

Electroproduction of Strangeness on $^{3,4}\text{He}$

Ê Motivation

Ë $^1\text{H}(e, e'K)L$

Ì $A(e, e'K), A=2,3,4$

Í Bound Λ Hypernuclei: $^3_{\Lambda}\text{H}, ^4_{\Lambda}\text{H}$

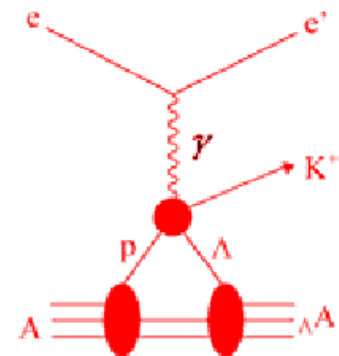
Î Outlook

HYP2003, Jefferson Lab

Frank Dohrmann, Forschungszentrum Rossendorf, Dresden

Motivation

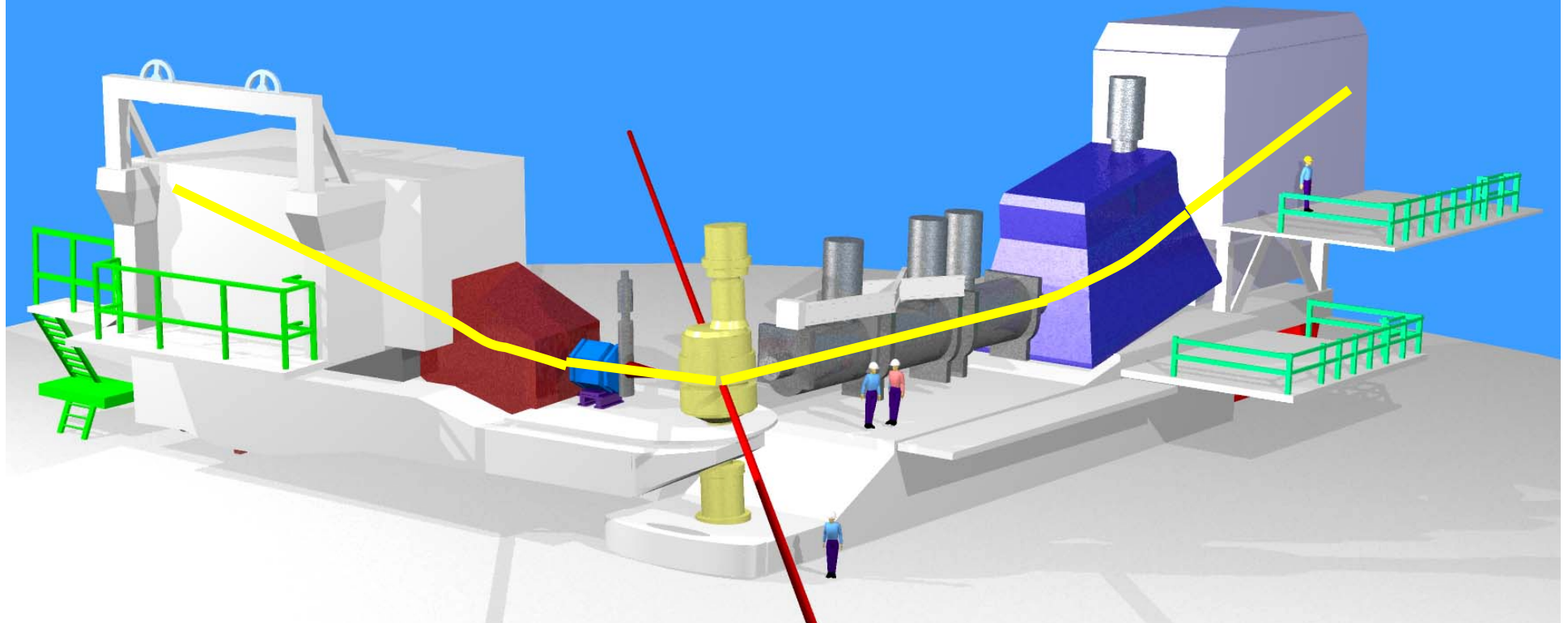
- ! Electroproduction of kaons is an efficient tool for investigating fundamental baryonic interactions in the nuclear medium
- ! New ${}^1\text{H}(e, e'K)Y$ data for new, improved models
- ! No data for electroproduction on ${}^3\text{He}$, ${}^4\text{He}$ exists
- ! $A(e, e'K)Y$ for ${}^1\text{H}$, ${}^2\text{H}$, ${}^3\text{He}$, ${}^4\text{He}$
 - ! A-dependence of kaon electroproduction cross section
- ! Bound hypernuclear states for $A=3,4$, è lightest hypernuclei



Hall C Spectrometers

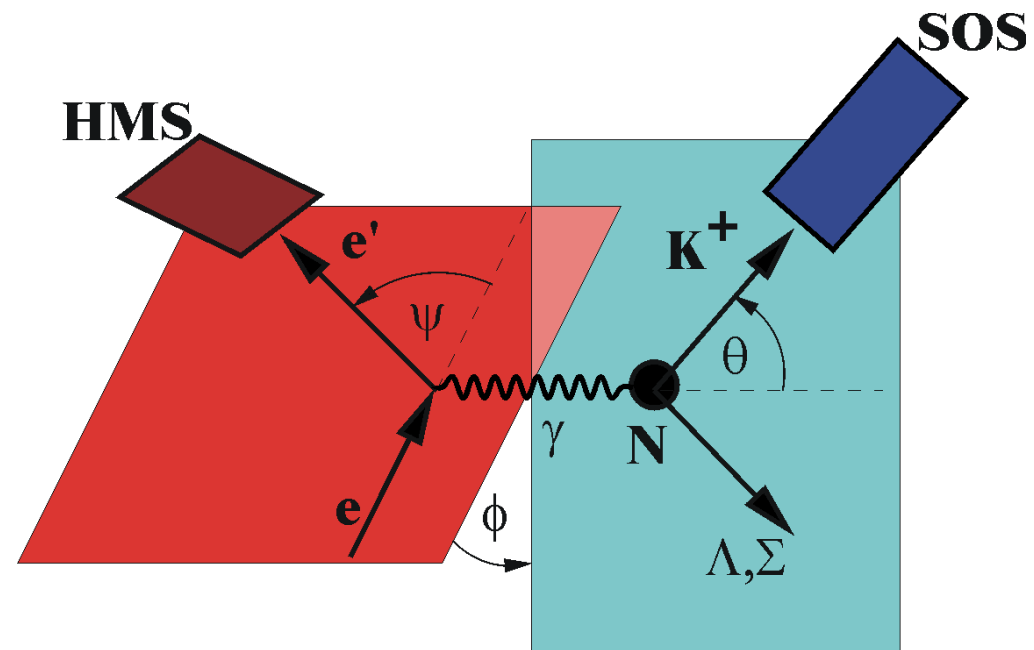
Short Orbit Spectrometer, SOS
kaon arm

High Momentum Spectrometer, HMS
electron arm



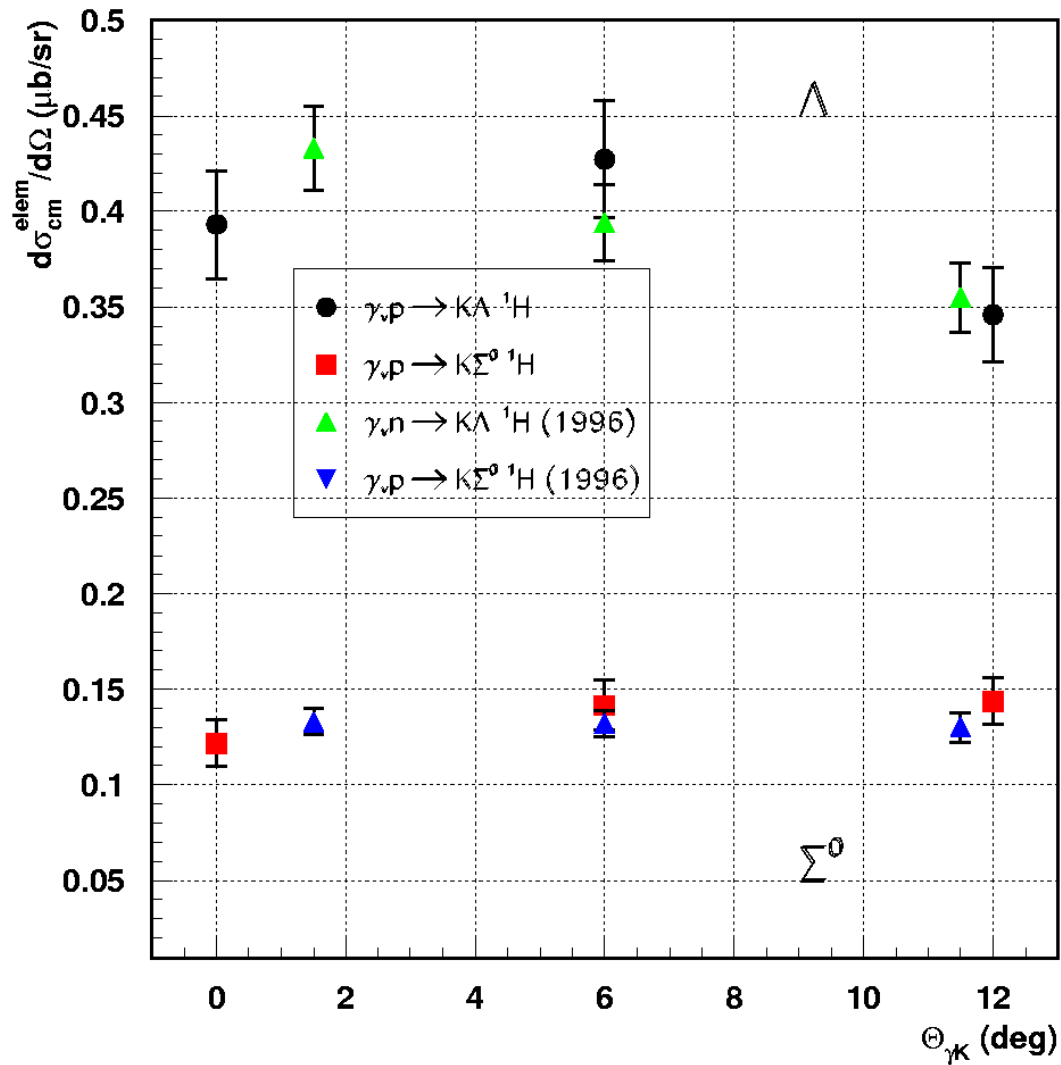
Electron beam

Experimental Program



| <u>Kinematics</u> | | |
|-------------------------------------|---|--|
| Targets 1996 | $^1\text{H}, ^2\text{H}, \text{C}, \text{Al}$ | |
| | $(Q^2(\text{GeV}^2), W(\text{GeV}))$ | $\theta_{\gamma K, \text{lab}} (^\circ)$ |
| $E_{\text{beam}}=3.245 \text{ GeV}$ | (0.38, 1.90) | (1.5, 6.0, 11.5) |
| Targets 1999 | $^1\text{H}, ^2\text{H}, ^3\text{He}, ^4\text{He}, \text{C}, \text{Al}$ | |
| | $(Q^2(\text{GeV}^2), W(\text{GeV}))$ | $\theta_{\gamma K, \text{lab}} (^\circ)$ |
| $E_{\text{beam}}=3.245 \text{ GeV}$ | (0.35, 1.91) | (0.0, 6.0, 12.0) |

H(e,e'K) Angular Distribution



Hydrogen Data '96 and '99

L Good agreement between independent measurements

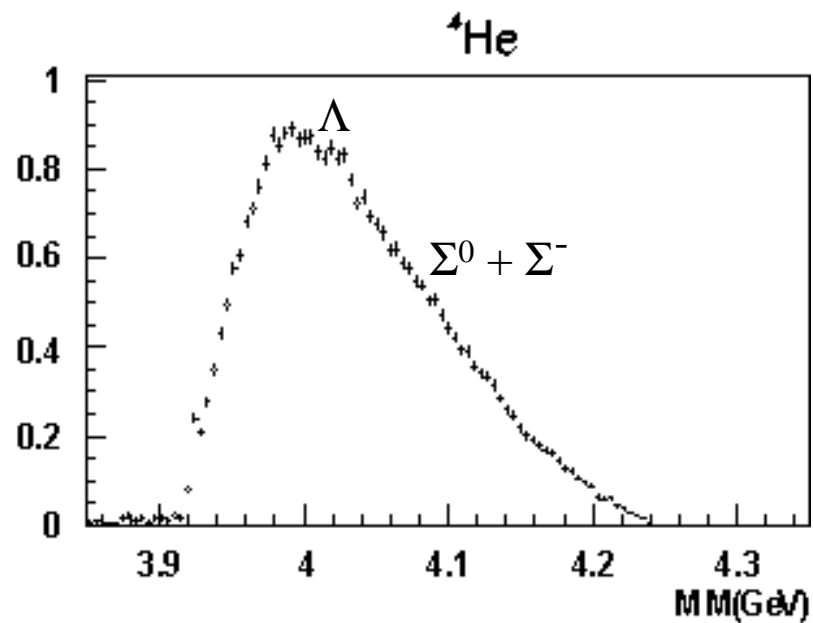
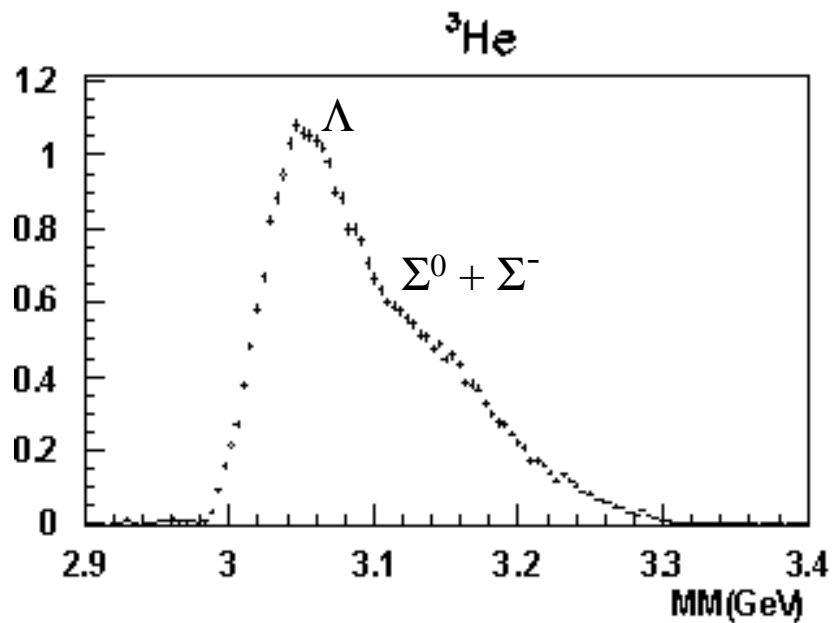
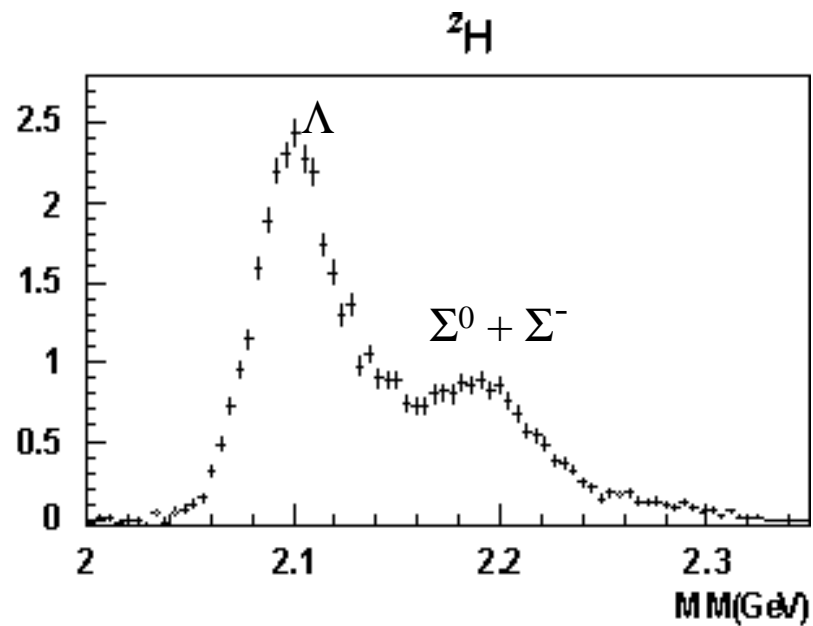
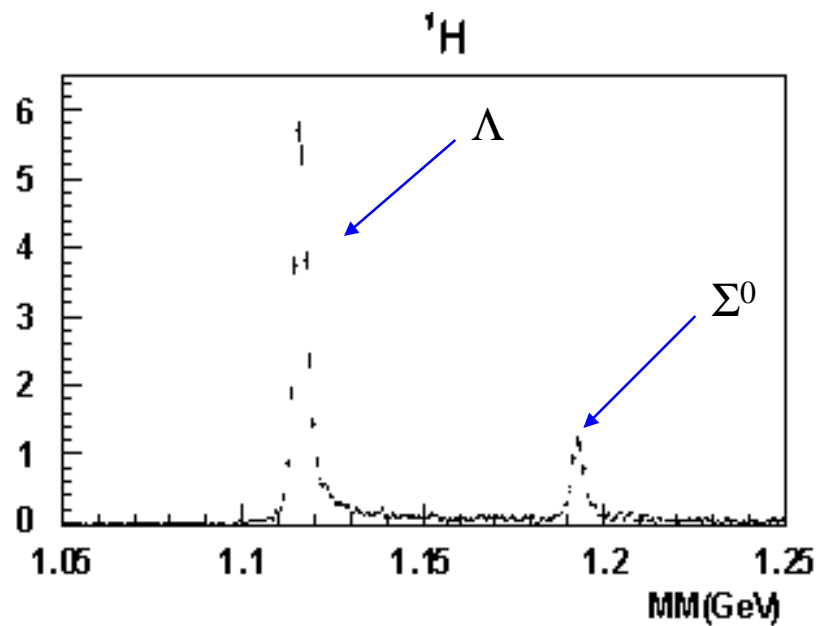
L He(e,e'K) data is used for

Oown model development

Onormalization

Ocalibration

L Data represents high statistics and high quality electro- production data on hydrogen.



Effective Range Ansatz

For $A(e,e'K)$, $A=2,3,4$:

Effective range Ansatz $k \cdot \cot d^0 = \frac{1}{a} + r_e \cdot \frac{k^2}{2} + O(k^4) \quad k = k_{\text{YN}}$

Calculate singlett and triplett s-wave enhancement factor I using inverse Jost Function

$$S = I \cdot S_0 \quad I = \frac{1}{J^* J} = \frac{k^2 + a^2}{k^2 + b^2} \quad J = \frac{k - ia}{k - ib}$$

$$\text{with } 0.5 \cdot r_e \cdot (b - a) = 1 \quad , \quad 0.5 \cdot r_e \cdot a \cdot b = \frac{1}{a}$$

Scattering length a and effective range r_e are given by hypernuclear potential such as Nijmegen 97f or Jülich A

Simulation of Production off Nuclear Targets (SIMC)

- ∅ Beam energy fixed (allow smearing)
- ∅ e' quantities and K⁺ angles are thrown flat
- ∅ production on free proton:
- ∅ |p_k| fixed by E and p conserv. E = M -
- ∅ production on nuclear target: **A=2**
- ∅ assume production on a single nucleon while the others are fixed (impulse approx.)
- ∅ throw p_{Fermi} using momentum space wave function (Bonn potential)
- ∅ For **A>2**: How to treat relative momentum of spectators ?
- ∅ Use spectral function P(k,E):
- ∅ P(k,E) represents the probability to find in a nucleus a nucleon of momentum k and removal energy E
- ∅ Use spectral function by Benhar et al (NPA 579(1994) 493)
- ∅ also tried Atti et al (PRC 53(1996) 1689) with similar results.

$$E_t = M_d - \sqrt{m_{spectator}^2 + p_{Fermi}^2}$$

SIMC created by Tom O'Neill, Naomi Makins, John Arrington and many others.

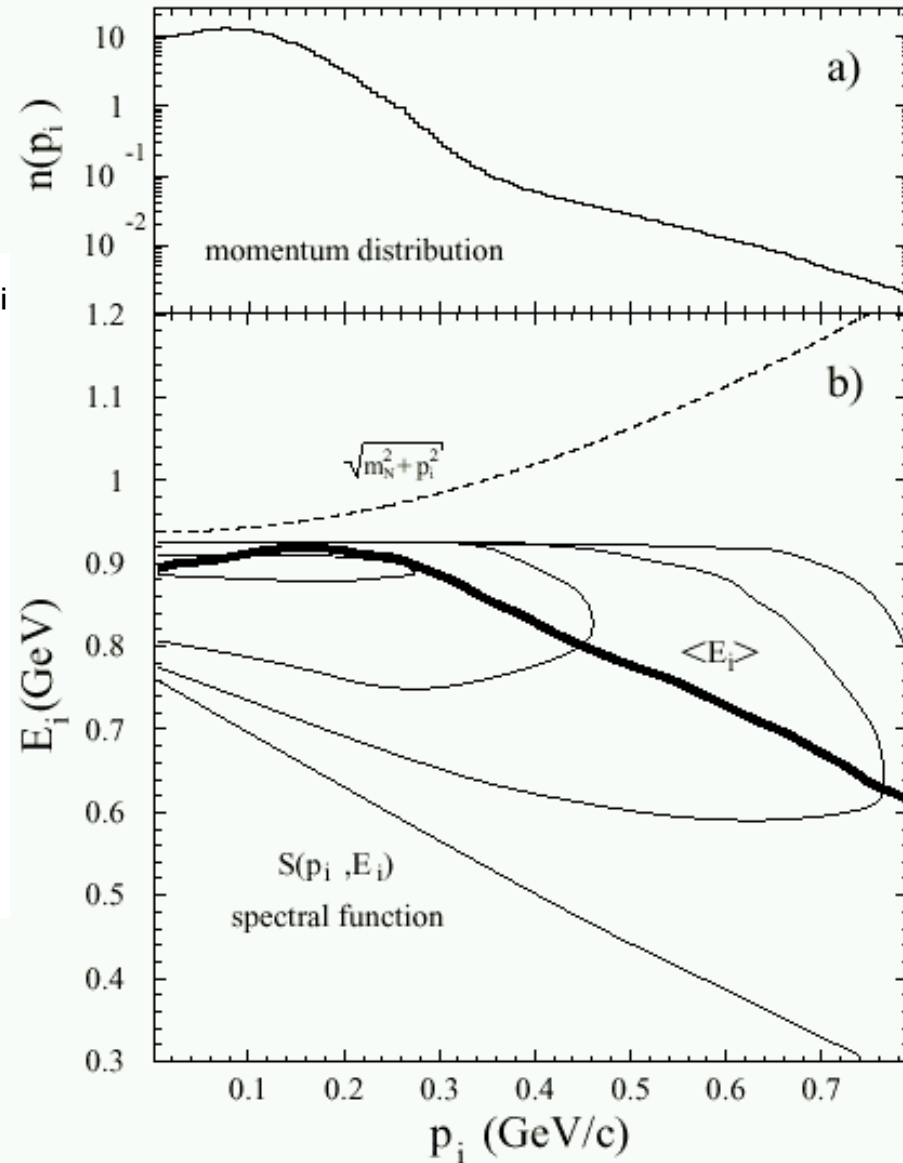
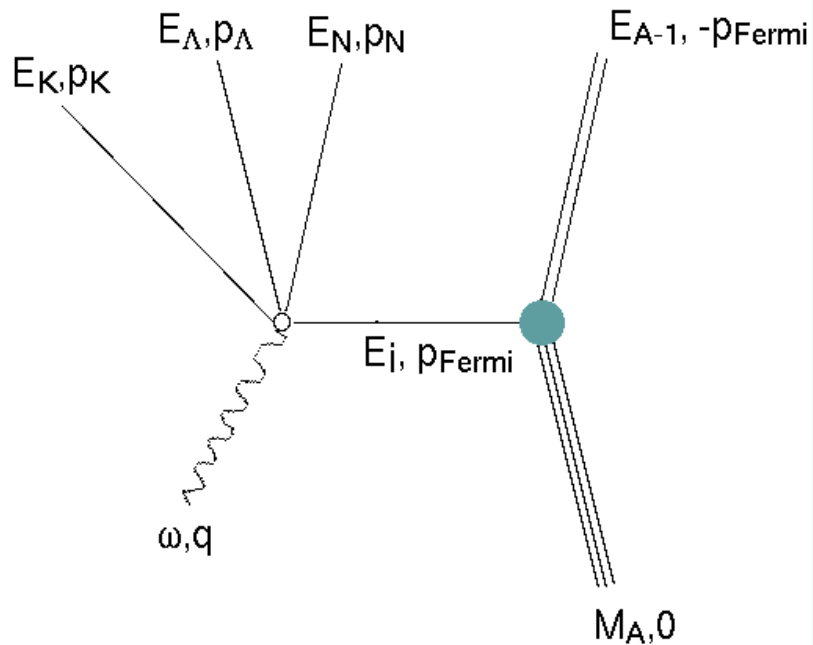
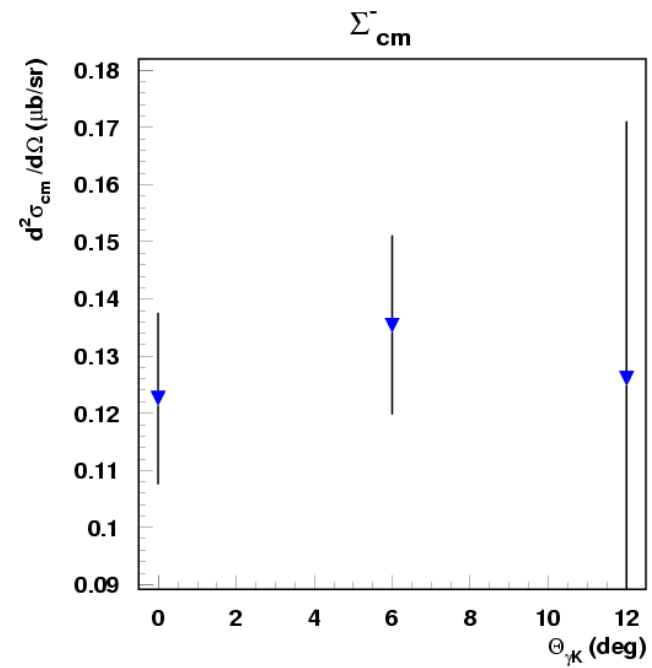
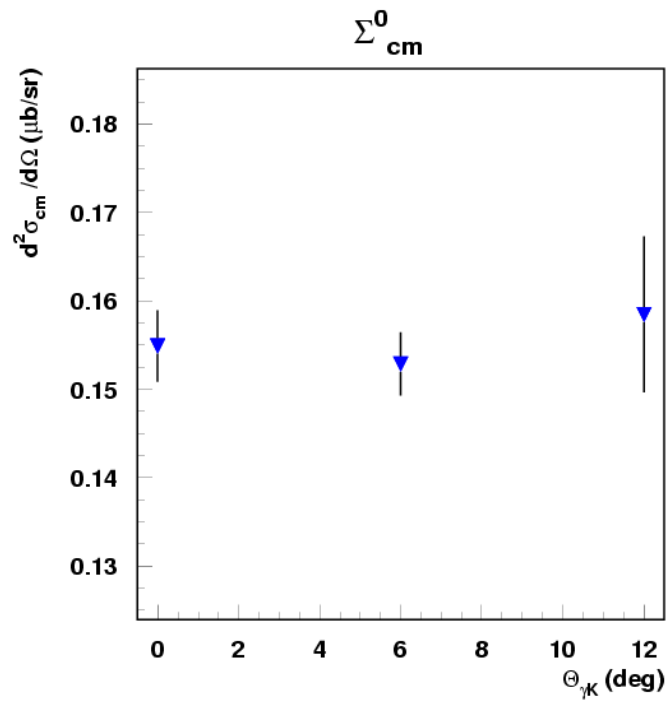
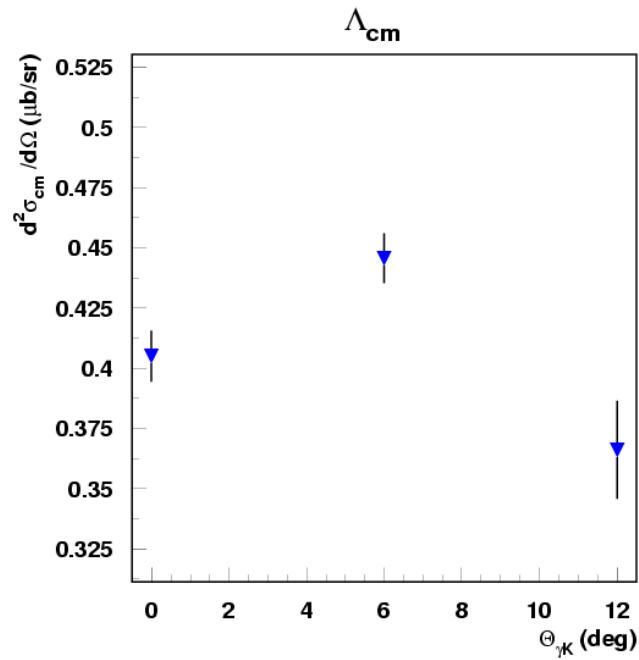
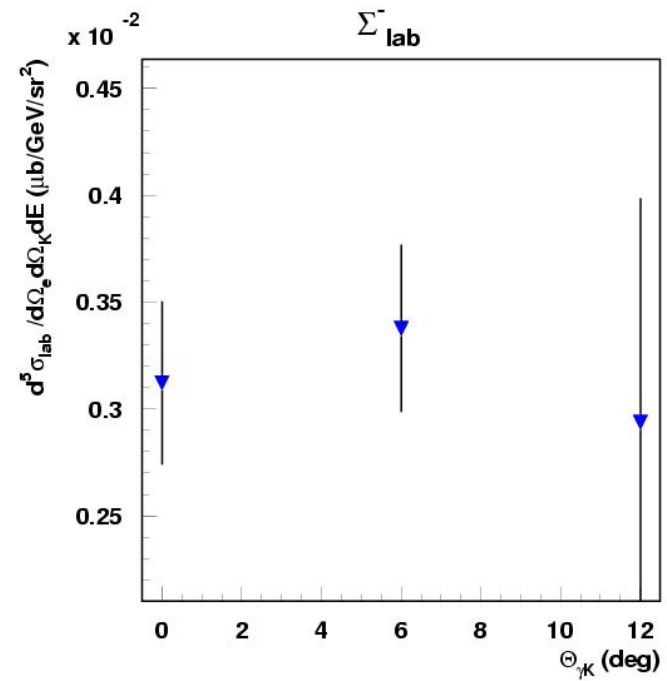
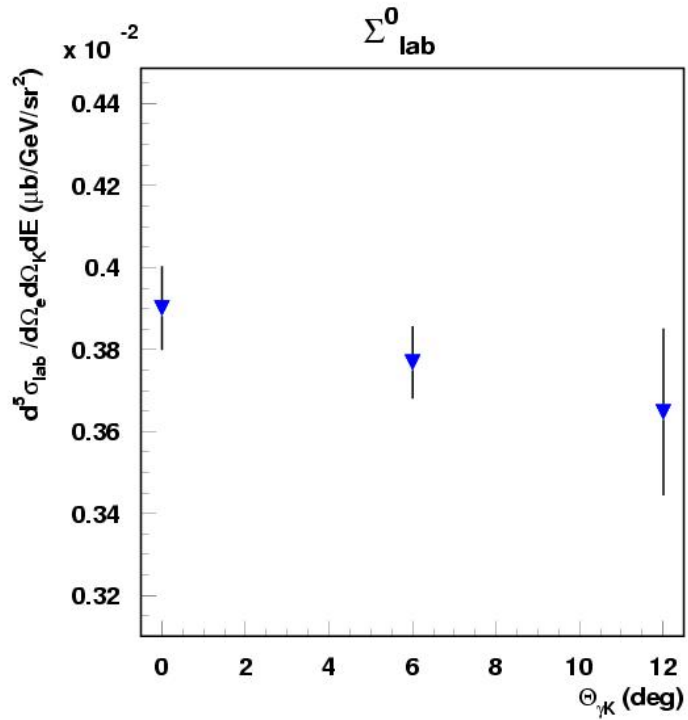
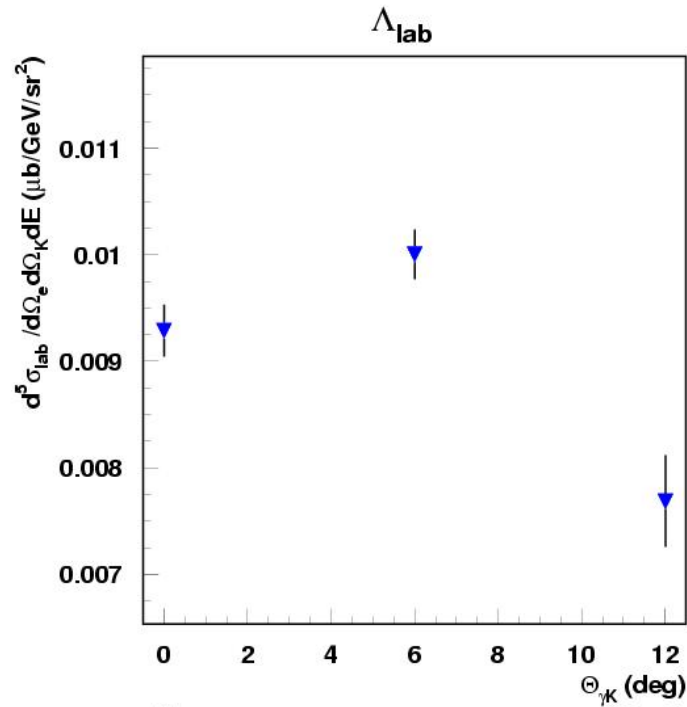


Fig. 16. Momentum distribution of nucleons in the nucleus (a) and the spectral function for the same nucleus (b), as calculated by Sick *et al.* [35]. The thick solid line is the ridge of the distribution. The dashed line correspond to the free space energy

Elementary Cross Sections for ^3He



Elementary Cross Sections for ^3He



Bound L-Hypernuclear States for He

A=3

⇒ ${}^3_{\Lambda}\text{H}$ (*Hypertriton*) $B(1/2)^+ = 130 \text{ keV}$

A=4

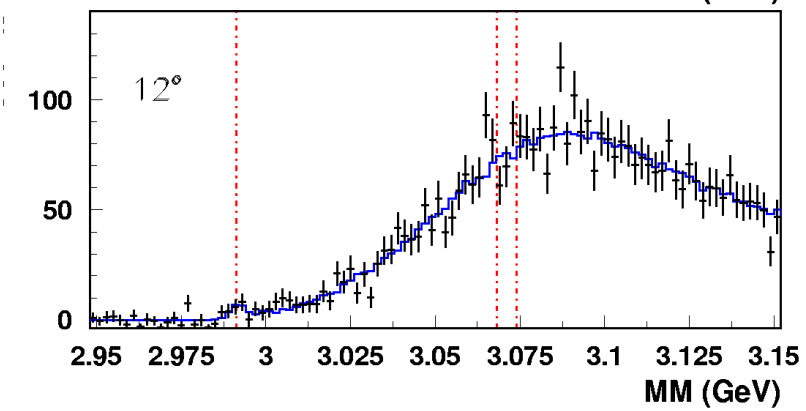
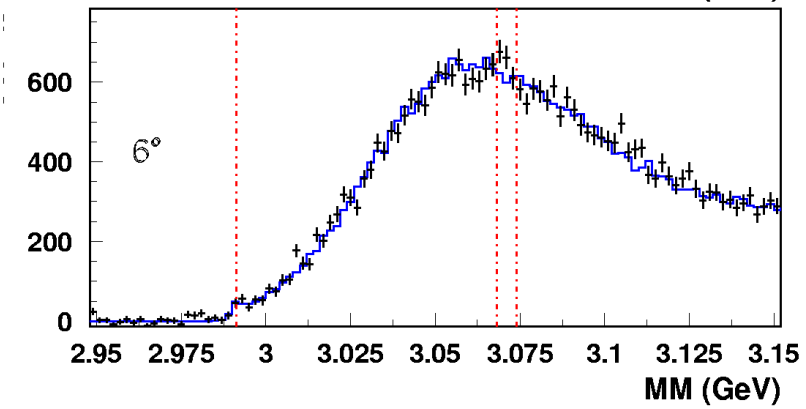
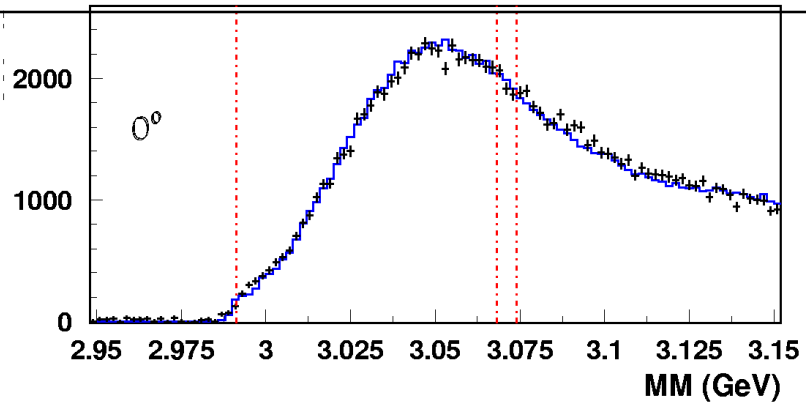
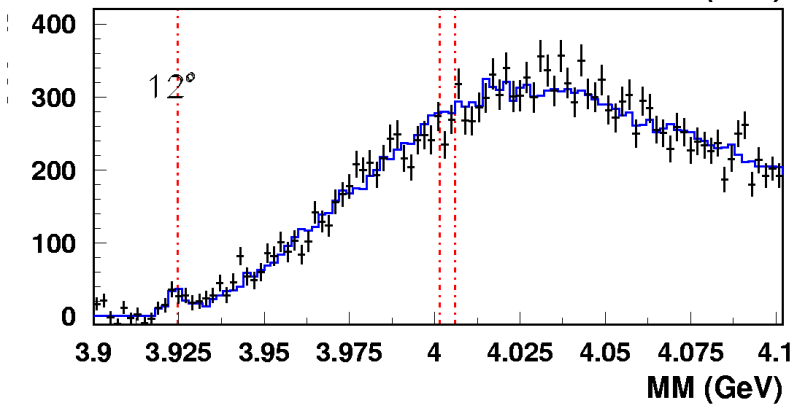
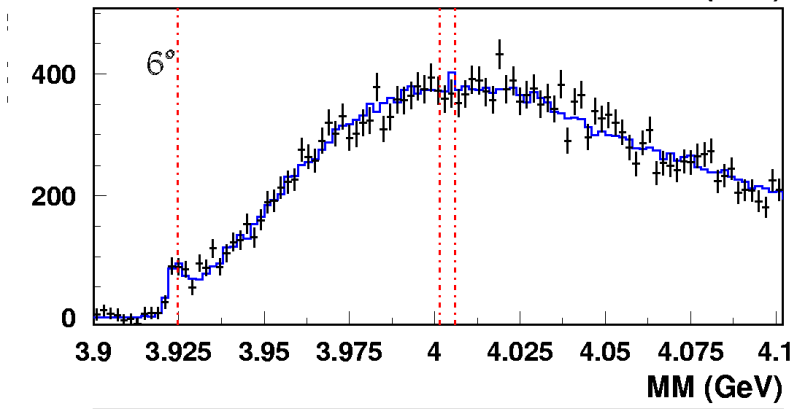
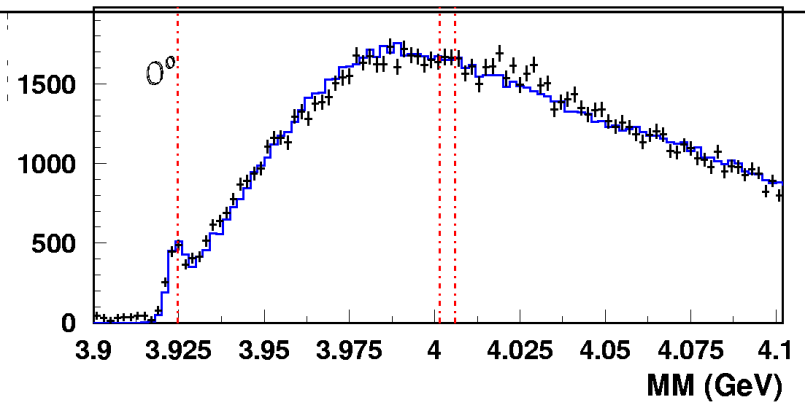
⇒ ${}^4_{\Lambda}\text{H}$ $B_{\text{ground}}(0^+) = 2.04 \pm 0.04 \text{ MeV}$
 $B_{\text{excited}}(1^+) = 1.00 \pm 0.06 \text{ MeV}$

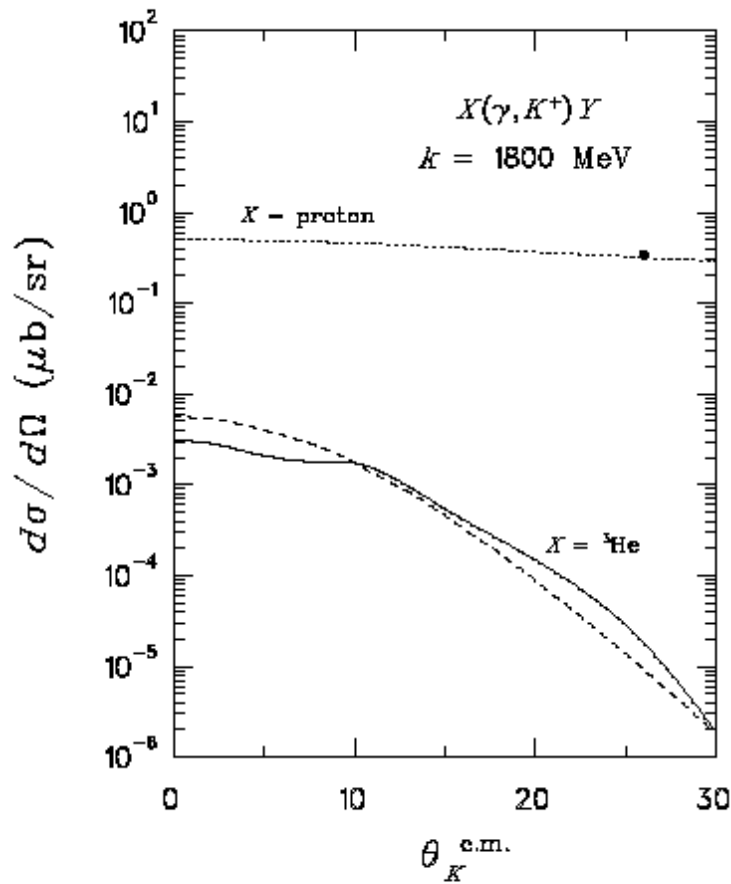
∅ Production Mechanism are different:

(K,p) negligible spin flip strength,
 good momentum matching,
 may populate substitutional states

(p,K) substantial momentum transfer,
 may excite higher spin states,
 medium spin flip strength

∅ (g*,K) **Large momentum transfer,**
 large spin flip strength

${}^3\text{He}$  ${}^4\text{He}$ 



$$\frac{d\sigma_T}{d\Omega_K} = \frac{1}{6} W_A^2 |F(Q)|^2 \left(\frac{d\sigma_T}{d\Omega_K} \right)_{\text{proton}},$$

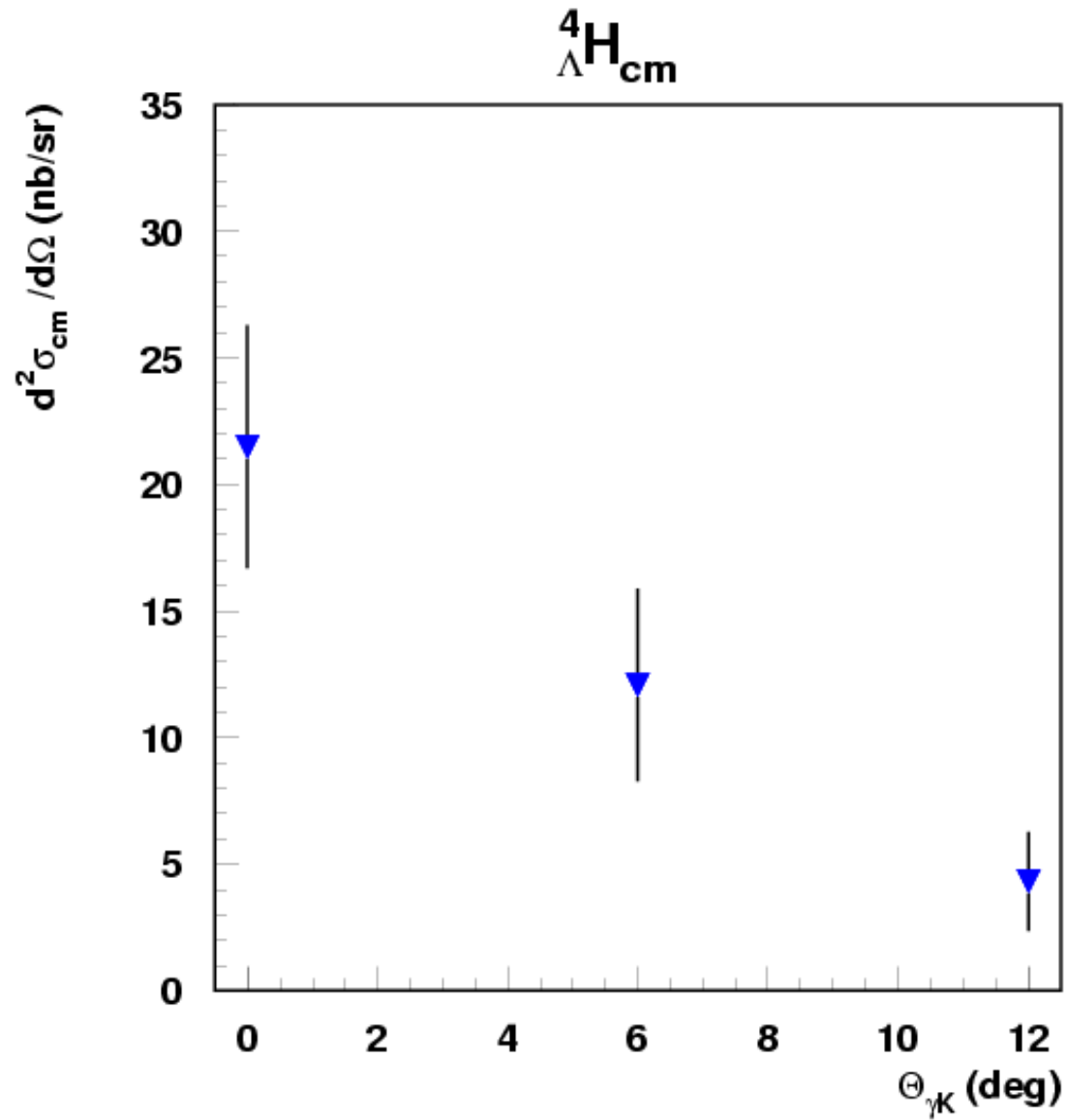
$$F(Q) = \int d^3q d^3p \Psi_{^3\text{H}}(p, q + \frac{2}{3}Q) \Psi_{^3\text{He}}(p, q)$$

$$W_A = \sqrt{\frac{|\mathbf{q}_K^{\text{c.m.}}|_{^3\text{He}} |\mathbf{k}^{\text{c.m.}}|_p M_{^3\text{He}} E_{^3\text{H}} W_p^2}{|\mathbf{k}^{\text{c.m.}}|_{^3\text{He}} |\mathbf{q}_K^{\text{c.m.}}|_p m_p E_A W_{^3\text{He}}^2}}$$

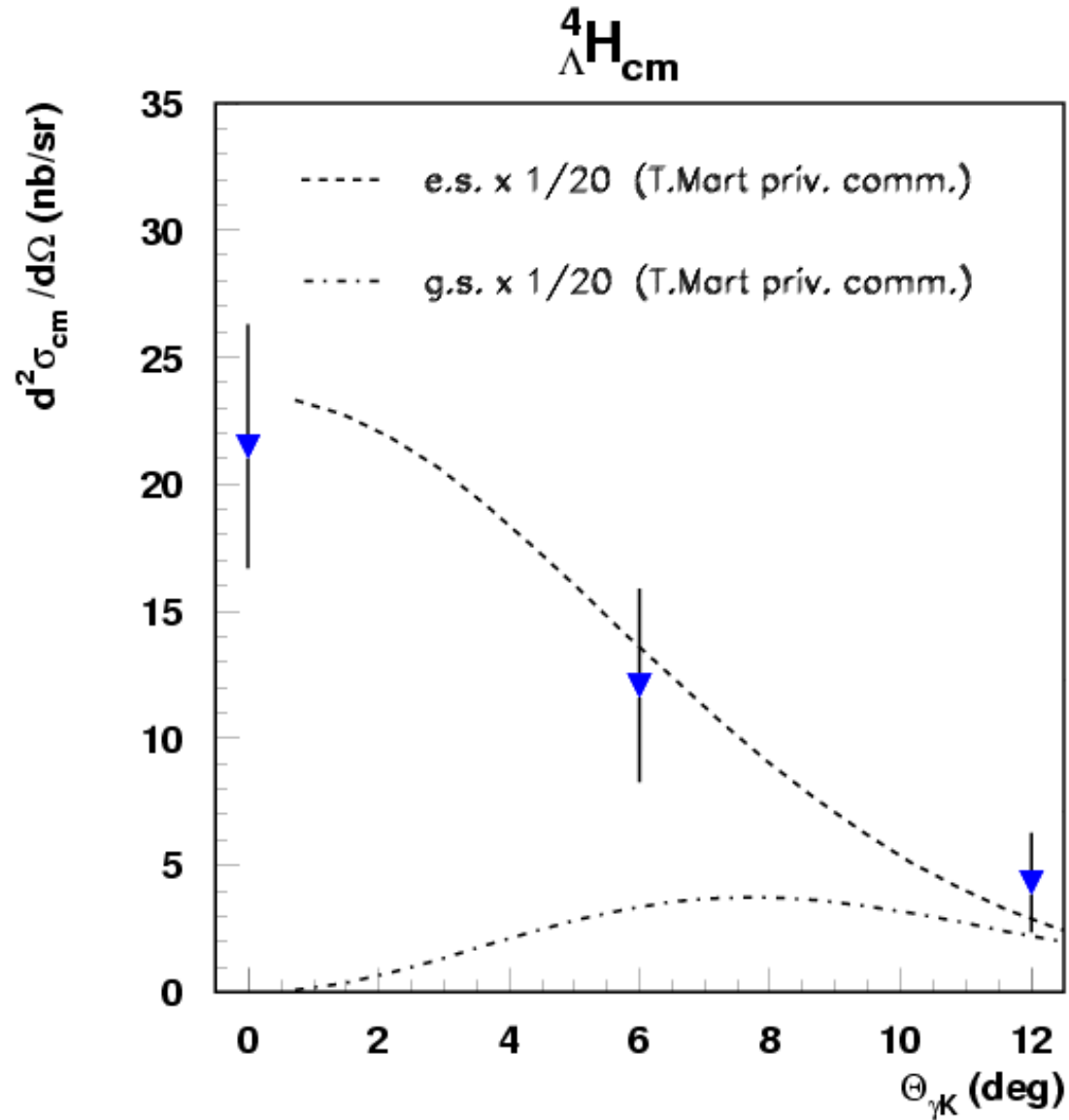
Cross section is kinematic factor times overlap integral times elementary cross section

Fig. 4. Differential cross section for kaon photoproduction off the proton and ^3He as function of kaon angle. The elementary reaction (dotted line) is taken from Ref. [32] and the corresponding experimental datum is from Ref. [35]. The dashed line shows the approximation for production off ^3He calculated from Eq. (32), the solid line represents the exact calculation using S-waves.

Bound Λ -Hypernuclei $A=4$

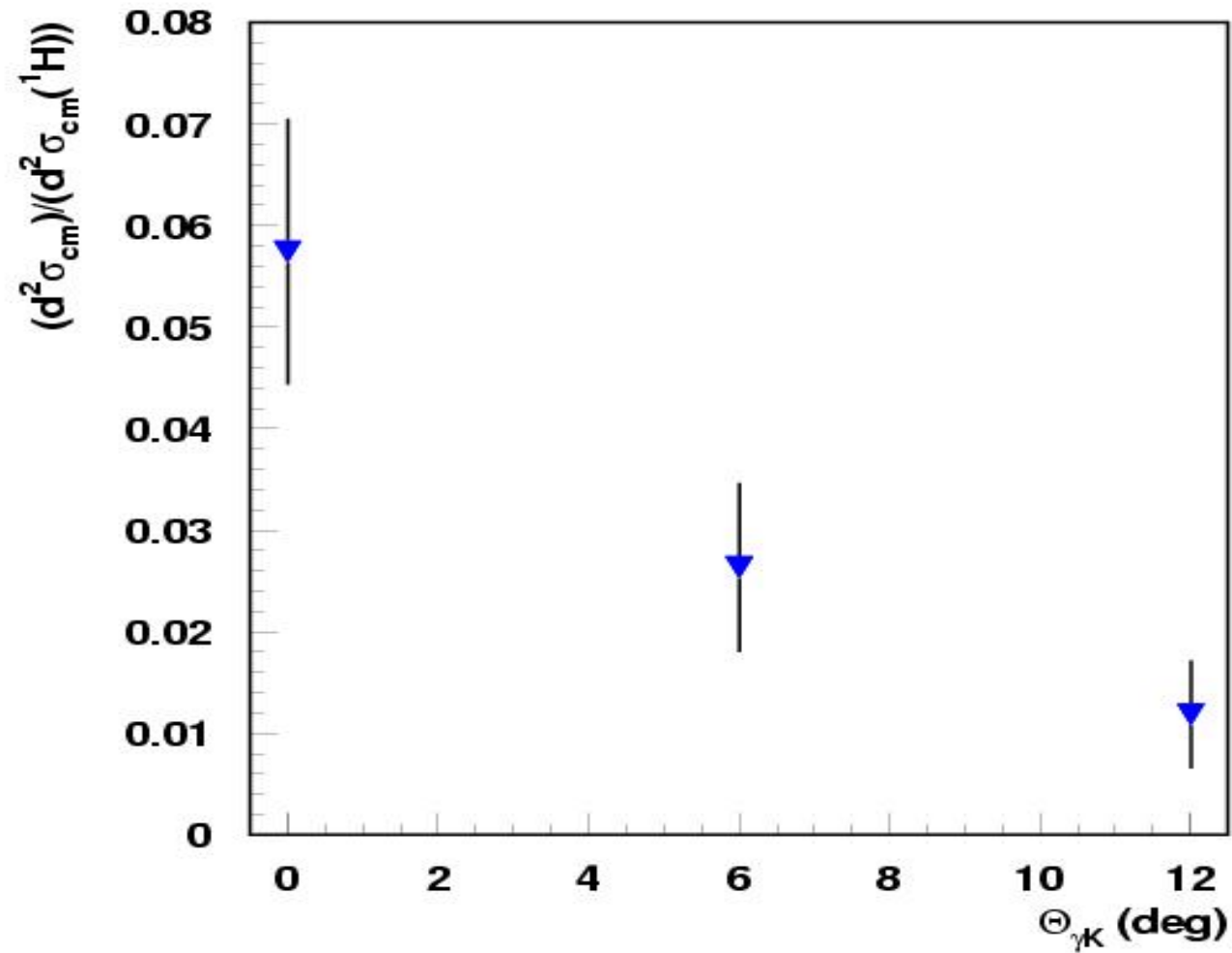


Bound Λ -Hypernuclei $A=4$

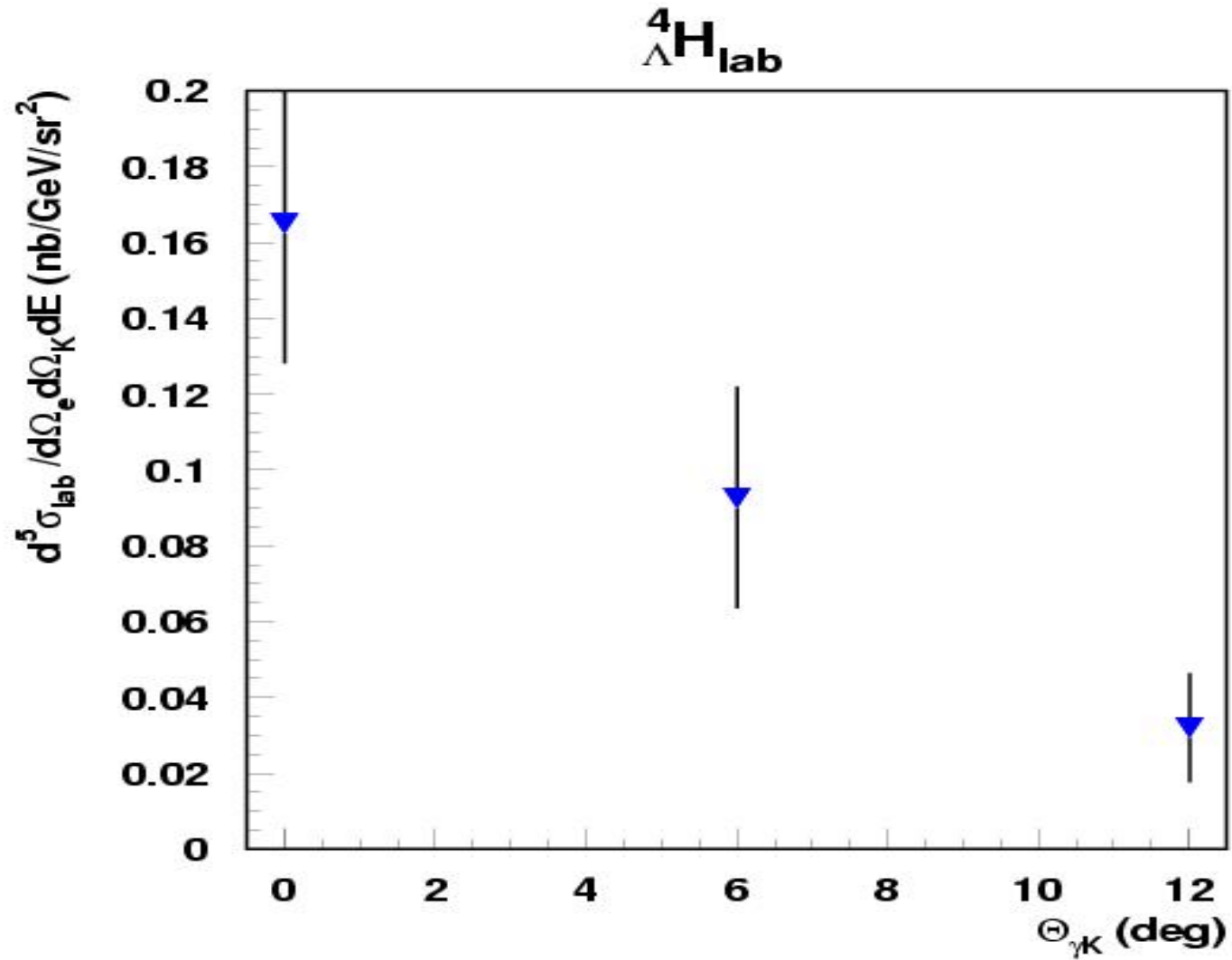


Bound Λ -Hypernuclei $A=4$

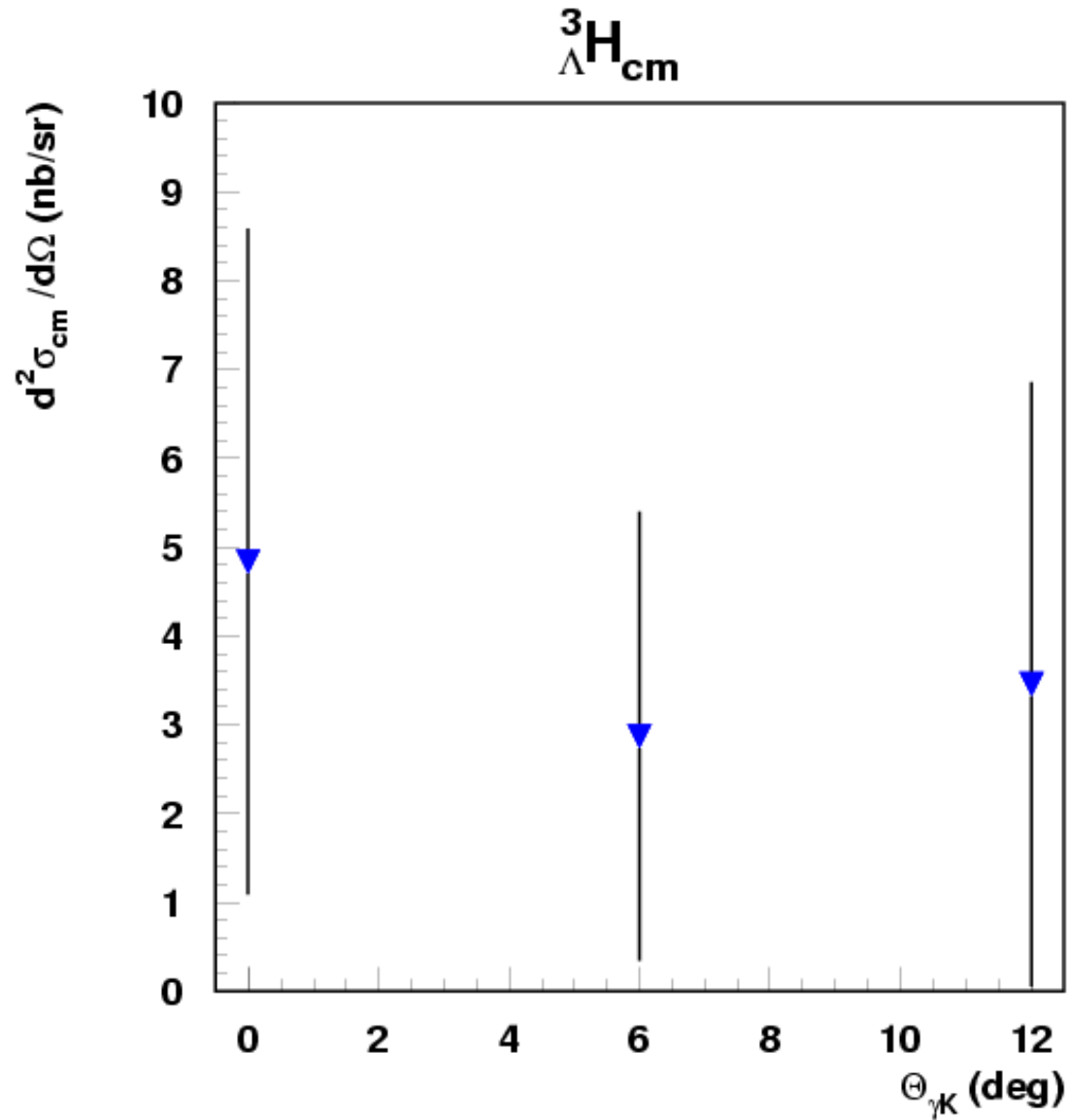
form factor CM



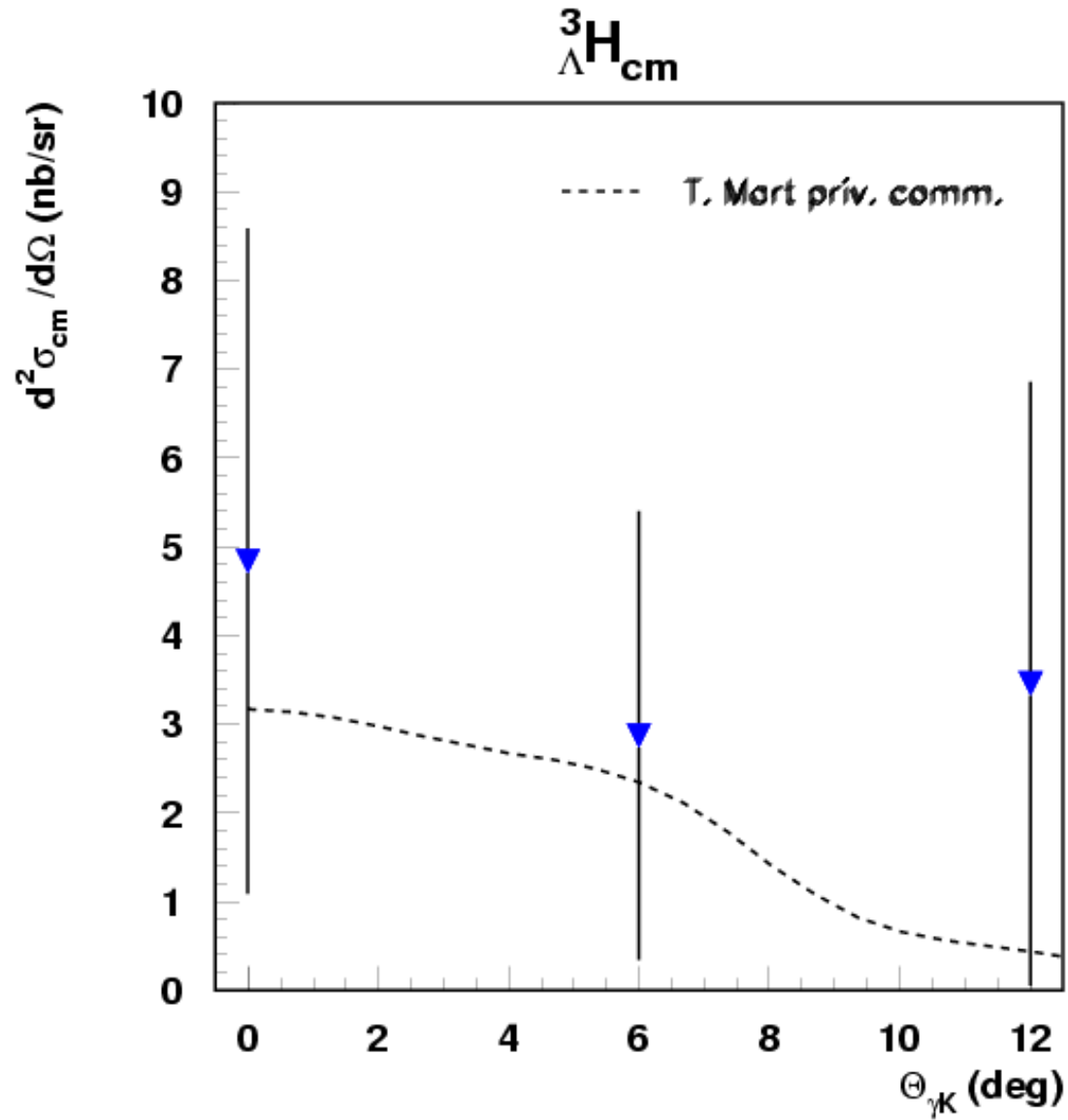
Bound Λ -Hypernuclei $A=4$



Bound Λ -Hypernuclei $A=3$

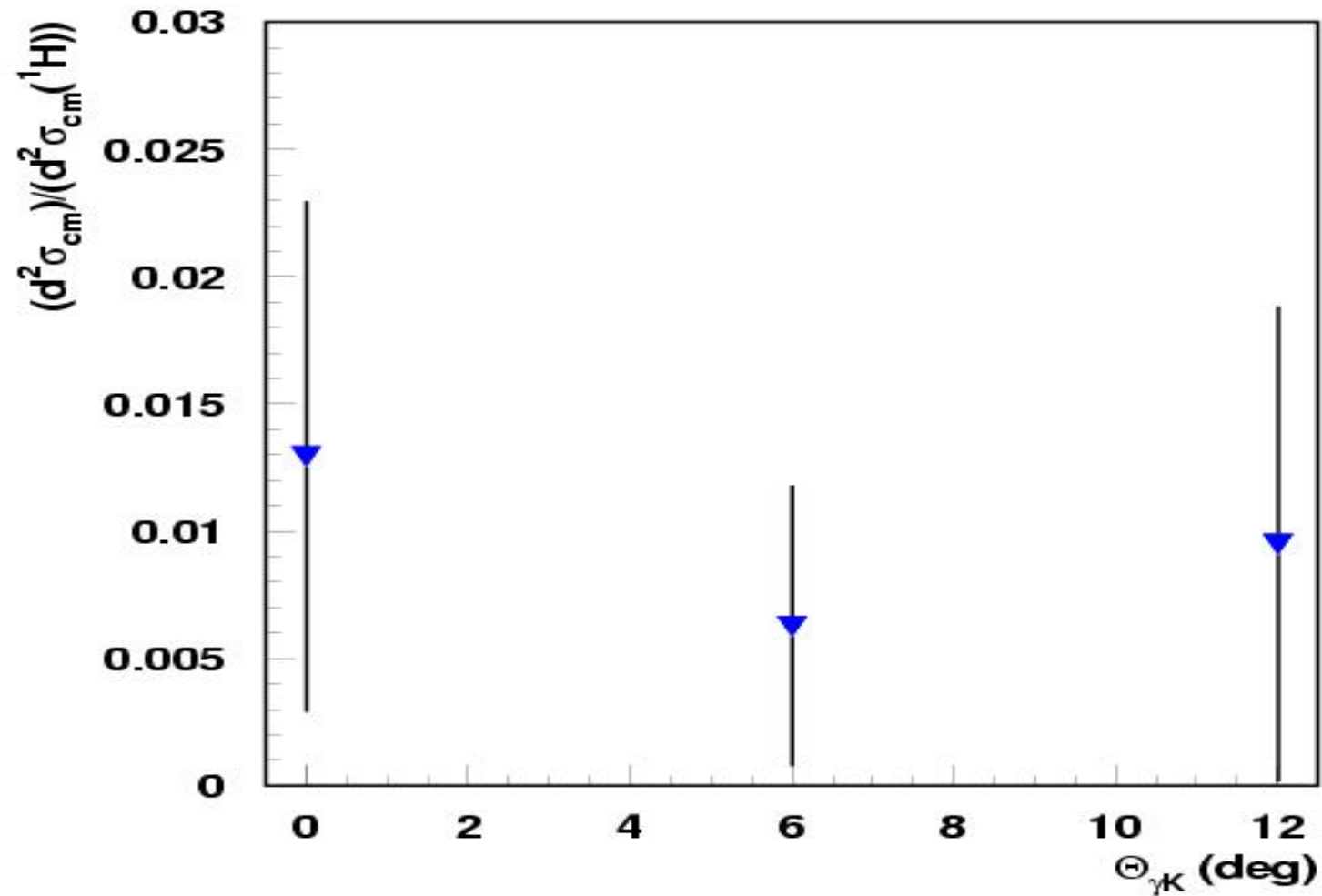


Bound Λ -Hypernuclei $A=3$

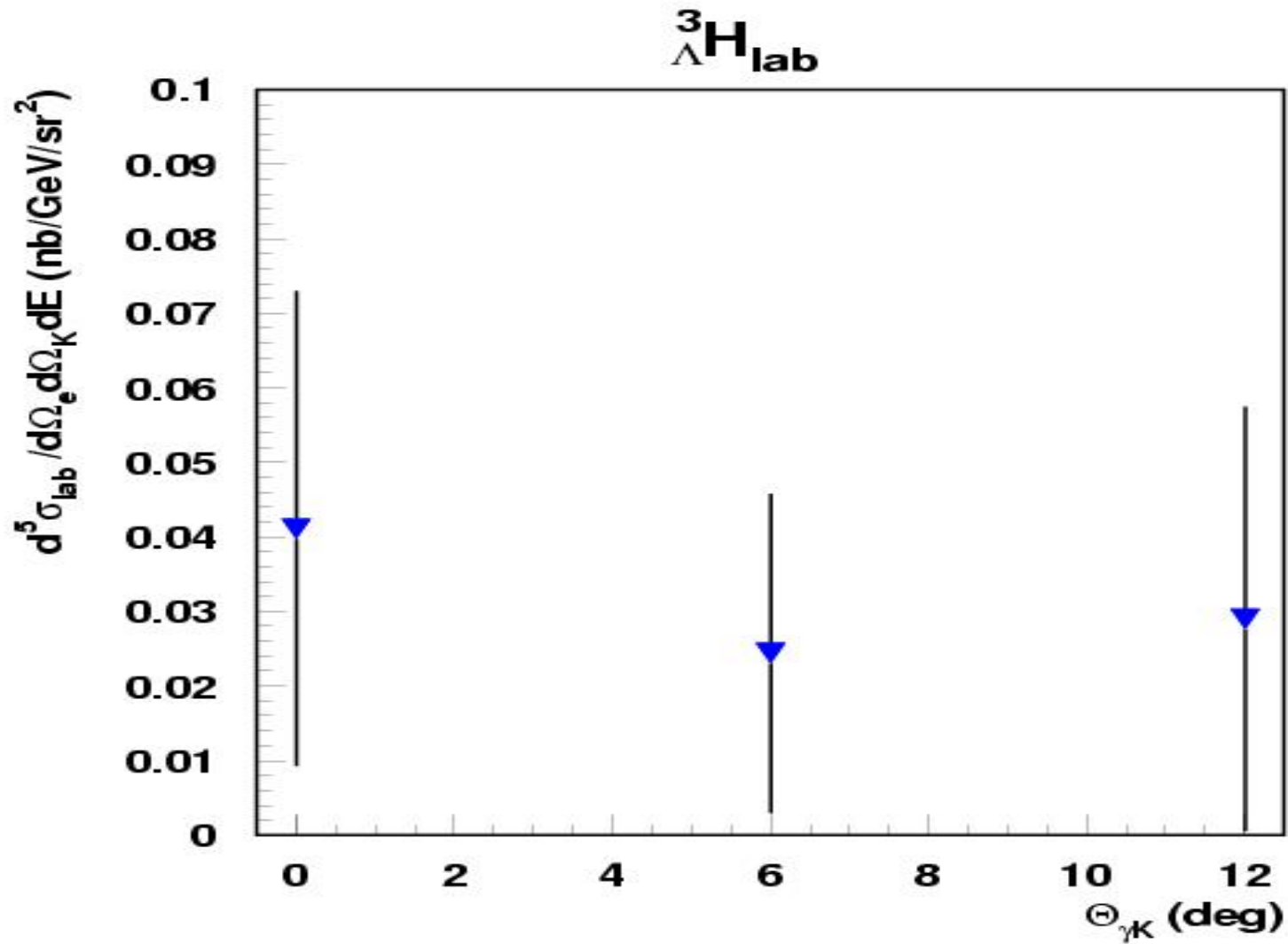


Bound Λ -Hypernuclei $A=3$

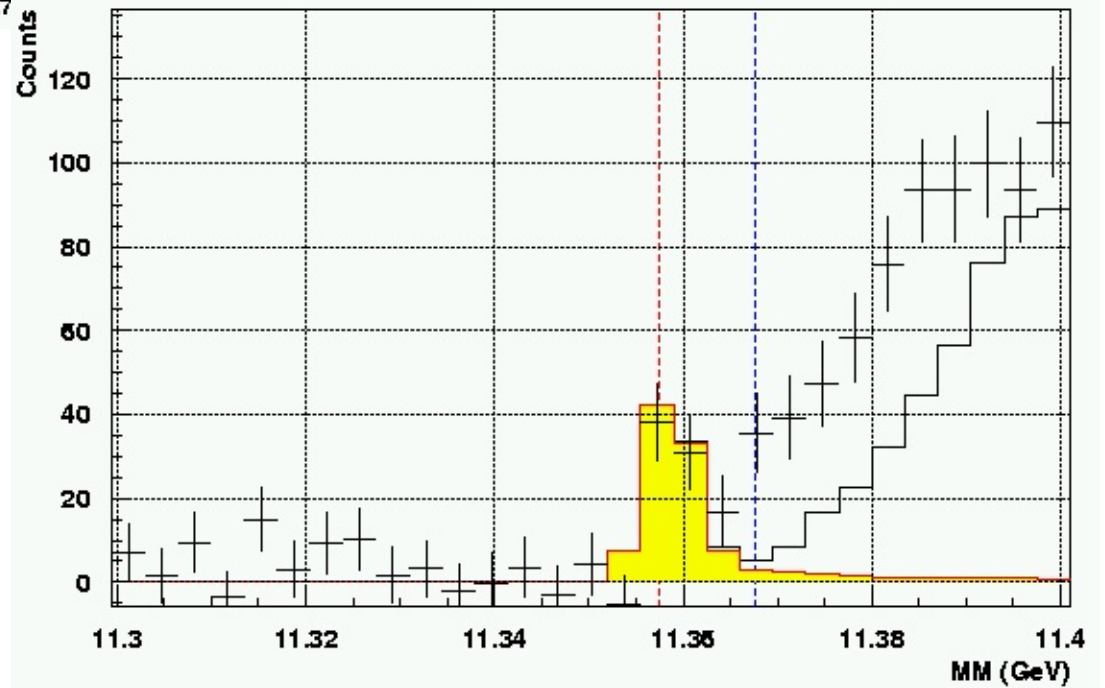
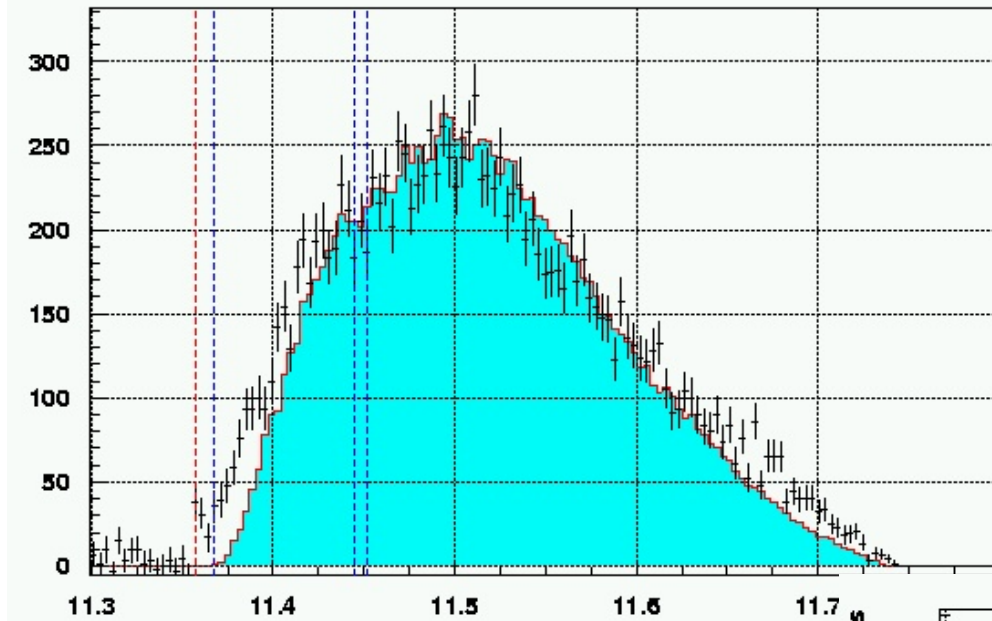
form factor CM



Bound Λ -Hypernuclei $A=3$



MMCarbon with ${}_{\Lambda}^{12}B$ Boundstate



Preliminary A-dependence
for $A(e,e'K)Y$ for
different
targets.

Thesis: AliciaUzzle,
Hampton Univ. 2001

...this is work in progress

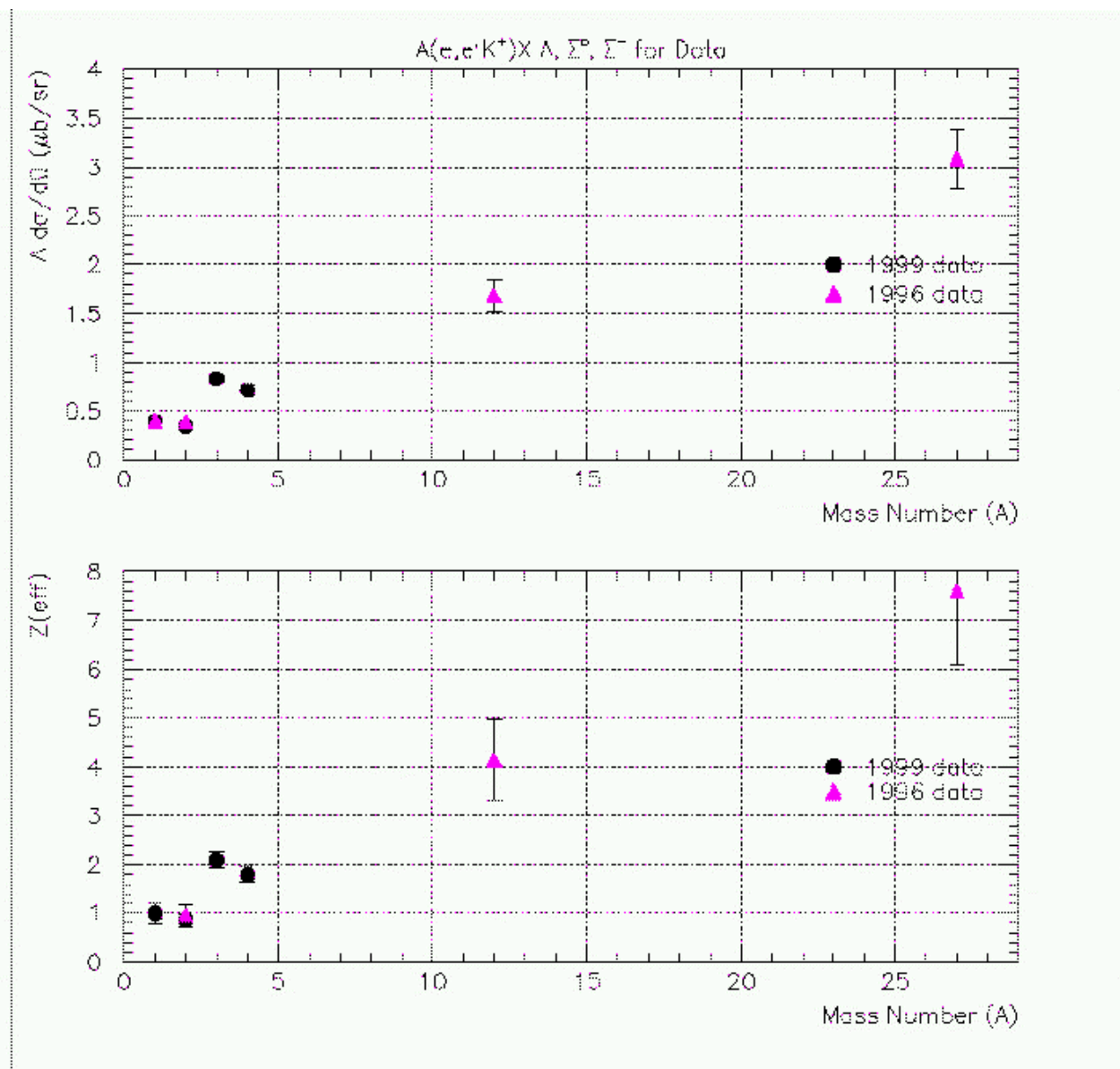


Figure 5.27: The A cross section and effective proton number for $\theta_{\gamma K} = 0^\circ$. Shown are the present data and 1996 kaon electroproduction data which includes the aluminum and carbon targets.

Summary and Outlook

- z $A(e,e'K)$ for $A=3,4,(12)$ has been measured at low Q^2 .
- z FSI (effective range approx.) yields satisfactory description using Nijmegen YN 97f.
- z Spectral functions (Benhar et al) are essential for the analysis of data on nuclear targets.
- z Statements on separated quasifree $\Sigma^{0,-}$ contributions require strong assumptions.
- z ${}^4\text{He}(e,e'K)\Lambda$ and ${}^3\text{He}(e,e'K)\Lambda$ data clearly indicate ${}^4\text{H}_\Lambda$ and ${}^3\text{H}_\Lambda$ bound states.
- z First measurement in electroproduction.
- z Indication of ${}^{12}\text{B}_\Lambda$ boundstate.

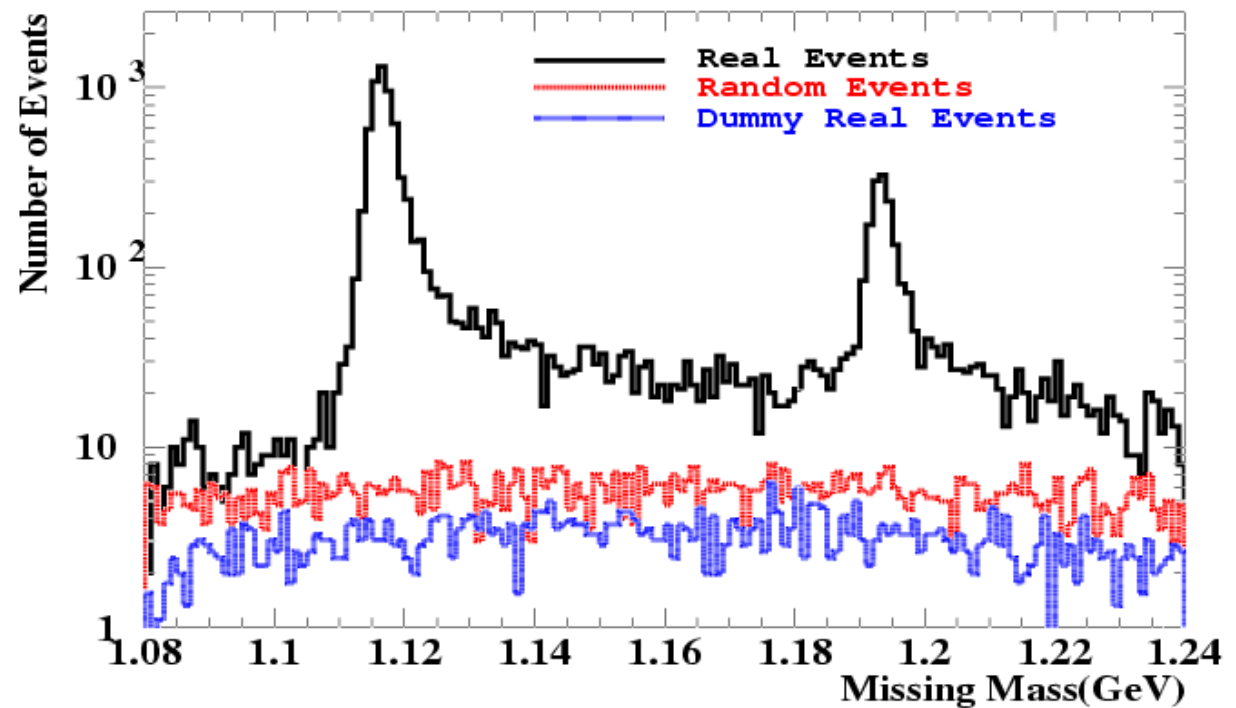
Jefferson Lab E91-016 Collaboration:

D. Abbott^b, A. Ahmidouch^{dfe}, **P. Ambrozewicz**^g, C.S. Armstrong^{bb}, J. Arrington^{c1},
R. Asaturyan^l, K. Assamagan^f, S. Avery^f, K. Bailey^c, O.K. Baker^f, S. Beedoe^{d1},
H. Bitao^f, W. Boeglin^{ab}, H. Breuer^k, D.S. Brown^k, R. Carlini^b, **J. Cha**^f, N. Chant^k,
E. Christy^f, A. Cochran^f, L. Cole^f, G. Collins^k, C. Cothran^l, J. Crowder^{mn},
W.J. Cummings^c, S. Danagoulian^{db}, F. Dohrmann^{cn}, F. Duncan^k, J. Dunne^b,
D. Dutta^o, T. Eden^f, M. Elaasar^p, R. Ent^b, L. Ewell^k, H. Fenker^b, H.T. Fortune^{q1},
Y. Fujii^r, L. Gan^f, H. Gao^c, K. Garrow^b, D.F. Geesaman^c, P. Gueye^f,
K. Gustafsson^k, K. Hafidi^c, J.O. Hansen^c, **W. Hinton**^f, H.E. Jackson^c, H. Juengst^g,
C. Keppel^f, A. Klein^t, **D. Koltenuk**^{q1}, Y. Liangⁿ, J.H. Liu^g, A. Lung^b, D. Mack^b,
R. Madey^{fe}, P. Markowitz^{ab}, C.J. Martoff^g, D. Meekins^b, J. Mitchell^b, T. Miyoshi^r,
H. Mkrtchyan^l, **R. Mohring**^k, S.K. Mtingwa^{d1}, B. Mueller^c, T.G. O'Neill^c,
G. Niculescu^{f v}, I. Niculescu^{f w}, D. Potterveld^c, J.W. Price^x, B.A. Raue^{ab},
P.E. Reimer^c, J. Reinhold^{abc}, J. Roche^b, P. Roos^k, M. Sarsour^r, Y. Sato^r,
G. Savage^f, R. Sawafte^{d1}, R.E. Segel^o, A. Semenov^e, S. Stepanyan^l, V. Tadevosian^l,
S. Tajima^g, L. Tang^f, B. Terburg⁺, **A. Uzzle**^f, S. Wood^b, H. Yamaguchi^r, C. Yan^{1 b},
C. Yan^{2 e}, L. Yuan^f, B. Zeidman^c, M. Zeier^l, and B. Zihlmann^l

^aFlorida International University, ^bThomas Jefferson National Accelerator Laboratory,
^cArgonne National Laboratory, ^dNC A&T State University, ^eKent State University,
^fHampton University, ^gTemple University, ^hCollege of William and Mary, ¹California
Institute of Technology, ^lYerevan Physics Institute, ^kUniversity of Maryland, ^lUni-
versity of Virginia, ^mJuniata College, ⁿForschungszentrum Rossendorf, ^oNorthwestern
University, ^pSouthern University at New Orleans, ^qUniversity of Pennsylvania, ^rTohoku
University, ^sUniversity of Minnesota, ^tOld Dominion University, ^uAmerican University,
^vOhio University, ^wThe George Washington University, ^xRensselaer Polytechnic Insti-
tute, ^yUniversity of Houston, ^zDuke University, ⁺University of Illinois

Missing Mass Analysis for $p(e, e' K^+)X$

$$\left. \begin{aligned} \vec{p}_e + \vec{p}_p &= \vec{p}_{e'} + \vec{p}_K + \vec{p}_{miss} \\ E_e + E_p^0 &= E_{e'} + E_K + E_{miss} \end{aligned} \right\} m_{miss}^2 = E_{miss}^2 - p_{miss}^2$$



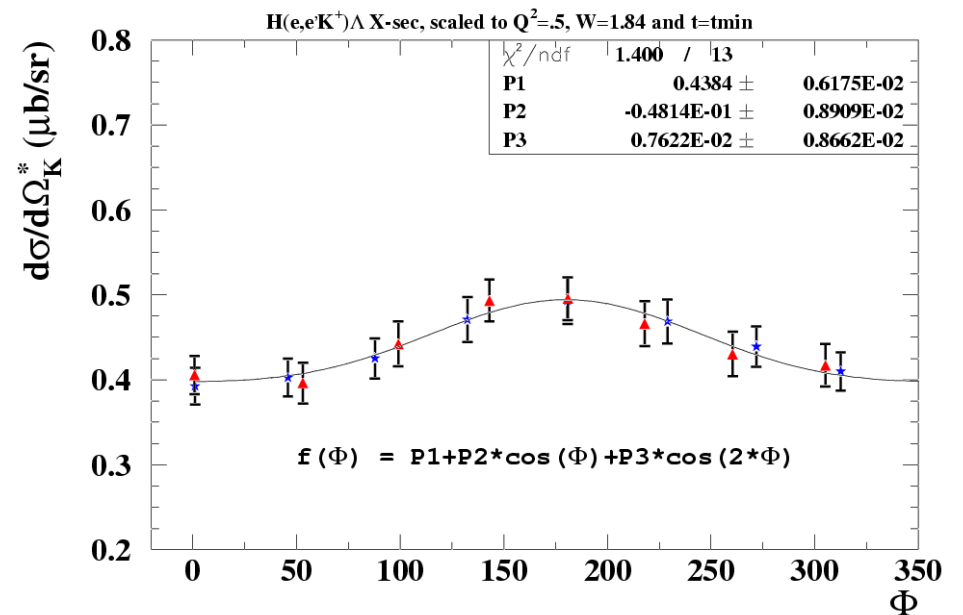
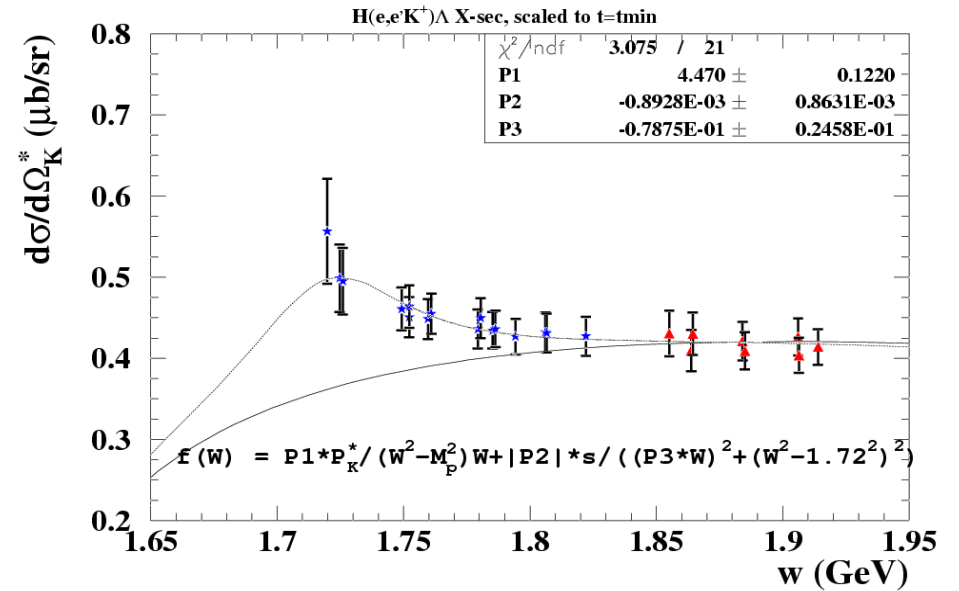
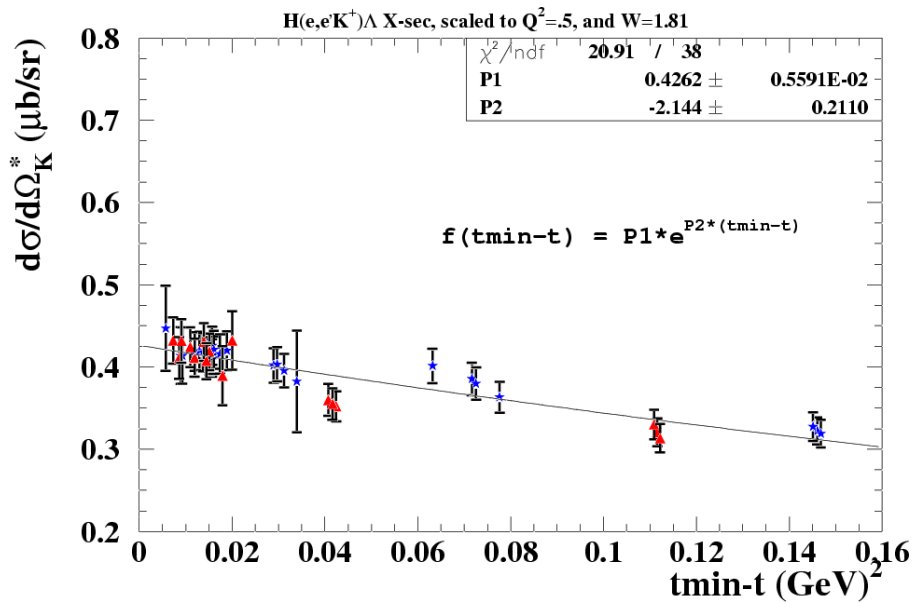
Cross-Section Parametrization

$$f(Q^2) = \text{Constant}$$

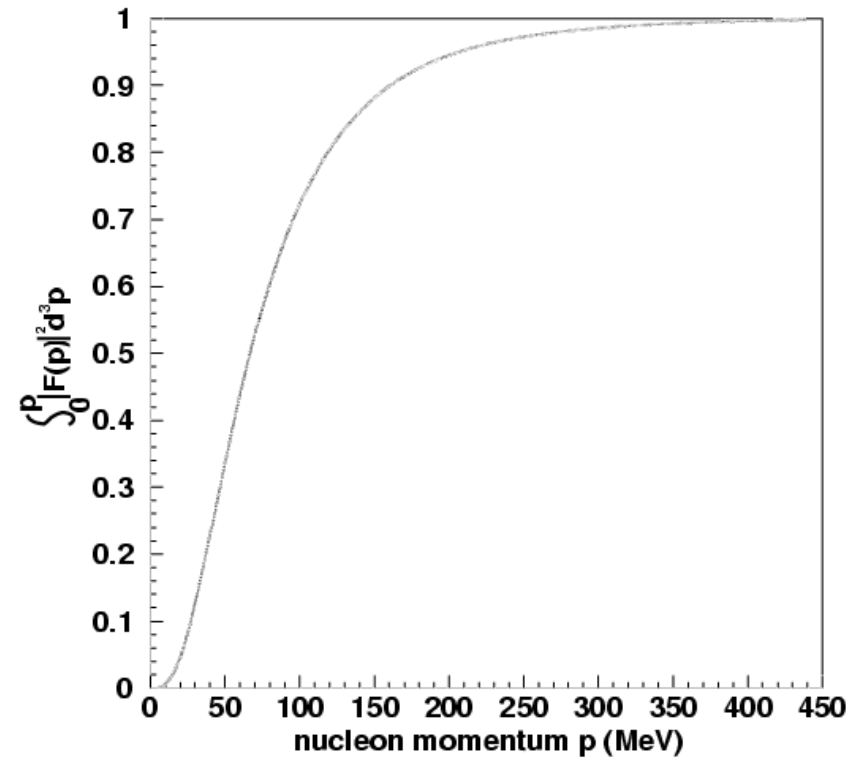
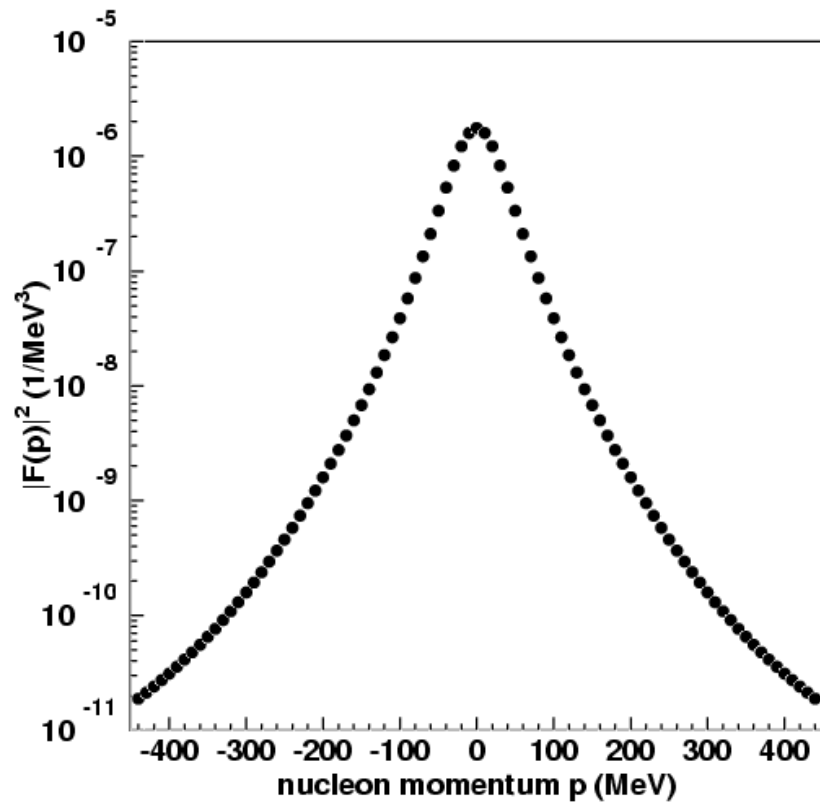
$$g(W) = \frac{P_1 P_k^{CM}}{(W^2 - M_p^2)W} + \frac{P_2 W^2}{(P_3 W)^2 + (W^2 - P_4^2)^2}$$

$$h(t_{min} - t) = P_1 e^{P_2(t_{min} - t)}$$

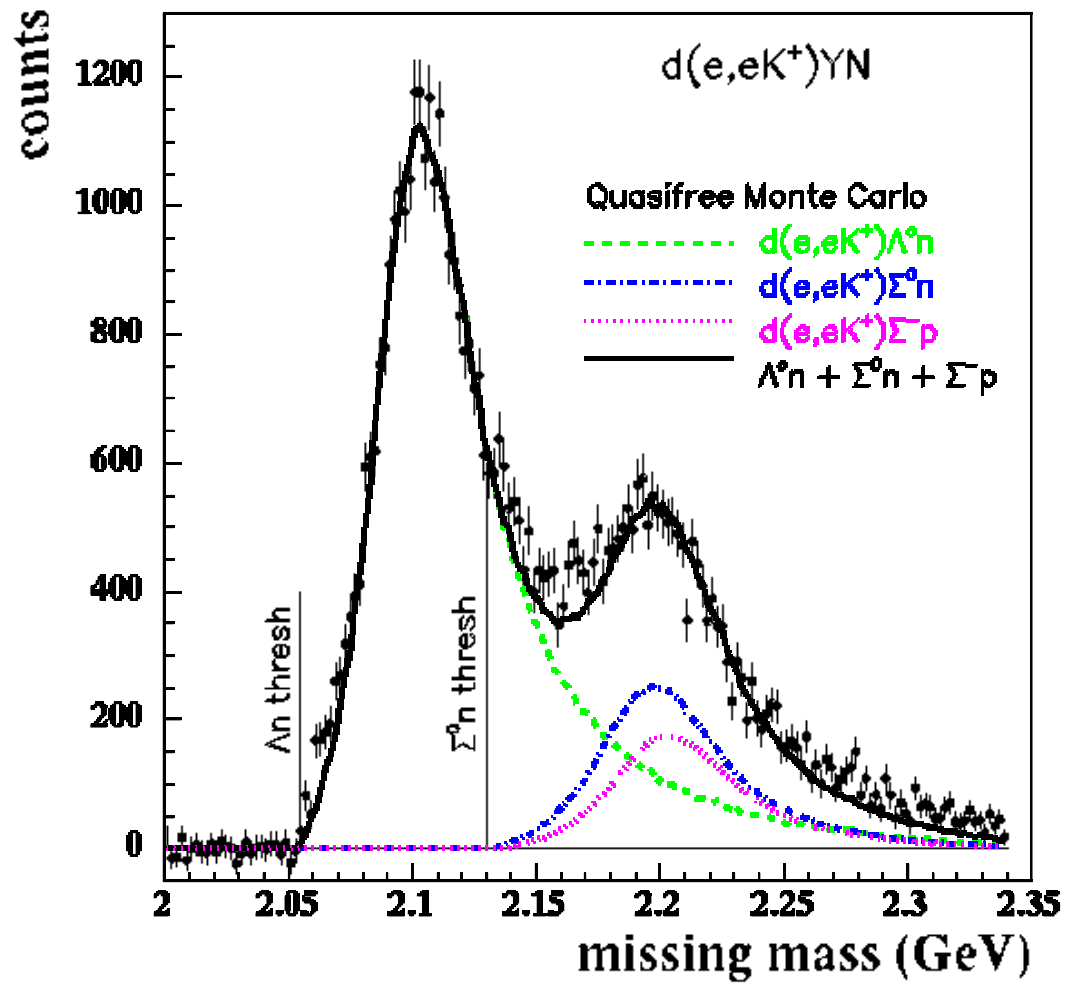
$$i(\phi) = P_1 + P_2 \cos(\phi) + P_3 \cos(2\phi)$$



Momentum Wavefunction (Bonn potential)



${}^2\text{H}(e, e' K^+)$



Modeling of FSI

$$\tilde{M}_{fi} = \frac{\psi(kr + \delta)}{\psi(kr)} M_{fi},$$

$$\left(\frac{d\sigma}{d\Omega}\right)_{FSI} = f_{P.S.} |\tilde{M}_{fi}|^2 = f_{P.S.} \left|\frac{\psi(kr + \delta)}{\psi(kr)}\right|^2 |M_{fi}|^2.$$

$$I = \left|\frac{\psi^*(kr + \delta)}{\psi(kr)}\right|^2$$

Λ n Potential Parametrization

$$V(r) = V_A e^{r^2/\beta_A^2} + V_R e^{r^2/\beta_R^2}$$

| Model | state | V_A (MeV) | β_A (fm) | a (fm) | r (fm) |
|----------|---------------------|-------------|----------------|--------|--------|
| Verma | Singlet (1S_0) | -167.34 | 1.100 | -2.29 | 3.15 |
| | Triplet (3S_1) | -132.42 | 1.100 | -1.77 | 3.25 |
| Jülich A | Singlet (1S_0) | -373.94 | 0.790 | -1.60 | 1.33 |
| | Triplet (3S_1) | -144.14 | 1.059 | -1.60 | 3.15 |
| Jülich B | Singlet (1S_0) | -131.49 | 1.095 | -0.57 | 7.65 |
| | Triplet (3S_1) | -189.60 | 0.964 | -1.94 | 2.42 |

Table 5.9: The parameters for the various potentials used in Eq. 5.11. The strength and range of the repulsive part are fixed for all three potentials at $V_R = 246.80$ MeV, $\beta_R = 0.82$ fm for the singlet state, and $V_R = 181.68$ MeV, $\beta_R = 0.82$ fm for the triplet state. From [8].

Λn FSI from ${}^2\text{H}(e, e'K^+)$

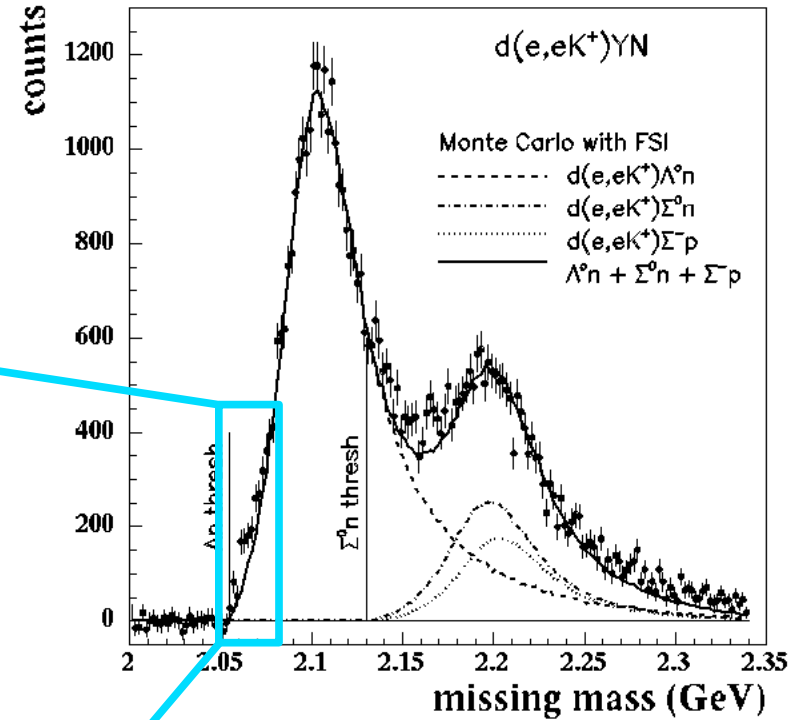
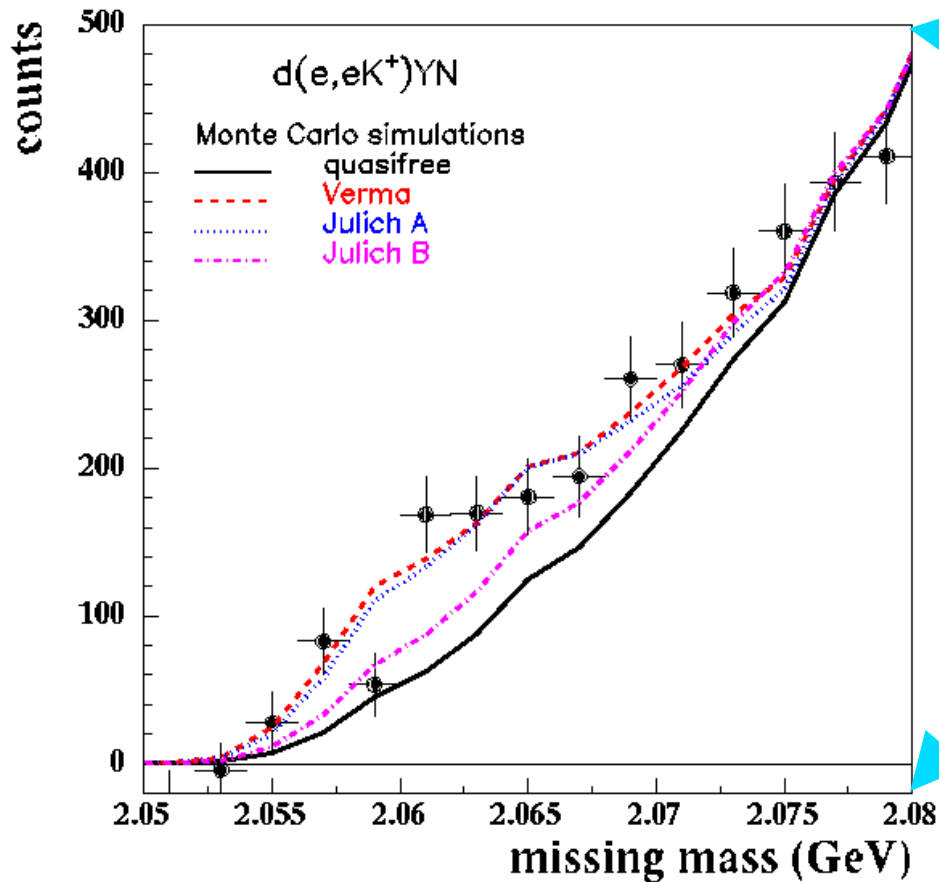
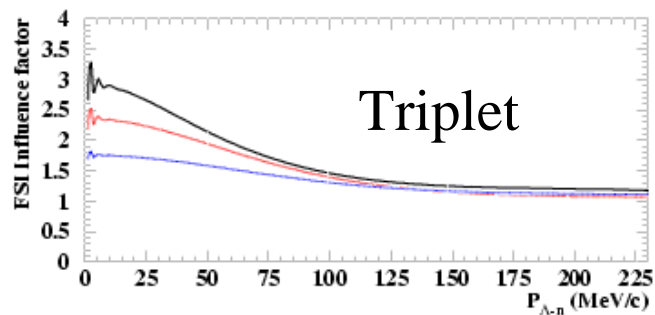
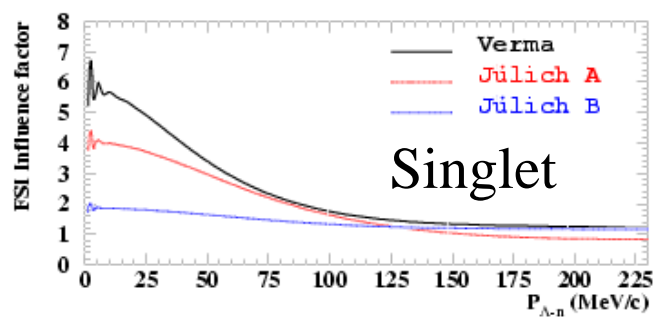
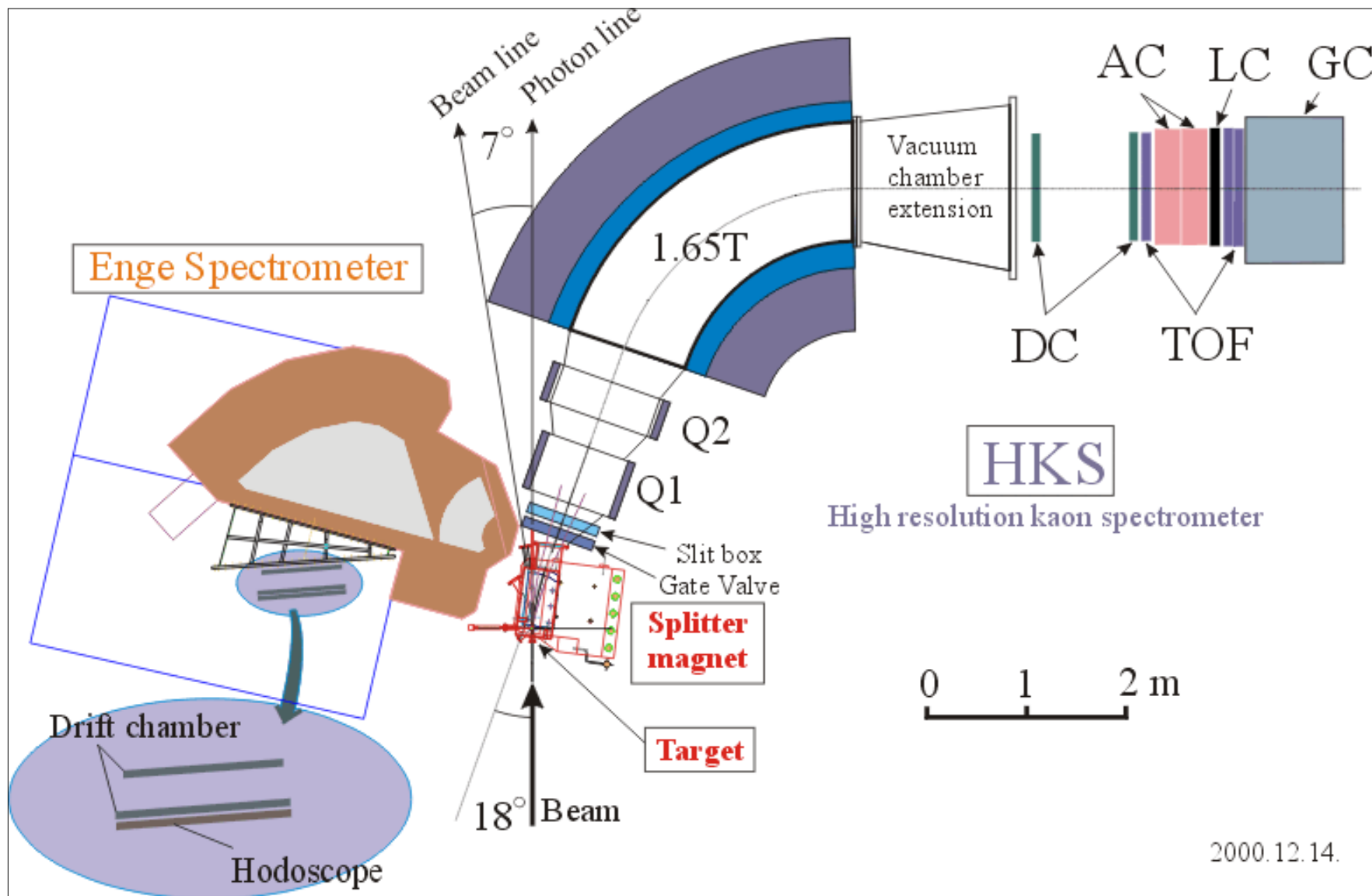


Table 1: Scattering length and effective range for the three hyperon-nucleon potentials used in the simulations.

| Model | State | a (fm) | r (fm) |
|----------|-----------|----------|----------|
| Verma | 1S_0 | -2.29 | 3.15 |
| | 3S_1 | -1.77 | 3.25 |
| Jülich A | 1S_0 | -1.60 | 1.33 |
| | 3S_1 | -1.60 | 3.15 |
| Jülich B | 1S_0 | -0.57 | 7.65 |
| | 3S_1 | -1.94 | 2.42 |

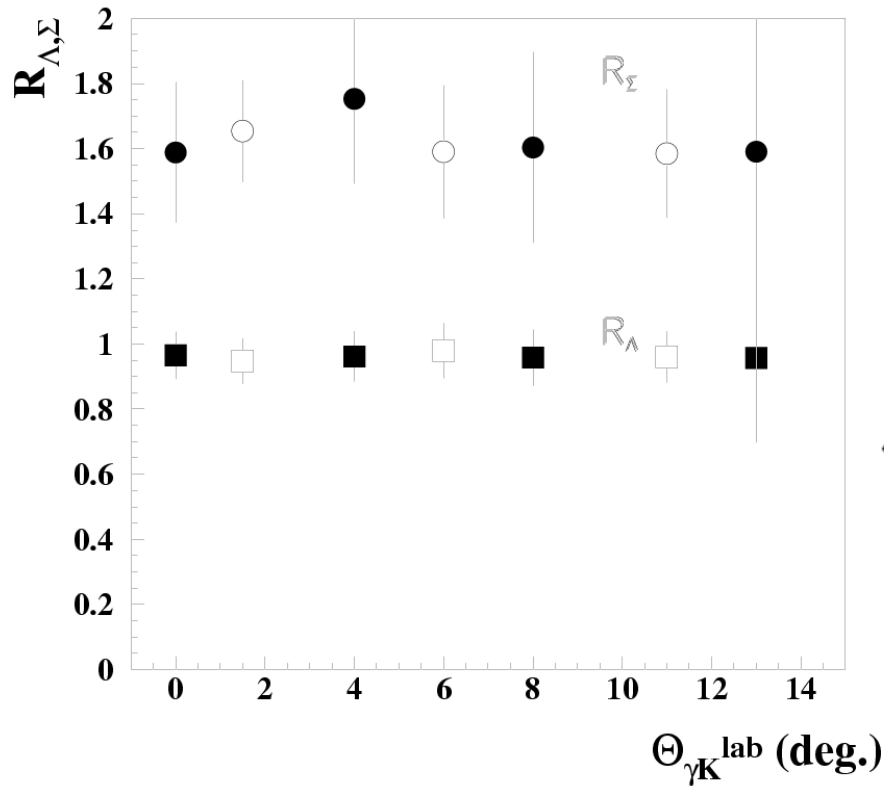


High Resolution Kaon Spectrometer (HKS)

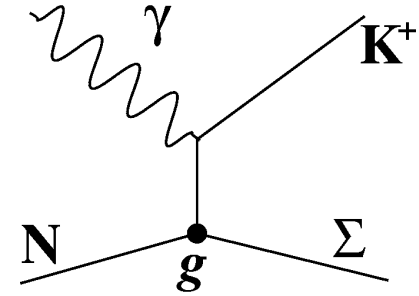


$$R_{\Lambda} = \frac{\frac{d\sigma}{d\Omega}(\gamma d \rightarrow K^+ \Lambda)}{\frac{d\sigma}{d\Omega}(\gamma p \rightarrow K^+ \Lambda)}$$

$$R_{\Sigma} = \frac{\frac{d\sigma}{d\Omega}(\gamma d \rightarrow K^+ \Sigma)}{\frac{d\sigma}{d\Omega}(\gamma p \rightarrow K^+ \Sigma^0)}$$

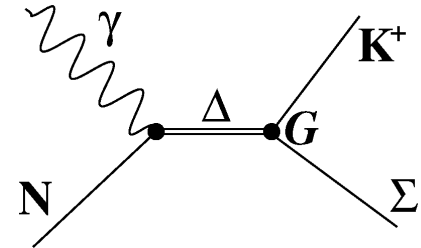


t-channel:

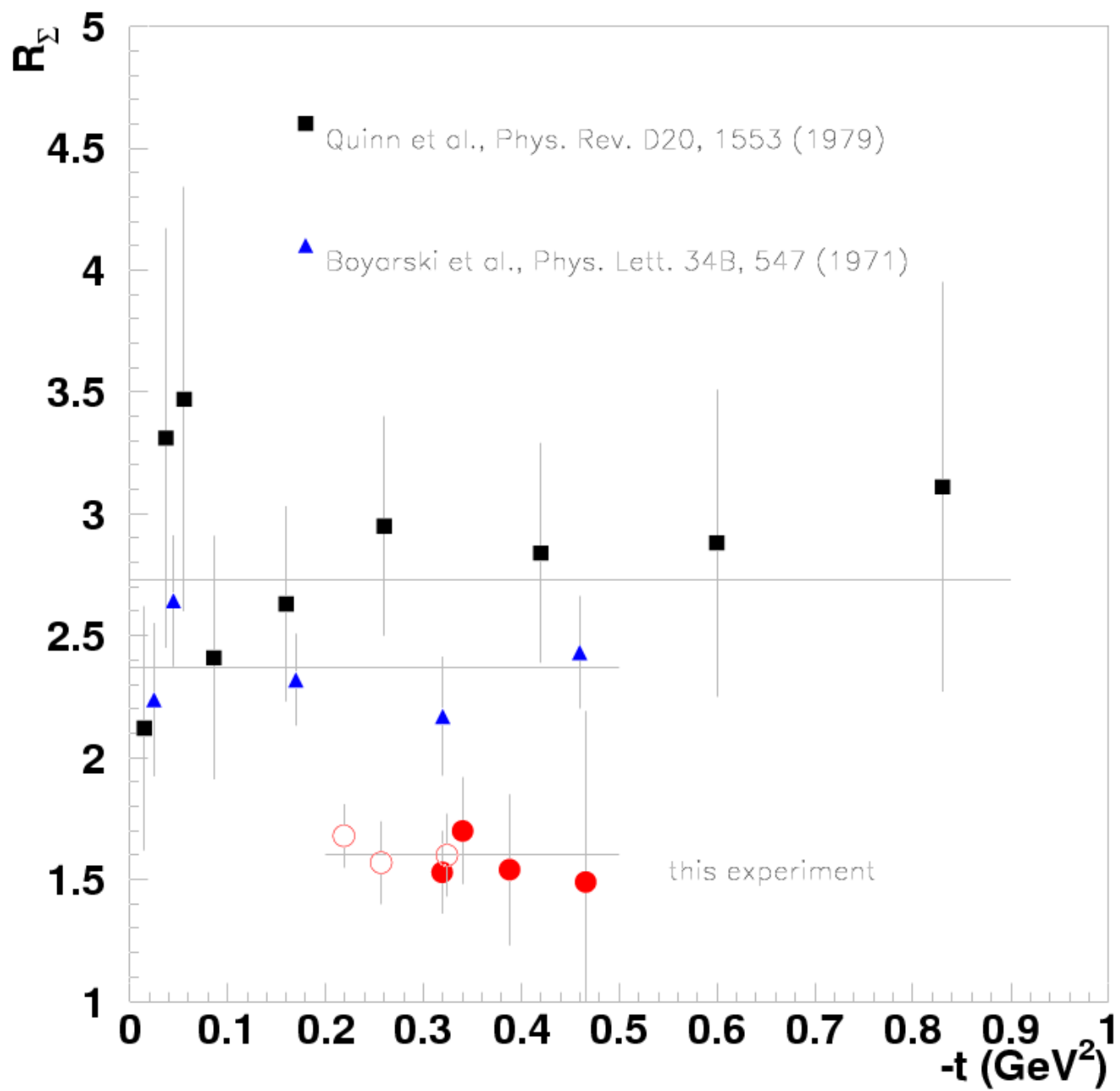


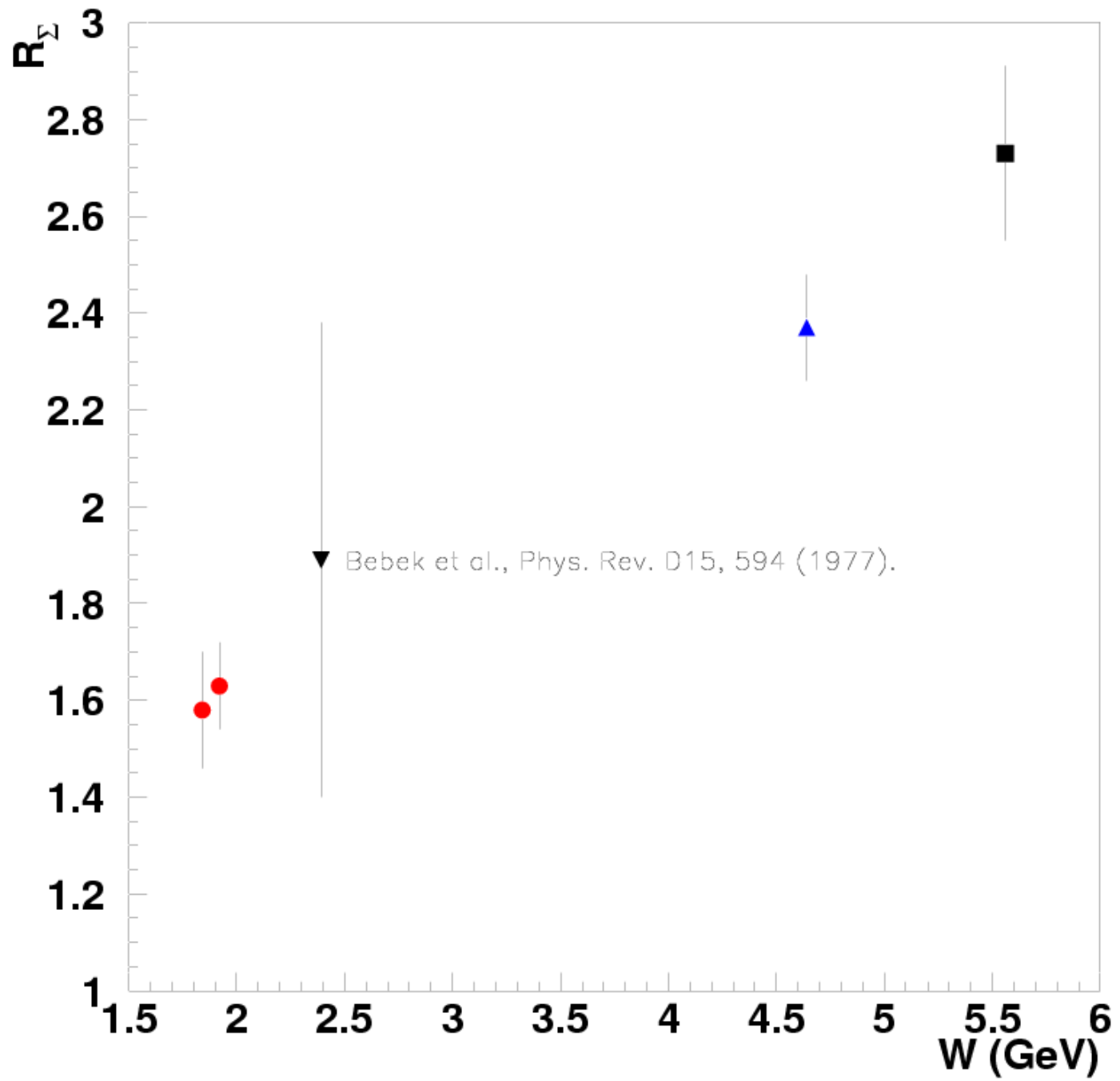
$$g_{K^+\Sigma^-n} = \sqrt{2}g_{K^+\Sigma^0p} \quad R_{\Sigma} = 3$$

s-channel:



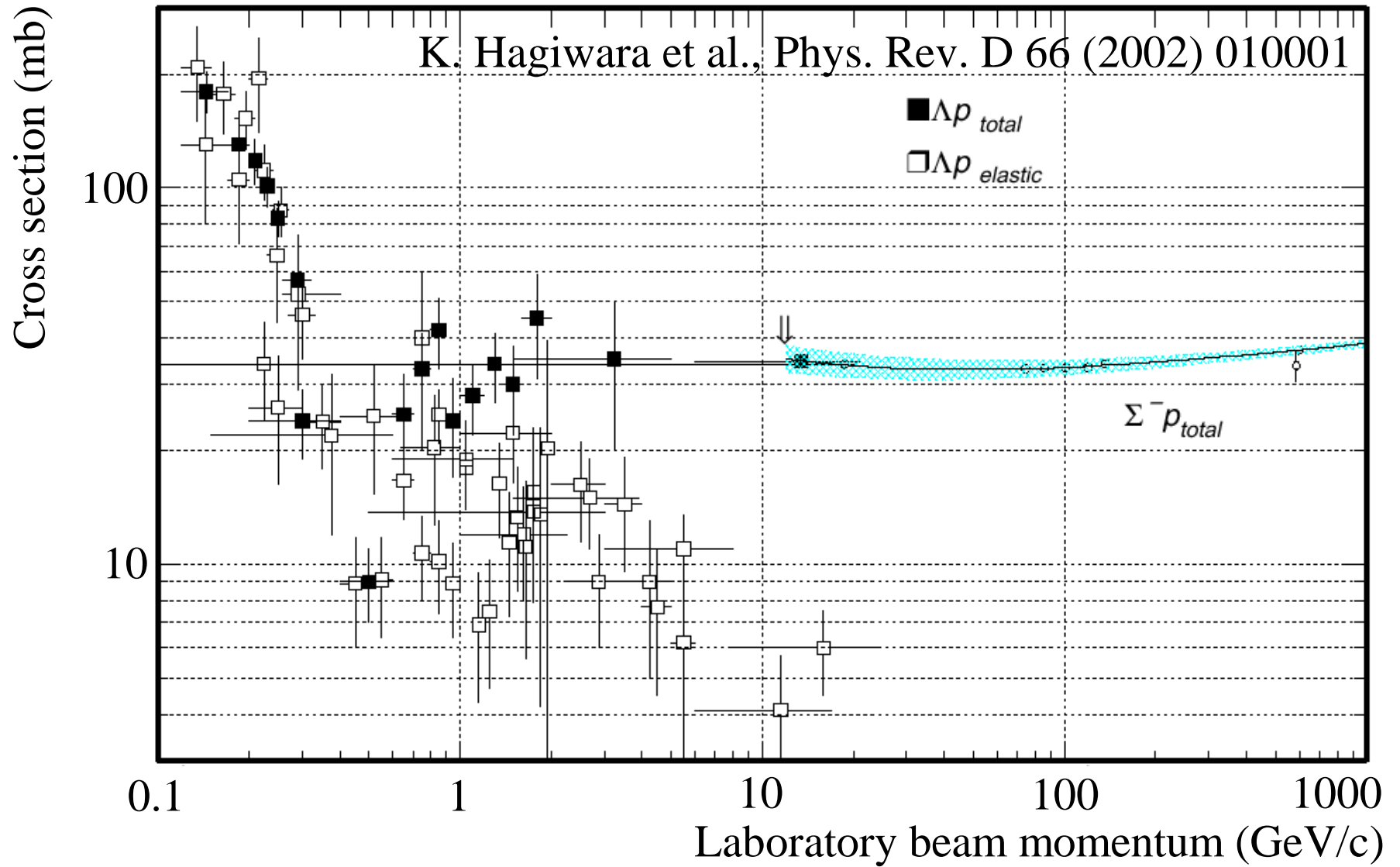
$$G_{K^+\Sigma^-\Delta^0} = G_{K^+\Sigma^0\Delta^+}/\sqrt{2} \quad R_{\Sigma} = 1.5$$



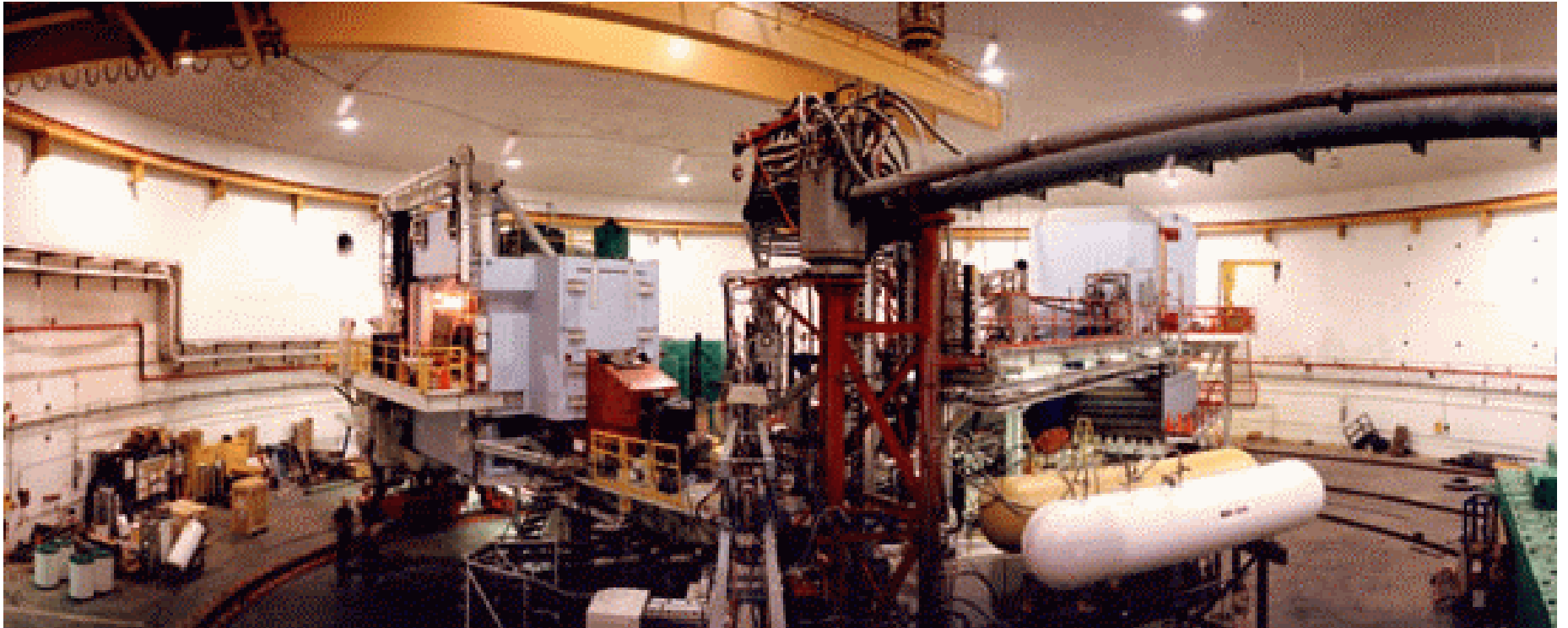


Lambda-Proton Cross Section

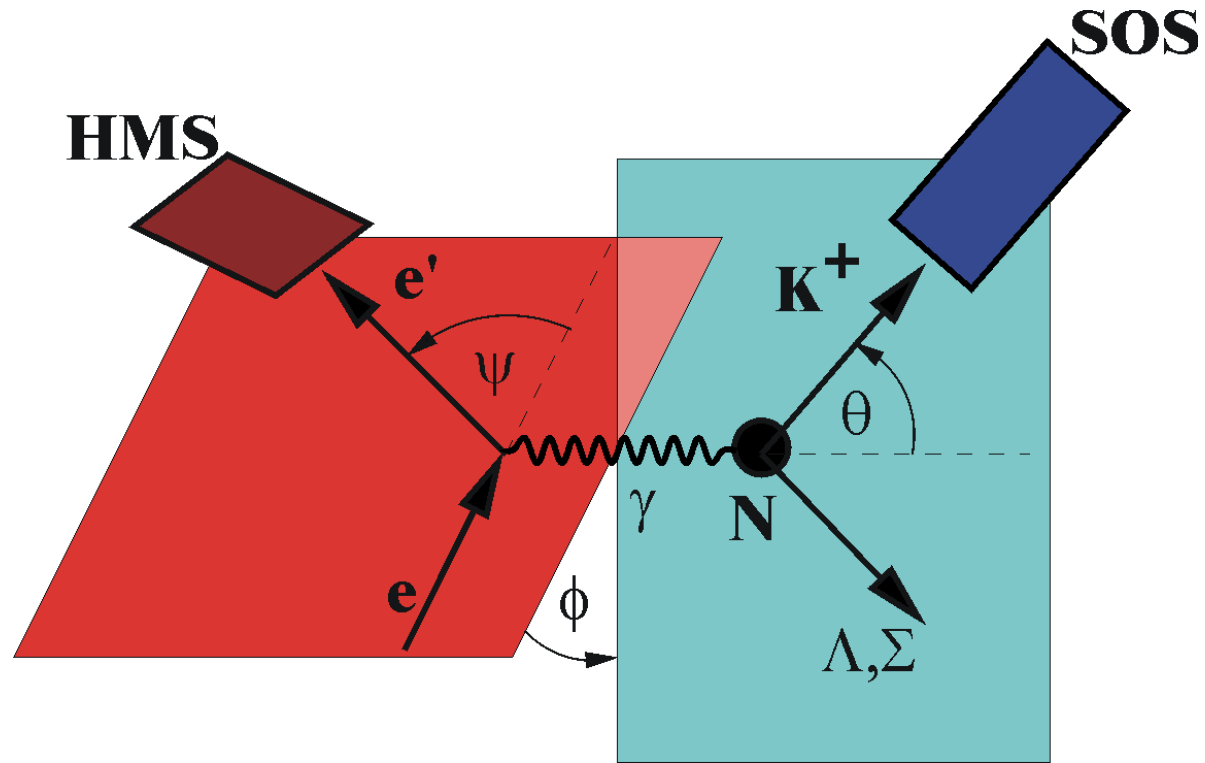
<http://pdg.lbl.gov/>



TJNAF Hall C



Kaon Electroproduction Cross Section



Virtual photon -
nucleon CM
cross section

$$\frac{d\sigma}{d\Omega_{CM}} = \sigma_T + \sigma_L + \sigma_{TT} \sin^2 \theta \cos(2\theta) + \sqrt{\frac{\sigma(\sigma+1)}{2}} \sigma_{LT} \sin \theta \cos \theta$$

Virtual photon
polarisation

$$e = \left[1 + 2 \frac{|q|^2}{Q^2} \tan^2 \left(\frac{y}{2} \right) \right]^{-1}$$

Particle Identification

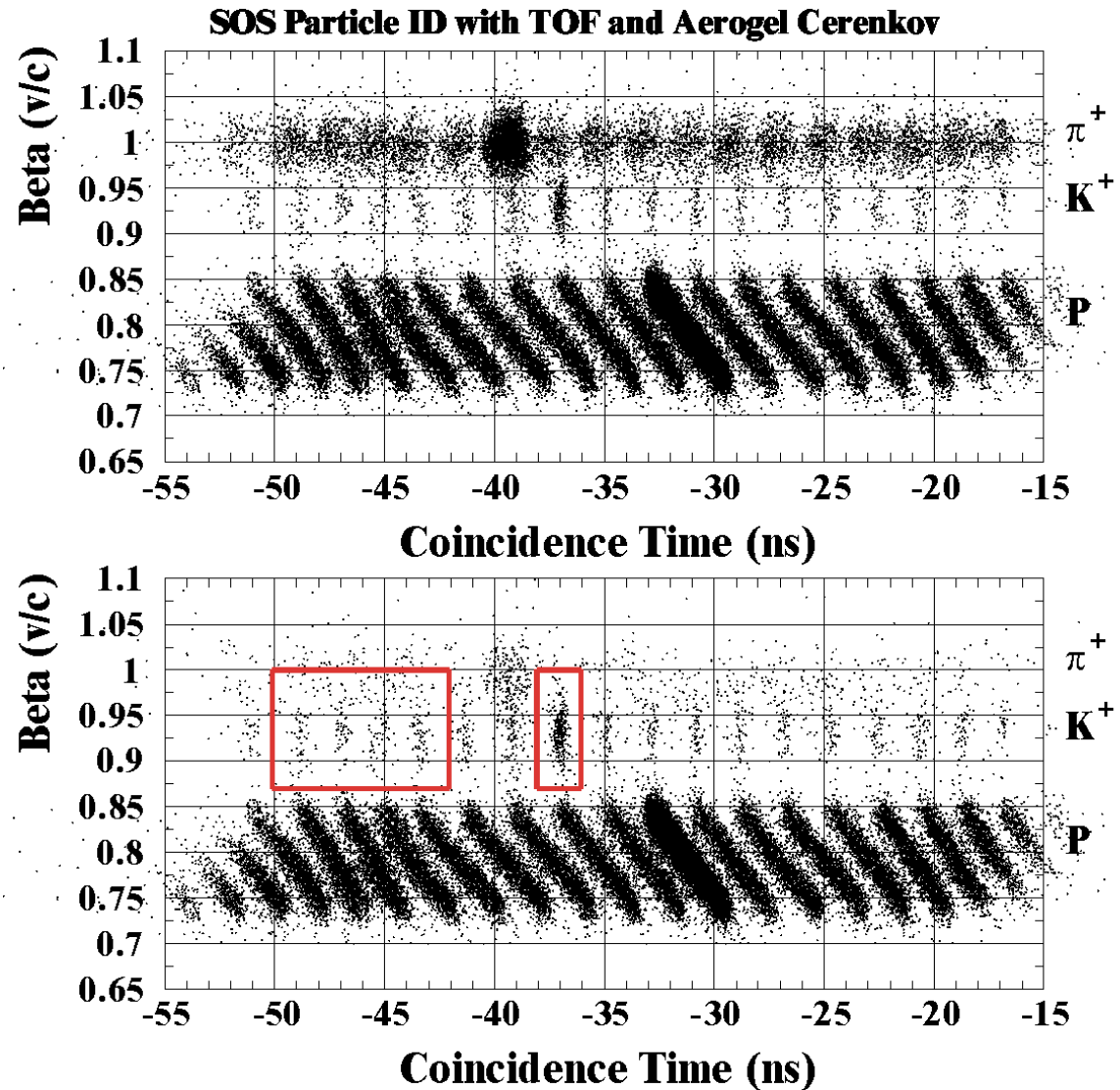
K⁺ PID:

Coincidence time cuts
separate 99.9 % real K⁺, π^+ , p

Aerogel cuts reject 98 % π^+
b cuts reject 99 % p

e⁻ PID

Cerenkov and calorimeter have
~ 99.8 % efficiency for electron
PID



Simple Model for Hydrogen Data

z Available models

y Adelseck, Saghai PRC42 (1990) 108

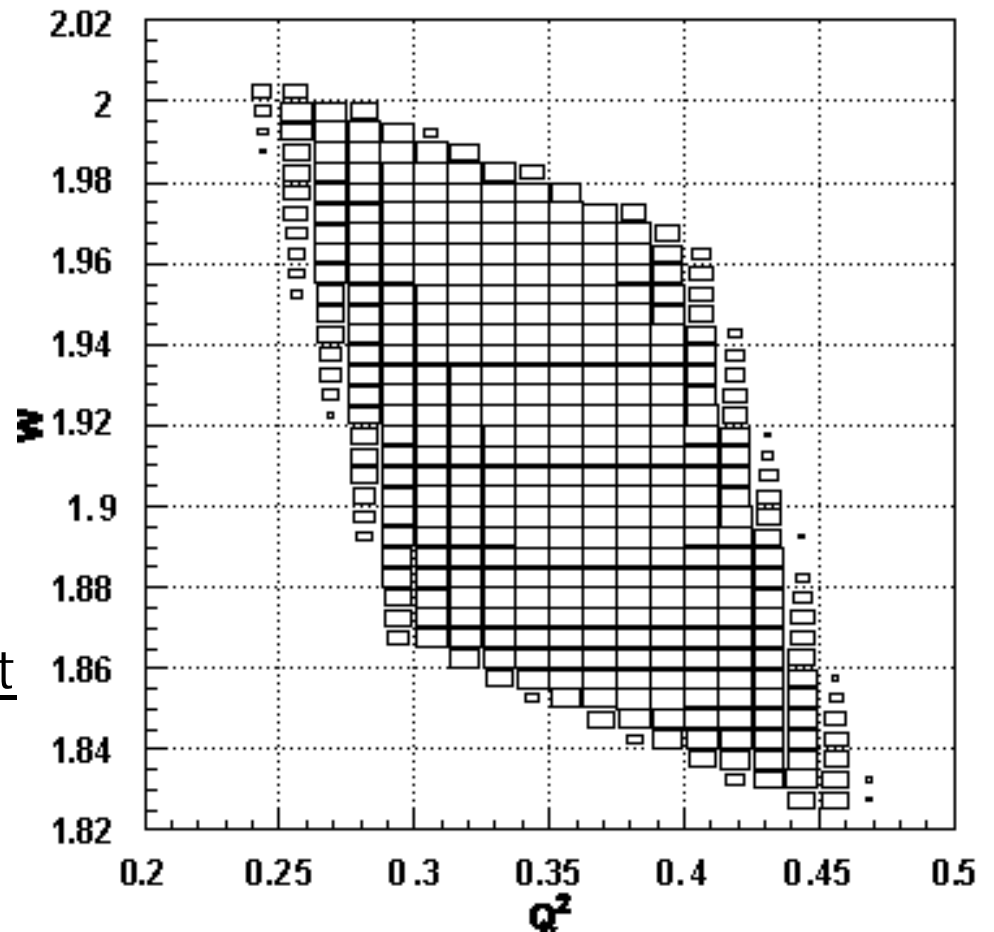
y Williams et al PRC46 (1992) 1617

do not describe the data

z Use instead simple factorization Ansatz which describes the data over our acceptance

$$\frac{dS}{d\Omega} \propto f(Q^2)g(W)h(t)i(j)$$

z Model describes 1996/1999 data well and is also used for all nuclear targets



W dependence of ${}^1\text{H}(e,e'K)\Lambda$ scaled to t_{min}

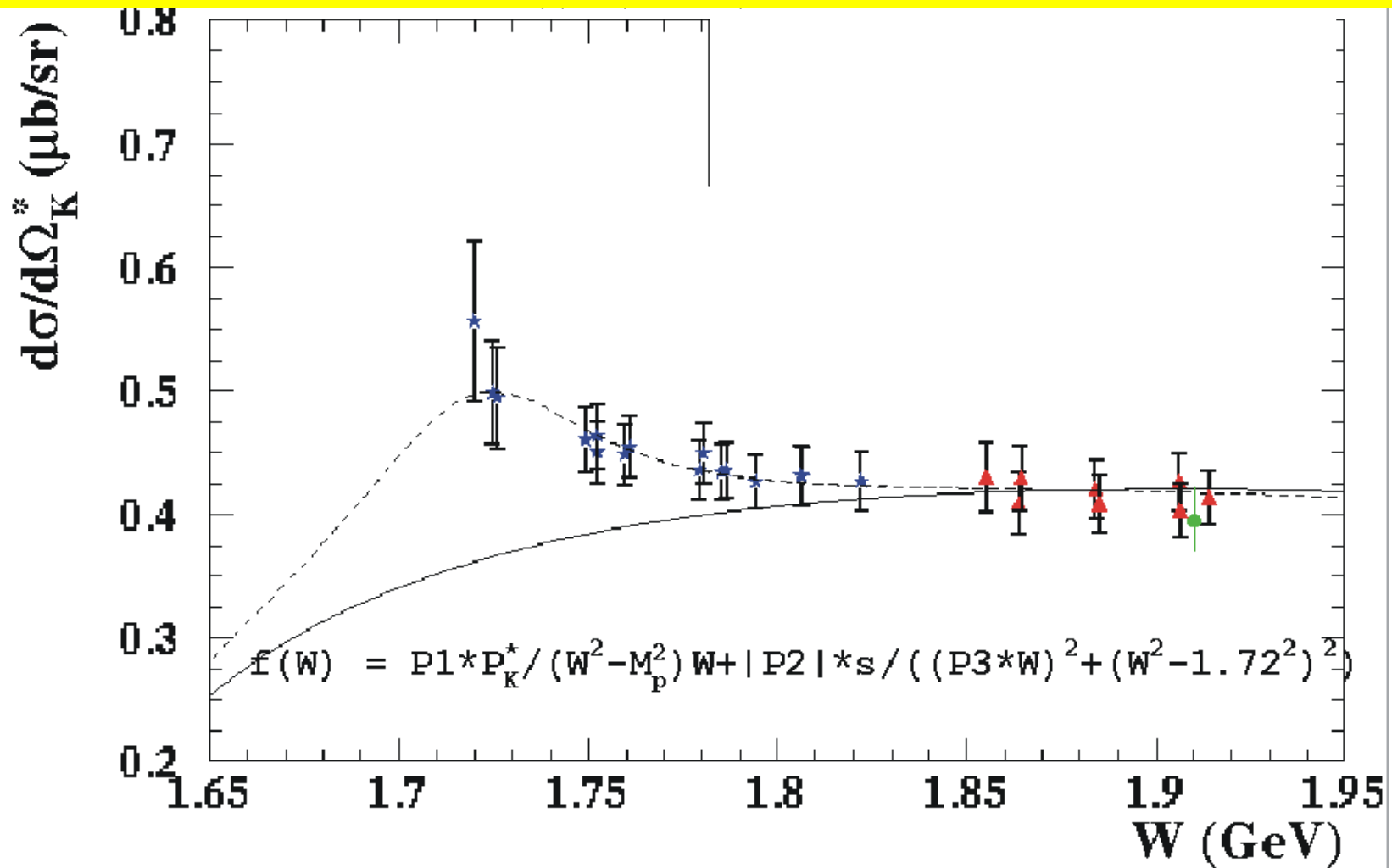
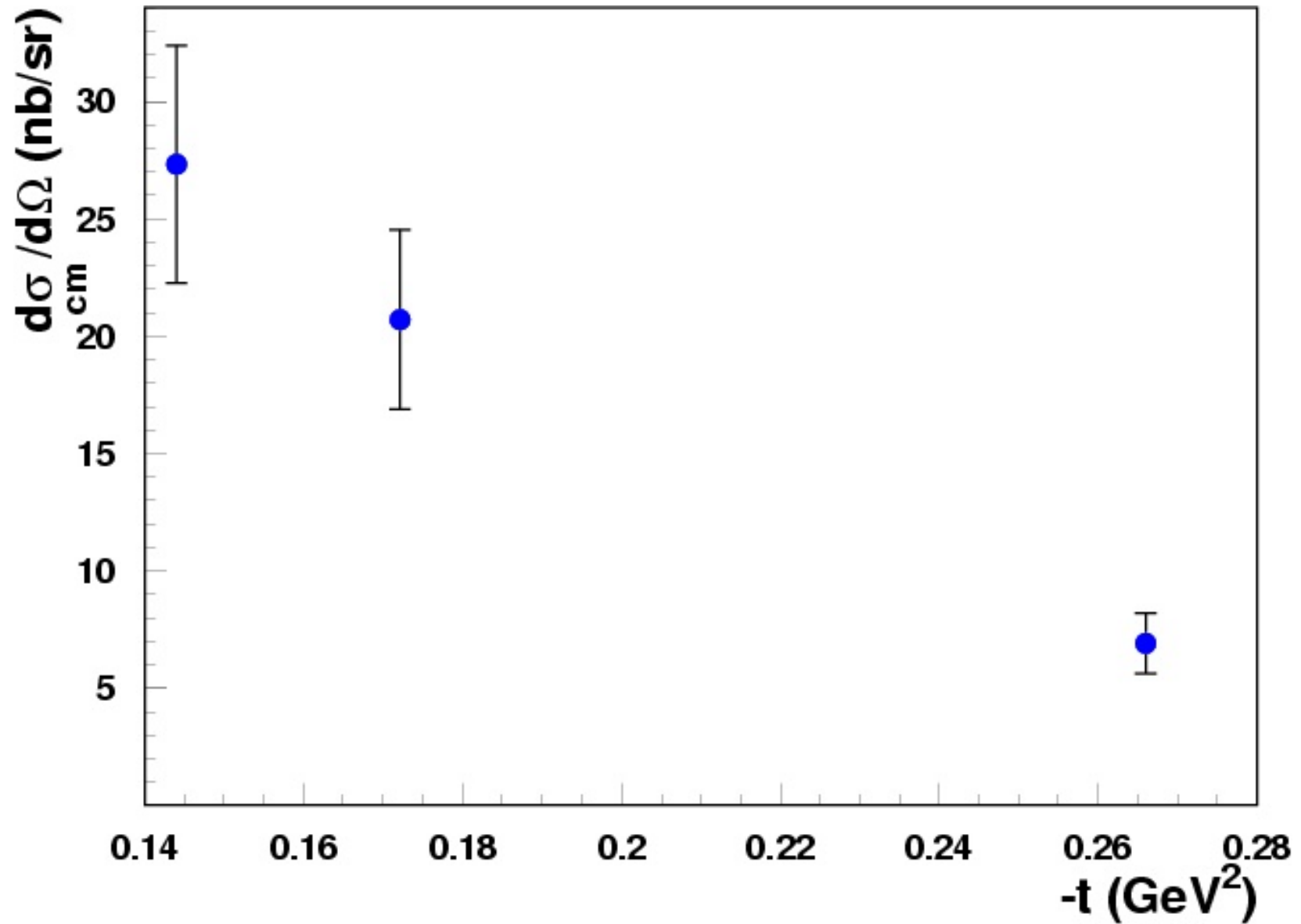


Figure from J. Cha Hampton Univ. 2000.

The data are scaled in $\theta_{\nu K.c.m}$ and Q^2 .

Bound Λ -Hypernuclei $A=4$



Elementary Cross Sections for ${}^3\text{He}$

