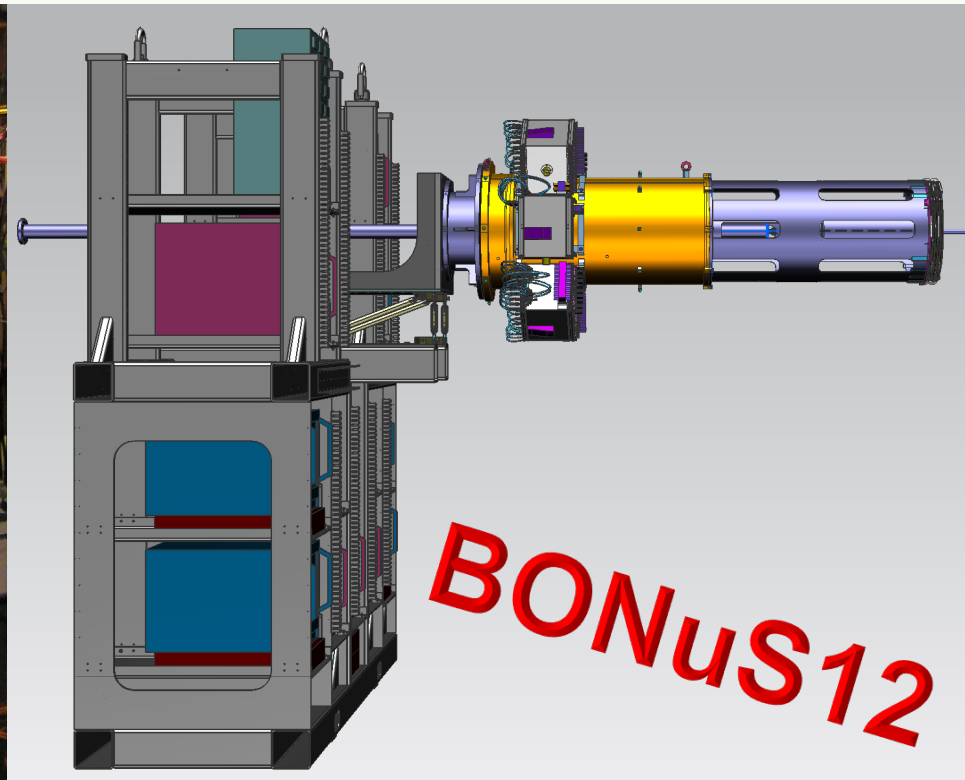
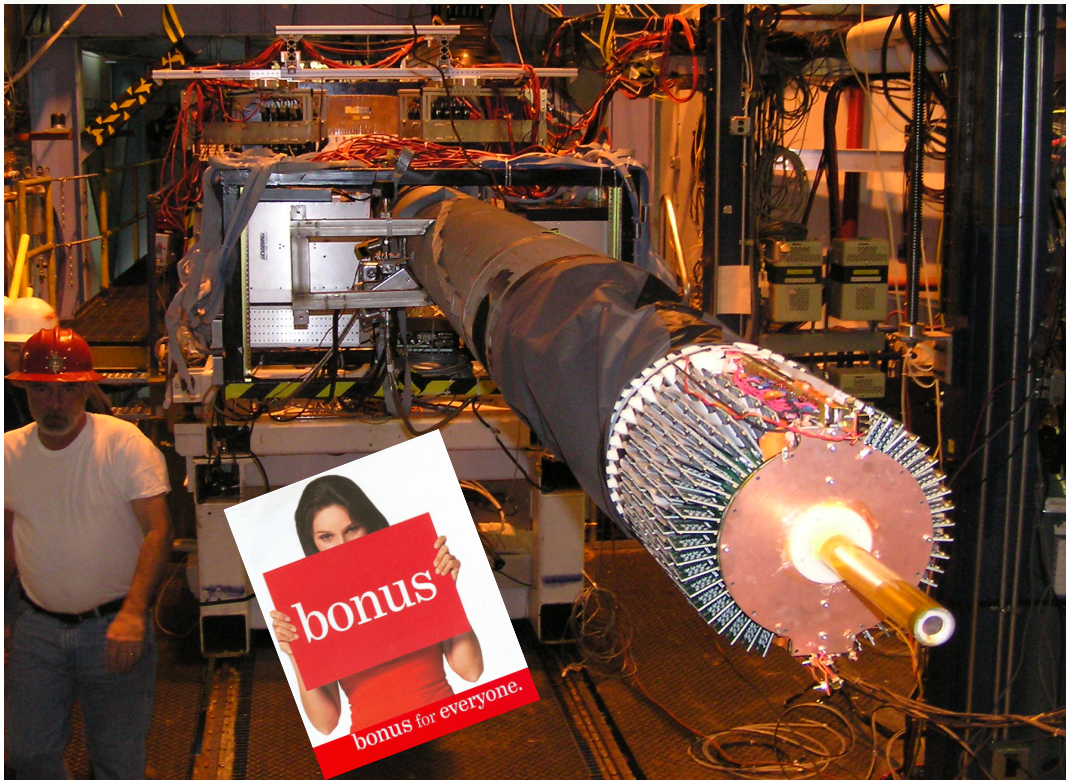


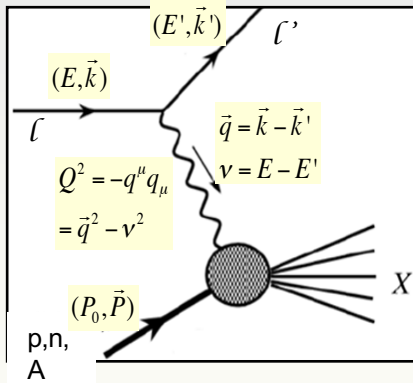
The BONuS measurements of the free neutron structure function



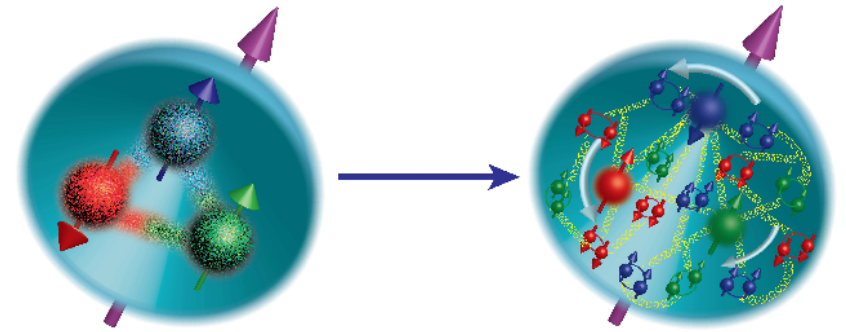
Sebastian Kuhn
Old Dominion University

Overview

- Neutron Structure Functions (esp. at large x) - Why?
- The Neutron - No Free ~~Lunch~~ Target
(Nucleon structure modifications in Nuclei)
- Spectator Tagging
(Principle and Experimental Realization - the RTPC)
- The “BONuS” experiment
- New recoil detectors
- The (11 GeV) Future of “BONuS”
- Conclusion and Outlook



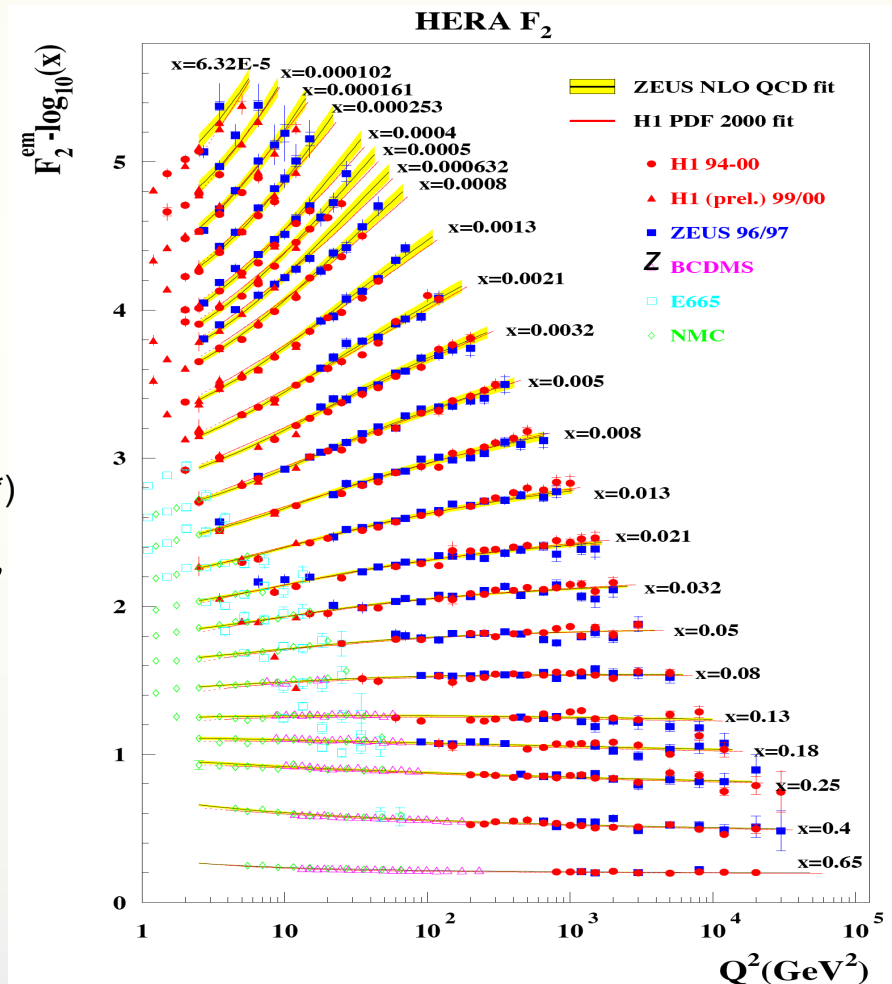
Introduction



- The familiar (?) 1D world of Nucleon longitudinal structure:

- Take a nucleon
- Move it real fast along z
 \Rightarrow light cone momentum
 $P_+ = P_0 + P_z (>>M)$
- Hit a “parton” (q, g, \dots) inside
- Measure **its** l.c. momentum
 $p_+ = p_0 + p_z (m \approx 0)$
- \Rightarrow Momentum Fraction $\xi = p_+ / P_+^*$
- In DIS: $\xi = (q_z - v)/M \approx x_{Bj} = Q^2/2Mv$
- Probability: $F_1(x) = \frac{1}{2} \sum_i e_i^2 q_i(x)$
- Because of spin-1/2: 2nd SF $F_2(x)$

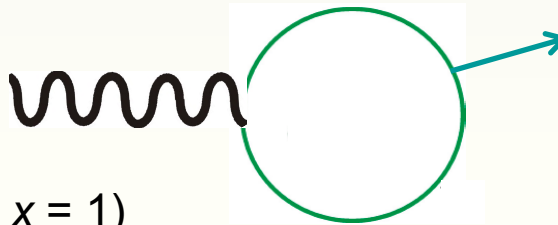
*) Advantage: Boost-independent



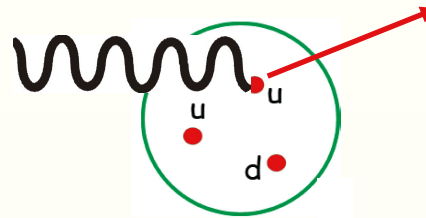
⇒ Our 1D View of the Nucleon

(also depends on the resolution of the virtual photon $\sim 1/Q^2$)

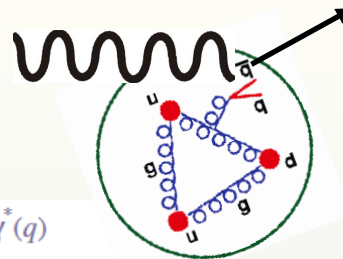
- Elastic scattering
(Whole system recoils, $x = 1$)



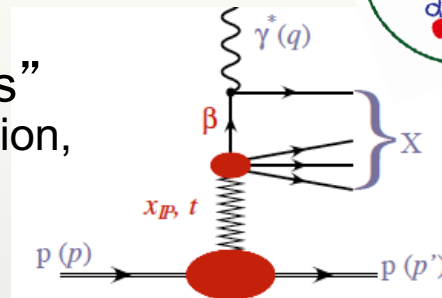
- Resonances
($x < 1$, $W < 2$ GeV)



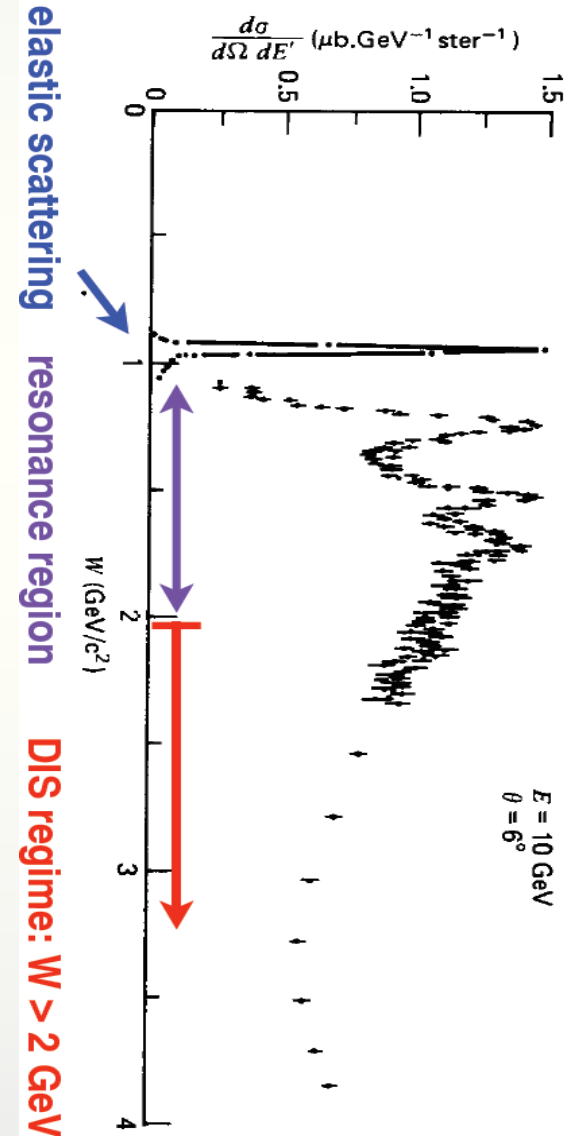
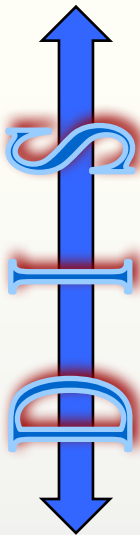
- Valence quarks
($x \approx 0.3 - 0.9$, $W > 2$ GeV)



- Sea quarks, gluons
($x < 0.3$)



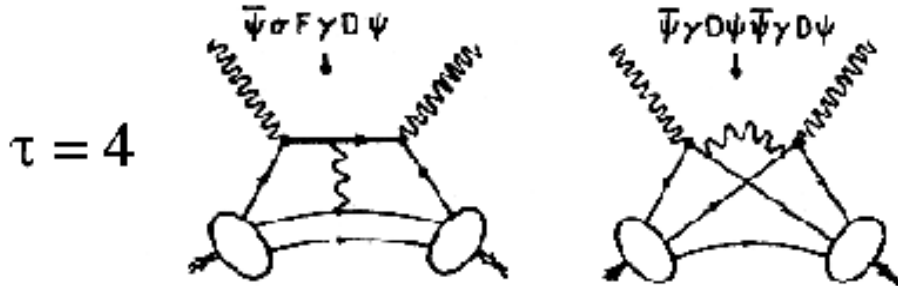
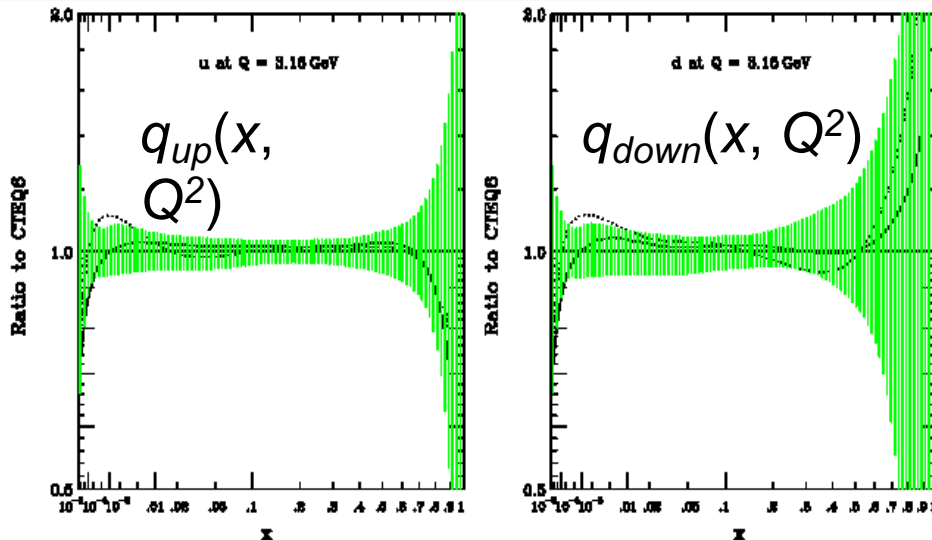
- “Wee Partons”
($x \rightarrow 0$, Diffraction, Pomerons)



Structure Functions and Moments:

Why large x?

$$\frac{d\sigma}{d\Omega dE'} = \sigma_{Mott} \left(\frac{F_2(x)}{\nu} + 2 \tan^2 \frac{\theta_e}{2} \frac{F_1(x)}{M} \right); \quad F_2(x, Q^2) = x \sum_{f=up, down, \dots} z_f^2 (q_f(x, Q^2) + \bar{q}_f(x, Q^2))$$



- $q_{down}/q_{up}(x \rightarrow 1)$ is a crucial test of valence quark models
 - SU(6) breaking, pQCD, DSE,...
- Precise PDFs at large x needed as input for LHC, ν experiments etc.
 - Large x, medium Q^2 evolves to medium x, large Q^2
 - Also: NUCLEAR structure functions
- Moments can be directly compared with OPE (twist expansion), Lattice QCD and Sum Rules
 - All higher moments are weighted towards large x
- Quark-Hadron Duality

$$M_n^{CN}(Q^2) = \int_0^1 dx \underline{x^{(n-2)}} F_2(x, Q^2) = \sum_{\tau=2k}^{\infty} E_{n\tau}(\mu, Q^2) O_{n\tau}(\mu) \left(\frac{\mu^2}{Q^2} \right)^{\frac{1}{2}(\tau-2)} + \text{TM corr.}$$

Why neutron?

$d(x)$ and $u(x)$ as $x \rightarrow 1$

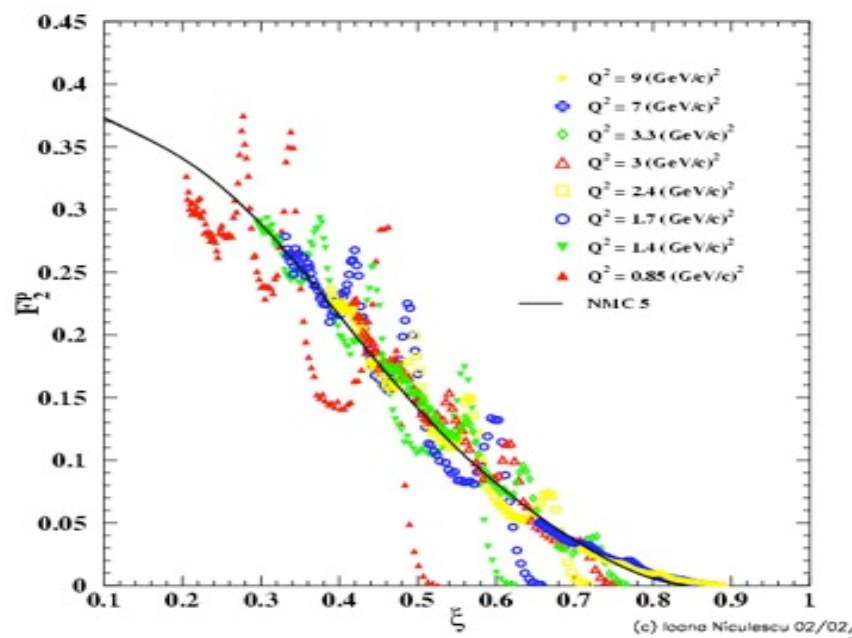
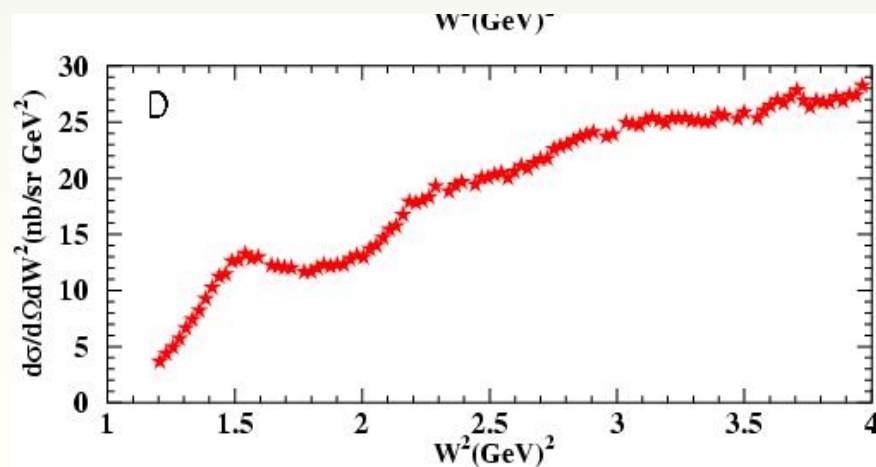
- Valence structure of the nucleon - sea quarks and gluons don't contribute
- SU(6)-symmetric wave function of the proton in the quark model:

$$|p \uparrow\rangle = \frac{1}{\sqrt{18}} \left(3u \uparrow [ud]_{S=0} + u \uparrow [ud]_{S=1} - \sqrt{2}u \downarrow [ud]_{S=1} - \sqrt{2}d \uparrow [uu]_{S=1} - 2d \downarrow [uu]_{S=1} \right)$$

- In this model: $d/u = 1/2$, $\Delta u/u^*) = 2/3$, $\Delta d/d = -1/3$ for all x
- Relativistic quark model: quark helicities reduced, orbital angular momentum introduced
- Hyperfine structure effect (1-gluon exchange): $S=1$ suppressed for small spectator pair mass $\Rightarrow d/u = 0$, $\Delta u/u = 1$, $\Delta d/d = -1/3$ for $x \rightarrow 1$
- pQCD: helicity conservation ($q \uparrow \uparrow p$) $\Rightarrow d/u = 2/(9+1) = 1/5$, $\Delta u/u = 1$, $\Delta d/d = 1$ for $x \rightarrow 1$
- Wave function of the neutron via isospin rotation:
replace $u \rightarrow d$ and $d \rightarrow u \Rightarrow$ using experiments with protons and neutrons one can extract information on u , d , Δu and Δd in the valence quark region.

*) spin dependent quark density $\Delta q = (q \uparrow - q \downarrow)$ for Nucleon $N \uparrow$

Structure Functions and Resonances



- Precise structure functions in Resonance Region constrain nucleon models
[Separate resonant from non-resonant background; isospin decomposition]
- Needed as input for spin structure function data, radiative corrections, ...
- Compare with DIS structure functions to test duality

Present Knowledge of d/u ($x \rightarrow 1$)

Assuming charge
independence

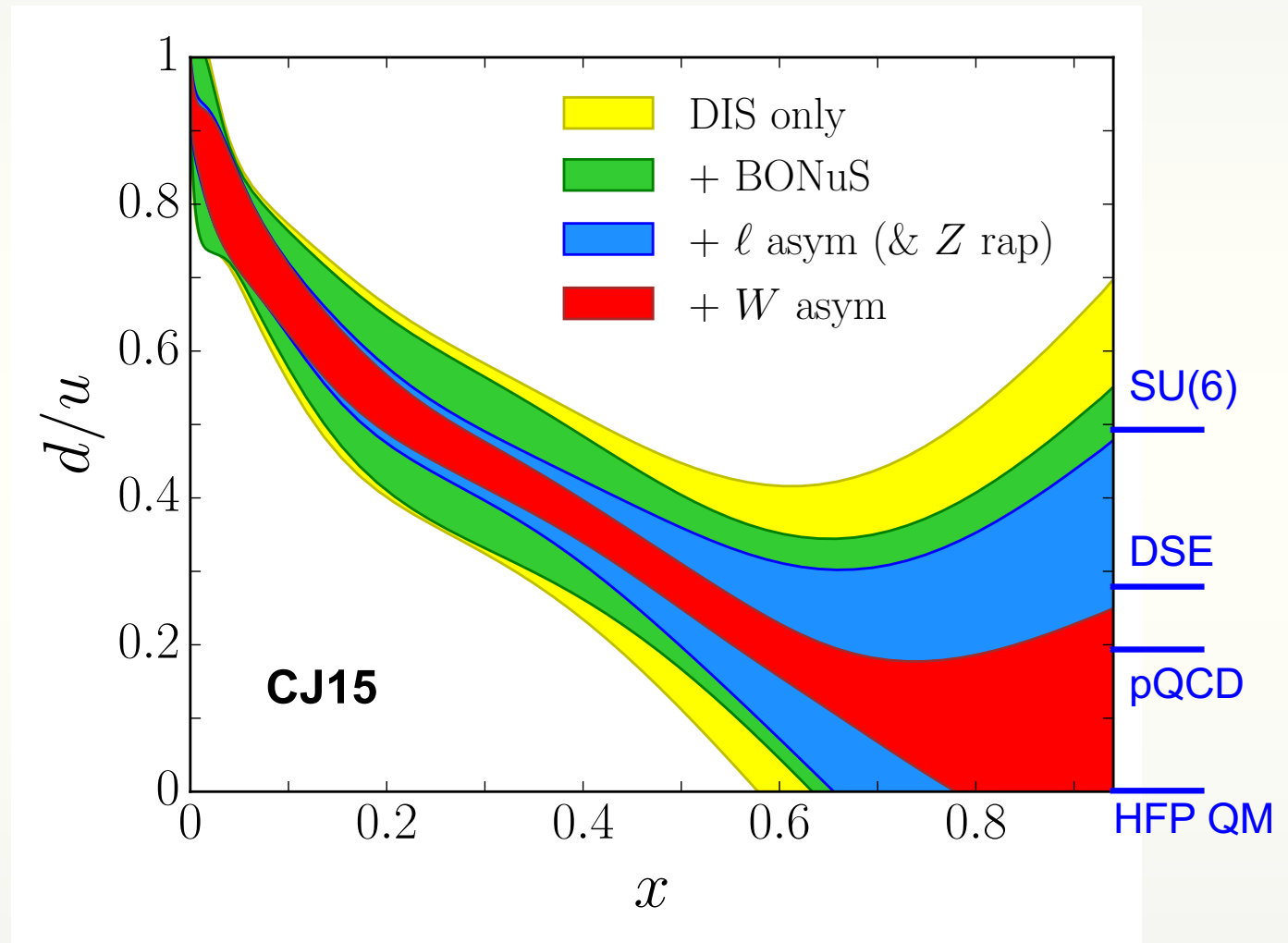
(= invariance under 180°
rotations in isospin space):

$$\frac{F_{2n}}{F_{2p}} \approx \frac{1 + 4d/u}{4 + d/u} \Rightarrow$$

$$\frac{d}{u} \approx \frac{4 F_{2n}/F_{2p} - 1}{4 - F_{2n}/F_{2p}}$$

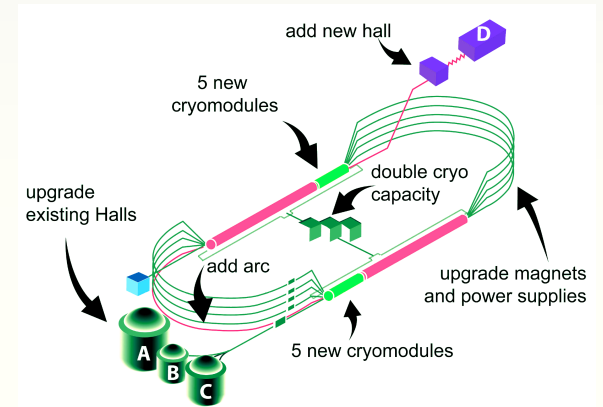
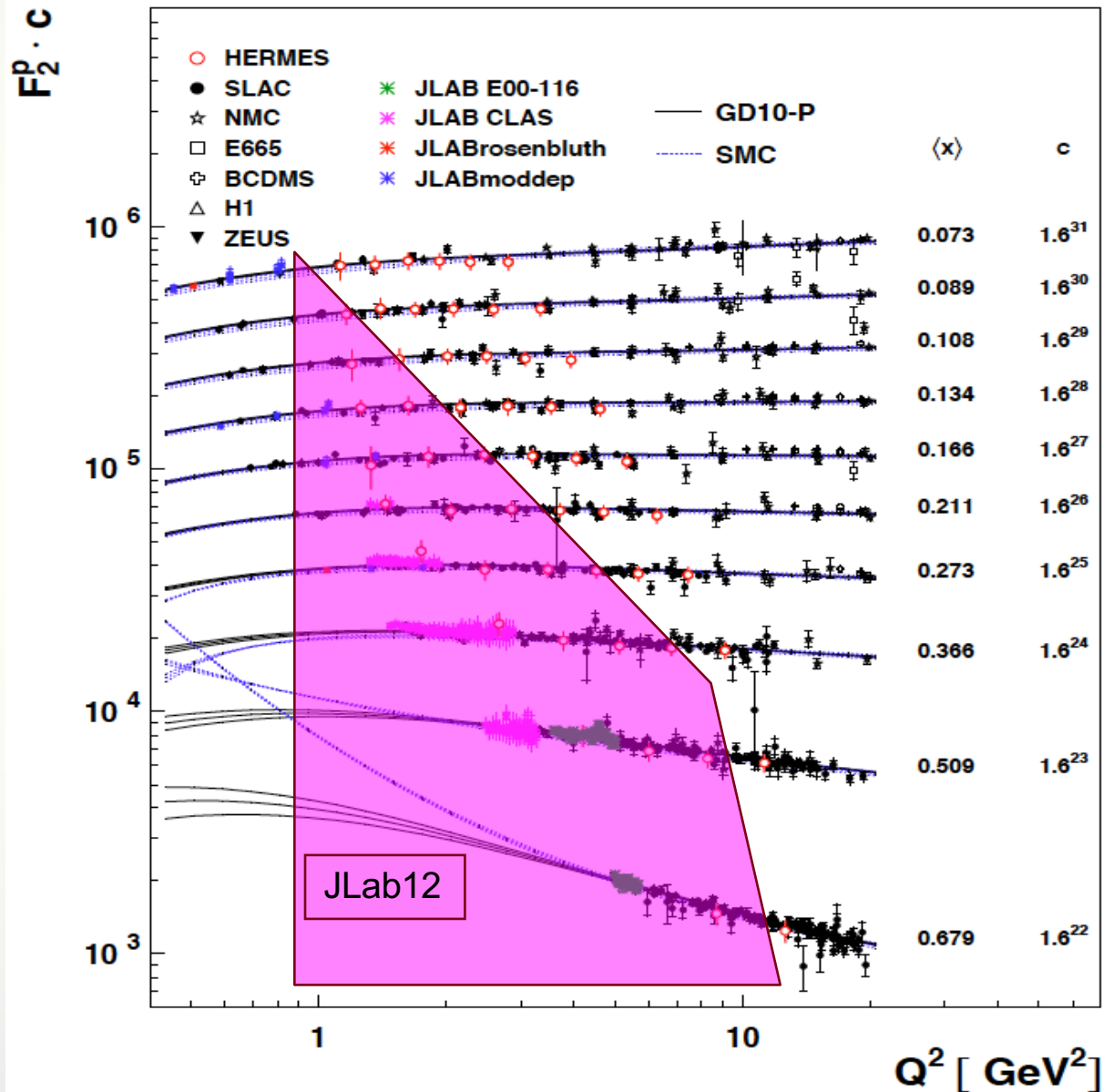
$$F_{2n}/F_{2p} = F_{2d}/F_{2p} - 1$$

???



- Neutron data limited by “Nuclear Binding Uncertainties”

Jefferson Lab in Context



JLab DIS

NOW: 12 GeV

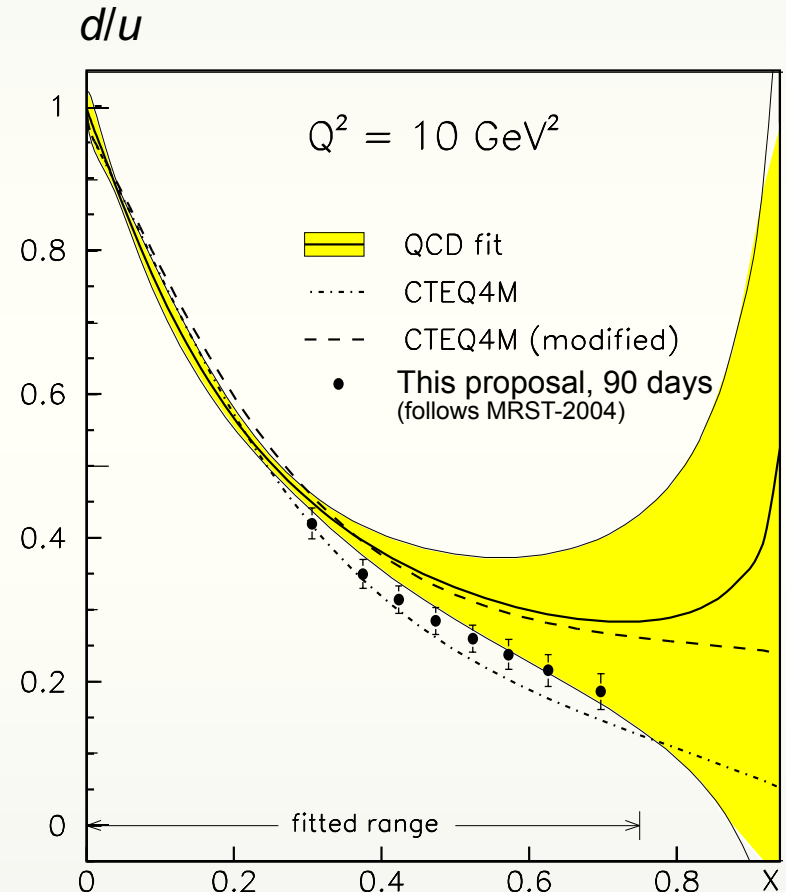
$Q^2 = 1 \dots 13 \text{ GeV}^2$

$x = 0.06 \dots 0.8$

$W = 0.94 \dots 4 \text{ GeV}$

Cleanest way to access d/u

- Exploit different “charge” ratios for weak and electromagnetic interaction.
- Possible processes: W/Z production, neutrino \rightarrow muon scattering, parity-violating lepton scattering (PV DIS).
- Advantage: Direct measurement on the proton; does not require assumptions about charge symmetry.
- Limitations in statistical precision.



PV DIS on p target, 90 days with SoLID. E12-10-007 in Jefferson Lab's Hall A, approved with A rating. Awaiting funding...

Neutron Data Are Important... ...but hard to get

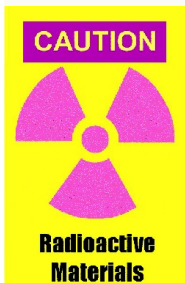
- Free neutrons decay in 15 min.

- Radioactivity!

- Zero charge makes it difficult to create a dense target

Magnetic bottle: $10^3 - 10^4$ n/cm² [TU München]

Typical proton target: $4 \cdot 10^{23}$ p/cm² [10 cm LH] – 10^{14} p/cm² [HERMES]



An Underground Nuclear Explosion as a Polarized Neutron Source

G. A. Keyworth and J. R. Lemley

**Polarization Phenomena
in Nuclear Reactions
Proceedings (1970)**

"Although an underground nuclear explosion is not the most conventional neutron source, it offers definite advantages..."

- => Alternative Solution: Deuterons, Tritons and Helium-3...

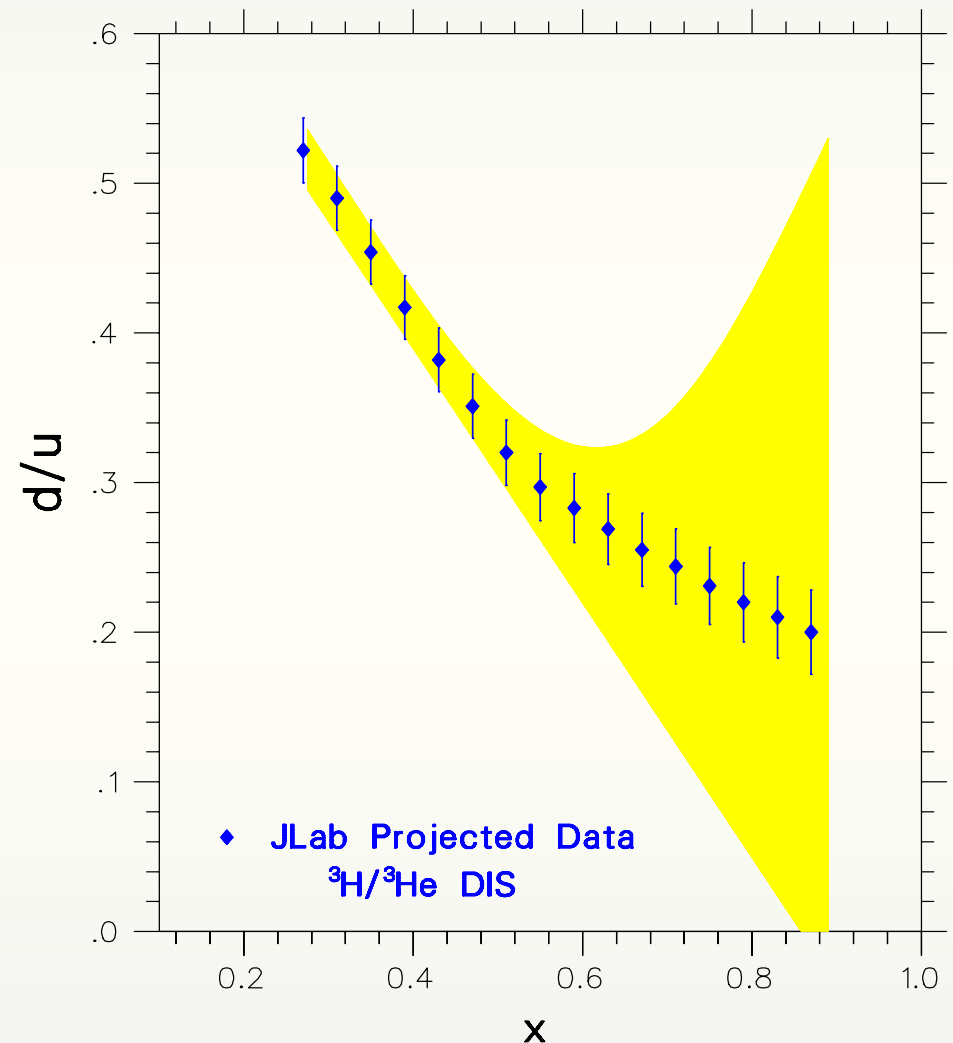
BUT: Nuclear Model Uncertainties:

Fermi motion, off-shell effects (binding), structure modifications (EMC effect), extra pions/Deltas, coherent effects, 6-quark bags...

One Solution: take ratio of nearly identical nuclei -> EMC effect largely cancels *)

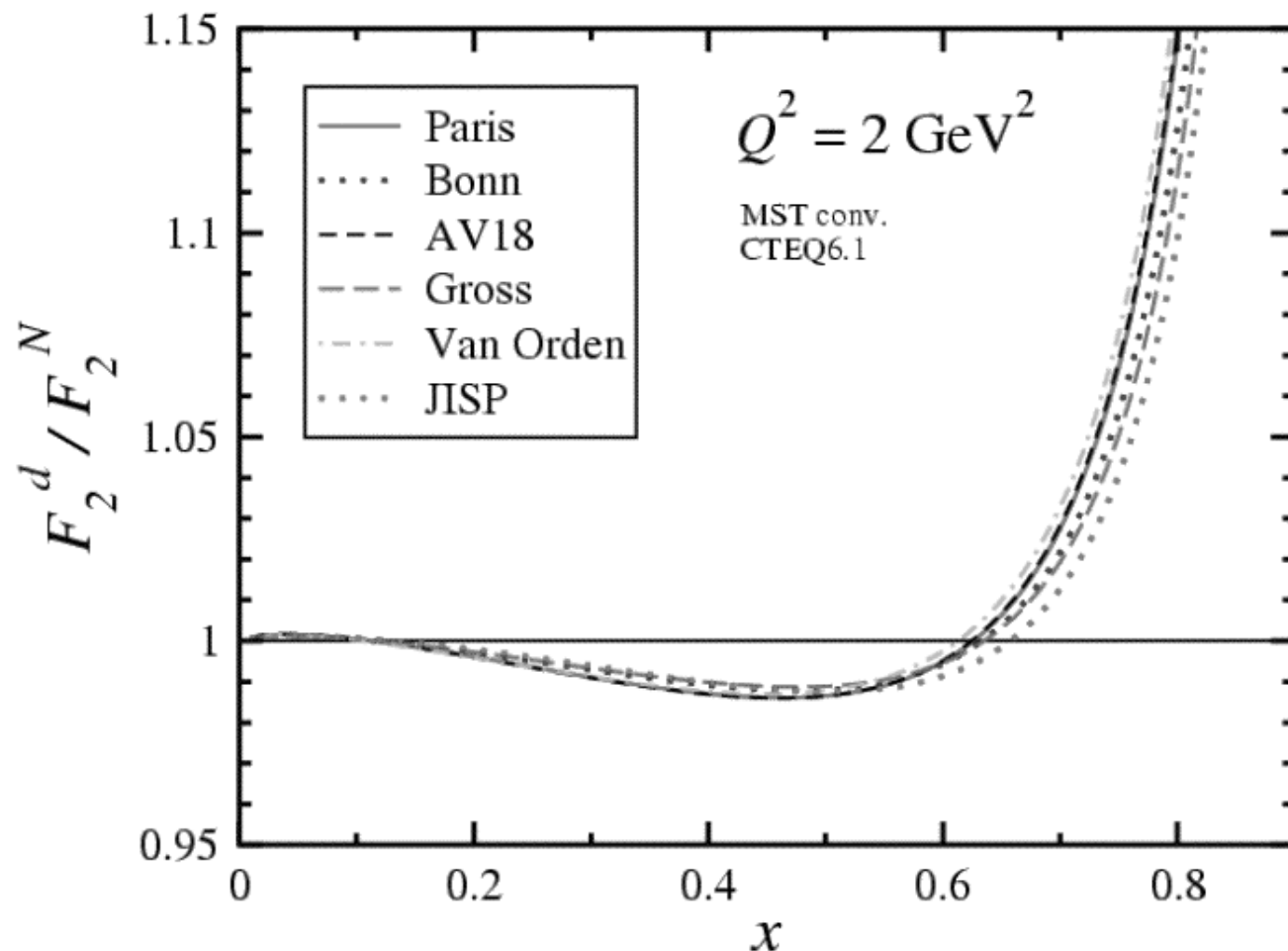
- Best case: Isospin doublet ${}^3\text{He}/{}^3\text{H}$.
- $$\frac{F_2^{3H}(x)}{F_2^{3He}(x)} = \frac{\left(\frac{4}{9} + 2\frac{1}{9}\right)u(x) + \left(\frac{1}{9} + 2\frac{4}{9}\right)d(x)}{\left(2\frac{4}{9} + \frac{1}{9}\right)u(x) + \left(2\frac{1}{9} + \frac{4}{9}\right)d(x)} = \frac{\frac{2}{3} + d(x)/u(x)}{1 + \frac{2}{3}d(x)/u(x)}$$
- Several experiments with tritium target planned for 2017 in Jefferson Lab's Hall A

*) But still measuring smeared SF, including smeared quasi-elastic!



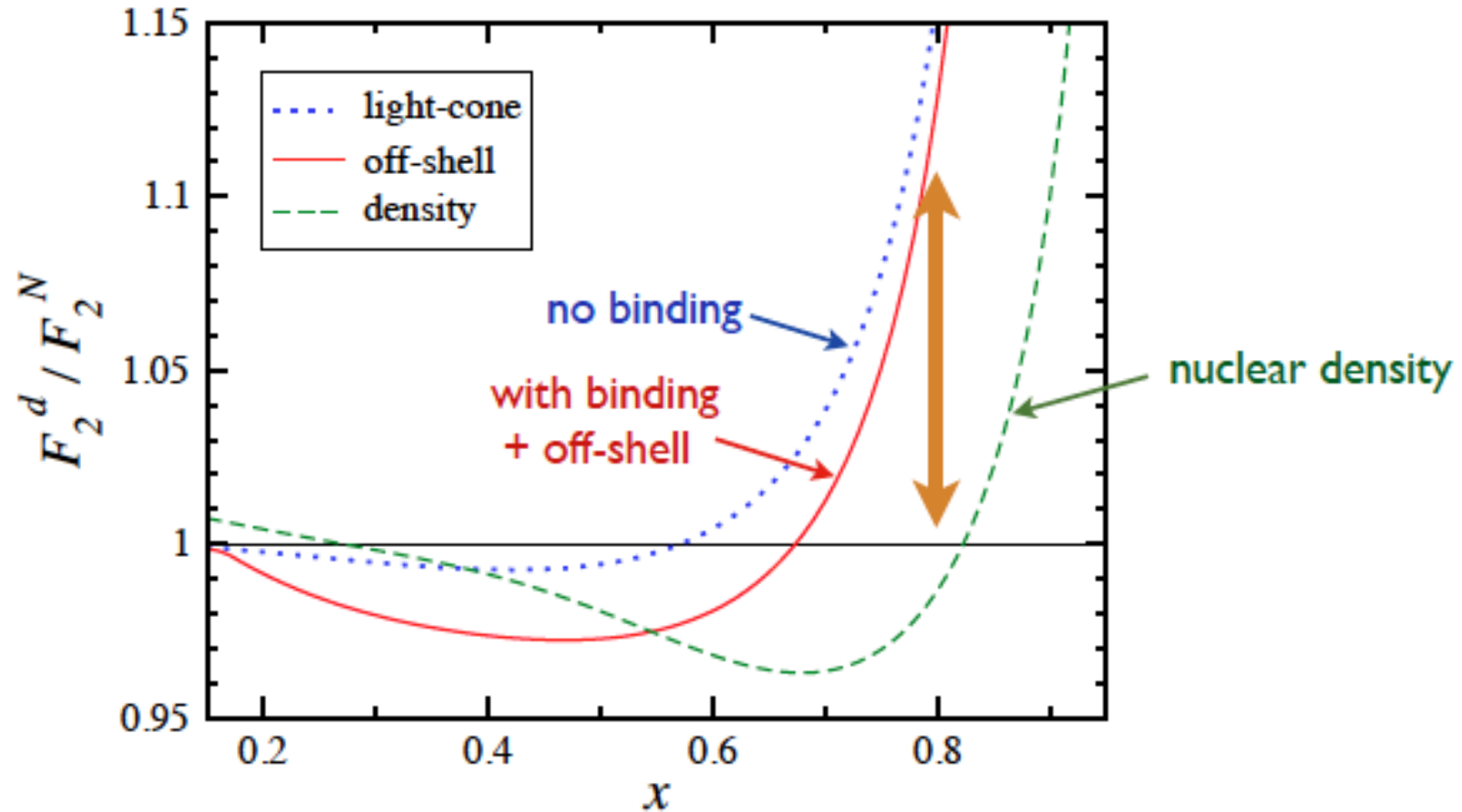
"Marathon" Experiment in Hall A. $W > 1.8$ GeV.
Experiment E12-10-103 42d, A rating, * from PAC41.
Scheduled to run (partially?) Fall/Winter 2017-18

Inclusive scattering on D: Large x - Large Nuclear Effects



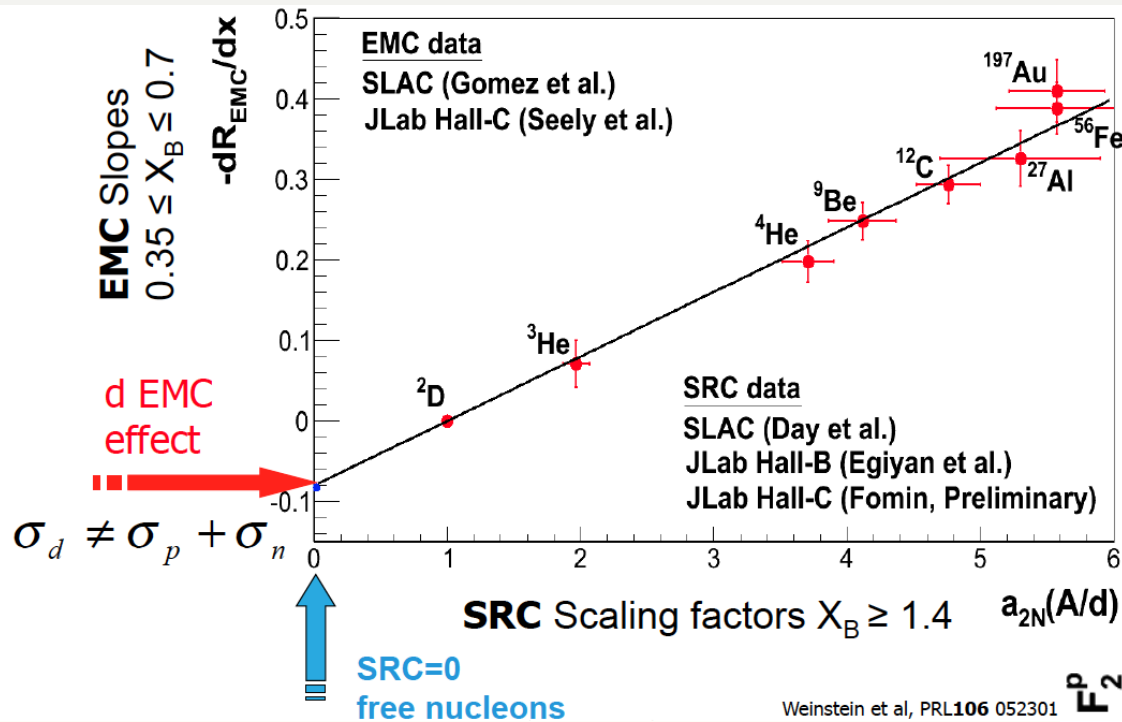
- Even simple “Fermi Smearing” leads to significant dependence on D wave function
- Different models for off-shell and “EMC” effects lead to large additional variations
- Contributions from MEC, $\Delta(1232)$ and “exotic” degrees of freedom unknown
- FSI?

EMC effect in deuteron

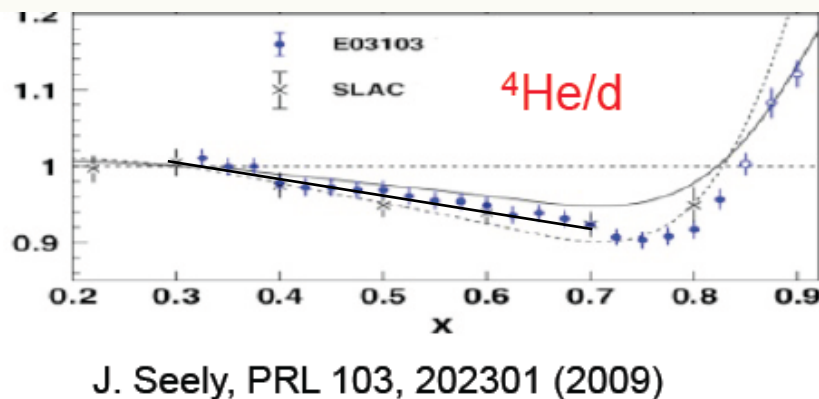


- using off-shell model, will get *larger* neutron *cf. light-cone* model
- but will get *smaller* neutron *cf. no nuclear effects* or *density* model

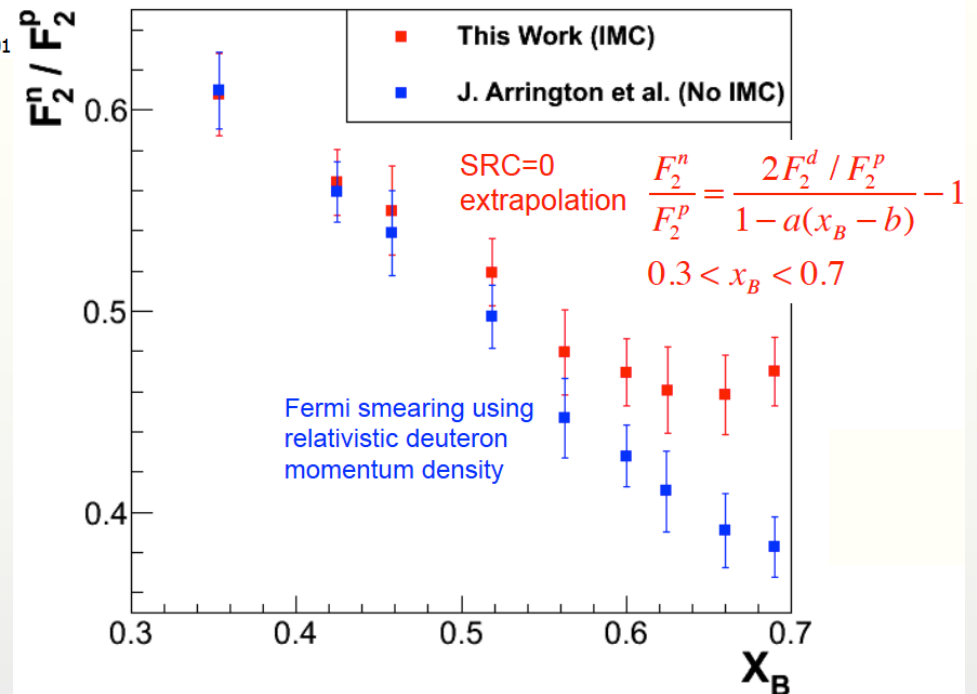
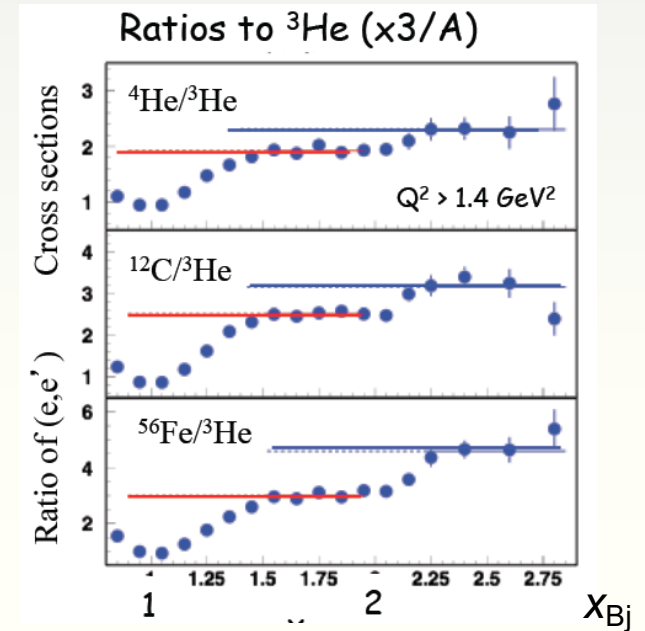
Estimating the EMC effect in Deuterium



Probability of a nucleon inside the nucleus to be in a “short-range” (tensor) correlation (dominated by pn correlations 10:1)

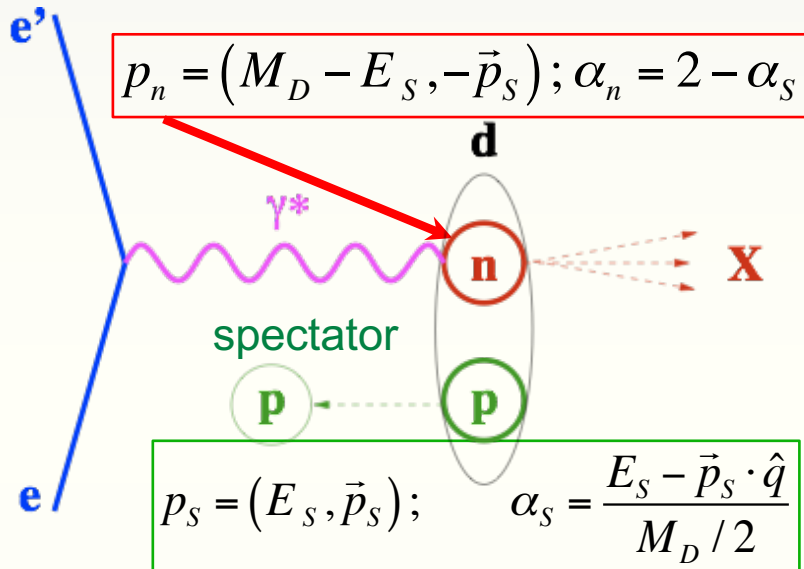


L. Weinstein, E.I. Piasetzky, D. Higinbotham, J. Gomez, O. Hen and R. Shneor, PRL106 052301 (2011)



The Solution: Spectator Tagging

$d(e, e' p_s) X$



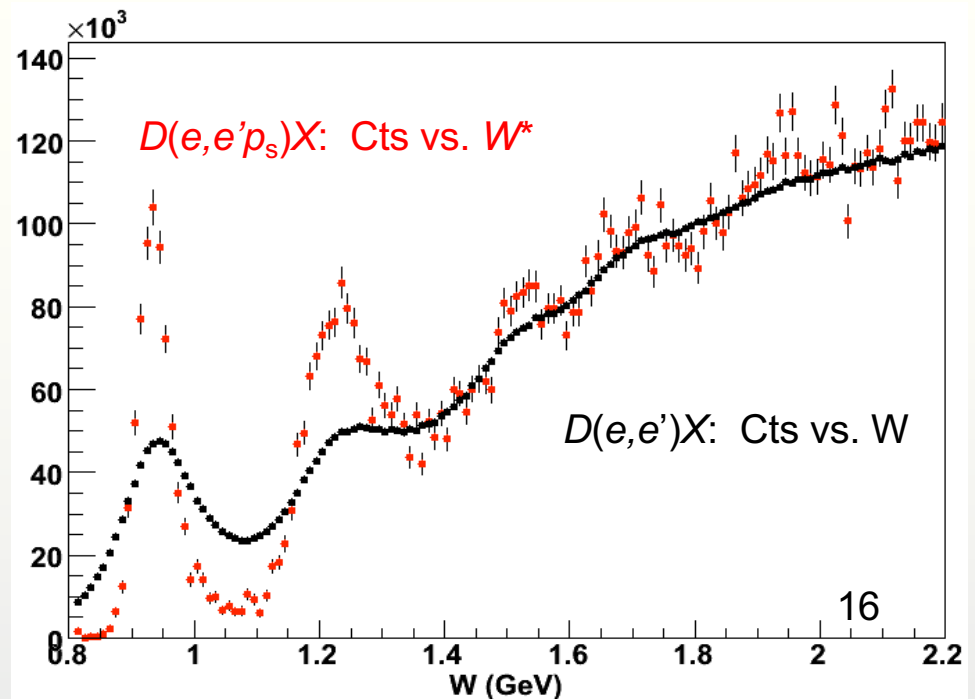
Relativistic Invariants

$$x = \frac{Q^2}{2p_n^\mu q_\mu} \approx \frac{Q^2}{2M\nu(2 - \alpha_s)} \quad Q^2$$

$$M^{*2} = p_n^\mu p_{n\mu} \approx \left(M_n - \varepsilon - \frac{\vec{p}_s^2}{M_n} \right)^2 \approx M_n^2 - 2M_n\varepsilon - 2\vec{p}_s^2$$

$$W^{*2} = (p_n + q)^2 = M^{*2} + 2((M_D - E_s)\nu - \vec{p}_n \cdot \vec{q}) - Q^2$$

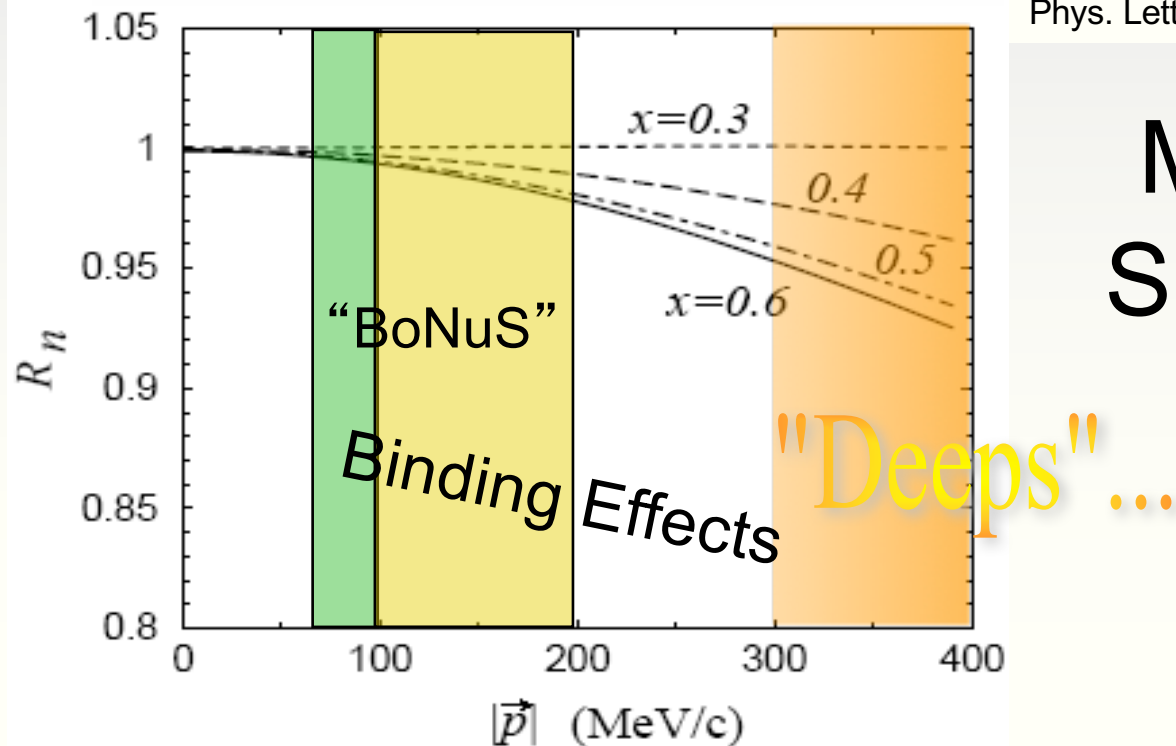
$$\approx M^{*2} + 2M\nu(2 - \alpha_s) - Q^2$$



$$R_n \equiv F_2^{n(eff)}(W^2, Q^2, p^2) / F_2^n(W^2, Q^2)$$

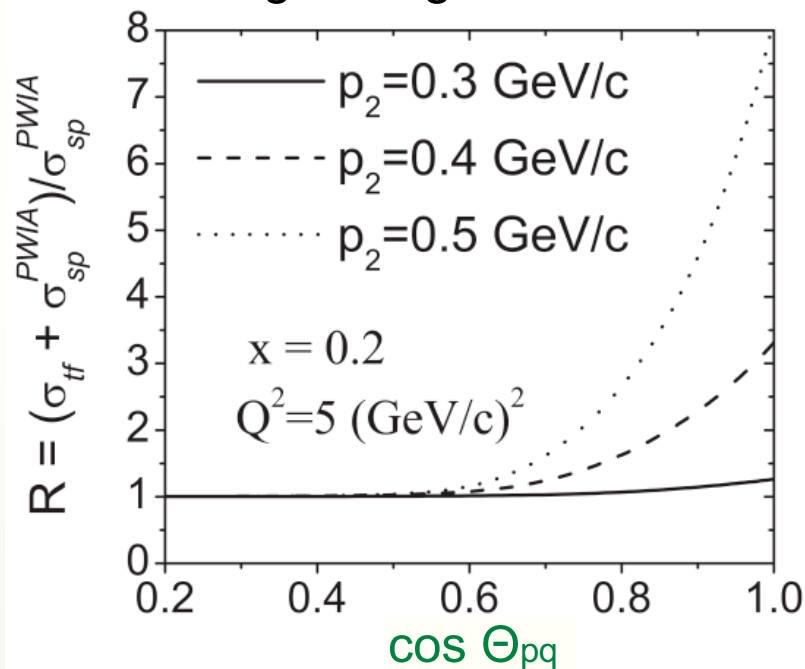
W. Melnitchouk, A.W. Schreiber and A.W. Thomas,

Phys. Lett. B335, 11 (1994); Phys. Rev. D 49, 1183 (1994).



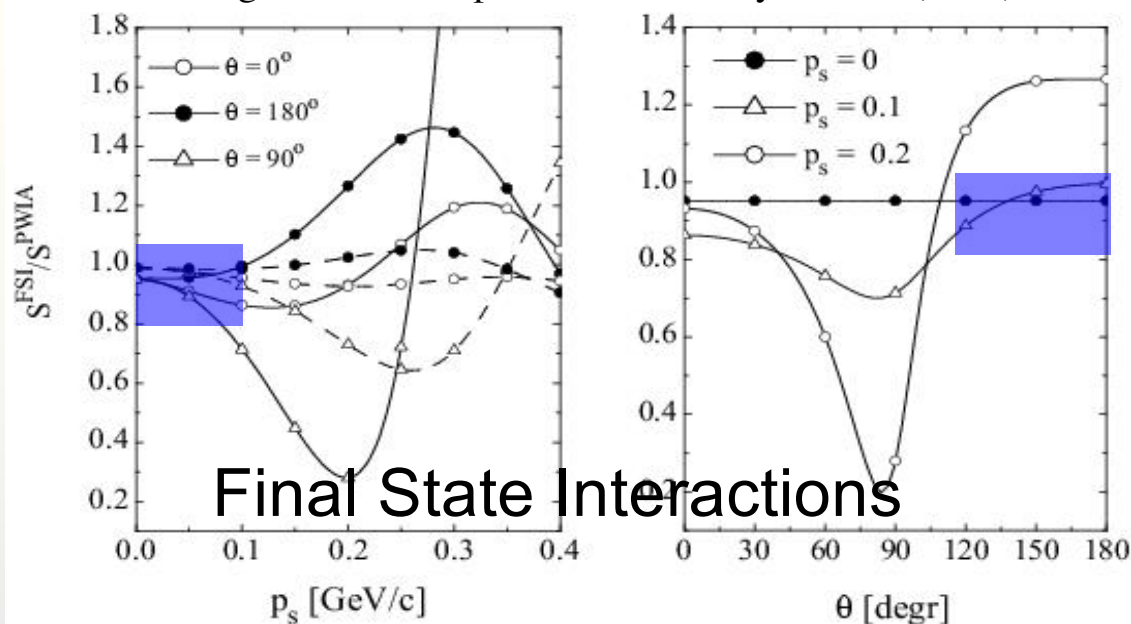
Modifications to Simple Spectator Picture

Target Fragmentation



Palli et al, PRC80(09)054610

Ciofi degli Atti and Kopeliovich, Eur. Phys. J. A17(2003)133



Final State Interactions

Alternative: Pole extrapolation

$$t' = M^2 - M_n^2 = (P_D - p_s)^2 - M_n^2 = \left(M_D - \sqrt{M_p^2 + \vec{p}_s^2} \right)^2 - \vec{p}_s^2 - M_n^2 \approx -2(M_n \varepsilon + \vec{p}_s^2); \varepsilon = 2.2 \text{ MeV}$$

Wim Cosyn (Gent U.), Misak Sargsian (Florida Intl. U.)

Mar 2, 2016 - 6 pages

EPJ Web Conf. 112 (2016) 03001
(2016-03-21)

DOI: [10.1051/epjconf/201611203001](https://doi.org/10.1051/epjconf/201611203001)

Conference: C15-09-07.1

Proceedings

e-Print: [arXiv:1603.00685](https://arxiv.org/abs/1603.00685) [nucl-th] | [PDF](#)

- Measure F_{2n} at fixed spectator angle, but varying momentum
- Extrapolate to on-shell neutron

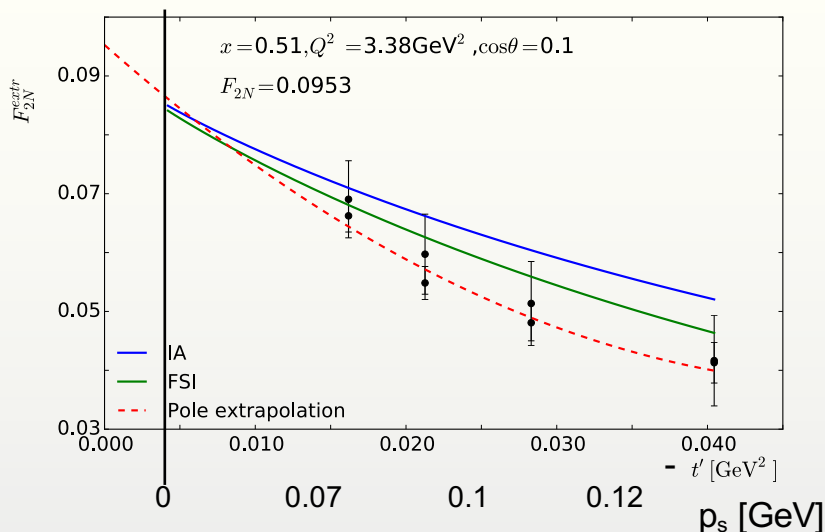


Figure 3. Example of the pole extrapolation method using the renormalized BONuS data (black circles) with the quadratic pole extrapolation curve (red dashed curve) as a function of $t'^2 = p_i^2 - m_n^2$. The IA (full blue curve) and FSI (full green curve) calculations are shown for comparison.

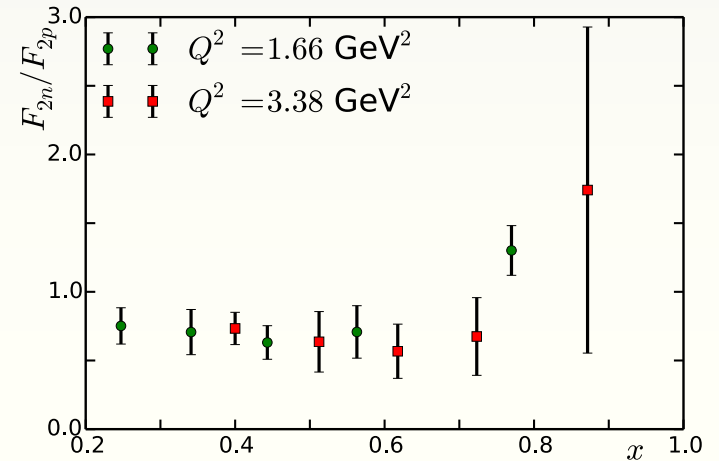
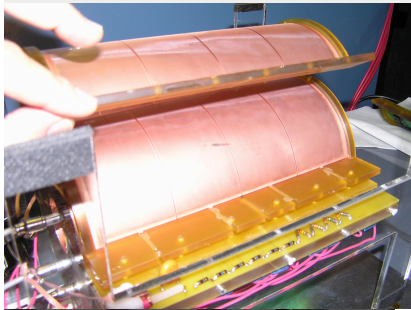
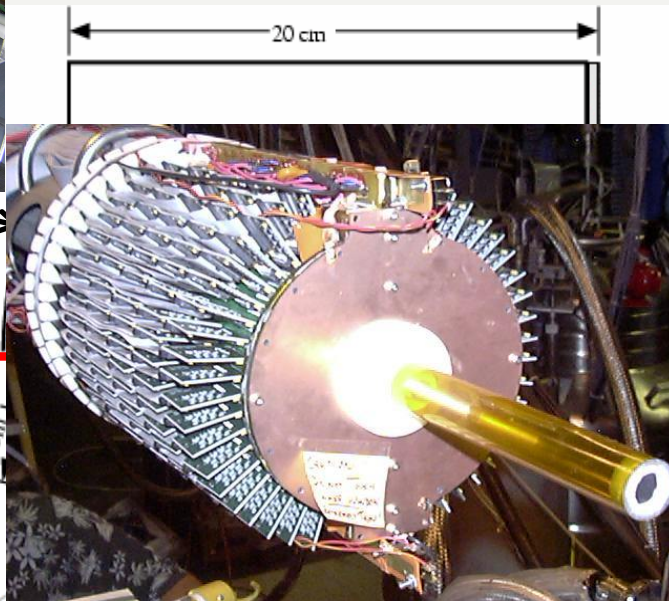


FIG. 4: (Color online) F_{2n} to F_{2p} ratio obtained using the pole extrapolation method on the renormalized BONuS data for $Q^2 = 1.66$ (green circles), 3.38 GeV^2 (red squares). The F_{2p} values are estimated using fit of Ref. [21].

BoNuS RTPC



Gas
Electron
Multiplier

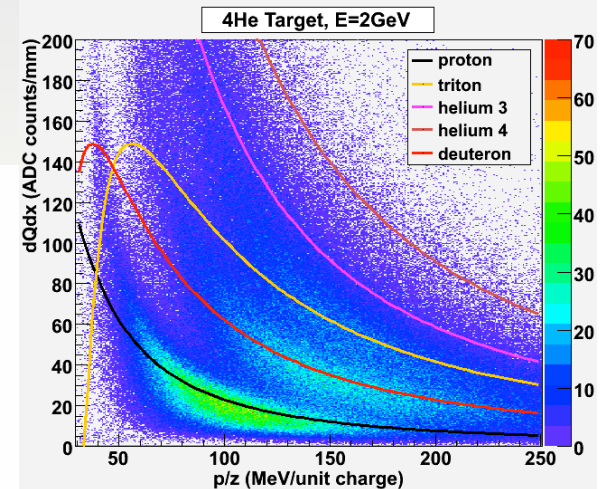


7 atm D₂ gas

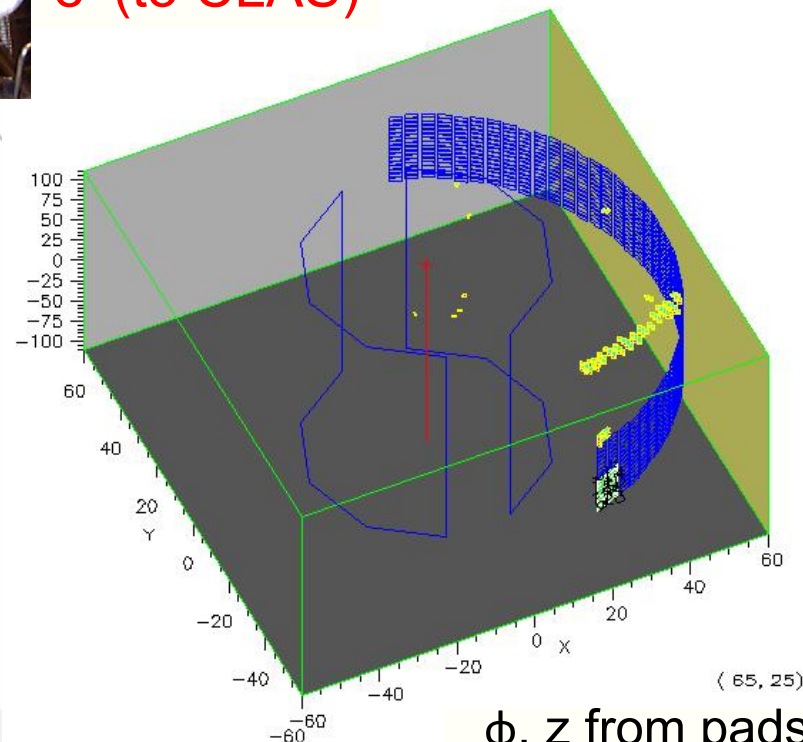
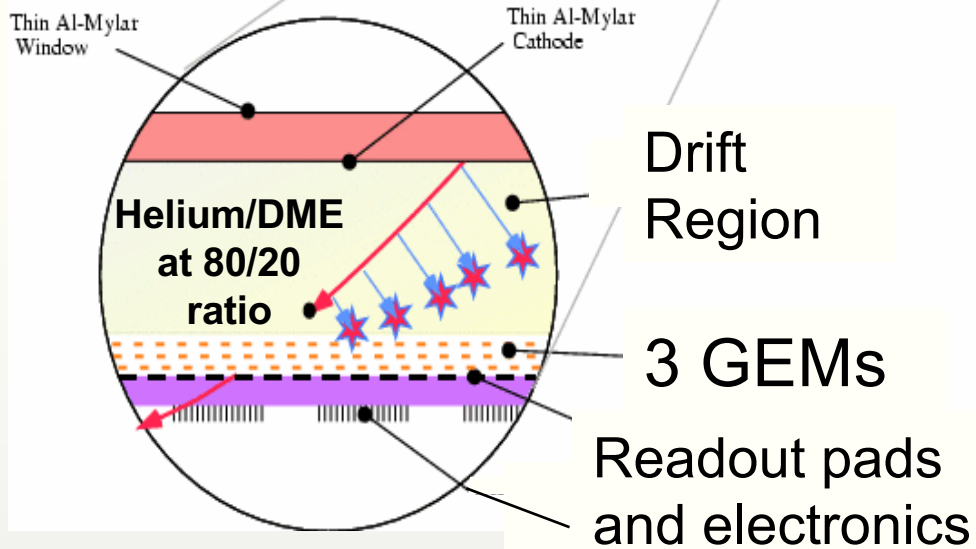
Thin-wall High Pressure
Gas Target

Møller el.

e⁻ (to CLAS)

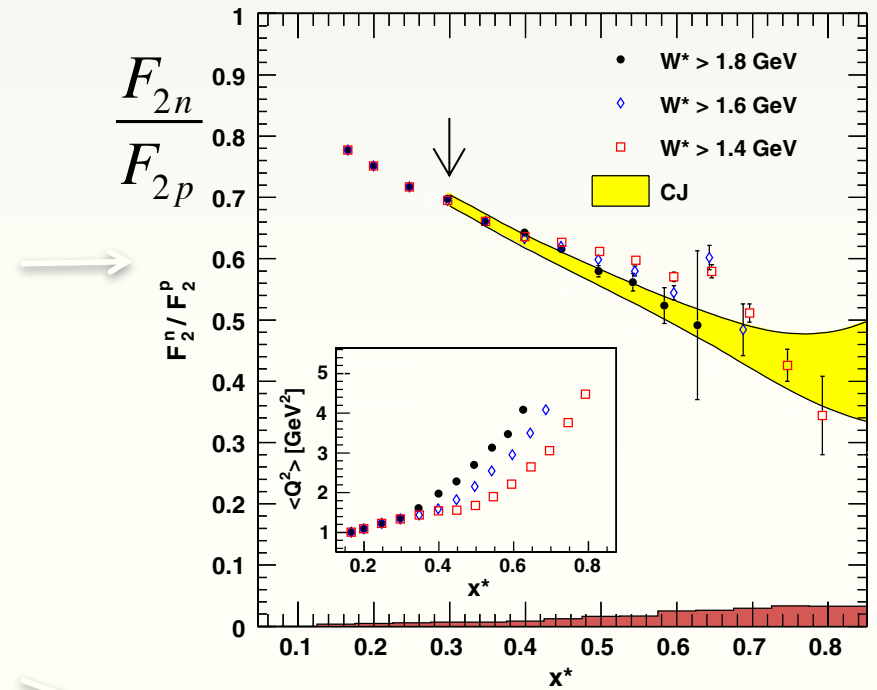
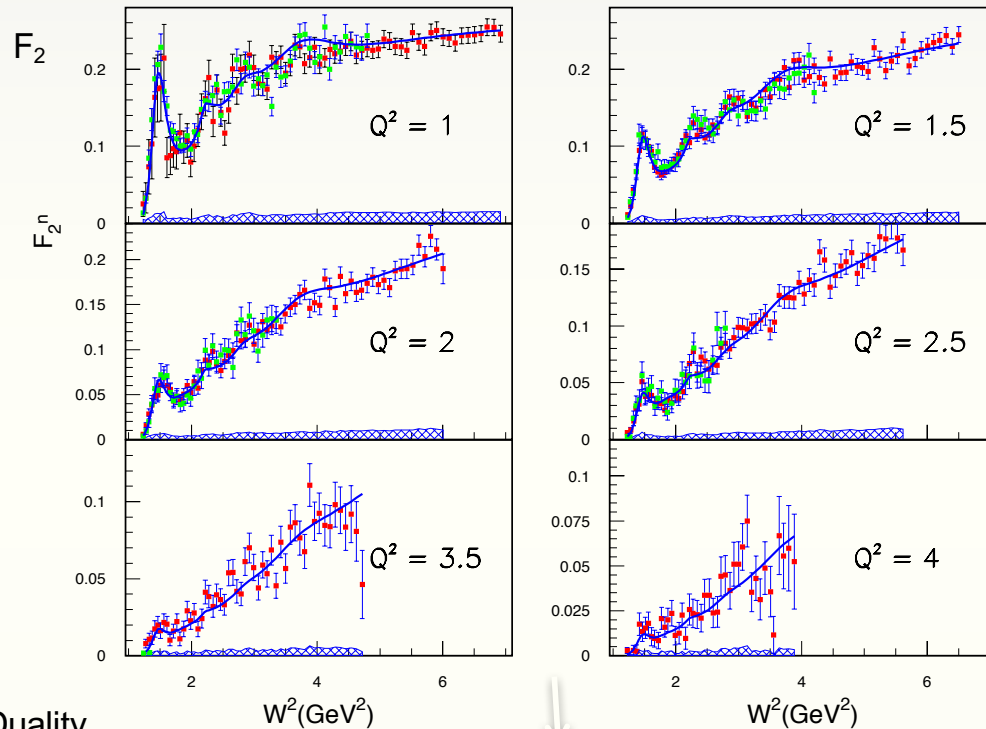


dE/dx from charge
along track (particle ID)



ϕ , z from pads
 r from time

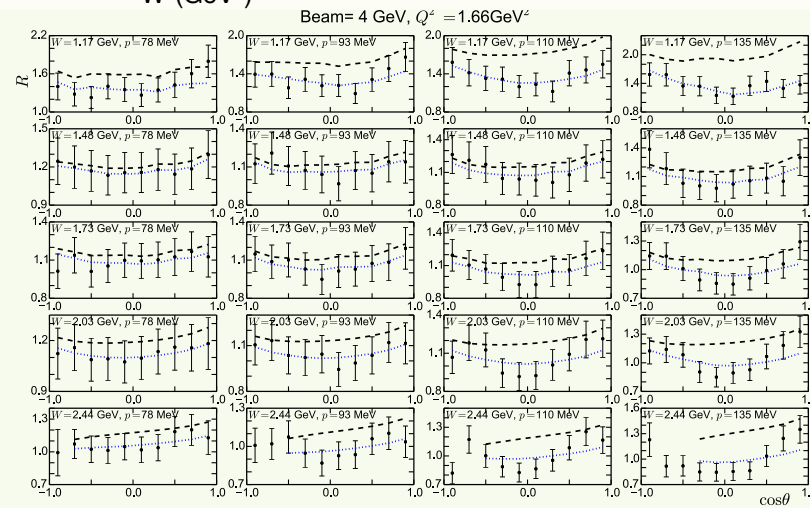
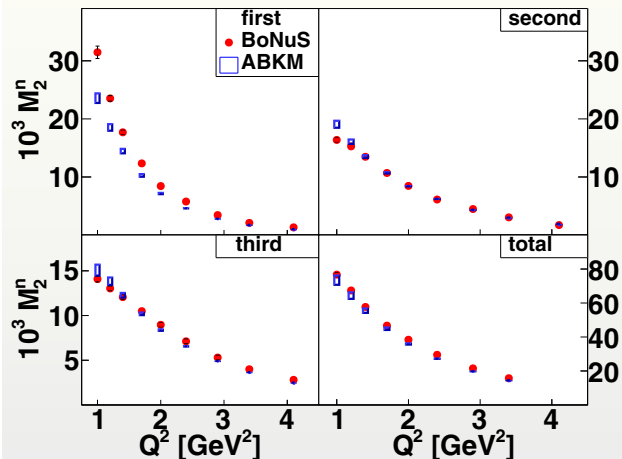
Spectator Tagging – BONuS6 Results



FSI: Cosyn et al.

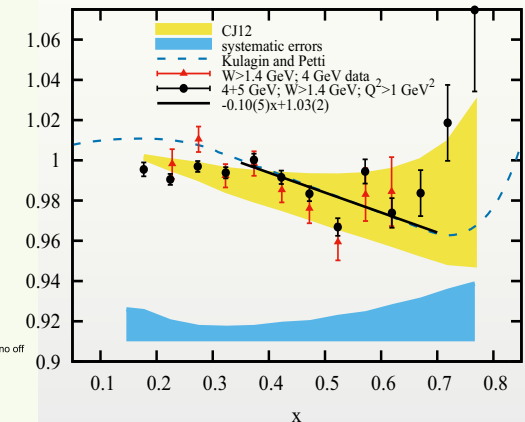
Duality

$W^2(\text{GeV}^2)$



EMC Ratio

PHYSICAL REVIEW C 92, 015211 (2015)

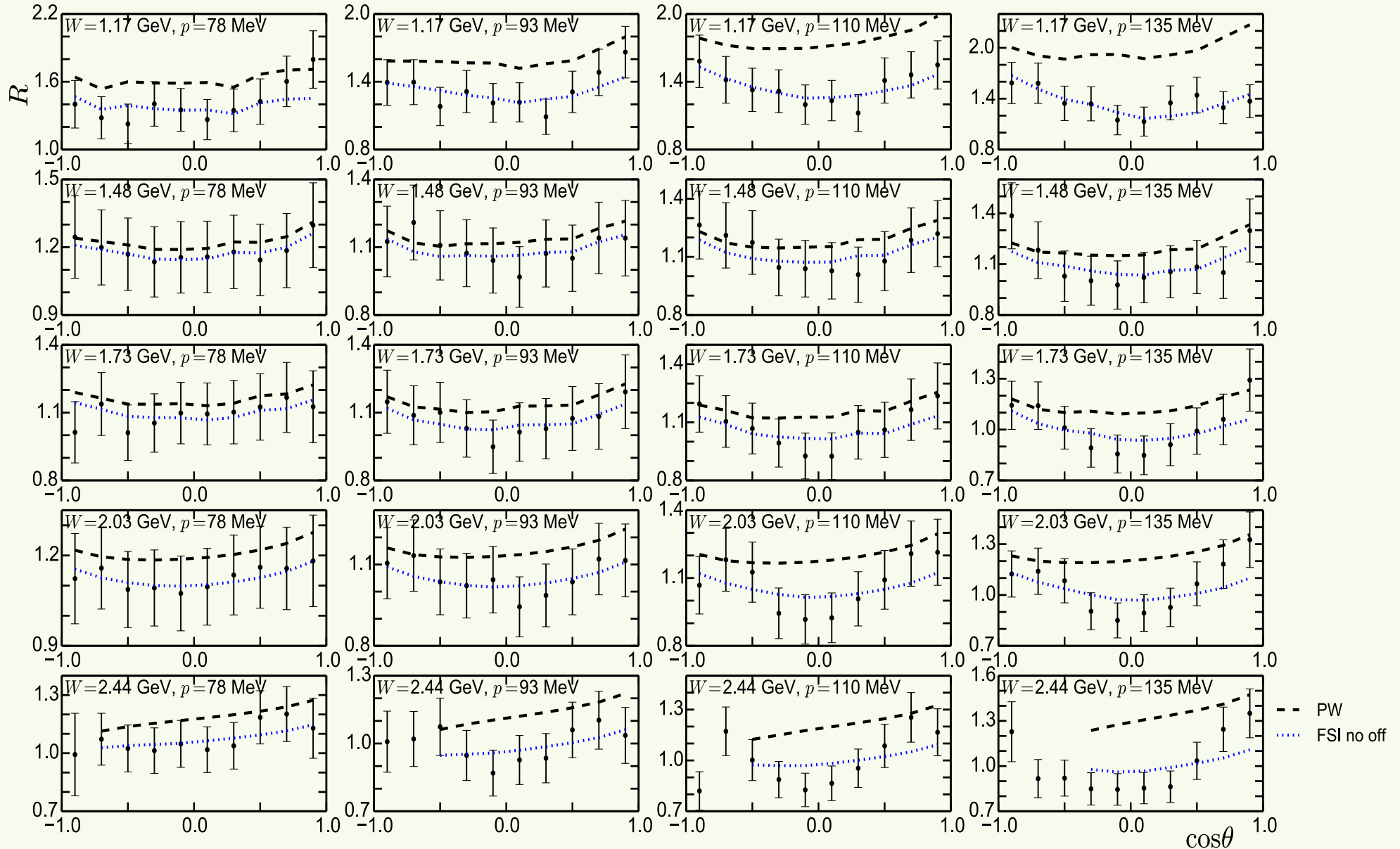


FSI - BONuS Results

FSI: Cosyn et al.

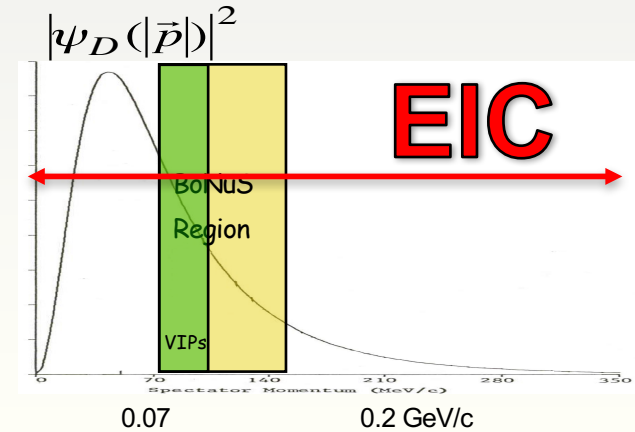
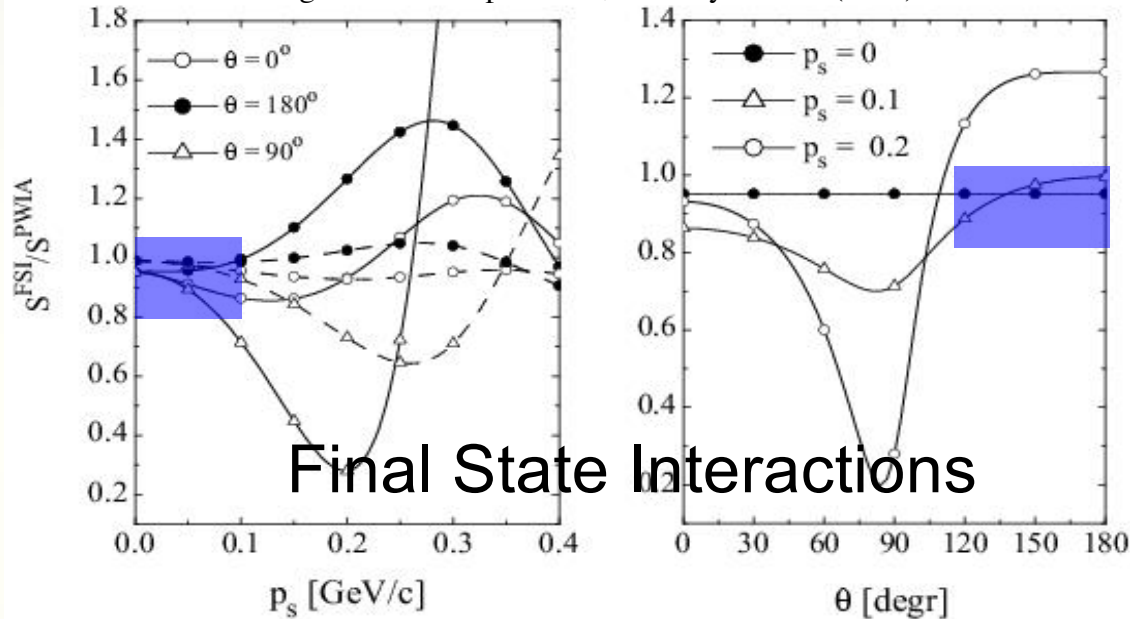
R = ratio of tagged SF
in $d(e,e'p_s)$ to "free" n
SF, vs. momentum and
angle (relative to \mathbf{q}
vector) of spectator p_s

Beam = 4 GeV, $Q^2 = 1.66 \text{ GeV}^2$



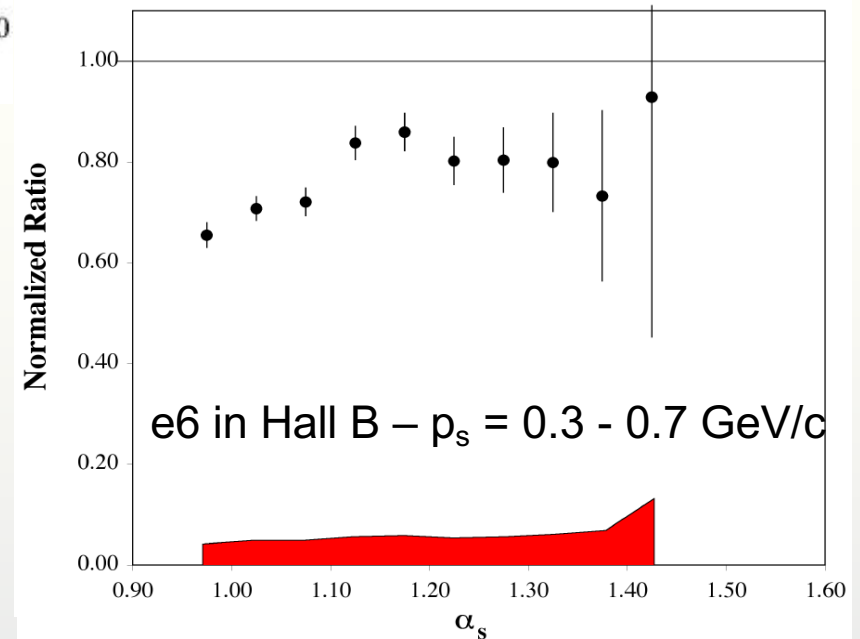
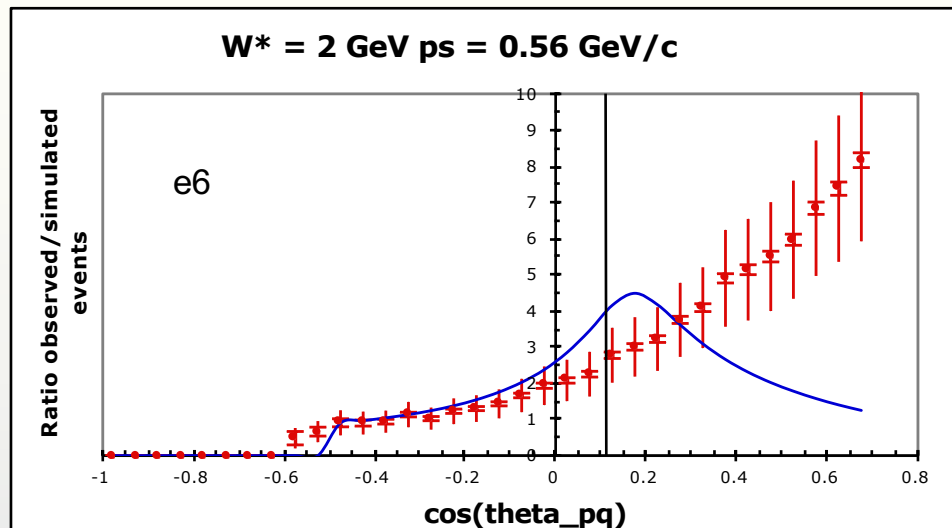
Spectator Tagging – E6 Results

Ciofi degli Atti and Kopeliovich, Eur. Phys. J. A17(2003)133



Finite coverage of WF

$$R_n \equiv F_2^{n(eff)}(W^2, Q^2, p^2) / F_2^n(W^2, Q^2)$$



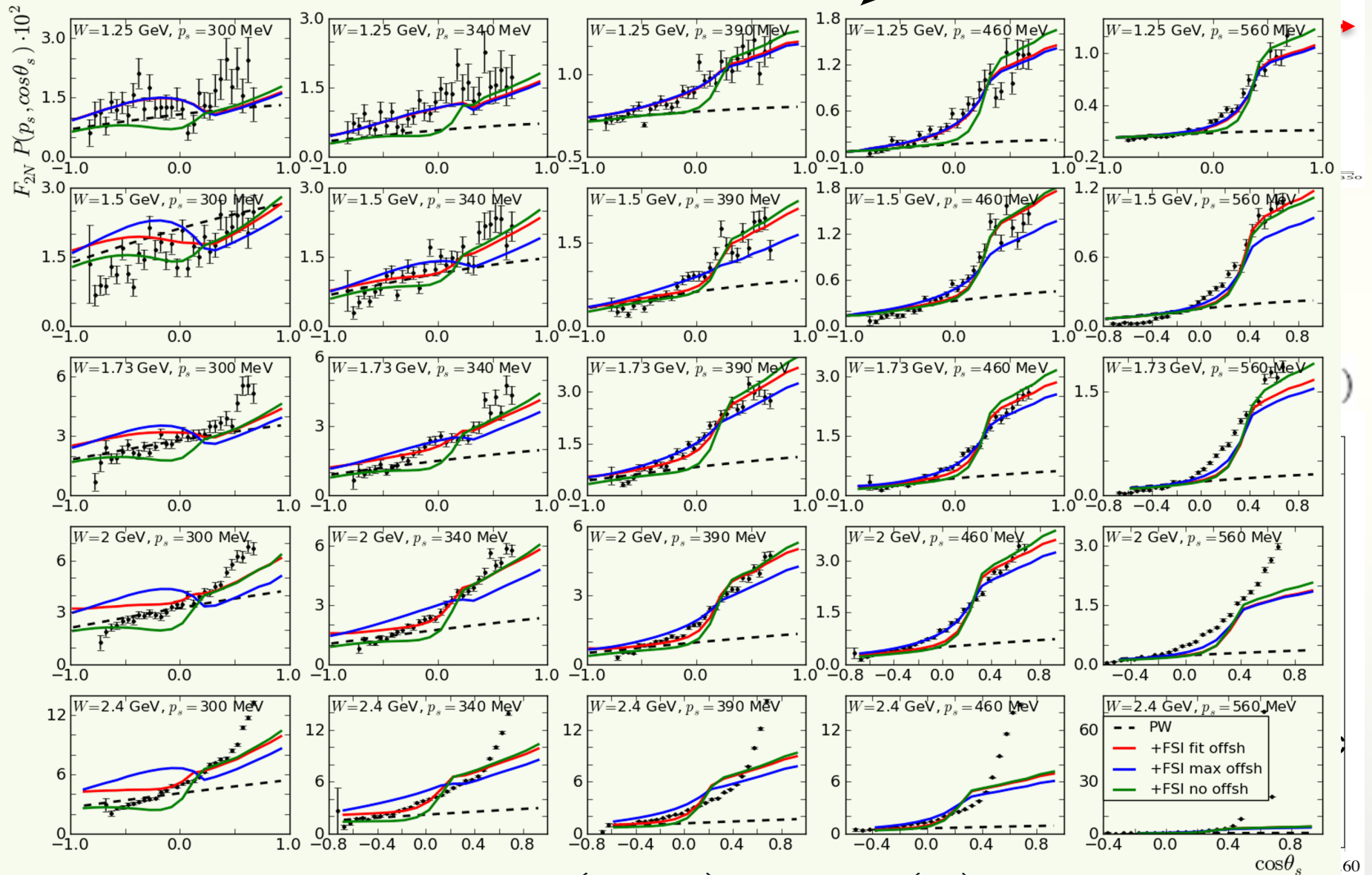
Spectator Tagging – E6 Results

$$b \cdot \vec{p}_s \propto |\vec{p}_s|^2$$

increasing p_s

W. Cosyn et al.

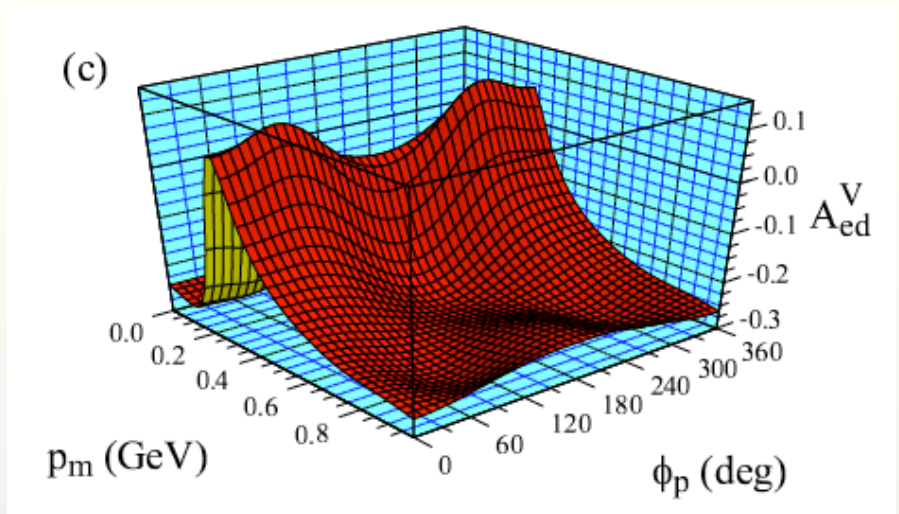
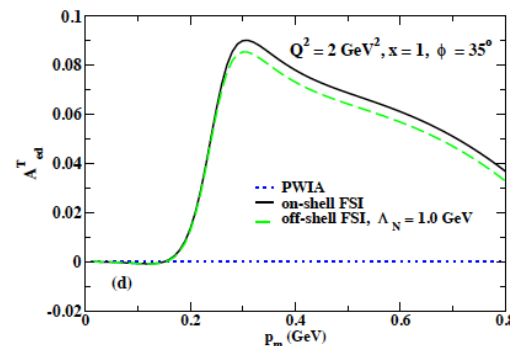
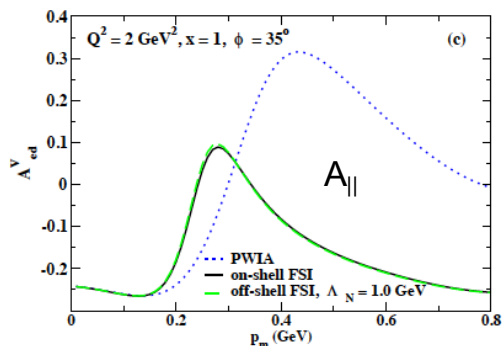
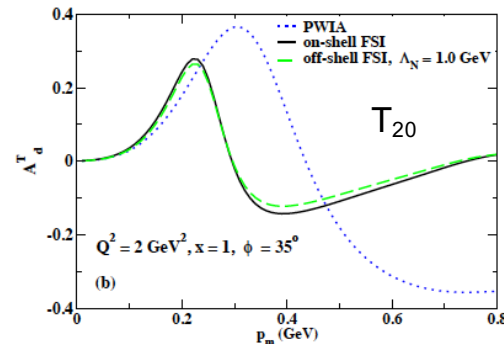
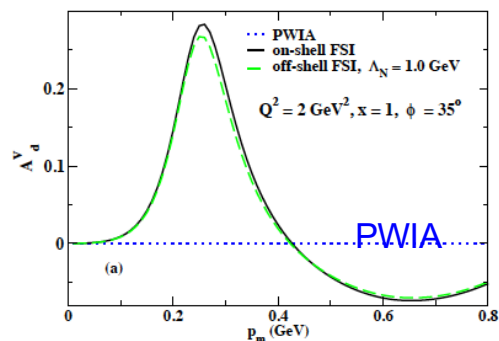
increasing invariant mass of X



data:Klimenko et al. (JLab CLAS), PRC73 035212 ('06)

Testing FSI Models in the quasi-elastic channel

- W. Van Orden and S. Jeschonnek have developed a fully relativistic description of cross sections, vector and tensor asymmetries for $D(e,e'p)n$, including (spin-dependent) FSI (based on known phase shifts)



Results from Jefferson Lab EG1b Experiment in CLAS

PHYSICAL REVIEW C **95**, 024005 (2017)

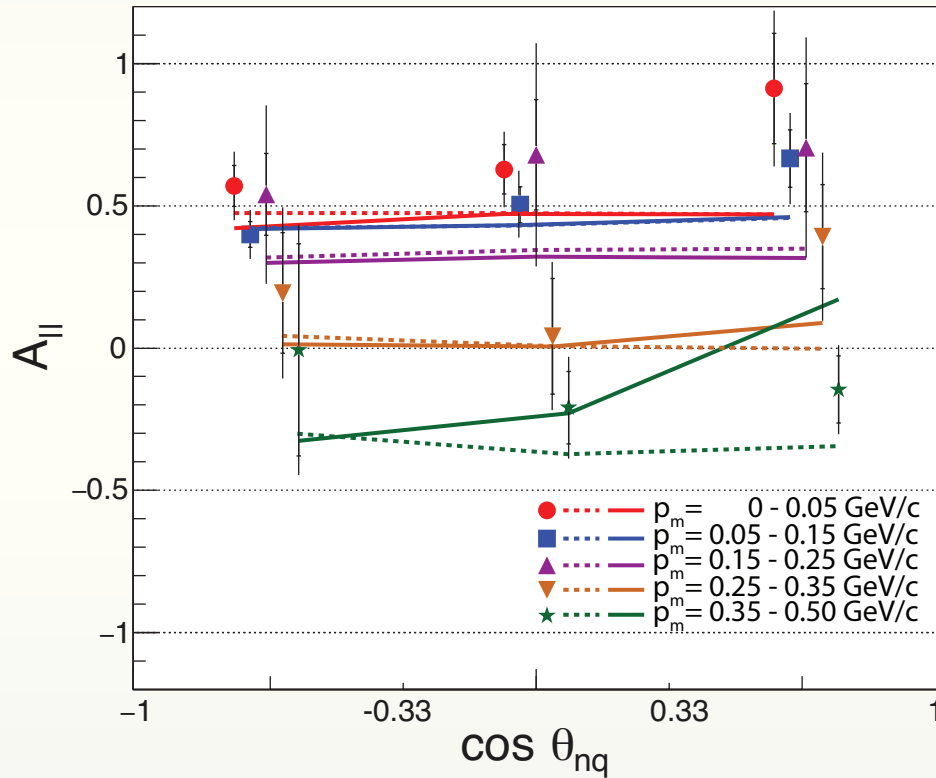


FIG. 11. $A_{||}$ for beam energies of 1.6 – 1.7 GeV and $0.38 \text{ GeV}^2/c^2 \leq Q^2 \leq 0.77 \text{ GeV}^2/c^2$,

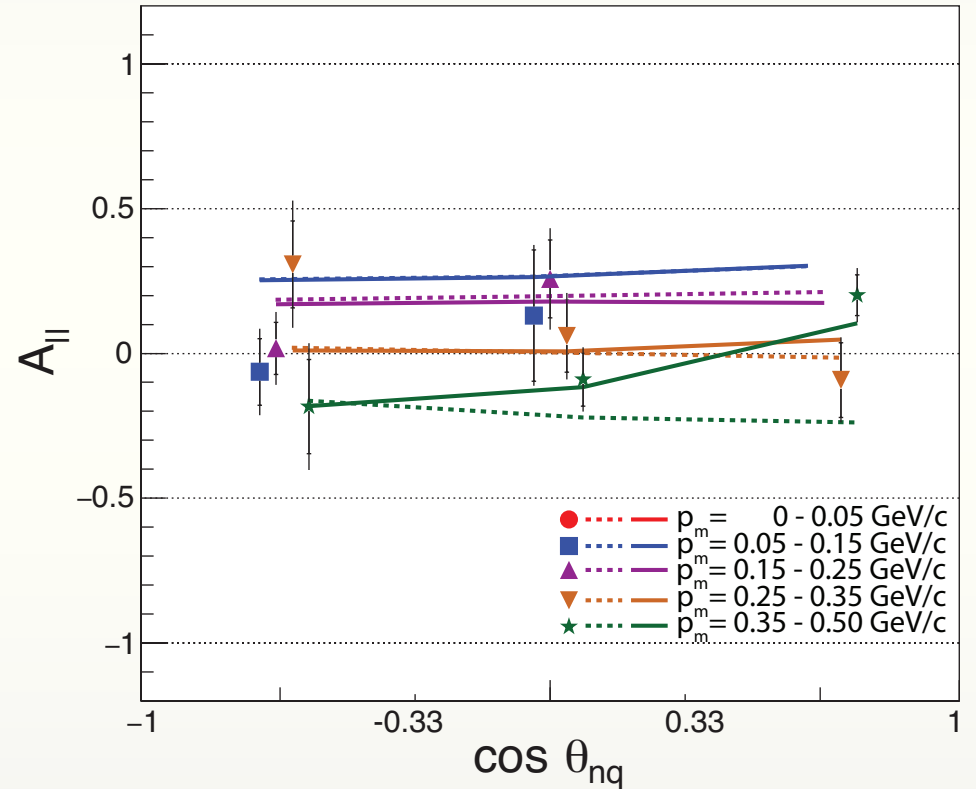
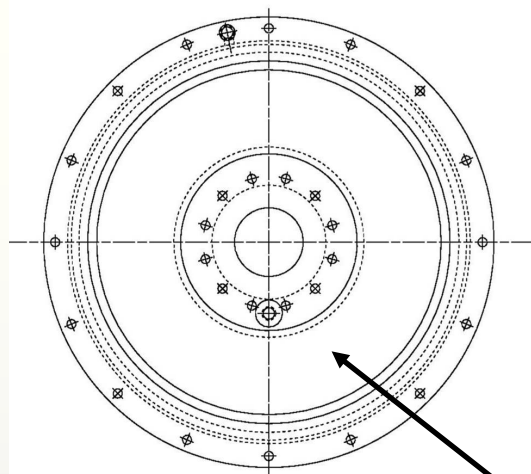
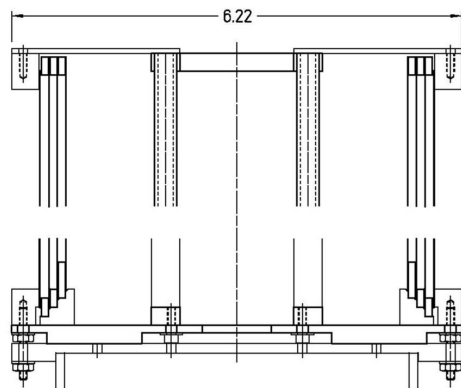
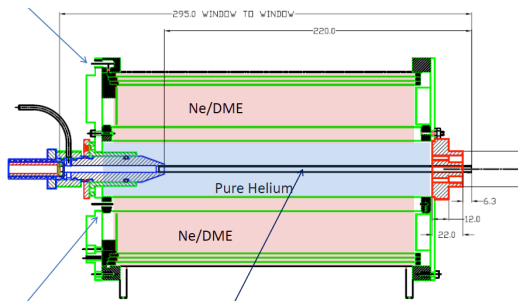


FIG. 12. $A_{||}$ for a beam energy of 2.5 GeV and the same Q^2

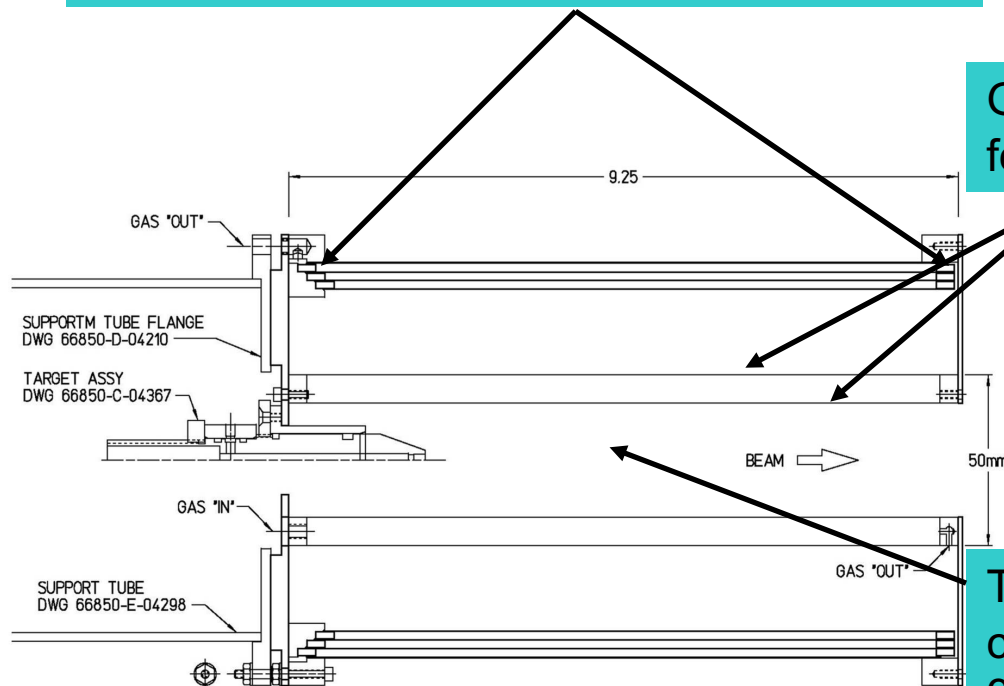
The total value for χ^2 , summed over all bins excluding the 5.x GeV data, is 165.6 for the PWIA model (dof=91, $p < 3 \times 10^{-6}$; or $\chi^2=182.3$ for dof = 103 when we include the 5.x GeV bins) and $\chi^2=121$ (dof = 91, $p \approx 0.02$) for the model with FSI included.

The 2nd RTPC (EG6)



BONUS - UPSTREAM END PLATES

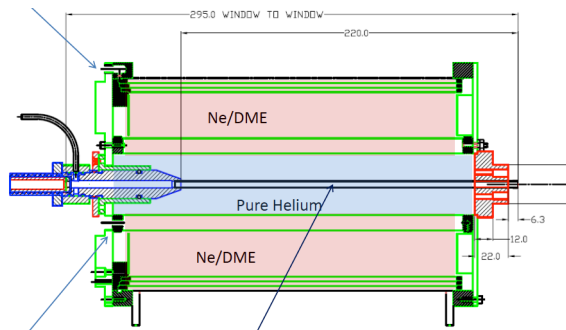
Lightweight (carbon-foam composite structure) rungs for stress-free, self support of GEMs – to avoid “wrinkles”



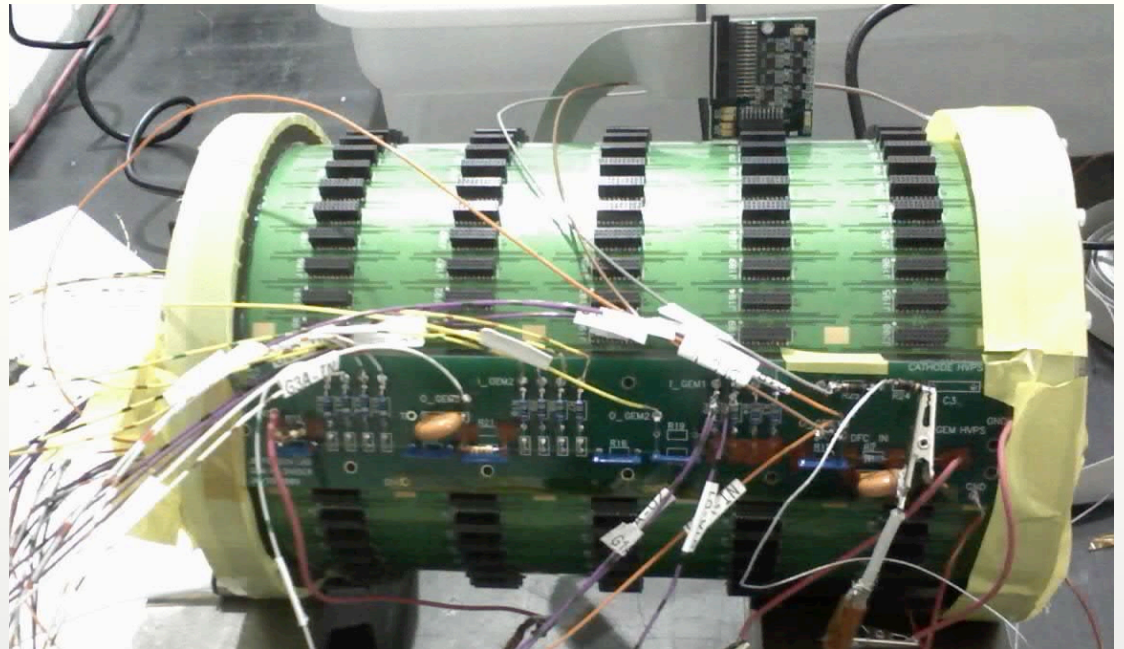
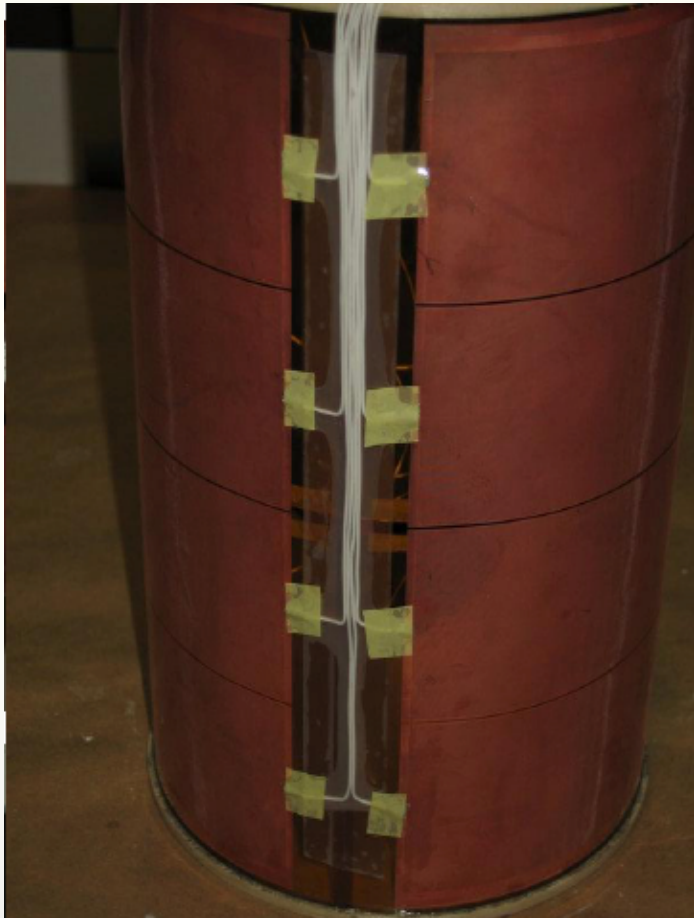
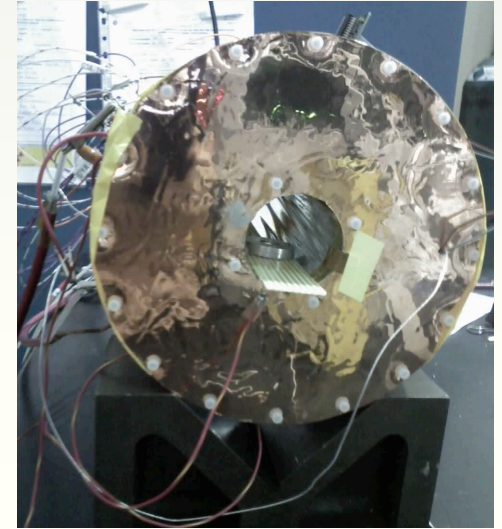
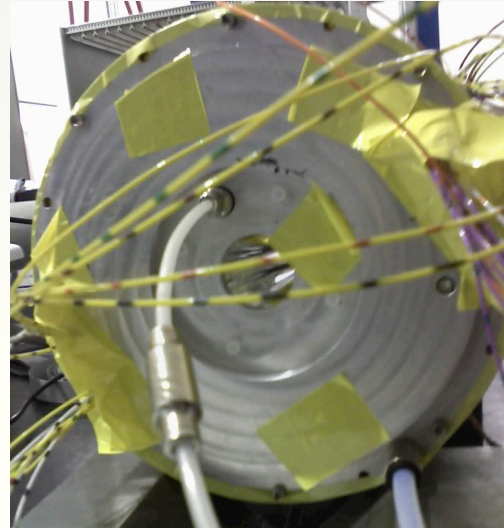
Ground and cathode foils, each 6μm thick

Target region - new cell – ID 4mm, 30μm thick wall

Open, 2π geometry – only 80% were accessible due to the GEM and readout pad sizes



The 2nd RTPC (ii)



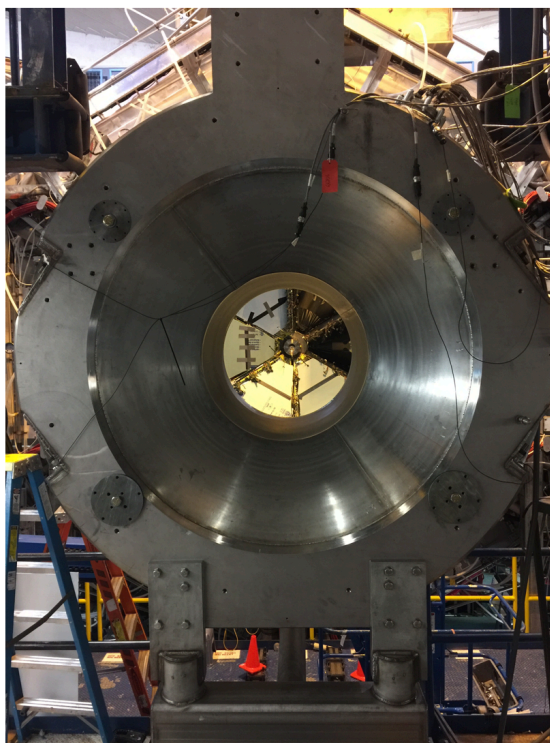
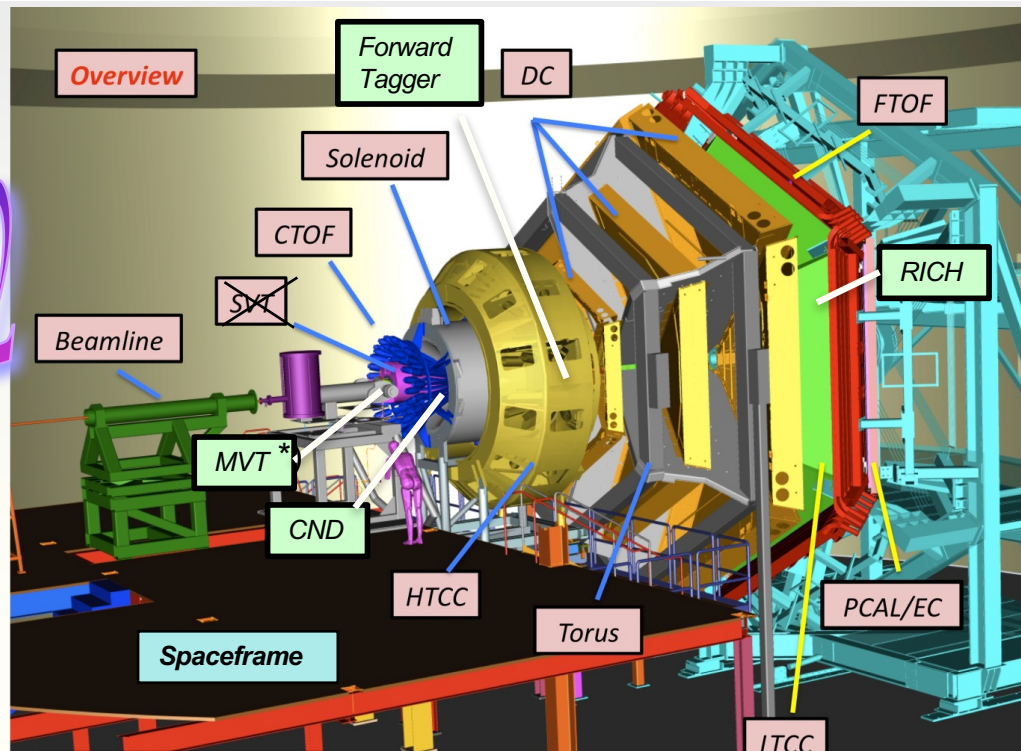
BONuS12 at 11 GeV

BoNuS12

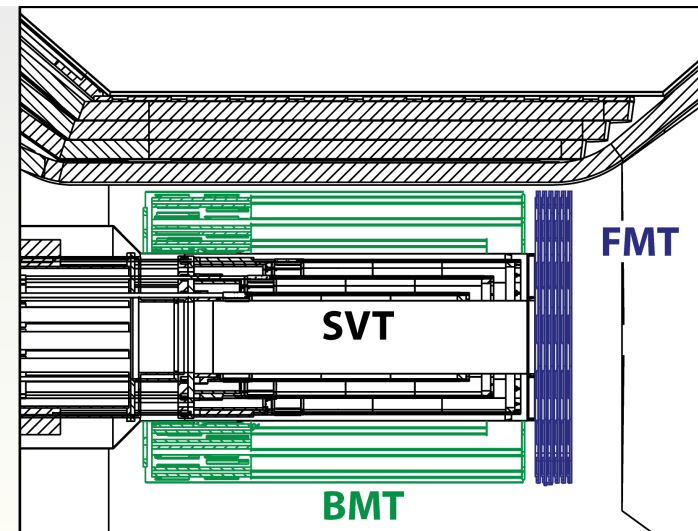
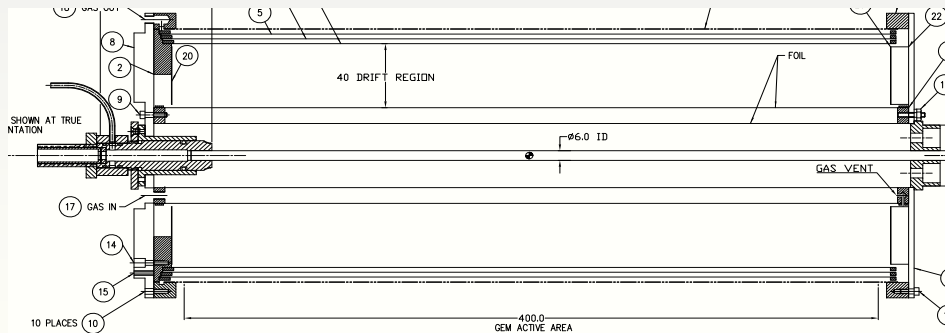
E12-06-113

CLAS12

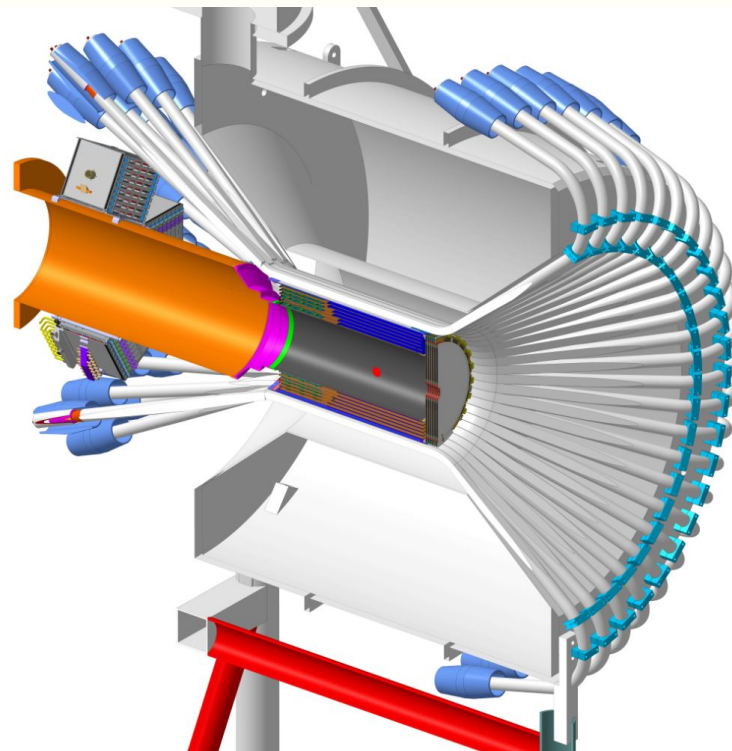
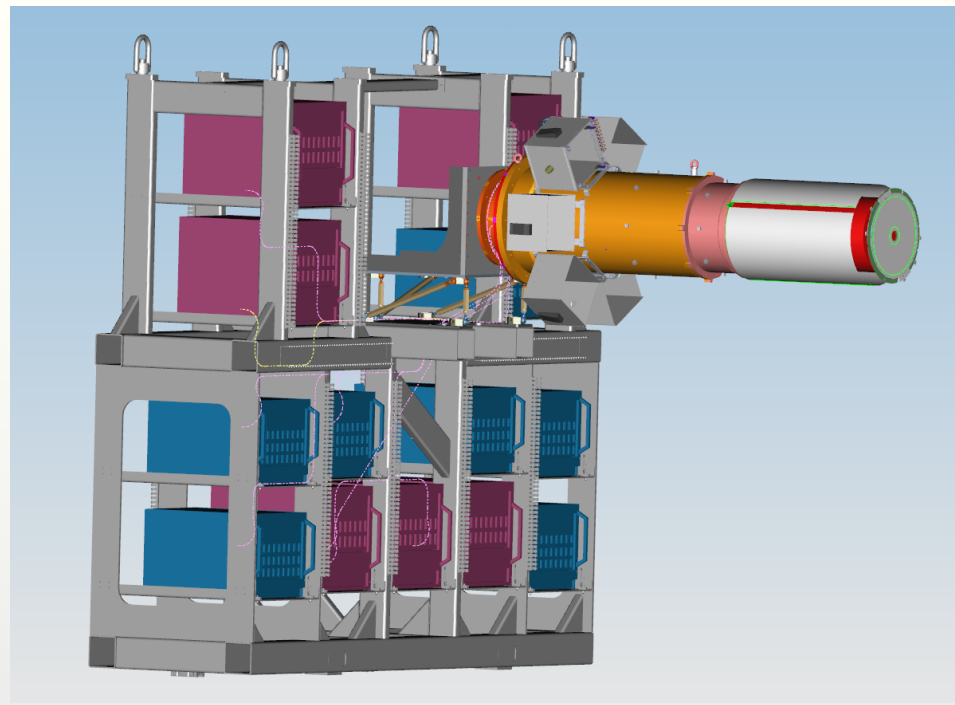
- Data taking for 35 days on D_2 ,
4 days on H_2 + 1 day aux.
with $\mathcal{L} = 2 \cdot 10^{34}$ nuclei/cm² s⁻¹
plus 2 days commissioning at 2.2 GeV
- **NEW** RTPC detector, DAQ
- DIS region with
 - $Q^2 > 1 \text{ GeV}^2/c^2$
 - $W^* > 2 \text{ GeV}$
 - $p_s > 70 \text{ MeV}/c$
 - $10^\circ < \theta_{pq} < 170^\circ$



BONuS12 at 11 GeV

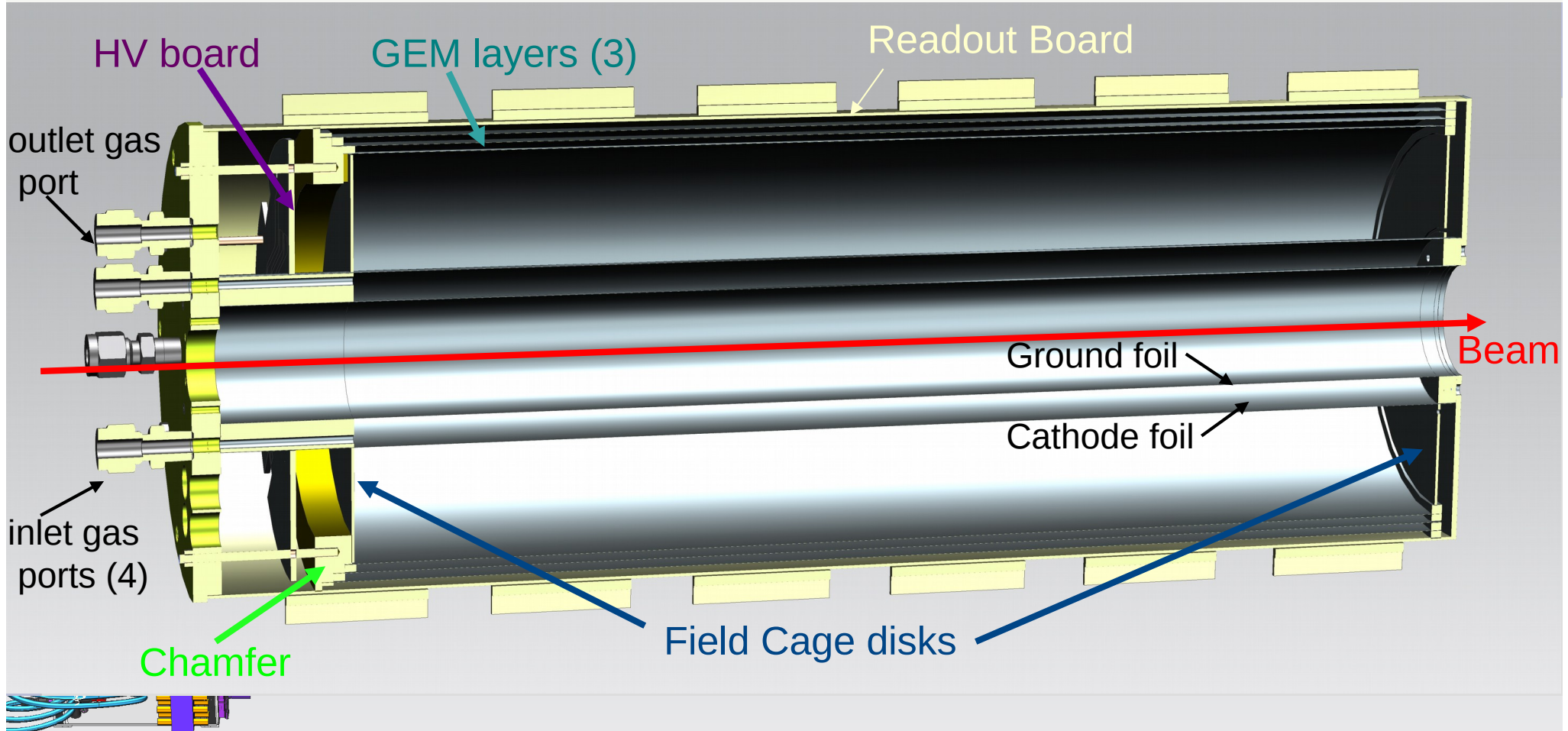


Central
Detector



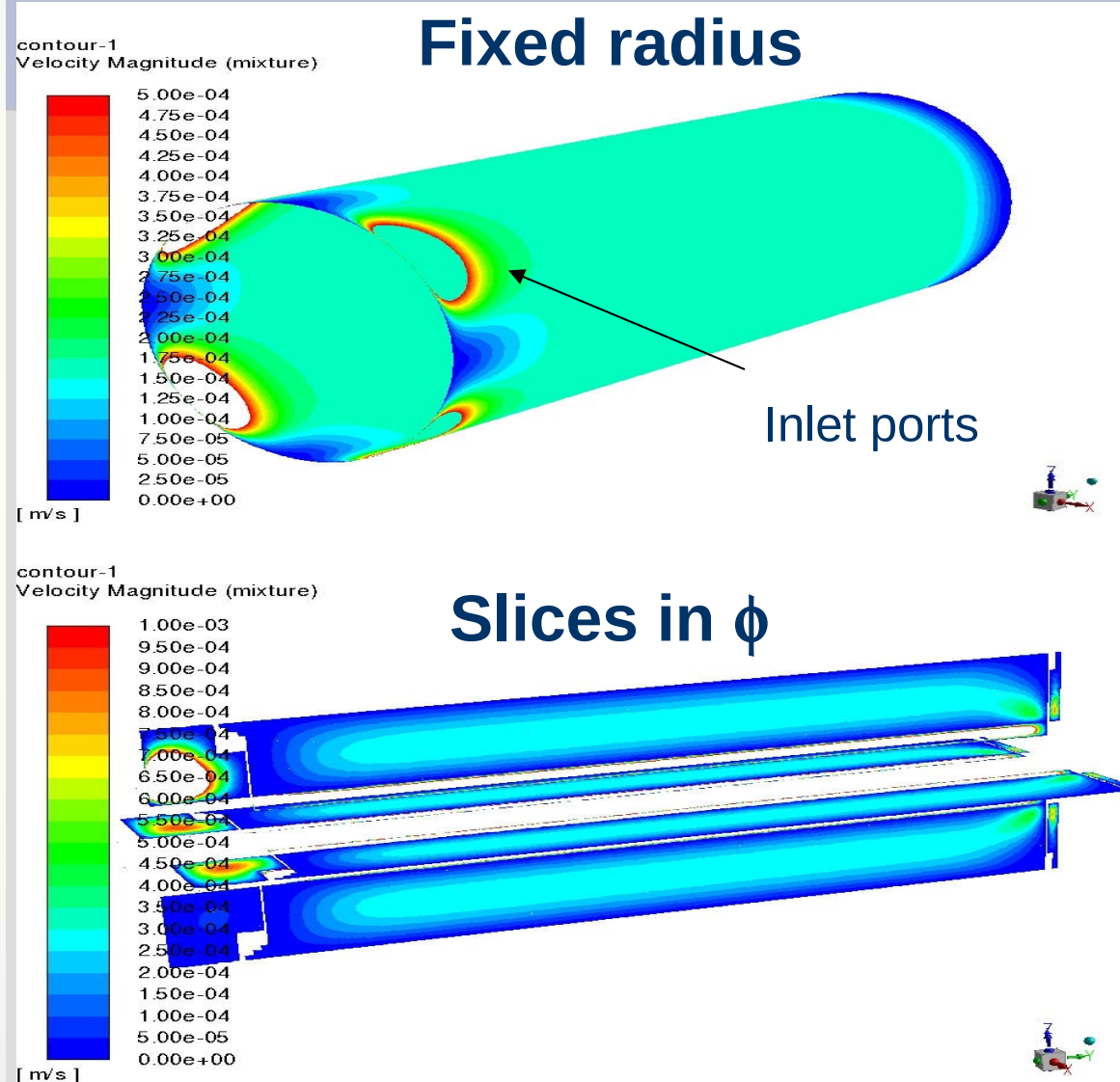
BONuS12 RTPC
replaces SiVtxT +
Barrel μ egas
(but forward Vtx
tracking needed!)

Completed RTPC w/ target



Velocity profiles from CFD

Silviu Covrig



Assumptions:

→ 0.2 L / min

→ Premixed

Conclusions:

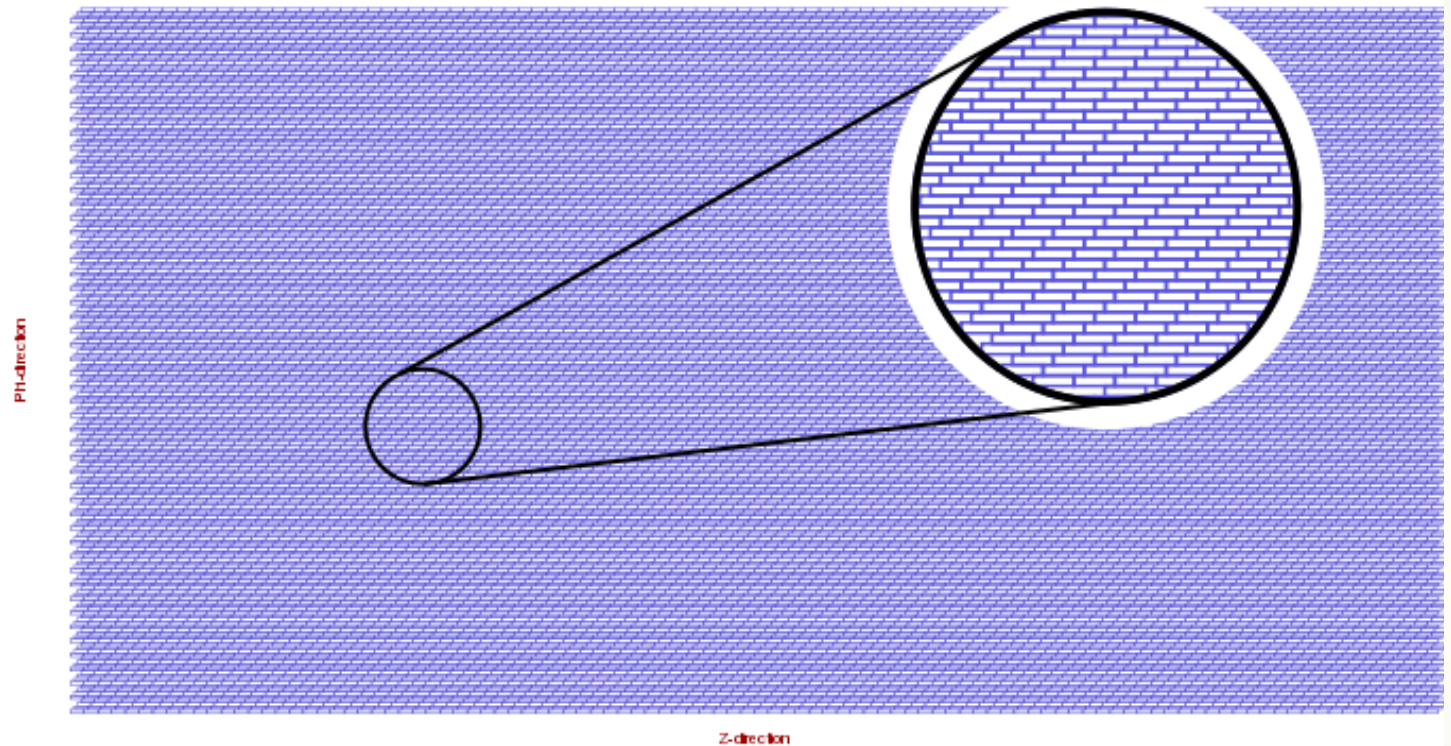
→ Current design
provides relatively
uniform flow

Basic Assembly of GEMs

- GEM layers will be constructed using similar process to KLOE2 / EG6
- Each GEM layer is wrapped on a cylindrical mandrel
- Epoxy inner (downstream) and outer (upstream) rings



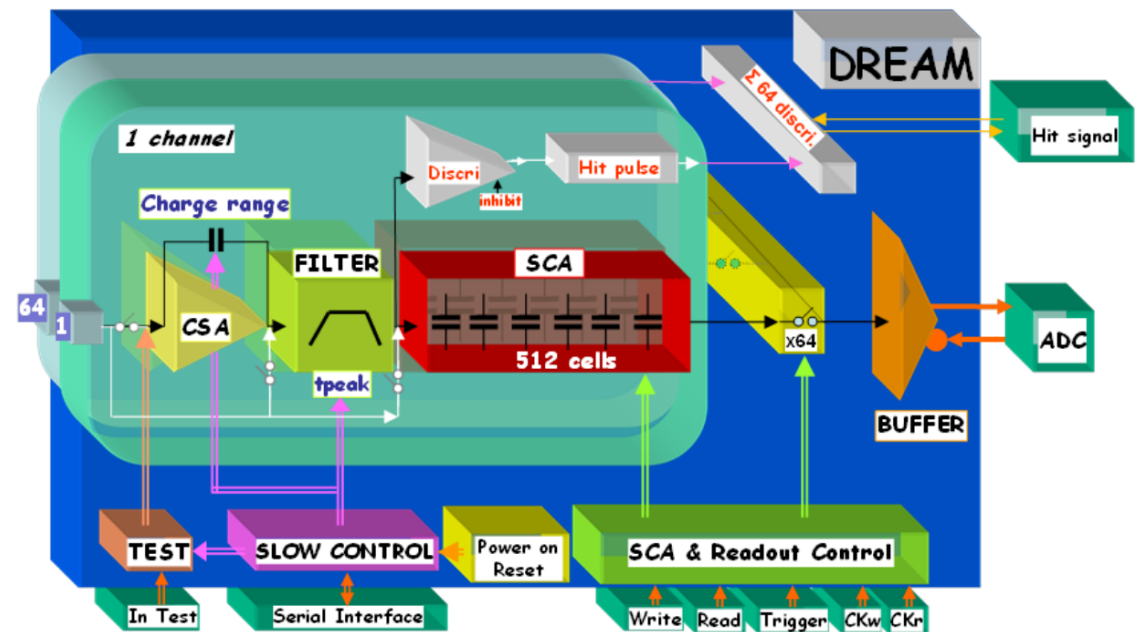
View of the readout board



180 rows, 96 columns for a total of 17280 pads
Each pad covers 4 mm in z and 2 degrees in phi
Rows shifted by $\frac{1}{4}$ pad size in z from one another
Read out by Micromegas DREAM electronics

DREAM electronics developed for the Micromegas of CLAS12

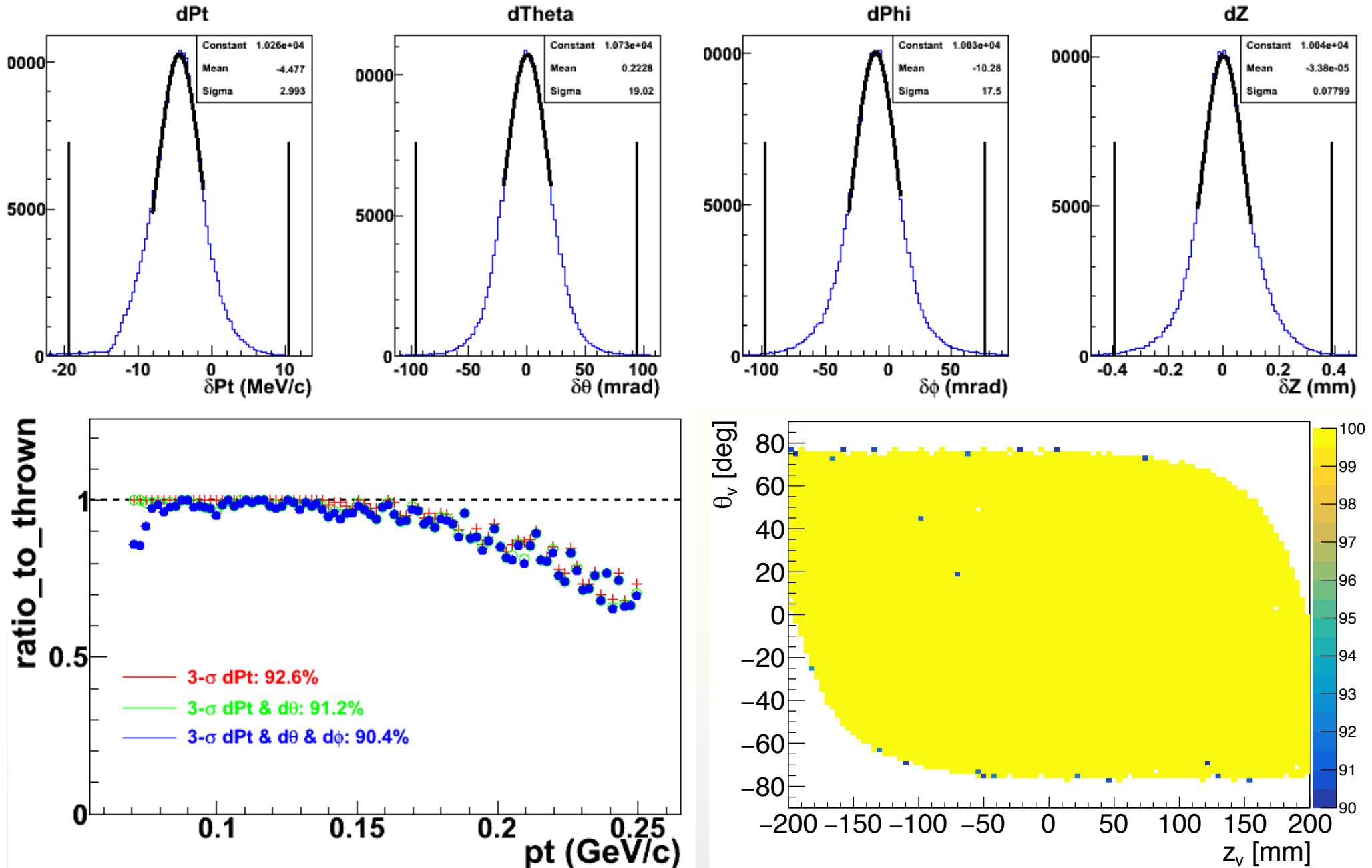
- 512 memory cells/channel
- read out selected cells after trigger
- Low noise
- Analogue multiplexed output
- Latency up to 16 μ s



BONuS12 will use DREAM electronics

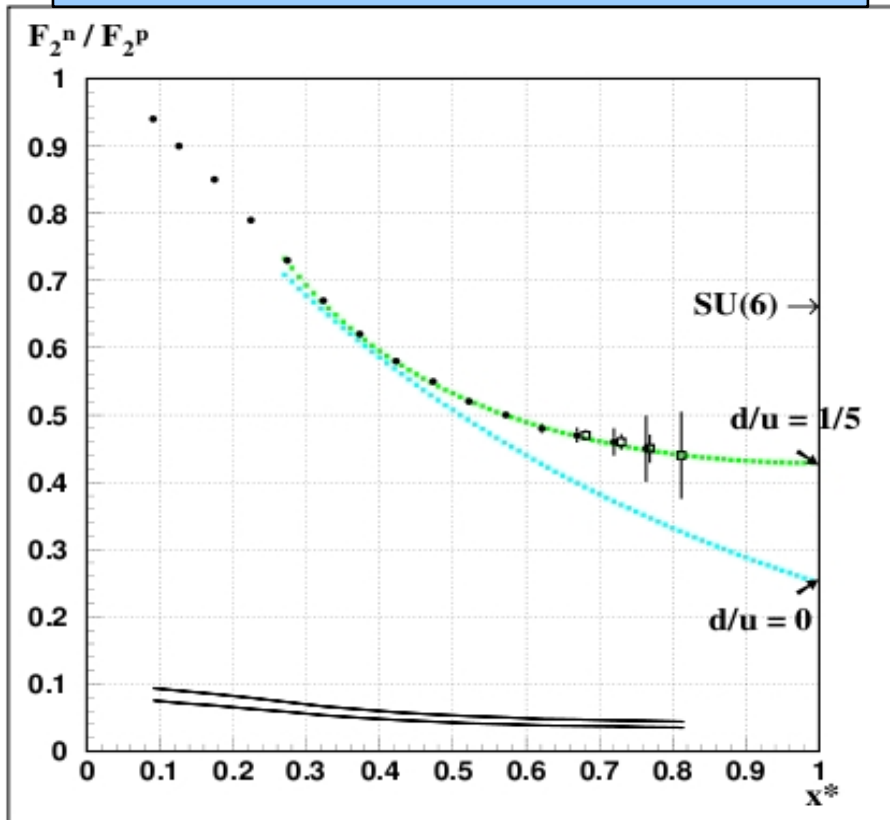
- Already available on site
- Fits BONuS12 needs
- Contract signed with Saclay (test bench + manpower)
- Need to update the firmware
- Test bench working at Old Dominion University
- BONuS12 will use available FEU and signal cables from barrel Micromegas
- Adaptation board to protect the electronics from over current

Simulation

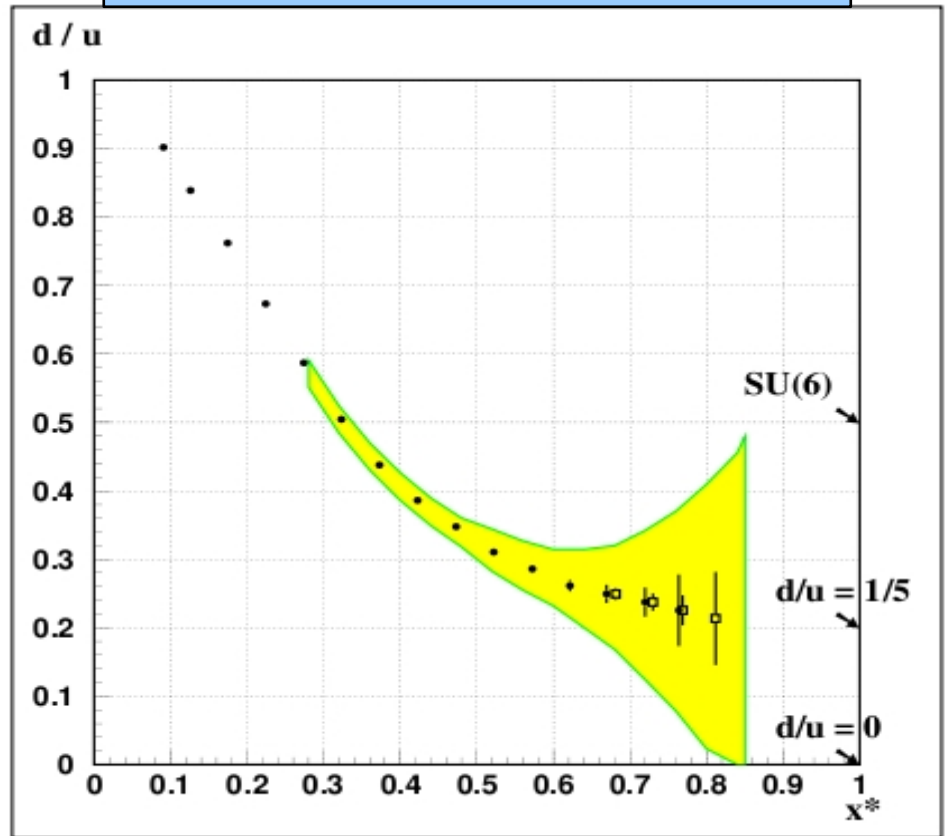


Expected Results

Neutron/Proton structure function



d/u

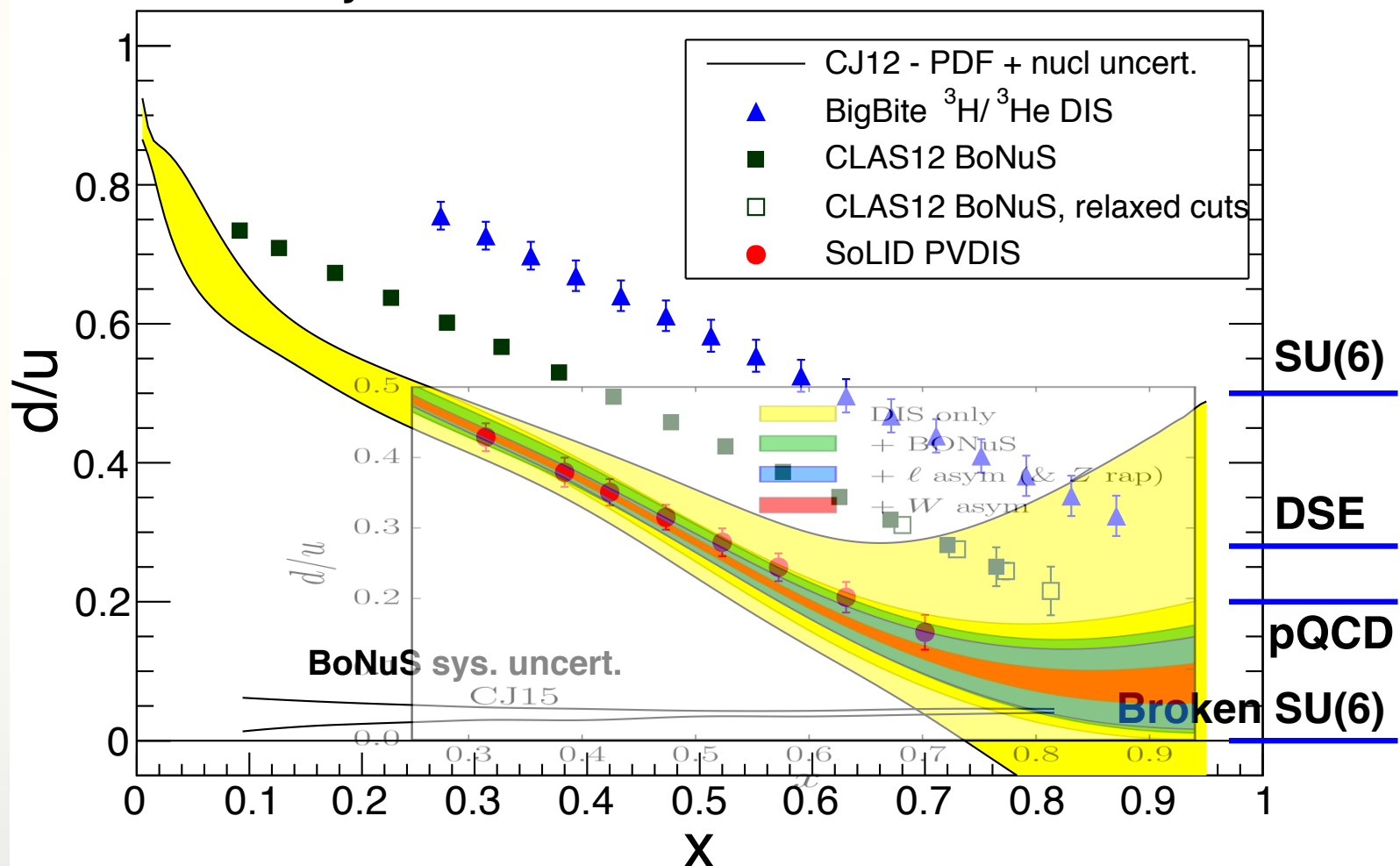


Dark Symbols: $W^* > 2$ GeV (x^* up to 0.8, bin centered $x^* = 0.76$)

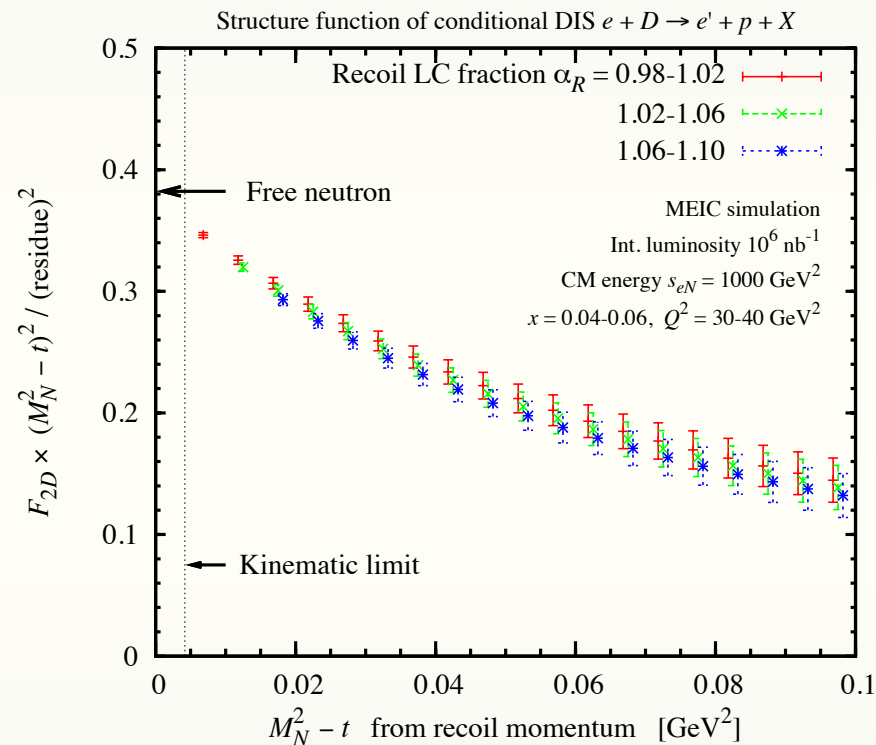
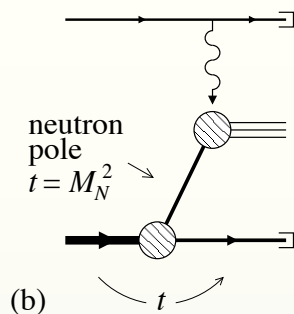
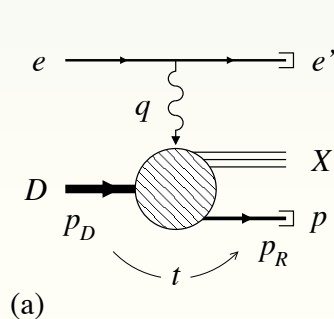
Open Symbols: "Relaxed cut" $W^* > 1.8$ GeV (x^* up to 0.83)

The future: JLab at 11 GeV

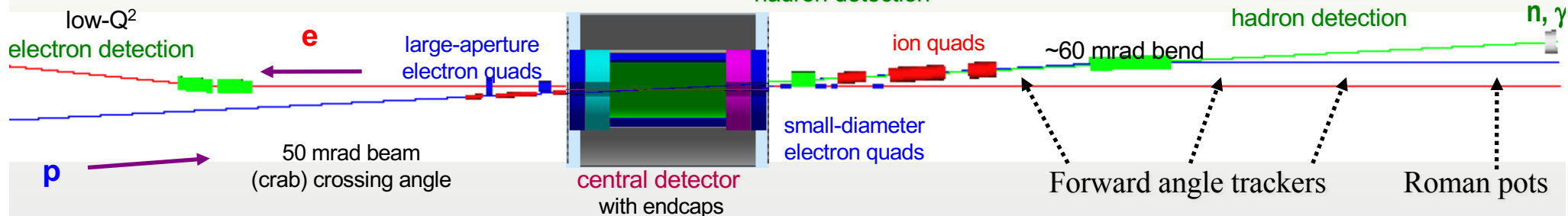
Projected 12 GeV d/u Extractions



The more distant future: EIC



(from GEANT4)

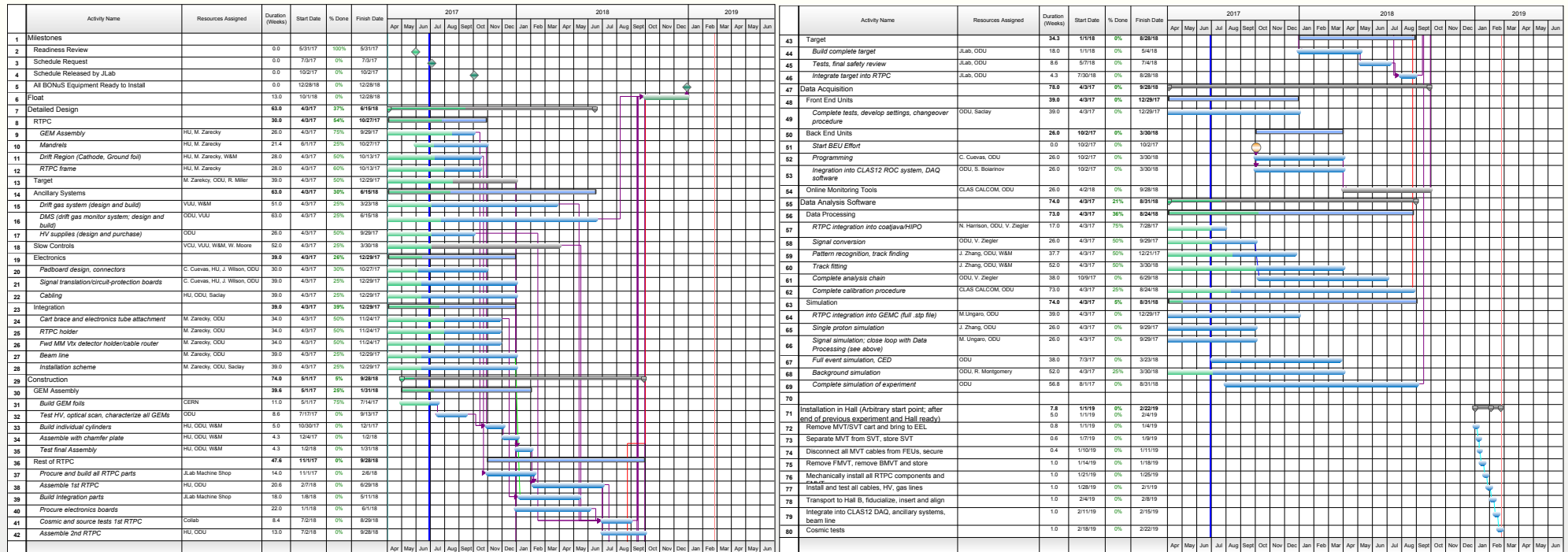


What else can we do with RG F?

- EMC effect in D
 - tag on slow/fast p; bench mark for heavier nuclei? FSI, d WF,...
- “LAND” experiment E12-11-003A
 - tag p structure in d with backward n (Approved Run Group proposal, PAC43; Or Hen, L. Weinstein, E. Piasetzky, H. Hakobyan)
- nDVCS?
 - At least calibrate CND and get first sample without nuclear distortions; fully exclusive!
- n Form Factors? Alternative method to cross check
- n resonances? (Use forward tagger)
- n SIDIS? (Flavor tagging, TMDs)
- phi-N bound state: $\text{Au}(\gamma^*, pK^+K^-)$

Summary

- BONuS12 will determine d/u out to $x = 0.8$ -> high priority
- Present goal: 3rd experiment after LH₂ and LD₂ target (Run Group F, Winter 2019? Passed initial ERR)
- Plethora of potential additional Physics topics
- Detector/target system design at advanced stage



Conclusion

- Few-body nuclei (D and ^3He) continue to be “neutron targets of choice”
 - Interpretation of results complicated by off-shell effects, possible structure modifications and final state interaction...
 - ...but we can also learn a lot about NN interaction and few-body nuclear structure by studying these effects
 - New, more precise theoretical calculations are becoming available and can be tested experimentally
 - Spectator tagging allows us to minimize binding effects or study them in detail
- BONuS12 will extract neutron valence quark distributions
 - Lots more experiments at 12 GeV! Tag polarized SFs?
 - Master of spectator tagging: EIC