

IN-MEDIUM EFFECTS ON THE K^-/K^+ RATIO AT GSI

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Outline of Talk

1. K^-/K^+ ratio: medium effects?
2. Thermal models: K^+/K^- ratio
3. In-medium properties of K^-
4. Results for K^+/K^- ratio at GSI
5. Conclusions

K^-/K^+ ratio: medium effects?

HEAVY-ION COLLISIONS at GSI (1-2 AGeV):

 K^- in hot and dense matter^a

- K^-/K^+ ratio: enhancement at energies below threshold (pp \leftrightarrow nucleus-nucleus)

$$NN \rightarrow K^+ \Lambda N \quad 1.58 \text{ GeV}$$

$$NN \rightarrow K^+ K^- NN \quad 2.5 \text{ GeV}$$

- $K^-/K^+ \sim ct$ for C+C, Ni+Ni, Au+Au

Absorption K^- via $K^- N \rightarrow Y \pi$ suppressed
and/or enhanced K^- production

- Centrality independence of K^-/K^+ ratio
 - K^+ and K^- yields strongly correlated^b
 - Thermal models: free particles and volume cancels out exactly.^c
 - “Broad-band equilibration” by Brown et al.^d

^a Oeschler, *JPG* 28, 1787 (2002); Senger, *NPA* 685, 312 (2001);
Sturm et al., *JPG* 28, 1895 (2002)

^b Hartnack et al., *PRL* 90, 102302 (2003)

^c Cleymans et al., *PRC* 60, 054908 (1999)

^d Brown et al., *NPA* 690, 184c (2001); *NPA* 698, 483c (2002)

Thermal models: K^+/K^- ratio

Hadronic Gas Model (μ_B, T)^a

using CANONICAL scheme for STRANGENESS

K^+/K^- independent of VOLUME



independent CANONICAL OR
GRAND-CANONICAL treatment

$$\frac{K^+}{K^-} = \frac{Z_{K^+}^1}{Z_{K^-}^1} \frac{Z_{K^-}^1 + Z_{B,S=-1}^1 + Z_{M,S=-1}^1}{Z_{K^+}^1 + Z_{M,S=+1}^1}$$

$$K^+/K^- \sim 33 \quad \text{at } 1.93 \text{ AGeV in Ni+Ni}^b$$

Data described:^c

$$T = 70 \pm 10 \text{ MeV}, \mu_B = 720 \pm 30 \text{ MeV}$$

^a Cleymans et al., PLB 485, 27 (2000) and PRC 59, 1663 (1999)

^b Menzel et al., PLB 495, 26 (2000)

^c Cleymans et al., PRC 57, 3319 (1998)

In-medium properties of K^-

$$\frac{K^+}{K^-} = 1 + \frac{Z_{\Lambda}^1 + Z_{\Sigma}^1 + Z_{\Sigma^*}^1}{Z_{K^-}^1} \sim e^{\frac{\mu_{B^+} + U_{K^-}}{T}}$$

- Mean-field for Λ and Σ , and Breit-Wigner for Σ^*

$$Z_{\Lambda, \Sigma}^1 = g_{\Lambda, \Sigma} V \int \frac{d^3 p}{(2\pi)^3} e^{-\frac{\sqrt{m_{\Lambda, \Sigma}^2 + p^2} - U_{\Lambda, \Sigma}(\rho)}{T}} e^{\frac{\mu_B}{T}}$$

$$Z_{\Sigma^*}^1 = g_{\Sigma^*} V \int \frac{d^3 p}{(2\pi)^3} \int ds e^{-\frac{\sqrt{p^2 + s}}{T}} \frac{1}{\pi} \frac{m_{\Sigma^*} \Gamma}{(s - m_{\Sigma^*}^2)^2 + m_{\Sigma^*}^2 \Gamma^2} e^{\frac{\mu_B}{T}},$$

- K^- meson in hot and dense matter

- ON SHELL

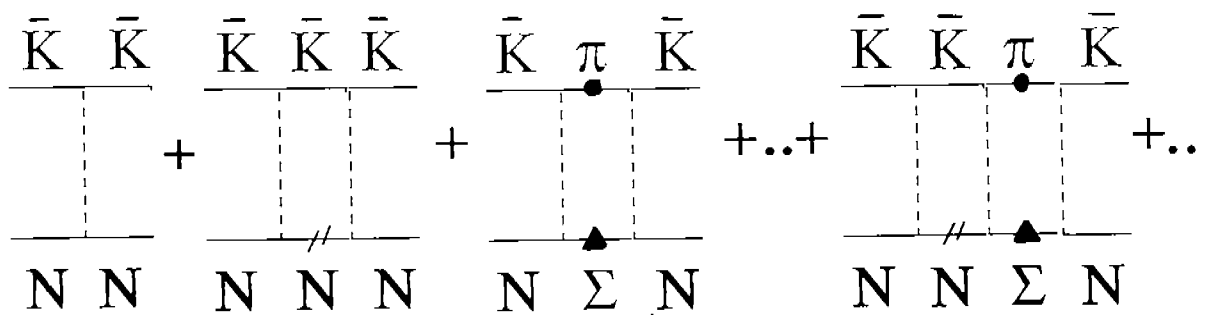
$$Z_{K^-}^1 = g_{K^-} V \int \frac{d^3 p}{(2\pi)^3} e^{-\frac{E_{K^-}}{T}}$$

$$E_{K^-} = \sqrt{m_{K^-}^2 + p^2} + \text{Re } U_{K^-}(p, E_{K^-}, \rho, T),$$

- SPECTRAL DENSITY

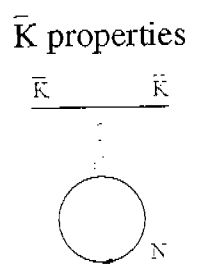
$$Z_{K^-}^1 = g_{K^-} V \int \frac{d^3 p}{(2\pi)^3} \int ds S_{\bar{K}}(p, \sqrt{s}, \rho, T) e^{-\frac{\sqrt{s}}{T}}$$

Microscopic calculation of the in-medium $\bar{K}N$ interaction (G -matrix) and the K^- optical potential from $T = 0$ to T finite in nuclear matter.



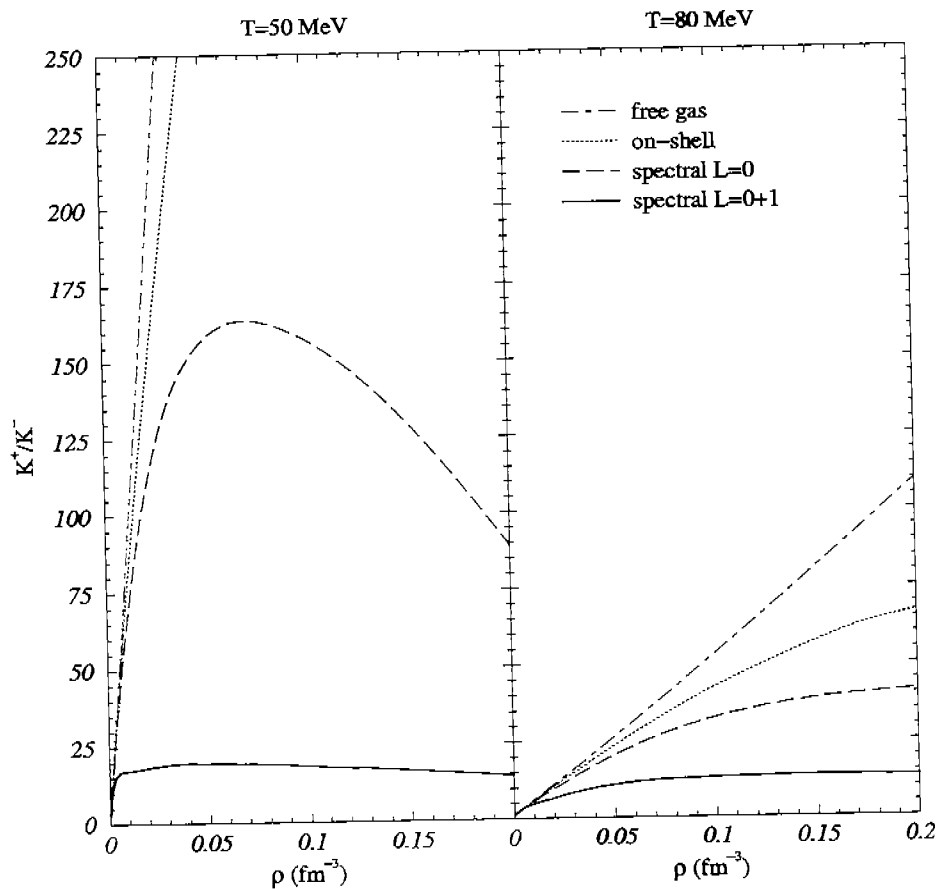
$$G = V + V \frac{Q}{\omega - H_0} V + V \frac{Q}{\omega - H_0} V \frac{Q}{\omega - H_0} V + \dots$$

$$G = V + V \frac{Q}{\omega - H_0} G \quad \text{Bethe-Goldstone eq.}$$

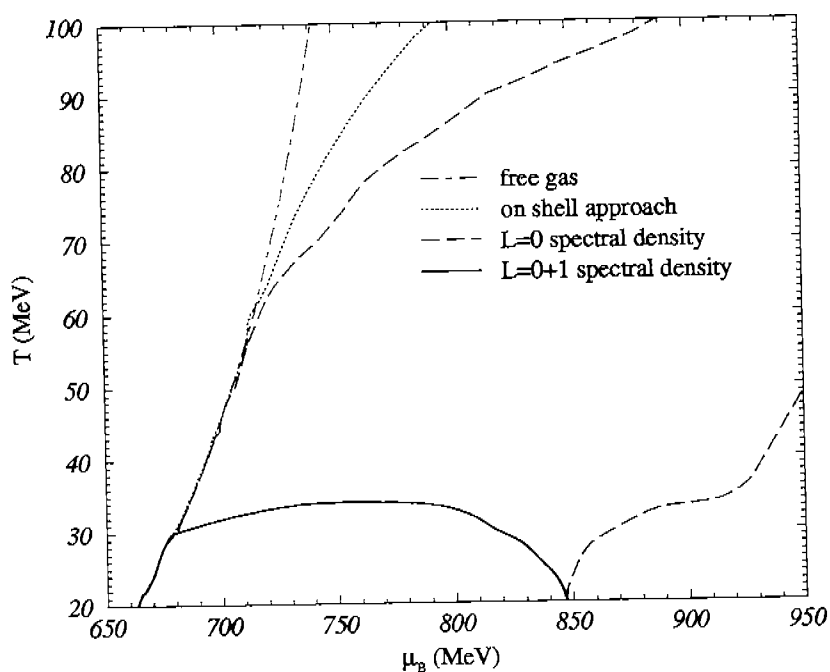


Jülich meson-exchange bare potential V

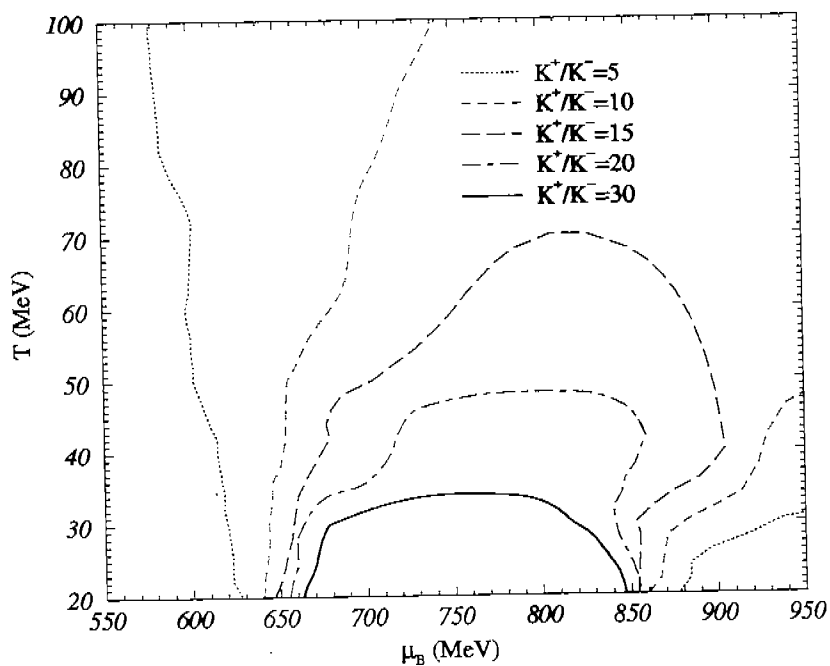
A. Müller et al., Nucl. Phys. A 513, 557 (1990)

Results for K^+/K^- ratio at GSI

- K^+/K^- ratio grows with $e^{\mu_B/T}$
- K^- properties in hot and dense matter make the curves bend down
- Inclusion of K^- spectral density implies additional attraction \rightarrow qualitative agreement with "Broad-band equilibration"



$T(\mu_B)$ for $K^+/K^- = 30$ within different approaches



Different K^+/K^- ratios using the K^- spectral density

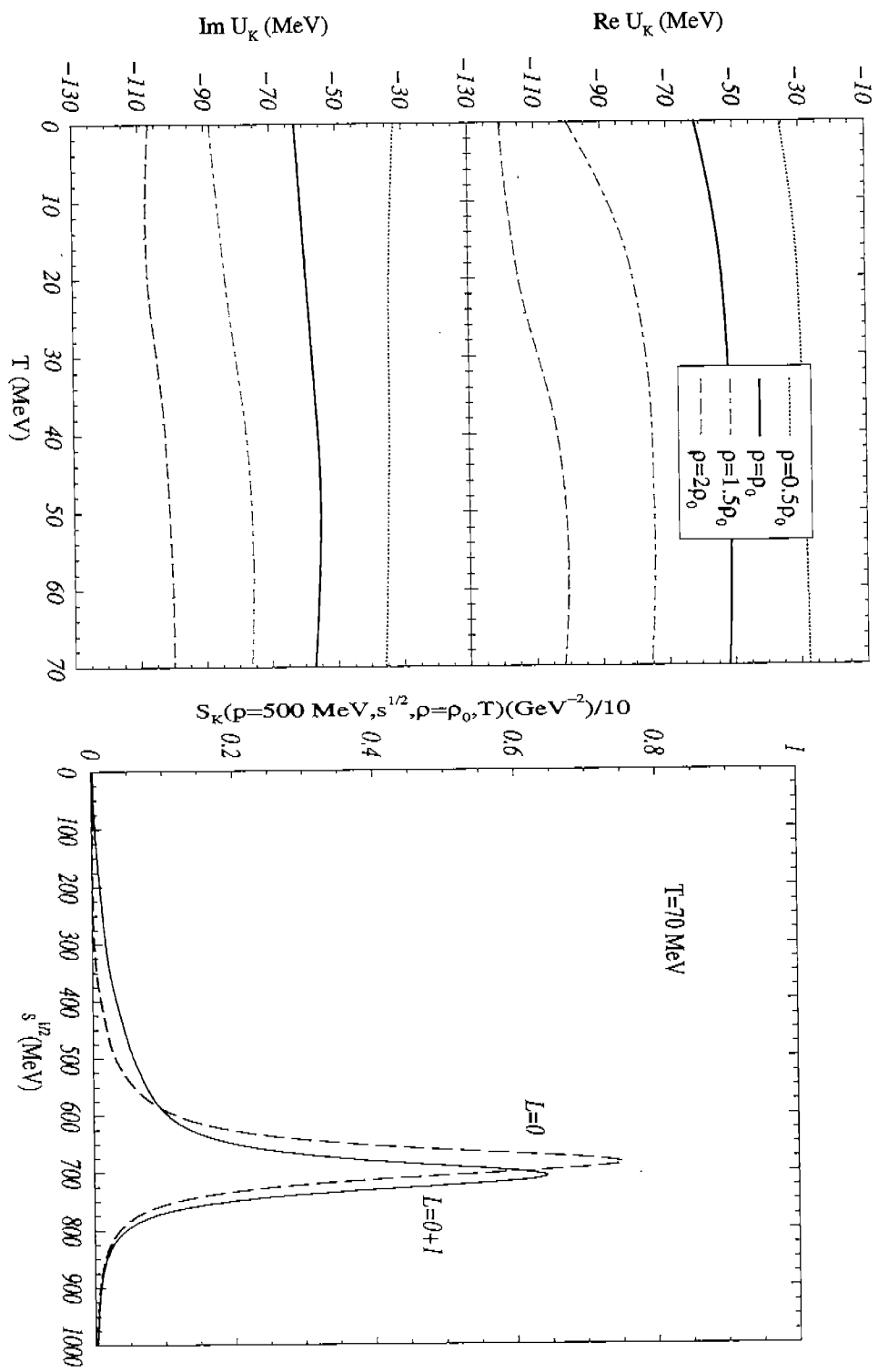
Conclusions

- $T(\mu_B)$ compatible with K^-/K^+ is very delicate depending very strongly on the approach ($T \approx 35$ MeV and $\mu_B \rightarrow 850$ MeV for $K^+/K^- \approx 30$ with the K^- spectral density).
- On-shell approach does not achieve Brown's "Broad-band equilibration" because of a stronger increase of μ_B with density and a less attractive U_{K^-} .
- "Broad-band equilibration" obtained only when the K^- spectral density is used. However, K^- is in excess \Rightarrow dynamical non-equilibrium effects to explain the smaller number of K^- detected.

Tolós, Polls, Ramos and Schaffner-Bielich,

Phys. Rev. C 68, 024903 (2003)

L. Tolós, A. Ramos, and A. Polls, *Phys. Rev. C* 65, 054907 (2002)



1. On-shell approximation(left panel): in-medium U_{K^-} at finite temperature.
2. K^- spectral density (right panel) using K^- self-energy.

