

Transverse Target Asymmetry in **Exclusive** Charged Pion Production at 12 GeV

Dipangkar Dutta

Mississippi State University

(with Dave Gaskell & Garth Huber)



Polarized Target Workshop:

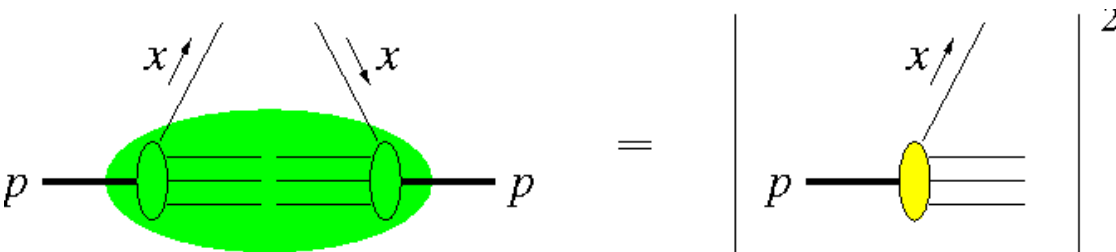
June 17-18, 2010

Outline

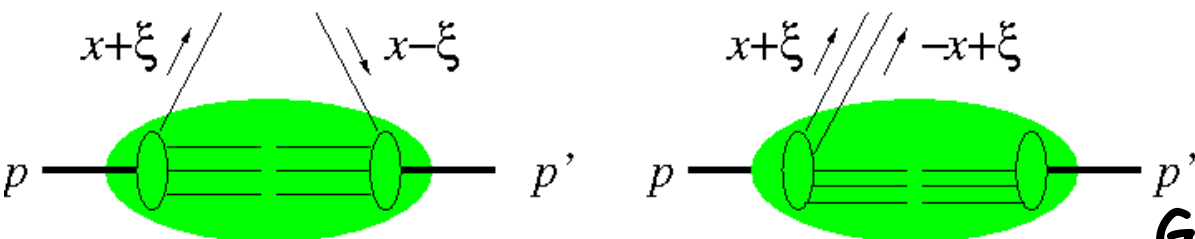
- Motivations
 - GPDs
 - Factorization Studies
- Transverse Target Asymmetry
- Possible Measurements
- Rate Estimates
- Summary

Generalized Parton Distribution

Over the last decade tremendous progress has been made on the theory of generalized parton distributions (GPD).



PDFs : Squared hadronic wavefunctions
= probability of finding a parton with specified longitudinal momentum fraction and polarization in a fast moving hadron

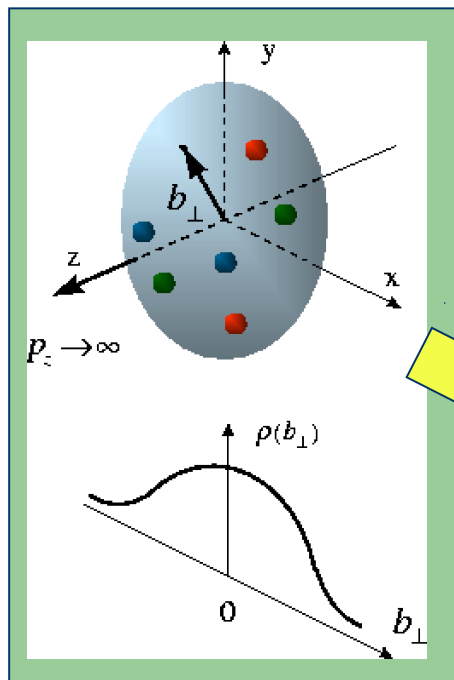


GPDs : interference between wavefns. of parton with momentum fraction $x+\xi$ and parton with momentum fraction $x-\xi$

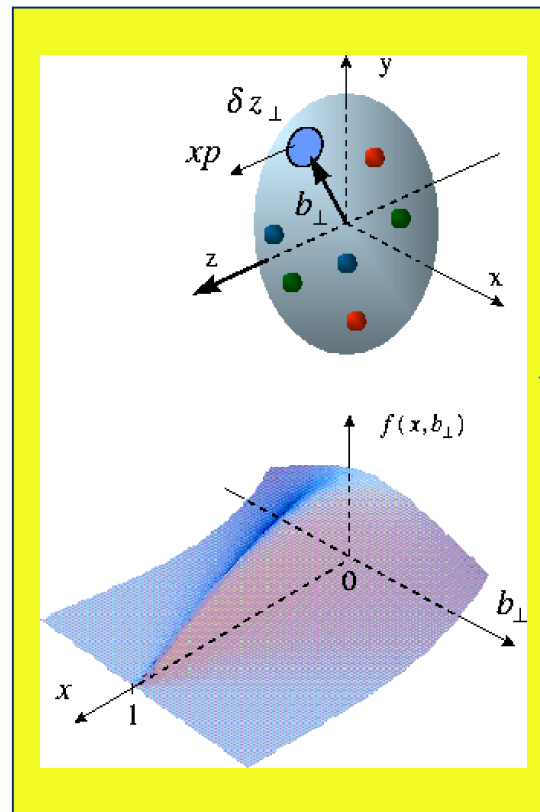
- GPDs may give us access to:
 - Orbital angular momentum in the nucleon
 - 3D pictures of the nucleon

Generalized Parton Distribution

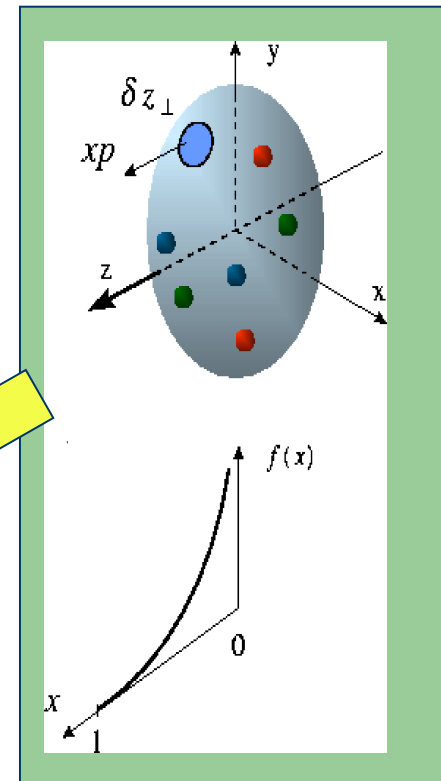
GPDs allow simultaneous determination of longitudinal momentum and transverse position of partons, they correlate different quark-gluon configurations in the hadron



Proton form factors,
transverse charge &
current densities



Correlated quark momentum
and helicity distributions in
transverse space - **GPDs**



Structure functions,
quark **longitudinal**
momentum & helicity
distributions

Leading Twist GPDs

- At leading twist, 4 independent GPDs for each quark, gluon type
- x is the light cone momentum fraction of struck parton ($x \neq x_{Bj}$)
- $t=\Delta^2$, momentum transfer to nucleon
- ξ , defined by

$$\Delta^+ = -2\xi(p+\Delta/2)^+$$

$H^{q,g}(x, \xi, t)$
 spin avg
 no hel. flip

$E^{q,g}(x, \xi, t)$
 spin avg
 helicity flip

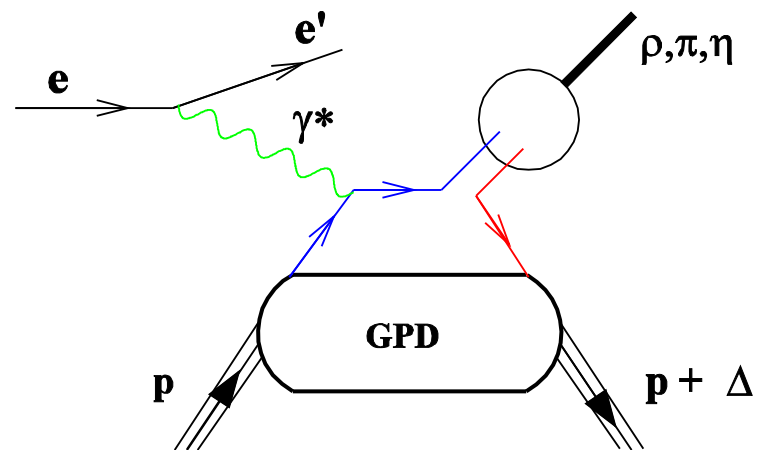
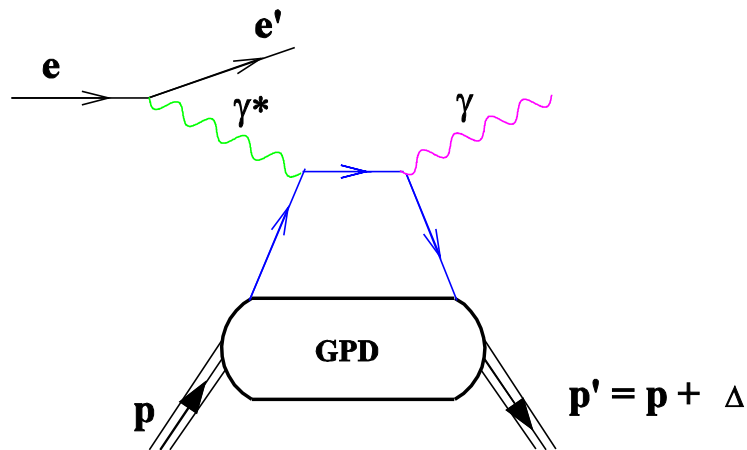
$\tilde{H}^{q,g}(x, \xi, t)$
 spin diff
 no hel. flip

$\tilde{E}^{q,g}(x, \xi, t)$
 spin diff
 helicity flip

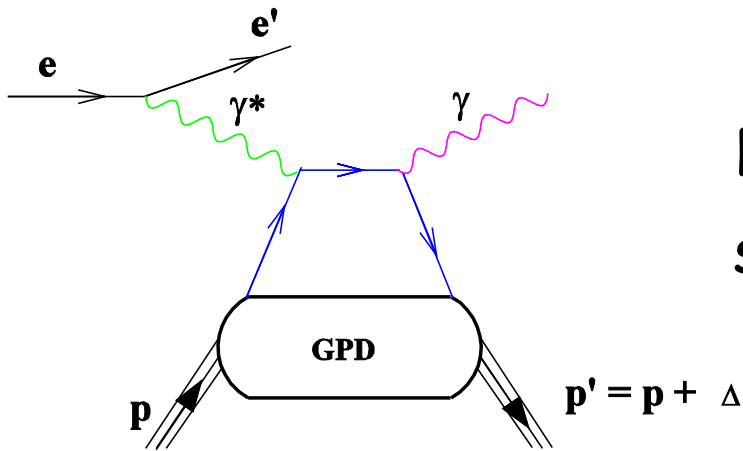
Parametrizes skewedness, or the Longitudinal fraction of the momentum transfer, t

Access to GPDs

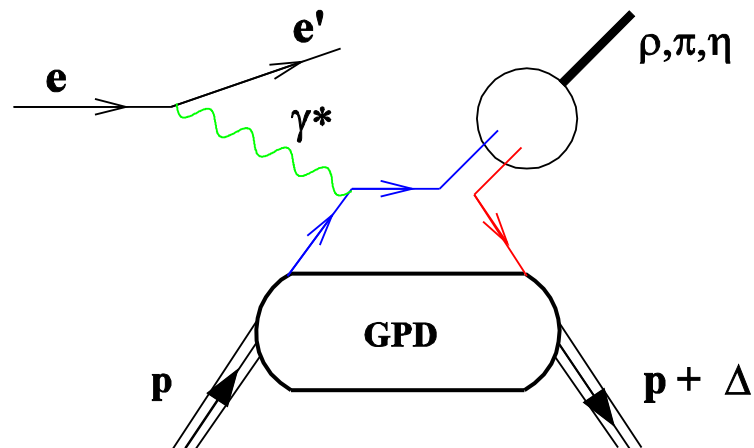
- GPDs are not observables – but, they are constrained by
 - DIS: $H(x,0,0) = q(x)$ and $\tilde{H}(x,0,0) = \Delta q(x)$ (PDFs, at vanishing t, ξ)
 - Elastic scattering: $\int dx x H(x, \xi, t) = F_1(t)$, (Form Factor, from model independent sum rules)
- To access GPDs, we need a program that will measure ...
 - a variety of **hard exclusive processes** (vector mesons, DVCS, pseudo-scalar mesons)
 - a broad range of phase space (t, x_B)



Access to GPDs



Deeply Virtual Compton Scattering
sensitive to a combination of all 4 GPDs



Vector mesons (ρ, ω, ϕ) sensitive to H and E
Pseudo scalar mesons sensitive to \tilde{H} and \tilde{E}

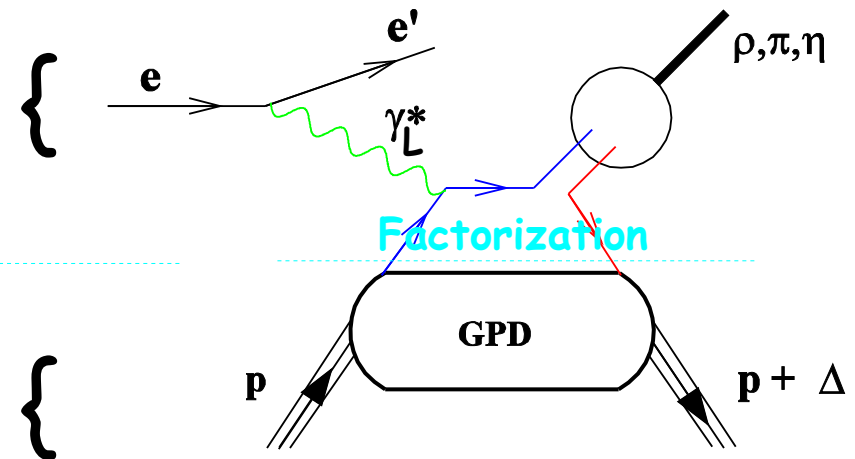
$$\tilde{E}^{ud}(x, \xi, t) = F_\pi(t) \frac{\theta(\xi > |x|)}{2\xi} \phi_\pi\left(\frac{x + \xi}{2\xi}\right),$$

\tilde{E} via hard pion electroproduction can provide new information on nucleon structure which is available from any other source.

- An extensive program of Hard Exclusive Measurements to constrain GPDs is a major part of the **12 GeV** upgrade.

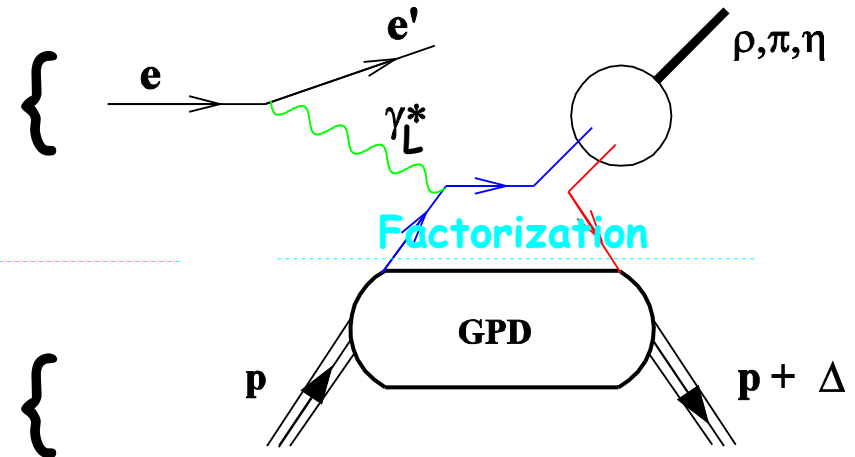
Factorization in Hard Exclusive Reactions

- Hard probe creates a small size configuration, its interactions can be described by pQCD (applies to longitudinally polarized virtual photon only)
- How the probe is transformed into the hadron etc can only be described non-perturbatively (parametrized by GPDs)



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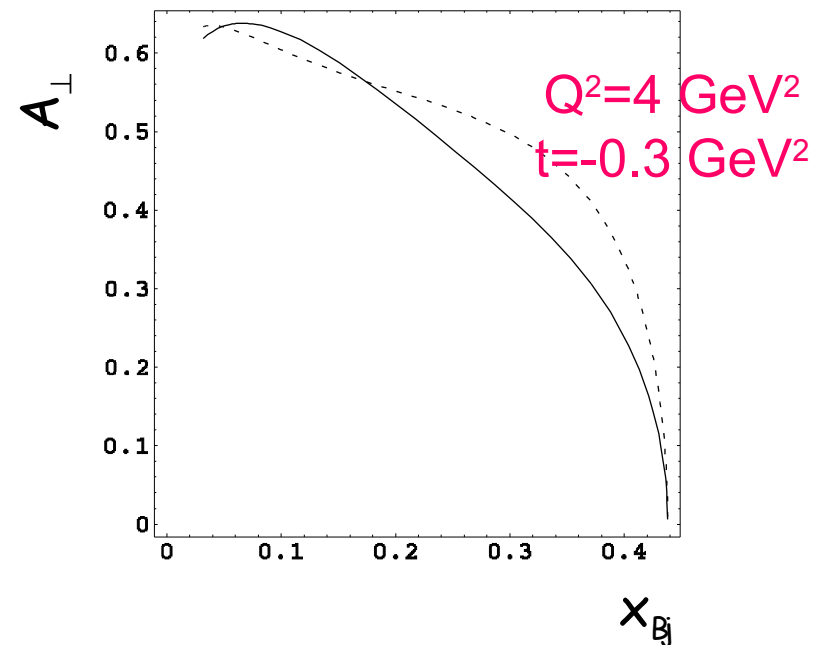
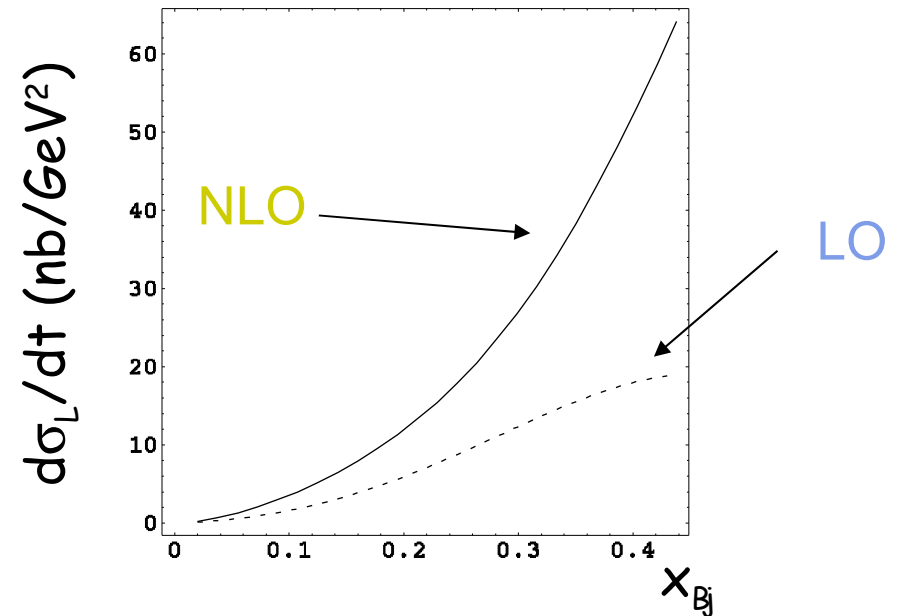


- How the probe is transformed into the hadron etc can only be described non-perturbatively (parametrized by GPDs)

- To say anything about GPDs, we must be confident we are in a regime where soft-hard factorization applies (large Q^2)
- Higher order corrections may be large for absolute cross sections for $Q^2 < 10 \text{ GeV}^2$
- Ratios have a better chance of exhibiting precocious factorization – higher order effects in numerator and denominator “cancel”

Exclusive π^+ Production at NLO

- Belitsky and Müller GPD based calc. of π^+ production to NLO (Phys Lett B 513, 349)
 - Even at $Q^2=10 \text{ GeV}^2$, NLO effects can be large, but cancel in the asymmetry, A_\perp
 - At $Q^2=4$, higher twist effects even larger in σ_\perp , but still cancel in asymmetry (CIPANP 2003)
- This cancellation of higher order effects known as precocious factorization



Exclusive π^+ Production with Target (or Recoil) Polarization

Unpolarized
Cross section

$$\frac{d\sigma}{d\Omega} = \sigma_T + \epsilon\sigma_L + \sqrt{\frac{1}{2}\epsilon(\epsilon+1)}\sigma_{LT}\cos\phi + \epsilon\sigma_{TT}\cos 2\phi$$

Exclusive π^+ Production with Target (or Recoil) Polarization

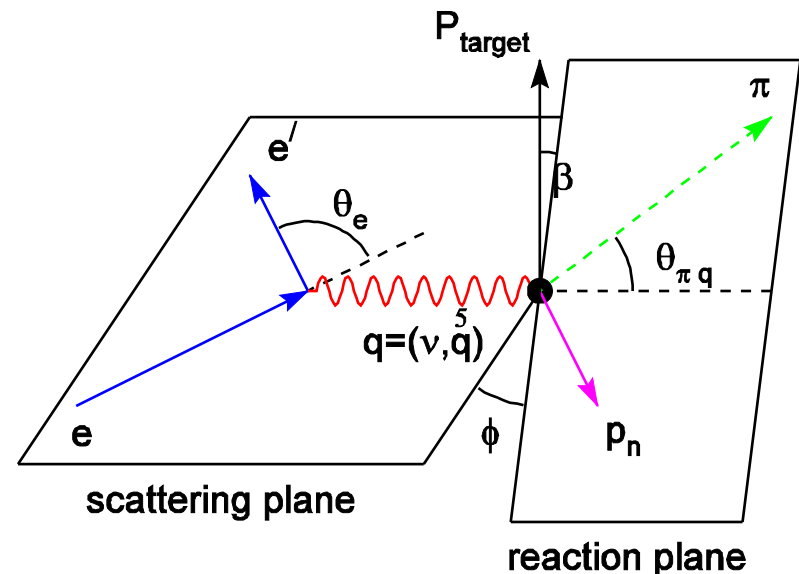
Unpolarized
Cross section

$$\frac{d\sigma}{d\Omega} = \sigma_T + \epsilon\sigma_L + \sqrt{\frac{1}{2}\epsilon(\epsilon+1)}\sigma_{LT}\cos\phi + \epsilon\sigma_{TT}\cos 2\phi$$

Polarized
cross section has
additional components

$$\begin{aligned} \sigma_t = & P_x \left[-\sqrt{2\epsilon(1+\epsilon)}\sin\phi\sigma_{LT}^x - \epsilon\sin 2\phi\sigma_{TT}^x \right] \\ & - P_y \left[\sigma_{TT}^y + \epsilon\cos 2\phi\sigma_{TT'}^y + 2\epsilon\sigma_L^y + \sqrt{2\epsilon(1+\epsilon)}\cos\phi\sigma_{LT}^y \right] \\ & + P_z \left[\epsilon\sin 2\phi\sigma_{TT}^z + \sqrt{2\epsilon(1+\epsilon)}\sin\phi\sigma_{LT}^z \right] \end{aligned}$$

- Target polarization components (P_x , P_y) are defined relative to the reaction plane
- β = azimuthal angle between (transverse) target polarization and reaction plane
 $P_x = P_\perp \cos\beta$ and $P_y = P_\perp \sin\beta$



π^+ Transverse Target Asymmetry

- Setting all transverse amplitudes to zero, the pion electroproduction cross section (with polarized target) is:

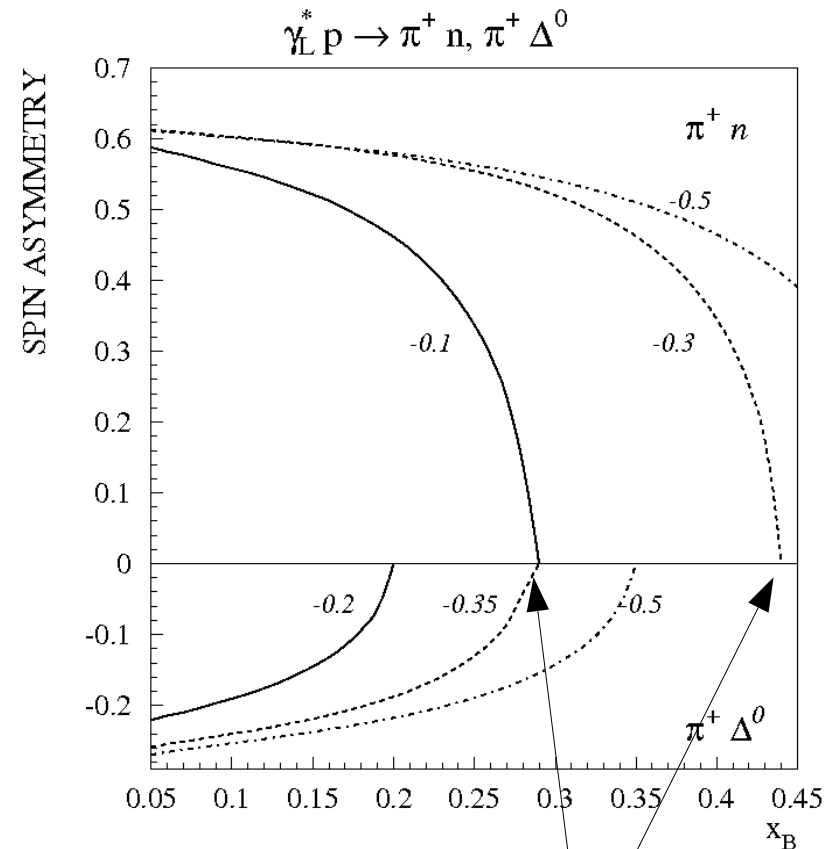
$$\sigma = \varepsilon \sigma_L - 2 \varepsilon \sigma_L^y P_\perp \sin \beta \quad (P_y = P_\perp \sin \beta)$$

- The transverse target asymmetry is typically defined [Frankfurt et al, PRD 60, 014010 (1999)]

$$A_\perp = \frac{\int_0^\pi d\beta \frac{d\sigma_L^{\pi^-}}{d\beta} - \int_\pi^{2\pi} d\beta \frac{d\sigma_L^{\pi^-}}{d\beta}}{\int_0^{2\pi} d\beta \frac{d\sigma_L^{\pi^-}}{d\beta}}$$

- The transverse target asymmetry then involves the ratio of two longitudinal cross sections

$$A_\perp = \frac{2}{\pi} \frac{2\sigma_L^y}{\sigma_L}$$



$A_T = 0$ at parallel kinematic limit, where P_y is not well defined.
PRL84, 2589 (2000)

Measurement of A_{\perp}

- At JLab energies, cannot ignore contributions from transverse photons (σ_T suppressed by $1/Q^2$ compared to σ_L)

Two alternatives to get "separated asymmetry"

- Two Rosenbluth separations and ratio of longitudinal cross sections

$$\sigma_A = \sigma_T^{\perp} + \epsilon \sigma_L^{\perp} \quad \text{where } \sigma(\epsilon) = \sigma_U + \sigma_A \sin\beta + \dots$$

$$\sigma_U = \sigma_T + \epsilon \sigma_L$$

$$A_{\perp} = \frac{2}{\pi} \frac{2\sigma_L^y}{\sigma_L}$$

- Rosenbluth separation of the asymmetry

$$A = A_T + \epsilon A_L$$

Measurement of A_{\perp}

- To cleanly extract A_{\perp} , we need
 - target polarized transverse to virtual photon direction
 - Large acceptance in π^+ azimuthal angle (i.e. ϕ and β)
 - Measurements at multiple beam energies and electron scattering angles $\rightarrow \epsilon$ dependence (L/T separation)
(advantage of focusing spectrometers in Hall-C)

for $r = \sigma_T/\sigma_L$, uncertainty in longitudinal cross sections is

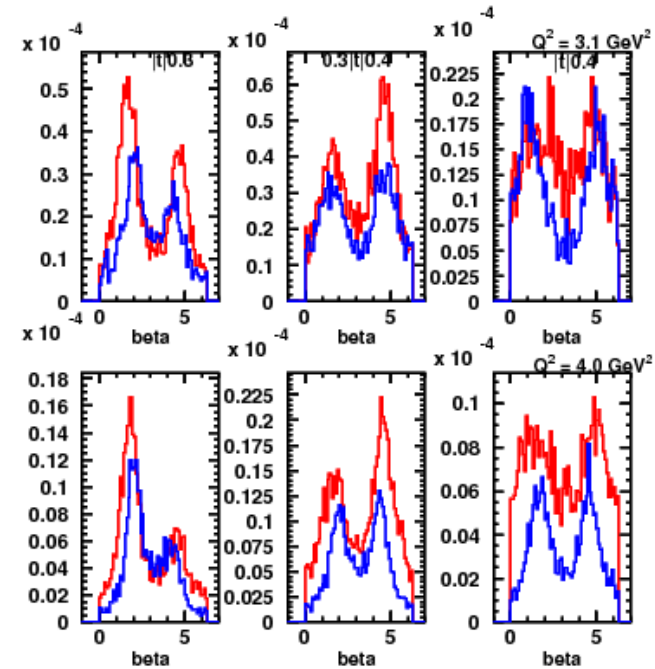
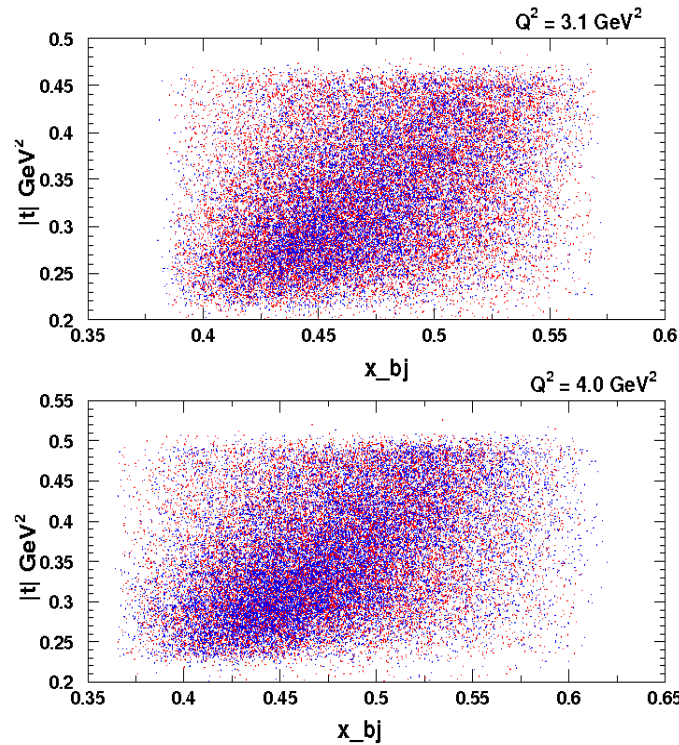
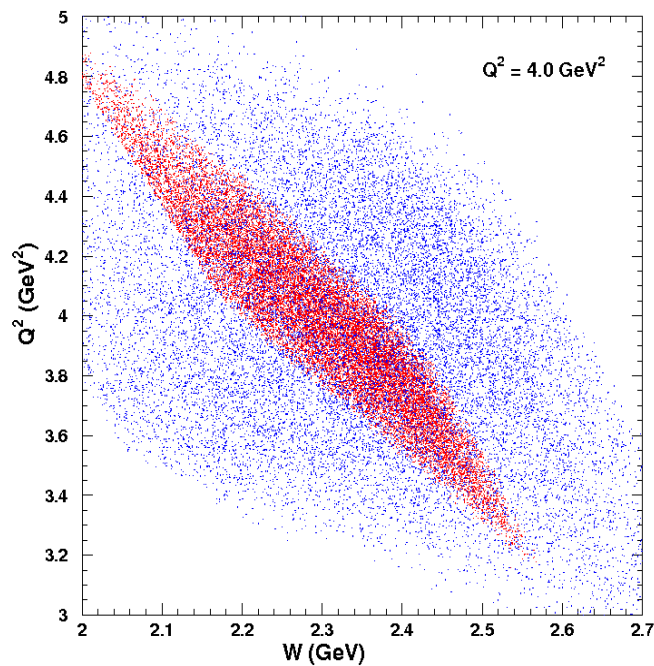
$$\Delta\sigma_L/\sigma_L = 1/(\epsilon_1 - \epsilon_2) \Delta\sigma/\sigma [(r + \epsilon_1)^2 + (r + \epsilon_2)^2]^{1/2}$$

Ratio of two L/T separations will have larger uncertainties than separated asymmetry method.

- Need $\Delta\epsilon$ as large as possible

HMS-SHMS Possible Kinematics

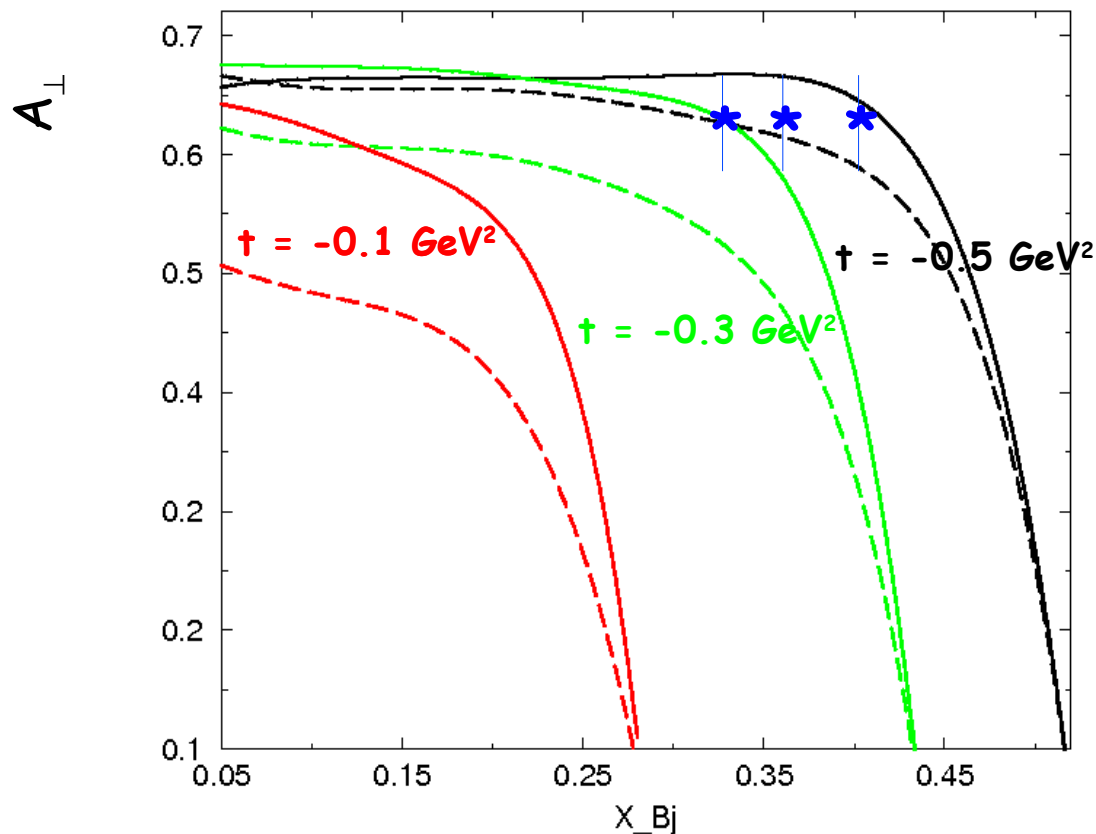
$n(e, e' \pi^-)p$ or $p(e, e' \pi^+)n$



high ε low ε $\Delta\varepsilon = 0.3$

Estimated Rates

Assuming a luminosity of $5 \times 10^{37} / \text{cm}^2/\text{s}$, using SIMC and the VGL model for the value of r and A_{\perp} etc, we estimate that a $\sim 6\%$ (stat) measurement of the transverse asymmetry can be performed in 500 hours.
(this rules out any measurement with the UVa NH_3 target.)



Solid: asymptotic pion
distribution amp.
Dashed: CZ pion dist. amp.

Issues to be Studied

- Compare figure of merit for recoil polarization measurement (with LD_2 target) vs high luminosity transverse pol. ^3He target...
- Compare backgrounds for different pol. target options
- Explore large acceptance devices BigCal, SBB etc

Summary

JLab 12 GeV upgrade will open the doors to a rich program of “measuring” GPDs

Deep Exclusive Pion Production (large Q^2 , low $-t$), will provide unique insight into certain class of GPDs and also on soft-hard factorization.

The transverse asymmetry in exclusive pion production is one of the important observables in these studies.

It is very important to separate the longitudinal cross-section and asymmetry.

High luminosity transversely polarized targets are critical for these type of experiments.