



BARYON CHIRAL DYNAMICS

Baryons 2002 @ JLab

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Feb. 2002

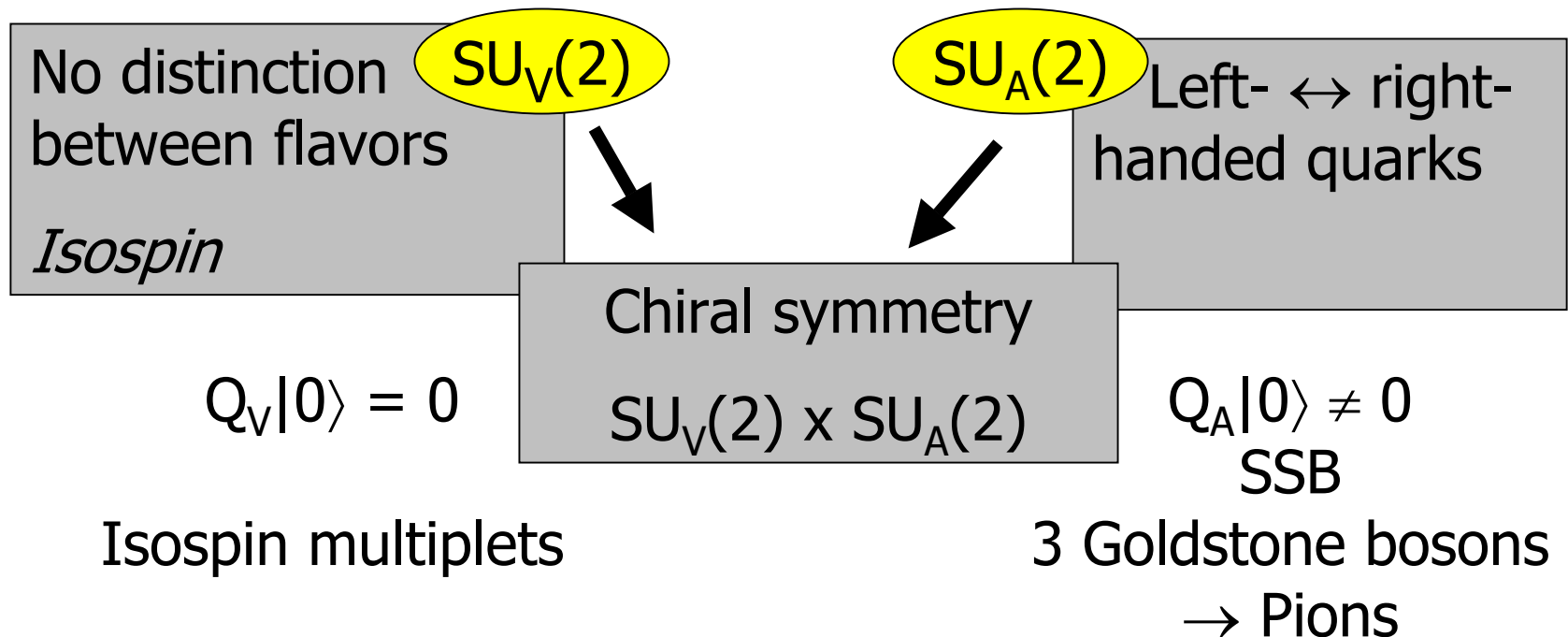


Overview

- Chiral dynamics ... with nucleons
- Higher, faster, stronger, ...
 - Formulation of the effective Theory
 - Full one loop results: $O(q^4)$ (Two loop result for m_N)
 - Isospin breaking, electromagnetism.
 - Pionic hydrogen
 - Two nucleon sector
 - Connection to lattice QCD: Quenched BCHPT
 - Photons: RCS, VCS, OMC, RMC,... →Merkel's talk
 - Higher energies: Resonances, dispersion relations
- ... but still puzzled: $g_{\pi N}$, σ -term, **SU(3)**

"QCD Lite" ... (QCD for $N=2$ massless flavors)

- One "parameter" $g_s \leftrightarrow \Lambda_{\text{QCD}}$
 - Parameter free predictions for dimensionless quantities
- High degree of symmetry: $SU_V(2) \times SU_A(2)$





Chiral Expansion

- Goldstone bosons decouple as $E_\pi \rightarrow 0$
 - Low energy singularities of the Green's functions from the propagation of π 's.
 - Account for those, expand amplitudes in external momenta.
- $\mathcal{L}_{\text{QCD}} = \mathcal{L}_{\text{Lite}} + \sum_i m_i \bar{q}_i q_i$
 - But m_u, m_d (and m_s ?) happen to be light
 - Treat mass term as a perturbation
- Simultaneous expansion in q, m_u, m_d



CHiral Perturbation Theory

- Chiral symmetry leads to relations between different amplitudes
- PCAC, Current Algebra
 - Work out structure of amplitudes and symmetry relations by hand. *Tedious!*
- CHPT: Use effective Lagrangian $\mathcal{L}_{\text{eff}}(\pi)$
 - most general \mathcal{L}_{eff} compatible with symmetries
 - QCD dynamics encoded in coupling constants.
 - Order \mathcal{L}_{eff} by the number of derivatives on the π -field and by powers of m_q



Goals

- Low energy QCD in terms of a small number of parameters
 - Connect different processes
 - Experiment vs. lattice simulations
- Extrapolation to QCD Lite
 - Predictions of the symmetry.
 - Mechanism of symmetry breaking?
 - E.g. nucleon mass in chiral limit?

Baryon CHPT

- Include nucleon
 - $\mathcal{L}_{\text{eff}} = \mathcal{L}_{\pi} + \mathcal{L}_N$
 - Lowest order

$$\mathcal{L}_N^{(1)} = -\frac{g_A}{2F_{\pi}} \bar{\psi} \gamma^{\mu} \gamma_5 (\partial_{\mu} \pi) \psi + \frac{1}{8F_{\pi}^2} \bar{\psi} \gamma^{\mu} i[\pi, (\partial_{\mu} \pi)] \psi + \dots$$

- $m_N \gg M_{\pi}$
 - $P_{\mu} \sim (m_N, 0, 0, 0)$ is $O(q^0)$
 - Interactions with soft pions: Nucleon remains nearly static, surrounded by cloud of π 's
- \mathcal{L}_N contains odd powers of derivatives

Low Energy Constants

- Pion sector

- Only even powers

$$\mathcal{L}_{\pi\pi} = \mathcal{L}_{\pi\pi}^{(2)} + \mathcal{L}_{\pi\pi}^{(4)} + \mathcal{L}_{\pi\pi}^{(6)}$$

(2) (7) (53)

- Nucleon sector

- Odd and even powers

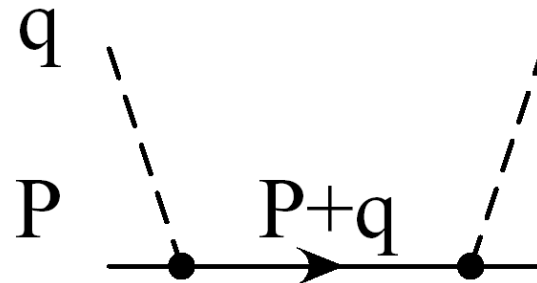
$$\mathcal{L}_{\pi N} = \mathcal{L}_{\pi N}^{(1)} + \mathcal{L}_{\pi N}^{(2)} + \mathcal{L}_{\pi N}^{(3)} + \mathcal{L}_{\pi N}^{(4)} + \dots$$

(1) (7) (23) (118)

Tree-level 1-loop 2-loop

Expansion of the N -Kinematics...

- Chiral expansion of the N -propagator



The diagram shows a horizontal solid line representing the N -propagator with an arrow pointing to the right. The momentum of this line is labeled $P+q$. Two dashed lines are attached to the ends of the solid line. The left dashed line is labeled q and the right dashed line is labeled P . The diagram is followed by an arrow pointing to the right, leading to the mathematical expression for the propagator.

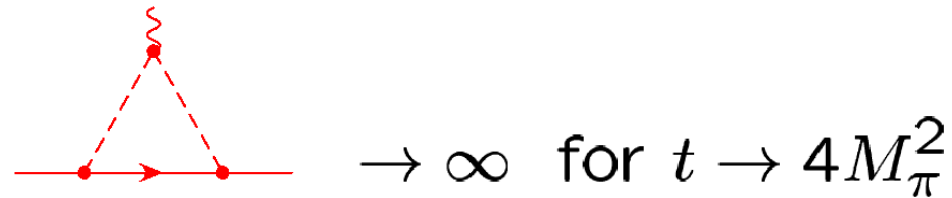
$$\rightarrow \frac{1}{(P+q)^2 - m_N^2}$$
$$= \frac{1}{2Pq} + O(q^0) \quad \text{for } P^2 = m_N^2$$

- Can be implemented into $\od�_{\text{eff}}$: **H**BCHPT



... is delicate ...

- Choice of kinematical variables
- Threshold singularities:
 - Breaks down at the threshold



- Slow convergence at threshold





... and can be avoided.

- Relativistic formulation
 - Dimensional regularization
 - Loop graphs are of the same order as tree level.
 - Problems to organize the perturbation series
 - Infrared regularization
 - Variant of dim. reg.
 - Well organized perturbation series
 - Avoids expansion of N -kinematics

Ellis & Tang; Becher & Leutwyler; Goity, Lehmann, Prezeau, Saez



Infrared Regularization

- Split dim. reg. loop graphs $L = R + I$
 - R : Large loop momentum $k \sim m_N$
 - Violates chiral counting
 - Trivial chiral expansion
 - I : Small loop momentum $k \sim M_\pi$
 - Contains all low energy singularities
- Absorb R into definition of \odot_{eff} and set $L = I$
 - Powercounting
 - Controlled expansion of kinematics
 - Manifest Lorentz invariance



$\pi N \rightarrow \pi N$ Scattering Amplitude

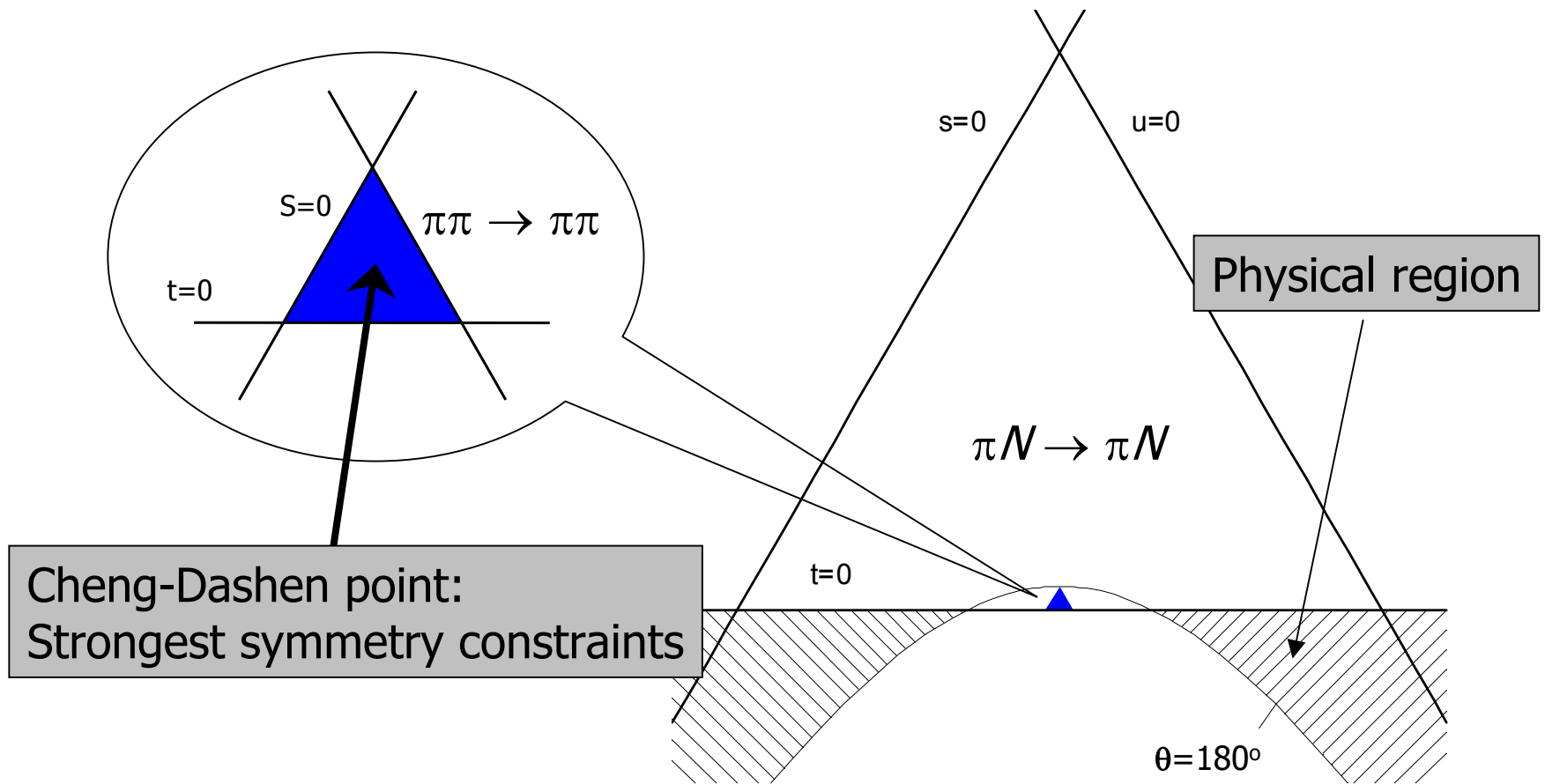
- $O(q^4)$ result in isospin limit
 - HBCHPT Meissner and Fettes
 - Infrared regularization Becher and Leutwyler
- Isospin breaking (strong and e.m.)
 - $O(q^3)$ Meissner and Fettes
- Inclusion of Δ -resonance
 - $O(\varepsilon^3)$ Ellis and Tang; Meissner and Fettes



$\pi\pi \rightarrow \pi\pi$ versus $\pi N \rightarrow \pi N$

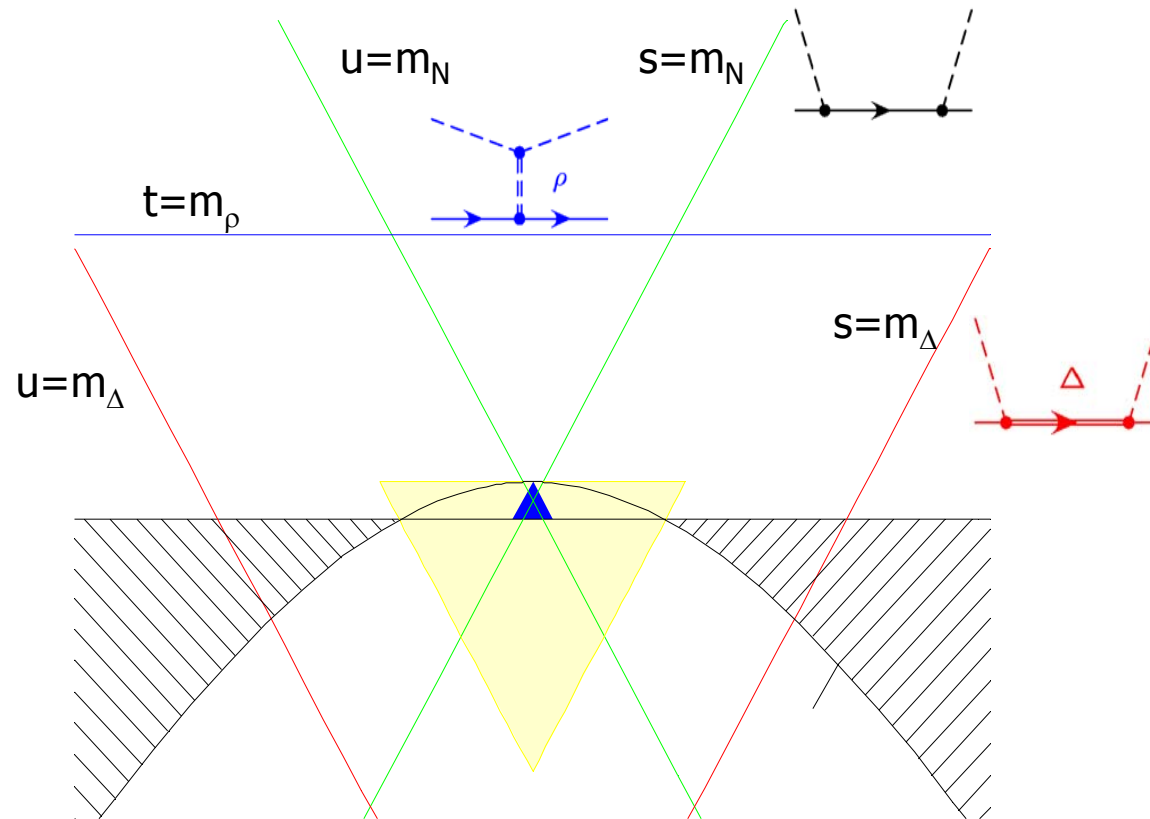
- *More LEC's* for πN -Scattering
 - 4 parameters for $\pi\pi \rightarrow \pi\pi$ to $O(q^4)$
 - 14 for $\pi N \rightarrow \pi N$ to $O(q^4)$
- But also *more data* !
- πN experimental region is at “higher” energies.

Mandelstam Triangle



Low Energy Region

- Low energy expansion breaks down, once resonances are produced!





Scattering amplitude at $O(q^4)$

- Simple parameterization

$$A(\nu, t) = B + P(\nu, t) + F(s) + F(u) + G(t), \quad \text{where} \quad \nu = (s - u) / 4m_N$$

- B : Nucleon pole terms
- P : Polynomial ← LEC's
- F : s -channel cut, linear in t .
- G : t -channel cut, linear in ν .
- Four (2 spin \times 2 isospin) amplitudes in terms of **9 functions of a *single variable***.
- Constraints from Chiral Symmetry
 - Goldberger-Treiman relation
 - Value of the amplitude at the CD-point

Goldberger Treiman Relation

- Relation between g_A and $g_{\pi N}$

$$g_{\pi N} = \frac{g_A m_N}{F_\pi} \{1 + \Delta_{GT}\}$$

- Δ_{GT} *vanishes* for $m_q = 0$. Note $M_\pi^2 \propto m_q$

$$\Delta_{GT} = c M_\pi^2 + O(M_\pi^4)$$

- No chiral logarithm, no M_π^3 term!
- Experimentally: $\Delta_{GT} = 2-4\%$

GWU

KH80



Isospin Odd Amplitude

Cheng Dashen point:

$$C = 2 F_\pi^2 \frac{\bar{D}^-(\nu, t)}{\nu} \Big|_{\nu=0, t=2M_\pi^2}$$

- $C=1$ for $m_q=0$

$$C = 1 - \frac{1 + 5 g_A^2}{24 \pi^2 F^2} M^2 \ln \frac{M}{\mu} + k_1 M^2 + k_2 M^3 + O(M^4)$$

- Contains chiral log: Expect O(10%) symmetry breaking
- "Experimentally" $C \sim 1.06-1.13$
- Constraint on subthreshold coefficients



Isospin Even Amplitude

- Scalar form factor \leftrightarrow Amplitude at CD

- Amplitude at CD: $\Sigma = F_\pi^2 \bar{D}^+(0, 2M_\pi^2)$

- And scalar form factor

$$\langle N' | m_u \bar{u}u + m_d \bar{d}d | N \rangle = \sigma(t) \bar{u}'u .$$

- Relation

$$\Sigma = \sigma(2M_\pi^2) + \Delta_{CD}$$

$$\Delta_{CD} = k_{CD} M^4 + O(M^5)$$

- No low E singularities in $\Delta_{CD} \approx 1\text{-}2\text{MeV}$.



σ -Term

- Quark mass dependence of the nucleon mass

$$\sigma \equiv \sigma(0) = m_u \frac{\partial m_N}{\partial m_u} + m_d \frac{\partial m_N}{\partial m_d}$$

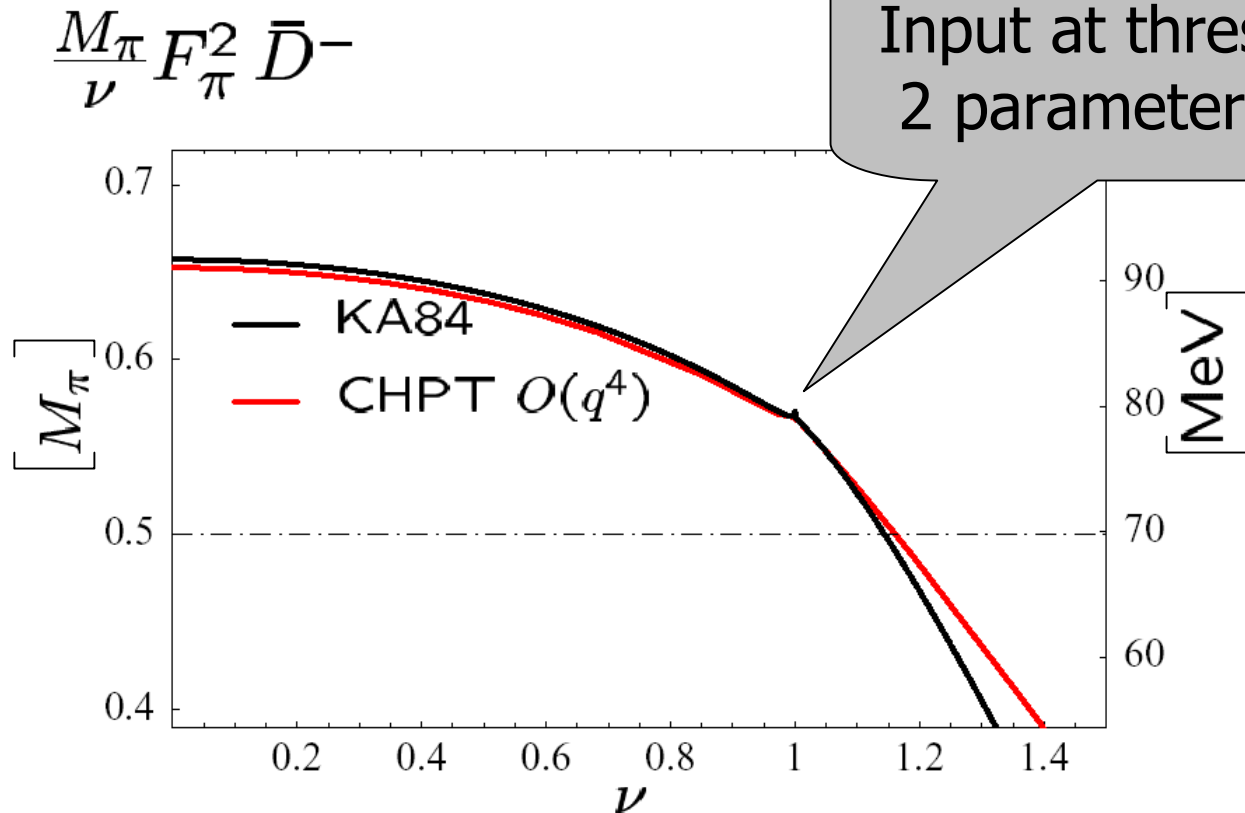
- Extrapolation $t = 2M_\pi^2 \rightarrow t=0$.

- CHPT confirms dispersive result

$$\sigma(2M_\pi^2) - \sigma = 15.2 \pm 0.4$$

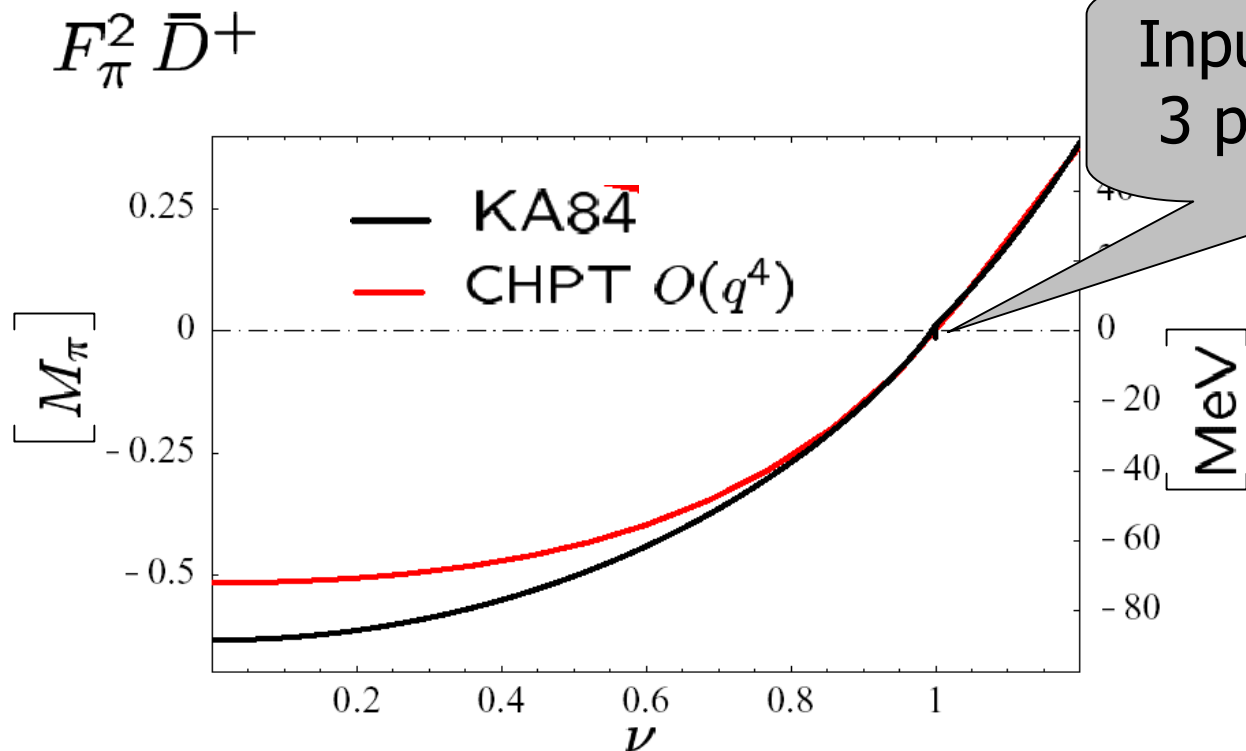
- Value of $\Sigma \rightarrow \sigma$

Amplitude for $t=0$



Starts to deviate soon above threshold...

Isospin Even Amplitude



Chiral representation is **not precise enough** to extrapolate experimental results to the CD region



Higher Energies

- Two issues
 - Unitarity is fulfilled only perturbatively in CHPT.
 - One loop imaginary parts are “tree level” squared and not very accurate.
 - Resonances
 - Δ -Resonance close to threshold and $g_{\Delta N\pi}$ is large.



Δ -Resonance

- Can be taken into account systematically
 - Jenkins & Manohar; Hemmert, Holstein & Kambor
 - Incorporate Δ into \odot_{eff} .
 - Count both q_π and $\Delta = m_\Delta - m_N$ as $O(\varepsilon)$.
- Scattering amplitude known to $O(\varepsilon^3)$.
 - Ellis & Tang, Fettes & Meissner
 - Pole term is dominant contribution.
 - Accuracy similar to $O(q^4)$ CHPT result
 - Unitarity remains an issue.
 - Nonrelativistic expansion of the Δ -propagator...



Implement Unitarity

- Unitarization
 - Unitarize model or CHPT result
 - E.g. Oller and Meissner
 - Good description of experimental data
 - Various unitarization prescriptions
 - Model dependent, distort structure of amplitude
- Dispersion relations
 - Use experimental imaginary parts, generate real part with dispersion relations
 - Integral equations



Dispersive approach

- Complicated! Function of *two* variables.
- Implement low energy structure found in CHPT (functions of *single* variable)
 - Leads to set of integral equations similar to the Roy-equations in $\pi\pi$ -scattering
 - Input:
 - Exp. data above elastic region
 - Four subtraction constants

Becher, Leutwyler; Mojzis; Stahov



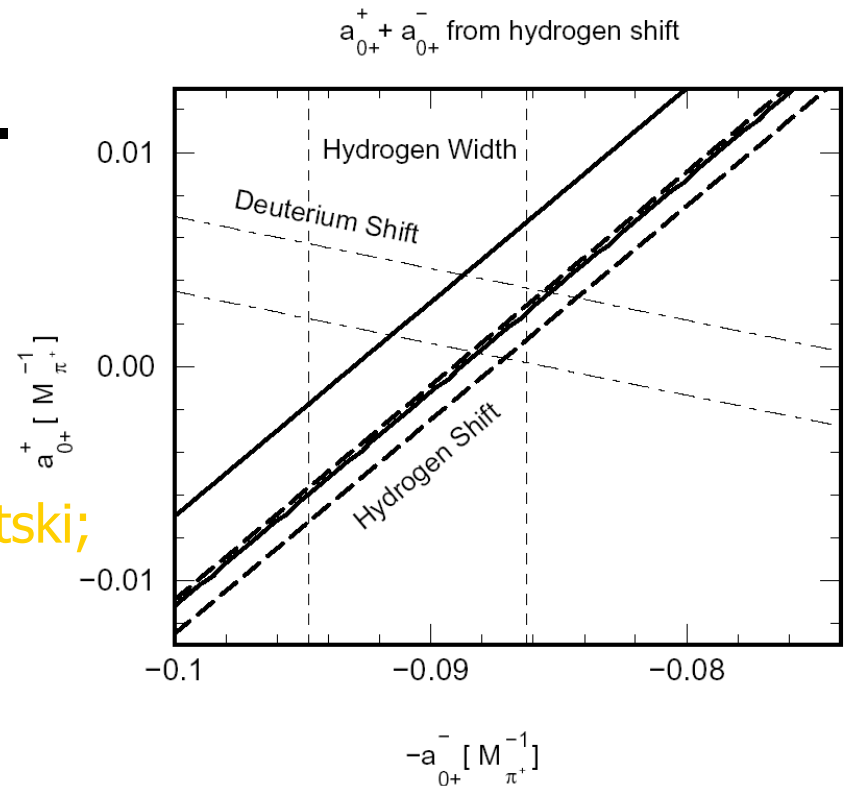
Isospin violation

- Two sources
 - Strong: Quark mass difference $m_u - m_d$
 - Electromagnetic
- Need to treat both on equal footing
 - Count e^2 as $O(q)$
 - $\pi N \rightarrow \pi N$ to $O(q^3)$ **Meissner and Fettes**
 - Isospin breaking small, mostly static
 - Dynamical effects only in S-wave $O(1\%)$.
 - Models incomplete

π^-p Bound State

- Scattering lengths from π^-p energy shift and decay width.
 - Theory developed.
 - Results for π^-p .
 - Level shift for π^-p .

Gasser, Lyubovitskij & Rusetski;
Eiras, Soto





Conclusion

- Have
 - Reached high precision
 - Studied wide range of processes
 - Good understanding of the role of chiral symmetry in the baryon sector.
- Chiral symmetry breaking is a *small effect*
 - Minimize model dependence
 - Push dispersive methods to connect with experimental results.
 - Resolve discrepancies in basic parameters