

Kaon electromagnetic production:

constraints set by new data

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INTRODUCTION

MODELS

CONCLUSION

INTRODUCTION

- $\gamma + p \rightarrow K^+ + \Lambda$
- good description of the elementary process
→ hypernuclear calculations [T. Motoiba]
- analysis of the models: Adelseck-Benkhold,
Adelseck-Wright, Adelseck-Saghai,
Williams-Si-Coland, Sadag-Lyon, Kaou-MAID,
and Regge in view of "new" data:
Mohring et al [E93-018, Phys. Rev. C67(2003)055205]
and Trau et al [SAPHIR(98), Phys. Lett. B445(1998)20]
can be found in nucl-th/0305039
- latest data on $\frac{d\Gamma}{d\Omega}$ and P for the photo-production
 - CLAS (J.W.C. McNabb et al) [nucl-ex/0305028]
 - SAPHIR (K.-H. Glauder et al) [nucl-ex/0308025]
 - inconsistent at small θ_K
- learn more about the reaction mechanism
Are the models able to conform to the data?

Isobaric models:

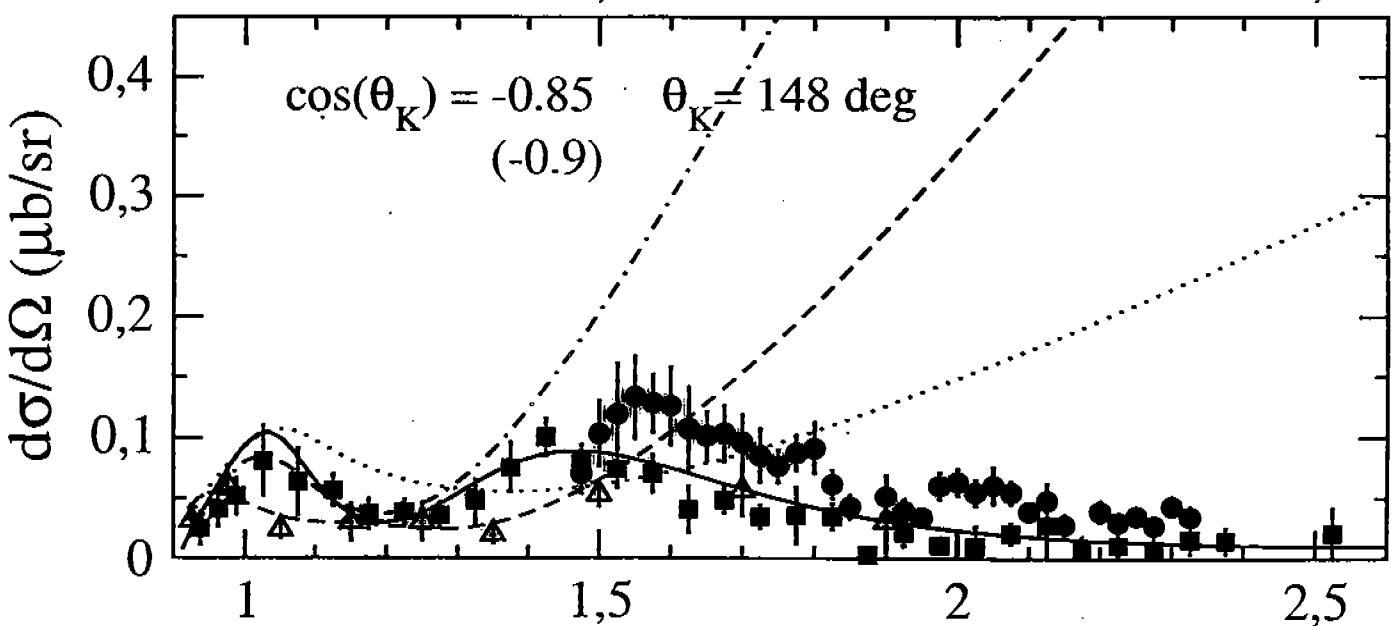
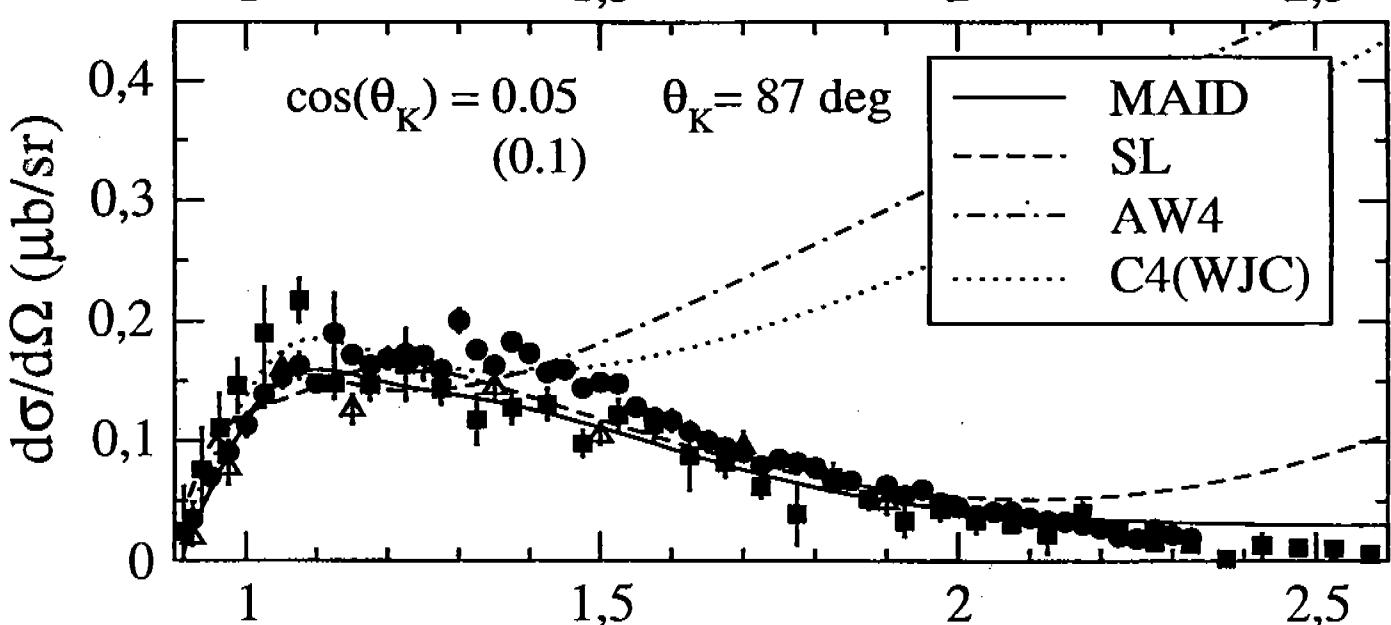
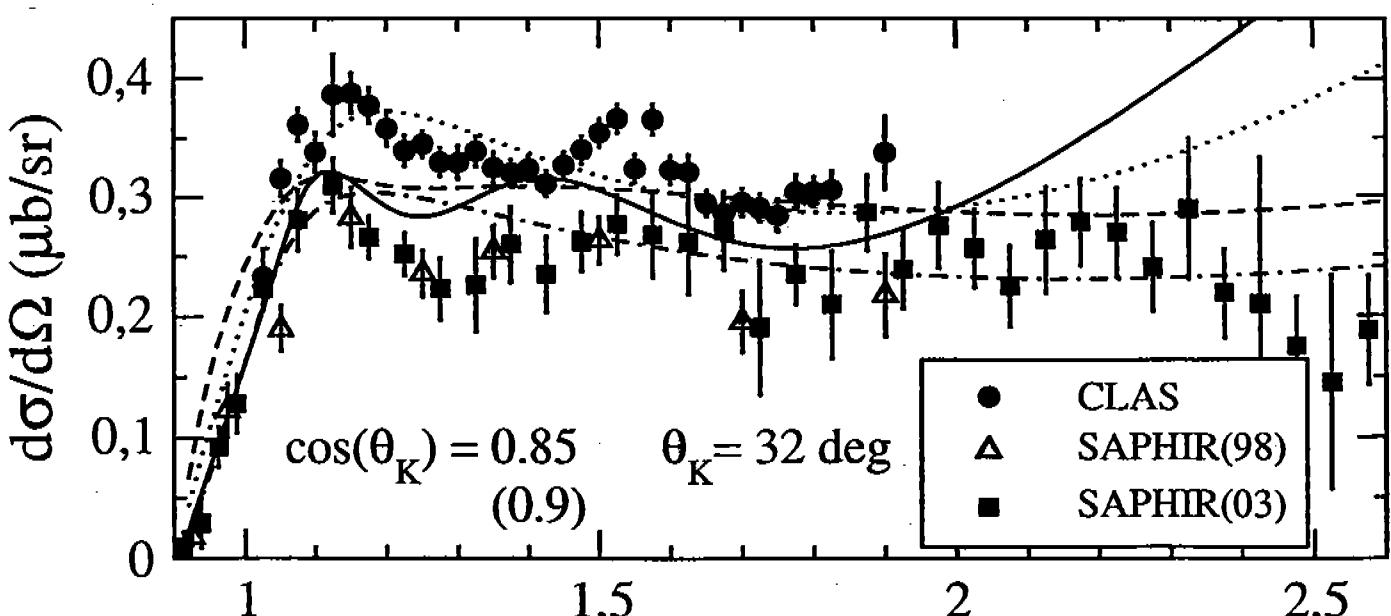
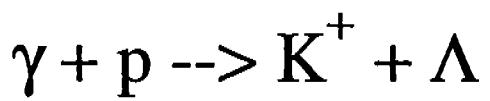
- tree-level effective Lagrangian approach
 - too many resonances: 20 - 30 [C. Bennhold]
 - SU(3) symmetry limits: $-4.4 \leq g_{K\Lambda}/\sqrt{4\pi} \leq -3.0$
 $0.8 \leq g_{K\Sigma}/\sqrt{4\pi} \leq 1.3$
 - crossing symmetry: $R_{JM} = \Gamma(K\bar{p} \rightarrow J/\psi) / \Gamma(K\bar{p} \rightarrow \text{all})$
- • structure of hadrons: hadronic form factors
X gauge invariance

→ method by Haberzettl; Davidson and Workman
[Phys. Rev. C63 (2001) 025210]

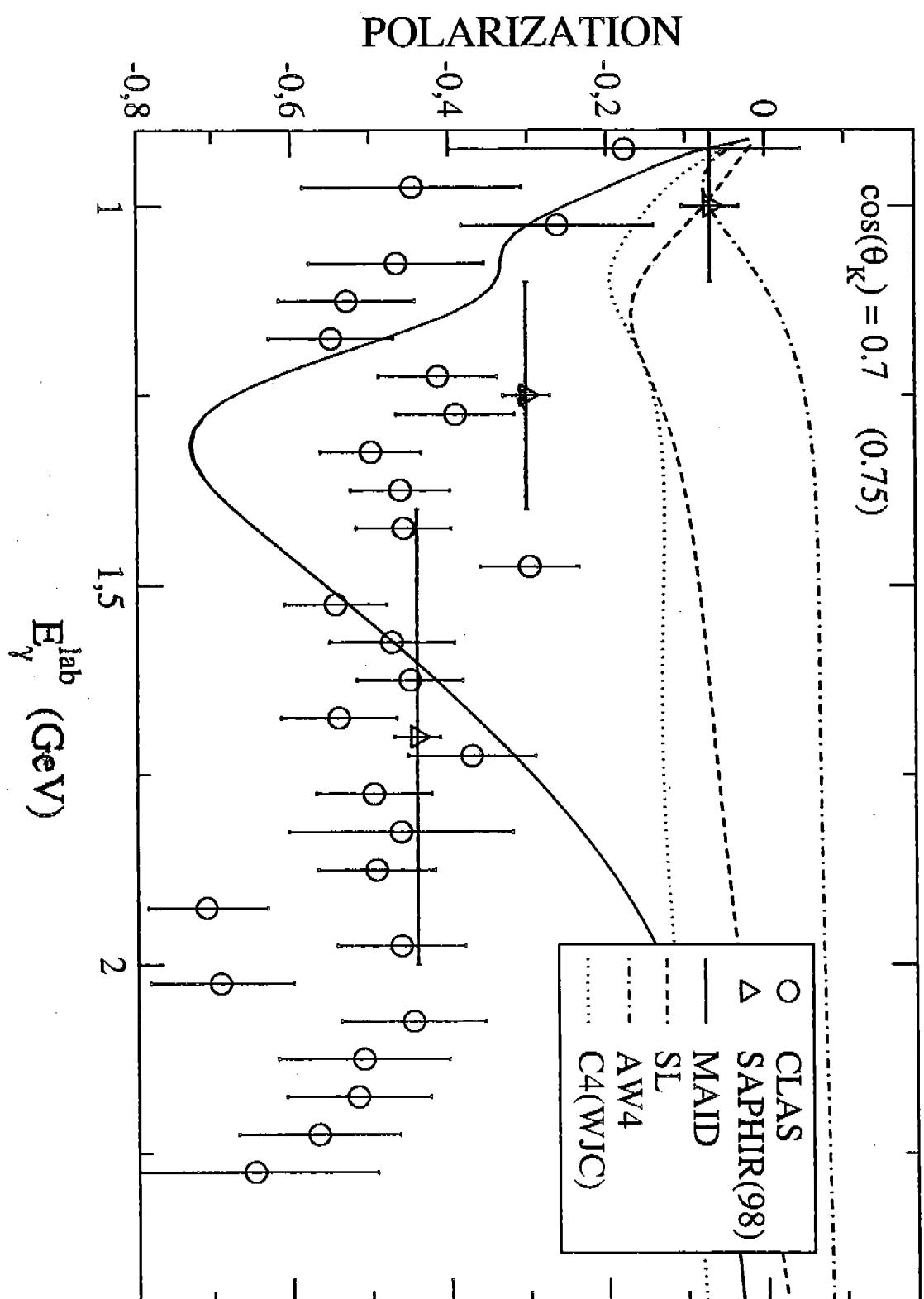
contact term: $F(s, u, t) = F_1(s) + F_2(u) + F_3(t) -$
 $- F_1(s)F_2(u) - F_1(s)F_3(t) - F_2(u)F_3(t) + F_1(s)F_2(u)F_3(t)$

with

$$F_i(x) = \frac{\Lambda^4}{\Lambda^4 + (x - m^2)^2} \quad (x = s, u, t)$$



$\gamma + p \rightarrow K^+ + \Lambda$



MODELS

Saclay-Lyon (SL)

[David et al, Phys. Rev. C53(1996)2613]

- no hadronic form factors

- resonances

(SL): $K^*, K_1, N(1440)(P_{11}), N(1720)(P_{13}), N(1675)(D_{15}),$

$\Lambda(1405)(S_{01}), \Lambda(1670)(S_{01}), \Lambda(1810)(P_{01}), \Sigma(1660)(P_{11})$

(SLA): $K^*, K_1, N(1720)(P_{13}),$

$\Lambda(1405)(S_{01}), \Lambda(1670)(S_{01}), \Lambda(1810)(P_{01}), \Sigma(1660)(P_{11})$

- coupling constants are re-fitted to $d\sigma/d\Omega$

CLAS: SL(CL) (918 data points)

SAPHIR(03): SL(S03) (701 data points)

with $SU(3)$ limits for $g_{K\Lambda N}$ and $g_{K\Sigma N}$

SL(CL)	SL(S03)	SLA(CL)	SLA(S03)
$\chi^2: 273 \rightarrow 2.23$	$171 \rightarrow 3.23$	$281 \rightarrow 3.65$	$171 \rightarrow 5.74$

- moderate change of the coupling constants with the exception of:

SL(CL): $N(1440)$, $N(1675) - 0.02 \rightarrow 2.40 \pm 0.8$

SLA(CL): $N(1720)$

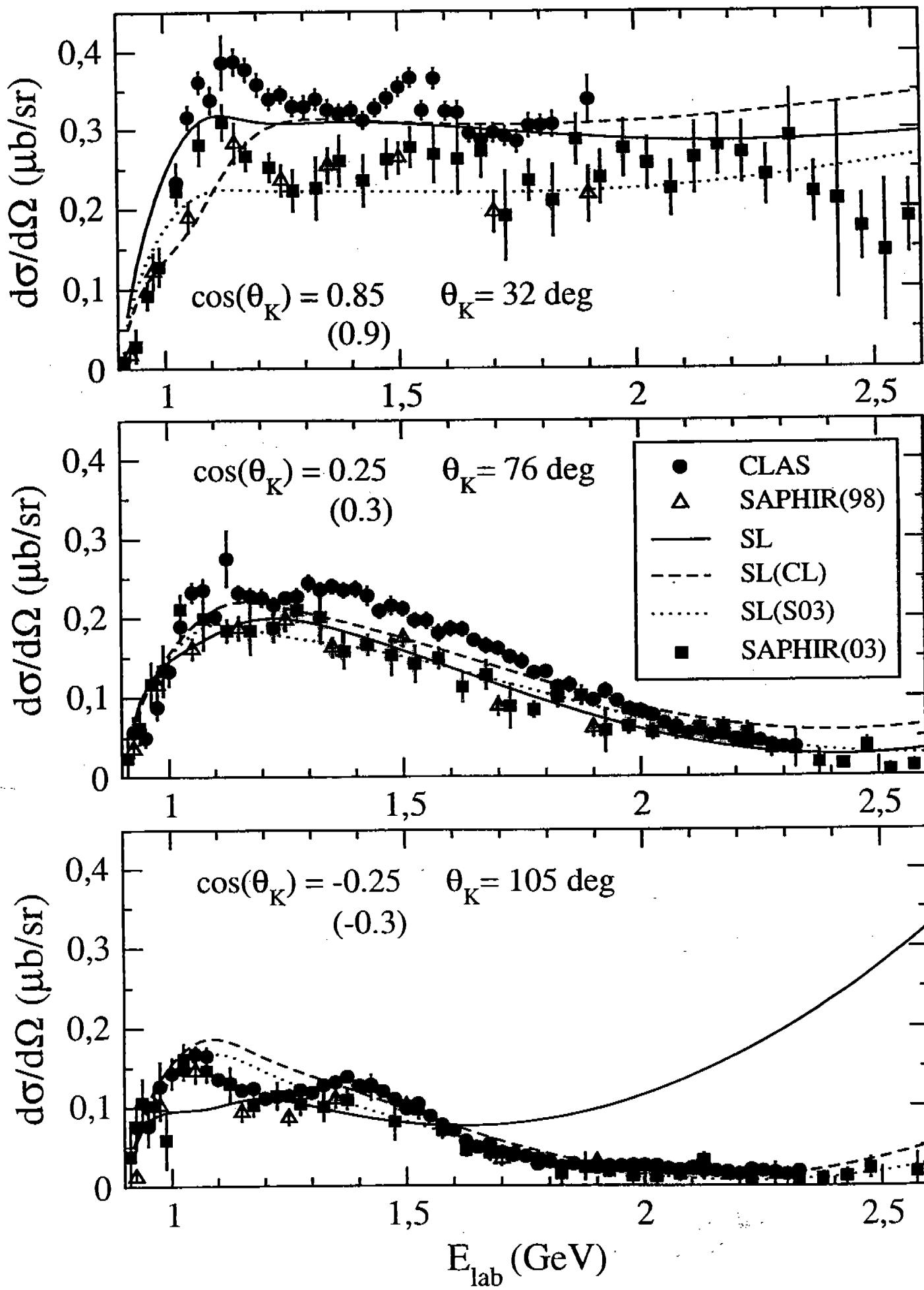
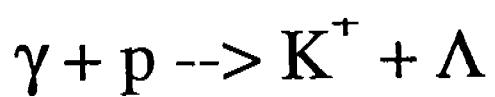
SL(S03): $N(1440), \underline{\Lambda(1810)} - 1.96 \rightarrow 6.26 \pm 15.$!

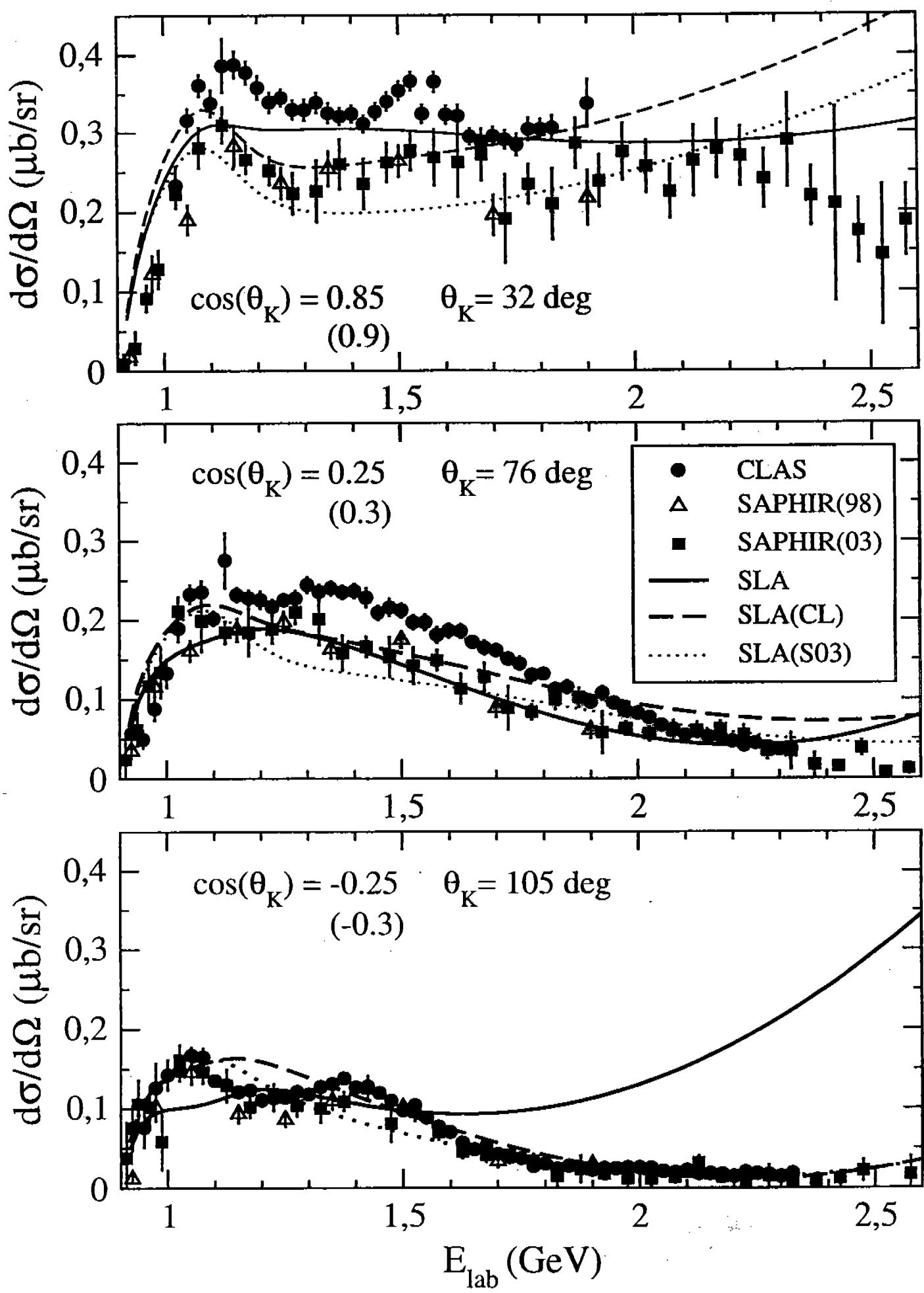
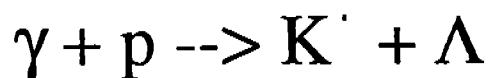
SLA(S03): $\Lambda(1405), \Lambda(1670), \Lambda(1810)$

- big errors for the hyperon resonances

→ the u-channel contributions are not well determined

- large non resonant contribution





Saclay-Lyon type model

- $K^*, K_1, N(1440)(P_{11}), N(1720)(P_{13}),$
 $\Lambda(1670)(S_{01}), \Lambda(1810)(P_{01})$
- *hadronic form factors* of Davidson and Workman (Λ_1)
- g 's and Λ_1 are fitted to CLAS data ($d\sigma/d\Omega$)
- SU(3) limits for $g_{K\Lambda N}$ and $g_{K\Sigma N}$; $0.6 \leq \Lambda_1 \leq 2.0$

results in $\boxed{\chi^2 = 2.08, \quad \Lambda_1 = 0.96 \pm 0.3 \text{ GeV}}$

big errors of g 's for

$$N(1440): -0.64 \pm \underline{5.2}$$

$$\Lambda(1670): 3.32 \pm \underline{7.6}$$

$$\Lambda(1810): -\underline{10.0} \pm \underline{18.7}$$

other choice of hyperon resonances leaves the uncertainty of g 's big

$\Lambda(1810) \rightarrow \Lambda(1600)$: $\boxed{\chi^2 = 1.66, \quad \Lambda_1 = 1.1 \pm 0.3 \text{ GeV}}$

$$\Lambda(1670): 5.79 \pm \underline{16.5}$$

$$\Lambda(1600): -\underline{10.0} \pm \underline{14.9}$$

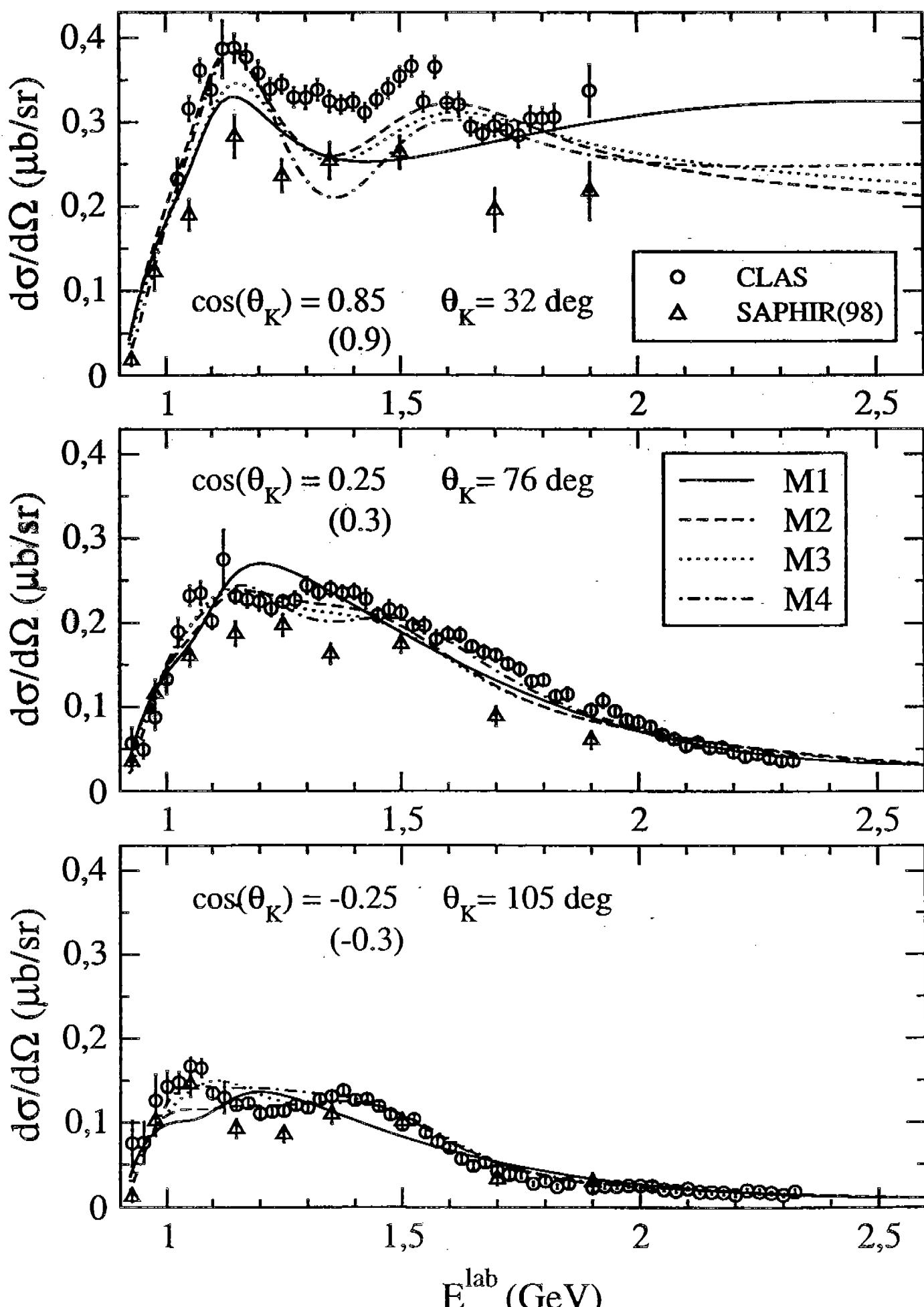
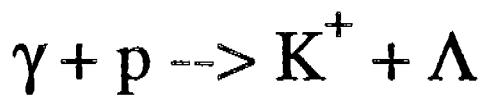
or
Kaon-MAID type model (Janssen et al.)

- *hadronic form factors* of Davidson and Workman (Λ_1)
- g 's and Λ_1 are fitted to CLAS data ($d\sigma/d\Omega$)
- SU(3) limits for $g_{K\Lambda N}$ and $g_{K\Sigma N}$; $0.6 \leq \Lambda_1 \leq 2.0$

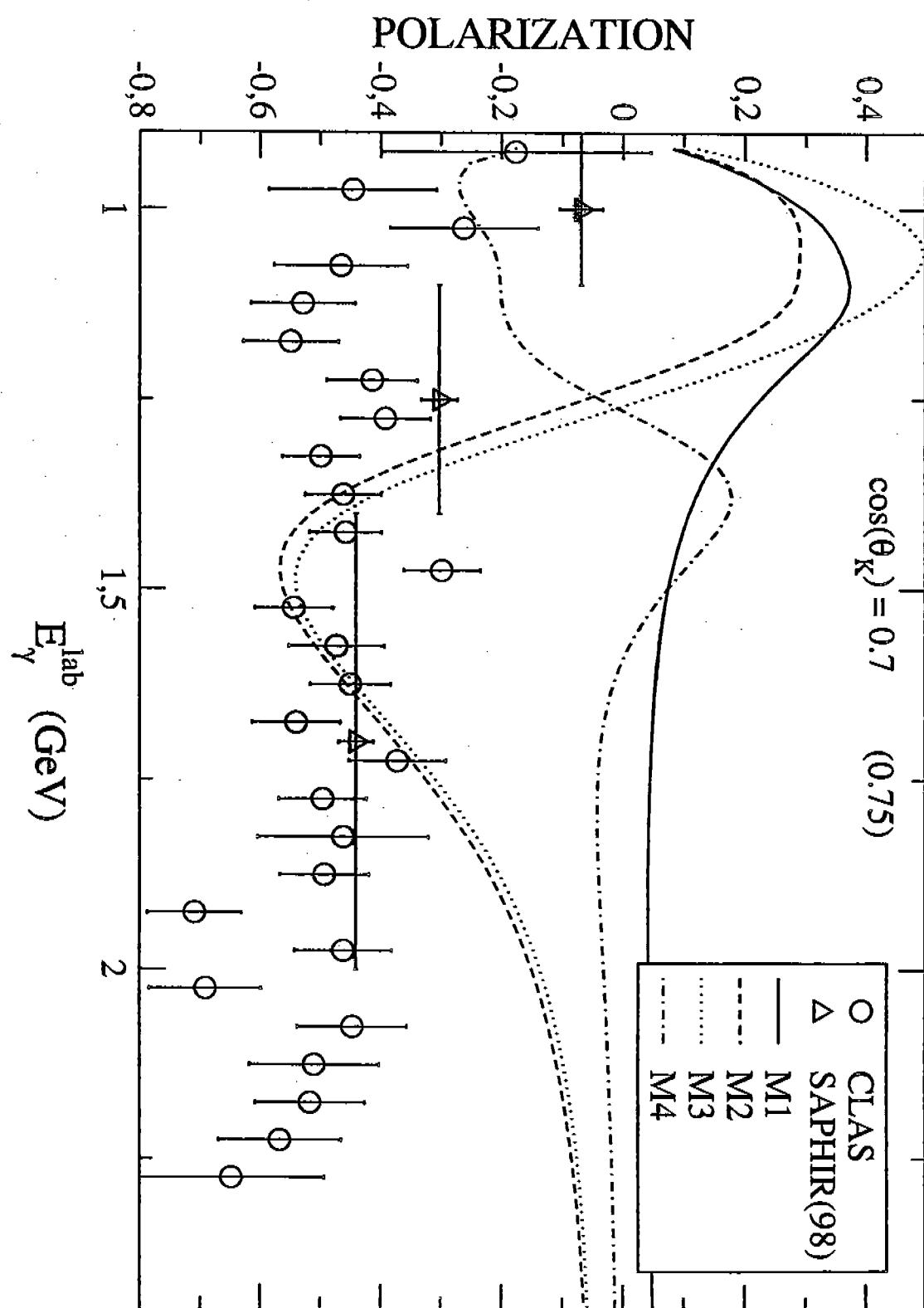
Coupling constants:

Model:	M1	M2	M3	M4
K^*	0.14 (0.6)	-0.01 (0.4)	-0.02 (0.6)	-0.20 (0.4)
	0.50 (1.0)	0.27 (0.9)	0.24 (0.7)	0.21 (0.8)
K_1	-0.38 (0.7)	0.28 (<u>1.3</u>)	0.30 (<u>1.7</u>)	—
	-2.14 (<u>1.9</u>)	-0.64 (<u>1.8</u>)	-0.69 (<u>2.0</u>)	—
$N^*(1650)$	0.12 (0.1)	0.14 (0.1)	0.13 (0.1)	-0.14 (0.1)
$N^*(1710)$	0.21 (1.0)	0.25 (<u>1.2</u>)	—	—
$N^*(1720)$	-0.11 (0.2)	-0.10 (0.2)	-0.13 (0.1)	0.16 (0.1)
	-0.55 (1.0)	-0.63 (<u>1.3</u>)	-0.88 (0.6)	0.92 (0.5)
$D_{13}(1895)$	—	-0.32 (0.8)	-0.22 (0.9)	-0.79 (0.6)
	—	-0.31 (0.5)	-0.25 (0.5)	-0.58 (0.3)
Λ_1 (GeV)	0.89 (0.2)	0.88 (0.2)	0.87 (0.2)	0.80 (0.1)
χ^2	1.37	0.83	0.92	0.99

Addition of $\Lambda^*(1600)$ and $\Lambda^*(1670)$ into the M1 model
 results in $\boxed{\chi^2 = 1.25 \text{ and larger } \Lambda_1: 1.00 \pm 0.06 \text{ GeV}}$
 but $g_{\Lambda(1670)} = 1.94 \pm \underline{3.2}$ and $g_{\Lambda(1600)} = \underline{-10.0 \pm 16.0}$!



$\gamma + p \rightarrow K^+ + \Lambda$



CONCLUSIONS (preliminary)

- "old" models (SL) being re-fitted can describe a global behaviour of $\frac{d\sigma}{du}$ up to 2.4 GeV (also for large θ_k) but cannot describe details of the data
- contributions from hyperon resonances are not well determined with the photo-production $\frac{d\sigma}{du}$
- the "missing" resonance $D_{13}(1895)$ is important for description of details of the structure of $\frac{d\sigma}{du}$; important also for polarizations