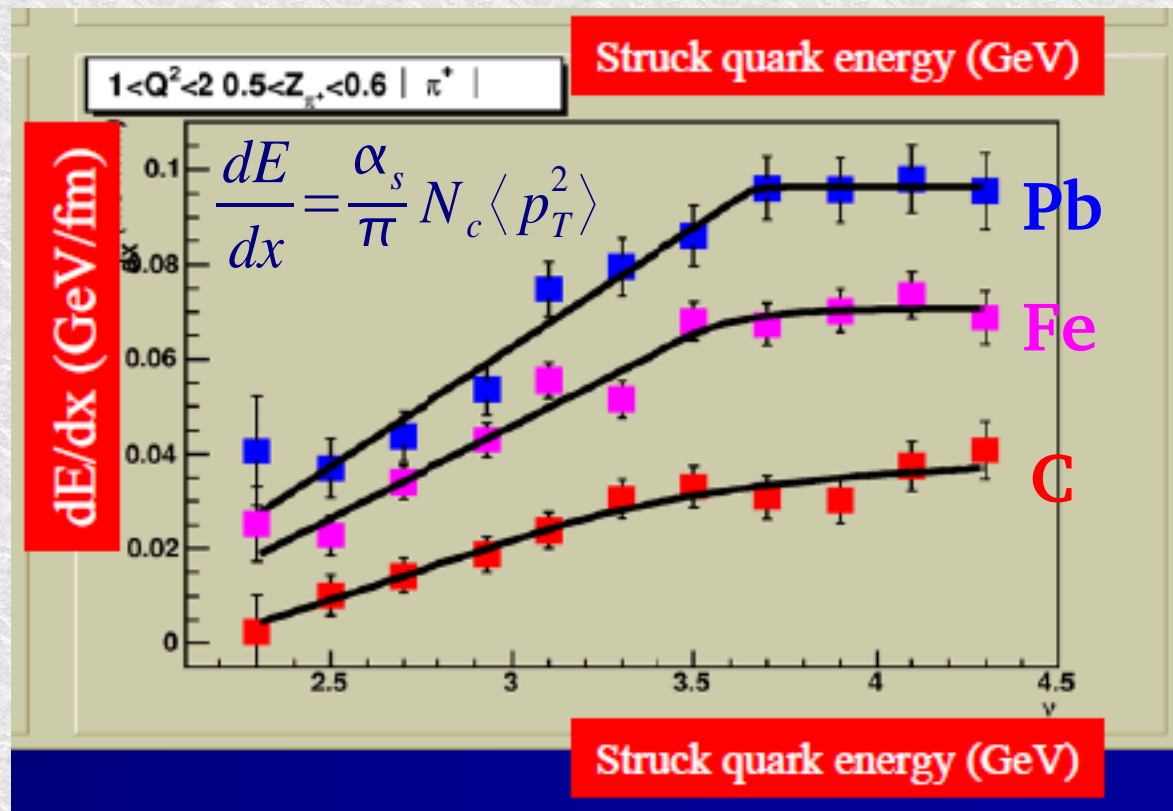
$$\Delta \langle p_T^2(L) \rangle = 2C(s) \int_0^L dz \rho_A(z), \quad \text{where:} \quad C(s) = \frac{d\sigma_{\bar{q}q}(r_T, s)}{dr_T^2} \Big|_{r_T=0}$$

dipole x-sect.

◆ Can be cross-checked by the scaling analysis of R_M

3) p_T – broadening

- If prehadron inside nucleus \Rightarrow quark path length varies with z_h, v
- Δp_T should:
 - 1) rise with v until $\langle t^* \rangle \sim R_A$, then level off
 - 2) decrease as $z_h \rightarrow 1$
 - 3) possibly, decrease as Q^2 increases



JLAB $E_{\text{lab}} = 5$ GeV – preliminary results
 Will Brooks, JLAB user's group workshop, 13 June 2006

3) p_T – broadening

➤ “Model independent” measurement of $\langle l^* \rangle = l_p$

[Kopeliovich, Nemchik, Schmidt, hep-ph/0608044]

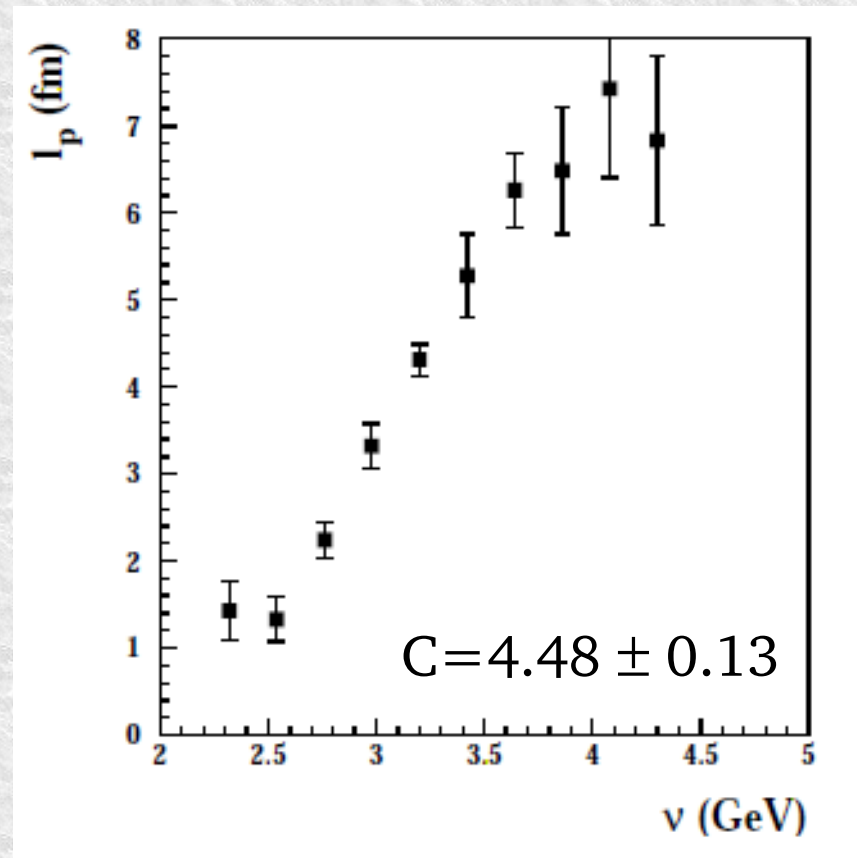
$$\Delta p_T^2 = \frac{2Cz_h^2}{A} \int d^2b \int_{-\infty}^{\infty} dz \rho_A(b, z) \int_z^{z+l_p} dz' \rho_A(b, z')$$

where

dipole cross-section

$$C(s) = \left. \frac{d\sigma_{\bar{q}q}(r_T, s)}{dr_T^2} \right|_{r_T=0}$$

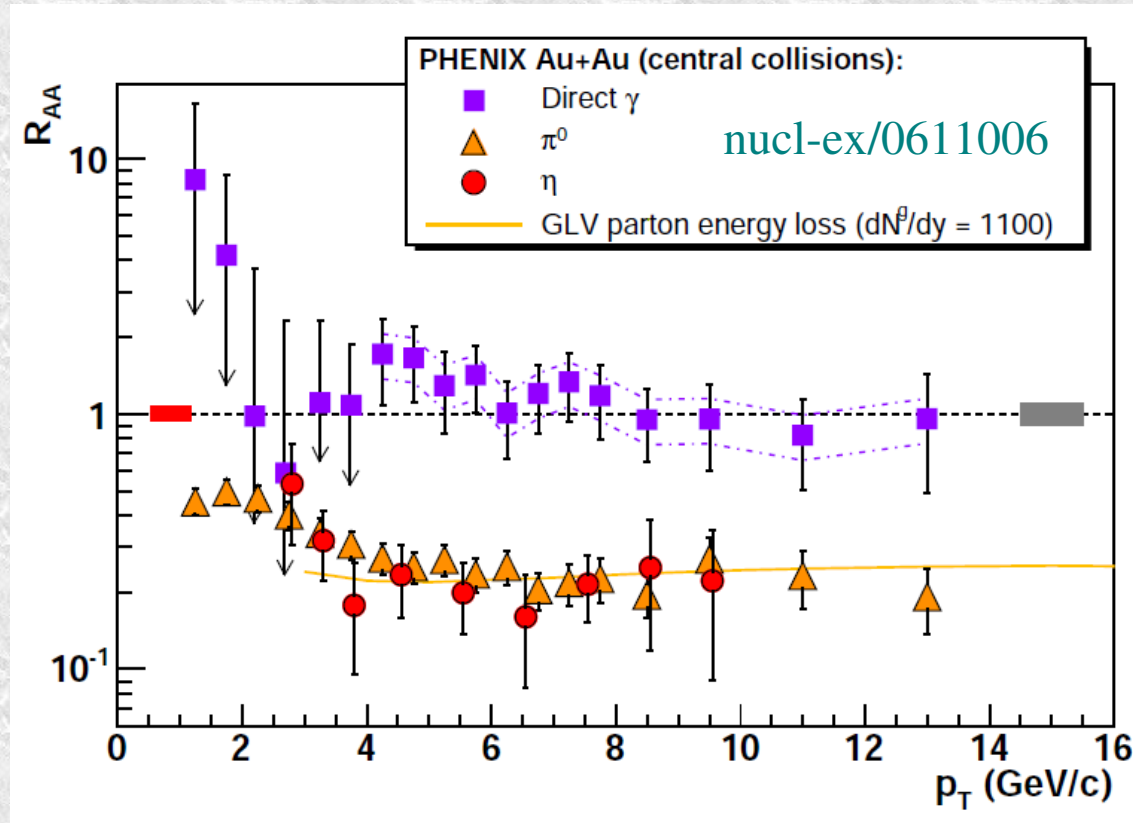
- 1) fit l_p to data for each nucleus
- 2) determine C by minimizing differences of l_p among nuclei



Perspectives

Perspectives 1 – mesons

- ★ Why is η as much suppressed as π in Au+Au collision?
 - ➔ points towards long lived quark
 - ➔ but scaling analysis suggests pions formed on short time scales



- ★ Is it so also in nDIS? [η is heavier \Rightarrow hadronizes earlier]
 - ➔ measurement possible at HERMES, CLAS @ JLAB

Perspectives 2 – baryons

- ★ Is the baryon anomaly in nDIS only for protons? \Rightarrow measure Λ !
(at RHIC R_{dAu} and R_{AuAu} similar for p and Λ)
- ★ possible at HERMES and CLAS @ JLAB

production rate @ CLAS++ 12 GeV (not folded with PID) [W.Brooks, FizikaB13(04)321]

hadron	$c\tau$	mass (GeV)	flavor content	detection channel	production rate per 1k DIS events
p	stable	0.94	ud	direct	1100
\bar{p}	stable	0.94	$\bar{u}\bar{d}$	direct	3
Λ	79 mm	1.1	uds	$p\pi^-$	72
$\Lambda(1520)$	13 fm	1.5	uds	$p\pi^-$	-
Σ^+	24 mm	1.2	us	$p\pi^0$	6
Σ^0	22 pm	1.2	uds	$\Lambda\gamma$	11
Ξ^0	87 mm	1.3	us	$\Lambda\pi^0$	0.6
Ξ^-	49 mm	1.3	ds	$\Lambda\pi^-$	0.9

- ★ More in general: clarify baryon production:
 - proton anomaly / antiproton normality
 - diquark content of nucleons
 - baryon transport

Perspectives 3 – heavy quarks

★ Heavy-quark energy loss and hadronization of D, B mesons in the spotlight at RHIC

➤ measure D, B in e+A !

★ At HERMES

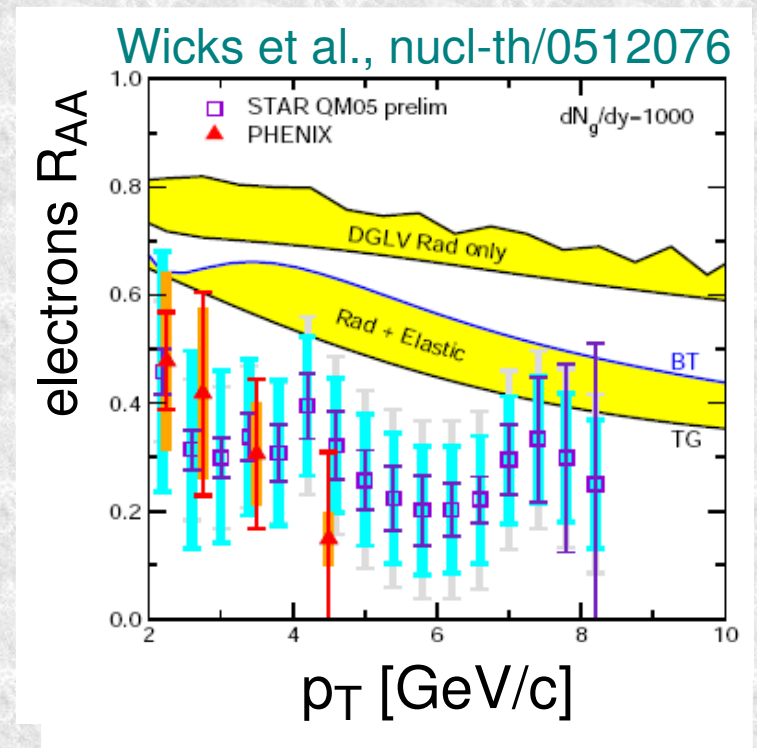
➤ luminosity is too low for D meson

★ At Jlab 12 GeV

➤ high luminosity may compensate for low- v and large- x (and PID)

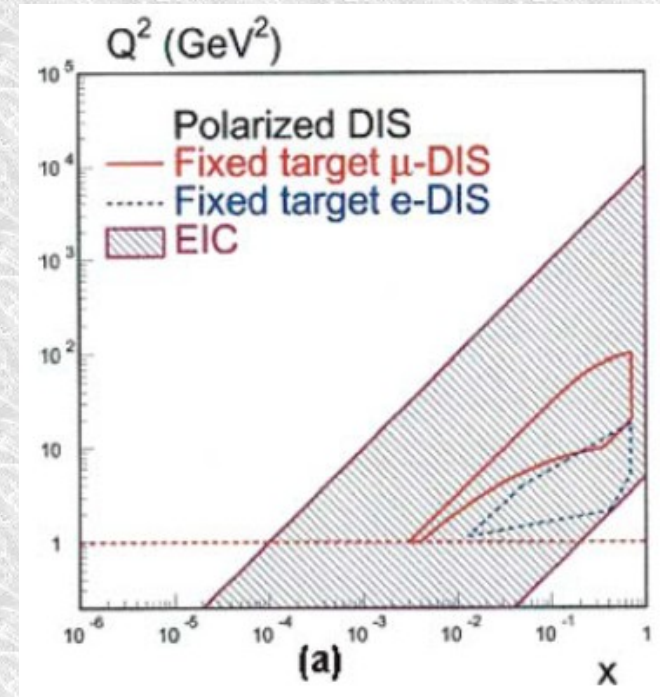
➤ chances for D meson measurement close to but not zero

★ Needs an Electron-Ion Collider!

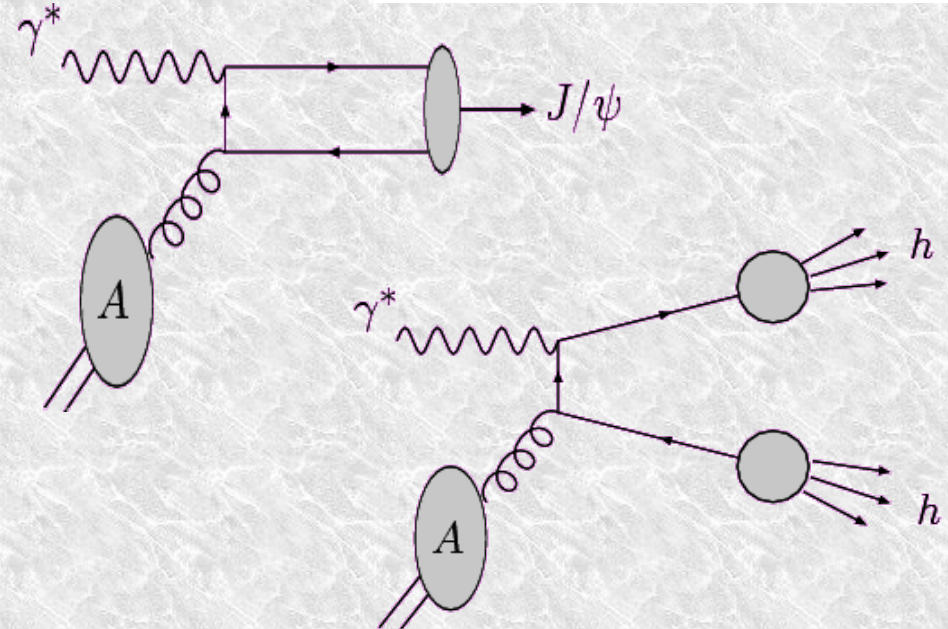


Perspectives 4 – eRHIC / ELIC

- ★ Can repeat HERMES / JLAB
- ★ Large ν -span: $10 \text{ GeV} < \nu < 1600 \text{ GeV}$
 - ➔ hadronization inside/outside target
 - ➔ test parton energy loss in cold nuke matter
 - ➔ test of factorization for FF
 - ➔ larger phase space for heavy hadrons
 - ➔ jet physics



- ★ Small x :
 - ➔ heavy quarks \Rightarrow D, B mesons
 - ➔ J/ ψ “normal suppression”
 - ➔ “away-side” correlations: hadron-hadron, γ -hadron



Conclusions

- ★ Space-time evolution of hadronization important:
 - ➔ on its own: understanding of color confinement
 - ➔ for correct interpretation of hadron suppression in A+A
- ★ τ -scaling analysis suggests small formation times for π
 - ➔ measures $t^*(\pi) \propto z^{0.6} (1-z) v$
- ★ pT-broadening confirms small t^*
 - ➔ measures $t^*(\pi) \propto 7 \text{ fm}$ at $v \sim 4 \text{ GeV}$ and $Q^2 \sim 1.5 \text{ GeV}^2$
- ★ A new challenge to the energy loss paradigm in A+A?
- ★ Rich hadronization physics to be studied at
 - ➔ HERMES @ DESY
 - ➔ CLAS, Hall-A @ JLAB, JLAB12
 - ➔ eRHIC/ELIC
- ★ Enough phenomenology in nDIS: can we address the space-time evolution of hadronization at a more fundamental level?

The end

Backup slides

Energy loss model - realistic geometry

- ➔ New: Full integration over γ^*q interaction point (b,y)

$$\frac{1}{N_A^{DIS}} \frac{dN_A^h(z)}{dz} = \frac{1}{\sigma^{lA}} \int d^2b dy \rho_A(b, y) \int_{\text{exp. cuts}} dx d\nu \sum_f e_f^2 q_f(x, Q^2) \frac{d\sigma^{lq}}{dx d\nu} \tilde{D}_f^h(z, Q^2; L(b, y))$$

- ➔ Realistic geometry: Woods-Saxon parametrization for $A > 2$
Reid's soft-core for 2D

$$L(b, y) = 2 \int_y^\infty dz (z - y) \rho(b, z) \Big/ \int_y^\infty dz \rho(b, z) \quad (= R(b) - y \text{ for Hard-Sphere})$$

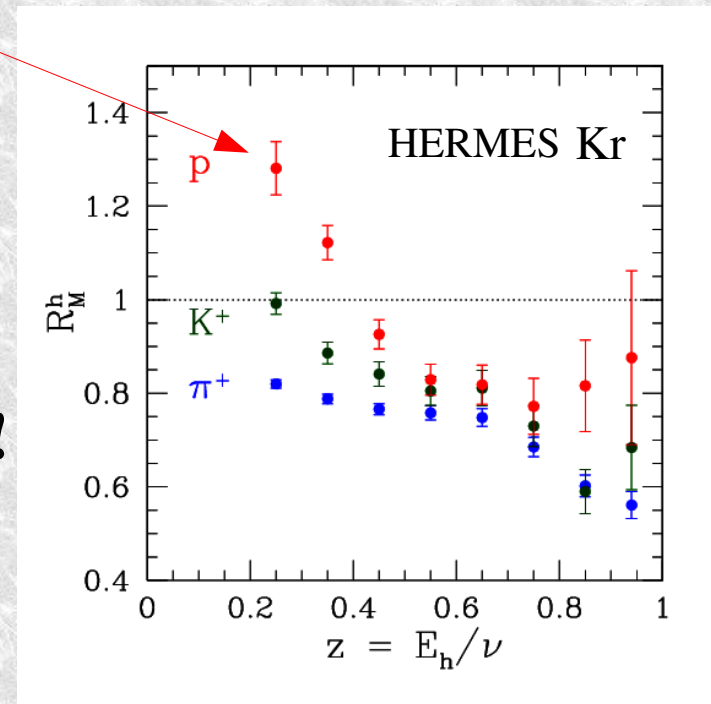
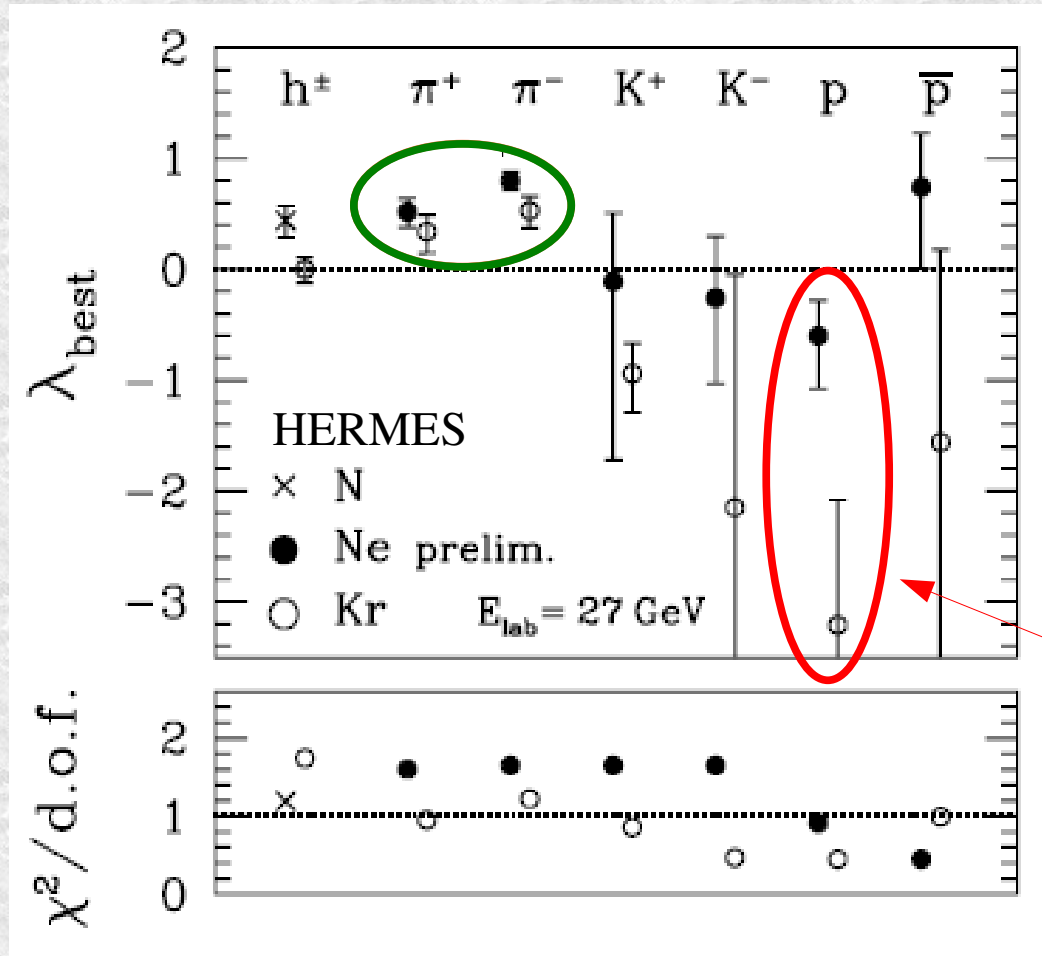
$$\langle \rho \rangle(b, y) = \int_y^\infty dz \rho(b, z) \Big/ L(b, y) \quad (= \rho_{HS} \text{ for Hard-Sphere})$$

- ➔ Transport coefficient

$$\langle \hat{q} \rangle(b, y) = \hat{q}_0 \frac{\langle \rho \rangle(b, y)}{\rho(0, 0)} \quad \text{where } \hat{q}_0 = \langle \hat{q} \rangle(0, 0)$$

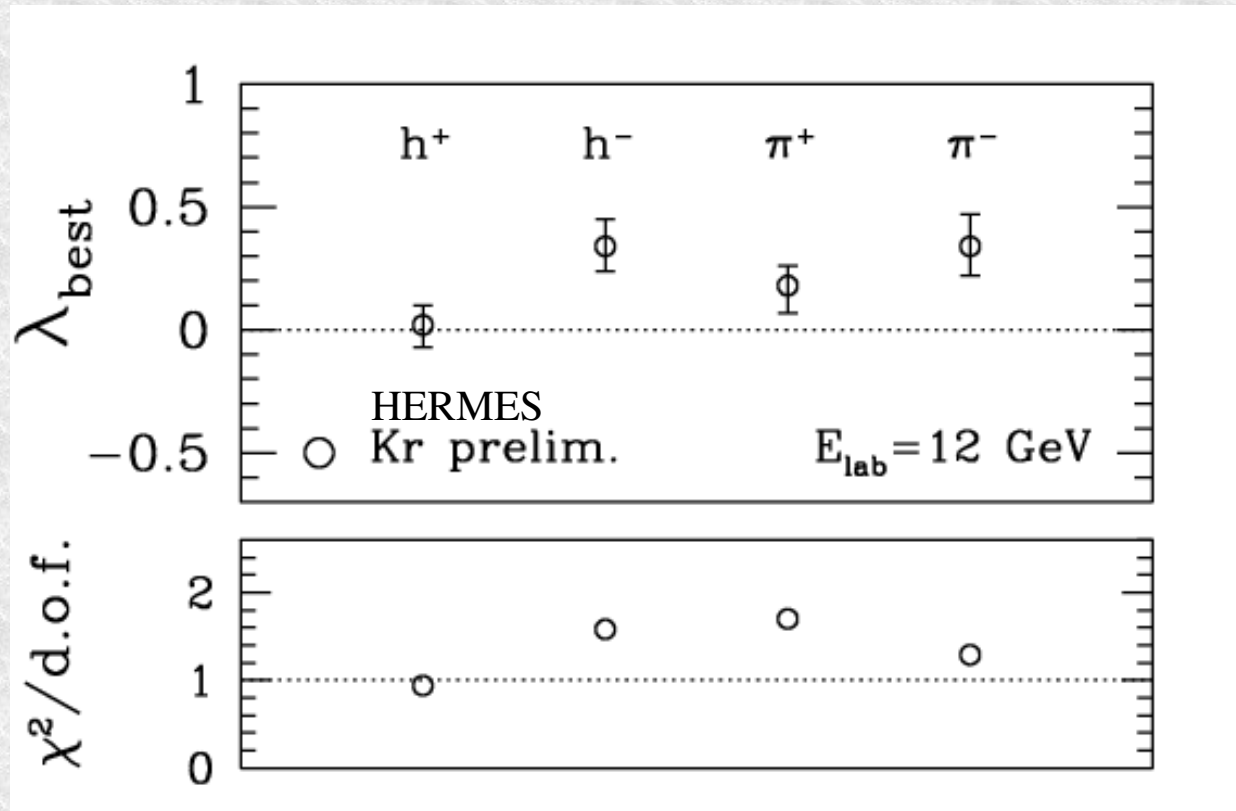
with $\hat{q}_0 = 0.5 \text{ GeV}^2/\text{fm}$ - fitted to $e^+ + \text{Kr} \rightarrow \pi^+ + X$

Results – $E_{\text{lab}} = 27 \text{ GeV}$



- $\lambda(\pi) > 0$: Formation-time scaling for pions!
- Why $\lambda_{\text{best}}(h^\pm) \sim 0$ on Kr?
- proton anomaly!

Results – $E_{\text{lab}} = 12 \text{ GeV}$



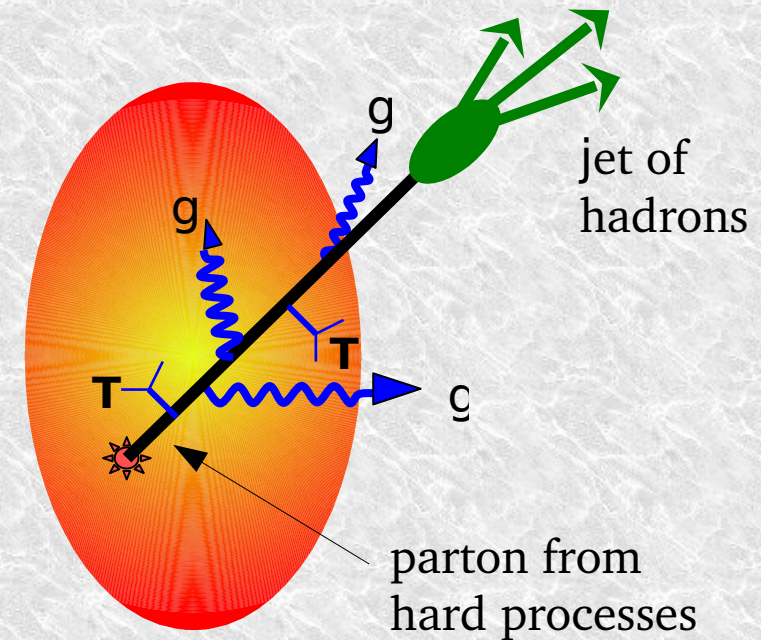
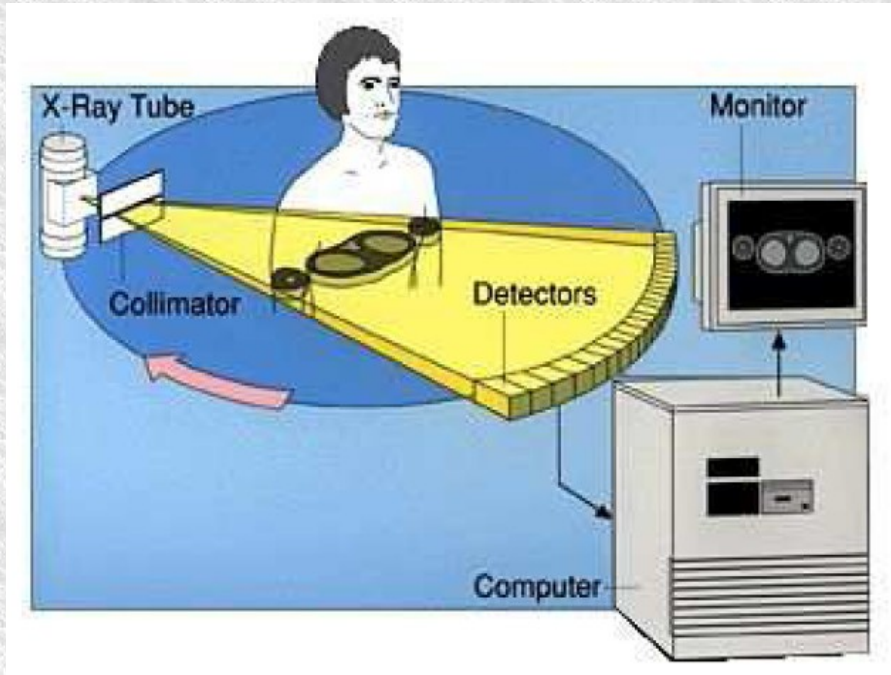
- pions are still positive! confirms results at 27 GeV
- $\lambda_{\text{best}}(h^+) \sim 0$ but $\lambda_{\text{best}}(h^-) > 0$
- proton anomaly hypothesis confirmed!

The energy loss paradigm

Review: Gyulassy, Vitev, Wang, Zhang, nucl-th/0302077

Baier, Dokshitzer, Levai, Mueller, Peigne, Schiff, Wiedemann, Zakharov, ...

- Jet tomography: QCD analog of Computed Axial Tomography (CAT)



Computed Axial Tomography

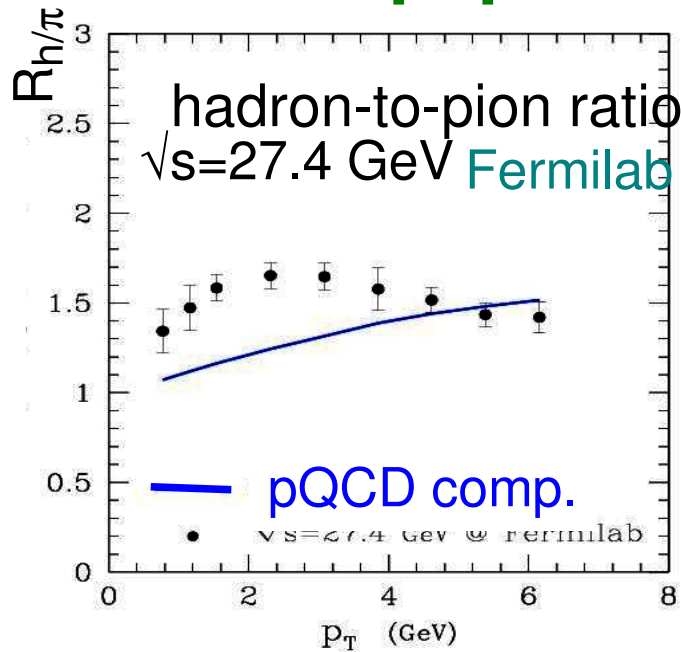
- Calibrated x-ray source
- x-ray absorption
- properties of the medium

Single hadron tomography

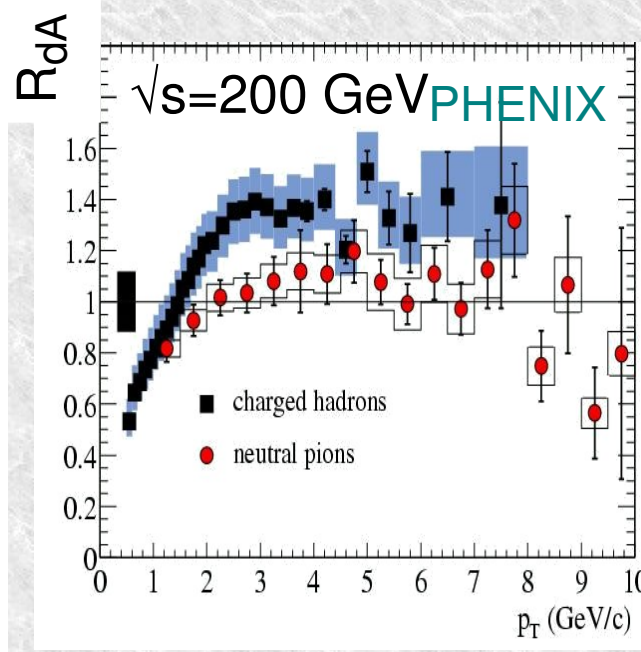
- Calibrated hard partons source
- energy loss (gluon bremsstrahlung)
computed in pQCD
- properties of the medium

Challenges: baryon anomaly

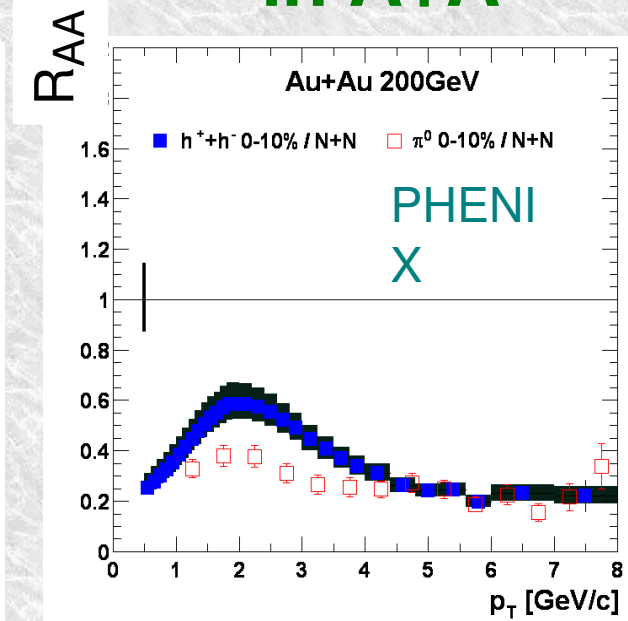
in p+p



in h+A



in A+A



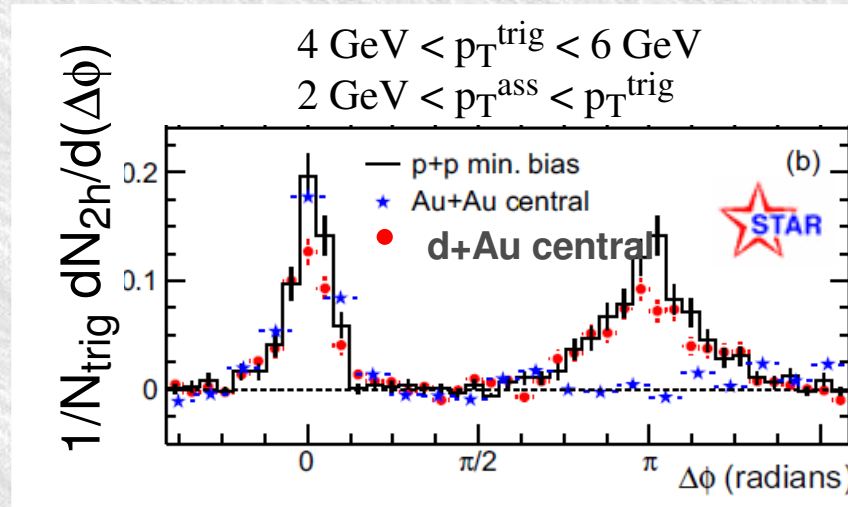
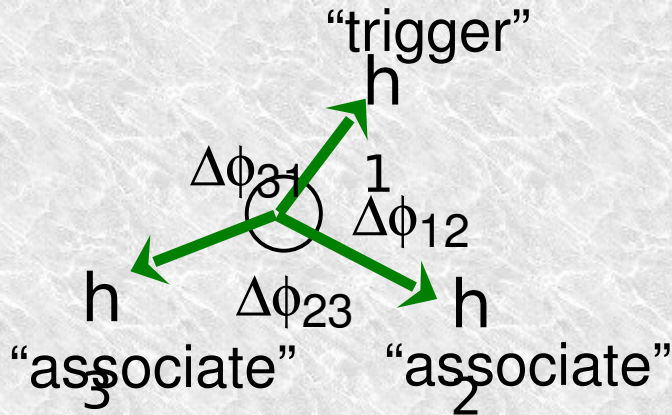
$$R_{h/\pi} = \frac{dN^h/d^2p_T}{dN^\pi/d^2p_T}$$

$$R_{BA} = \frac{1}{T_{BA}(b)} \frac{(dN^h/d^2p_T)_{d+A}}{(d\sigma^h/d^2p_T)_{p+p}}$$

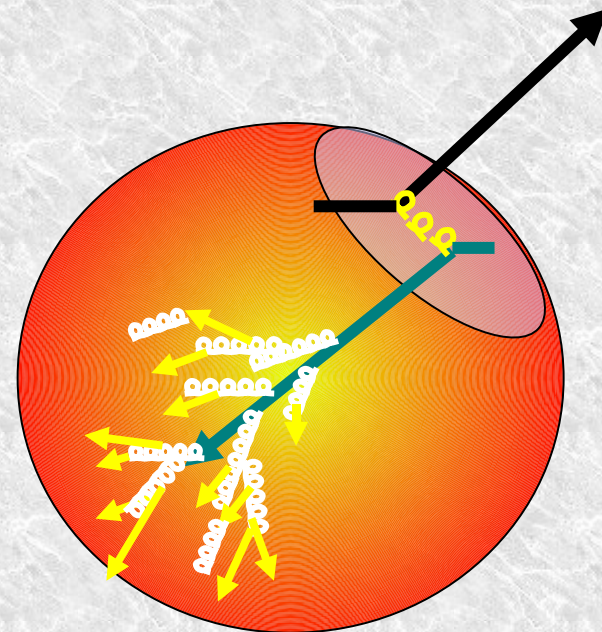
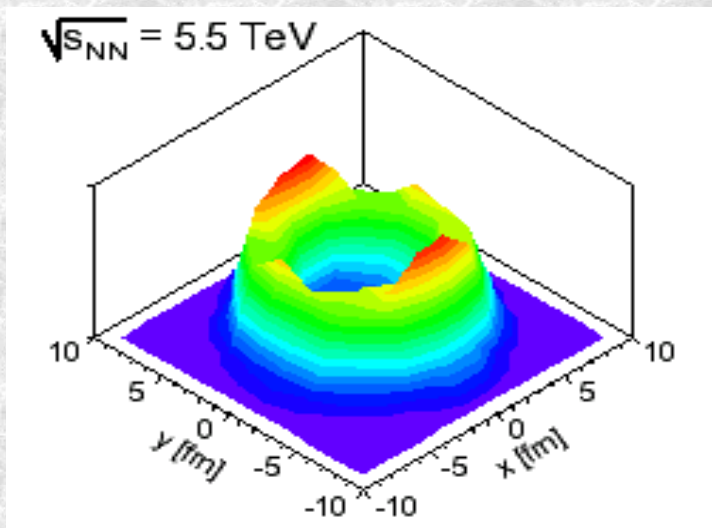
- p+p, p+A: difference in meson and baryon production at medium p_T not understood in pQCD + quark fragment.
- A+A : baryon anomaly persists – h/ π ratio increases

Challenges 3: surface emission

- Disappearance of the back-to-back jet

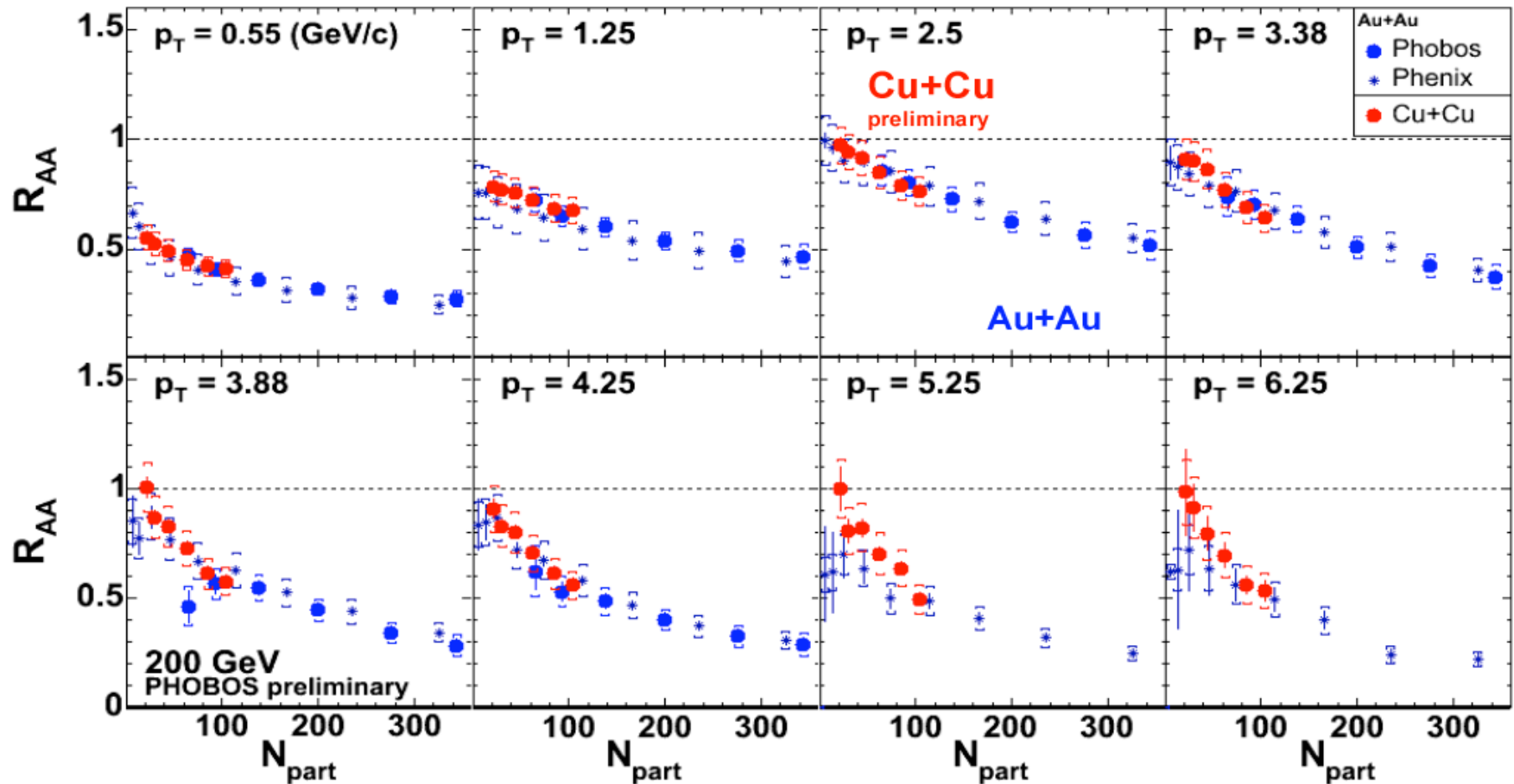


- Montecarlo simulations [Dainese et al.]



Challenges 3: surface emission

Yields vs N_{part} , 200 GeV



$$N_{part}^{CuCu} (0-3\%) \sim N_{part}^{AuAu} (35-45\%) \sim 100$$

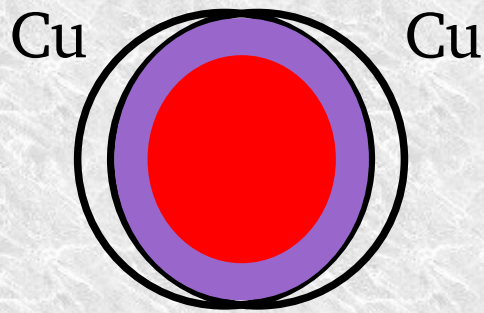
Au+Au: PRL 94, 082304 (2005), PLB 578, 297 (2004)

See poster by Gerrit van Nieuwenhuizen

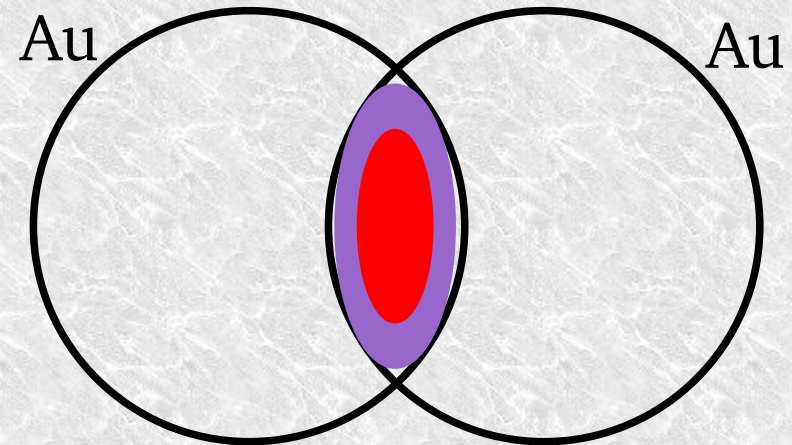
Challenges 3: surface emission

➔ But... Cu+Cu should quench more than Au+Au!

$$(S/V)_{\text{CuCu}} < (S/V)_{\text{AuAu}}$$



$$N_{\text{part}}^{\text{CuCu}} (0-3\%) \sim 100$$

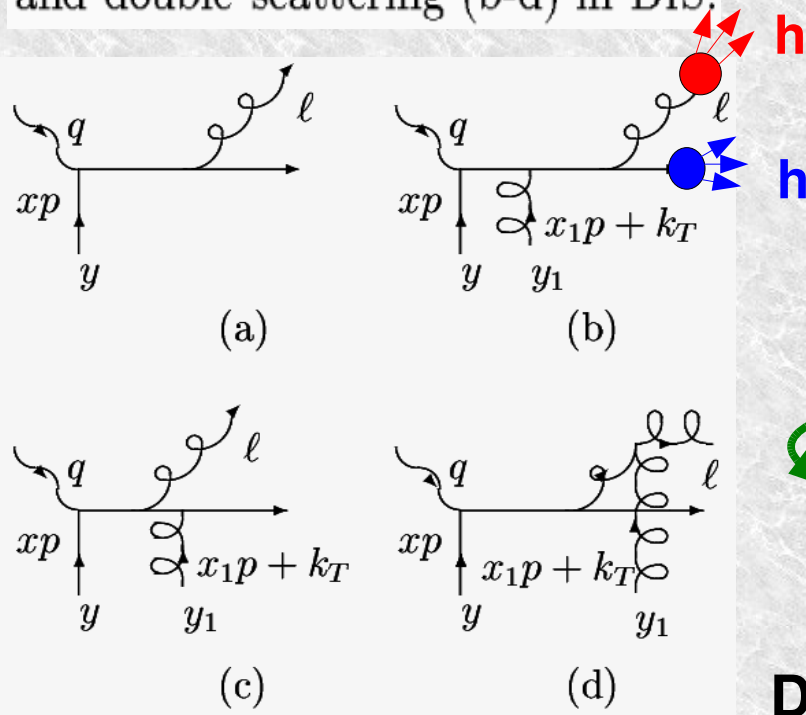


$$N_{\text{part}}^{\text{AuAu}} (35-45\%) \sim 100$$

Twist-4 effects: quark rescattering and gluon bremsstrahlung

Guo-Wang '00, Wang-Wang '02

Gluon radiation from a single scattering (a) and double scattering (b-d) in DIS.



- Soft gluon emission
- quark rescattering + induced radiation

$$\tilde{D}_{q \rightarrow h}(z_h, Q^2) \equiv D_{q \rightarrow h}(z_h, Q^2) + \Delta D_{q \rightarrow h}(z_h, Q^2)$$

LO

NLO + twist-4

$$= \int_0^{Q^2} \frac{d\ell_T^2}{\ell_T^2} \frac{\alpha_s}{2\pi} \int_{z_h}^1 \frac{dz}{z} [\Delta\gamma_{q \rightarrow qg}(z, x, x_L, \ell_T^2) D_{q \rightarrow h}(z_h/z) + \Delta\gamma_{q \rightarrow gq}(z, x, x_L, \ell_T^2) D_{g \rightarrow h}(z_h/z)]$$

Def: quark fractional energy loss

$\langle Z_g \rangle$ = average energy carried by radiated gluon

- Result effectively approximated by:

$$\tilde{D}_{a \rightarrow h}(z) \approx \frac{1}{1 - \Delta z} D_{a \rightarrow h} \left(\frac{z}{1 - \Delta z} \right)$$

= 0.006 @ $Q^2 = 3 \text{ GeV}^2$ fitted to HERM
0.013 @ $Q^2 = 40 \text{ GeV}^2$ fitted to DY

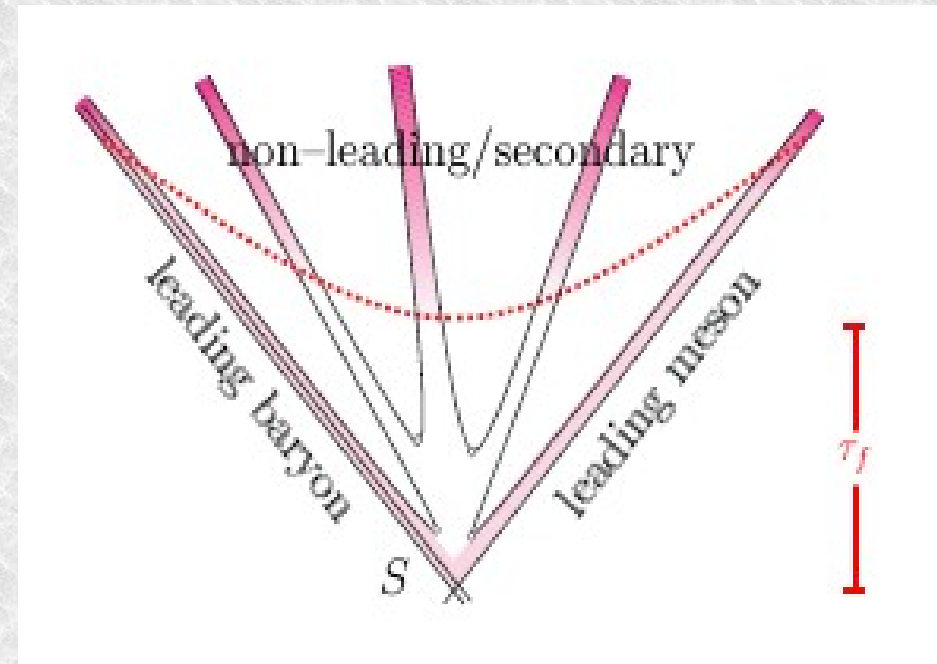
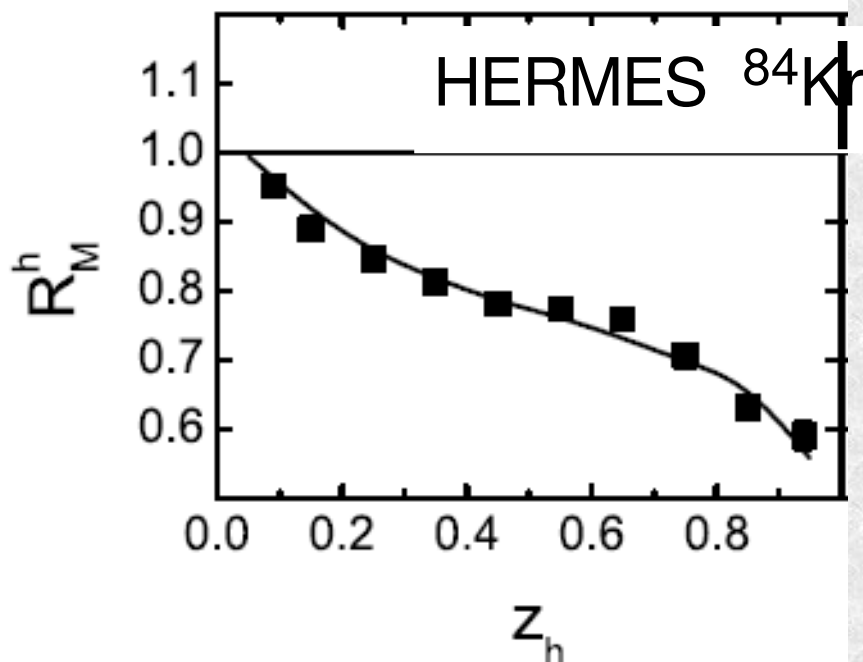
Where
:

$$\Delta z|_{\text{Wang}} = 0.6 \langle z_g \rangle \propto A^{2/3} \alpha_s^2(Q^2) \tilde{C}(Q^2) \frac{1}{\nu} \ln \left(\frac{m_N \nu}{Q^2} \right)$$

An absorption+transport model for Au+Au

Falter et al., PRC 70(04)054609

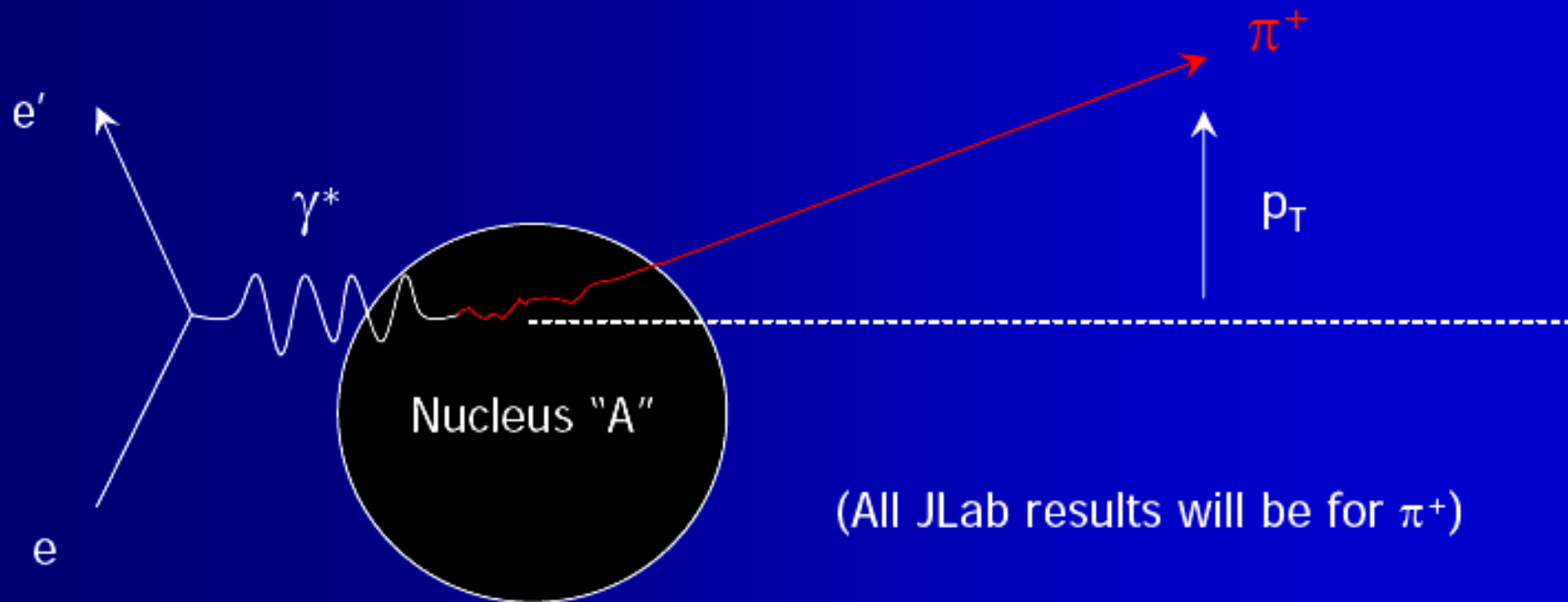
- ★ Formation times: $t_* = 0$ fm and $t_h = (E_h/m_h) \tau_F$ with $\tau_F = 0.5$ fm
- Cross sections: leading h: $\sigma_* = 1/2 \sigma_h$ (mesons) $1/3 \sigma_h$ (barions)
- subleading h: $\sigma_* = 0$ mbarn
- ★ Final state X
 - by PYTHIA and FRITJOF
 - Fermi motion, Pauli blocking, shadowing
 - evolved by BUU transport equations



Observable Number 1 – “ p_T Broadening”

Definition of “Transverse Momentum Broadening”

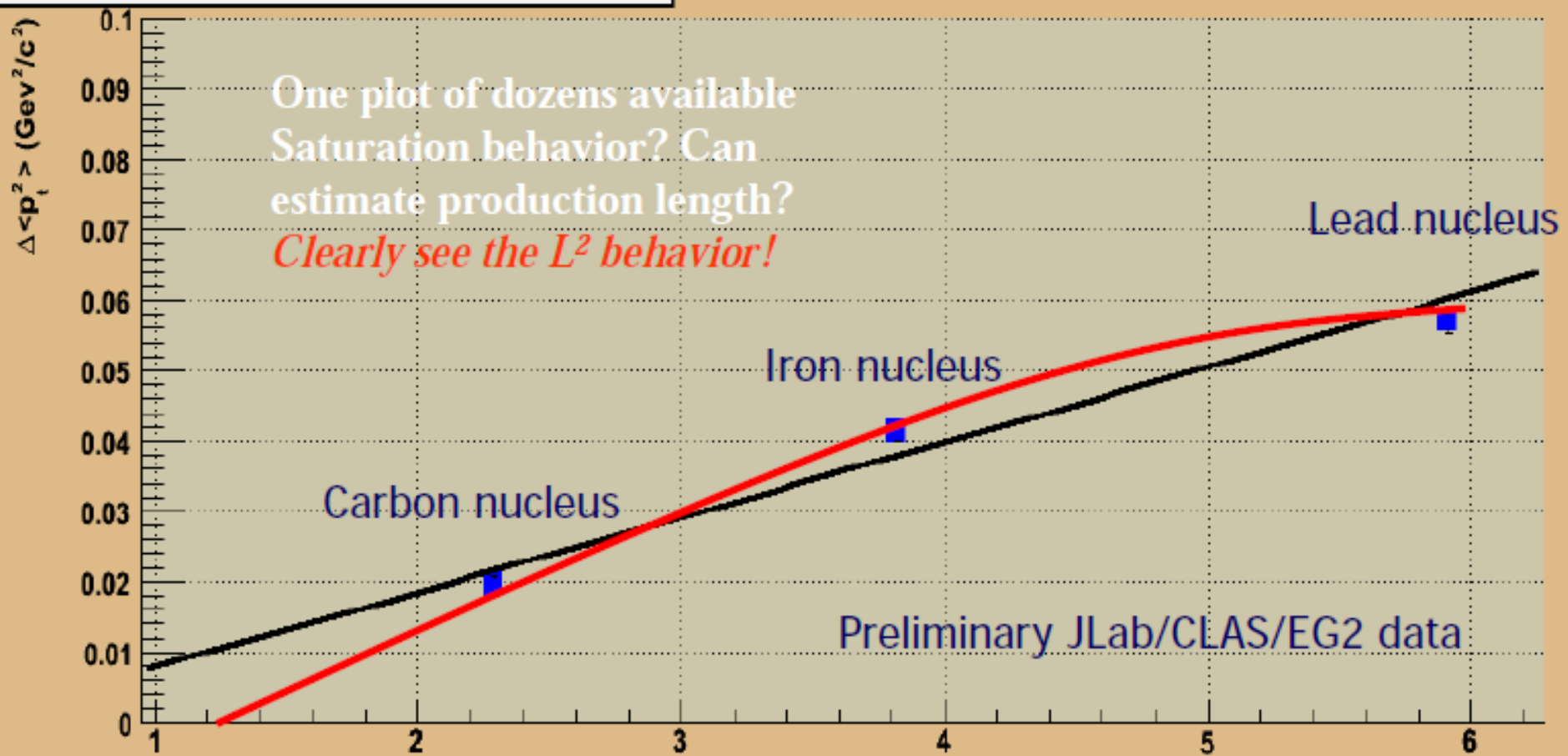
$$\Delta(p_T^2) = p_T^2(A) - p_T^2(^2H)$$



p_T broadening: “A” dependence, large v , mid- z , π^+

- Transverse momentum broadening vs. nuclear radius is \sim linear

$1 < Q^2 < 2$ $3 < v < 4$ $0.500000 < Z_{\pi^+} < 0.600000$ | π^+



Medium-Induced Quark Energy Loss

assuming perturbative formula

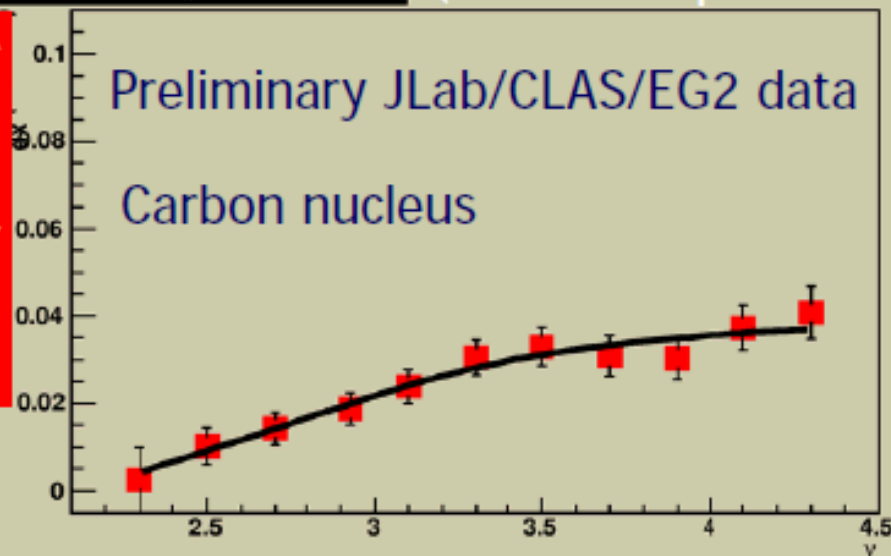
$$dE/dx \approx \frac{\alpha_s}{\pi} N_c \langle p_T^2 \rangle_L$$

$1 < Q^2 < 2 \quad 0.5 < Z_{\pi^+} < 0.6 \quad | \quad \pi^+ \quad | \quad C$

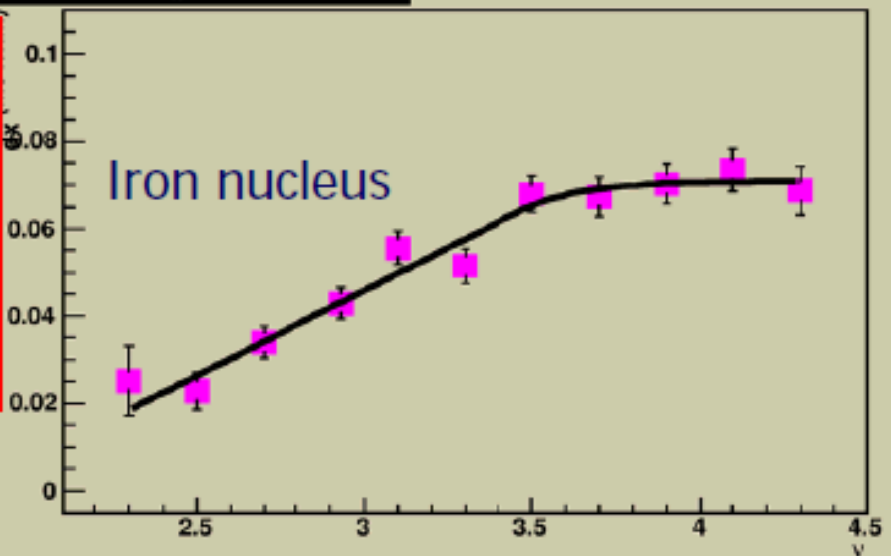
(Dozens of plots like this)

$1 < Q^2 < 2 \quad 0.5 < Z_{\pi^+} < 0.6 \quad | \quad \pi^+ \quad | \quad Fe$

dE/dx (GeV/fm)



dE/dx (GeV/fm)

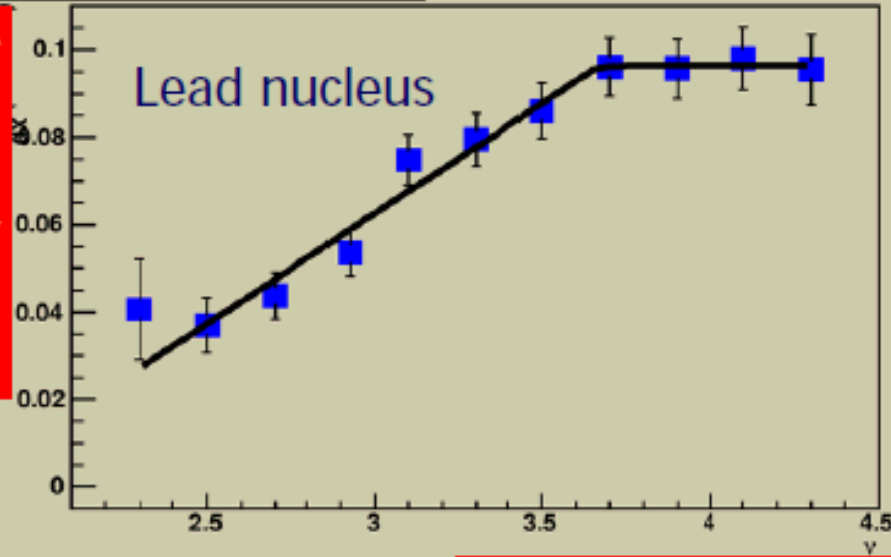


Struck quark energy (GeV)

Struck quark energy (GeV)

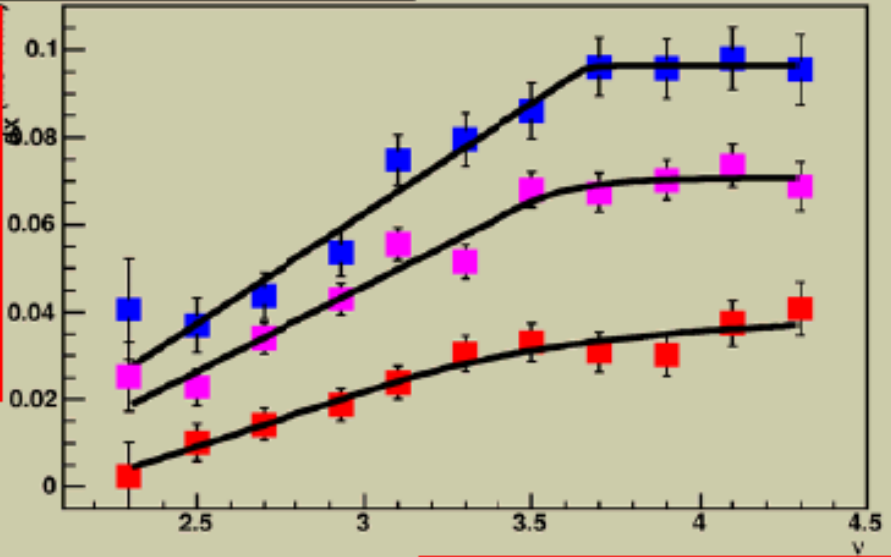
$1 < Q^2 < 2 \quad 0.5 < Z_{\pi^+} < 0.6 \quad | \quad \pi^+ \quad | \quad Pb$

dE/dx (GeV/fm)



dE/dx (GeV/fm)

$1 < Q^2 < 2 \quad 0.5 < Z_{\pi^+} < 0.6 \quad | \quad \pi^+ \quad |$



Struck quark energy (GeV)

Struck quark energy (GeV)