

Heavy quark physics on the lattice

BARYONS2002

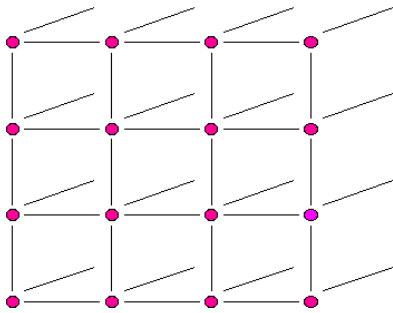
Christine Davies



UNIVERSITY
of
GLASGOW

Lattice QCD =

Euclidean space-time lattice



+ QCD Lagrangian (discretised)

$$\begin{aligned}\mathcal{L}_{QCD} &= \mathcal{L}_g + \mathcal{L}_q \\ &= \frac{1}{2g^2} \text{Tr} F_{\mu\nu}^2 + \bar{\psi}(\gamma \cdot D + m)\psi\end{aligned}$$

Parameters are those of QCD:

Bare gauge coupling $\beta = 6/g^2$

Quark masses, $m_i a$.

Lattice spacing (a) is implicit u.v. cutoff

Determine a and fix m_i from $1 + n_f$ hadron masses.

Difficulties

Systematic errors:

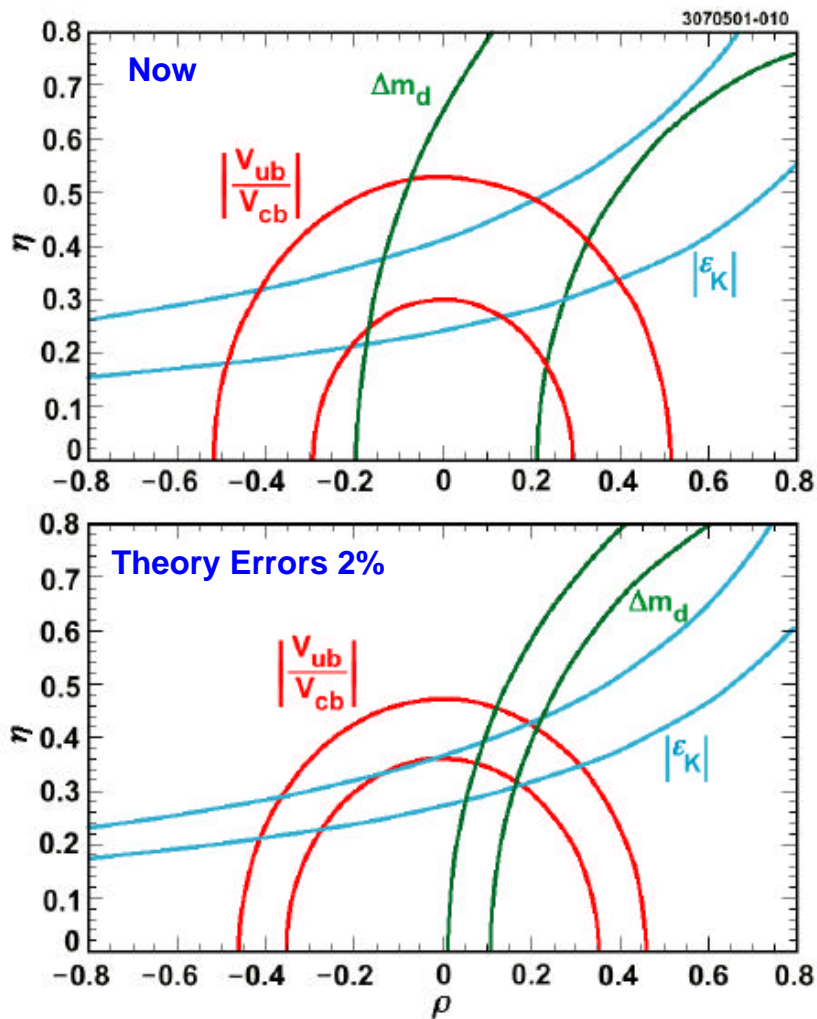
- Discretisation errors - physical results are dependent on a . Reduce the dependence and/or extrapolate to $a = 0$.
- Matching errors - must renormalise lattice matrix elements to obtain continuum results. Requires pert. or nonpert. matching calculation.
- Quenching errors - v. expensive to include light dynamical (sea) quarks. Error from using quenched approx. = 10-20% (?).

A short history of lattice QCD

- Invented 1974
slow progress
- Renaissance in 1990s
 - Much improved understanding of systematic errors
 - 10-20% errors possible on spectrum, form factors etc
- Second lattice revolution - early 2000s
 - Teraflop computing power will enable simulations with 'real' dynamical quarks
 - carry improvement of systematic errors further
 - 2-3% errors will be possible

Why do you care?

Improved theoretical precision will add huge value to experiment, e.g. from B factories.



(CLEO-c - new expts will check lattice errors.)

Heavy Quark Physics

Heavy (b and c) quarks ($m \gg \Lambda_{QCD}$) present special challenges to lattice QCD because

$m_Q a \geq 1$. $m_b a \sim 2 - 3$, $m_c a \sim 0.5 - 1$.

$\vec{p} \approx m_Q$ very distorted and naive use of relativistic quark formulations (Wilson, clover) give large errors ($\propto m_Q a, (m_Q a)^2$).



BUT, b and c are non-relativistic in their bound states.

m_Q and $\vec{p} \approx m_Q$ are irrelevant dynamical scales. Can treat b and c quarks accurately on the lattice with non-relativistic techniques.

Several ways to proceed:

1. **Static Quarks** : $m_Q = \infty$ limit. Spinless and flavourless. Quark prop = string of gluon fields in time dirn. Useful limit for understanding HQET.

2. **NRQCD** : Non-relativistic effective theory.

$$\mathcal{L}_Q = \bar{\psi} \left(D_t - \frac{\vec{D}^2}{2m_Q a} - c_B \frac{\vec{\sigma} \cdot \vec{B}}{2m_Q a} + \dots \right) \psi$$

ψ a 2-component spinor.

$m_Q a$ fixed by requiring one heavy hadron mass correct. $E_h(p) = E_h(0) + p^2/2m_h$.

c_i fixed by pert. or nonpert. matching to QCD.

$m_Q \rightarrow \infty$ is static.

Cannot take a to 0 but improve until a -dependence small enough.

3. Heavy Wilson quarks (FNAL method). :

$$\mathcal{L}_Q = \bar{\psi}(\not{D} + m_Q a - \frac{iac_{sw}}{4}\sigma_{\mu\nu}F^{\mu\nu})\psi$$

but interpret non-relativistically to fix $m_Q a$. Match to QCD with $m_Q a$ -dependent coefficients.

Small $m_Q a$ limit is light quarks.

Large $m_Q a$ limit is NRQCD.

4. **Wilson/clover quarks** : Same action as above but work only at small ma (OK for m_c ?).

Extrapolate to large $m \Rightarrow$ large errors *and* expensive. Anisotropic lattices may be better.

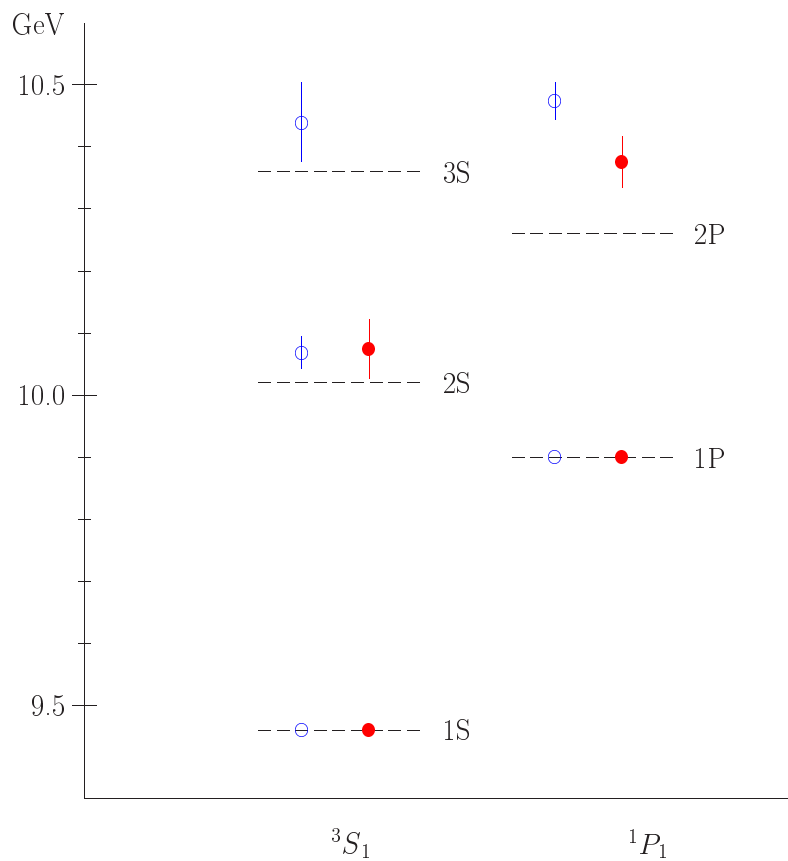
(fine $a_t \Rightarrow$ small ma_t even with large a_s). See

X. Liao parallel talk.

Results on the spectrum

1. $b\bar{b}$ (Υ) spectrum (UKQCD collaboration - Marcantonio et al)

Radial and Orbital Splittings in the $b\bar{b}$ system

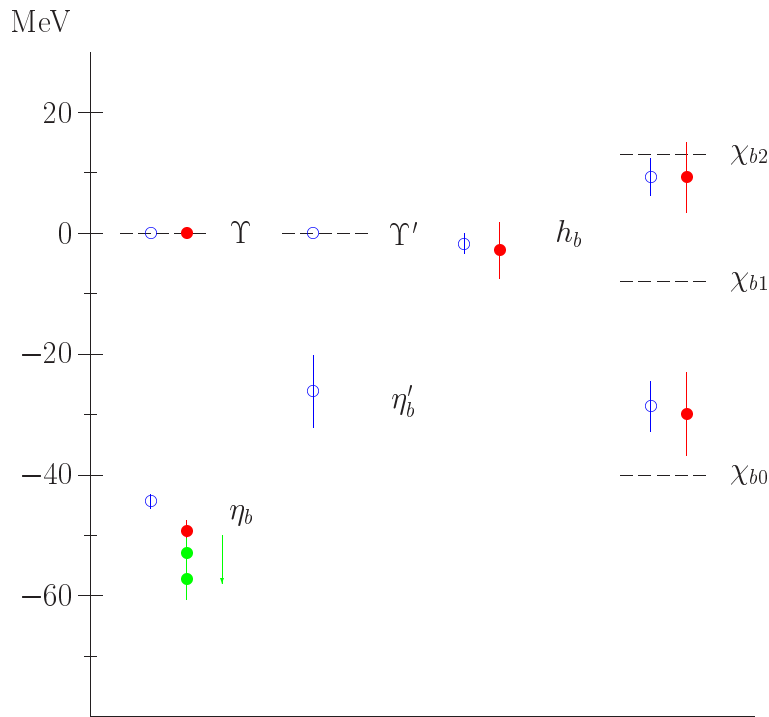


--- : Experiment

○ : NRQCD for the b , glue: $n_f = 0, \beta = 6.0$.

● : NRQCD for the b , glue: $n_f = 2, \text{clover}, \kappa = 0.135, \beta = 5.2$, UKQCD ensemble.

Fine Structure in the $b\bar{b}$ system



--- : Experiment

○ : NRQCD for the b , glue: $n_f = 0, \beta = 6.0$.

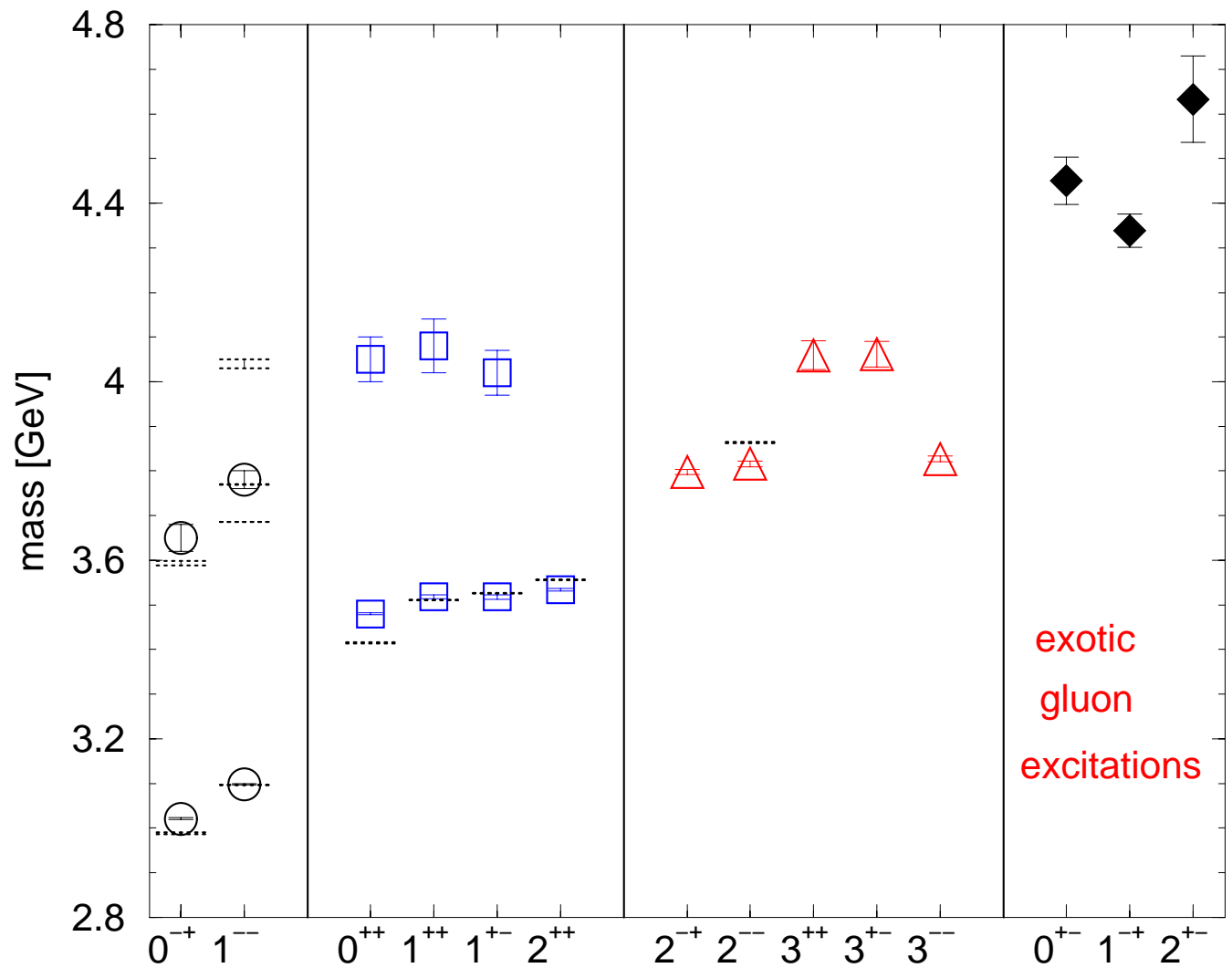
● : NRQCD for the b , glue: $n_f = 2, \text{clover}, \kappa = 0.135, \beta = 5.2$, UKQCD ensemble.

● ↓ : extrapolate to light dynamical mass and to $n_f = 3$.

Hyperfine splitting sees dynamical quarks. Predict $m(\Upsilon) - m(\eta_b) = 60 \pm 15$ MeV. Aim for 1 – 2 % error for CLEO-c.

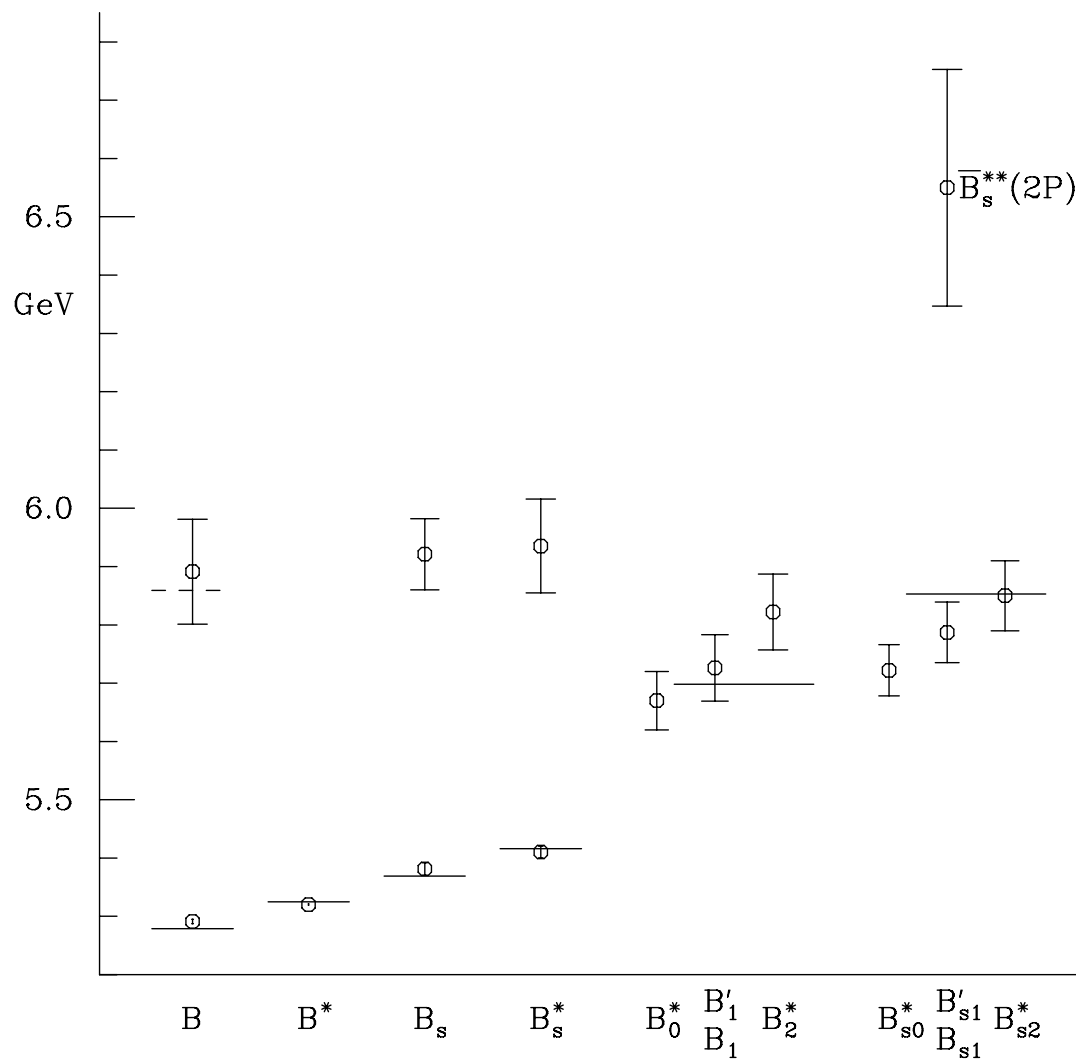
2. $c\bar{c}$ (ψ) spectrum

Columbia results (Chen et al) on anisotropic lattices in the QA, including $c\bar{c}g$ hybrids.



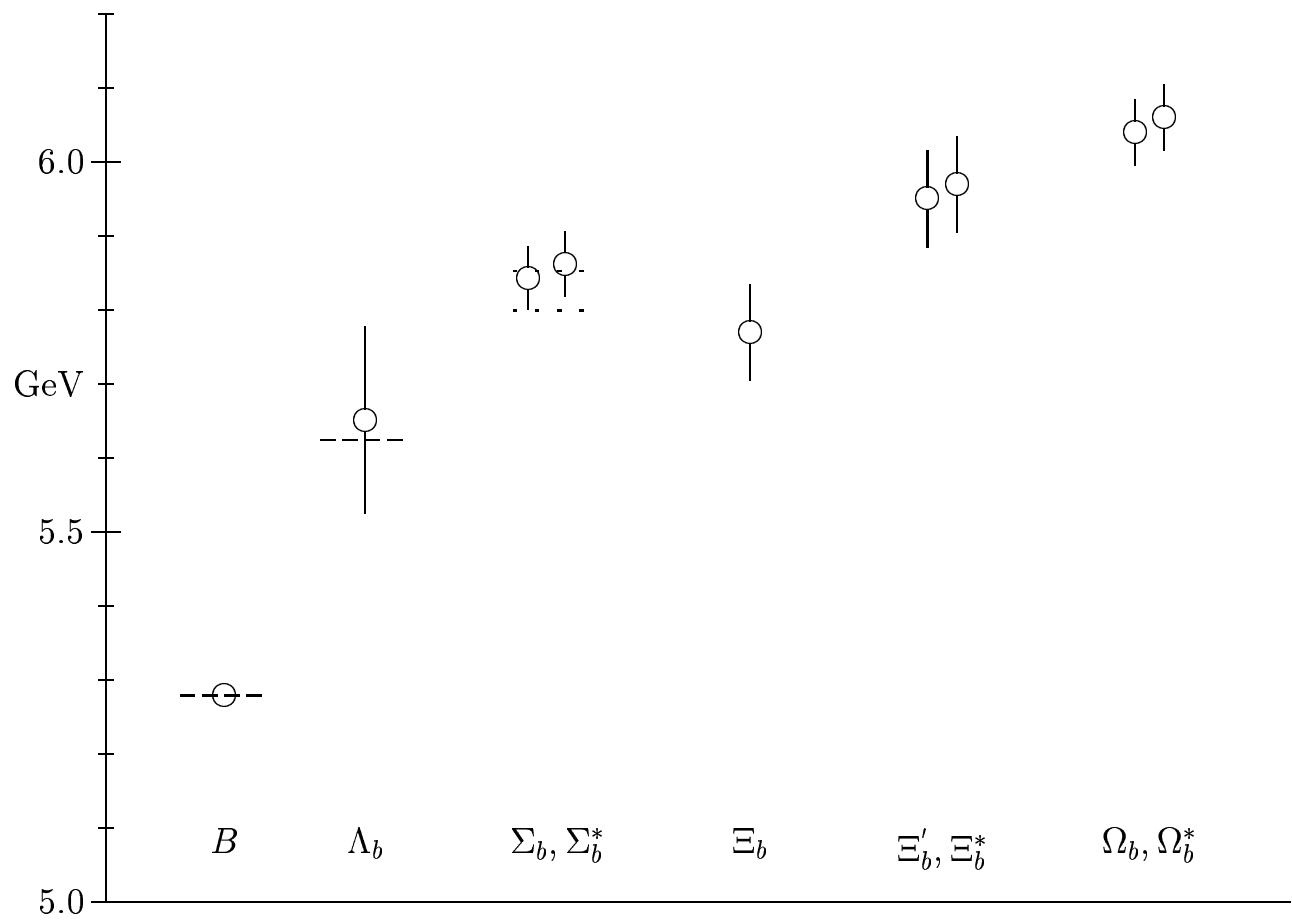
3. b -light (B) spectrum

NRQCD results in the QA (Hein et al).



4. b -light-light baryon spectrum

NRQCD results in the QA (Ali Khan et al) for udb , usb , ssb states.



Quark masses

m_b fixed s.t. B or Υ mass correct. Best determination is from B in static limit in QA.

$$\bar{m}_b(\bar{m}_b) = Z_{cont}(m_B - E_B(\vec{p} = 0) + E_0)$$

Z_{cont} and E_0 known in pert. th. through α_s^3 .

'World average' is 4.30(10) GeV. (Ryan, LAT01)

$m_b \approx 50$ MeV smaller unquenched (?)

New non-pert methods in progress (Sommer et al).

m_c from ψ and α_s^2 lattice mass renorm for non-rel. case (LAT01: Juge et al)

$$m_c(M_c) = 1.28(4) \text{ GeV, QA}$$

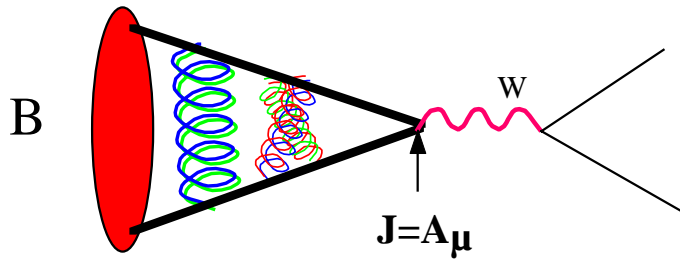
m_c from D_s and non-pert renorm. for rel. case (Becirevic et al)

$$m_c(m_c) = 1.26(13) \text{ GeV, QA.}$$

Determining B matrix elements

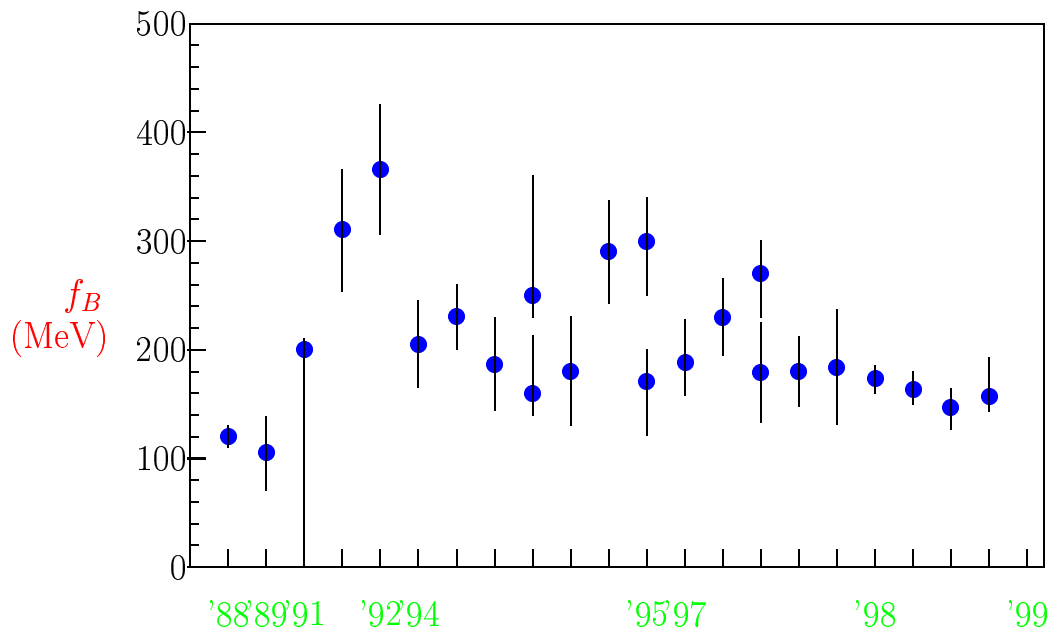
1. f_B

Simplest 2pt m.e. for B leptonic decay.



$$\langle 0 | A_\mu | B \rangle = p_\mu f_B$$

Has improved a lot with time (QA).



Match A_μ on lattice to continuum. Done to $\mathcal{O}(\alpha_s, 1/m_Q, a)$ for NRQCD and heavy Wilson.
 Typical error ‘budget’:

Source	<i>percent</i>
statistical + interp.	3
disc. $O((a\Lambda)^2)$	4
pert. $O(\alpha_s^2, \alpha_s^2/(aM))$	7
NRQCD $O((\Lambda/M)^2, \alpha_s\Lambda/M)$	2
light quark mass	+4
$a^{-1}(m_\rho)$	4
Total	10

World averages (Ryan, LAT01)

$$f_B^{(QA)} = 173 \pm 23 \text{ MeV, } 20\% \text{ larger unq. ?}$$

$$f_{D_s}^{(QA)} = 203 \pm 14 \text{ MeV}$$

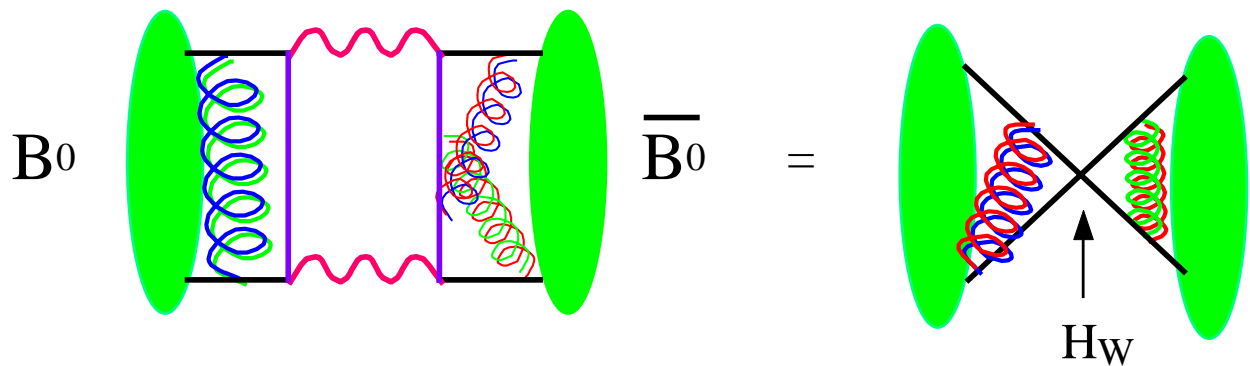
$$f_{B_s}/f_B = 1.15(5); f_{D_s}/f_D = 1.16(4)$$

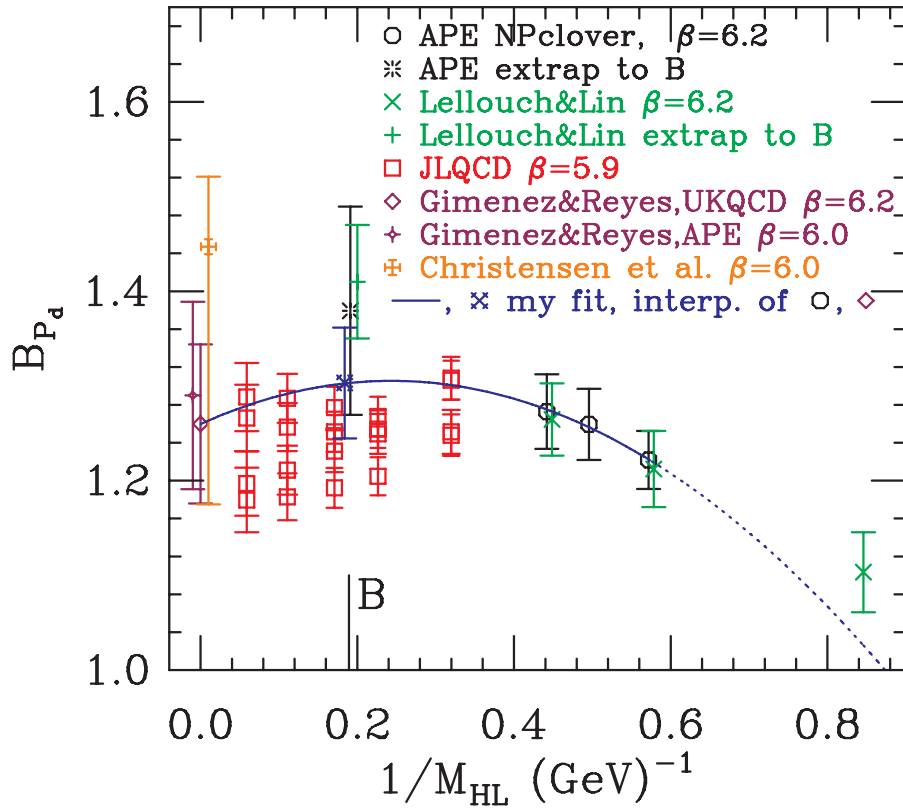
Need to reduce pert. errors by $\mathcal{O}(\alpha_s^2)$ calcs/
 non-pert. techniques.

2. B_B

Matrix element for 'box' diagram for $B^0 - \bar{B}^0$ mixing – becomes m.e. of 4-q operator from H_W . Convenient to take ratio to f_B^2 , call result B_B .

$$\Delta m_B \propto |V_{tb}^* V_{tq}|^2 f_B^2 B_B.$$





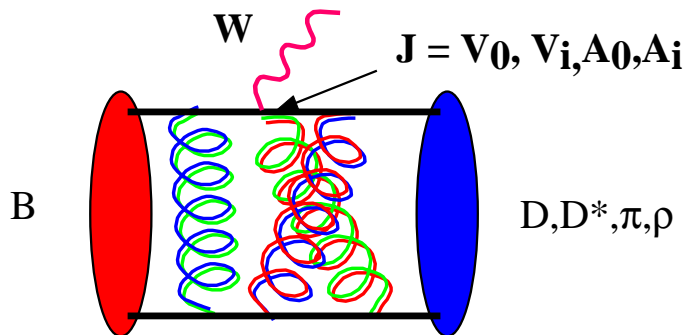
World averages (Ryan, LAT01, Bernard, LAT00):

$$\hat{B}_{B_d} = 1.30(12)(13); f_{B_d} \sqrt{\hat{B}_{B_d}} = 230(40) \text{ MeV.}$$

$$\hat{B}_{B_s} / \hat{B}_{B_d} = 1.01(3);$$

$$\xi \equiv \frac{f_{B_s} \sqrt{\hat{B}_{B_s}}}{f_{B_d} \sqrt{\hat{B}_{B_d}}} = 1.16(5)$$

3. B SL decay.

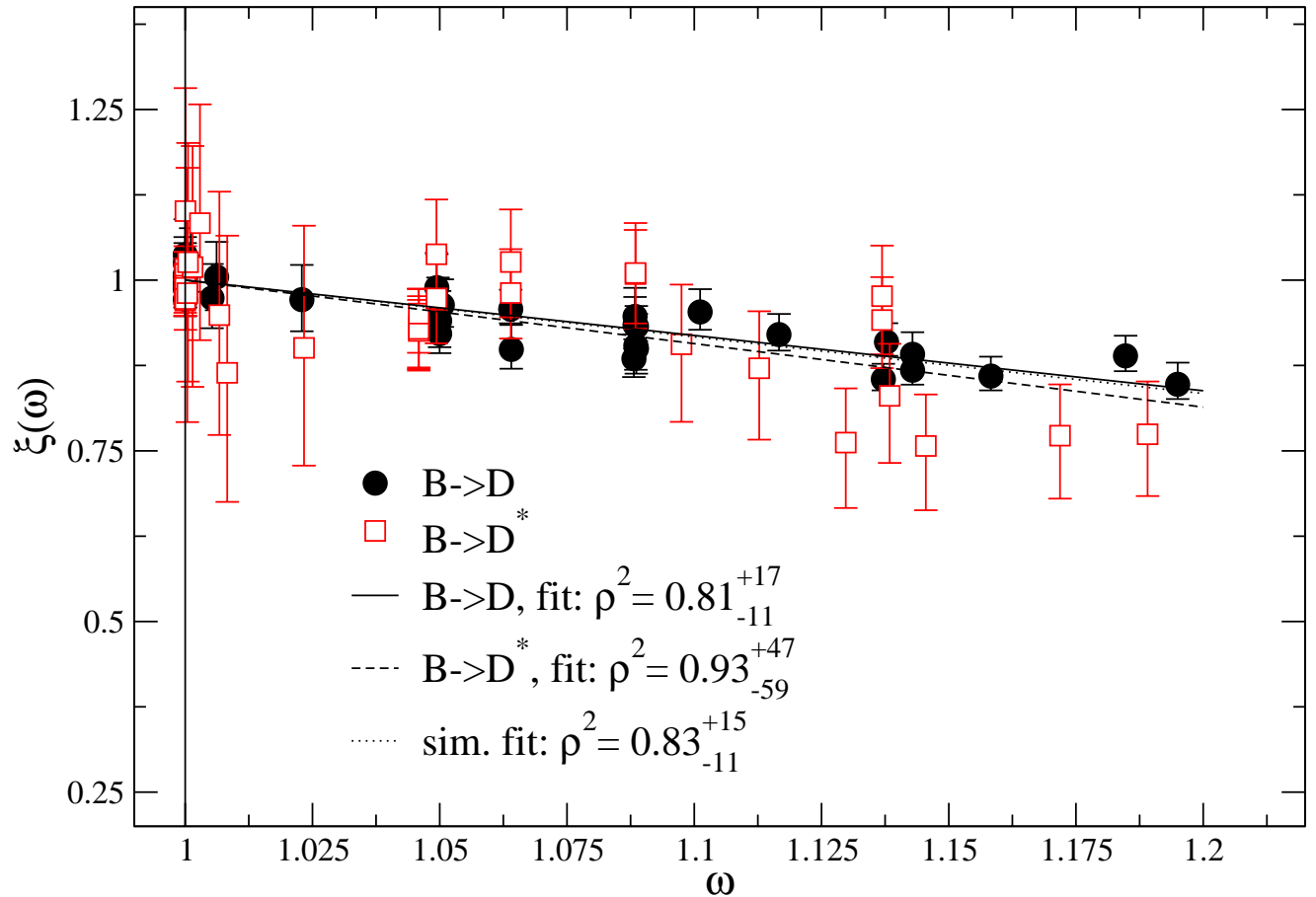


Matrix elements needed for determination of V_{ub}, V_{cb} .

$B \rightarrow D^{(*)}$ decay.

Heavy quark symmetry very useful here \rightarrow study matrix element as function of $\omega = v_B \cdot v_D$.

Form factor then has universal shape to a good approx. = Isgur-Wise function.



(UKQCD - rel. c quarks)

exptl rate($v_B \cdot v_D = 1$) = $|V_{cb}|^2 \mathcal{F}(1)$.

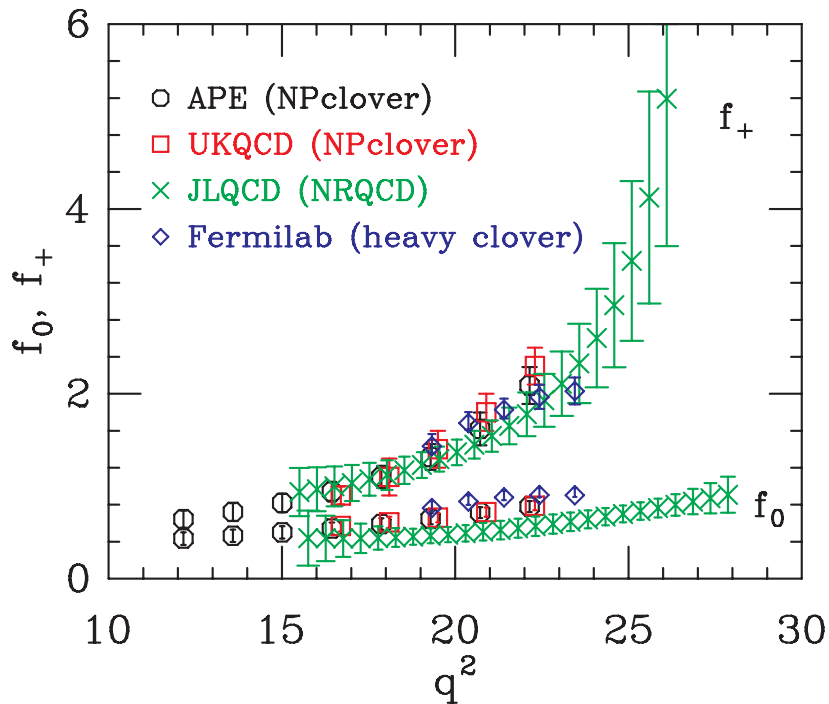
FNAL LAT01(Hashimoto et al):

$\mathcal{F}_{B \rightarrow D^*}(1) = 0.913(30)$. (QA)

gives $V_{cb} \times 10^3 = 38.7 \pm 1.8 \pm 1.5$ (LEP).

$B \rightarrow \pi, \rho$ decay.

Lattice calculations (all QA) work at small \vec{p}_π , far from physical region. Require extrapolation, interpolation, etc. Smooths out v. rough raw data. More work is needed.



Soft pion theorem $f_0(q_{max}^2) = f_B/f_\pi$ doesn't work well. Chiral extrapolations to light quark masses very important.

Future

- MILC collab. using improved staggered quarks → can simulate light dyn. quarks with Tflop computers in next few years.
- CLEO-c will determine Υ, ψ, D physics to high precision.

Lattice calcs must improve systematic errors to 2-3%.

- Improve heavy quark actions, higher rel. corrns and better match to continuum.
- Improve matching of matrix elements using automated pert. th. (Trottier, Horgan)
- Simulate with light dynamical quarks and match to chiral pert. theory.