Excitations of the Nucleon - $\mathbf{N}^{*}$ spectroscopy with clos

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## $\mathbb{N}^{*}$ resonances in the early universe

Dramatic chiral crossover at about $t_{0}+1 \mu \mathrm{~s} \Leftrightarrow T_{c}=154 \pm 9 \mathrm{MeV}$


- Chiral symmetry is broken
- quarks acquire mass

- color confinement emerges
- copious production of hadronic resonances
- Paolo Neruda (paraphrased):
" Make you choice in life, but then embrace the consequences. "
$\Leftrightarrow$ leaves little room for desperation or depression !


## Where are these $N^{*}$ resonances?



- ...up to a decade ago: only the lowest few in each band correspond to 4* or 3* PDG states
- Vintage explanations: \{eg.Anselmino et al, Rev. Mod. Phys. 65 (93) 1199\}
- 2 quarks in a baryon quasibound in a color isotriplet $[d i Q+q \Leftrightarrow$ isosinglet]
- internal "diQuark" excitations frozen out in spin $=0$, isospin $=0$ states $\Leftrightarrow$ fewer degrees of freedom $\Leftrightarrow$ fewer states


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## $N^{*}$ s that just can't go away ...

## DSE \& LQCD calculations of $\mathbf{N}^{*}$ spectrum:

- axial-vector color-triplet q-q correlations are attractive
- diQuarks correlations must exist
- BUT, they are not point-like, eg.r $[u d]_{1} \sim$ pion radius $\{$ Few Body Sys 35 (04); PRL97(00)\}
$\Leftrightarrow$ internal diQuark excitations are NOT frozen out
$\Leftrightarrow$ diQuark correlations are already observed in the LQCD calculations


## LQCD calculations of the $T_{c}$ phase transition in the $\sim 1 \mu$ universe:

- PDG states alone are insufficient
- full suite of $Q M / L Q C D$ states needed $\Leftrightarrow \sim 25 \%$ baryon pressure increase needed from as yet unobserved $N^{*}$ S
\{ Bazavov et al., PRL 113 (2014) 072001 \}



## goals of the $N^{*}$ program with CLAS at Jefferson Lab

- dressings of strongQCD (non-perturbative) generate a running quark-mass function $\Leftrightarrow$ "constituent-like" correlations at low $\boldsymbol{p}$
that generate the $N^{*}$ spectrum
$\Leftrightarrow$ account for $\sim 98 \%$ of the visible mass (Higgs mechanism is the other $\sim 2 \%$ )


## - CLAS goals:

- elucidate the structure of $N^{*}$ states that are observed, and find the ones that aren't!
- clarify the role of complex correlations:

"constituent"
range
Higgs
- meson cloud
- dynamical meson-baryon "molecules"

sQCD emerged $a t \sim t_{0}+1 \mu s$ when $N^{*} s$ filled the Universe


## Challenge \#1 - $N^{*}$ resonances are broad and overlapping

- $\pi N \rightarrow \pi N \Leftrightarrow$ chief source of pre-2008 PDG states
- 2 complex spin-dependent amplitudes $\Leftrightarrow$ requires 3 observables to define the full amplitude, within a phase
- there are only 4 observables $(\sigma, P, \boldsymbol{R}, \boldsymbol{A})$
- data base: $\sim 40,000$ pts on $\boldsymbol{\sigma} ; \sim 8300$ pts on $\boldsymbol{P}$

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- almost no data on $\boldsymbol{R} \& \boldsymbol{A}(\sim 30$ pts above the $\Delta)$ $\Leftrightarrow$ amplitude is under-determined
$\Leftrightarrow$ very difficult to isolate weaker states

$$
\sigma_{\text {tot }}(\pi+N)
$$



## Challenge \#1 - $N^{*}$ resonances are broad and overlapping

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- $\boldsymbol{\gamma N} \boldsymbol{\rightarrow} \boldsymbol{\pi} \mathbf{N}, \boldsymbol{\eta} \mathbf{N}, \boldsymbol{K} \boldsymbol{Y}, \ldots$
- 4 complex amplitudes
$\Leftrightarrow$ requires 7 (8) observables to define the full amplitude, within a phase
- there are 16 observables:
( $\sigma, \Sigma, T, P, E, G, F, H, O_{x^{\prime}}, O_{z^{\prime}}, C_{x^{\prime}}, C_{z^{\prime}}, L_{x^{\prime}}, L_{z^{\prime}}, T_{x^{\prime}}, T_{z^{\prime}}$ )
$\Leftrightarrow$ possible to over-determine the amplitude



NP B41(1972) 445
[SHKL, J Phys G38 (11) 053001]

| Photon beam |  | Target |  |  | Recoil |  |  | Target - Recoil |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $x$ | $y$ | $z$ | $x^{\prime}$ | $y^{\prime}$ | $z$ ' | $\begin{aligned} & x^{\prime} \\ & x \end{aligned}$ | $x^{\prime}$$y$ | $\begin{gathered} x^{\prime} \\ z \end{gathered}$ | $\begin{gathered} y^{\prime} \\ x \end{gathered}$ | $\begin{array}{\|l\|} \hline y^{\prime} \\ y \end{array}$ | $\begin{gathered} y^{\prime} \\ z \end{gathered}$ | $\begin{gathered} z^{\prime} \\ x \end{gathered}$ | $\begin{aligned} & z^{\prime} \\ & y \end{aligned}$ | $\begin{gathered} z^{\prime} \\ \hline z \end{gathered}$ |
|  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| unpolarized | $\sigma_{0}$ |  | $T$ |  |  | $\boldsymbol{P}$ |  | $\boldsymbol{T}_{x}$, |  | $L_{x}$, |  | $\Sigma$ |  | $\boldsymbol{T}_{z}$, |  | $L_{z}$, |
| $\boldsymbol{P}_{L}^{\gamma} \sin \left(2 \phi_{\gamma}\right)$ |  | H |  | G | $\boldsymbol{O}_{\boldsymbol{x}}$, |  | $\boldsymbol{O}_{z}$, |  | $C_{z}$, |  | E |  | F |  | ${ }_{-} C_{x}$, |  |
| $\boldsymbol{P}_{L}^{\gamma} \cos \left(2 \phi_{\gamma}\right)$ | $-\Sigma$ |  | $-P$ |  |  | $-T$ |  | $-\boldsymbol{L}_{z}$, |  | $\boldsymbol{T}_{\boldsymbol{z}}$, |  | $-\sigma_{0}$ |  | $L_{x}$, |  | $-T_{x}$, |
| circular $\boldsymbol{P}_{\boldsymbol{c}}^{\gamma}$ |  | F |  | -E | $C_{x}$, |  | $C_{z}$, |  | $-\boldsymbol{O}_{z}$, |  | G |  | $-H$ |  | $\boldsymbol{O}_{\boldsymbol{x}}$, |  |



- 16 different observables
- combine asymmetries for different final states in a coupled-channel PWA
$\Leftrightarrow$ identify $N^{*}$ resonances
$\Leftrightarrow$ extract $\gamma N^{*}$ couplings


## Challenge \#2: the dressings of sQCD

## $N^{*}$ resonance $\Leftrightarrow$ s-channel pole



- meson-loop "dressings" of the Electromagnetic vertex affect the dynamical properties (excitation mechanism) and determine $Q^{2}$ evolution, but not spectral properties
- coupled-channel "dressings" of the strong vertex determine the $N^{*}$ spectral properties (mass/pole positions, widths)


## Challenge \#2: the dressings of sQCD - eg. the $P_{11}$ Roper clos*

$N^{*}$ resonance $\Leftrightarrow$ s-channel pole


- coupled-channel "dressings" of the strong vertex determine the $N^{*}$ spectral properties
- dynamic coupled-channel model of $\pi N, \gamma N \rightarrow \pi N, \pi \Delta, \eta N, K Y$
[EBAC/AO, PRL 104 (2010) 042302]
$\Leftrightarrow$ "bare" $N^{*}$ excitation at 1763 evolves to doublet of poles at ~1360
$\Leftrightarrow$ no PWA of a single channel can be sufficient with such couplings


## Challenge \#2: the dressings of sQCD - eg. the $P_{11}$ Roper clos

## $N^{*}$ resonance $\Leftrightarrow$ s-channel pole



- meson-loop "dressings" of the Electromagnetic vertex affect the dynamical properties (excitation mechanism) and determine $Q^{2}$ evolution, but not spectral properties
- $Q^{2}$ evolution demonstrates the basic character of the second $J^{\pi}=1 / 2^{+}$ state of the nucleon as a radial excitation of a dressed $3 q$ core
[ Chen, El-Bennich, Roberts, et al., arXiv:1711.03142]



## ©EBAF Large Acceptance Spectrometer (CLAS): 1997-2012

- tagged photon beams

circular polarization from bret of polarized $e^{-}$
linear polarization from $e^{-}$prem in diamond

- longitudinally polarized $e^{-}$beams


## FROST - frozen-spin proton target


$\gamma p \rightarrow \pi^{+}(n)$
g9a: Strauch et al., PL B750 (2015) 53


- target: $15 \mathrm{~mm} \varnothing \times 50 \mathrm{~mm}$
- material: $\mathrm{C}_{4} \mathrm{H}_{9} \mathrm{OH}$ (butanol)
- p-dilution: $10 / 74$
- $P(H)=83 \%$
- $T_{1}(1 / e$ spin relaxation $)=115 d(+h)$

$$
=65 \mathrm{~d}(-h)
$$

- repolarize ~ weekly


FROST: NIM A684 (2012) 27

## HDice - frozen-spin target for neutrons



$$
\gamma n \rightarrow \pi^{-} p
$$

g14: Ho et al., PRL 118 (2017) 242002


- target: $15 \mathrm{~mm} \varnothing \times 50 \mathrm{~mm}$
- material: solid HD
- p-dilution: $1 / 2$; $n$-dilution: $1 / 1$
- $P(H)=60 \%$ or $P(D)=30 \%$
- $T_{1}$ (1/e spin relaxation) $\sim$ years


$$
\boldsymbol{\Sigma}\left(\vec{\gamma} \boldsymbol{p} \rightarrow \pi^{0} n\right)
$$

| Photon beam |  | Target |  |  | Recoil |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $x$ |  |  | $x^{\prime}$ | $y^{\prime}$ | $z '$ |
|  |  |  | $y$ | $z$ |  |  |  |
| unpolarized | $\sigma_{0}$ |  | $T$ |  |  | $\boldsymbol{P}$ |  |
| $\boldsymbol{P}_{L}^{\gamma} \sin \left(2 \phi_{\gamma}\right)$ |  | H |  | G | $\boldsymbol{O}_{x}$ |  | $\boldsymbol{O}_{z}$, |
| $P_{L}^{\gamma} \cos \left(2 \phi_{\gamma}\right)$ | - $\Sigma$ |  | $-P$ |  |  | -T |  |
| circular $\boldsymbol{P}_{c}^{\gamma}$ |  | F |  | -E | $C_{x}$, |  | $C_{z}$, |



| Photon beam |  | Target |  |  | Recoil |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $x$ |  |  | $x^{\prime}$ | $y^{\prime}$ | $z^{\prime}$ |
|  |  | $y$ | $z$ |  |  |  |
| unpolarized | $\sigma_{0}$ |  |  | $T$ |  |  | $\boldsymbol{P}$ |  |
| $\boldsymbol{P}_{L}^{\gamma} \sin \left(2 \phi_{\gamma}\right)$ |  | H |  | G | $O_{x}$, |  | $\boldsymbol{O}^{\prime}{ }^{\prime}$ |
| $\boldsymbol{P}_{L}^{\gamma} \boldsymbol{\operatorname { c o s }}\left(2 \phi_{\gamma}\right)$ | $-\Sigma$ |  | -P |  |  | -T |  |
| circular $\boldsymbol{P}_{\text {c }}^{\gamma}$ |  | $F$ |  | -E | $C_{x}$, |  | $C_{z}$, |



| Photon beam |  | Target |  |  | Recoil |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $x$ |  |  | $x^{\prime}$ | $y^{\prime}$ | $z$ ' |
|  |  |  | $y$ | $z$ |  |  |  |
| unpolarized | $\sigma_{0}$ |  | $T$ |  |  | $\boldsymbol{P}$ |  |
| $\boldsymbol{P}_{L}^{\gamma} \sin \left(2 \phi_{\gamma}\right)$ |  | H |  | G | $\boldsymbol{O}_{\boldsymbol{x}}$, |  | $\boldsymbol{O}_{z}$, |
| $P_{L}^{\gamma} \cos \left(2 \phi_{\gamma}\right)$ | - $\Sigma$ |  | $-P$ |  |  | -T |  |
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| Photon beam |  | Target |  |  | Recoil |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  | $x^{\prime}$ | $y^{\prime}$ | $z^{\prime}$ |
|  |  | $x$ | $y$ | $z$ |  |  |  |
| unpolarized | $\sigma_{0}$ |  | $T$ |  |  | $P$ |  |
| $\boldsymbol{P}_{L}^{\gamma} \sin \left(2 \phi_{\gamma}\right)$ |  | H |  | G | $O_{x}$, |  | $\mathrm{O}_{\mathbf{z}}$, |
| $\boldsymbol{P}_{L}^{\gamma} \cos \left(2 \phi_{\gamma}\right)$ | $-\Sigma$ |  | -P |  |  | $-T$ |  |
| circular $\boldsymbol{P}_{c}^{\boldsymbol{\gamma}}$ |  | $F$ |  | -E | $C_{x}$, |  | $C_{z}$, |



## $P_{\Lambda}\left(\gamma \boldsymbol{p} \rightarrow K^{+} \vec{\Lambda}\right)$

CLAS-g11 : Phys Rev C81 (2010) 025201 clos


$C_{x}, C_{z}\left(\vec{\gamma} p \rightarrow K^{+} \vec{\Lambda}\right)$
CLAS-g1c : Phys Rev C75 (2007) 035205 clos



$O_{x}, O_{z}\left(\vec{\gamma} p \rightarrow K^{+} \vec{\Lambda}\right)$
CLAS-g8 : Phys Rev 993 (2016) 065201 clos



Final states and observables measured in CLAS

|  | $\boldsymbol{\sigma}$ | $\boldsymbol{\Sigma}$ | $\mathbf{T}$ | $\mathbf{P}$ | $\mathbf{E}$ | $\mathbf{F}$ | $\mathbf{G}$ | $\mathbf{H}$ | $\mathbf{T}_{\mathrm{x}}$ | $\mathbf{T}_{\mathbf{z}}$ | $\mathbf{L}_{\mathrm{x}}$ | $\mathbf{L}_{\mathbf{z}}$ | $\mathbf{O}_{\mathrm{x}}$ | $\mathbf{O}_{\mathbf{z}}$ | $\mathbf{C}_{\mathrm{x}}$ | $\mathbf{C}_{\mathbf{z}}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | Proton target: $\gamma p \rightarrow X$


| $\mathrm{p} \mathrm{r}^{0}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{n} \mathrm{x}^{+}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |  |  |  |  |  |  |  |  |
| pM | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |  |  |  |  |  |  |  |  |
| p ${ }^{\prime}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |  |  |  |  |  |  |  |  |
| $\mathrm{pr} \mathrm{x}^{+} \pi$ | $\checkmark$ | $\checkmark$ |  |  | $\checkmark$ |  | $\checkmark$ | $\checkmark$ VNU +8 more лл兀 observables UNUN |  |  |  |  |  |  |  |  |
| p $\omega$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | SDME |  |  |  |  |  |  |  |
| $\mathbf{K}+\boldsymbol{\Lambda}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| $\mathbf{K}^{+} \mathbf{\Sigma}^{0}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| $\mathbf{K}^{0+} \mathbf{\Sigma}^{+}$ | $\checkmark$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| $\mathbf{K}^{0} \mathbf{\Sigma}^{+}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |  |  |  |  |  |  |  |  |

"Neutron" target: $\gamma n \rightarrow X$


Published

Confirmation of new states near $W=1900$

## - Bonn-Gatchina + Zagreb PWA:

- CLAS $\gamma p \rightarrow K^{+} \Lambda$ used to fix $L=0,1$ multipoles
- used in a coupled-channel search for poles
$\Leftrightarrow$ reveals new $N^{*}$ s that couple strongly to $K \Lambda$ (but weakly to $\pi N$; not evident in $\pi N \rightarrow \pi N$ )


$N(1860) 1 / 2^{+}$



Confirmation of 1 -star $\mathbf{\Delta ( 2 2 0 0 )} 7 / 2^{-}$in coupled-channel PWA clos

- well established $\Delta(1950) 7 / 2^{+}[P D G * * * *$ missing a parity-partner
$\Leftrightarrow$ possible weak $\Delta(2200) 7 / 2^{-}\left[P D G^{*}\right]$ ?
- Bonn-Gatchina coupled-channel PWA of CLAS and CBELSA/TAPS data from many channels [ Phys Lett B766 (2017) 357]
$\Leftrightarrow$ requires $\Delta$ (2176) 7/2-
- small $\pi N$ branch $\Leftrightarrow$ very weak in $\pi N$ scattering
- but reflected in the $\gamma N \rightarrow \pi N$ " $E$ " asymmetries
- no evidence of mass-degenerate partners near 1950 (arguing against Chiral restoration)



## Photo-production from neutrons

- the electromagnetic interactions do not conserve isospin
$\mathcal{A}_{\gamma p \rightarrow \pi^{+} \boldsymbol{n}}=\sqrt{2}\left\{\mathcal{A}_{p}^{I=1 / 2}-\frac{1}{3} \mathcal{A}^{I=3 / 2}\right\}$
$\Leftrightarrow$ proton data determine $\boldsymbol{\mathcal { A }}^{I=3 / 2}$
$\mathcal{A}_{\gamma n \rightarrow \pi^{-} p}=\sqrt{2}\left\{\mathcal{A}_{n}^{I=1 / 2}+\frac{1}{3} \mathcal{A}^{I=3 / 2}\right\}$
$\Rightarrow$ both proton and neutron target data needed for the $I=1 / 2$ amplitudes
- $\quad \gamma+n$ data base is very sparse
$\Leftrightarrow \gamma n N^{*}$ couplings very poorly determined
$\Leftrightarrow$ CLAS run periods g10 $(\gamma+D)$,

$$
\begin{aligned}
& g 13(\vec{\gamma}+D), \\
& \text { g14 }(\vec{\gamma}+\vec{D})
\end{aligned}
$$

| Photon beam |  | Target |  |  | Recoil |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $x$ |  |  | $x^{\prime}$ | $y^{\prime}$ | $z^{\prime}$ |
|  |  | $y$ | $z$ |  |  |  |
| unpolarized | $\sigma_{0}$ |  |  | $T$ |  |  | $\boldsymbol{P}$ |  |
| $\boldsymbol{P}_{L}^{\gamma} \sin \left(2 \phi_{\gamma}\right)$ |  | H |  | G | $\boldsymbol{O}_{\boldsymbol{x}}$, |  | $\boldsymbol{O}_{z}$, |
| $\boldsymbol{P}_{L}^{\gamma} \cos \left(2 \phi_{\gamma}\right)$ | - $\Sigma$ |  | -P |  |  | -T |  |
| circular $\boldsymbol{P}_{c}^{\gamma}$ |  | F |  | -E | $C_{x}$, |  | $C_{z}$, |



## Deuteron reactions

## restricted to create an effective neutron target

- select events for which the proton in Deuterium is a passive "spectator" $\Leftrightarrow$ key variable is its momentum, eg. equivalently, the momentum of the undetected proton in $\gamma+n(p) \rightarrow \pi^{-} p(p)$ $\Leftrightarrow$ use the data itself to determine the kinematic region in which a measured observable is stable
- eg. the beam-target helicity asymmetry " $E$ ":
[ PRL 118 (2017) 242002 ]

$$
\left|\mathrm{P}_{\text {miss }}\right|<0.1 \mathrm{GeV} / \mathrm{c}
$$

- with these tight requirements, the D-state gives no contribution

- NB: stable region is observable dependent




## Evidence for an N(1975)3/2+ from $E\left(\vec{\gamma} \vec{n} \rightarrow \pi^{-} p\right)$



## A decade of advances in mapping the $N^{*}$ spectrum

| $N(W) J^{\pi}$ | PDG＇08 | PDG＇17 | ＋recent | $\gamma \mathrm{N}$ | $\pi N$ | KY |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N(1710) 1 / 2^{+}$ | $\star \star \star$ | $\star \star \star \star$ |  | $\checkmark \checkmark \checkmark \checkmark$ | $\checkmark \checkmark \checkmark \checkmark$ | $\checkmark \checkmark \checkmark \checkmark$ |
| N（1860） $1 / 2^{+}$ |  |  | $+v$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| $N(1860) 5 / 2^{+}$ |  | $\star \star$ |  |  | $\checkmark \checkmark$ |  |
| $N(1875) 3 / 2^{-}$ |  | 大 $\star$ 大 |  | $\checkmark \checkmark \checkmark$ | $\checkmark$ | $\checkmark \checkmark \checkmark$ |
| $N(1880) 1 / 2^{+}$ |  | 大 |  | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| $N(1895) 1 / 2^{-}$ |  | 大 $\star$ | $+v$ | $\checkmark \checkmark$ | $\checkmark$ | $\checkmark \checkmark$ |
| $N(1900) 3 / 2^{+}$ | $\star \star$ | $\star \star \star$ | $+v$ | $\checkmark \checkmark \checkmark$ | $\checkmark$ | $\checkmark \checkmark \checkmark$ |
| N（1975）3／2＋ |  |  | $+v$ | $\checkmark$ | $\checkmark$ |  |
| $N(2040) 3 / 2^{+}$ |  | $\star$ |  |  | $\checkmark$ |  |
| $N(2060) 5 / 2^{-}$ |  | 大 $\star$ |  | $\checkmark \checkmark$ | $\checkmark \checkmark$ | $\checkmark \checkmark$ |
| N（2100） $1 / 2^{+}$ | $\star$ | $\star$ | $+v$ | $\checkmark$ | $\checkmark$ |  |
| $N(2120) 3 / 2^{-}$ |  | 大 |  | $\checkmark \checkmark$ | $\checkmark \checkmark$ | $\checkmark$ |
| $N(2300) 1 / 2^{+}$ |  | 大 $\star$ |  |  | $\checkmark$ |  |
| $N(2570) 5 / 2^{-}$ |  | 大 $\star$ |  |  | $\checkmark \checkmark$ |  |
| $\Delta(1940) 3 / 2^{-}$ | $\star$ | $\star \star$ |  | $\checkmark$ | $\checkmark$ |  |
| $\Delta(2200) 7 / 2^{-}$ | $\star$ | $\star$ | $+v$ | $\checkmark$ | $\checkmark$ |  |

## Probing the dynamics of $N^{*}$ excitation with $Q^{2}$ in $\left(e, e^{\prime}\right) \quad$ clos $\$$




- $N^{*}$ excitation depend on $Q^{2}=-\left(\boldsymbol{k}_{e}-\boldsymbol{k}_{e^{\prime}}\right)^{2}$
$\Leftrightarrow$ different responses to changes in $Q^{2}$
- mapped out in large CLAS data sets eg. PRC 77 (2008)015208 ~ 35,000 data pts PRC 91 (2015) 045203 ~ 37,000 data pts



## Published CLAS data on

 exclusive meson electroproduction from protons| Hadronic final <br> state | Covered $\boldsymbol{W}$ range <br> $(\mathrm{GeV})$ | Covered $\boldsymbol{Q}^{2}$ range <br> $\left(\mathrm{GeV}^{2} / \mathrm{c}^{2}\right)$ | Measured <br> observables |
| :--- | :--- | :--- | :--- |
| $\pi^{+} \mathrm{n}$ | $1.1-1.38$ | $0.16-0.36$ | Observables <br> • cross section |
|  | $1.1-1.55$ | $0.3-0.6$ |  |
| angular distributions |  |  |  |

$Q^{2}$ evolution of photo $\rightarrow$ electro-couplings probe the $N^{*}$ excitation mechanisms

V. Burkert \& C. Roberts, arXiv:1710.02549

V. Burkert, NSTAR'2017

LF RQM: I. Aznauryan, V.B. arXiv:1603.06692
DSE: J. Segovia, C.D. Roberts et al., PRC94 (2016) 042201
$\Rightarrow$ Non-quark contributions are significant at $\mathrm{Q}^{2}<2.0 \mathrm{GeV}^{2}$.
$\rightarrow$ The $1^{\text {st }}$ radial excitation of the $q^{3}$ core emerges as the probe penetrates the MB cloud
> "Nature" of the Roper - its core is the $1^{\text {st }}$ radial excitation of the nucleon.
eg. $N(1535) 1 / 2^{-}$, the parity partner of the nucleon


- consistent couplings extracted from different decay channels (again)
- non-quark contributions are significant for $Q^{2}<1 \mathrm{GeV}^{2}$
- $\quad L F R Q M$ describes data for $Q^{2}>1.5 \mathrm{GeV}^{2}$
[I. Aznauryan \& V. Burkert, arXiv:1603.06692]
$\rightarrow N(1535) 1 / 2^{+}$is consistent with the $1^{\text {st }}$ orbital excitation of the nucleon
eg. $N(1675) 5 / 2^{-}$, a cloud-dominated resonance
K. Park et al, PRC 91 (2015) 045203

- $\quad \gamma p N^{*}\left(Q^{2}=0\right) \gg R Q M$
$\Leftrightarrow$ RQM is suppressed by selection rules, if only a single quark is excited [ Moorhouse, Phys Rev Lett 16 (1966) 772] [ Burkert et al, Phys Rev C67 (2003) 035204 ] ( NOT a meson-Baryon molecule )
- BUT, non-quark (meson-baryon cloud) contributions are significant for all $Q^{2}$


## $\gamma^{*} N \rightarrow N^{*}$ reveals a running quark-mass



- $\quad$ CLAS $G_{M}(N \rightarrow \Delta)$ (normalized to dipole)
"frozen" momentum-independent $M_{q}$ [Wilson et al., Phys Rev C85 (2012)025205] $\Leftrightarrow M_{q} \sim 300 \mathrm{MeV}$, dynamically generated by contact interactions btw current quarks
- dressed quark mass-function $M_{q}\left(p_{q}\right)$ [Roberts, J. Phys. Conf. 706 (2016) 022003] [Segovia et al., Few Body Phys 55 (2014) 1185]

- LQCD has confirmed Quark Model predictions for large numbers of $N^{*}$ states
$\Leftrightarrow$ no reduction in the effective degrees of freedom within the Nucleon
$\Leftrightarrow$ full LQCD/QM range of states required to provide the baryon pressure at $T_{c}$
- polarization in photo-production reactions can over-determine the amplitude $\Leftrightarrow$ extensive data on large numbers of polarization observables and final states have been collected and are in various stages of analysis
$\Leftrightarrow$ coupled-channel PWA have been essential in disentangling the $N^{*}$ spectrum $\Leftrightarrow$ large numbers of new candidate states have been identified
- $Q^{2}$ dependence of electro-production couplings provide insights to the role of the meson cloud and of the $N^{*}$ excitation mechanism
$\Leftrightarrow$ large data sets have been collected and analysis is ongoing ...
$\Leftrightarrow$ meson cloud effects are generally very strong below $\sim 1-2 \mathrm{GeV}^{2}$
$\Leftrightarrow$ transitions to $N^{*}$ s confirm a running quark mass-function

