## Spin density matrix elements in $\Lambda(1520)$ photoproduction at CLAS

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 $\gamma p \rightarrow K^+ \Lambda(1520)$ 

## (1) $\gamma p \rightarrow K^+ \Lambda(1520)$ : Theoretical background and motivation

## 2 Experiment and analysis procedure

## (3) $\gamma p \rightarrow K^+ \Lambda(1520)$ : Decay distributions and spin density matrix elements

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- $\Lambda(1520)$  is  $\frac{3}{2}^{-}$  baryon
- Decay modes
  - $N\bar{K}$  ( $pK^-$ ,  $n\bar{K}^0$ ): 45%
  - $\Sigma \pi (\Sigma^+ \pi^-, \Sigma^0 \pi^0, \Sigma^- \pi^+)$ : 42%
  - Λππ: 10%

 $\bullet\,$  Narrow resonance ( $\Gamma=15\,{\rm MeV})$  compared to other excited baryons

## $\gamma p \rightarrow K^+ \Lambda(1520)$ : Polarization observables

- Polarization of  $\Lambda(1520)$  expressed by spin density matrix, measured by angular distribution of decay products
- Polarization reveals information about production mechanism
- Use Gottfried-Jackson (t-channel helicity) frame



## $\gamma p \rightarrow K^+ \Lambda(1520)$ : Polarization observables

• For decay  $\frac{3}{2}^- \rightarrow \frac{1}{2}^+ 0^-$  with unpolarized target, unpolarized beam, parity-conserving production and decay: seven independent observables:

$$\begin{pmatrix} \frac{1}{2} - \rho_{\frac{1}{2}\frac{1}{2}} & \operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) + i\operatorname{Im}(\rho_{\frac{3}{2}\frac{1}{2}}) & \operatorname{Re}(\rho_{\frac{3}{2}-\frac{1}{2}}) + i\operatorname{Im}(\rho_{\frac{3}{2}-\frac{1}{2}}) & i\operatorname{Im}(\rho_{\frac{3}{2}-\frac{3}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}} & i\operatorname{Im}(\rho_{\frac{1}{2}-\frac{1}{2}}) & \operatorname{Re}(\rho_{\frac{3}{2}-\frac{1}{2}}) - i\operatorname{Im}(\rho_{\frac{3}{2}-\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}} & \rho_{\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) + i\operatorname{Im}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) + i\operatorname{Im}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) + i\operatorname{Im}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}} & \rho_{\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}\frac{1}{2}} & \rho_{\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}\frac{1}{2}} & \rho_{\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) \\ \rho_{\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}} & -\operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}\frac{$$

• Only three of these observables are measureable in decay distribution:

$$W(\theta, \phi) = \frac{3}{4\pi} \left\{ \left( \frac{1}{3} + \cos^2 \theta \right) \rho_{\frac{1}{2}\frac{1}{2}} + \sin^2 \theta \left( \frac{1}{2} - \rho_{\frac{1}{2}\frac{1}{2}} \right) - \frac{1}{\sqrt{3}} \sin 2\theta \cos \phi \operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) - \frac{1}{\sqrt{3}} \sin^2 \theta \cos 2\phi \operatorname{Re}(\rho_{\frac{3}{2}-\frac{1}{2}}) \right\}$$

Jackson, High Energy Physics, Les Houches 1965

- Can generalize to case of polarized photon beam
  - Unpolarized beam: can measure 3 independent observables
  - Linearly polarized beam: 6 additional observables
  - Circularly polarized beam: 2 additional observables

## Possible production mechanisms

- t-channel
  - K exchange
    - Pure scalar meson exchange implies  $\rho_{\frac{1}{2}\frac{1}{2}} = \frac{1}{2}$ , all other  $\rho = 0$  (in GJ frame)
  - K<sup>\*</sup> exchange
- contact term
  - Needed to preserve gauge invariance
  - Absent for photoproduction off neutron
  - Contact term dominance could explain suppressed cross-section off neutron LEPS, PRL 103, 012001 (2009)
- s-channel
  - Prediction of  $N^* \to K\Lambda(1520)$  decays from  $N^*(2120)\frac{3}{2}^-$  (formerly called  $N^*(2080)$ ), missing  $\frac{1}{2}^-$  and  $\frac{5}{2}^-$  states Capstick and Roberts, PRD 58, 074011 (1998)
- u-channel









- Nam and Kao predict decay angular distributions as function of production angle and energy PRC 81, 055206 (2010)
- Model includes:
  - Reggeized t-channel (K and K\*) exchange
  - contact term
  - s-channel (ground-state N and  $N^*(2120) \frac{3}{2}^-$ ) exchange
  - u-channel (ground-state Λ) exchange



## Previous measurements of decay distributions

#### Barber et al (LAMP2), Z. Physik C 7, 17-20 (1980)



- LAMP2 (Daresbury):  $E_{\gamma} = 2.8 - 4.8 \text{ GeV}$
- LEPS:  $E_{\gamma} = 1.75 2.4$  GeV, 2 angular bins
- SAPHIR: 4 bins from  $E_{\gamma} = 1.69 2.65 \text{ GeV}$
- All previous results averaged over wide energy bins, coarse (or no) binning in production angle

## $\Lambda(1520)$ cross-section bump

- Bump in  $\Lambda(1520)$  differential cross-section at  $\sqrt{s} = 2.1 \,\mathrm{GeV}$
- Origin unknown
  - Resonance?
  - Other?



LEPS, PRL 104, 172001



CLAS, PRC 88, 045201

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- CEBAF: 6 GeV (now 12 GeV)
   e<sup>-</sup> accelerator at Jefferson Lab (Newport News, Virginia)
- CLAS (CEBAF Large Acceptance Spectrometer) electroproduction and photoproduction experiments
- Tagged bremsstrahlung photon beam



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## Analysis overview

- g11a dataset
  - photon beam on liquid  $H_2$  target
  - unpolarized beam, unpolarized target
  - $E_e = 4.019$  GeV electron beam energy
  - 20 billion triggers
- $\Lambda(1520) \rightarrow pK^-$  decay mode ( $pK^+K^-$  final state)
  - 3-track:  $pK^+K^-$
  - 2-track: pK<sup>+</sup>(K<sup>-</sup>): 10x more statistics, wider acceptance, background difficulties (results not presented today)
- Bin in 60 MeV wide  $\sqrt{s}$  bins
  - 13 bins from  $\sqrt{s} = 2.04 2.82 \,\mathrm{GeV}$
- Standard fiducial, PID cuts
- Kinematic fit with 5% confidence level cut
  - 4C fit:  $\gamma p \rightarrow K^+ K^- p$
- Cut out  $\mathit{IM}(\mathit{K}^+\mathit{K}^-) < 1.040~{\rm GeV}$  to remove  $\phi$

#### After all cuts:



 $\Lambda(1520)$  peak on top of non- $\Lambda(1520)$   $pK^+K^-$  events

- Can we separate  $\Lambda(1520)$  events?
- Not possible if processes interfere

### Q-value method:

- For each event, calculate probability that given event is signal
  - Find N=100 nearest neighbor events in phase space
  - Fit mass distribution of nearest neighbor events to signal (Breit-Wigner) + background (polynomial) function
  - Williams et al JINST 4 P10003 (2009)
- Assumes non-interfering background
- Use signal probability as weight in event-based maximum likelihood fit

## Background subtraction



For SDME extraction, only consider events in the center of the peak (1500-1540 MeV)

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## Decay distributions



- Look at (acceptance-corrected) decay distributions in Gottfried-Jackson frame  $(\cos \theta_{GJ})$  weighted by  $\Lambda(1520)$ probability.
  - Bin in production angle  $\theta_{K^+,CM}$
  - Correct for acceptance
  - Not how we will extract spin density matrix elements!
  - Just a check
- After integrating over  $\phi$ , decay distribution should have form  $\alpha + \beta \cos^2 \theta_{GJ}$ . Distribution is even in  $\cos \theta$ !

= 900

Can compare with models of Nam and Kao at two energies/angles.



$$W(\theta, \phi) = \frac{3}{4\pi} \left\{ \left(\frac{1}{3} + \cos^2 \theta\right) \rho_{\frac{1}{2}\frac{1}{2}} + \sin^2 \theta \left(\frac{1}{2} - \rho_{\frac{1}{2}\frac{1}{2}}\right) - \frac{1}{\sqrt{3}} \sin 2\theta \cos \phi \operatorname{Re}(\rho_{\frac{3}{2}\frac{1}{2}}) - \frac{1}{\sqrt{3}} \sin^2 \theta \cos 2\phi \operatorname{Re}(\rho_{\frac{3}{2}-\frac{1}{2}}) \right\} \right\}$$

What we measure is not the true decay distribution,  $W(\rho, \vec{x})$ , but decay distibution times acceptance:  $W(\rho, \vec{x})\eta(\vec{x})$  ( $\rho$  is the spin density matrix,  $\vec{x}$  is the kinematics of the reaction,  $\eta$  is acceptance).

Construct a PDF for probability of detecting event with kinematics  $\vec{x}$  given SDM  $\rho$ :

$$\mathcal{P}(
ho, ec{x}) = rac{W(
ho, ec{x})\eta(ec{x})}{\int W(
ho, ec{x}')\eta(ec{x}')\,dec{x}'}$$

Denominator is easy to calculate using Monte Carlo method:

$$\mathcal{N}(\rho) = \int \mathcal{W}(\rho, \vec{x}') \eta(\vec{x}') \, d\vec{x}' = C \sum_{i \in accepted} \mathcal{W}(\rho, \vec{x}_i)$$

Construct likelihood

$$L \propto \prod_{i \in data} rac{W(
ho, ec{x}_i)}{N(
ho)}$$

Maximize L to find best values of  $\rho$ 

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## Spin density matrix elements



• 
$$\rho_{\frac{1}{2}\frac{1}{2}} = .25 \implies \rho_{\frac{3}{2}\frac{3}{2}} = .25$$
:  $S_z = \pm \frac{3}{2}$ ,  $S_z = \pm \frac{1}{2}$  equally populated

•  $\rho_{\frac{3}{2}\frac{1}{2}}$  consistent with zero •  $\rho_{\frac{3}{2}-\frac{1}{2}}$  non-zero (non-flat  $\phi_{GJ}$  distribution)

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## Spin density matrix elements



Gottfried-Jackson frame, statistical errors only

- No strong energy dependence
- $\rho_{\frac{1}{2}\frac{1}{2}} = .20$ -.40 in most regions
- $\rho_{\frac{3}{2}\frac{1}{2}}$  consistent with zero
- $\rho_{\frac{3}{2}-\frac{1}{2}}$  non-zero (non-flat  $\phi_{GJ}$  distribution)

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## Spin density matrix elements vs. energy



Gottfried-Jackson frame, statistical errors only

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- Polarization of  $\Lambda(1520)$  expressed in spin density matrix formalism
- Spin density matrix elements extracted from CLAS photoproduction data
  - Much finer binning in energy, production angle than previous measurements
  - $ho_{rac{1}{2}rac{1}{2}}$  measurement shows neither  $S_z=\pmrac{3}{2}$  or  $S_z=\pmrac{1}{2}$  dominates
  - Non-flat  $\phi_{GJ}$  distribution measured for first time (non-zero  $\rho_{\frac{3}{2}-\frac{1}{2}}$ )
- $\bullet\,$  More statistics, wider angular coverage coming soon with missing  $K^-$  analysis

# Backup slides

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## $\phi_{GJ}$ Decay Distributions

What about  $\phi_{GJ}$ ? Irregular acceptance makes difficult to compare.



Some deviation from flat distribution

Need to input functional form of  $MM(K^+)$  to do background fit:



Breit-Wigner w/ mass-dependent width, convoluted with Gaussian. Quadratic background.

## Q-value method

Choose kinematic variable, M, whose distribution can be described by a sum of background and signal functions:

$$F(M,\vec{\alpha}) = S(M,\vec{\alpha}) + B(M,\vec{\alpha})$$

where  $\vec{\alpha}$  is a set of unknown parameters, *S* is the signal distribution (e.g. Breit-Wigner), *B* is the background distribution (e.g. polynomial) For each event *i*, find *N* nearest neighbors, with distance to event *j* as:

$$d_{ij} = \sum_{k} \left[ rac{ heta_k^i - heta_k^j}{R_k} 
ight]^2$$

where  $\vec{\theta}$  are kinematic variables other than M (e.g.  $\cos \theta_{\text{production}}$ ,  $\cos \theta_{\text{decay}}$ ),  $R_k$  is range of  $\theta_k$ Fit M distribution of nearest neighbors to F to determine  $\vec{\alpha}_i$ . Calculate signal probability:

$$Q_i = \frac{S(M_i, \vec{\alpha_i})}{S(M_i, \vec{\alpha_i}) + B(M_i, \vec{\alpha_i})}$$

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## two-track vs three-track



## $K^+\Lambda(1520)$ - $\phi p$ coupled-channel effects

- $K^+\Lambda(1520)$  intermediate state studied in  $\gamma p \rightarrow \phi p$ 
  - Proposed to explain bump in  $\gamma p \rightarrow \phi p$  cross-section near  $K\Lambda(1520)$  threshold
  - Ozaki et al, PRC 80, 035201 (2009)
  - Ryu et al, arxiv:1212.6075
- Understanding  $K^+\Lambda(1520)$ production mechanism may help understanding of  $\phi$ photoproduction





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## $\Lambda(1520)/\phi$ overlap



- $\phi p$  and  $K^+ \Lambda(1520)$  can decay to the same final state  $(K^+ K^- p)$ , overlap in phase space
- Interfering background
- No acceptance in overlap region for 3-track topology
- Hard cut on  $K^+K^-$  mass to cut out  $\phi$
- No overlap at higher energies



 $\sqrt{s} = 2340 - 2400$  GeV Black is data (not acceptance corrected). Blue is accepted Monte Carlo weighted by full fit to decay angular distribution.