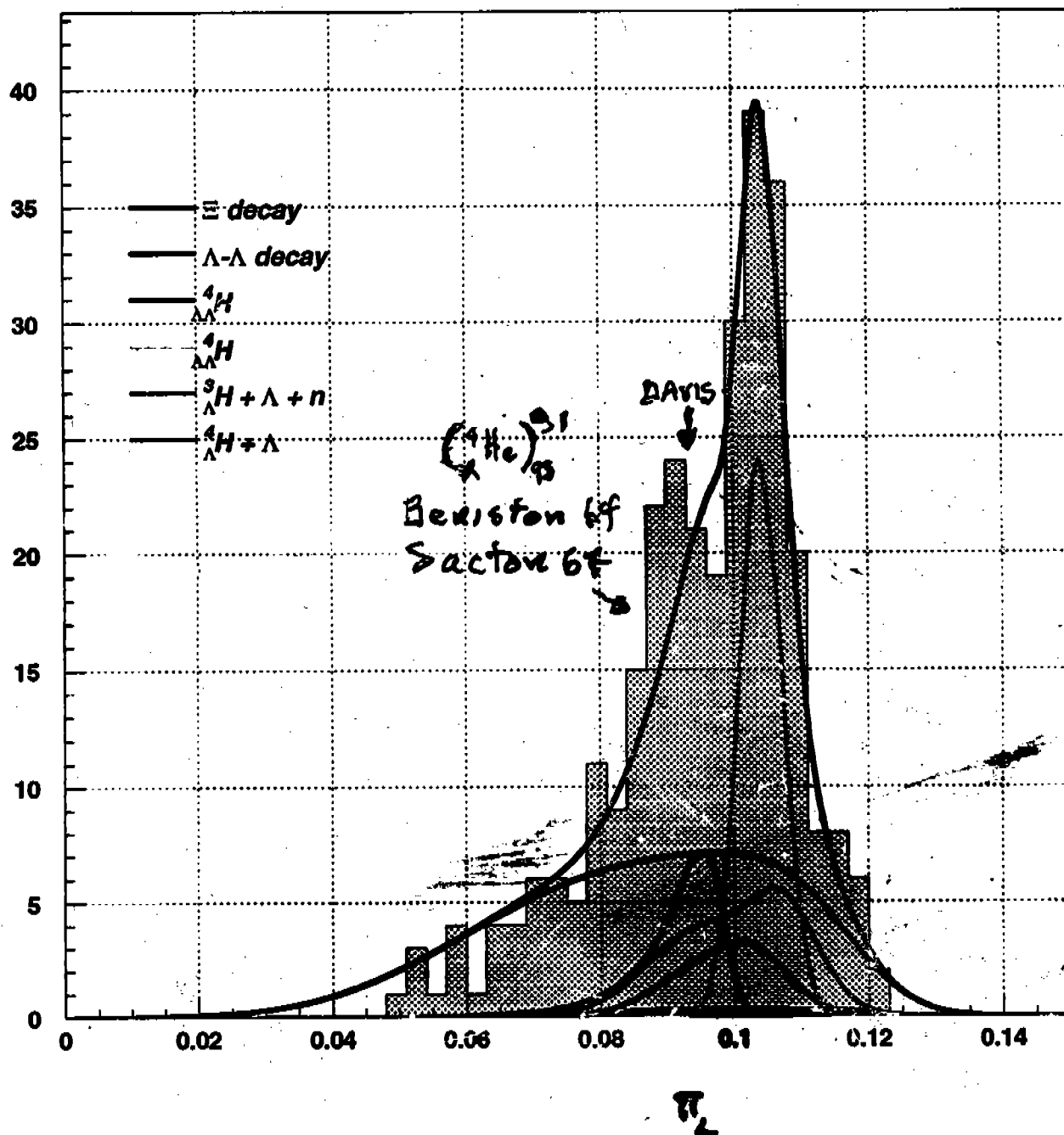


# E906 and Resonant Mode of Decay

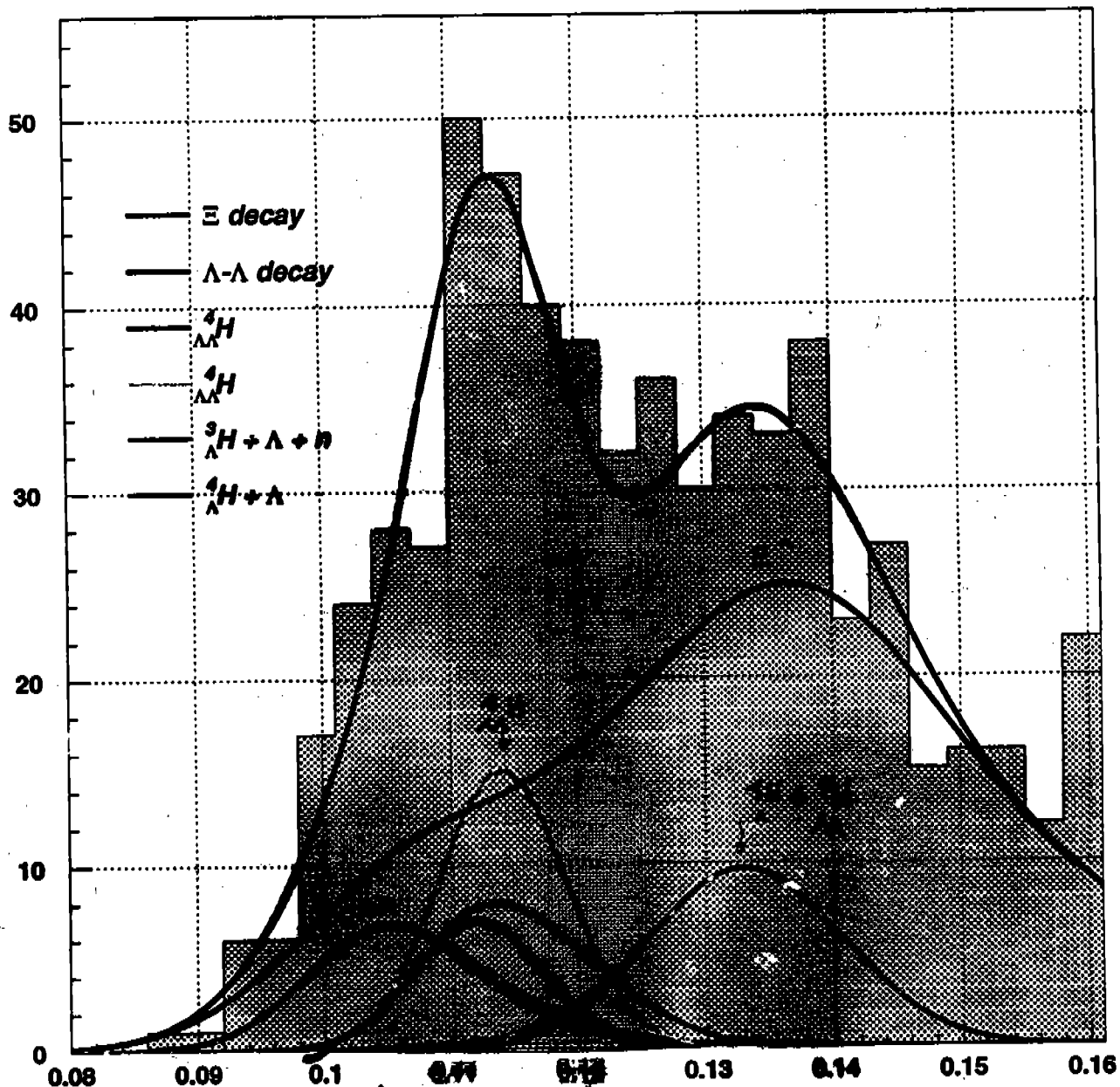
w. D.E.Kahana and D.J.Millener

- In *E906* at the BNL-AGS, strong evidence for the nuclide  ${}_{\Lambda\Lambda}^4\text{H}$  was found.
- The most striking components of this data was the appearance of a narrow low-momentum  $\pi^-$  line at  $k_\pi = 104 - 105$  MeV/c, ascribed to the decay into a resonant state in  ${}_{\Lambda}^4\text{He}$ .
- The existence of such a state is shown to be plausible and its characteristics delineated.



(Correlated pion pair  $\pi_{+}, \pi_{-}$ )

Gen. 10'



↑  
cut.

↑

2

## Best Evidence for $4_{\Lambda\Lambda}H$

- Structure at 104-105 MeV/c narrow, unreported.
- Structure at 114-116 broad, **two components**

*Correlated*



$\alpha$ - $\Lambda$  potential on the basis of a nonrelativistic quart-cluster model. This would stimulate theoretical investigations about the deeper origin of the central repulsion.

### 3.1.2 Decay spectrum

The energy spectrum of  $\pi^-$  is shown in Fig. 12 for the  ${}^5\text{He}$  case together with the experimental data by Gajewski *et al.* (1969). Again a good agreement is obtained between the theoretical result and the data. In the Green's function calculation a smearing with  $\Delta E = 0.5$  MeV was done to compare the result with the spectrum data which was obtained by counting emulsion events with 0.5 MeV energy interval.

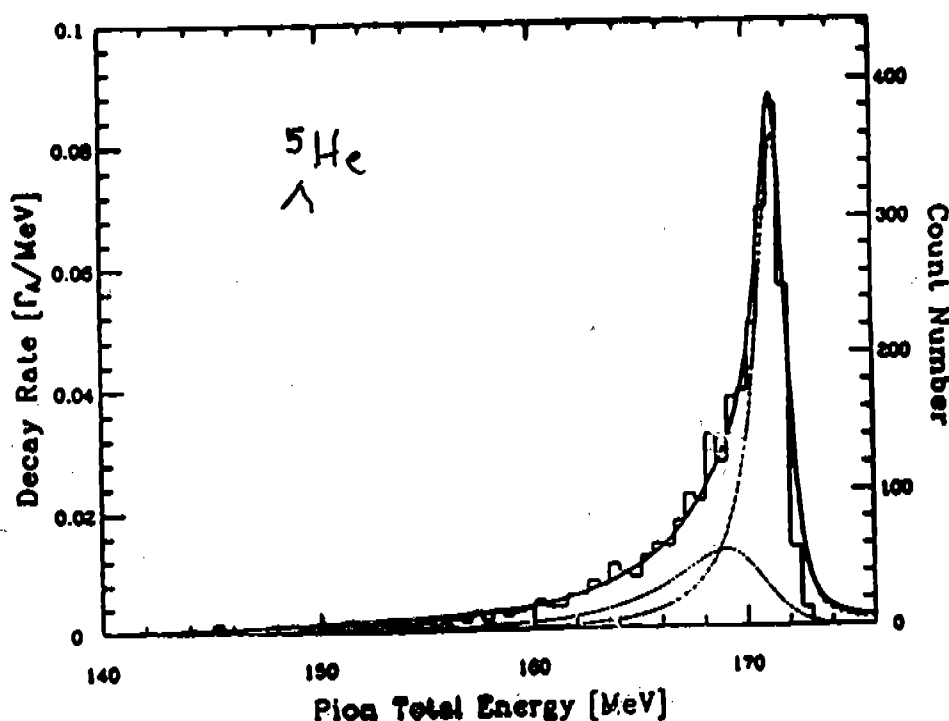


Fig. 12. The  $\pi^-$  energy spectrum of  ${}^5_\Lambda\text{He}$ . The dash-dotted and dotted curves are the  $p_{3/2}$  and  $p_{1/2}$  contributions, respectively.

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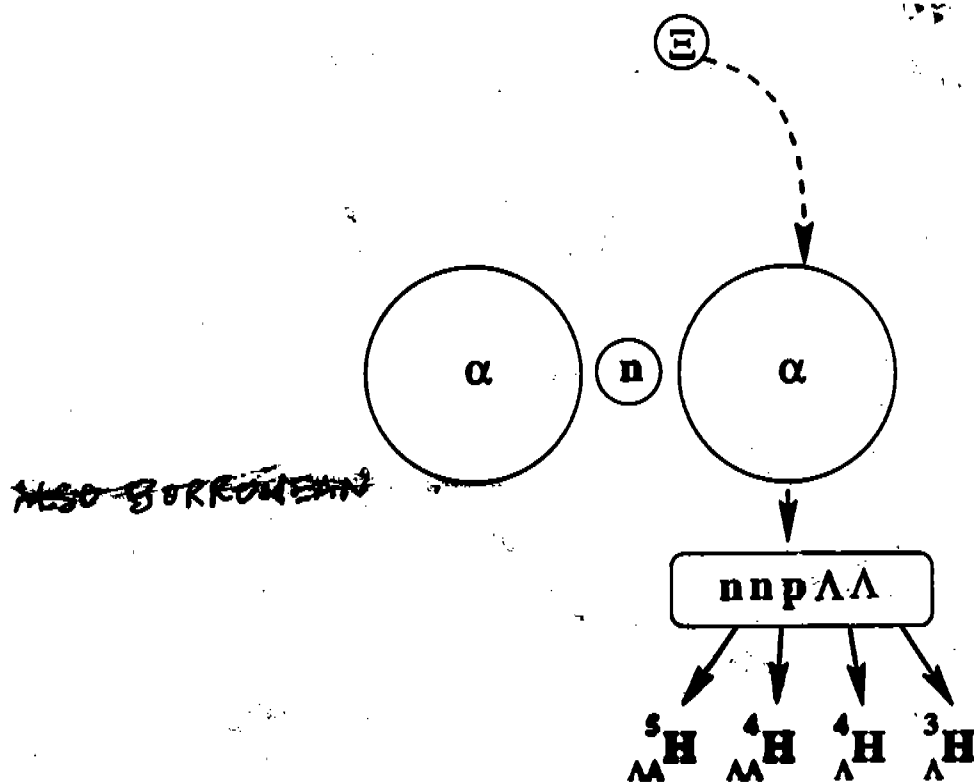
Progress in  
Particle & Nuclear  
Phys.

(1997)

# REACTION MECHANISM

12  
17

**Likely scenario, consistent with observed species:**



Nuclear events  
initiated  
with appreciable  $T_{\text{eff}}$   
(initially slower)

**Other possible species:**



↓  
non-meson  
decays

hurt

## ${}^4_{\Lambda\Lambda}H$ Decay Modes

$${}^4_{\Lambda\Lambda}H \rightarrow {}^4_{\Lambda}He + \pi_H^- \quad (\sim 112 - 118 \text{ MeV}/c) \quad (1)$$

$${}^4_{\Lambda}He \rightarrow {}^3H + p + \pi_L^- \quad (\sim 85 - 95 \text{ MeV}/c) \quad (2)$$

and also, a decay into a possible excited state of  ${}^4_{\Lambda}He$ ,

$${}^4_{\Lambda\Lambda}H \rightarrow {}^4_{\Lambda}He^* + \pi_L^- \quad (\sim 104 - 105 \text{ MeV}/c) \quad (3)$$

$${}^4_{\Lambda}He^* \rightarrow {}^3_{\Lambda}H + p \quad (4)$$

$${}^3_{\Lambda}H \rightarrow {}^3He + \pi_H^- \quad (114.3 \text{ MeV}/c), \quad (5)$$

where  $\pi_H^-$  and  $\pi_L^-$  refer to the high and low momentum members of a correlated pair seen in the experiment.

CONCENTRATE





$\sim 8.9$  \_\_\_\_\_ **(0-2) Resonance**

7.78 - - - - -  $p + d + \Lambda$

7.75 - - - - -  ${}^3_\Lambda\text{H} + p$

2.39 - - - - -  ${}^3\text{He} + \Lambda$

1.15 \_\_\_\_\_  $1^+$

0.00 \_\_\_\_\_  $0^+$

**${}^4_\Lambda\text{He}$  Spectrum**

## Resonance Candidates and Weak Decay

$$\sqrt{\frac{\mu}{3+\mu}} (s^2 p)[3]1/2 \times s_\Lambda - \sqrt{\frac{3}{3+\mu}} s^3 \times p_\Lambda \quad (1)$$

$$(s^2 p)[21]1/2 \times s_\Lambda \quad (2)$$

$$(s^2 p)[21]3/2 \times s_\Lambda \quad (3)$$

$$\begin{aligned} {}^4_{\Lambda}\text{H} &\equiv S^2(1) S^2_\Lambda(0) \quad S=1 \\ \therefore \text{ONLY } {}^4_{\Lambda}\text{H} & \quad L=1, S=1 \quad J=0, 1, 2 \end{aligned}$$

Matrix elements for the production of single- $\Lambda$  hypernuclear configurations in the  $\pi^-$  weak decay of  ${}^4_{\Lambda\Lambda}\text{H}$  are given in the first row in units of  $s^2_\pi - \langle l_N | j_{l_N}(k_\pi r) | s_\Lambda \rangle^2 / 4\pi$  where  $\mu/(3+\mu) = 0.284$ . The remaining rows specify the breakup of the single- $\Lambda$  hypernuclear configurations under the assumptions described in the text. Combining production and decay in this simple approach shows that  ${}^3_\Lambda\text{H}+p$  is favored over  $d+p+\Lambda$  by a factor of 11/3 to 4/3.

Final state	${}^4_\Lambda\text{He}(1^+)$	Eq. (1)	Eq. (2)	Eq. (3)
Production	1	$\mu/(3+\mu)$	1	4
${}^4_\Lambda\text{He}(1^+)$	1			
${}^3\text{He}+\Lambda$		1		
${}^3_\Lambda\text{H}+p$		$(\frac{1}{2})$	1/9	8/9
$d+p+\Lambda$		$(\frac{1}{2})$	8/9	1/9

$$\mu = m_\Lambda / m_N$$

## Average (cluster) ${}^3_{\Lambda}\text{H}$ -Proton Potential

$$V^{Surf}(r) = -V_0 \left[ \frac{4e^{(r-R)/a}}{(1 + e^{(r-R)/a})^2} \right], \quad (1)$$

with  $R$  the radius and  $a$  the diffusivity. A Coulomb potential is included, taken as that of a uniform charge distribution with radius parameter  $R_C = R$ .

Only the strength of  $V^{Surf}(r)$  is varied, to reproduce the correct ground state separation energy 7.75 MeV, yielding a depth:

$$V_0 = 28.09 \text{ MeV}, \quad (2)$$

with radius parameter  $R = 2.07$  fm, and diffusivity  $a = 0.5$  fm.

In this completely specified well, a  $p$ -wave resonance appears at;

$$\epsilon = [1.10 - (1/2)1.00] \text{ MeV}. \quad (3)$$

↓ N.B. SAME DEPTH FOR  ${}^3_{\Lambda}\text{H}^{0.5}$  &  ${}^3_{\Lambda}\text{H}^*$  ↓

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REPRODUCE (WEAK DECAY) RESULTS OF KUNAGAI-FUSE & OKADA

TO  $\underline{\lambda_{H_2}(1)} \propto$  TO  $\underline{{}^3H + p}$  :  $\underline{0.681\pi_\lambda \text{ \& } 0.80\pi_\lambda}$

USING  
ESSENTIALLY  $\underline{S_{\pi} = \lambda_\lambda(k_{\pi} r) Y_\lambda}$

Dominant piece, no nucleon  
spin

SUFFICIENT  $\underline{{}^3H}$  produced for scenario to work

STATE IN Q5 DOESN'T MIX MUCH WITH (2)

(1) OFF-DIAGONAL MATRIX ELEMENTS  $\sim$  spin-spin of  $\Lambda N$   
 $\sim$  singlet-triplet  
 (both attractive)

(2)  $(S^2 p) [3]$  purely spurious: no NN overlap with  
 $(S^2 p) [2]$  by transl. invariant force

## Decay Momentum

With the resonance energy calculated the  $\Lambda\Lambda$  pairing energy  $\Delta B_{\Lambda\Lambda}$  can be estimated from the position of the narrow  $\pi^-$  peak ascribed to the weak decay from  ${}^4_{\Lambda\Lambda}H$  to  ${}^4_{\Lambda}He^*$ . The  $\pi^-$  momentum  $k^*$  is given by

$$\epsilon_{\pi} + k^{*2}/2M({}^4_{\Lambda}He^*) = M({}^4_{\Lambda\Lambda}H) - M({}^4_{\Lambda}He^*). \quad (1)$$

To a good approximation here,

$$k_{\pi}^* = (107.466 - 1.6391\Delta) \text{ MeV}/c, \quad (2)$$

where  $\Delta = B^* + \epsilon_R$  and

$$B^* = 2\bar{B}_{\Lambda}({}^3_{\Lambda}H) + \Delta B_{\Lambda\Lambda}$$

$$\left. \vphantom{\begin{matrix} B^* \\ \Delta B_{\Lambda\Lambda} \end{matrix}} \right\} \rightarrow \epsilon_R + \Delta B_{\Lambda\Lambda}$$

is the full binding energy of the  $\Lambda$  pair in  ${}^4_{\Lambda\Lambda}H$ .

## $\Delta B_{\Lambda\Lambda}$ -Pairing Energy

Taking

$$\epsilon_R = 1.18 \text{ MeV}, \quad (1)$$

$$\Delta B_{\Lambda\Lambda} = 0.34 \text{ MeV}, \quad (2)$$

puts the centroid of the decay momentum at

$$k^* = 104.5 \text{ MeV}/c. \quad (3)$$

Alternatively,  $\epsilon_R = 0.97 \text{ MeV}$  and the same decay momentum results in  $\Delta B_{\Lambda\Lambda} = 0.55 \text{ MeV}$ .

Assuming 300 KeV uncertainty in  $k^*$  produces range:

$$0.17 \text{ MeV} < \Delta B_{\Lambda\Lambda} < 0.55 \text{ MeV}. \quad (4)$$

Interestingly, the two body decay of  ${}^4_{\Lambda\Lambda}H$  for  $\Delta B_{\Lambda\Lambda} = 0.34$  now appears at

$$k_\pi = 116.5 \text{ MeV}/c, \quad (5)$$

Easily distinguishable in a future (better resolution) experiment.

## Comments

(1). We see that the existence of states above threshold in  ${}^3_\Lambda\text{H} + \text{proton}$  system can straightforwardly describe the narrow, low momentum,  $\pi^-$  feature at  $104 - 105 \text{ MeV}/c$  observed in the BNL experiment E906.

(2). The inferred energy of the resonance and energy range possible for  $\Delta B_{\Lambda\Lambda}$ , the latter likely somewhat less than  $0.5 \text{ MeV}$ , are not unreasonable. Takahashi *et al.* found  $\Delta B_{\Lambda\Lambda} = 1.0 \pm 0.38 \text{ MeV}$  for  ${}^6_\Lambda\text{He}$ . One certainly expects a smaller value for the more extended mass 4 system, but the suggested values and their likely errors allow for consistency.

(3). For the meantime, it is also of importance to perform many-body calculations with  $\Sigma - \Lambda$  coupling, and with forces constrained by as much data as is available, including E906. The latter experiment has pointed to the existence of two interesting nuclides, both mass 4 objects, one a very light  $S = -2$  hypernucleus, the other an unusual, if not completely unexpected, resonance in an  $S = -1$  daughter nucleus.

(4).  $B(S = -2) < 4.5 \text{ MeV}$ ,  $B(H^0) \leq 1.5 \text{ MeV}$ . No bound  $H^0$ ?