Deeply virtual Compton scattering on Proton with CLAS12

CLAS collaboration

Maxime DEFURNE

June 23rd 2020



The nucleon: a formidable lab for QCD

- The nucleon is a dynamical object made of quarks and gluons.
- This dynamics is ruled by the strong interaction.
- A perturbative approach from first principles to unravel this dynamics is impossible due to the large size of the strong coupling constant.



Although non-perturbative approaches (DSE, lattice QCD) starts making progress, the experimental approach remains more convenient to get complex information about this dynamics.

A set of distributions encoding the nucleon structure



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DVCS and GPDs



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$$Q^2 = -q^2 = -(k - k')^2$$
.
• $x_B = \frac{Q^2}{2p \cdot q}$

- x longitudinal momentum fraction carried by the active quark.
- $\xi = \frac{x_B}{2-x_B}$ the longitudinal momentum transfer.
- $t = (p p')^2$ squared momentum transfer to the nucleon.

The GPDs enter the DVCS amplitude through a complex integral. This integral is called a *Compton form factor* (CFF).

$$\mathcal{H}_{++}(\xi,t) = \int_{-1}^{1} H(x,\xi,t) \left(\frac{1}{\xi - x - i\epsilon} - \frac{1}{\xi + x - i\epsilon}\right) dx .$$
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The generalized parton distributions and the nucleon

At leading twist there are 8 GPDs:

- 4 chiral-even GPDs: H, E, H̃ and Ẽ.
 4 chiral-odd GPDs: H_T, E_T, H̃_T and Ẽ_T.

Using the GPDs, we can determine the total angular momentum of quarks in the nucleon.

$$\int_{-1}^{1} x \left[H^{f}(x,\xi,0) + E^{f}(x,\xi,0) \right] dx = J^{f} \qquad \forall \xi$$

By Fourier transform of the GPD H at $\xi=0$ (need extrapolation), we obtain the distribution in the transverse plane of the partons as a function of their longitudinal momentum.



Photon electroproduction

We use leptons beam to generate the γ^{\ast} in the initial state... not without consequences.

Indeed, experimentally we measure the cross section of the process $ep \rightarrow ep\gamma$ and not strictly $\gamma^* p \rightarrow \gamma p$.



Photon electroproduction and GPDs

The interference term allows to access the phase of the DVCS amplitude, *i.e* allows to isolate imaginary and real parts of CFFs.

$$\begin{split} c^{DVCS}_{0,UU} &\sim \quad 4(1-x_B) \left(\mathcal{H}\mathcal{H}^* + \widetilde{\mathcal{H}}\widetilde{\mathcal{H}}^* \right) \,, \\ c^{\mathfrak{I}}_{1,UU} &\sim \quad F_1 \; \text{Re}\mathcal{H} + \xi(F_1 + F_2) \; \text{Re}\widetilde{\mathcal{H}} \,, \\ s^{\mathfrak{I}}_{1,LU} &\sim \quad F_1 \; \text{Im}\mathcal{H} + \xi(F_1 + F_2) \; \text{Im}\widetilde{\mathcal{H}} \,, \\ s^{\mathfrak{I}}_{1,LU} &\sim \quad F_1 \; \text{Im}\widetilde{\mathcal{H}} + \xi(F_1 + F_2) \; \text{Im}\widetilde{\mathcal{H}} \,, \end{split}$$

For these observables, at leading-order and leading-twist:

• there is only $c_{0,UU}^{DVCS}$ in $DVCS^2$.

• there are only
$$c_{1,UU}^{\mathcal{I}}$$
 and $s_{1,UL}^{\mathcal{I}}$ for \mathcal{I} .



Figure: Unpolarized (top) and beam-helicity dependent cross sections at Q²=2.3 GeV², $x_B=0.36$ and t=-0.3 GeV².

What have we collected so far?

Getting the GPDs: Cover the entire phase space with as many observables as possible!



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CLAS12: A unique place to study GPDs

The 12-GeV upgrade of the accelerator opens an access to a terra incognita in Q^2/x_B .

- CLAS to CLAS12 upgrade to reach higher luminosity and to explore the terra incognita!
- Disentangle the different GPD contributions with :
 - \rightarrow Polarization of beam and targets for the different channels. RG-A/K (LU/Rosenbluth) , RG-C (Long. pol. target), RG-H (Trans. pol. target).
 - \rightarrow Neutron data:

RG-B (See Silvia Niccolai's talk).

→ Have lepton/anti-lepton beam Proposal by Voutier, Niccolai *et al.*.

The most complete GPD experimental program ever run!

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Focus on CLAS12: Beam from right to left.



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A typical DVCS event in CLAS12

All CLAS12 detectors are necessary to reconstruct all particles of a DVCS final state:

- The electron is going through Cerenkov detector, drift chambers and electromagnetic calorimeter.
- The photon is either detected in a sampling calorimeter or a small PbWO₄-calorimeter close to the beamline.
- The recoil proton goes in the Silicium and Micromegas detector.



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DVCS event selection

Exclusivity is enforced by cutting on 5 variables:



However some exclusive π^0 's contaminate your DVCS sample: $\exists \cdot \exists \cdot \neg \land \bigcirc$ CLAS collaboration (Maxime DEFURNE)pDVCS @ CLAS12June 23rd 202012 / 20

Run group K: Unpolarized proton @ 7.5 GeV

After exclusivity cuts, here are the momentum/angle distribution.



- Unpolarized protons with 7.5 GeV longitudinally polarized electron beam.
- π^0 subtraction not performed on this BSA.
- Stay tuned for next DNP results with pass-1.





Run group A: Unpolarized proton @ 10.6 GeV

After exclusivity cuts, here are the momentum/angle distribution.



Figure 2: Momentum (GeV) vs θ angle (°) for electrons, photons and protons after selection and exclusivity cuts

Some exclusive π^0 's contaminate your DVCS sample.



Run group A: Unpolarized proton @ 10.6 GeV

The π^{0} -contamination has been studied and subtracted:

- BH/DVCS statistics is high for low Q^2 and low x_B .
- π^0 contamination grows larger with x_B .





Figure 6: Missing mass squared $ep \rightarrow ep\gamma X$ for bins with $-t/Q^2 < 0.25$ and estimation of π^0 contamination

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Courtesy G. CHRISTIAENS.

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Run group A: Unpolarized proton @ 10.6 GeV

The asymmetry is fitted with the following functionnal form:

$$\mathcal{A}^{DVCS}(\phi_{Trento}) = \frac{1}{P_{el}} \frac{N^+/Q^+ - N^-/Q^-}{N^+/Q^+ + N^-/Q^-} = \frac{p_0 \sin \phi_{Trento}}{1 + p_1 \cos \phi_{Trento}} \qquad P_{el} \sim 85\%$$
(1)

Courtesy G. CHRISTIAENS.





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Conclusion

- The results in this presentation were shown at last DNP 2019.
- Pass-1 for RG-A data provides a much larger statistics... on-going analysis.
- π^0 subtraction procedure validated.
- Refining binning, computation of bin center and radiative corrections remain to be performed.
- CLAS12 may (will) carry out the most complete DVCS program ever.
- A great thanks to the PhD students leading the analysis effort:
 - \rightarrow Guillaume CHRISTIAENS (RGA/BSA),
 - \rightarrow Joshua ARTEM TAN (RGK),
 - \rightarrow Katheryne PRICE (RGB),
 - \rightarrow Sangbaek LEE (RGA/Cross sections),

Back-Up

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Many channels, many observables provided by different facilites and each of them holds a specific piece of the puzzle.

Need to work hand-by-hand with phenomenologists and theorists .

We will need to develop global analysis tools in order to:

- combine all data and thus strongly constrain fits or models.
- test systematically the impact of diverse assumptions:
 - LO, NLO, NNLO,...
 - the numbers of flavours,
 - the numbers of GPDs,

I want to use this opportunity to make some advertisement for a project of major importance.

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GPD Computing made simple.

Differential studies: physical models and numerical methods.

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Phenomenology of Generalized Parton Distributions

Full processes Experimental data and phenomenology

Small distance Computation of amplitudes

Large distance First principles and fundamental parameters



- Many observables.
- Kinematic reach.
- Perturbative approximations.
- Physical models.

Fits.

- Numerical methods.
- Accuracy and speed.

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H. Moutarde Recent results

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What have we learned so far with data collected before 2016?

Several phenomenological fits of Compton Form factors has been performed at leading-twist:

- Assuming that colliders are sensitive to gluons and fixed-target experiments to quarks.
- From either "empirical" models or Neural networks.

First sketch of proton profile is drawn, especially in the valence region. Top: Kumericki *et al.*, JHEP07 (2011), 073 Bottom: Moutarde *et al.*, Eur.Phys.J.C78(2018)no.11,890



A glimpse of gluons through DVCS





- \rightarrow First data set at fixed kinematics but mutiple beam-energy.
- \rightarrow First phenomenological analysis including kinematical power corrections.
- NLO: Gluon transversity GPDs.
- HT: Q-G-Q correlations.

M. Defurne *et al.*, Hall A collaboration, arXiv:1703.0944

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