

# Unraveling confinement forces with DVCS

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USC Columbia

# Outline

- 1 Theoretical Context
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- 3 Future Measurements
- 4 Projected Impact on GPD Extractions
- 5 Summary and Outlook

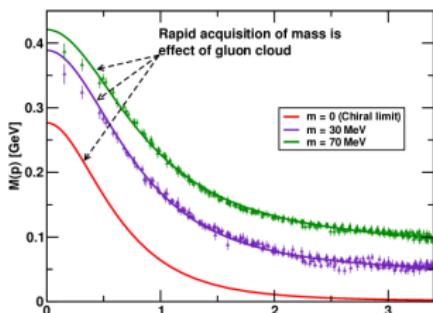


# Theoretical Context



# Confinement Mechanism(s?)

Hadrons are singlets under  $SU(3)_{\text{color}}$ : No net color charge in asymptotic particle states

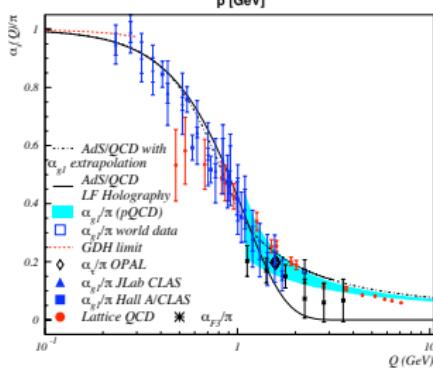


- Linear growth of the static quark-antiquark pair  
Area-law falloff for the Wilson loop
- Gribov Confinement for light quarks  
Analytical properties of the propagators in the infrared  
Instability of the vacuum above a supercritical charge

$$\alpha_{\text{QED}}^{\text{crit}} = 137 \text{ for a point-like nucleus}$$

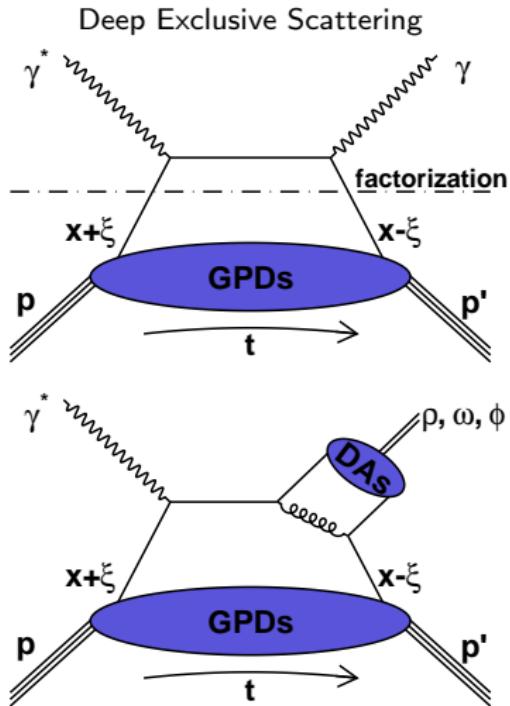
$$\approx 180 \text{ for a finite size nucleus}$$

$$\frac{\alpha_{\text{QCD}}^{\text{crit}}}{\pi} = C_F^{-1} \left[ 1 - \sqrt{\frac{2}{3}} \right] \approx 0.137$$



- Light-Front AdS/QCD  
quark and gluon chiral condensates confined!  
→condensates contribution to the cosmological constant already included in hadron mass
- Mass-Gap Millenium problem and Yang-Mills existence  
\$1M from the Clay Mathematical Institute

# Generalized Parton Distributions



$$\gamma^* p \rightarrow \gamma p', \rho p', \omega p', \phi p'$$

Bjorken regime :

$$Q^2 \rightarrow \infty, x_B \text{ fixed}$$

$$t \text{ fixed} \ll Q^2, \xi \rightarrow \frac{x_B}{2-x_B}$$

$$\begin{aligned} & \frac{P^+}{2\pi} \int dy^- e^{ixP^+y^-} \langle p' | \bar{\psi}_q(0) \gamma^+ (1 + \gamma^5) \psi(y) | p \rangle \\ &= \bar{N}(p') \left[ H^q(x, \xi, t) \gamma^+ + E^q(x, \xi, t) i \sigma^{+\nu} \frac{\Delta_\nu}{2M} \right. \\ & \quad \left. + \tilde{H}^q(x, \xi, t) \gamma^+ \gamma^5 + \tilde{E}^q(x, \xi, t) \gamma^5 \frac{\Delta^+}{2M} \right] N(p) \end{aligned}$$

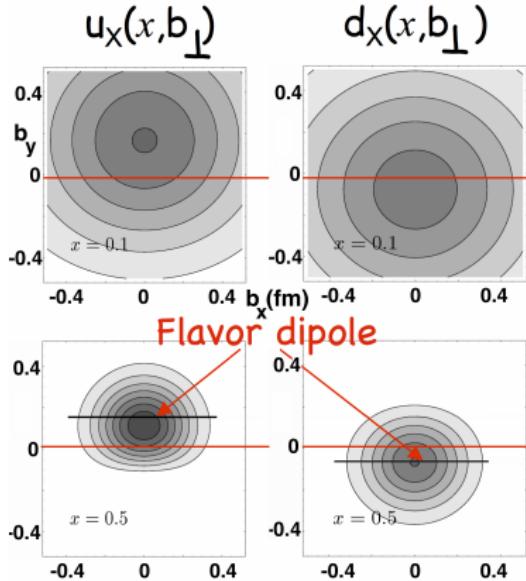
spin	N no flip	N flip
q no flip	$H$	$E$
q flip	$\tilde{H}$	$\tilde{E}$

3-D Imaging conjointly in transverse impact parameter **and** longitudinal momentum

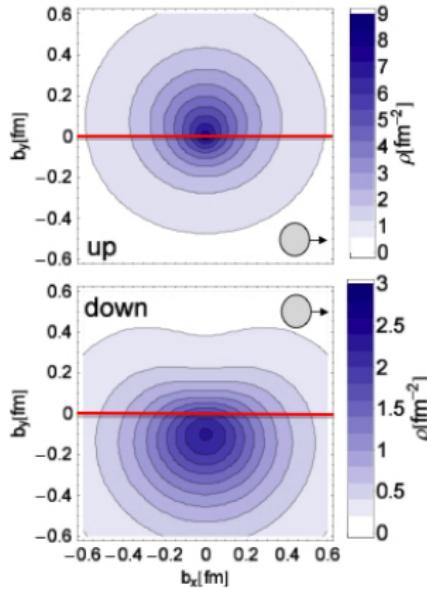
# GPDs and Transverse Imaging ( $x_B, t$ )

$$q_X(x, \vec{b}_\perp) = \int \frac{d^2 \vec{\Delta}_\perp}{(2\pi)^2} \left[ H(x, 0, t) - \frac{E(x, 0, t)}{2M} \frac{\partial}{\partial b_y} \right] e^{-i \vec{\Delta}_\perp \cdot \vec{b}_\perp}$$

Target polarization



Lattice calculation



# Energy Momentum Tensor ( $x, \xi$ )

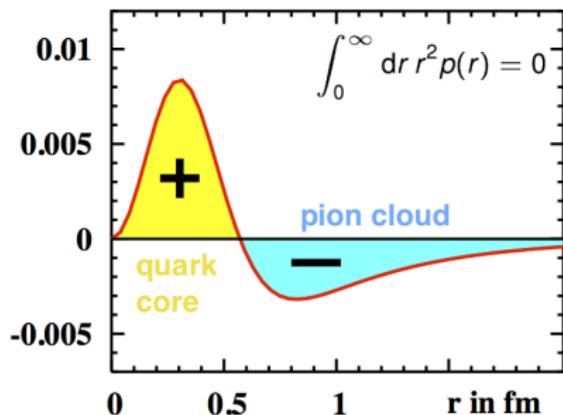
Form Factors accessed via second x-moments :

$$\langle p' | \hat{T}_{\mu\nu}^q | p \rangle = \bar{N}(p') \left[ M_2^q(t) \frac{P_\mu P_\nu}{M} + J^q(t) \frac{i(P_\mu \sigma_{\nu\rho} + P_\nu \sigma_{\mu\rho}) \Delta^\rho}{2M} + d_1^q(t) \frac{\Delta_\mu \Delta_\nu - g_{\mu\nu} \Delta^2}{5M} \right] N(p)$$

Angular momentum distribution

$$J^q(t) = \frac{1}{2} \int_{-1}^1 dx x [H^q(x, \xi, t) + E^q(x, \xi, t)]$$

Distribution of pressure  
 $r^2 p(r)$  in  $\text{GeV fm}^{-1}$



Mass and force/pressure distributions

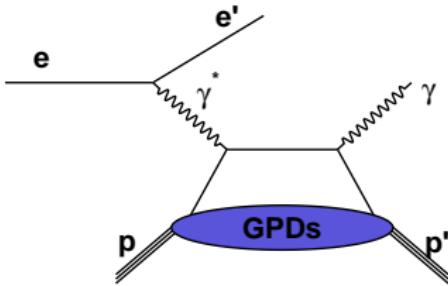
$$M_2^q(t) + \frac{4}{5} d_1(t) \xi^2 = \frac{1}{2} \int_{-1}^1 dx x H^q(x, \xi, t)$$

$$d_1(t) = 15M \int d^3 \vec{r} \frac{j_0(r\sqrt{-t})}{2t} p(r)$$



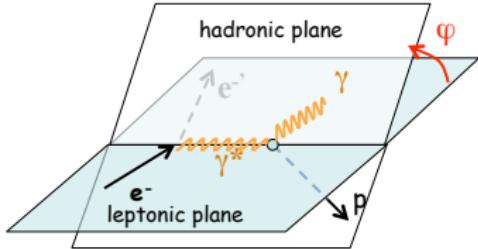
# Deeply Virtual Compton Scattering

The cleanest GPD probe at low and medium energies



$$ep \rightarrow e p' \gamma$$

$$\sigma(ep \rightarrow e p' \gamma) \propto \left| \begin{array}{c} \text{DVCS} \\ + \\ \text{(a)} \\ + \\ \text{BH} \\ + \\ \text{(b)} \\ + \\ \text{(c)} \end{array} \right|^2$$



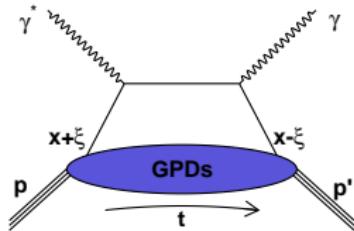
$$A_{LU} = \frac{d^4\sigma^\rightarrow - d^4\sigma^\leftarrow}{d^4\sigma^\rightarrow + d^4\sigma^\leftarrow} \stackrel{\text{twist-2}}{\approx} \frac{\alpha \sin \phi}{1 + \beta \cos \phi}$$

$$\alpha \propto \text{Im} \left( F_1 \mathcal{H} + \xi G_M \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E} \right)$$

$$\mathcal{H}(\xi, t) = i\pi H(\xi, \xi, t) + \mathcal{P} \int_{-1}^1 dx \frac{H(x, \xi, t)}{x - \xi}$$

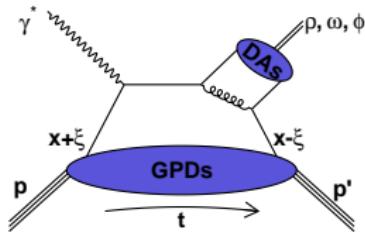
$$A_{UL} \propto \text{Im} \left( F_1 \tilde{\mathcal{H}} + \xi G_M \mathcal{H} + G_M \frac{\xi}{1 + \xi} \mathcal{E} + \dots \right) \sin \phi$$

# Observables sensitivities to GPD



DVCS

	$\mathcal{I}m$	$\mathcal{R}e$
$\mathcal{H}$	$A_{LU}$	$\sigma$
$\tilde{\mathcal{H}}$	$A_{UL}$	$A_{LL}, A_{LT}$
$\mathcal{E}$	$A_{UT}$	



DVMP

	Meson	Flavor
$\mathcal{H}, \mathcal{E}$	$\rho^+$	$u - d$
	$\rho^0$	$2u + d$
	$\omega$	$2u - d$
	$\phi$	$g$

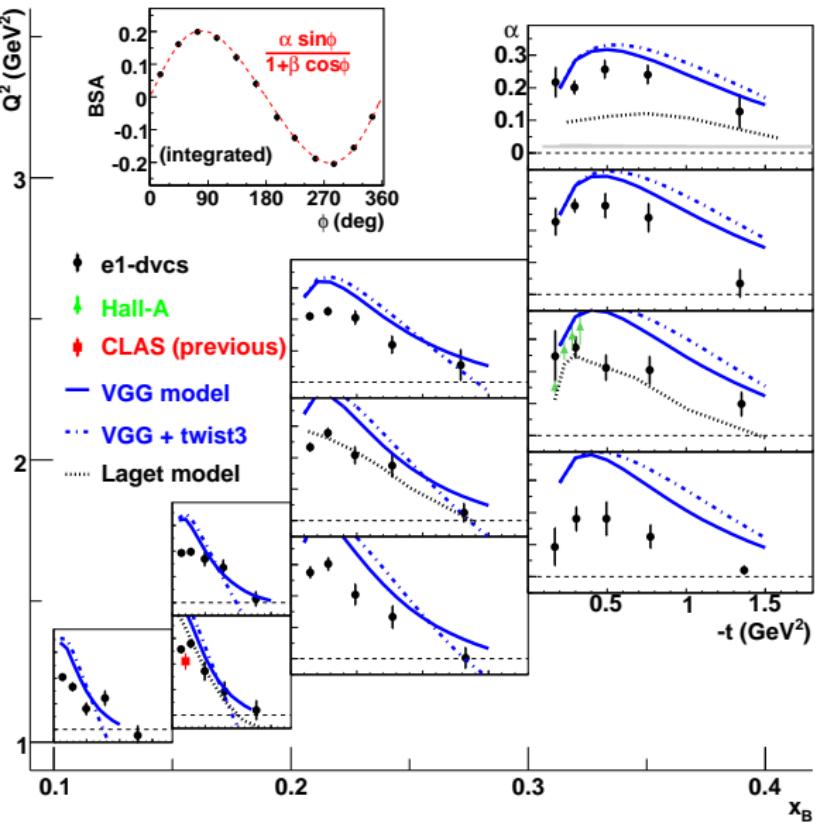
*A global analysis is needed to fully disentangle GPDs  
All channels can also be measured in transition GPDs to excited Nucleon states*

# Published Data



# DVCS Beam Spin Asymmetry

6 GeV



$$F_1 \mathcal{H} + \xi G_M \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$

Precision in a large phase-space ( $x_B, Q^2, t$ )

Qualitative model agreement

quantitative constraints on parameters



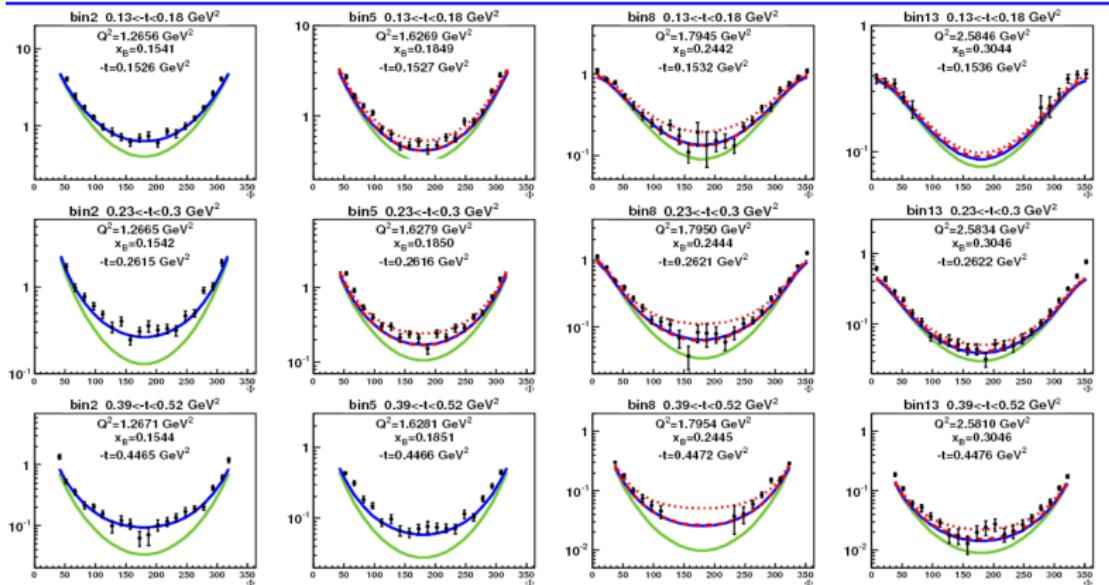
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# DVCS Unpolarized Cross-Sections 6 GeV



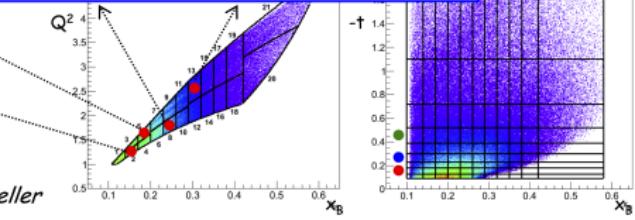
•  $\frac{d^4\sigma_{ep \rightarrow e\gamma}}{dQ^2 dx_B dt d\Phi} (\text{nb}/\text{GeV}^4)$

— BH — VGG (H only)

..... KM10 - - KM10a

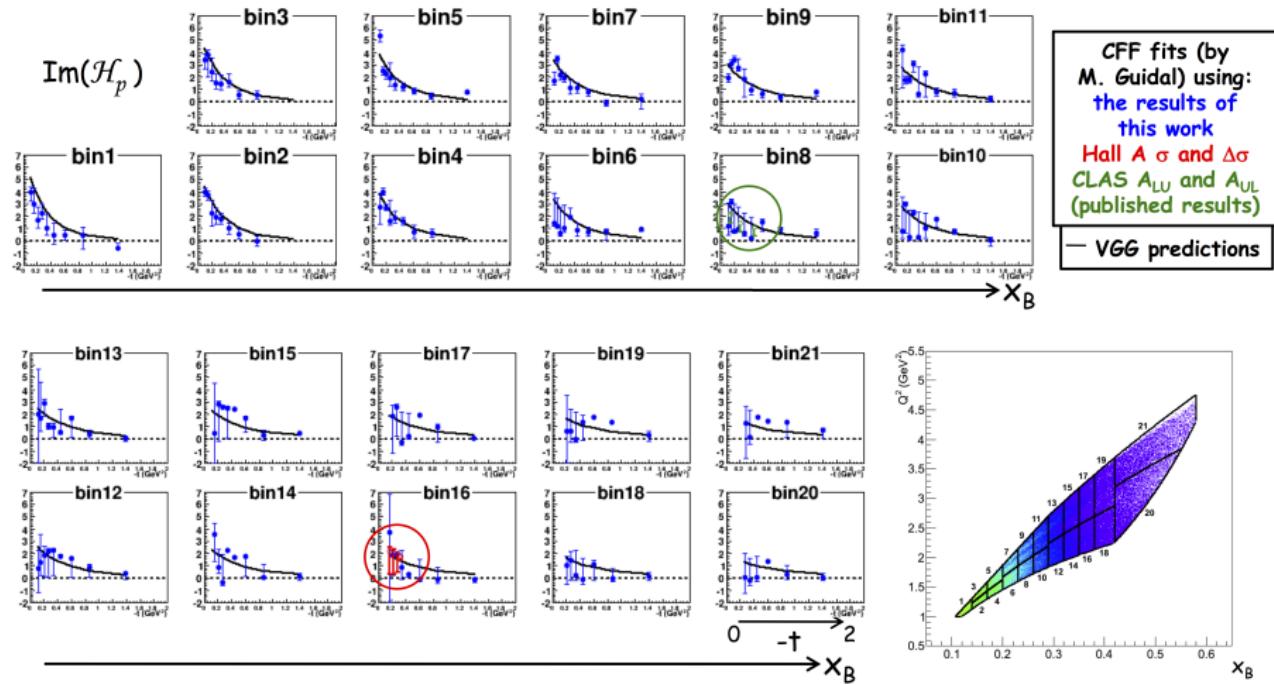
VGG : Vanderhaeghen, Guichon, Guidal

KM : Kumericki, Mueller



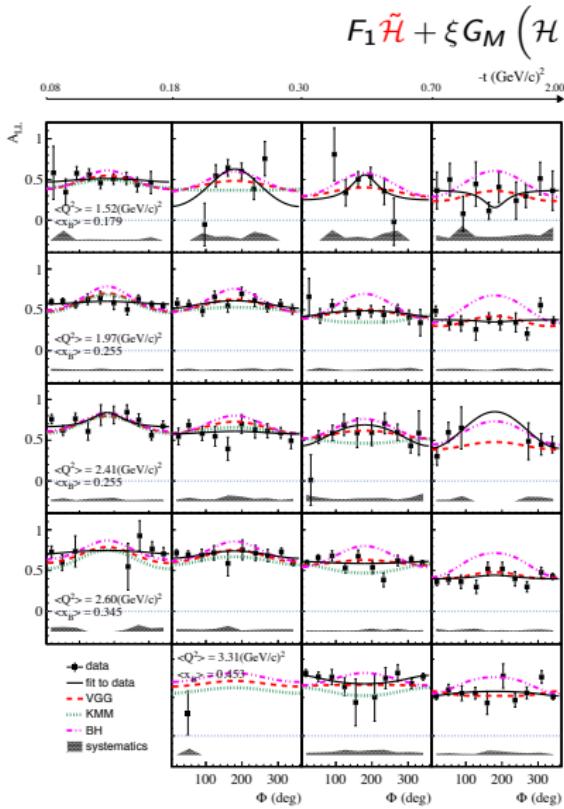
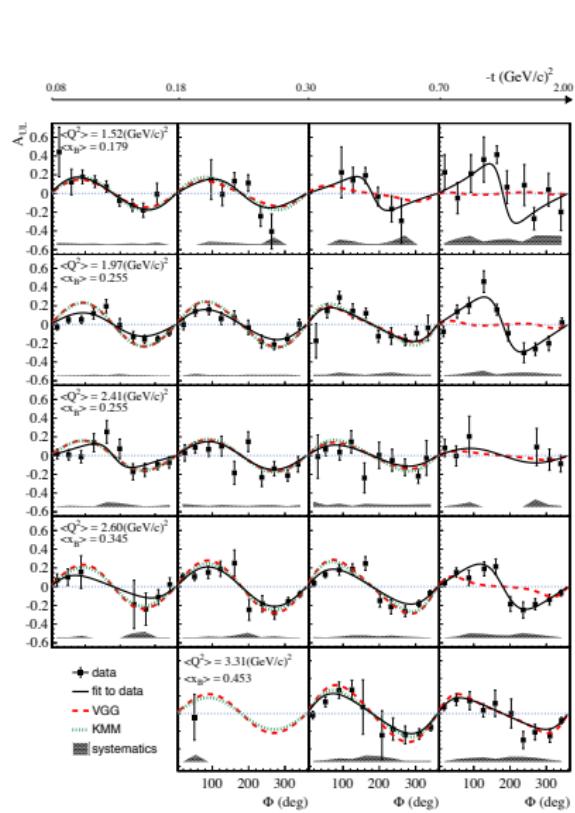
# Compton Form Factors

6 GeV



The  $t$ -slope becomes flatter with increasing  $x_B$ :  
valence quarks (higher  $x_B$ ) at the center of the nucleon and sea quarks (small  $x_B$ ) at its periphery

# Target Longitudinal Spin DVCS 6 GeV



# Model independent extraction

6 GeV

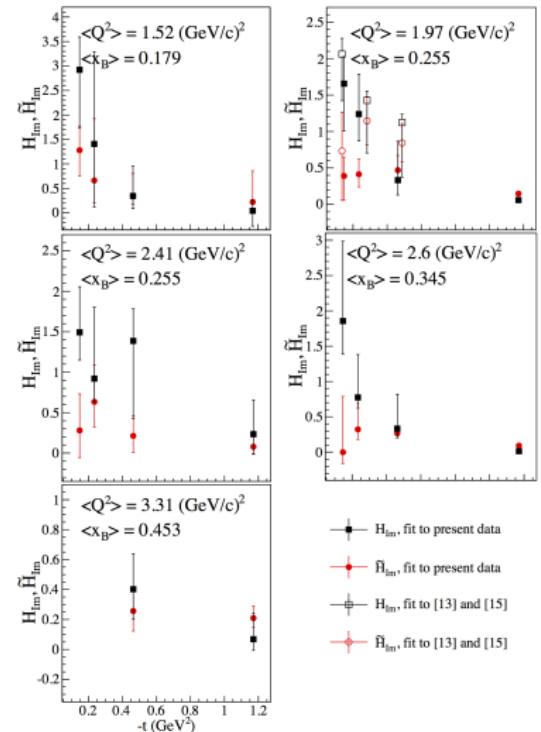
GPD dependencies versus  $x_B$  mirror their respective ordinary PDFs

$\tilde{H}$  and  $H \leftrightarrow \Delta q(x)$  and  $q(x)$

Change of  $\Delta q(x)$  t-slope vs  $x_B$  less pronounced than  $q(x)$

Axial charge more concentrated than

EM charge



# Future Measurements



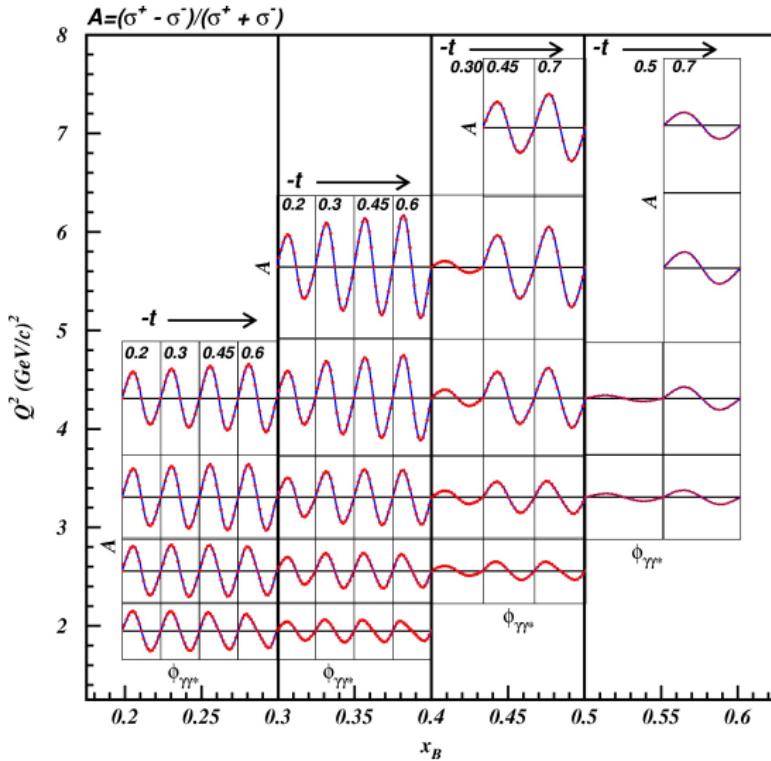
# CLAS12 GPD program

Number	Title	Contact	Days	Energy	Target
E12-06-108	Hard Exclusive Electroproduction of $\pi^0$ and $\eta$	Kubarovski	80	11	IH <sub>2</sub>
E12-06-119	Deeply Virtual Compton Scattering	Sabatie	80	11	IH <sub>2</sub>
E12-12-001	Timelike Compton Scat. & J/ψ prod. in e <sup>+</sup> e <sup>-</sup>	Nadel-Turonski	120	11	IH <sub>2</sub>
E12-12-007	Exclusive $\phi$ meson electroproduction	FXG	60	11	IH <sub>2</sub>
E12-11-003	DVCS on Neutron Target	Niccolai	90	11	ID <sub>2</sub>
E12-06-119	Deeply Virtual Compton Scattering	Sabatie	120	11	NH <sub>3</sub>
C12-12-010	DVCS with a transverse target	Elouadrhiri	110	11	HD-ice
E12-16-010	DVCS with CLAS12 at 6.6 GeV and 8.8 GeV	Elouadrhiri	50+50	6.6 & 8.8	IH <sub>2</sub>



80 days @  $\mathcal{L} = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  with 85% polarized beam

$$A_{LU} \propto F_1 \mathcal{H} + \xi G_M \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$



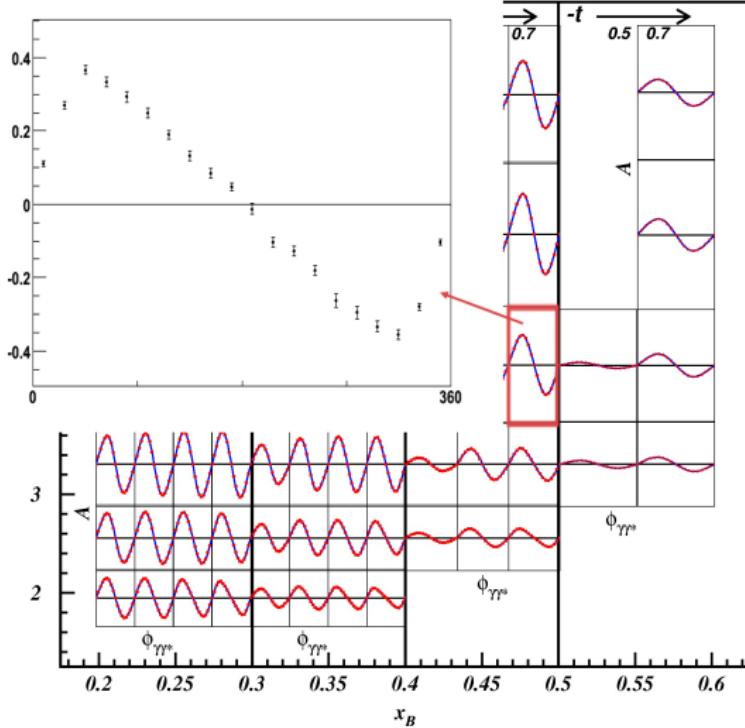
Beam Spin Asymmetries  
 $\phi$  dependence

Statistical uncertainties :  
from 1 % (low  $Q^2$ )  
to 10 % (high  $Q^2$ )

Unprecedented statistics  
over the full  $\phi$  range  
up to high  $x = 0.6$

80 days @  $\mathcal{L} = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  with 85% polarized beam

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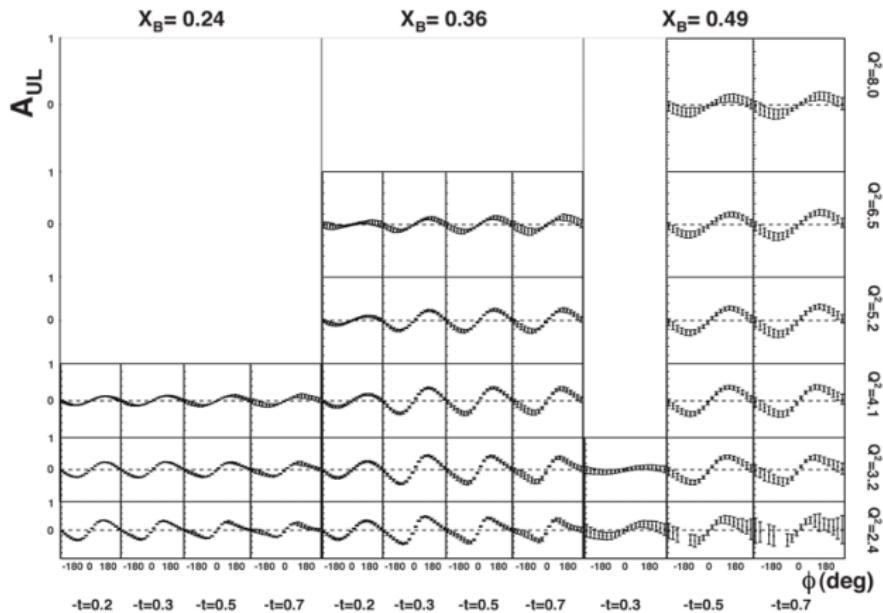
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120 days @  $\mathcal{L} = 2 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  with 80% polarized NH<sub>3</sub>

$$A_{UL} \propto F_1 \tilde{\mathcal{H}} + \xi G_M \left( \mathcal{H} + \frac{\xi}{1+\xi} \mathcal{E} \right) - \dots$$



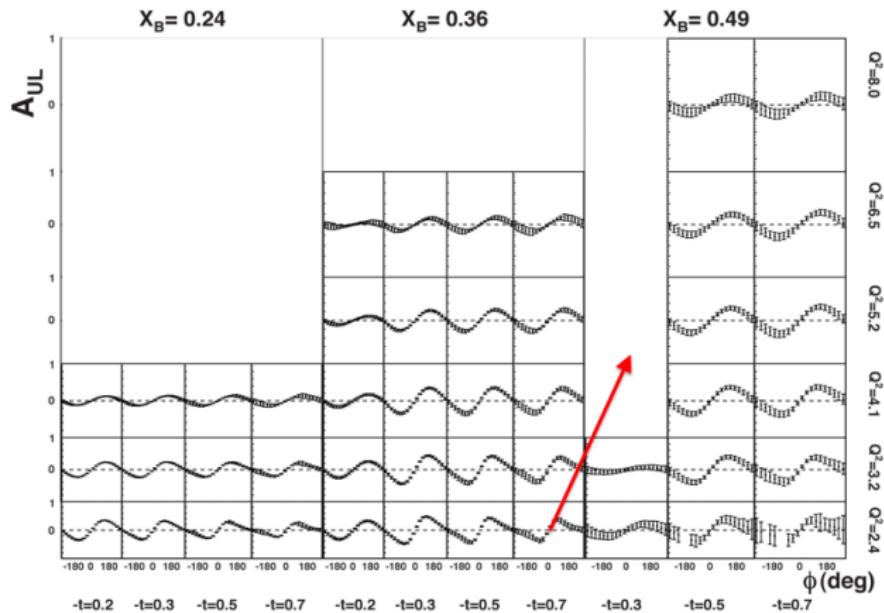
Target Spin Asymmetries  
 $\phi$  dependence

Statistical uncertainties :  
from 2 % (low  $Q^2$ )  
to 30 % (high  $Q^2$ )

Unprecedented statistics  
over the full  $\phi$  range  
up to high  $x = 0.6$

120 days @  $\mathcal{L} = 2 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  with 80% polarized NH<sub>3</sub>

$$A_{UL} \propto F_1 \tilde{\mathcal{H}} + \xi G_M \left( \mathcal{H} + \frac{\xi}{1+\xi} \mathcal{E} \right) - \dots$$



Target Spin Asymmetries  
 $\phi$  dependence

Statistical uncertainties :  
from 2 % (low  $Q^2$ )  
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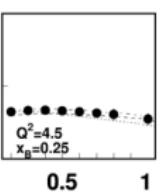
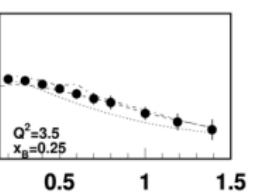
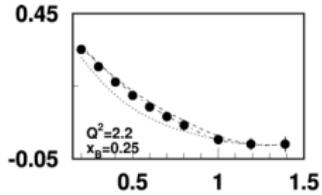
Unprecedented statistics  
over the full  $\phi$  range  
up to high  $x = 0.6$

# Proton DVCS TSA $A_{UL}$

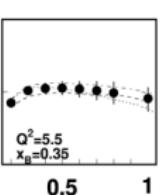
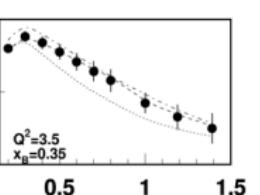
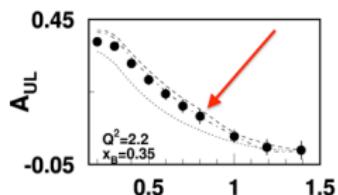
E12-06-009

120 days @  $\mathcal{L} = 2 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  with 80% polarized  $\text{NH}_3$

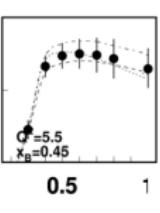
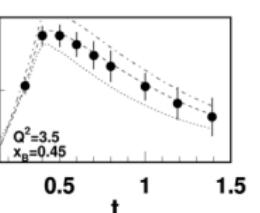
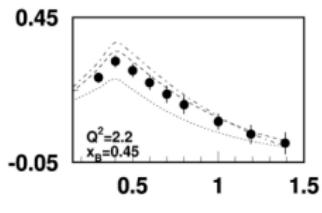
$$A_{UL} \propto F_1 \tilde{\mathcal{H}} + \xi G_M \left( \mathcal{H} + \frac{\xi}{1+\xi} \mathcal{E} \right) - \dots$$



TSA t-slopes



Sample kinematics  
for target asymmetry



Change of  $t$ -slope with  $x_B$   
 $\leftrightarrow$   
 imaging  $\Delta q(x_B, b_\perp)$

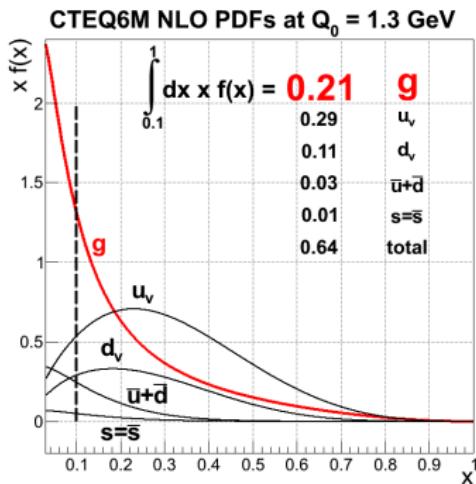


JSA FX Girod

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- Large glue density at  $x > 0.1$

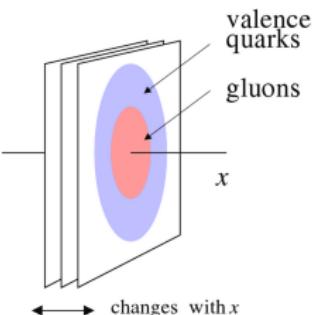
PDF from global fits  
( $F_2$  evolution,  $\nu_{\text{DIS}}$ , jets)

Gluons carry more than 30%  
of the momentum for  $0.1 < x$

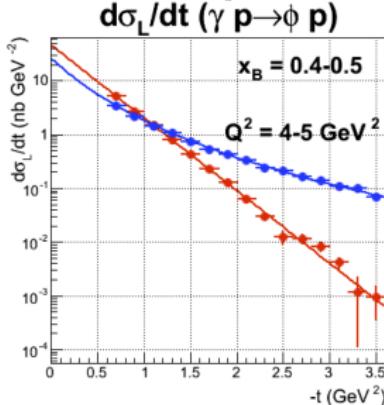
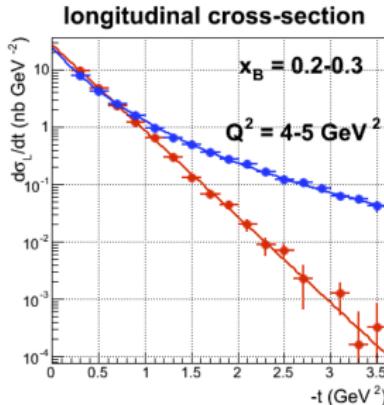
- 3D imaging of the nucleon

spatial distribution of valence quarks :  
elastic scattering, DVCS, ...

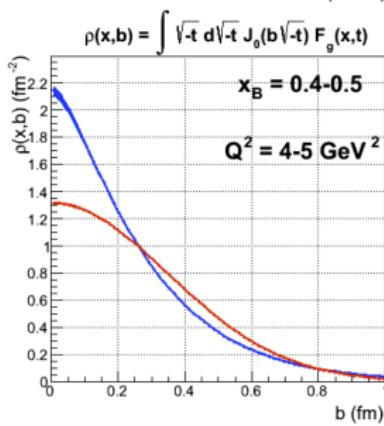
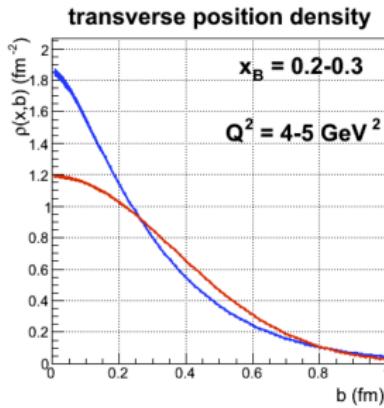
Nucleon gluonic radius ?  
exclusive  $\phi$



# Extraction of gluonic profiles



Longitudinal cross-section



Corresponding sensitivity in transverse position space

$$b = 1/\sqrt{-t}$$

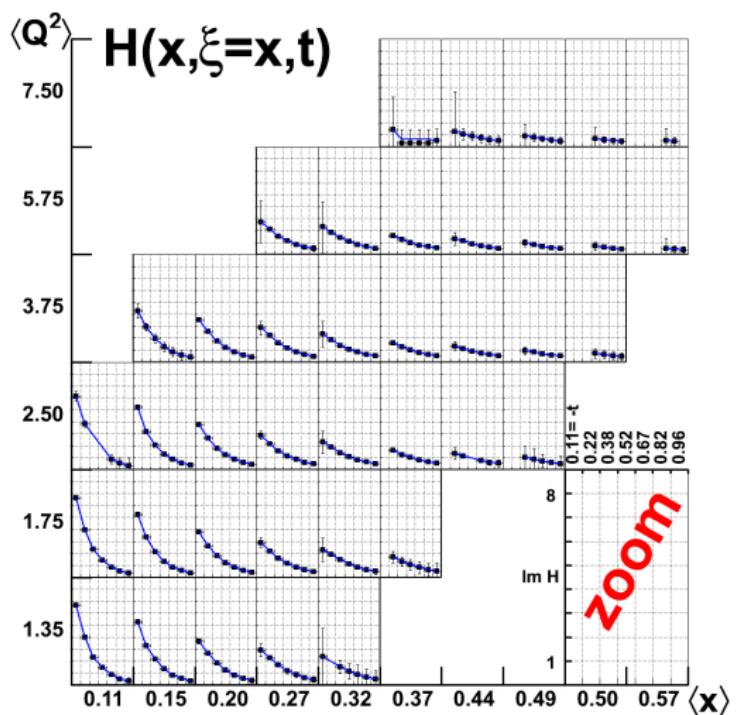
Error propagation study  
Skewness  $\xi \neq 0$  neglected

# Projected Impact on GPD Extractions



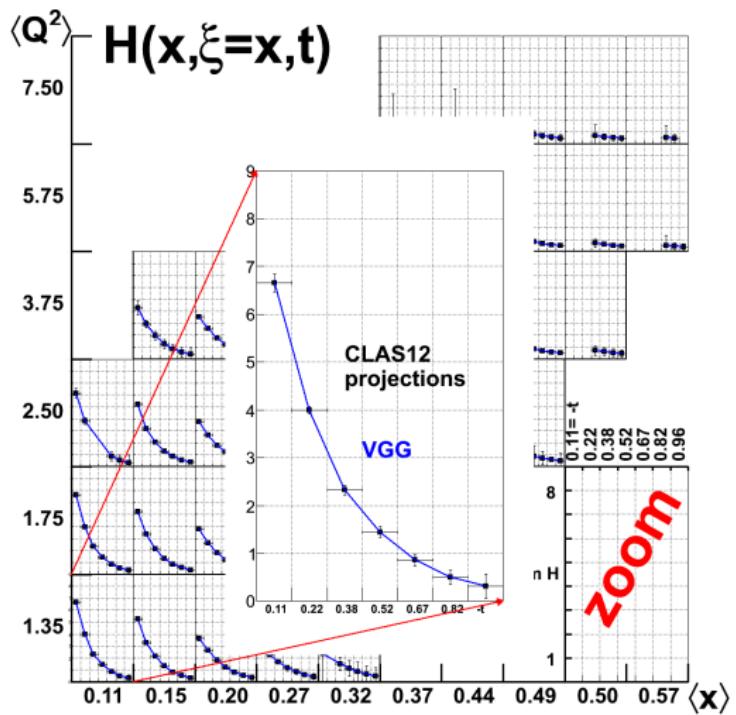
# Projected Impact on GPD Extraction

Using simulated data  
based on VGG model.  
Input GPD H extracted  
with good accuracy



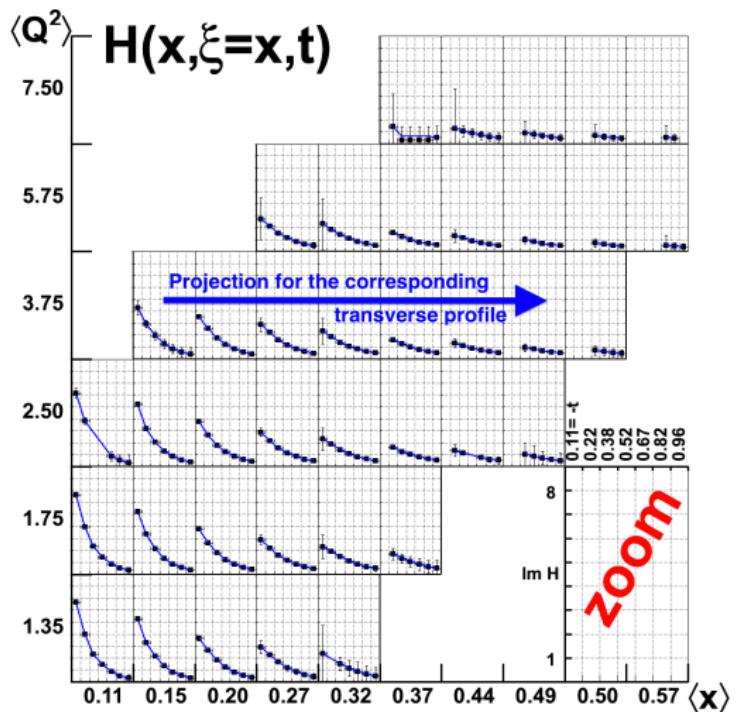
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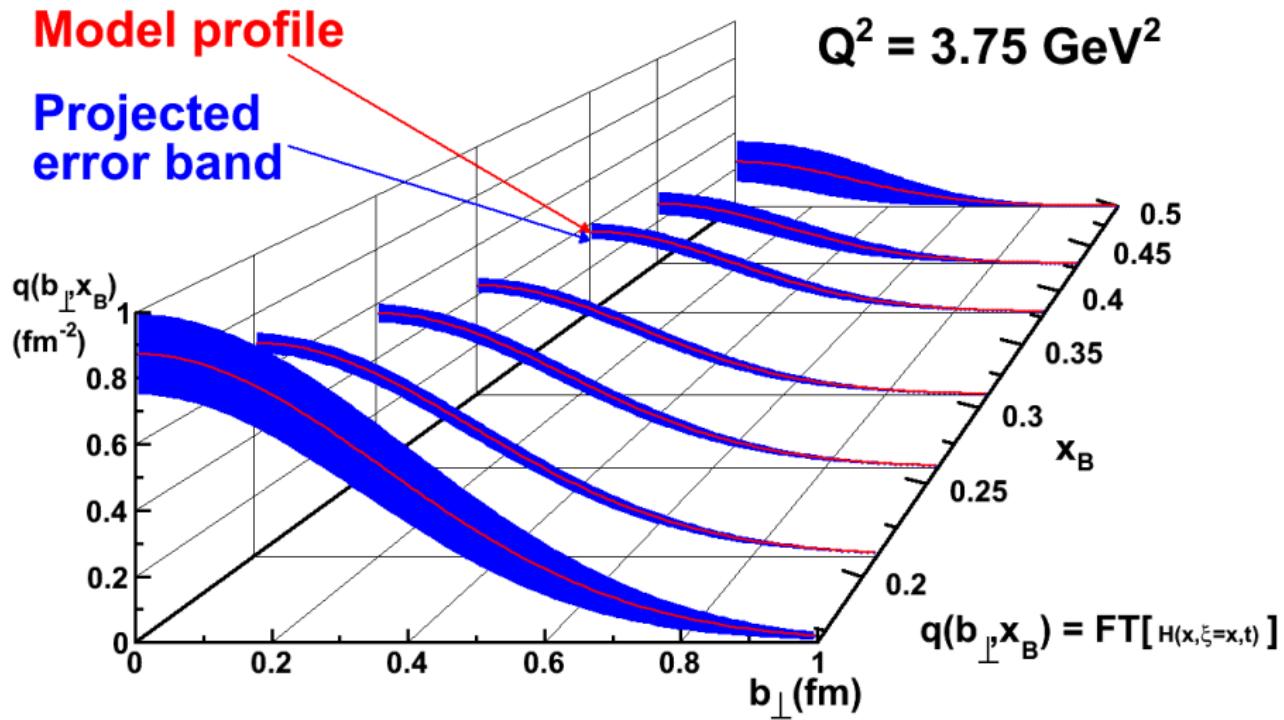


# Projected Impact on GPD Extraction

Using simulated data based on VGG model.  
Input GPD H extracted with good accuracy

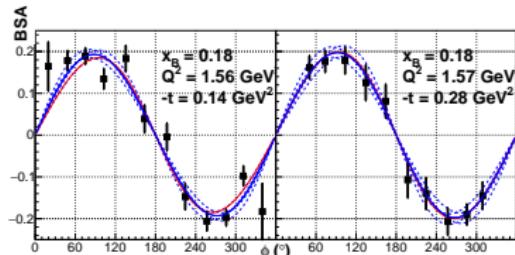


# Projection for the Nucleon transverse profile



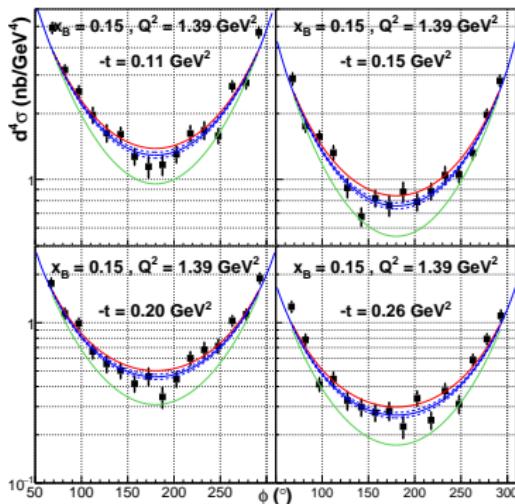
Precision tomography in the valence region

# Global Fits to extract the D-term



Beam Spin Asymmetries

$$\text{Im}\mathcal{H}(\xi, t) = \frac{r}{1+x} \left( \frac{2\xi}{1+\xi} \right)^{-\alpha(t)} \left( \frac{1-\xi}{1+\xi} \right)^b \left( \frac{1-\xi}{1+\xi} \frac{t}{M^2} \right)^{-1}$$



Unpolarized cross-sections

Use dispersion relation:

$$\text{Re}\mathcal{H}(\xi, t) = D + \mathcal{P} \int dx \left( \frac{1}{\xi - x} - \frac{1}{\xi + x} \right) \text{Im}\mathcal{H}(\xi, t)$$

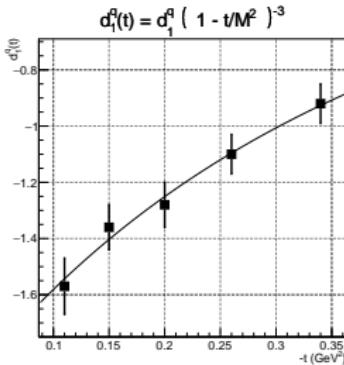
pure Bethe-Heitler

local fit + uncertainty range

resulting global fit



# D-term and Pressure distribution



$$D^q\left(\frac{x}{\xi}, t\right) = \left(1 - \frac{x^2}{\xi^2}\right) \left[ d_1^q(t) C_1^{3/2}\left(\frac{x}{\xi}\right) + d_3^q(t) C_3^{3/2}\left(\frac{x}{\xi}\right) + \dots \right]$$

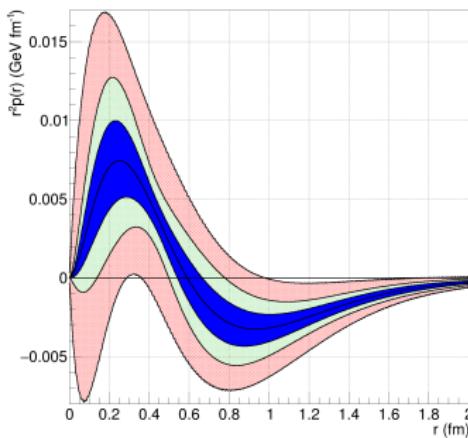
t-dependence of the D-term :

Dipole gives singular pressure at  $r = 0$

Quadrupole implied by counting rules?

Exponential?

...



Resulting pressure distribution

$$\text{Stability condition : } \int_0^\infty dt r^2 p(r) = 0$$

World data fit

CLAS 6 GeV data

Projected CLAS12 data

# Summary and Outlook

- Exclusive reactions offer a partonic approach to the problem of confinement
- Transverse Imaging, Energy Momentum Tensor
- First Generation Experiments successful
- Extraction frameworks already producing publications
- Formalism also applies to Nucleon Resonances
- Exciting times at the beginning of the 12 GeV high precision era!

