

Electromagnetic Production of Pions in the Resonance Region - Theoretical Aspects

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Motivation

Understand the hadron and nuclear dynamics within QCD

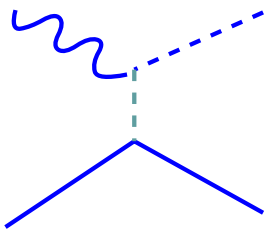
Probe the quark-substructure of hadrons with GeV photons and electrons investigating N^* excitations

Consider photo-pion production mechanism

Δ and N^* in Meson-Baryon continuum

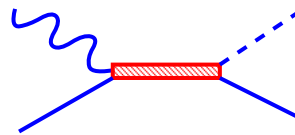
Two reaction mechanisms:

Fast



Particle Exchange

Time delayed



Resonance

- Reaction amplitudes must have the following form

$$T(E) = t_{bg}(E) + t_{res}(E)$$

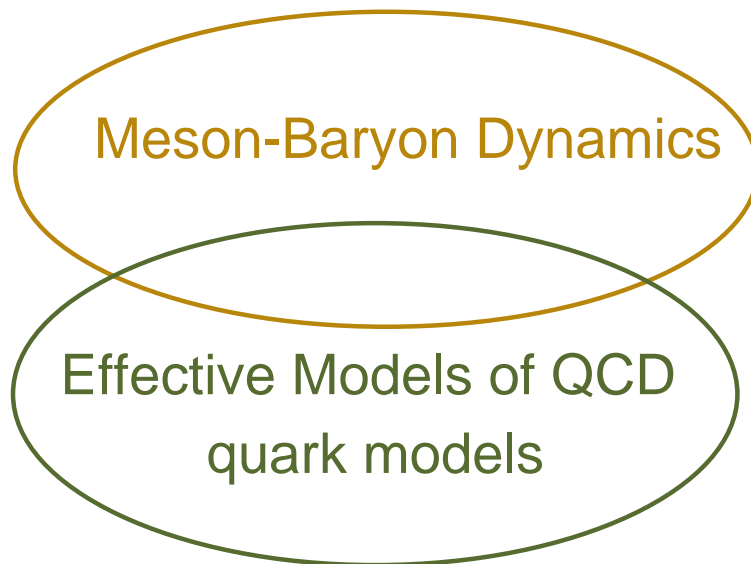
$t_{bg}(E)$: amplitude due to direct processes

- Unitary Condition $T^\dagger T \sim Im[T]$

t_{bg} and t_{res} are related

Example: $\pi N t_{res} \sim e^{i(\delta+\delta_{bg})} |t_{res}|$

$N(e, e\pi)$, $N(\gamma, \pi)$ Data



QCD

- Brief review of theoretical approaches
- $\gamma N \rightarrow \pi N$ and $N(e, e'\pi)$ Reaction at Δ Region
- Summary

Theoretical Approaches:

- Dispersion Relations

Introduced by Chew, Goldberg, Low Nambu in 1950's

Recently Mainz Group for $\gamma N \rightarrow \pi N$ for $E_\gamma < 400 MeV$

- Effective Lagrangian Approaches

RPI, Mainz(MAID), Giessen-GW(multi-channel)

t_{bg} : tree-diagram of effective Lagrangian

t_{res} : Isobar Model

Unitarization : K-matrix approaches

- Dynamical Approaches

Tanabe-Ohta, Yang, Nozawa-Blenkleider-Lee, Gross-Surya, Julich Group, Sato-Lee(SL), Kamalov-Yang (DMT)

Solve dynamical scattering equations with interactions defined by effective Hamiltonian and/or hadron models

Effective Lagrangian Approaches

$$T = V + VG_0T \quad K = V + VRe(G_0)K$$

$$T_{fi}^{on} = \sum_n \left(\frac{1}{1 - iK_{on}} \right) f_n K_{ni}^{on}$$

$K \sim V$ and first order EM

$$T_{\pi,\gamma}^{on} = \sum_n \left(\frac{1}{1 - iV_{on}} \right) \pi_{,n} V_{n\gamma}^{on}$$

$$V_{\gamma\pi} = \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \text{[diagram 4]}$$

- Giesen-GW

Multi channels and resonances

$$n = \pi N, \pi\pi(\zeta)N, \eta N, K\Lambda \quad + \quad \Delta, N^*$$

- MAID

$$T_{\pi,\gamma}^{on} = (1 - iT_{on})_{\pi,\pi} V_{\pi\gamma}^{BG} + \sum_r T^r(BW)$$

$$T^r(BW) = A^r(Q^2) \frac{f_{\gamma r}(E) \Gamma_r M_r f_{\pi r}(E)}{M_r^2 - E^2 - iM_r \Gamma_r} e^{i\phi_r(E)}$$

$$n = \pi N, \quad \eta_{\pi N} e^{2i\delta_{\pi N}} \rightarrow T_{\pi N}^{on}$$

Dynamical Approach (SL)

- Define a Hamiltonian for $M + B \leftrightarrow B$ transitions

M: $\gamma, \pi, \rho, \omega..$ B: $N, \Delta, N^*, ..$

Coupling constants calculated from a hadron model

- Apply unitary transformation method to eliminate 'off shell process'

$$H_{eff} = H_0 + v_{bg} + \gamma_{\Delta \leftrightarrow MN}$$

v_{bg} : energy independent non-resonant interactions

- Derive a reaction model from H_{eff} for meson nucleon reactions

Dynamical equations in a "finite" Fock space: a solvable few-body problem

Pion Production Amplitude

$$\begin{array}{c}
 \text{---} \\
 | \\
 T_{\gamma\pi} \\
 | \\
 \text{---}
 \end{array}
 =
 \begin{array}{c}
 \text{---} \\
 | \\
 t_{\gamma\pi} \\
 | \\
 \text{---}
 \end{array}
 + \frac{\Gamma_{\pi}(W) \Gamma_{\gamma}(W)}{W - m^0 - \Sigma(W)}$$

with

$$\begin{array}{c}
 \text{---} \\
 | \\
 \text{---}
 \end{array}
 =
 \begin{array}{c}
 \text{---} \\
 | \\
 \text{---}
 \end{array}
 +
 \begin{array}{c}
 \text{---} \\
 | \\
 \text{---}
 \end{array}$$

$$\Gamma_{\gamma} = \gamma_{\gamma} + \int \Gamma_{\pi} G_0 v_{\gamma\pi}$$

Dressed
Bare

Compare with other works

$$T_{\gamma\pi} = e^{i\delta} \cos \delta [v_{\gamma\pi} + \int K_{\pi N} \text{Re}(G_0) v_{\gamma\pi}] + \frac{\Gamma_{\pi} \gamma_{\gamma}}{E - m^0 - \Sigma}$$

- DMT 0 0 $T(BW)$
- MAID 0 $T(BW)$

$N(\gamma, \pi)$ and $N(e, e'\pi)$ at Δ_{33} Region

Main issues of $\gamma + N \rightarrow \Delta$

- Magnetic dipole coupling and Q^2 evolution
- Is Δ deformed?

Dynamical approach of SL: Can describe extensive data

- πN scattering [Fitted]
- $\gamma N \rightarrow \pi N$ LEGS, Mainz
[Extract G_M, G_E of $\gamma N \rightarrow \Delta$]
- $N(e, e'\pi)$ Jlab, MIT-Bates, Mainz, NIKHEF [Q^2]

$(e, e'\pi)$ Differential Cross Section

$$\frac{d\sigma^v}{d\Omega} = A + B \cos \phi_\pi + C \cos 2\phi_\pi + h_e D \sin \phi_\pi$$

$$A = \frac{d\sigma_T}{d\Omega} + \epsilon \frac{d\sigma_L}{d\Omega} \qquad B = \sqrt{2\epsilon(\epsilon + 1)} \frac{d\sigma_{LT}}{d\Omega}$$

$$C = \epsilon \frac{d\sigma_{TT}}{d\Omega} \qquad D = \sqrt{2\epsilon(1 - \epsilon)} \frac{d\sigma_{LT'}}{d\Omega}$$

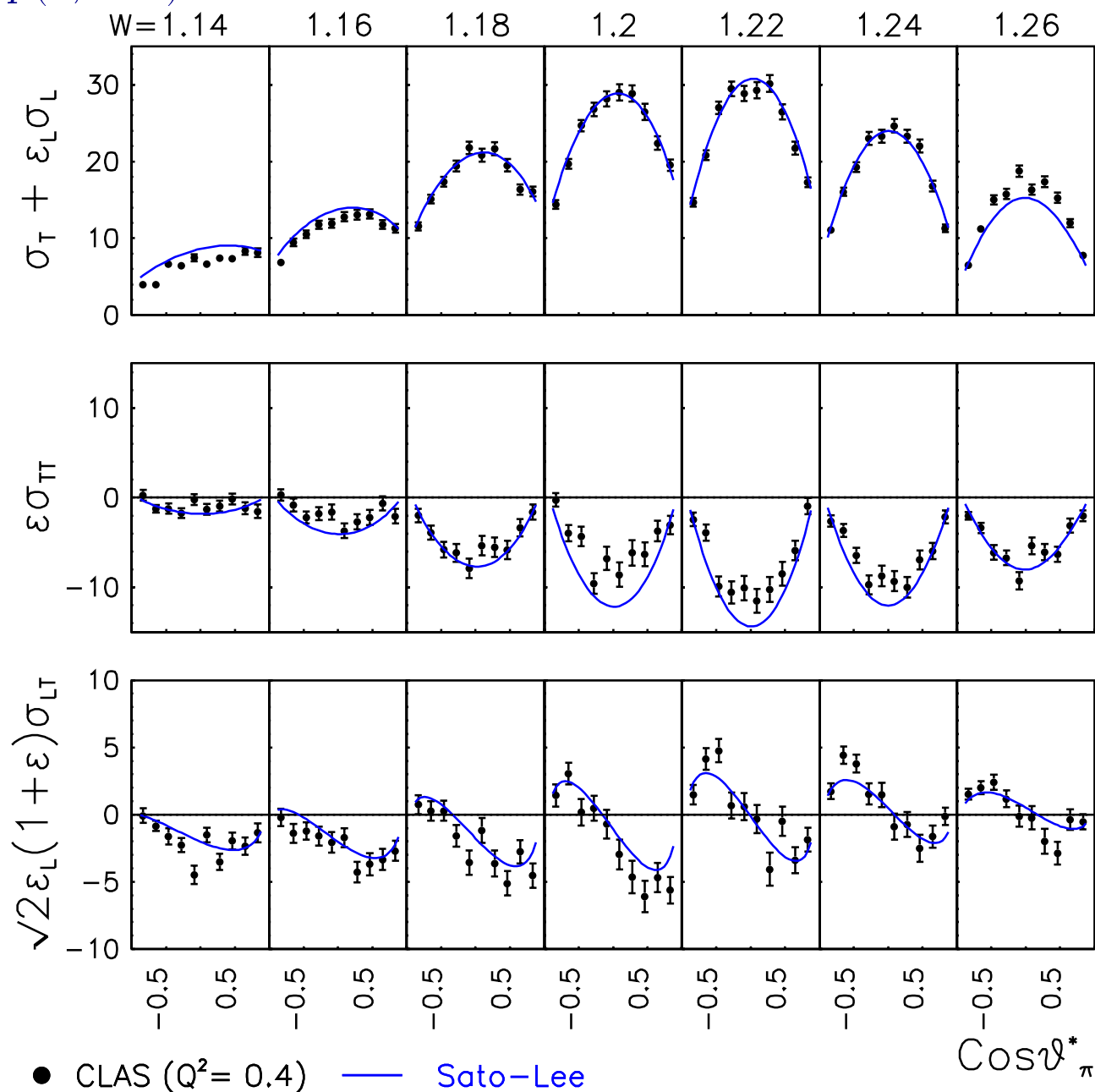
$(e, e'\pi)$

- A, B, C
- $A_{LT} \sim \sigma_{LT} \sim \text{Re}(S_{1+} M_{1+}^*)$

$(\vec{e}, e'\pi)$

- $A_{LT'} \sim \sigma_{LT'} \sim \text{Im}(S_{1+} M_{1+}^*)$

$p(e, e'\pi^0)$



Jlab from Joo

$p(e, e'\pi^+)$

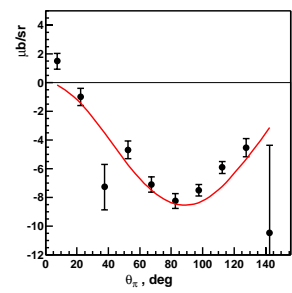
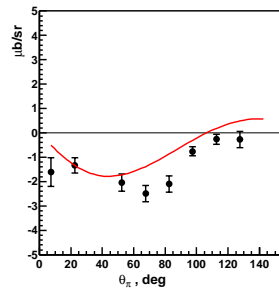
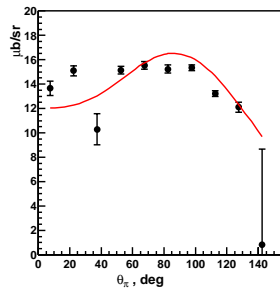
$W = 1230\text{MeV}$

$\sigma_T + \epsilon\sigma_L$

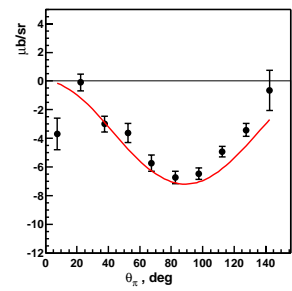
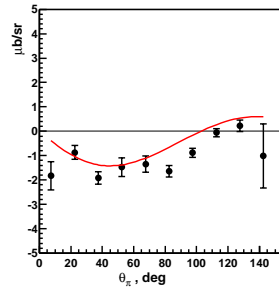
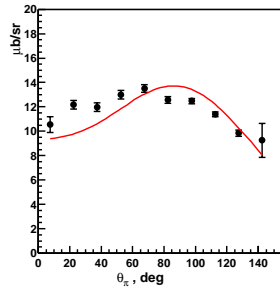
σ_{LT}

σ_{TT}

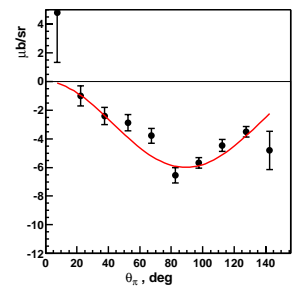
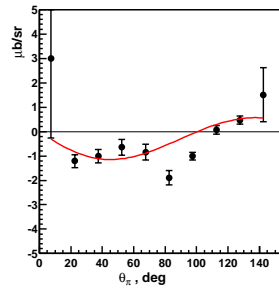
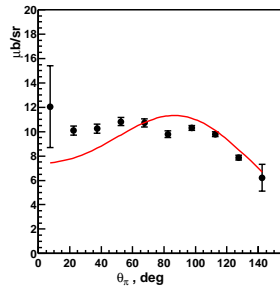
$Q^2 = 0.3$



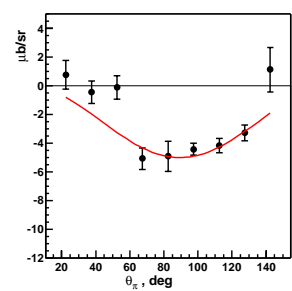
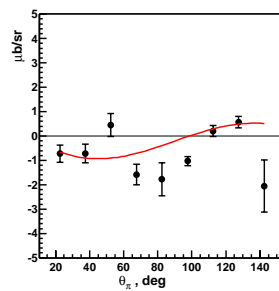
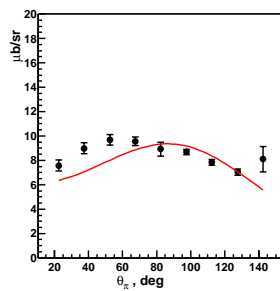
0.4



0.5



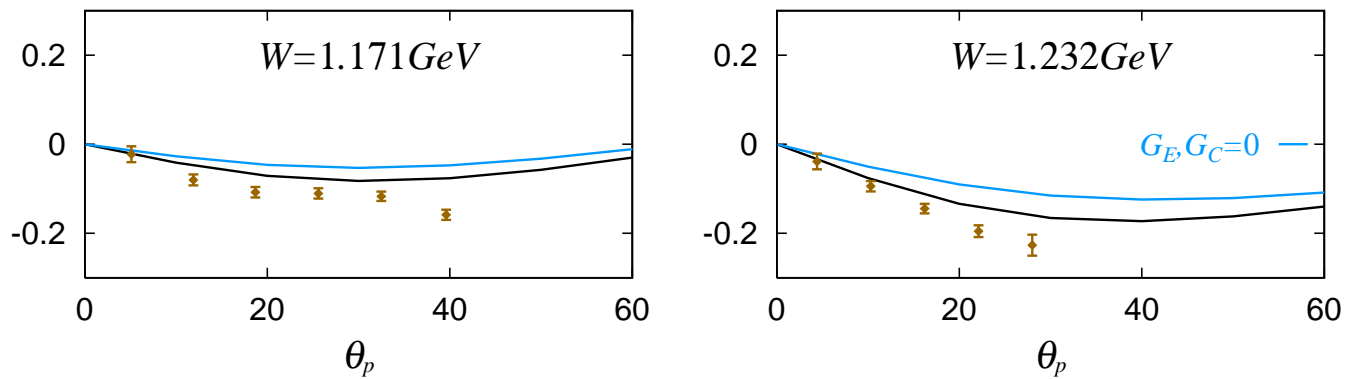
0.6



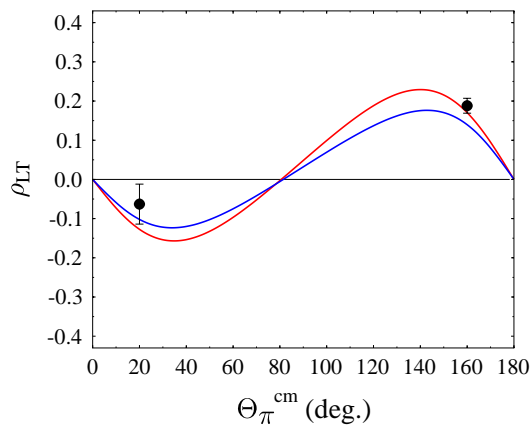
JLAB from Egiyan

A_{LT}

$Q^2 = 0.126(\text{GeV}/c)^2$ MIT-Bates C.Mertz et al.



$Q^2 = 0.2(\text{GeV}/c)^2$ $W = 1232 \text{ MeV}$ Mainz
from H. Schmieden (PRELIMINARY)

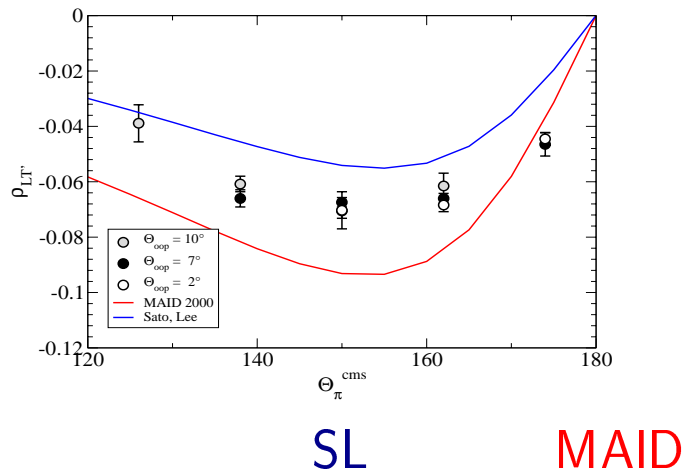


SL

MAID

$A_{LT'}$

$Q^2 = 0.2(\text{GeV}/c)^2$ $W = 1232\text{MeV}$ Mainz
from H. Schmieden (PRELIMINARY)



K. Joo (NSTAR 2001)

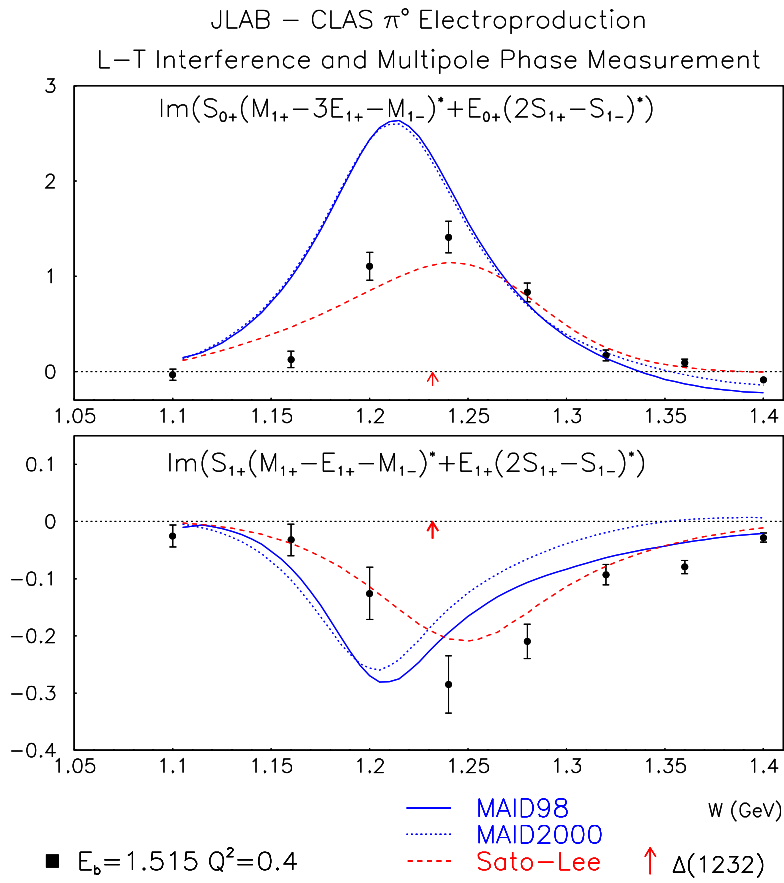


Figure 2. Top figure: D_0' in μb as a function of W (GeV). Bottom figure: D_1' in μb as a function of W (GeV) at $Q^2=0.4$ GeV². Error bars are statistical only. Results are very preliminary.

MAID

SL

$$\gamma N \rightarrow \Delta$$

$$\Gamma_{\gamma N \rightarrow \Delta} = \gamma_{\gamma N \rightarrow \Delta} + \int \Gamma_{\pi N \rightarrow \Delta} G_0 v_{\gamma\pi}$$

dressed and bare form factors have been extracted

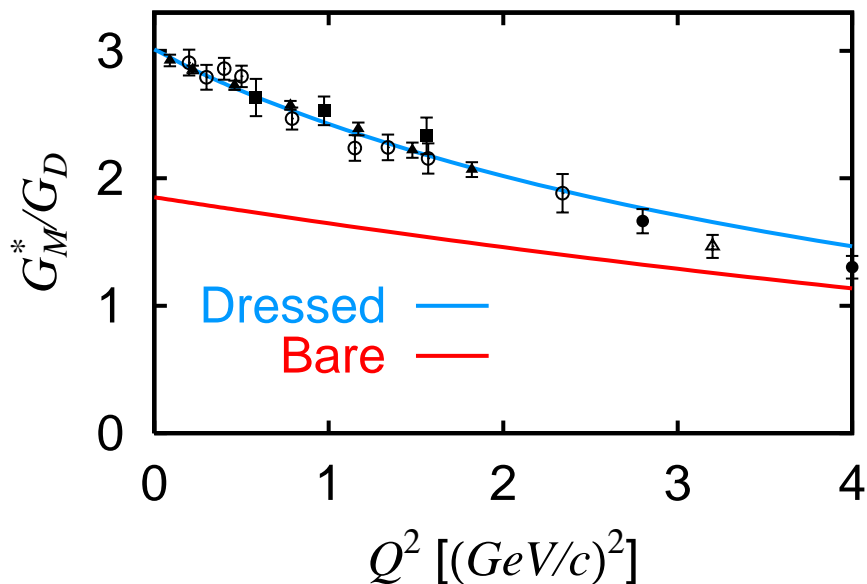
- Magnetic M1 A_M, A_M
Electric E2 A_E, A_E
Coulomb C2 A_C, A_C
- A_α compared with empirical amplitudes
- Naive Quark model $A_E = A_C = 0$

Magnetic Dipole

at $Q^2 = 0$

$$A_M = 271 \cdot 10^{-3} \text{GeV}^{-1/2}$$

$$A_M = 173$$



- $G_M^*(Q^2)$ drops faster than proton form factor
- Non-resonant process contributes about 40% at $Q^2 = 0$
- soft component of form factor due to meson cloud

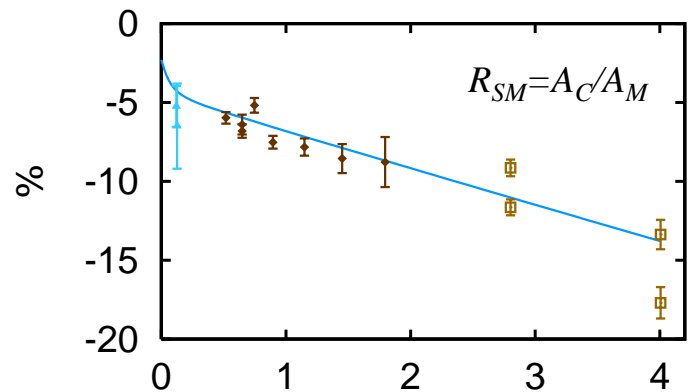
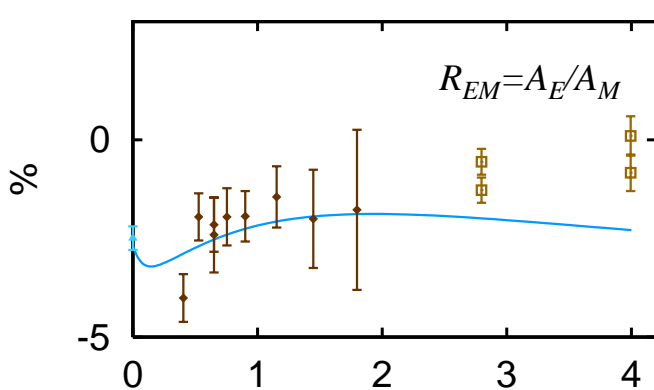
Electric and Coulomb Quadrupole

at $Q^2 = 0$

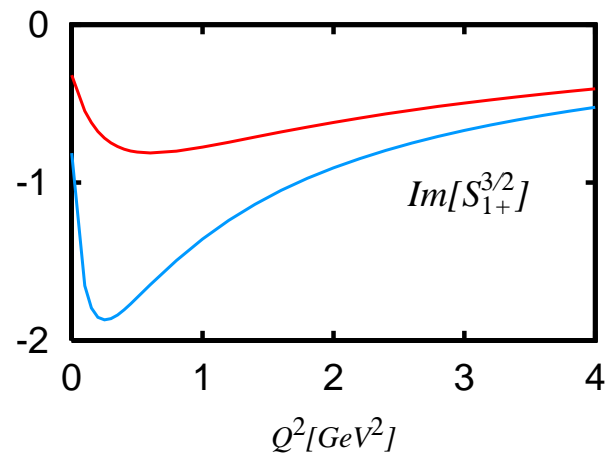
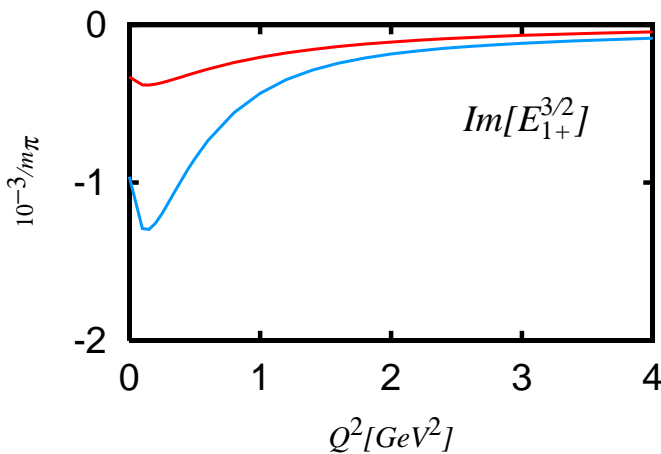
$$A_E/A_M = -2.7 \%$$

$$A_E/A_M = -1.3$$

BRAG report -2.38 ± 0.27



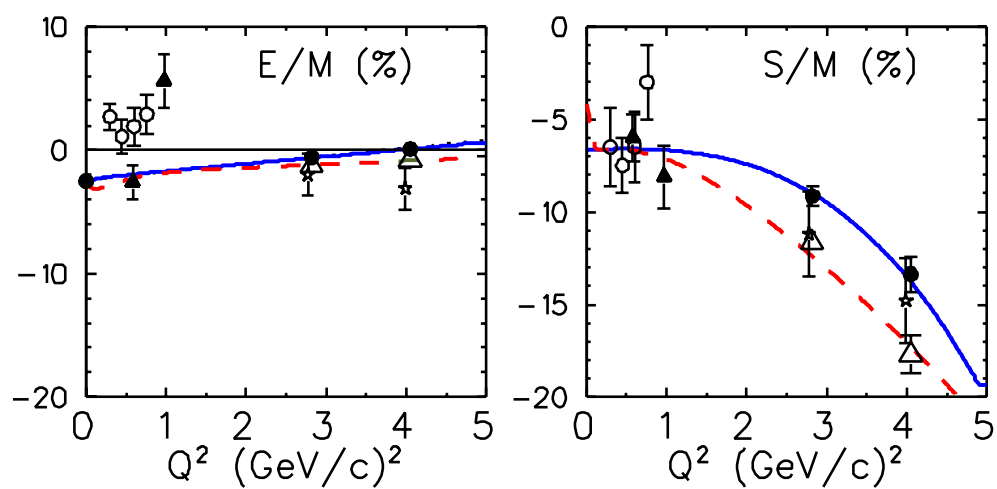
Data from DNP Town Meeting 2000



- E2,C2 components exist for both **bare** and **dress**
- **Dressed** R_{EM} , R_{SM} agree well with empirical amplitude analyses
- Low Q^2 enhancement due to pion cloud

DMT and MAID

nucl-th/0106027

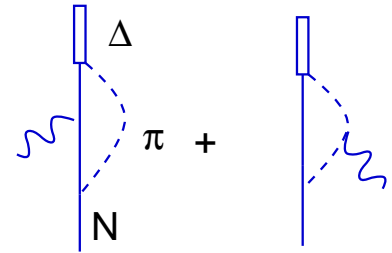


Role of Pion on $\gamma N \rightarrow \Delta$ Quadrupole (C2) Form Factor

Within SL model (Unitary transformation method)

- 'real' pion

$$A_C^{resc} = \int \Gamma_{\pi N \rightarrow \Delta} G_0 v_{\gamma\pi} =$$



- 'virtual' pion

$$|B\rangle = \frac{1}{\sqrt{Z_B}} [|qqq\rangle + C |qqq\pi\rangle + \dots]$$

→ renormalize $\gamma N \rightarrow \Delta$ vertex

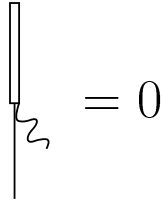
$$A_C^{Bare} =$$

~ Chiral Bag Model

(Lu et al., Bermuth et al., Dong et al.)

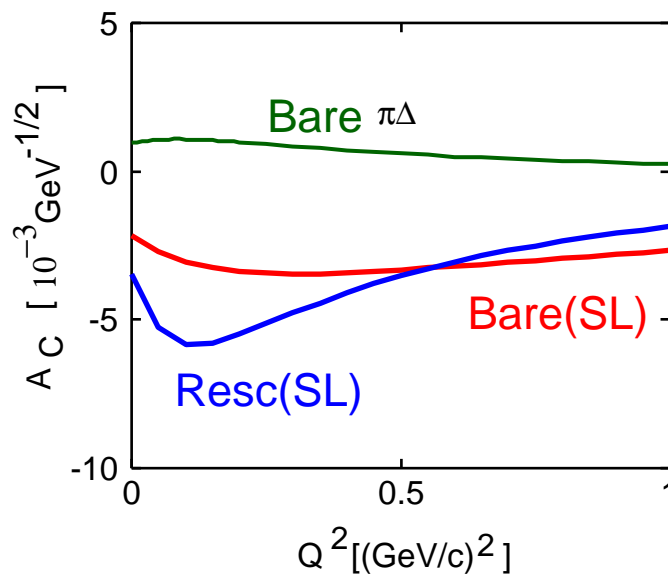
Estimation using chiral constituent quark model

$(0s)^3$ configuration



$$= 0$$

$$f_{\pi NN} \rightarrow f_{\pi N\Delta}, f_{\pi\Delta\Delta}$$



- $A_C(\pi\Delta - loop)$ does not explain $A_C^{Bare}(SL)$

Limit $m_\Delta \rightarrow m_N$

$$A_C(\pi\Delta - loop) = -A_C(\pi N - loop)$$

- Low Q^2 enhancement of S_{1+} is due to 'real' pion process

Summary

- The dynamical approach can describe extensive data of $\gamma N \rightarrow \pi N$ and $N(e, e'\pi)$ in the Δ region
 - Can be used to test predictions of baryon resonances from hadron models
 - Extracted $\gamma N \rightarrow \Delta$ form factors:
 - M1 form factor $G_M^*(Q^2)$ drops faster than proton form factor
 - Dressed $E2/M1$ and $C2/M1$ ratios agree well with empirical amplitude analyses
 - “dynamical” pion cloud has about 50% effect on the form factors at $Q^2 = 0$ and decreases as Q^2 increases
 - “static” $\pi - \Delta$ loop contribution to bare $E2$ and $C2$ disagree with extracted values
 - Beyond Δ : Three-body treatment of $\pi\pi N$ channels must be developed
- A necessary step for a investigation of higher mass N^*