

Studies of Baryon Structure and Baryon Interactions at Thomas Jefferson National Accelerator Facility

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Thomas Jefferson National Accelerator Facility

Premier U.S. Facility to study Quantum Chromodynamics: quantitative understanding of the internal structure of nucleons and nuclei



JLab in Newport News, VA

- Polarized c.w. electron
 beam: P_e ~ 86%
- Beam energies up to
 E₀ = 6 GeV (recently
 upgraded to 12 GeV)
- Four experimental Halls
 A, B, C, and D

The CEBAF Large Acceptance Spectrometer CLAS



Performance

- $L = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- $\int B dI = 2.5 T m$
- $\Delta p/p \sim 0.5-1 \%$
- ~ 4π acceptance
- Best suited for charged multiparticle final states
- Bremsstrahlung Photon Tagger ($\Delta E_{\gamma}/\,E_{\gamma}\,{\sim}10^{-3})$







Outline

Transition from Hadronic to Partonic Degrees of Freedom Photodisintegration of Few-Nucleon Systems

Hyperon-Nucleon Interaction Final-State Interactions in Strangeness Photoproduction

Future Perspectives: Studies at 12 GeV Near-Threshold J/ψ Photoproduction off Deuteron

Transition from Hadronic to Partonic Degrees of Freedom

Photodisintegration of Few-Nucleon Systems at Medium Energies

Large Momentum Transfer Exclusive Processes (Hard Scattering)

Short-range dynamics (quark content of nuclei)









Studying the Strong Interaction through Hard Scattering

- Identify large-momentum-transfer (hard) subprocess
- Apply factorization (separation of scales):
 - -Apply pQCD to hard subprocess
 - –Non-perturbative contributions expressed through measurable/calculable quantities

Why Hard Photodisintegration?

- Real photon: $Q^2=0$. Factorization not proved.
- Hard scale is given by the overall momentum transfer in the reaction (size of interaction volume): search for onset of partonic dynamics in nuclei
 - experimental tool: dimensional scaling of invariant cross sections

Dimensional Scaling Laws in Nuclear Physics

Brodsky, Farrar (1973): from dimensional analysis and perturbative QCD

• At high t and high s, power-law behavior of the invariant cross section of an exclusive process $A + B \rightarrow C + D$ at fixed CM angle:

$$\frac{d\sigma}{dt} = \frac{1}{s^{n-2}}f(t/s)$$

where n is the total number of the initial and final elementary fields.

• The energy dependence of the scattering amplitude given by the 'hard-scattering amplitude' T_H for scattering collinear constituents from the initial to the final state

$$pp \rightarrow pp \equiv 3q 3q \rightarrow 3q 3q$$

$$p \rightarrow f_{H} \qquad p \rightarrow f_{H} \qquad f_$$

S.J. Brodsky and G.R. Farrar, Phys. Rev. Lett 31, 1153 (1973); S.J. Brodsky and J.R. Miller, Phys. Rev. C 28, 475 (1983)

Previous Experimental Findings in Two-Nucleon Systems



• Scaling at p_T > 1.1 GeV/c

• Scaling at p_T > 1.4 GeV/c

he 2 — $2\sqrt{s'}$ $\sum_{\substack{n=1\\ n=1}}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{i=1}^$ rk he $h_{otodisintegration}^{\alpha_c x_1 + \frac{1}{s'}}$ in the Scaling Regime P on le- $\stackrel{ ext{on}}{ ext{he}} = \left[-igT_c^F \gamma^{
u} \right]$ **XGSM)** $\psi_N(x_1, p_{1\perp}, k_{1\perp}) \psi_N^{\dagger}(x'_1, p_{F\perp}, k_{1\perp})$ (HRM) Three quark exchange with arbitrary – Photor th(QGSM) ge)_ Photon absorbed by a quark in one ere th Nonlinear Regge trajectories $\psi_N(x_2, p_{2\perp}, x_2) = \frac{1}{x_2} \frac{\psi_N(x_2, p_{2\perp}, x_2)}{x_2}$ nucleon, followed by a hard scattering off a quark in the second nucleon Requires phenomenological input on pn $\gamma \times \frac{dx_1}{\sqrt{2(2\pi)^3}} \frac{d^2k_{1\perp}}{1 - x_2} \frac{dx_2}{2(2\pi)^3} \frac{d^2k_{2\perp}}{\frac{\varphi}{2(2\pi)^3}} \frac{\Psi_d(\alpha_c, p_\perp)}{\alpha_c(1 - \varphi_c)} \frac{d^2p}{2(2\pi)^3} \frac{d^2p}{2(2\pi)^3} \frac{d^2p}{\alpha_c(1 - \varphi_c)} \frac{d^2p}{2(2\pi)^3} \frac{d^2p}{\alpha_c(1 - \varphi_c)} \frac{d^2p}{\alpha_c(1 - \varphi_c)}$ elastic scattering where $\epsilon^{\pm} \equiv \frac{1}{2} (\epsilon_x \not\equiv i\epsilon_y)$ The dominant contrib ing from the soft component of the deuteron of p_A **p**₁ $k_1 + q$ $\alpha_c \sim \frac{1}{2}$. Thus we may substitute $\alpha_c = \frac{1}{2}$. Afte p_d out electromagnetic term, one can identify the - k₂ integral in Eq.(3) (up to a scaling factor $f(l^2/s)$ \mathbf{p}_2 **p**_B ticular contribution to the quark interchange V. Y. Grishina, L. Kondratyuk, W. Cassing, E. De Sanctis for the wide angle nucleon-nucleon scattering M. Mirazita, F. Ronchetti, and P. Rossi ed L.L. Frankfurt, G.A. Miller, M.M. Sargsian, and M.I. Strikman, Phys. Rev. Lett. 84, 3045 (2000) em Eur Alor 1. (\$9,127) (2054) en summing over the struck chtributions from photon scattering off neutron and proton ^{of}PreviousⁿCross-Section and Polarization-Transfer cannot select between the models $\underline{i(a_u New)}($ Spine Observable is expected to make an impact.

$d\sigma = \sigma_0 \{P_{lin}, Cos 2\phi + \alpha \cos \theta_x (-P_{lin}, Cos 2\phi + \alpha \cos \theta_x (-P_{lin}, Cos 2\phi + \alpha \cos \theta_x (-P_{lin}, Cos 2\phi) - \alpha \cos \theta_x (-P_{lin}$



Figure from Nick Zachariou

Linearly Polarized Photons: $\vec{\gamma}$ $\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega_0} \left[1 - P_{lin} \sum \cos 2\varphi \right]$

Linearly Polarized Photons (E06-103)

- CLAS Detector
- E_e = 3.3 5.2 GeV
- coherent edge at: 1.3, 1.5, 1.7, 1.9,
 2.1, 2.3 GeV
- **-** Pγ = 70% 90%
- ~30×10⁹ triggers

Dynamics of Hard Deuteron Photodisintegration in the Scaling Regime

Beam Spin Asymmetry: New Results



N. Zachariou, Y. Ilieva, B.L. Berman, N.Ya. Ivanov, M.M. Sargsian, R. Avakian, G. Feldman, P. Nadel-Turonski et al., Phys. Rev. C **91**, 055202 (2015)

- Σ at θ_{cm}= 90°
 - σ_{stat} = 6% 15%
 - σ_{syst} < 6%
- The beam spin asymmetry is positive at $E_{\gamma} = 0.75 2.3$ GeV.
- Data suggest a continuous increase up to E_γ of 2 GeV.
- Current QGSM calculations do not reproduce the shape of the energy distribution.
- HRM reproduces the shape, underestimates the magnitude.

Dynamics of Deuteron Hard Photodisintegration in the Scaling Regime



•Rich structures at low energies evolving to a single peak at 90° at higher E_{γ} .

•Kinematics at the limit of applicability of HRM.

Complex dynamics QGSM needs more
 work.

N. Zachariou, Y. Ilieva, B.L. Berman, N.Ya. Ivanov, M.M. Sargsian, R. Avakian, G. Feldman, P. Nadel-Turonski et al., Phys. Rev. C **91**, 055202 (2015)



Scaling in A > 2 expected at energies much higher than 1 GeV, however...

Two-Body Photodisintegration of ³He

Scaling in A > 2 expected at energies much higher than 1 GeV, however...



I. Pomerantz, Y. Ilieva, R. Gilman, D. Higginbotom, E. Piazetski, S.Strauch et al., Phys. Rev. Lett. 110, 242301 (2013).

Two-Body Photodisintegration of ³He

Scaling in A > 2 expected at energies much higher than 1 GeV, however...



 Extracted value from fits to JLab data:

N = 17 ± 1

 |t|_{thr} and p_{⊥thr} are too low to support hard scattering hypothesis:

$$|t|_{thr} = 0.64 (GeV/c)^2$$

 $p_{\perp thr} = 0.95 \text{ GeV/c}$

- What is the phenomenology of the observed scaling?
- I. Pomerantz, Y. Ilieva, R. Gilman, D. Higginbotom, E. Piazetski, S.Strauch et al., Phys. Rev. Lett. 110, 242301 (2013).

Dynamics of ³He Hard Photodisintegration in the Scaling Regime

Why does it scale at such low energies?



Photodisintegration of Few-Nucleon Systems: Summary

Hard Deuteron Photodisintegration

- New data on the beam-spin asymmetry is in general consistent with the hardrescattering mechanism.
- Data can be used to constrain the helicity amplitudes of the underlying pn elastic scattering (at kinematics where no data exist).

Hard ³He Photodisintegration

- Unexpected dimensional scaling observed at E_{γ} > 0.7 GeV.
- The origin of dimensional scaling is a puzzle (regime of conformal window or hard rescattering?).

The phenomenology of QCD transition between scales will be further explored at 12 GeV.

Studies of Hyperon-Nucleon Interaction

Why Study?

YN interaction not well known

- not all free parameters of the YN potential can be obtained from the NN potential via flavor SU(3) symmetry
- example: large uncertainties of YN scattering lengths: $a({}^{1}S_{0}) = -0.7 - -2.6 \text{ fm},$ $a({}^{3}S_{1}) = -1.7 - 2.15 \text{ fm}$
- YN elastic scattering database poor
- alternative approaches:
 - hypernuclear spectroscopy
 - studies of FSI in production reactions:

 $\gamma d \to K^+ \Lambda n \qquad pp \to K^+ \Lambda p$ $\gamma d \to K^0 \Lambda p$

YN Interaction and Exclusive Hyperon Photoproduction off Deuteron



New approach

- Hyperon Beam produced in first step
- Hyperons scatter off neutrons in a second step

Theoretical Studies

- Observables sensitive to YN potentials at certain kinematics at a level of ~10% (K. Miyagawa et al., Phys. Rev. C 74, 034002 (2006); A. Salam et al., Phys. Rev. C 74, 044004
 (2006); H. Yamamura et al., Phys. Rev. C 61, 014001 (1999))
- Spin-averaged scattering length can be extracted from data close to threshold (A. Gasparian et al., Phys. Rev. C 69, 034006 (2004))

Exclusive Hyperon Photoproduction off Deuteron Background Mechanisms





Exclusive Hyperon Photoproduction off Deuteron Polarization Observables

 $\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d\Omega_0} \Big[1 - P_{lin} \sum \cos 2\varphi - \alpha \cos \theta_x (P_{lin} O_x \sin 2\phi + P_{circ} C_x) - \alpha \cos \theta_z (P_{lin} O_z \sin 2\phi + P_{circ} C_z) + \dots \Big]$



 Λ self-analysing power: $\alpha = 0.642 \pm 0.013$

Relevant Kinematic Variables:

- From $\Lambda n \rightarrow \Lambda n$: $\Theta_{\Lambda'}$, $W_{\Lambda n} = IM_{\Lambda n}$
- From $\gamma p \rightarrow K^{+} \Lambda$: E_{γ} , p_{K} , θ_{K}

 $\Sigma(E_{\gamma}, p_{K}, \theta_{K}, \theta_{\Lambda}', W_{\Lambda n})$ $C_{x}(E_{\gamma}, p_{K}, \theta_{K}, \theta_{\Lambda}', W_{\Lambda n})$ $C_{z}(E_{\gamma}, p_{K}, \theta_{K}, \theta_{\Lambda}', W_{\Lambda n})$

Experimental Facility: CLAS at Jefferson Lab Experiment E06-103 (g13)

Circularly Polarized Photons (g13a)

- E_e = 2 GeV; 2.65 GeV
- electron polarization: ~ 80%
- triggers: ~20×10⁹ triggers

Linearly Polarized Photons (g13b)

- E_e = 3.3 5.2 GeV
- coherent edge at: 1.3, 1.5, 1.7, 1.9,
 2.1, 2.3 GeV
- $P\gamma = 70\% 90\%$
- ~30×10⁹ triggers



K⁺ and ∧: measured n: reconstructed from 4-momentum vector conservation

P. Nadel-Turonski, B. Berman, Y. Ilieva, D. Ireland, A. Tkabladze et al., E06-103: "Kaon Production on the Deuteron Using Polarized Photons"

Selection of Quasi-Free and FSI Mechanisms

Event Distribution over Spectator Momentum Comparison with Model Distribution P_n ($\gamma d \rightarrow K^+ \Lambda n$)



The removal of events with $P_n < 0.2$ GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure. Paris Potential describes well low P_n data. High-momentum tail drops off at ~0.6 GeV/c: effect on data interpretation.

Preliminary Results **Observables for Circularly-Polarized Photons** $\frac{d\sigma^{\pm}}{d\Omega} = \frac{d\sigma}{d\Omega_{unp}} \Big[1 \pm P_{circ}C_x\alpha\cos\theta_x \pm P_{circ}C_z\alpha\cos\theta_z + \alpha P_y\cos\theta_y \dots \Big]$

Work by Tongtong Cao









Model Comparison: Perspectives



- one-, ..., four-fold
 differential observables
 will be extracted
- experimental observables are integrated over the CLAS acceptance
- direct comparison with models is not trivial
 - collaboration with theorists
 - use the model as event generator and process through CLAS simulation
 - provide a sample of FSI events that can be binned in any way

Linearly Polarized Photons: $\frac{d\sigma}{d\Omega} = \frac{dP}{d\Omega_0} \left[1 + \frac{P_{lin}}{2} + \frac$



Figure from Nick Zachariou

- Single-polarization observable, i.e. smallest statistical uncertainties
- For the K⁺Λn final state, φ is not uniquely defined: φ_K, φ_Λ, φ_n, φ_{KΛ}, φ_{Λn}, φ_{Kn}
- Observable can be used as A probe for dominance of various FSI mechanisms

Theoretical Prediction



K. Miyagawa et al., Phys. Rev. C 74, 034002 (2006);

Preliminary Results: Beam Spin Asymmetry

Work by Nick Zachariou



Preliminary Results

Observables for Linearly-Polarized Photons $\vec{\gamma}d \rightarrow \frac{d\sigma}{d\Omega} = \frac{d\sigma}{\vec{\Lambda}(\Omega_{u})_{pp}} \left[1 - P_{lin} \sum \cos 2\varphi - P_{lin}O_{x}\alpha \cos \theta_{x} - P_{lin}O_{z}\alpha \cos \theta_{z} + ... \right]$



Hyperon-Nucleon Interaction: Summary

First ever estimates of a large set of polarization observables for FSI in the reaction $\gamma d \rightarrow K^+ \Lambda n$ have been obtained from JLab g13 experiment.

- One-, two-, three-, and four-fold differential estimates are obtained with adequate statistical uncertainty.
- Beam-spin asymmetry provides indication that $\Lambda n \rightarrow \Lambda n$ mechanism may be dominant at large θ_{Λ}' .
- Beam-spin asymmetry can be binned sufficiently fine at low $W_{\Lambda n}$ to be suitable for extraction of the Λn scattering length.
- Data have sufficient coverage and statistical significance to impact YN.

Future Perspectives: Studies at 12 GeV

Near-Threshold J/W Production off Deuteron Why Study?

- Determination of the elementary $J/\psi N$ total cross section ($I_F \sim 1 \text{ fm}$).
- Studies of the gluonic structure of deuteron
 - gluonic structure of short-range correlations
 - deuteron gluonic form factor
- Studies of the hidden color component of deuteron?



J.-M. Laget, Nucl. Phys. A **581**, 397 (1995) M. Sargsian, private communication

Near-Threshold Exclusive Incoherent Photoproduction: Theoretical Studies



(a) Impulse



(b) NN Re-scattering



(c) J/ ψ N Re-scattering

J.-J. Wu, T.-S. H. Lee, Phys. Rev. C 88, 015205 (2013)



J/ψ Photoproduction off Deuteron with CLAS 12

Fully Exclusive Measurement of Incoherent Photoproduction of J/ψ

Types of data samples:

- Quasi-real Photoproduction:
 e', p, J/ψ measured
 n reconstructed
 Q² ≥ 0.05 (GeV/c)²
- Untagged real photoproduction
 n, p, J/ψ measured
 v reconstructed



Summary

The CLAS detector at 6 GeV produced a wealth of **unpolarized and polarized data** that have been critical for the quantitative description of baryon structure and baryon dynamics

- Phenomenology of the transition from long to short scales in fewnucleon systems
- Final-State Interactions: key to access YN properties

The 12 GeV upgrade will allow to study charm baryons near threshold

- Gluonic structure of deuteron
- Elementary $J/\psi N$ elastic cross section

The End

Dynamics of Hard Photodisintegration in the Scaling Regime

Both, QGSM and HRM, models for $\gamma d \rightarrow pn$ describe well

measured experimental observables.

E_γ=2.0 GeV



Dynamics of Hard Photodisintegration in the Scaling Regime

Both, QGSM and HRM, models for $\gamma d \rightarrow pn$ describe well

measured experimental observables.





X. Jiang et al., Phys. Rev. Lett. 98, 182302 (2007)

AdS/CFT: Conformal Window

- At short distances, dimensional scaling laws reflect the scale independence of a_s (asymptotic freedom)
- At large distances, dimensional scaling laws reflect the existence of infrared fixed point of QCD: as is large but scaleindependent
- Scale-invariance is broken in the transition between these two dynamical regimes

S.J. Brodsky and G.F. de Teramond, Phys. Rev. D **77**, 056007 (2008); J. Polchinski and G.R. Strassler, Phys. Rev. Lett. **88**, 031601 (2002).



Two-Body Photodisintegration of ³He

Scaling of invariant cross sections

³He



$$s^{17} \frac{d\sigma}{dt} \sim \text{const.}$$

- Indication that above
 ~ 0.7 G e V d a t a
 consistent with scale
 invariance for all CM
 angles
- Onset of dimensional scaling depends on the momentum transfer to i n d i v i d u a l constituents: supports AdS/CFT hypothesis

Event Distribution over Missing Momentum

 $P_x (\gamma d \to K^+ \Lambda X)$



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Event Distribution over Missing Momentum

 $P_x (\gamma d \to K^+ \Lambda X)$



The removal of events with Px < 0.2 GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure.

⁴⁰⁰ Selection of Exclusive 200 Events

Event Distribution over Missing Momentum $P_{x} (\gamma d \rightarrow K^{+}\Lambda X)$ Event Distribution over Missing Mass M_x ($\gamma d \rightarrow K^+ \Lambda X$) mm(GeV/c²)



The removal of events with Px < 0.2 GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure.

counts 300





Figure from Nick Zachariou

⁴⁰⁰ Selection of Exclusive 200 Events

Event Distribution over Missing Momentum P_{χ} ($\gamma d \rightarrow K^+ \Lambda X$) Event Distribution over Missing Mass M_x ($\gamma d \rightarrow K^+ \Lambda X$) mm(GeV/c²)



The removal of events with Px < 0.2 GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure.

³⁰⁰



200

Figure from Nick Zachariou

⁴⁰⁰ Selection of Exclusive 200 Events

Event Distribution over Missing Momentum $P_{x} (\gamma d \rightarrow K^{+}\Lambda X)$ Event Distribution over Missing Mass M_x ($\gamma d \rightarrow K^+ \Lambda X$) mm(GeV/c²)



The removal of events with Px < 0.2 GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure.

counts 300





Figure from Nick Zachariou

⁴⁰⁰ Selection of Exclusive 200 Events

Event Distribution over Missing Momentum $P_{x} (\gamma d \rightarrow K^{+}\Lambda X)$ Event Distribution over Missing Mass M_x ($\gamma d \rightarrow K^+ \Lambda X$) mm(GeV/c²)



The removal of events with Px < 0.2 GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure.





200

Figure from Nick Zachariou

⁴⁰⁰ Selection of Exclusive 200 Events

Event Distribution over Missing Momentum $P_r (\gamma d \rightarrow K^+ \Lambda X)$ Event Distribution over Missing Mass M_x ($\gamma d \rightarrow K^+ \Lambda X$) mm(GeV/c²)



The removal of events with Px < 0.2 GeV/c provides a sample that is by far dominated by FSI events. Standard analysis procedure.

Removal of physics background based on ³⁰⁰ realistic simulation of reactions, detector, and accidentals, followed by histogram or event-by-event fits for each kinematic bin. ²⁰⁰

Figure from Tongtong Cao

Figure from Nick Zachariou

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