



STRONG2020 Hadron Spectroscopy (HaSP) General Workshop

MesonEx : Meson Spectroscopy with CLAS12

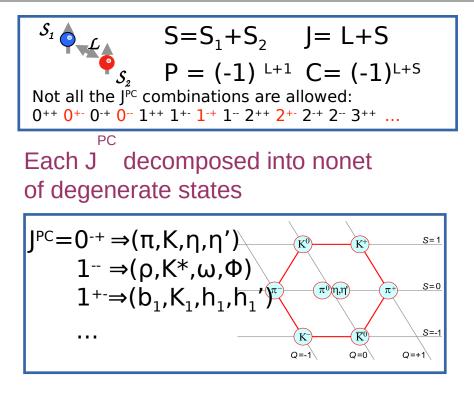
Derek Glazier University of Glasgow

Task 3: Meson Spectroscopy analysis of new and exotic states

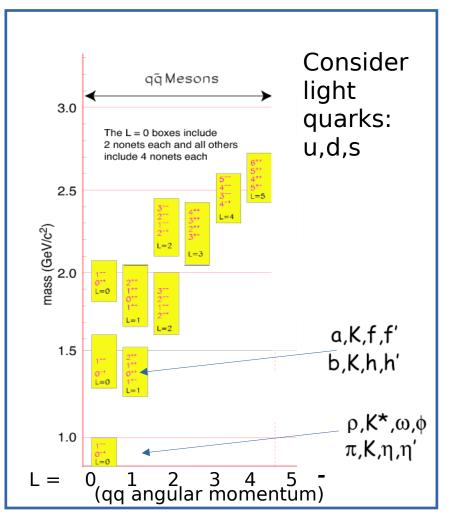
 Task 3.1: Search for and study of light exotic mesons, charmonium and strangeonium Coordinator: Vincent Mathieu



Light Quark Meson Spectroscopy



- Quark model explains much of the observed states
- Some states not well established or not definitively assigned
- Some additional unassigned states
- Particular ?? with 0++ scalars





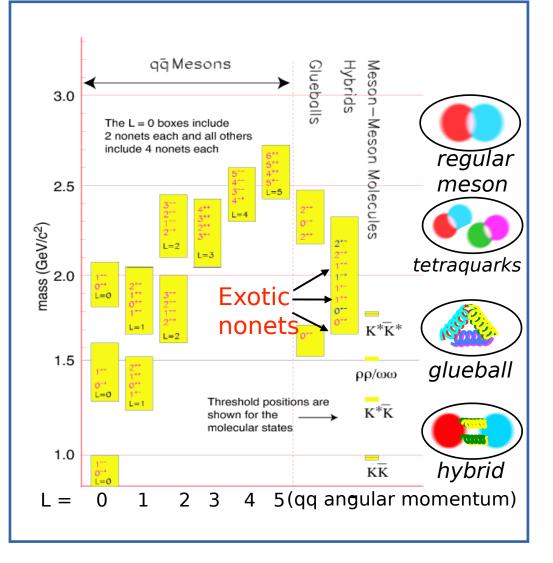
What else might exist

QCD does not forbid other compositions

Some states predicted by theory

Evidence for some states found experimentally

Hybrids include states of exotic quantum numbers



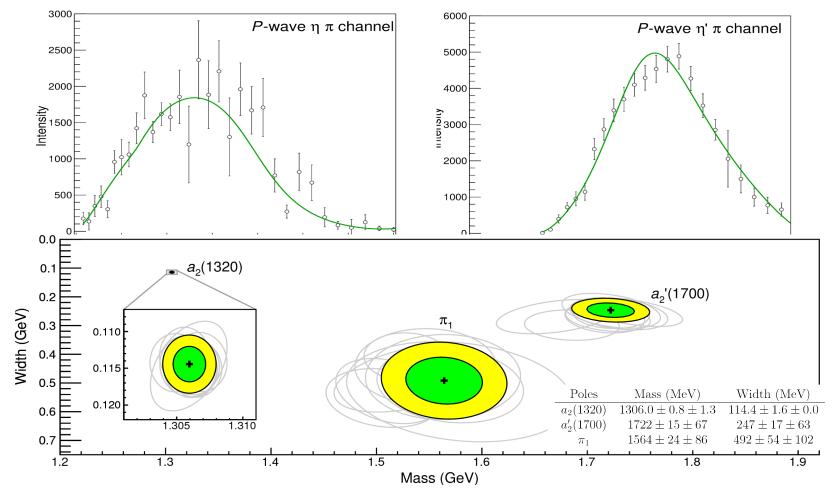


$π_1$ (1600) from COMPASS η(')π

Determination of the Pole Position of the Lightest Hybrid Meson Candidate

A. Rodas, A. Pilloni, M. Albaladejo, C. Fernández-Ramírez, A. Jackura, V. Mathieu, M. Mikhasenko, J. Nys, V. Pauk, B. Ketzer, and A. P. Szczepaniak (Joint Physics Analysis Center) Phys. Rev. Lett. **122**, 042002 – Published 29 January 2019

Coubled analysis of $\eta\pi$ and $\eta'\pi$



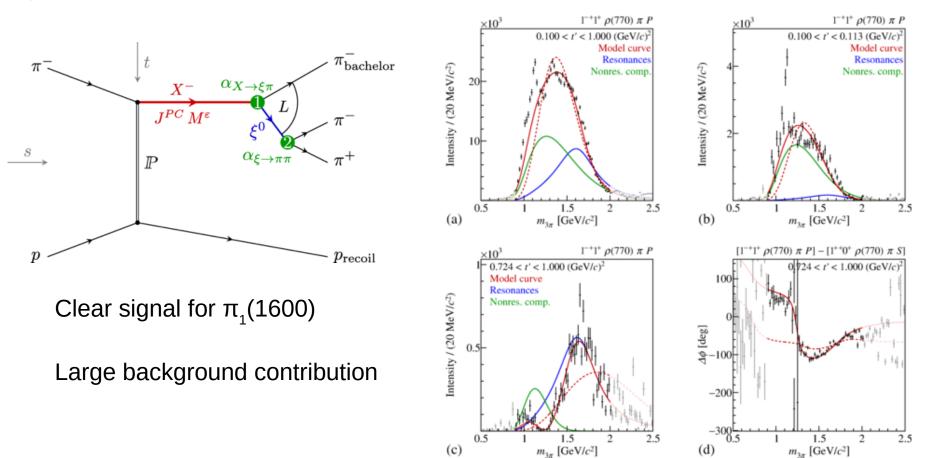


$\pi_1(1600)$ from COMPASS 3π

Light isovector resonances in $\pi^-p o \pi^-\pi^-\pi^+p$ at $190~{ m GeV}\,/c$

M. Aghasyan et al.

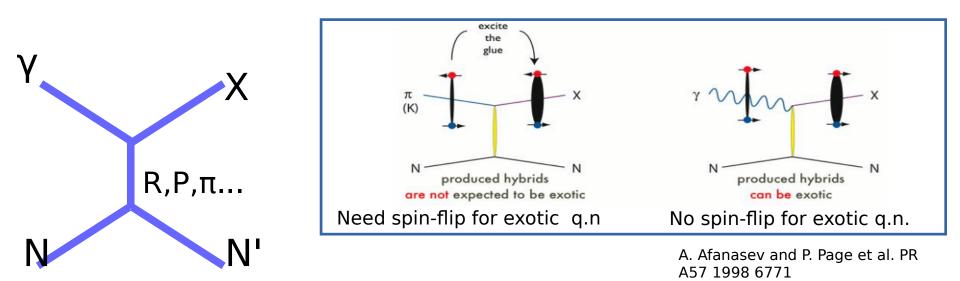
Phys. Rev. D 98, 092003 - Published 2 November 2018





Why photoproduction ?

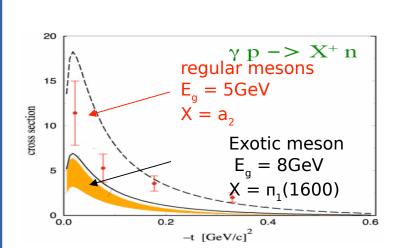
Exotic J^{PC} are more likely produced by S=1 probe



Linear polarization acts like a filter to disentangle the production mechanisms

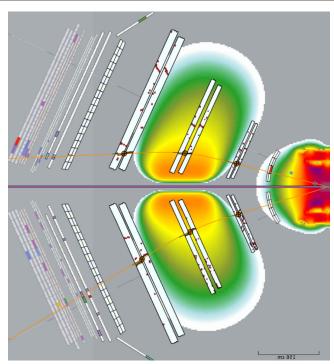
Production rate for exotics is expected comparable as for regular mesons

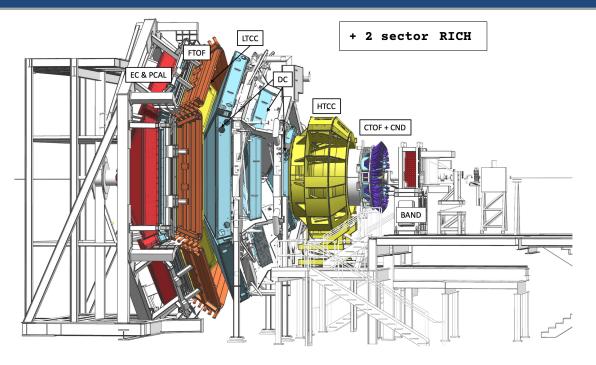
A. Szczepaniak and M. Swat PLB 516 2001 72





CLAS12 at Jefferson Lab





High luminosity electron scattering $(10^{35} \text{ cm}^{-2}\text{s}^{-1})$ produces high flux of nearly real photons.

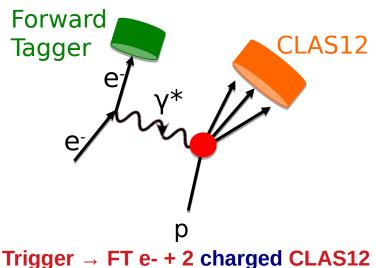
High resolution tracking spectrometer, (1% momentum, 1 mrad angle) Excellent PID e-, K, p, $\pi,~n,\gamma$

Can make measurements with missing particles

Can run MesonEx simultaneously with other experiments

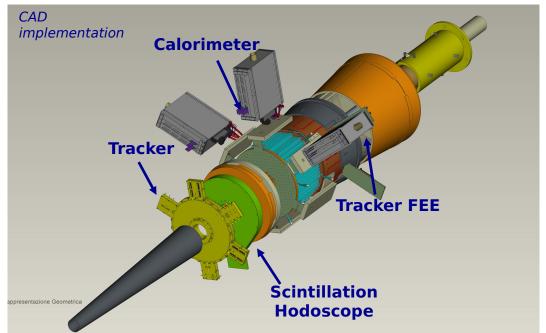


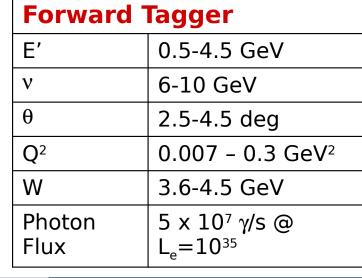
Quasi-real photoproduction with CLAS12



Quasi-real photoproduction:

- Detection of multiparticle final state from meson decay in the large acceptance spectrometer CLAS
 - Detection of the scattered electron for the tagging of the quasi-real photon in the CLAS12 FT
 - High-intensity and high-polarization tagged "photon" beam; degree of polarization determined event-by-event from the electron kinematics



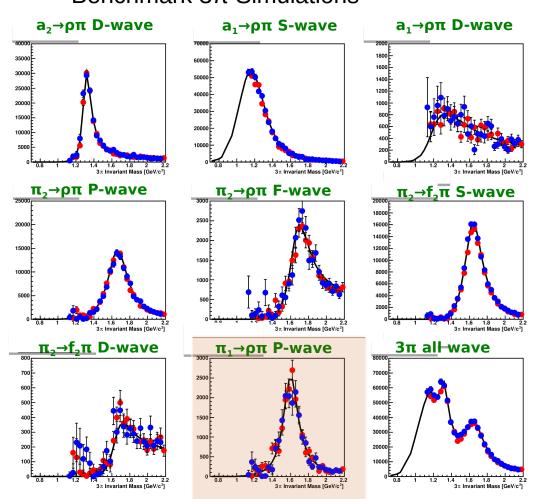




MesonEx program

Meson spectroscopy in the light-quark sector:

- Detailed mapping of the meson spectrum up to masses of 2.5 GeV
- Search for rare or poorly known states (strangeness-rich, scalars, ...)
- Search states with unconventional quarkgluon configurations



Benchmark 3π Simulations

Note poor acceptance below 1.3GeV (half field)



Approximately 20% of expected data ready for analysis

Focus on charged decay products (better resolution)

First extract two pseudoscalar ($\pi^+\pi^-$, K⁺K⁻)

Fourier Analyse angular distributions == extract moments

- more general expansion than just partial waves
- check acceptance corrections
- check distortions from backgrounds
- model independent formalism

Extract partial waves from moments or directly fit partial waves

Expand to vector-pseudoscalar final states



MesonEx : Polarised 2 meson production

arXiv.org > hep-ph > arXiv:1906.04841

Search... Help | Advand

High Energy Physics - Phenomenology

Moments of angular distribution and beam asymmetries in $\eta \pi^0$ photoproduction at GlueX

V. Mathieu, M. Albaladejo, C. Fernández-Ramírez, A. W. Jackura, M. Mikhasenko, A. Pilloni, A. P. Szczepaniak (JPAC collaboration)

(Submitted on 11 Jun 2019)

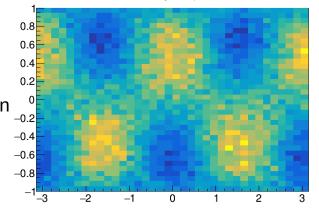
$$I(\Omega, \Phi) = I^{0}(\Omega) - P_{\gamma}I^{1}(\Omega)\cos(2\Phi) - P_{\gamma}I^{2}(\Omega)\sin(2\Phi)$$
$$I^{0}(\Omega) = \sum_{L}\sum_{M=0}^{M \leq L} \sqrt{\left(\frac{2L+1}{4\pi}\right)} (2 - \delta_{M,0})H^{0}(L, M) \Re\left[Y_{L}^{M}(\Omega)\right]$$

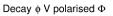
$$I^{1}(\Omega) = -\sum_{L} \sum_{M=0}^{M \leq L} \sqrt{\left(\frac{2L+1}{4\pi}\right)} (2 - \delta_{M,0}) H^{1}(L, M) \Re \left[Y_{L}^{M}(\Omega)\right]$$
$$I^{2}(\Omega) = 2\sum_{L} \sum_{M=0}^{M \leq L} \sqrt{\left(\frac{2L+1}{4\pi}\right)} \Im \left[H^{2}(L, M)\right] \Im \left[Y_{L}^{M}(\Omega)\right]$$

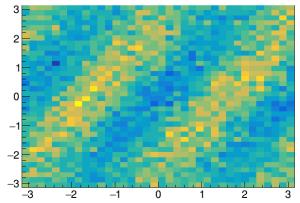
Moments relate directly to partial wave amplitues

$$\begin{split} H^0(11) &= H^1(11) + 2\sqrt{\frac{2}{5}}\operatorname{Re}(P_1^{(+)}D_2^{(+)*}) \ , \\ H^1(11) &= \frac{2}{15} \left[3\sqrt{5}\operatorname{Re}(P_0^{(+)}D_1^{(+)*}) - \sqrt{15}\operatorname{Re}(P_1^{(+)}D_0^{(+)*}) + 5\sqrt{3}\operatorname{Re}(S_0^{(+)}P_1^{(+)*}) \right] \ , \\ H^0(20) &= H^1(20) - \frac{2}{35} \left[7|P_1^{(+)}|^2 - 5|D_1^{(+)}|^2 + 10|D_2^{(+)}|^2 \right] \ , \\ H^1(20) &= \frac{4}{35} \left[7|P_0^{(+)}|^2 + 5|D_0^{(+)}|^2 + 7\sqrt{5}\operatorname{Re}(S_0^{(+)}D_0^{(+)*}) \right] \ , \\ H^0(21) &= H^1(21) + \frac{2}{7}\sqrt{6}\operatorname{Re}(D_1^{(+)}D_2^{(+)*}) \ , \\ H^1(21) &= \frac{2}{35} \left[7\sqrt{5}\operatorname{Re}(S_0^{(+)}D_1^{(+)*}) + 7\sqrt{3}\operatorname{Re}(P_0^{(+)}P_1^{(+)*}) + 5\operatorname{Re}(D_0^{(+)}D_1^{(+)*}) \right] \ , \end{split}$$

Decay θ V φ



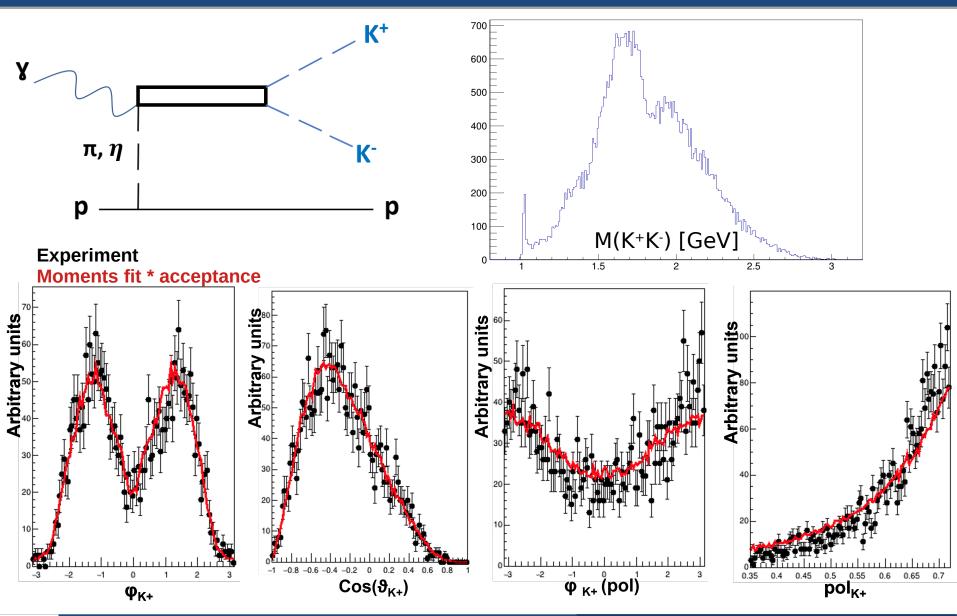




Moments can be determined from Fourier analysis of decay angles



MesonEx : K⁺K⁻ p Preliminary Data



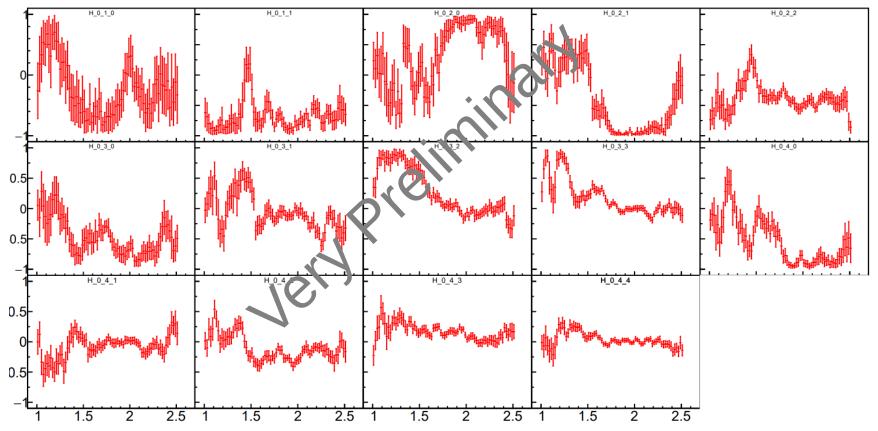
From Matthew Nicol, University of York



MesonEx : K⁺K⁻ p Preliminary Data

Allowed partial waves $J^{PC}(I^G)$: $0^{++}(0^+,1^-)$, $1^{--}(0^+,1^-)$, $2^{++}(0^+,1^-)$, $3^{--}(0^+,1^-)$, ...





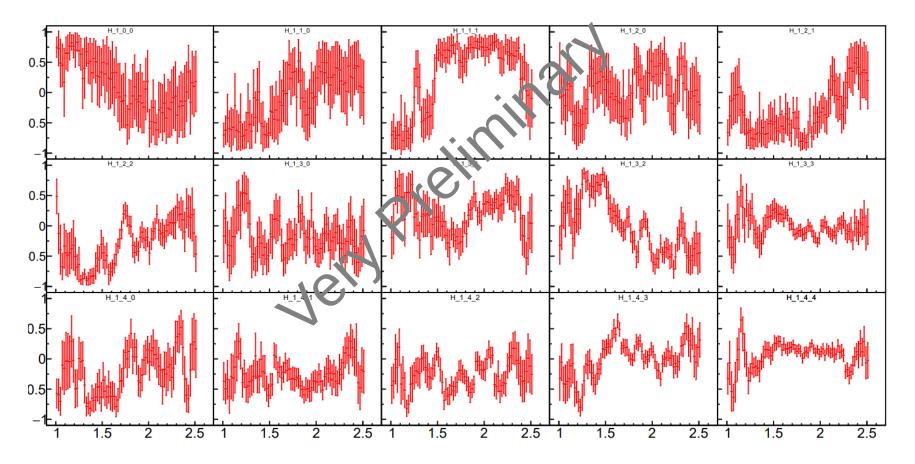


From Matthew Nicol, University of York

MesonEx : K⁺K⁻ p Preliminary Data

Allowed partial waves $J^{PC}(I^G)$: $0^{++}(0^+,1^-)$, $1^{--}(0^+,1^-)$, $2^{++}(0^+,1^-)$, $3^{--}(0^+,1^-)$, ...

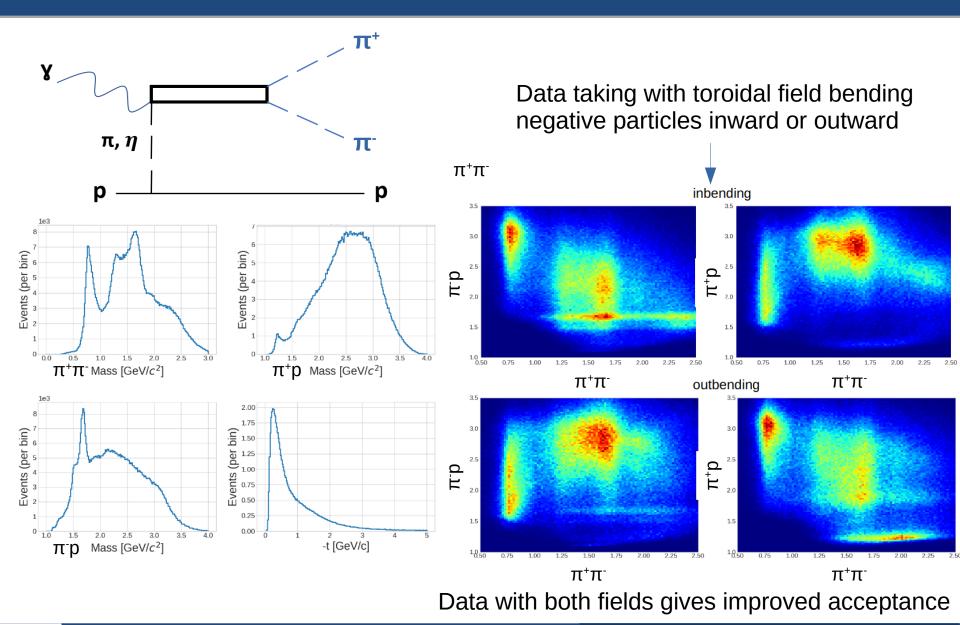
Polarised Moments





From Matthew Nicol, University of York

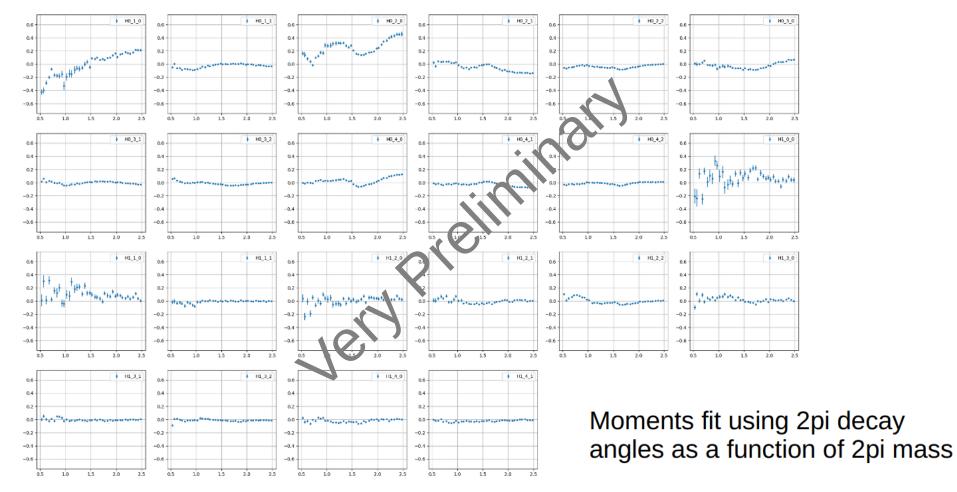
MesonEx Trigger : $\pi^+\pi^-p$ Final State



From Adam Thornton, University of Glasgow

MesonEx : $\pi^+\pi^-$ p Preliminary Moments

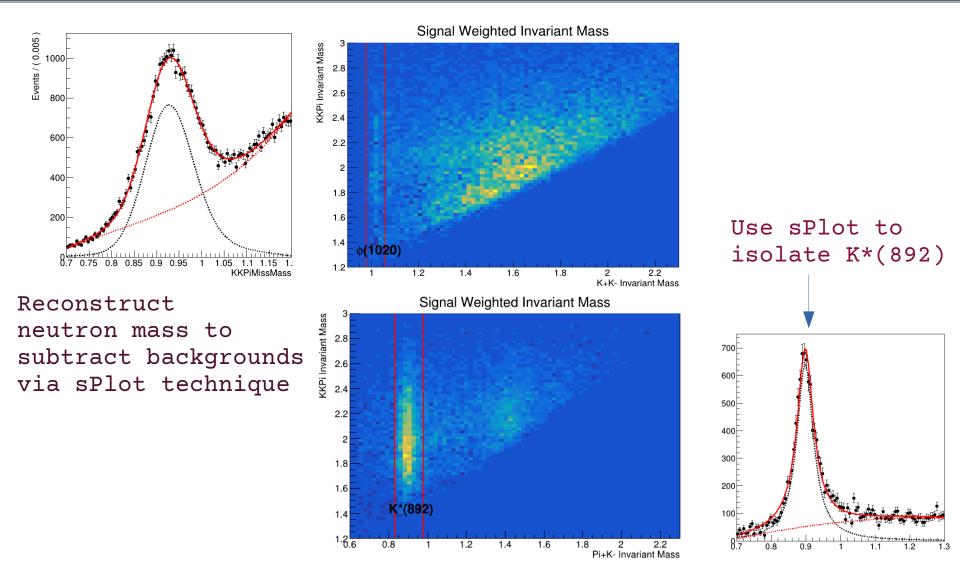
Allowed partial waves $J^{PC}(I^G)$: : $0^{++}(0^+)$, $1^{--}(1^+)$, $2^{++}(0^+)$, $3^{--}(1^+)$, ...





From Adam Thornton, University of Glasgow

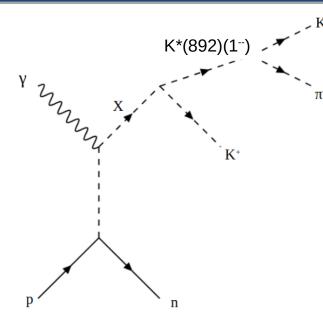
MesonEx : $\pi^+K^+K^-$ n Preliminary Data



class

From Robert Wishart, University of Glasgow

Vector-Scalar Formalism



Write intensity in terms of decay angles and linear polarisation $\mathcal{I}(\Omega, \Phi) = \mathcal{I}_0(\Omega) - \mathcal{P}_{\gamma} . \mathcal{I}_1(\Omega) \cos 2\Phi - \mathcal{P}_{\gamma} . \mathcal{I}_2(\Omega) \sin 2\Phi.$

Expand polarised intensity components in moments of D-Wigner function products (depend on 4 decay angles)

$$\mathcal{I}^{\alpha}(\Omega) = \sum_{JMS\Lambda} \left(\frac{2J+1}{4\pi}\right) \left(\frac{2S+1}{4\pi}\right) H^{\alpha}(JMS\Lambda) \\ \times D^{J*}_{M,\Lambda}(\phi_{GJ}, \theta_{GJ}, 0) D^{S*}_{\Lambda,0}(\phi_{HF}, \theta_{HF}, 0).$$

 $\begin{array}{l} \text{Relate moments to production Spin Density Matrix Elements} \\ b = \{j = J_X, l = L_X, m = M_X, s = S_X\} \\ = 1 \\ \\ \sum_{b,b'} \sum_{\lambda,\lambda'} \left(\frac{\sqrt{(2l'+1)(2l+1)}}{2j'+1} \right) \left(\sqrt{\frac{2s+1}{2s'+1}} \right) (l,0;s,\lambda|j,\lambda) (l',0;s',\lambda'|j',\lambda') \\ \\ \times (s,\lambda;S,\Lambda|s',\lambda') (s,0;S,0|s',0) (j,m;J,M|j',m') (j,\lambda;J,\Lambda|j',\lambda') \\ \\ \times \rho_{bb'}^{\alpha} \times R_{Yi}(m_Y) R_{Yi'}^{*}(m_Y), \end{array}$

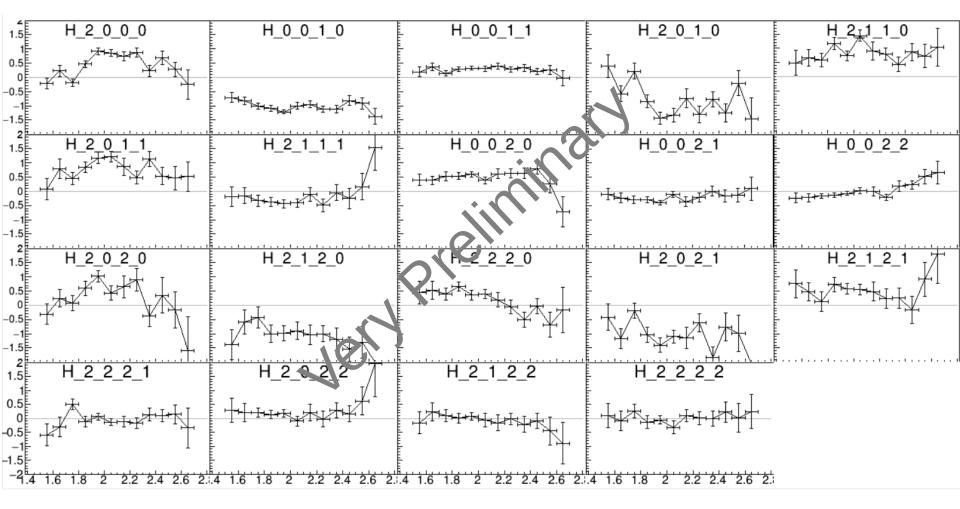
Relate Spin Density Matrix Elements to Partial Waves dependent On meson X spin and parity





MesonEx : $\pi^+K^+K^-$ n Preliminary Data

Vector-Scalar Moments fit results (unpolarised only)





From Robert Wishart, University of Glasgow

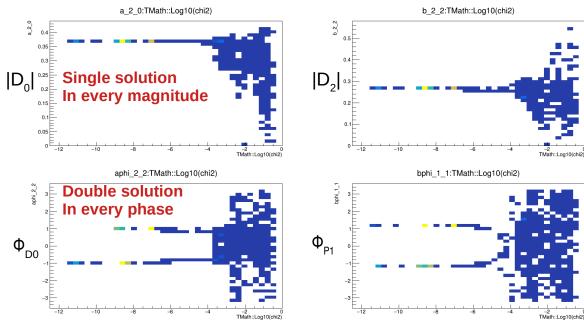
Moments to Partial Waves

$$\begin{split} H^{0}(LM) &= \sum_{\substack{\ell\ell' \\ mm'}} \left(\frac{2\ell'+1}{2\ell+1} \right)^{1/2} C^{\ell 0}_{\ell' 0 L 0} C^{\ell m}_{\ell'm' LM} \rho^{\alpha,\ell\ell'}_{mm'}, \end{split}$$

$$\begin{aligned} &\stackrel{(\epsilon)}{=} \rho^{0,\ell\ell'}_{mm'} &= \kappa \sum_{k} \left([\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} \right), \\ &\stackrel{(\epsilon)}{=} \rho^{1,\ell\ell'}_{mm'} &= -\epsilon \kappa \sum_{k} \left((-1)^{m} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{m';k} + (-1)^{m'} [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{-m';k} \right), \\ &\stackrel{(\epsilon)}{=} \rho^{2,\ell\ell'}_{mm'} &= -i\epsilon \kappa \sum_{k} \left((-1)^{m} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{m';k} - (-1)^{m'} [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{-m';k} \right), \\ &\stackrel{(\epsilon)}{=} \rho^{3,\ell\ell'}_{mm'} &= \kappa \sum_{k} \left([\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{m';k} - (-1)^{m'} [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{-m';k} \right), \\ &\stackrel{(\epsilon)}{=} \rho^{3,\ell\ell'}_{mm'} &= \kappa \sum_{k} \left([\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{m';k} - (-1)^{m'} [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{-m';k} \right). \end{aligned}$$

The moments can be expanded in partial waves via the SDMEs (p) => a system of simultaneous equations

=> if sufficient moments can numerically extract the partial waves (e.g. with Minuit)



Need to run many times!

For linearly polarised photoproduction find we can invert equations with 2 (trivial) complex conjugate solution sets of partial waves.

i.e. there are 2 solutions,Or 1 if 1 partial wave is forced+ve real part



CLAS12 has successfully taken data for ~ 4 years

Several analyses on beam spin asymmetries already published

Data with forward tagger allow large rate of meson photoproduction events

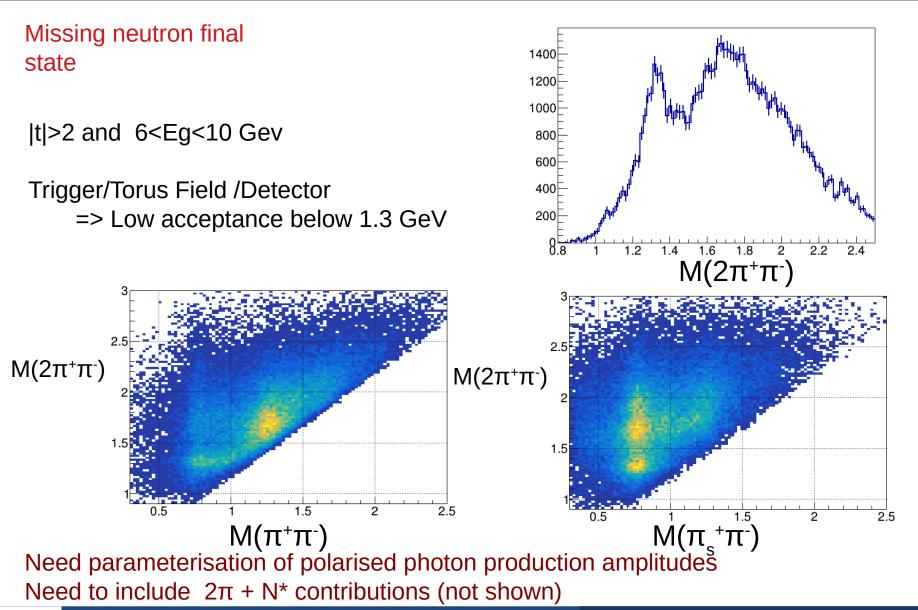
Currently analysing several two-body final states

Working on extracting moments of angular distributions then partial waves

Extending moments and partial waves formalisms as part of Strong2020 work package



MesonEx : $\pi^+\pi^+\pi^-$ n Preliminary Data





MesonEx : $\pi^+\pi^- p$ Preliminary Data

