Spin Observables and Spin Density Matrix Elements for $\gamma p \rightarrow \rho^0 p$

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QNP, Valparaiso, Mar 15

ToC

The Motivation

Experimental Details

Analysis

QCD and its shortcomings. Need different models for the resonant region.

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- Spin of final state particles affected by resonances.
- Polarization of initial spin states increase sensitivity.

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- Spin of final state particles affected by resonances.
- Polarization of initial spin states increase sensitivity.
- \blacktriangleright $\pi^+\pi^-$ channel why? stronger coupling?
- Preceeding analyses dominated by single pion channels and 1.7GeV two pion channel becomes dominant.

Theory

$$\mathscr{F} \equiv \langle \mathbf{q} \lambda_V \lambda_2 \mid T \mid \mathbf{k} \lambda \lambda_1 \rangle \tag{1}$$

$$\langle \boldsymbol{q}\lambda_{V}\lambda_{2}|T|\boldsymbol{k}\lambda\lambda_{1}\rangle \rightarrow H_{a\lambda_{V}}(\theta)$$
 (2)

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The vector meson amplitude can be expressed in helicity space by the following matrix:

$$\mathscr{F} = \begin{bmatrix} H_{21} & H_{11} & H_{3-1} & -H_{4-1} \\ H_{41} & H_{31} & -H_{1-1} & H_{2-1} \\ H_{20} & H_{10} & -H_{30} & H_{40} \\ H_{40} & H_{30} & H_{10} & -H_{20} \\ H_{2-1} & H_{1-1} & H_{31} & -H_{41} \\ H_{4-1} & H_{3-1} & -H_{11} & H_{21} \end{bmatrix}$$

$$(3)$$

Generally spin observables, Ω can be expressed as:

$$\Omega = \frac{Tr[\mathscr{F}(A_{\gamma}A_{N})\mathscr{F}^{\dagger}(B_{V}B_{N'}]}{Tr[\mathscr{F}\mathscr{F}^{\dagger}]}$$
(4)

Where the trace is over the helicity quantum numbers.

$$\Sigma = \frac{Tr[\mathscr{F}\sigma_X^{y}\mathscr{F}^{\dagger}]}{Tr[\mathscr{F}\mathscr{F}^{\dagger}]}$$
 (5)

Usually seen expressed as an asymmetry.

$$\Sigma = \frac{\sigma^{\parallel} - \sigma^{\perp}}{\sigma^{\parallel} + \sigma^{\perp}} \tag{6}$$

So much for the beam asymmetry.

SDMEs

The helicity state is dependent on the spin-helicity relationship and we need the SDM for the decaying vector meson which is associated with the SDM of the photon:

$$\rho(V) = T\rho(\gamma)T^{\dagger} \tag{7}$$

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$$\rho(\gamma) = \frac{1}{2}I + P_{\gamma} \cdot \sigma$$
(8)

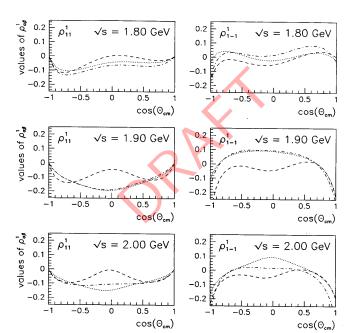
Using the above relations and the helicity-amplitude formalism we can show the dependence of the polarisation vector ${m P}_{\gamma}$ of the density matrix $\rho(V)$:

$$\rho(V) = \rho^{o} + \sum_{i=1}^{3} P_{\gamma}^{\alpha} \rho^{\alpha}$$
 (9)

Where P^{α}_{γ} are the components of the polarisation vector, \mathbf{P}_{γ} , and the ρ^{α} are hermitian matrices.



SDMEs Resonance Dependency



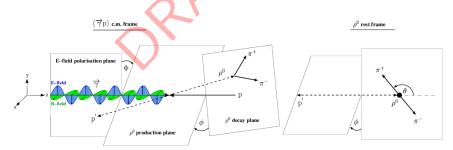
Angular Decay Distribution

The density matrix is related to the decay angular distribution:

$$W(\cos\theta,\phi) = M\rho(V)M^{\dagger} \tag{10}$$

M is the decay amplitude and θ and ϕ are the polar and azimuthal angles of the detected π^+ in the helicity frame. And so we get the following for ρ^0 -meson decay distribution:

$$W(\cos\theta,\phi,\rho) = W^0(\cos\theta,\phi,\rho) - P_{\gamma}\cos 2\Phi W^1(\cos\theta,\phi,\rho) - P_{\gamma}\sin 2\Phi W^2(\cos\theta,\phi,\rho)$$
(11)



$$\begin{split} W^0(\cos\theta,\phi,\rho) &= \frac{3}{4\pi} (\frac{1}{2} (1-\rho^0_{00}) + \frac{1}{2} (3\rho^0_{00} - 1)\cos^2\theta - \sqrt{2}Re\rho^0_{10}\sin2\theta\cos\phi \\ &- \rho^0_{1-1}\sin^2\theta\cos2\phi \\ W^1(\cos\theta,\phi,\rho) &= \frac{3}{4\pi} (\rho^1_{11}\sin^2\phi + \rho^1_{00}\cos^2\theta - \sqrt{2}\rho^1_{10}\sin2\theta\cos\phi - \rho^1_{1-1}\sin^2\theta\cos2\phi) \\ W^2(\cos\theta,\phi,\rho) &= \frac{3}{4\pi} (\sqrt{2}Im\rho^2_{10}\sin2\theta\sin\phi + Im\rho^2_{1-1}\sin^2\theta\sin2\phi) \end{split}$$

To simplify our task we integrate over two of the angles to leave us with a single angle function. We are left with the five following equations:

$$W(\cos\theta) = \frac{3}{4}[1 - \rho_{00}^0 + (\rho_{00}^0 - 1)\cos^2\theta]$$
 (12)

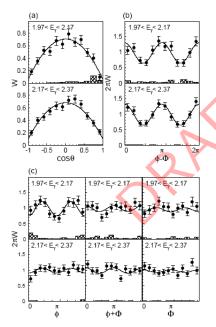
$$W(\phi) = \frac{1}{2\pi} [1 - 2Re\rho_{1-1}^{0} \cos 2\phi]$$
 (13)

$$W(\phi - \Phi) = \frac{1}{2\pi} [1 + P_{\gamma}(\rho_{1-1}^{1} - Im\rho_{1-1}^{2})\cos 2(\phi - \Phi)]$$
 (14)

$$W(\phi + \Phi) = \frac{1}{2\pi} [1 + P_{\gamma}(\rho_{1-1}^{1} + Im\rho_{1-1}^{2})\cos 2(\phi + \Phi)]$$
 (15)

$$W(\Phi) = \frac{1}{2\pi} [1 - P_{\gamma}(2\rho_{11}^{1} + \rho_{00}^{1})\cos 2\Phi]$$
 (16)

And what these distributions look like.



Distributions for photoproduced ϕ . T. Mibe et al, 2005 from Salamanca-Bernal thesis.

Experimental Details

- ► Linearly polarized beam using the bremsstrahlung tagging facility in Hall B: 70 -80%.
- g8b data. Photon energy runs at 1300,1500,1700,1900,2100MeV.
- Unpolarized 40cm hydrogen target: simplifies analysis.
- Exclusive detection of all final state particles.

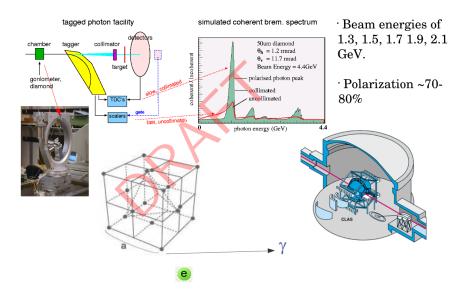
Jefferson Lab and CEBAF



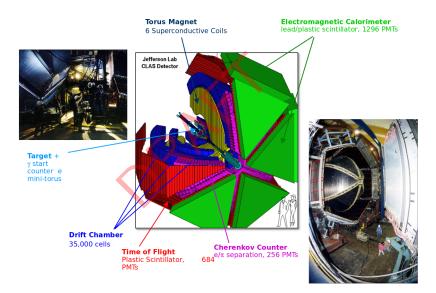
Continuous Electron Beam Accelerator Facility

- \rightarrow E: 0.75 –6 GeV
- $\rightarrow I_{\text{max}}: 200 \text{ ÅA}$
- \rightarrow RF: 1499 MHz
- \rightarrow Duty Cycle: 100%
- \rightarrow (E)/E: 2.5x10⁻⁵
- $\rightarrow \ Polarization; 80\%$
- → Simultaneous distribution to 3 experimental Halls

Photon Tagging Facility



CEBAF Large Acceptance Detector



Particle Identification Stage

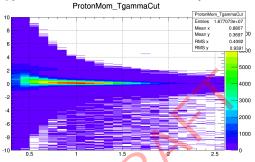
- ► First stage is to identify charged particles and remove all events with less than two charged particles.
- ▶ This means we can have more than one topology: e.g. π^+p
- Fudicial cuts around the regions near the torus coils.
- Momentum and energy loss corrections.
- Then apply cuts such as missing mass cuts.

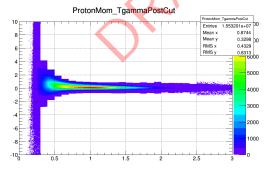
Z Vertex and Production Angle Cuts



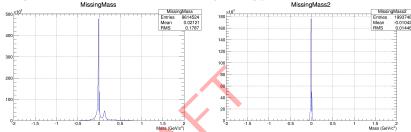


Timing Cuts: Momentum Dependent ProtonMom_TgammaCut

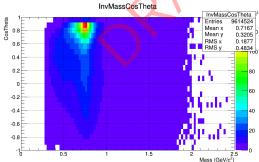




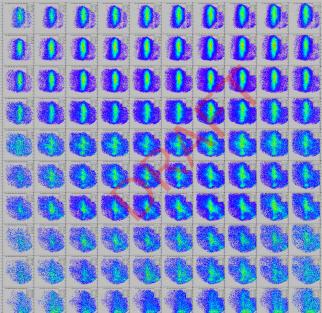
Missing mass for exclusive topology



Then we split the events into kinematic bins: W and $\cos \theta$.

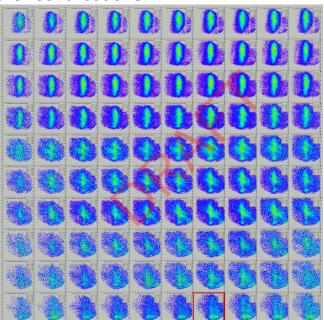


Separating Other Contributions



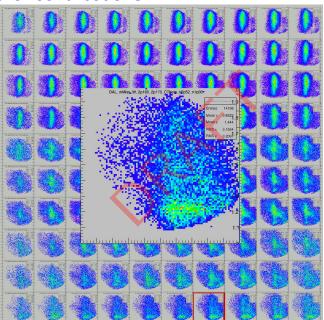
- Can try and separate other contributions to the final state.
- E.g. Δ⁺⁺ is produced with a negative pion.
- Reconstruct it the same as with the two pions for the ρ^0 .
- Separate using Dalitz plots. Not necessarily viable for every kinematic bin.

Other contributions



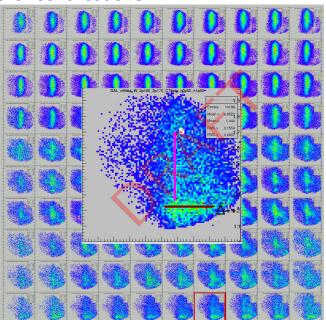
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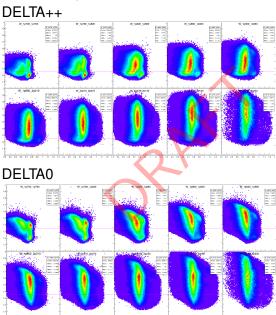
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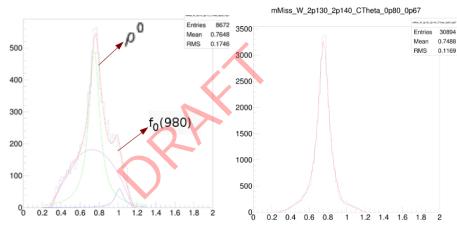


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Delta dalitz plots for different W ranges



ρ^0 Signal Extraction



Currently using a binned fit but it isn't robust enough over all kinematic bins.

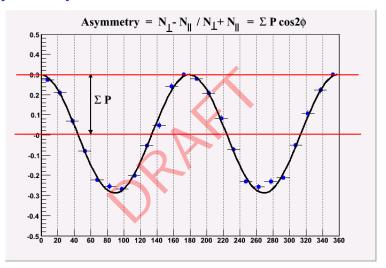
ρ^0 Signal Extraction

Decreasing $cos\theta$ (1 to -1)



W (2.07 to 2.23GeV).

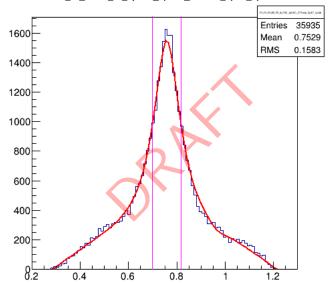
Asymmetry Extraction



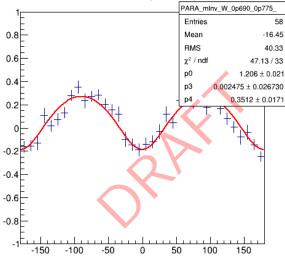
- Systematics of detector acceptance cancel out.
- Only need to know P_{lin}.



PI_PI_INVM_W_2p150_2p160_CTheta_0p67_0p48



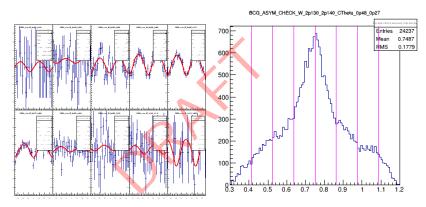
PARA_mInv_W_0p690_0p775_



- Yield not equal for PARA and PERP.
- Photon polarization not equal.
- Small offset left as free parameter.

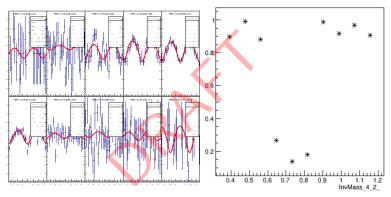
$$\frac{\sigma_{\perp} - \sigma_{||}}{\sigma_{\perp} + \sigma_{||}} = \frac{\left(\frac{N_{\perp}}{N_{||}} - 1\right) - \left(\frac{N_{\perp}}{N_{||}} P_{\perp} + P_{||}\right) \sum \cos\left(2(\phi - \phi_{0})\right)}{\left(\frac{N_{\perp}}{N_{||}} + 1\right) - \left(\frac{N_{\perp}}{N_{||}} P_{\perp} - P_{||}\right) \sum \cos\left(2(\phi - \phi_{0})\right)}$$

Dilution Factor



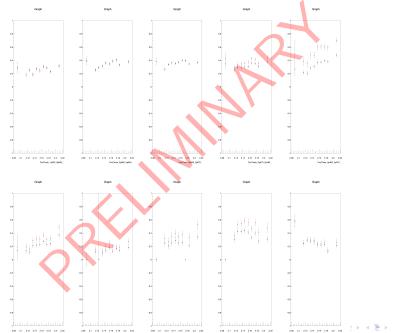
Checking what is considered background doesn't have a contribution to the asymmetry.

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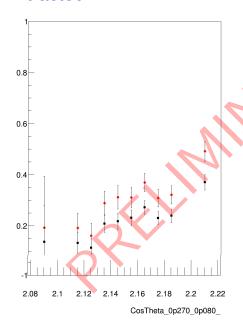


Checking what is considered background doesn't have a contribution to the asymmetry.

Extracted Σ



Extracted Σ



- Red markers are with dilution factor.
- Black markers are without the dilution factor.

Summary

- Managed to extract beam asymmetry for one beam energy setting with a non-trivail form.
- ▶ Fitting for extracting the ρ^0 signal still needs work.
- Hope to use simulated data to describe the shape of the bcg contribution to aid fit (phase space and deltas projection).
- Need to finish acceptance correction in order to extract the SDMEs.
- Also working on the simulation for that purpose.