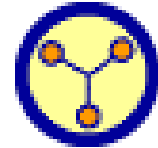


Proton Structure Results from the HERA Collider



R. Yoshida

Argonne National Laboratory, USA



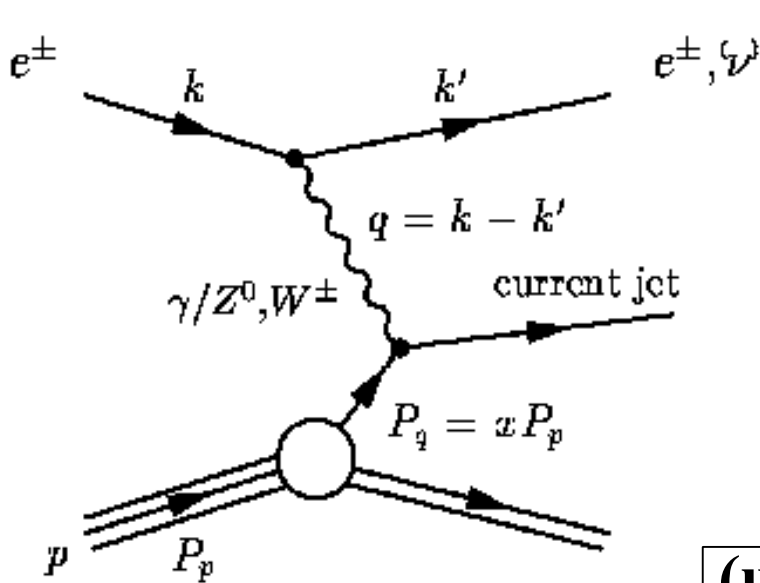
9th International Conference on Structure of Baryons

- 1. Proton Structure Functions and pQCD analyses**
- 2. DIS cross-section at low x and low Q^2**
- 3. DIS diffraction and elastic VM production**
- 4. Conclusions**



Structure function and parton distributions:

Deep inelastic scattering



$$Q^2 = -q^2$$

Virtuality ("size" of probe)

x

Mom. fraction of the struck parton.

Factorization: $\sigma_{DIS} \sim f_p(x) \otimes \hat{\sigma}$

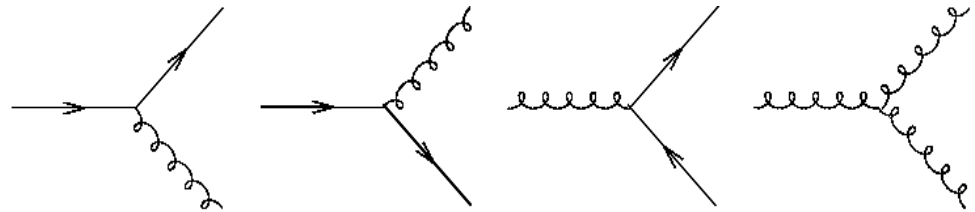
(universal) parton densities

PQCD cross-sec.

DGLAP evolution equations:

$$\frac{\partial f_p}{\partial \ln Q^2} \sim f_p \otimes P$$

P 's are splitting functions:



Str. fcn. and parton dist.:

Neutral Current (γ, Z exchange) interaction

$$\frac{d\sigma_{e^\pm p}^2}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} (Y_+ F_2 - y^2 F_L \mp Y_- xF_3)$$

New HERA (H1 and ZEUS) measurements have ca. 3% precision.

$y = Q^2/xs$, the inelasticity parameter, $Y_\pm = (1 \pm (1-y)^2)$

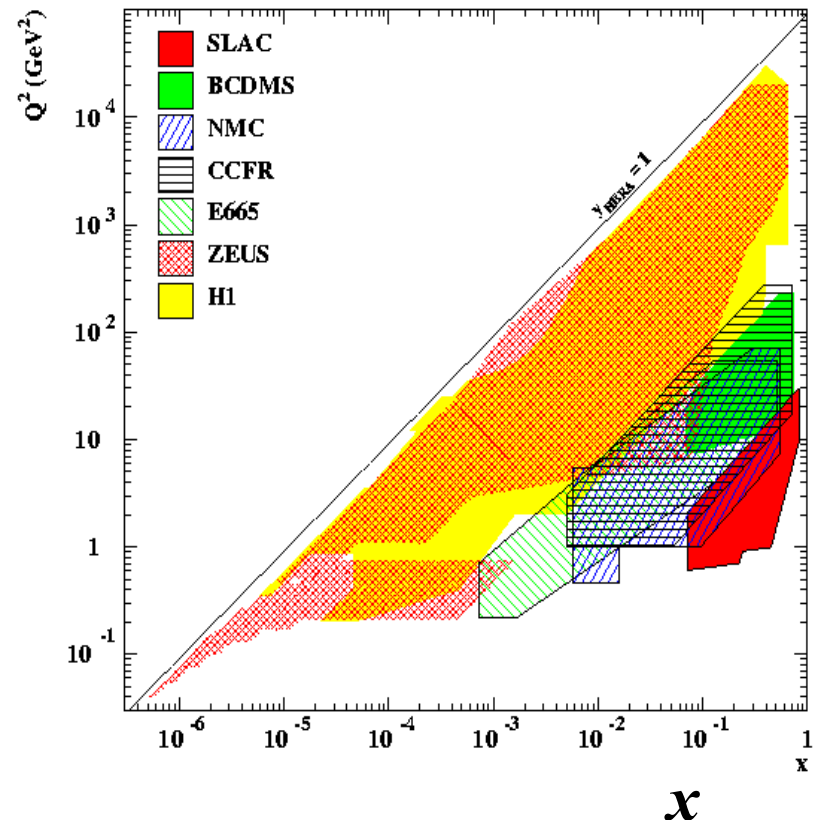
F_2, F_L , and xF_3 are structure functions of the proton.

- F_L : longitudinal component, damped by y^2 .
- xF_3 : Small at $Q^2 \ll M_Z^2$,

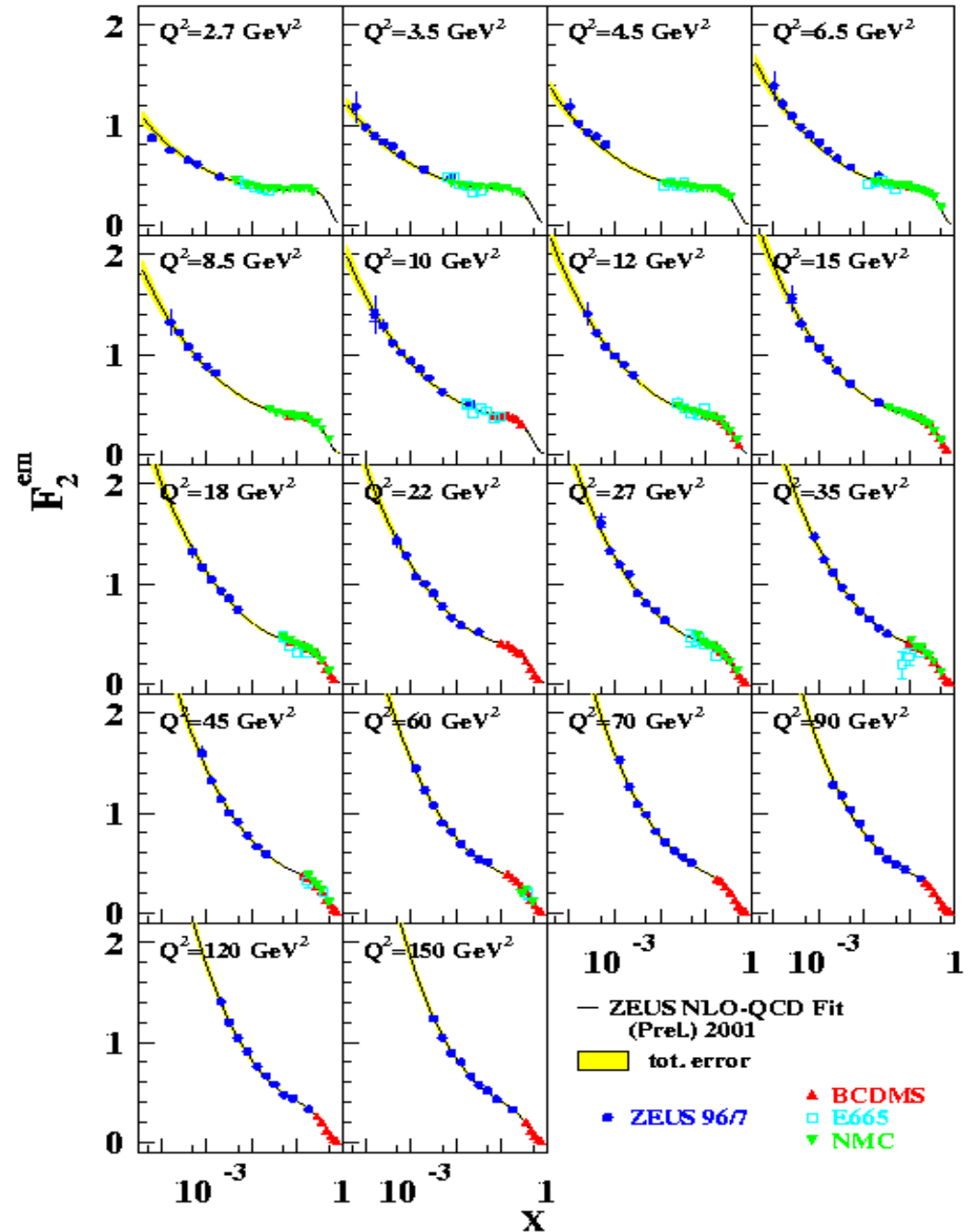
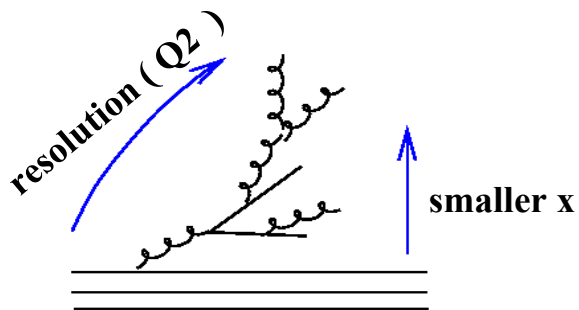
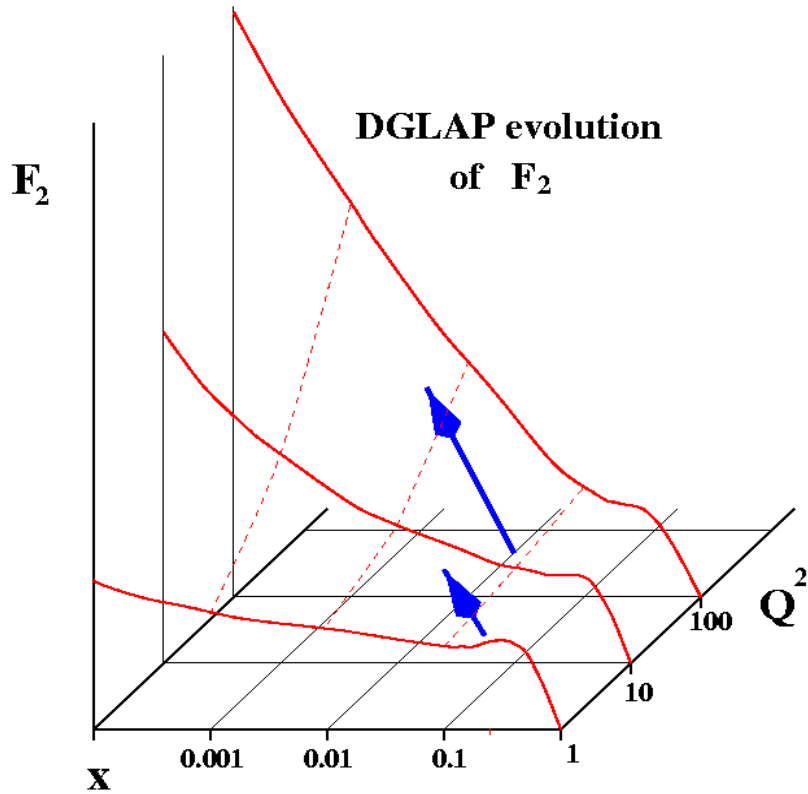
And (LO (or in DIS scheme...))

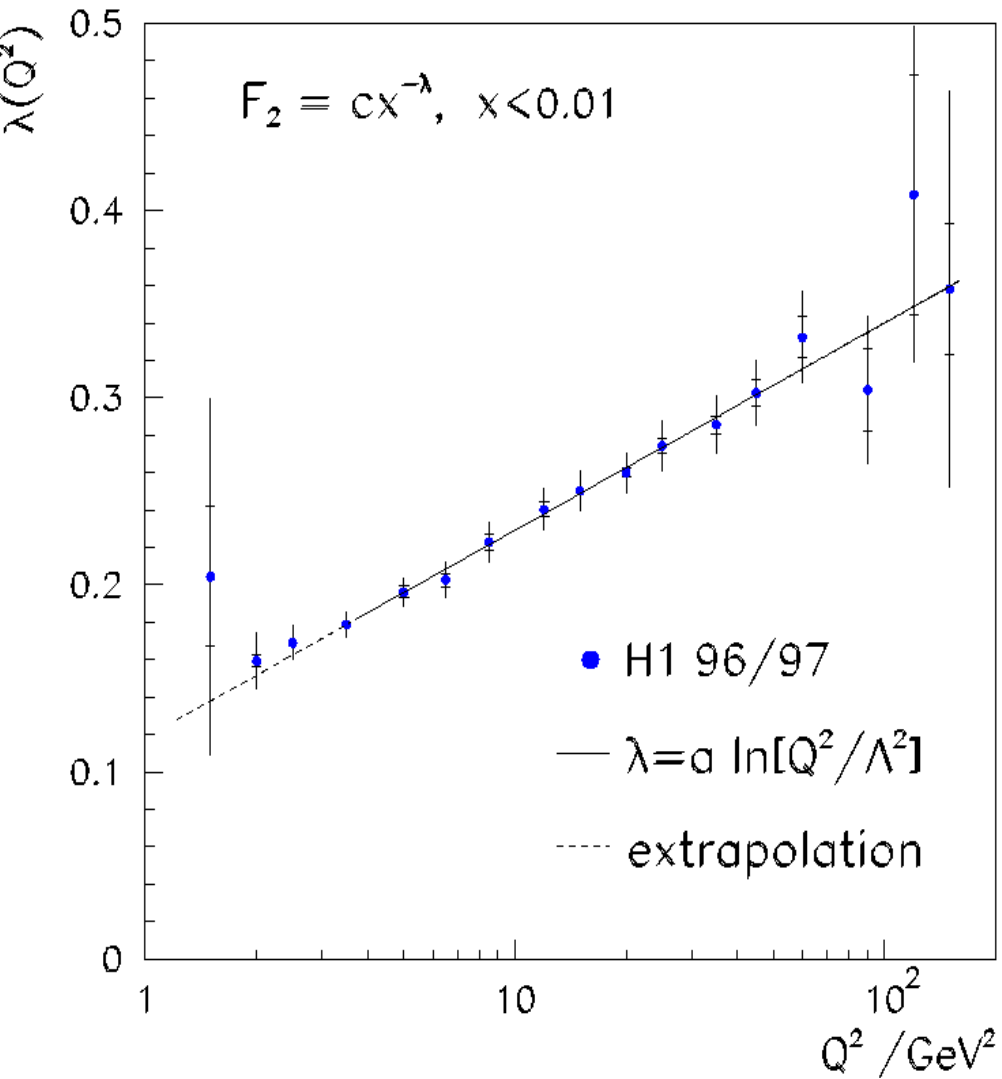
$$F_2(x, Q^2) = x \sum_q e_q^2 (q(x, Q^2) + \bar{q}(x, Q^2))$$

where q, \bar{q} are quark, antiquark densities in the proton.



Str. fcn. and parton dist.:

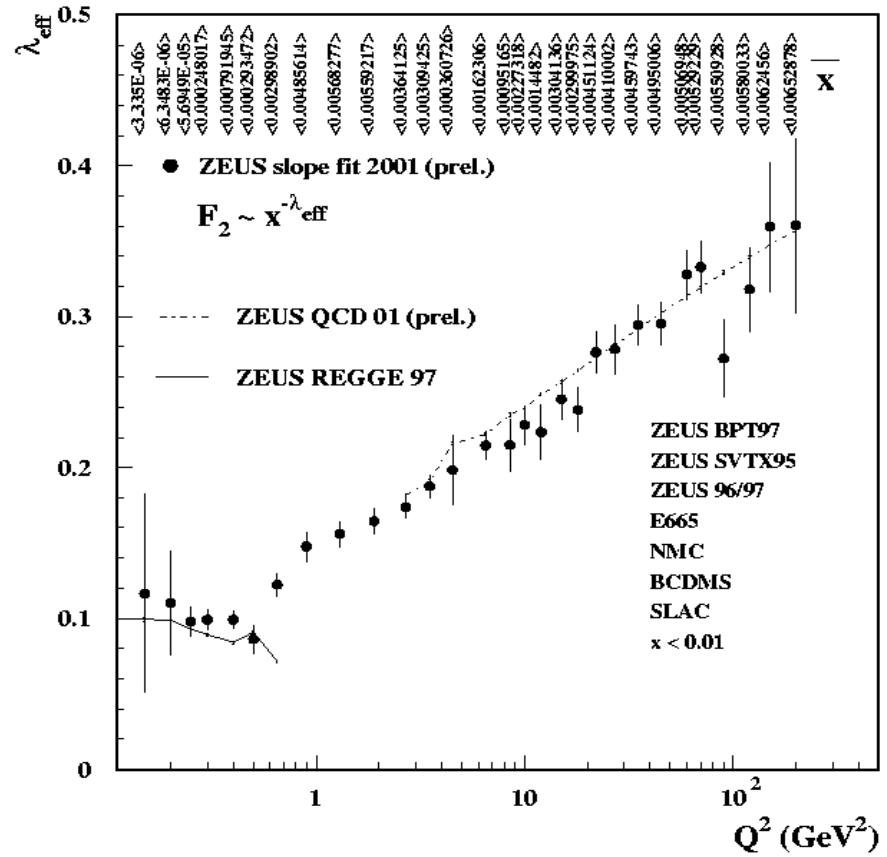




H1 Collaboration

$F_2 \propto x^{-\lambda}$
 $\lambda \approx 0.2$
 $(Q^2 \approx 10 \text{ GeV}^2)$

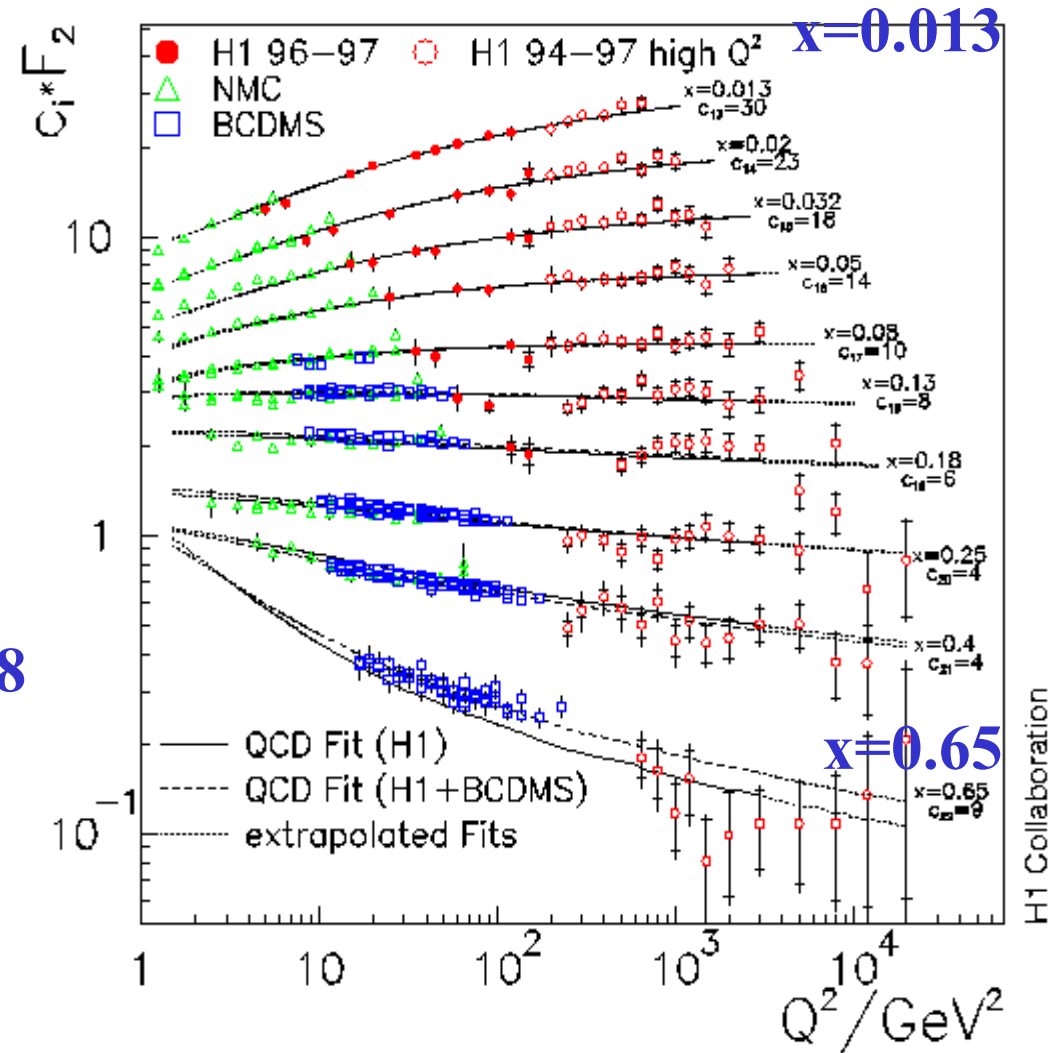
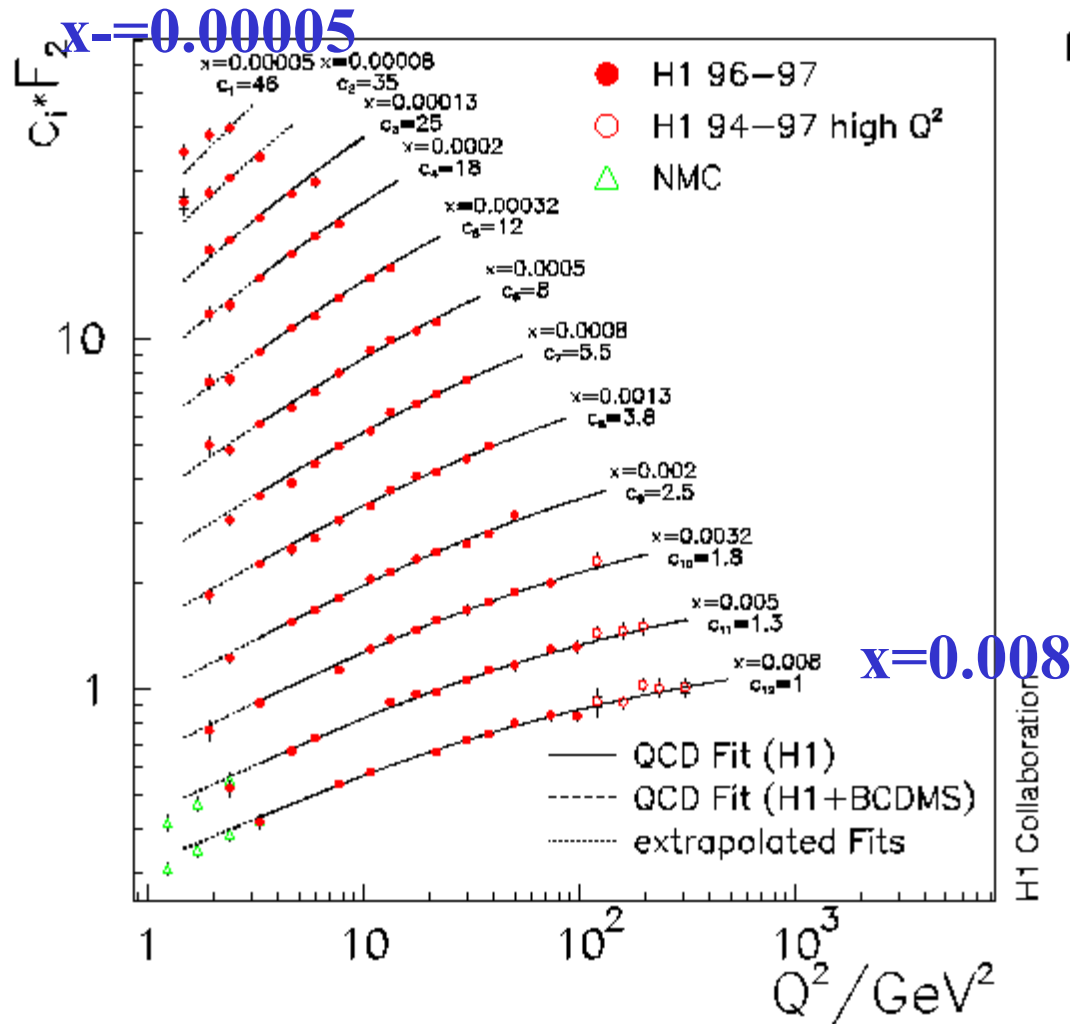
ZEUS



Tends to flatten out at lower Q^2 →



Str. fcn. and parton dist.:



To LO: $\frac{\partial F_2}{\partial \ln Q^2} \sim \alpha_s xg$

NLO DGLAP fit with q/g parameterized

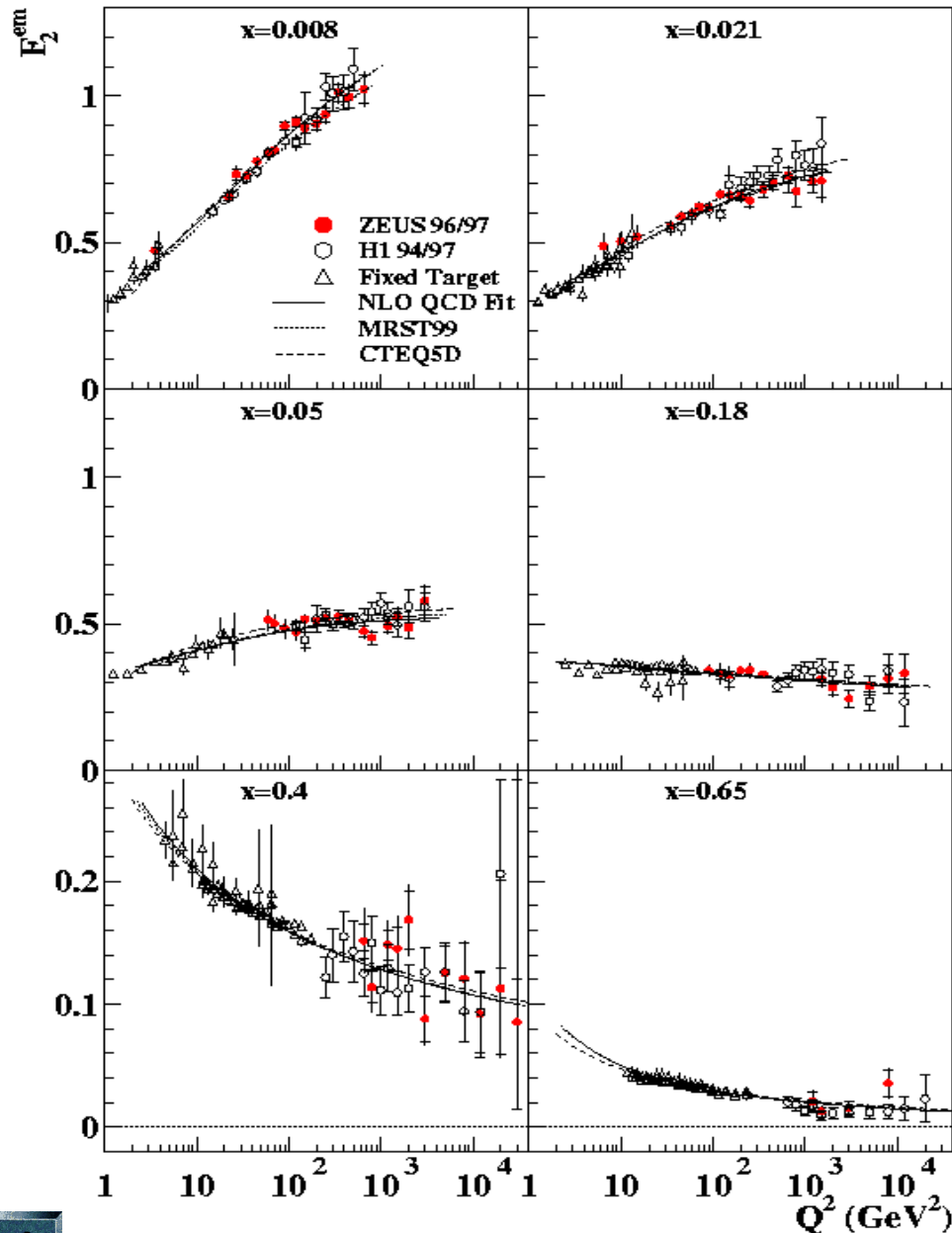
$$xq(x) = ax^b(1-x)^c[1 + d\sqrt{x} + ex]$$



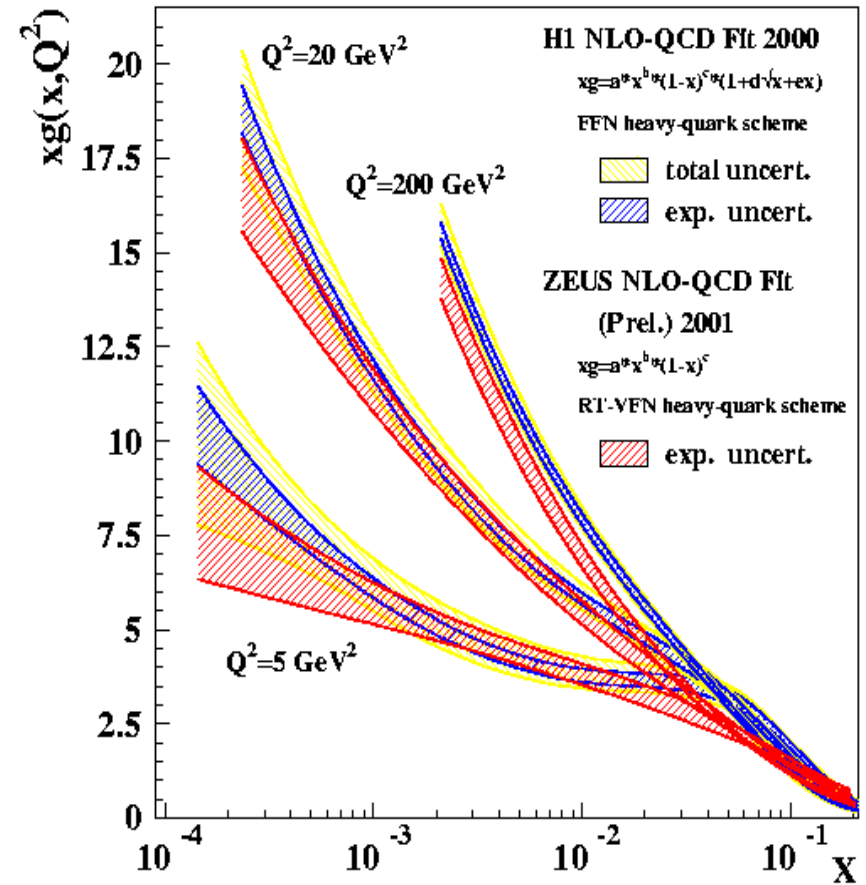
Str. fcn. and parton dist.:

GLUONS

Good agreement H1 and ZEUS



H1+ZEUS



H1

$$\alpha_s(M_Z^2) = 0.1150 \pm 0.0017(\text{exp})_{-0.0005}^{+0.0009} (\text{model}) \pm 0.005(\text{theory})$$

ZEUS (prel.)

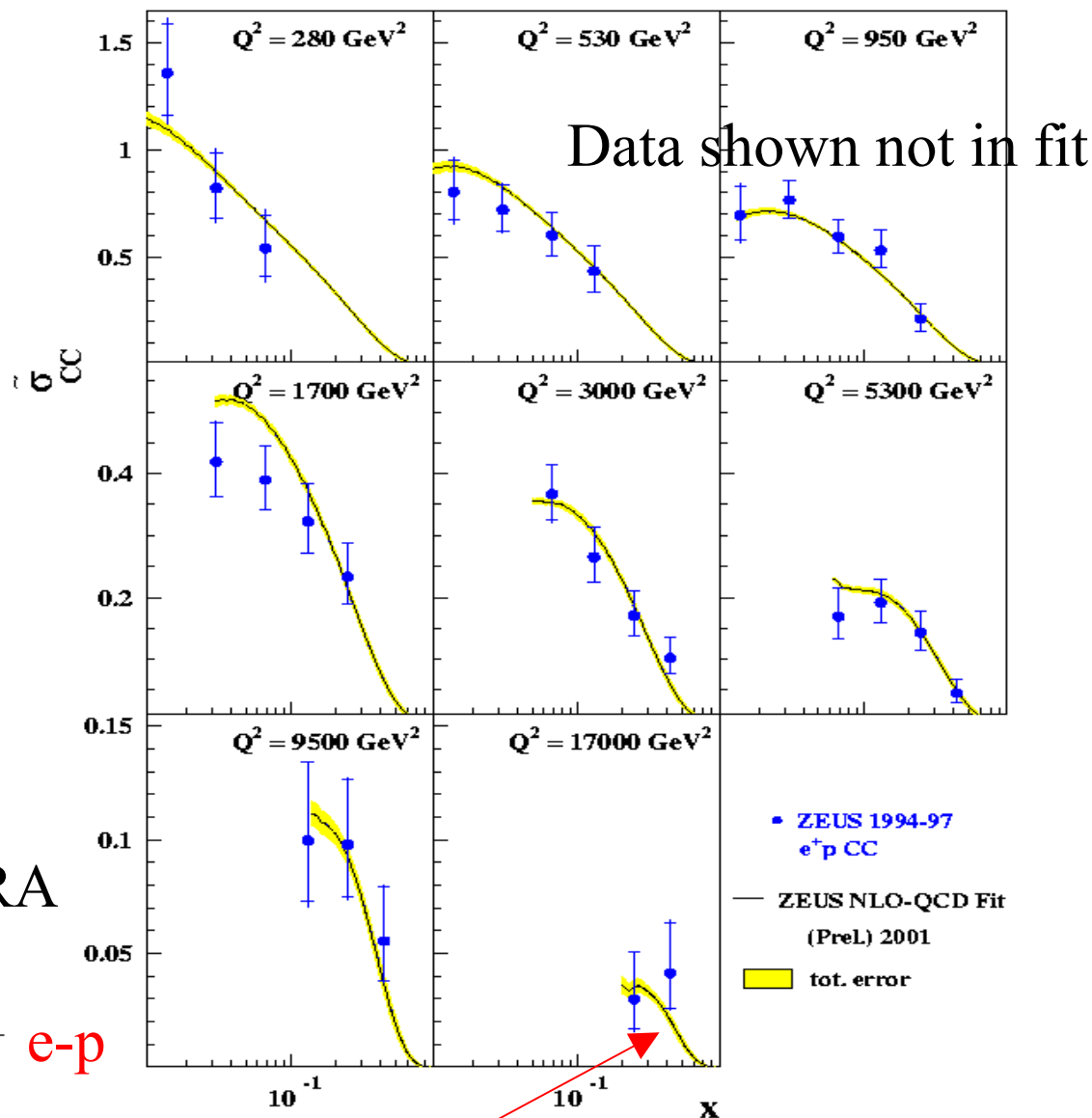
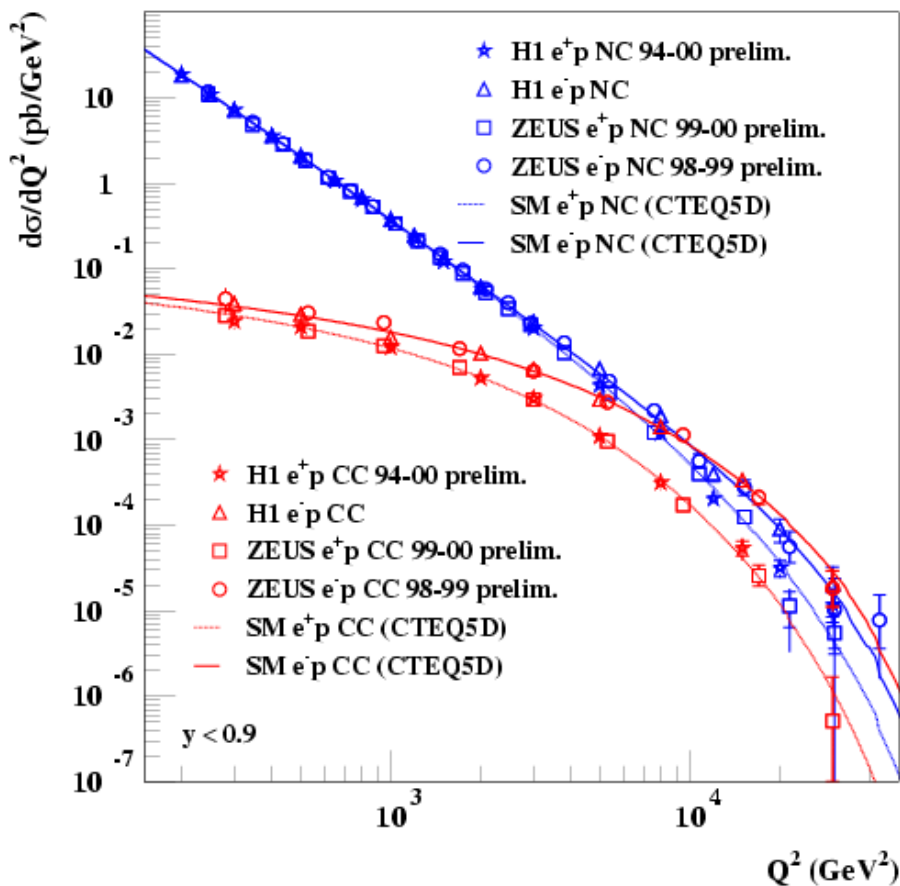
$$\alpha_s(M_Z^2) = 0.117 \pm 0.001(\text{stat} + \text{uncorr}) \pm 0.005(\text{corr})$$

theory error to be evaluated



Str. fcn. and parton dist.: Charged Current cross-section

$$\sigma_{cc}(e^+p) \sim G_F^2 \left(\frac{M_W^2}{M_W^2 + Q^2} \right)^2 [(1-y)^2(d+s) + (\bar{u} + \bar{c})]$$



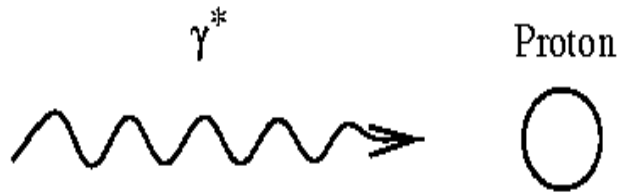
Measurements of NC and CC at HERA

Results ca. $100 \text{ pb}^{-1} e+p$ and $15 \text{ pb}^{-1} e-p$

HERA upgrade will bring
x10 data by 2005-6.

Precision of DGLAP fit
from Fixed Target exp.results





Note: Small- $x \Rightarrow$ low Q^2 (Well below 1 GeV^2 at lowest x)

DIS IN THE HADRONIC PICTURE

$W^2 \approx Q^2/x$: the γ^*p cms energy squared ("s")

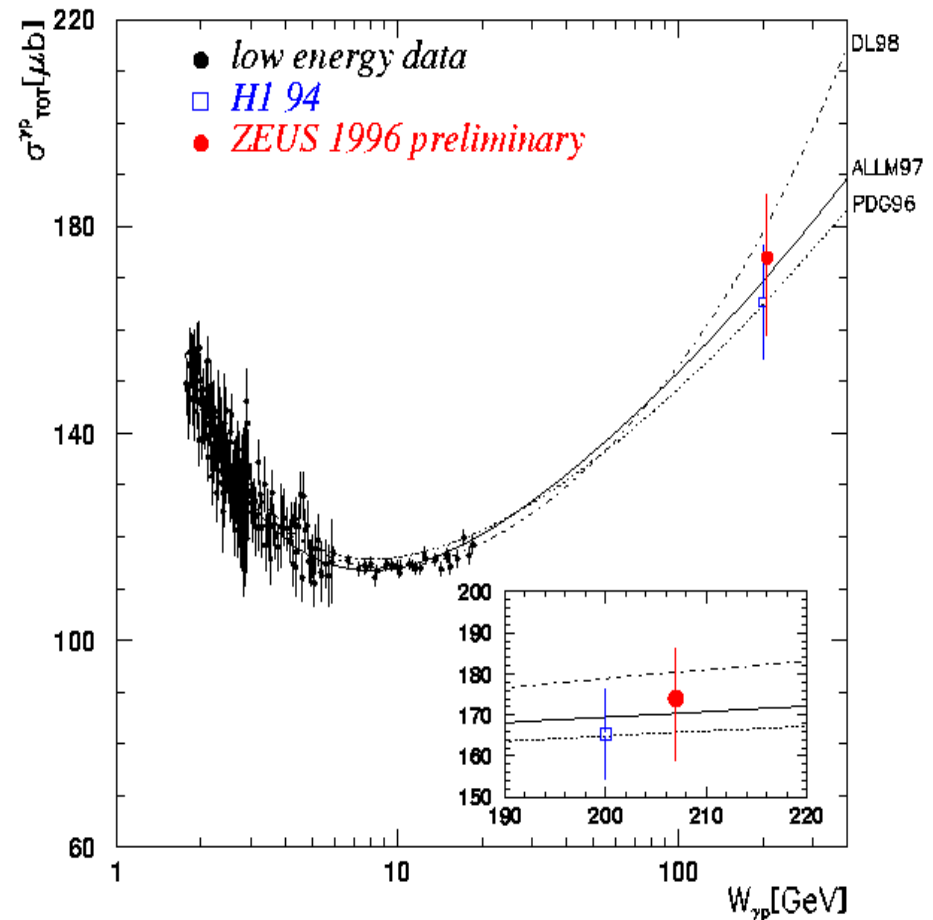
For fixed Q^2 : $x^{-\lambda} \rightarrow W^{2\lambda}$ ($F_2 \sim W^{0.4}$, $Q^2 \sim 10 \text{ GeV}^2$)

And at small- x :

$$\sigma_{tot}^{\gamma^*P}(W^2, Q^2) \approx \frac{4\pi^2\alpha}{Q^2} F_2(x \approx Q^2/W^2, Q^2)$$

- F_2 vanishes like Q^2 at low Q^2 (conservation of EM current):

- $\sigma_{tot}^{\gamma P}(W^2)$ described by Regge [$\sim W^{2(\alpha_P - 1)}$]
($\alpha_P \approx 1.08$, the "soft Pomeron") i.e. $W^{0.16}$



F_2 at $Q^2 < 1 \text{ GeV}$
 $10^{-3} > x > 10^{-6}$

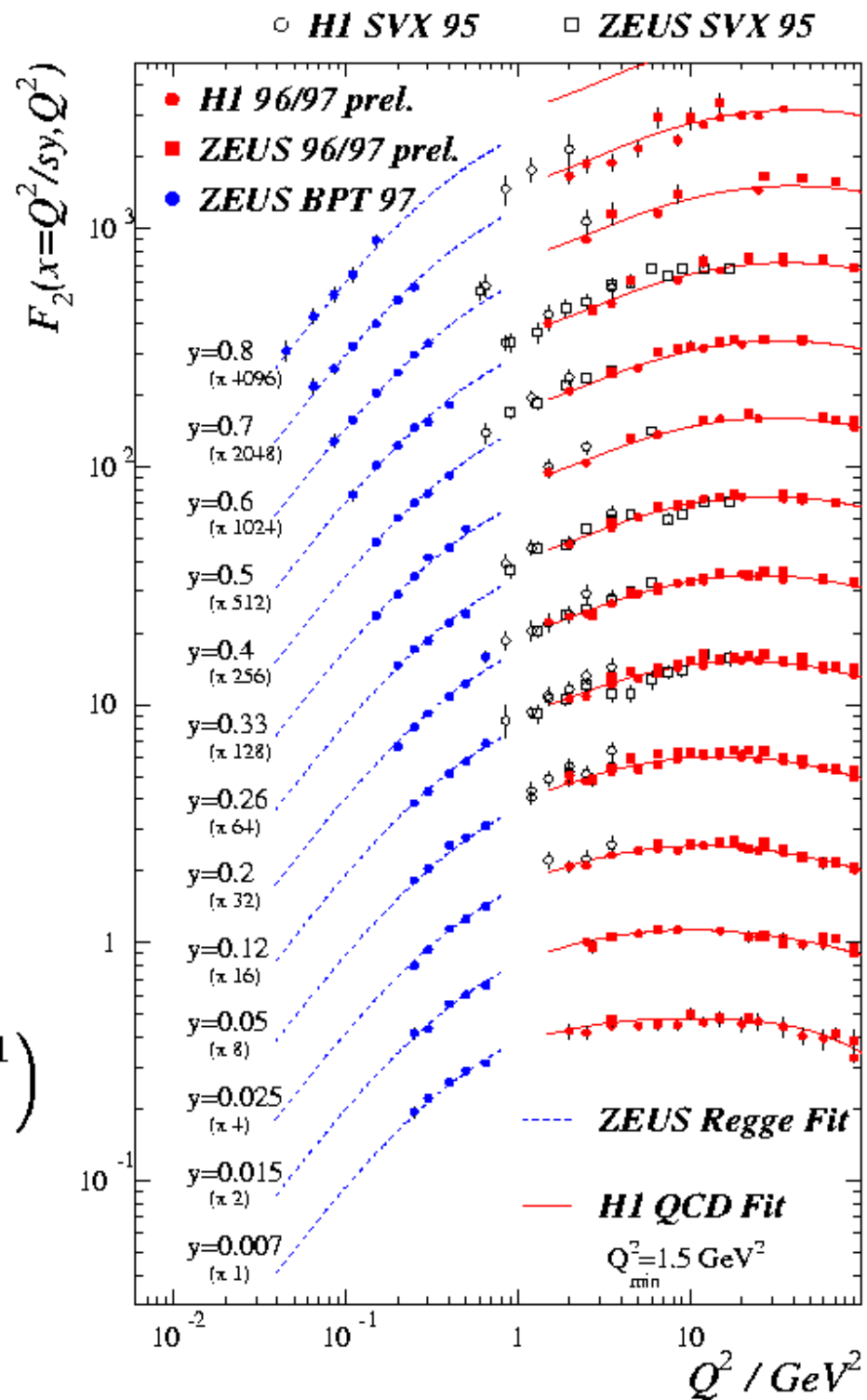
Recall $y = W^2/s$

F_2 vanishes like Q^2
at fixed W .

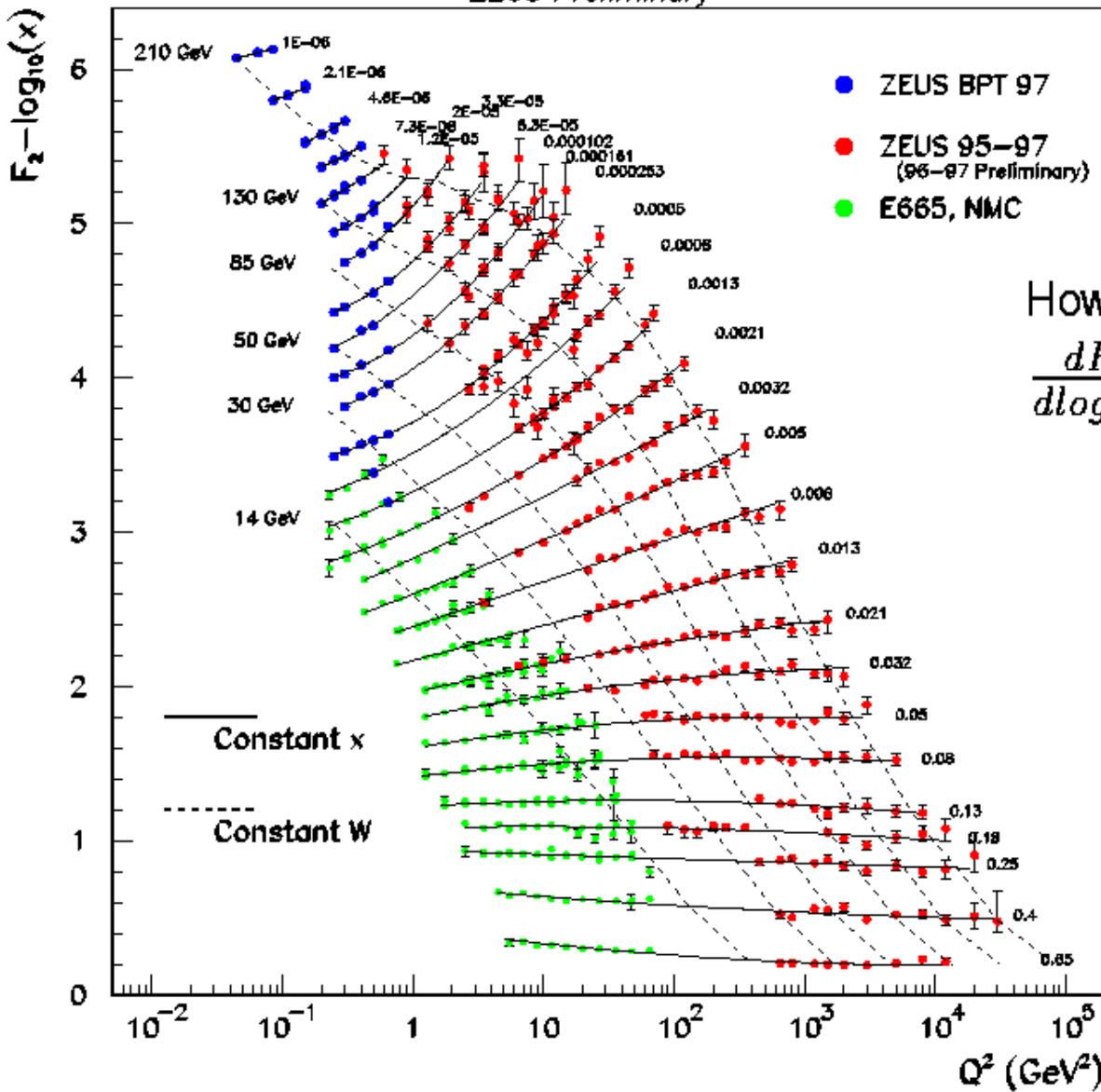
Regge fit:

$$F_2(x, Q^2) = \left(\frac{Q^2}{4\pi^2\alpha} \right) \cdot \left(\frac{M_0^2}{M_0^2 + Q^2} \right) \cdot \left(A_{\mathbb{R}} \cdot (W^2)^{\alpha_{\mathbb{R}} - 1} + A_{\mathbb{P}} \cdot (W^2)^{\alpha_{\mathbb{P}} - 1} \right)$$

$\alpha_{\mathbb{P}} \approx 1.1$ and $M_0 \approx 0.7 \text{ GeV}$.



ZEUS Preliminary



How to make the transition more visible:

$$\frac{dF_2}{d\log Q^2} \Big|_x = B + 2C(\log Q^2) \text{ at fixed } W$$

What is happening at low x and medium Q^2 ?

note: offset constant is $\log x$ i.e. no distortion

Lines are fits $A(x) + B(x)\log Q^2 + C(x)(\log Q^2)^2$



In the low Q^2 region: $F_2 \sim Q^2 \sigma_0$, i.e. vanishing like Q^2 . Therefore:

$$dF_2/d\log(Q^2) \sim Q^2 \sigma_0$$

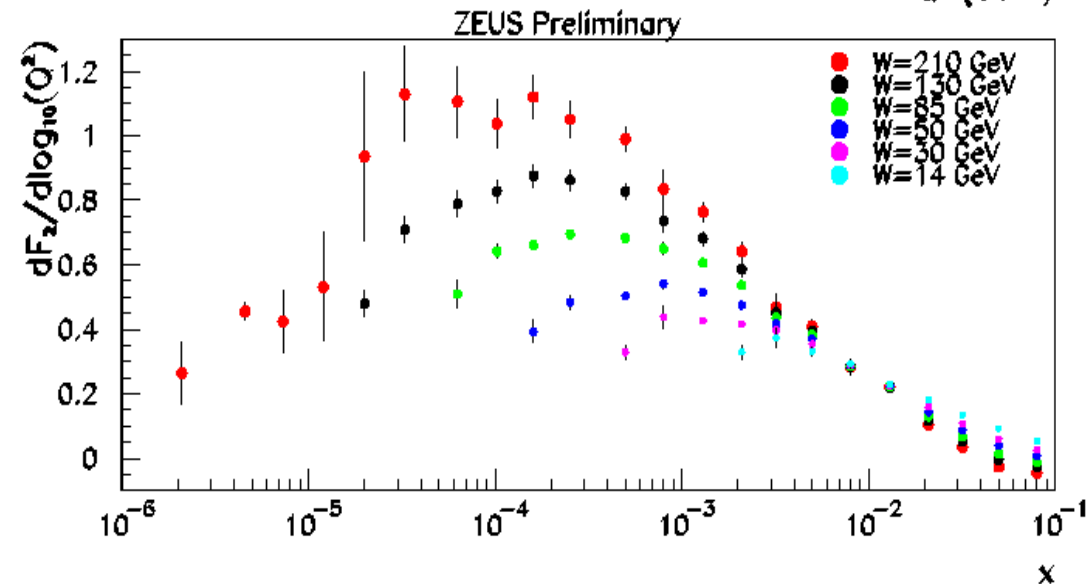
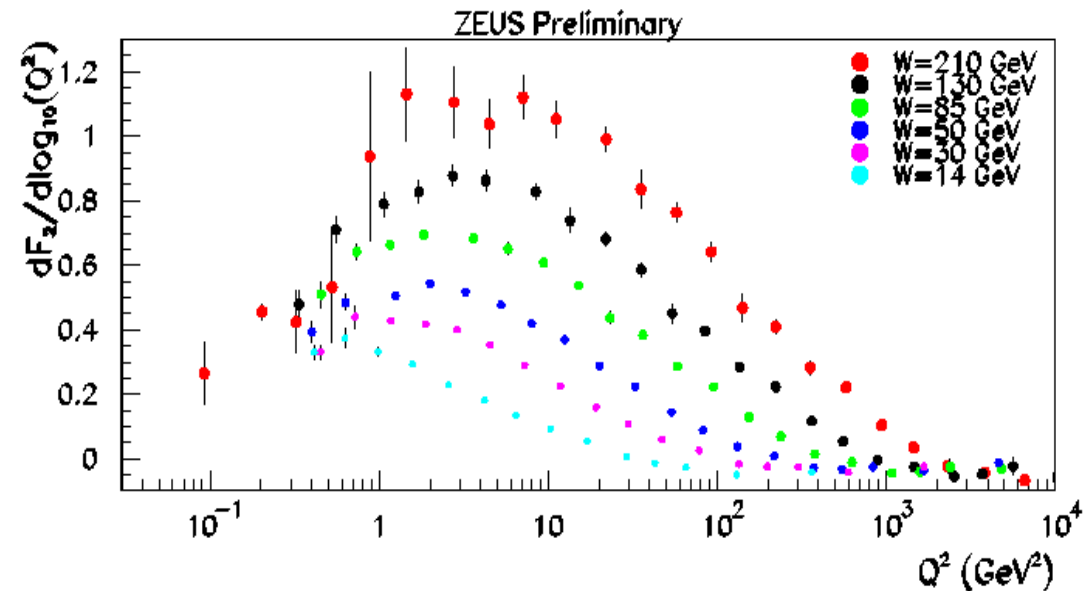
$\Rightarrow x$ is not an important variable. Q^2 is the important variable in the left hand side of the plots.

Naively, in the region where GLAP applicable (Q^2 large):

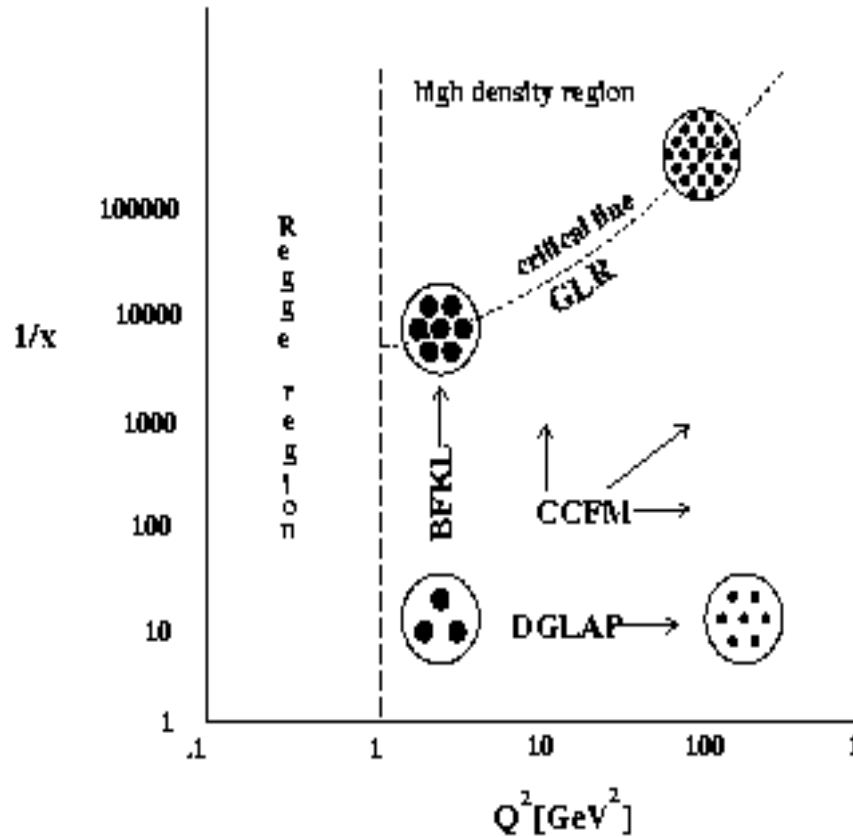
$$dF_2/d\log(Q^2) \sim xg \sim x^{-\lambda}$$

IF λ is a slow function of x and Q^2 .

$\Rightarrow Q^2$ is not an important scale. x is the important variable in the right hand side of the plots.



Beyond DGLAP: pert. QCD considerations



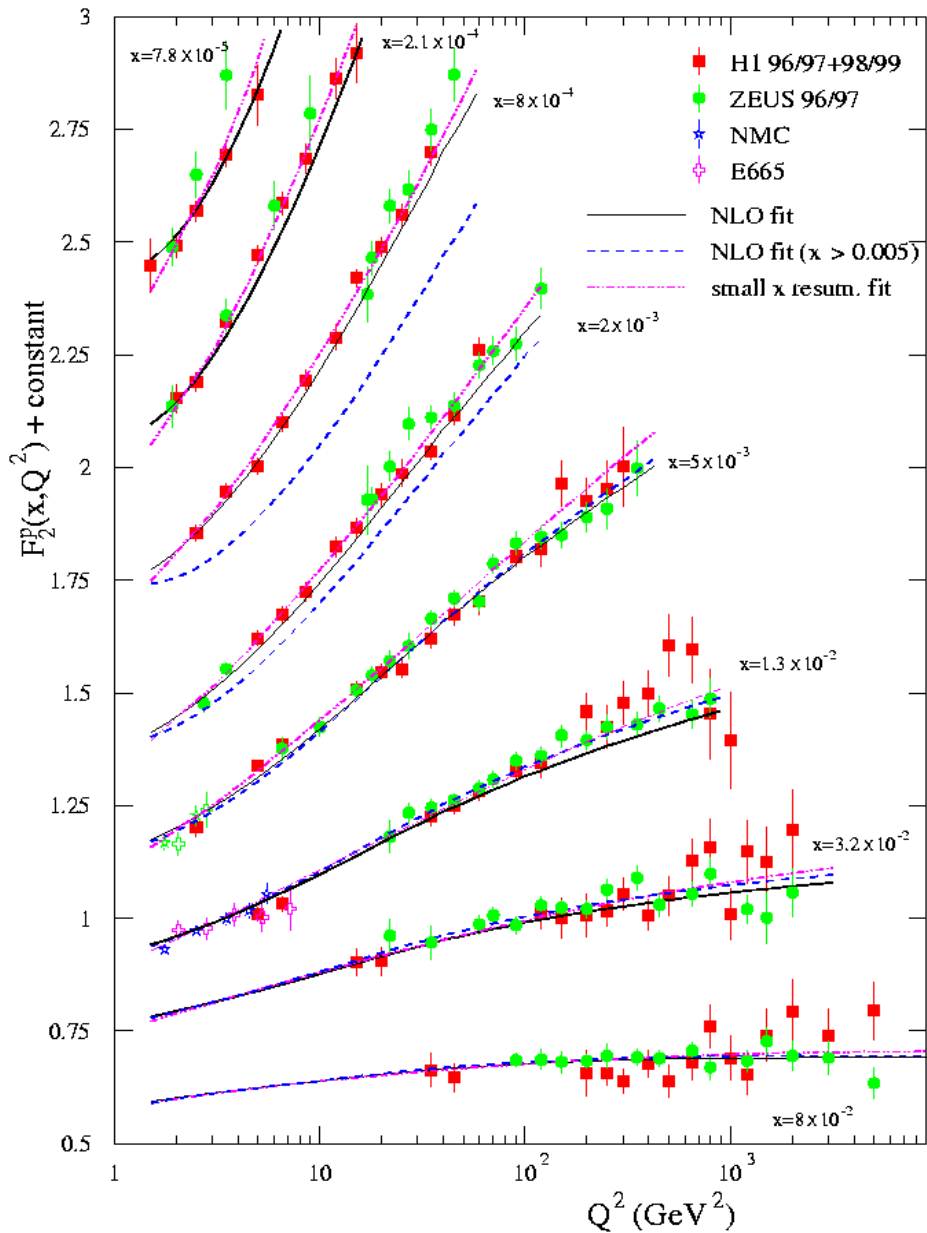
- GLAP fits ($\ln Q^2$ resummation of QCD) can describe data well above $Q^2 > 1 \text{ GeV}^2$
- Regge fits can describe data well $Q^2 < 1 \text{ GeV}^2$

- What about $\ln 1/x$ (BFKL,CCFM), higher twist, terms etc.?
- What about shadowing or saturation (GLR): i.e. effects from high density of partons at low-x?



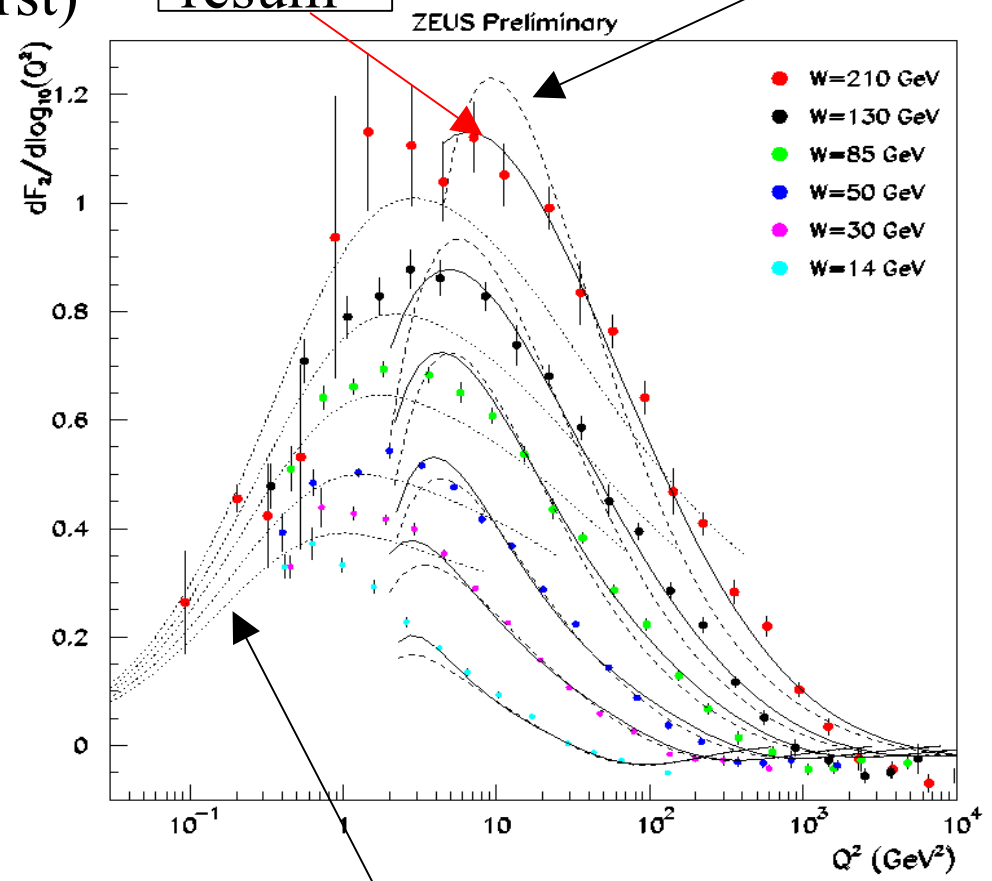
Small x resummation (not full BFKL):

Martin, Roberts, Sterling, Thorne(mrst)



small-x
resum

DGLAP



Dipole model



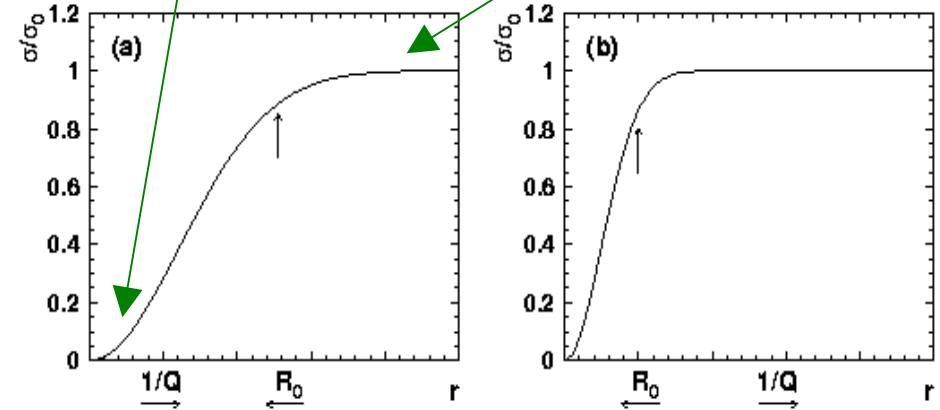
DIPOLE models: a class of phenomenological models

example: Golec-Biernat & Wuesthoff

- Proton dissociates to a $q\bar{q}$ pair (dipole) upstream of the protons

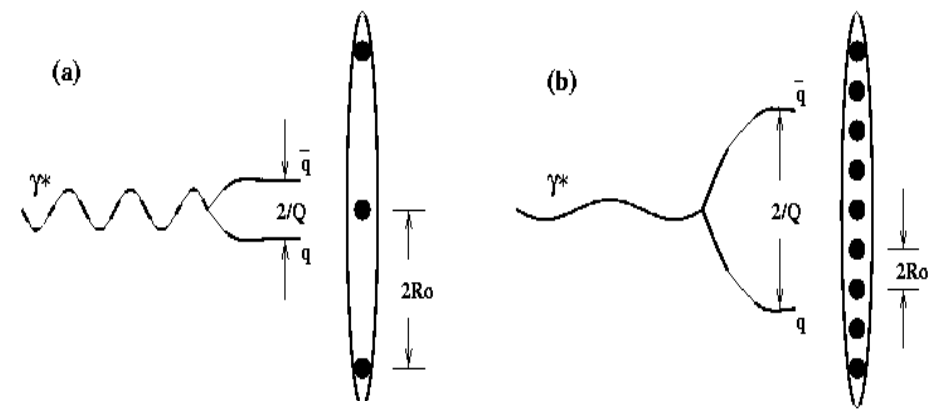
DGLAP-like

saturation



- $\hat{\sigma}_{dipole}$ increases as r^2 at small r . \Rightarrow Color transparency.
- $\hat{\sigma}_{dipole}$ becomes constant (saturates) at large r .

- At what r the cross section saturates depends on the density of the partons, i.e. x . $\rightarrow \hat{\sigma}_{dipole}(x, r)$
 $[R_0 \sim 1/gluon \sim x^\lambda]$

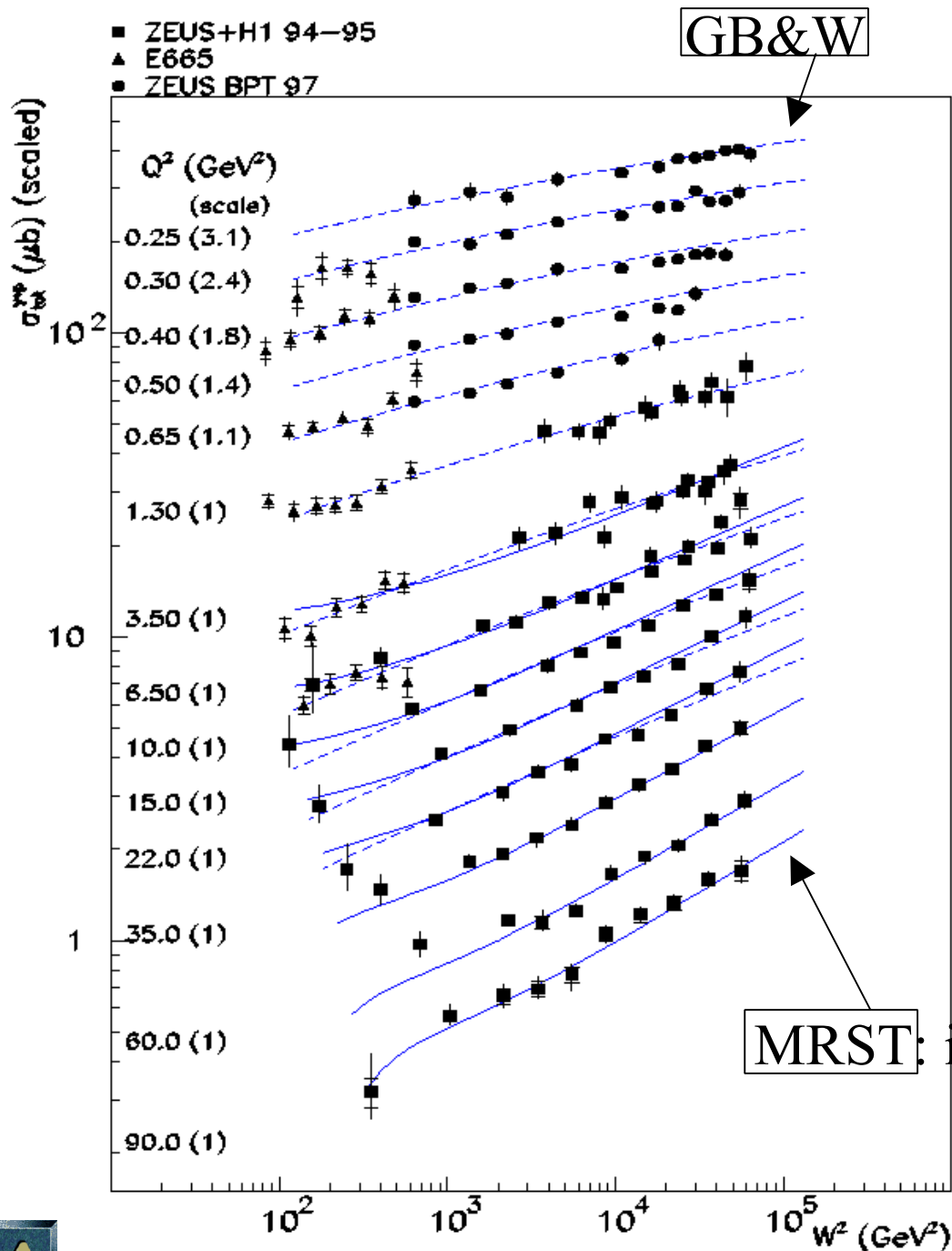


DIPOLE formulation of DIS

$$\sigma = \int d^2r \int dz |\Psi(z, r)|^2 \hat{\sigma}_{dipole}(x, r)$$



GB&W Description of $\sigma_{tot}^{\gamma^* p}$

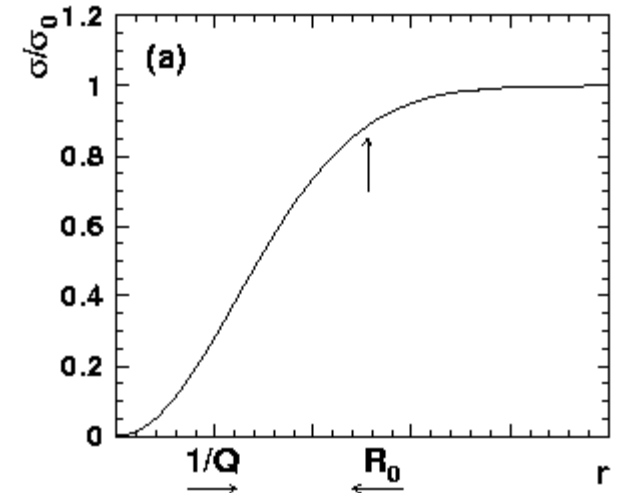


Recall:

$$W^2 \approx Q^2 / x$$

$$\sigma_{tot}^{\gamma^* p}(W^2, Q^2) \approx \frac{4\pi^2\alpha}{Q^2} F_2(x \approx Q^2/W^2, Q^2)$$

At low x (high W) medium Q^2 , dipole cross-section away from DGLAP-like region is being probed.

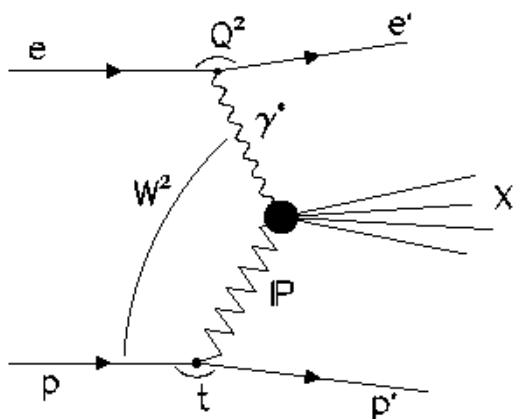


On the other hand, pure DGLAP fits also describe the data.

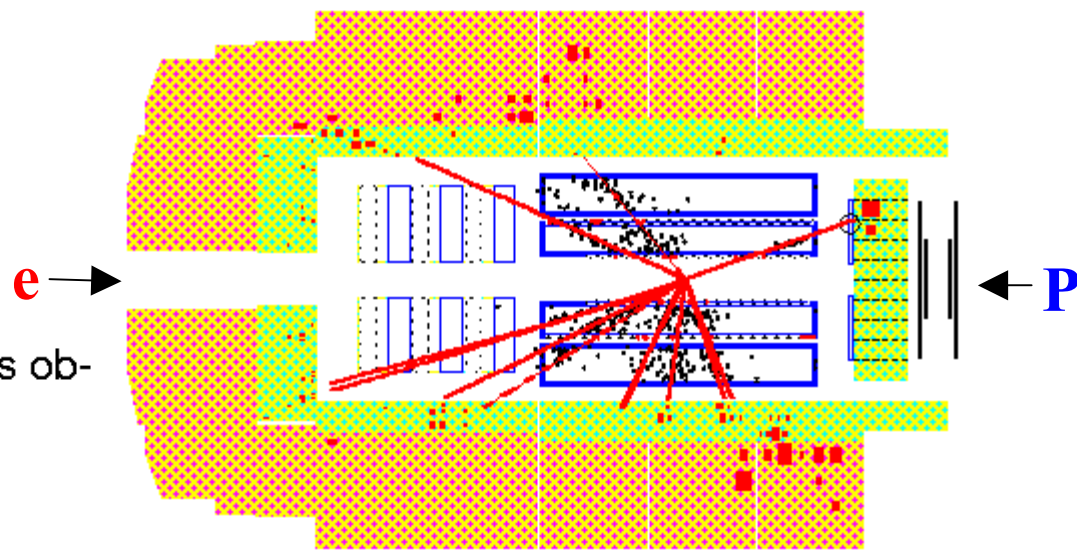


DIS diffraction at HERA:

At HERA diffractive dissociation of virtual photons is observed ($\approx 10\%$): can we understand this in terms of pQCD?



Exchange of a colorless object. (Pomeron)



Define 3 more variables:

- x_P : momentum fraction taken by the pomeron.
- β : momentum fraction of the struck parton in the pomeron.
- t : momentum transfer at the proton vertex.

- Beyond the simple GLAP picture : Need to interact with 2 gluons, at least!
- What is the relationship of diffraction with F_2 at small- x ?



The diffractive cross section is written as:

$$\frac{d^3\sigma^D}{d\beta dQ^2 dx_{\text{IP}}} = \frac{4\pi\alpha^2}{\beta Q^2} \left(1 - y + \frac{y^2}{2}\right) F_2^{D(3)}(\beta, Q^2, x_{\text{IP}})$$

(integrated over t)

Hard factorization has been proven (Collins 1997):

$$F_2^D \sim f^D \otimes \hat{\sigma}$$

Where f^D is the diffractive parton densities.

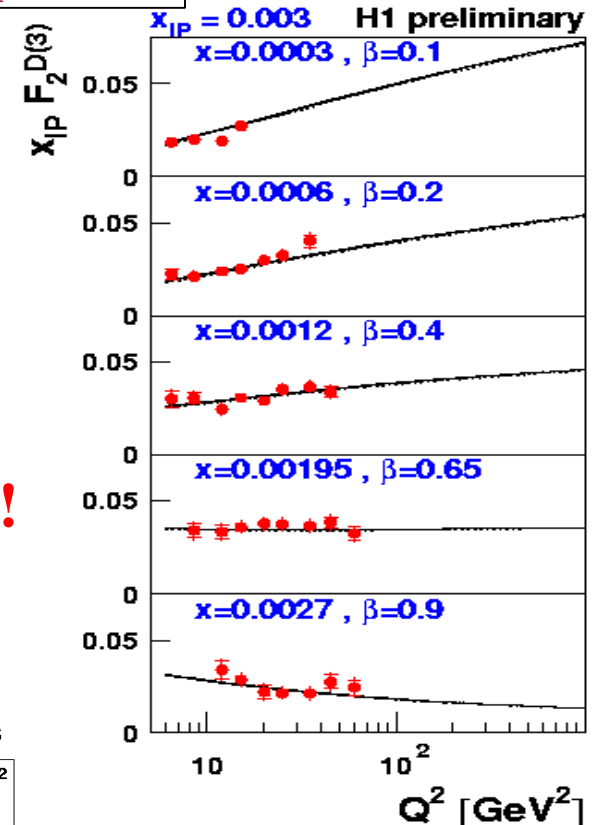
If Regge factorization is assumed:

$$F_2^D(x_{\text{IP}}, t, Q^2, \beta) = f(x_{\text{IP}}, t) \cdot F_2^{\text{IP}}(\beta, Q^2)$$

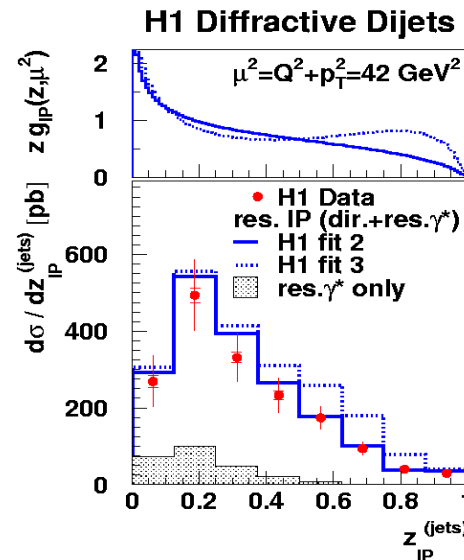
⇒ Can do a GLAP analysis of F_2^{IP} , the **universal** (for DIS) Pomeron structure function. (More from Max.)

$f(x_{\text{IP}}, t) \approx 1/x_{\text{IP}}$, the Pomeron flux factor.

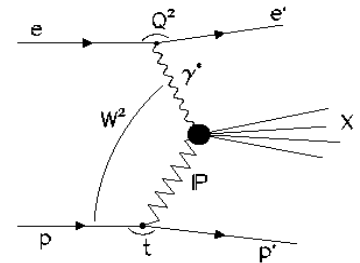
The flux factor must be telling us about gluon correlation in the proton!



Predictive power!



● H1 (prel.)
 — QCD fit (IP+IR)
 - - - QCD fit (IP)



However....

⇒ The origin of the Pomeron flux unexplained.

⇒ What is the relationship of σ_{diff} to σ_{tot} ?

• Diffraction has the same W^2 (or x) dependence as the total cross section! ($\sigma_{diff} \sim W^{0.4}$, $Q^2 \sim 10 \text{ GeV}^2$)

• Contradicts naive optical theorem expectation

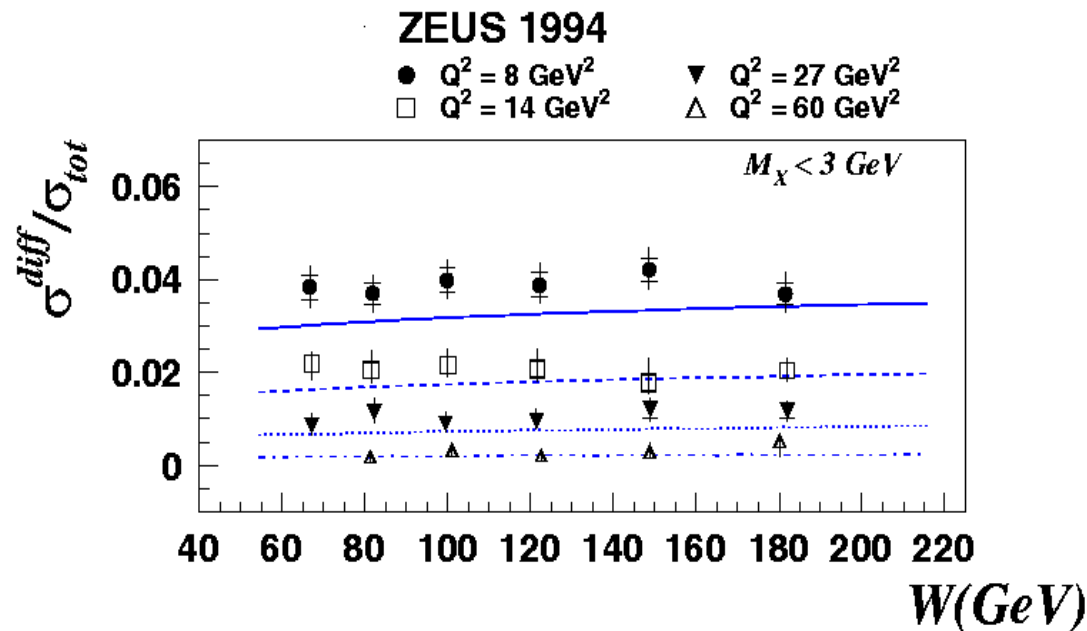
⇒ $\sigma_{\gamma^*p}^{tot} \sim W^a$ then $\sigma^{diff} \sim W^{2a}$

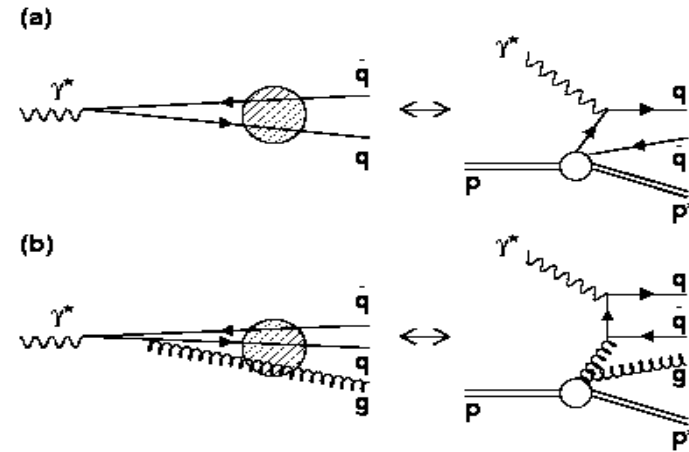
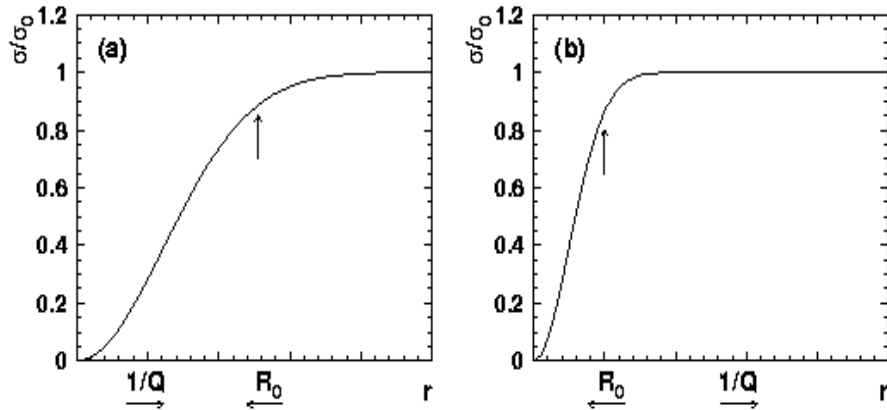
• Contradicts expectation—if $\sigma_{\gamma^*p}^{tot} \sim \text{Gluon}$, then $\sigma^{diff} \sim \text{Gluon}^2$??

• ⇒ Also contradicts expectation of Regge:

$\alpha_P = 1.08$ (i.e. $\sim W^{0.16}$)

The lines in the figure are from the saturation model of Golec-Biernat and Wüsthoff

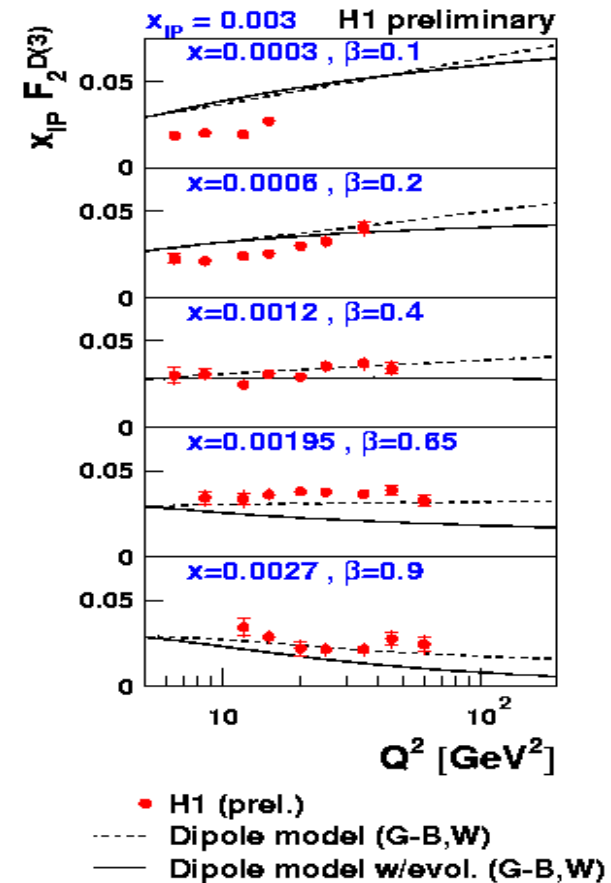




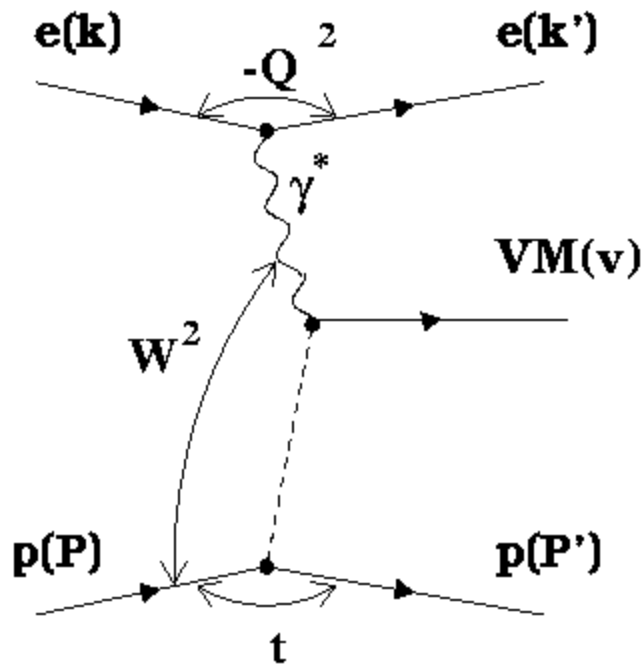
$$\sigma_{diff} \sim \int d^2r \int dz |\Psi_\gamma|^2 \hat{\sigma}_{dipole}^2$$

The dipole model qualitatively describe the data: i.e. give the right behavior for the "pomeron structure"

→ **Simple physical picture :
effective correlation of partons
due to interplay of photon size
and parton separation.**



Elastic Vector Meson Production at HERA:



- Very similar process to the inclusive diffraction

- What is the W dependence?

Recall:

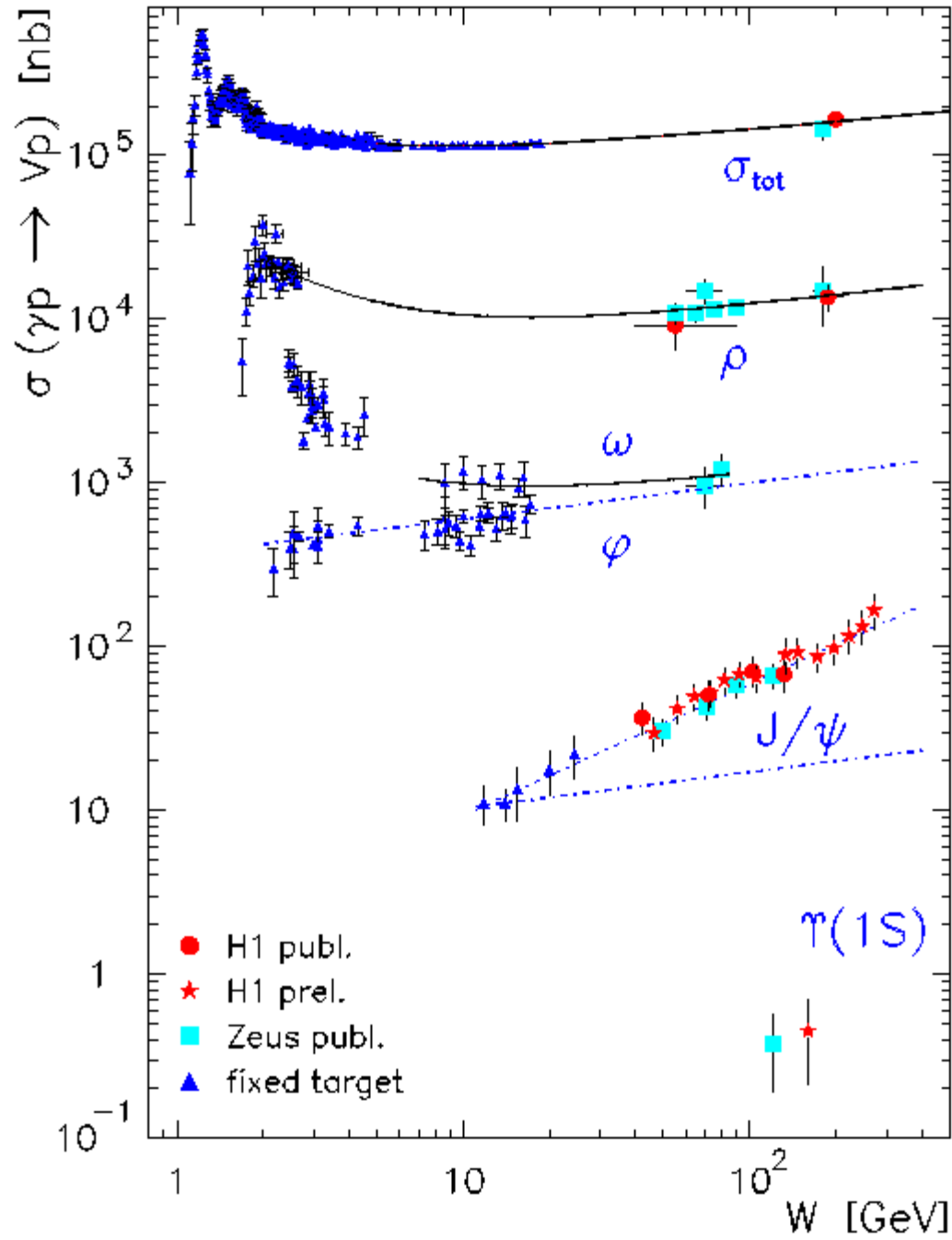
$$\sigma_{\gamma p}^{tot} \sim W^{2(\alpha_{\mathbb{P}} - 1)} \approx W^{2(0.08)} = W^{0.16}$$

$$F_2^P \text{ (or } \sigma_{\gamma^* p}) \text{ (at } Q^2 \approx 10 \text{ GeV}^2) \sim x^{-0.2}$$

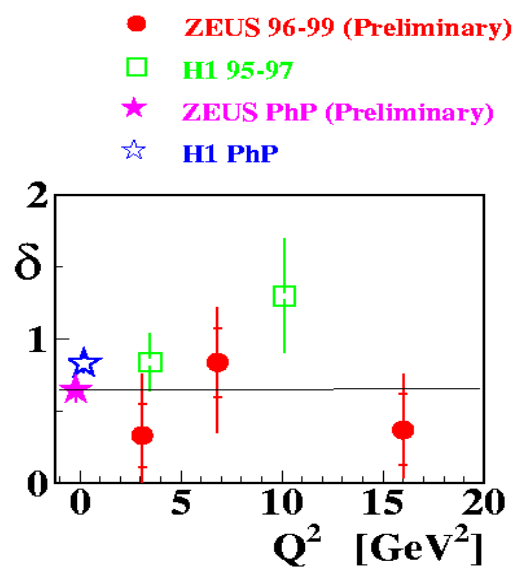
$$\Rightarrow W^{2(0.2)} = W^{0.4}$$

$$\sigma^{diff} / \sigma_{\gamma^* p} \text{ constant.}$$

- What about Vector Meson production \Rightarrow



J/ψ rises like $W^{0.8}$! (i.e. twice F_2 at 10 GeV^2)

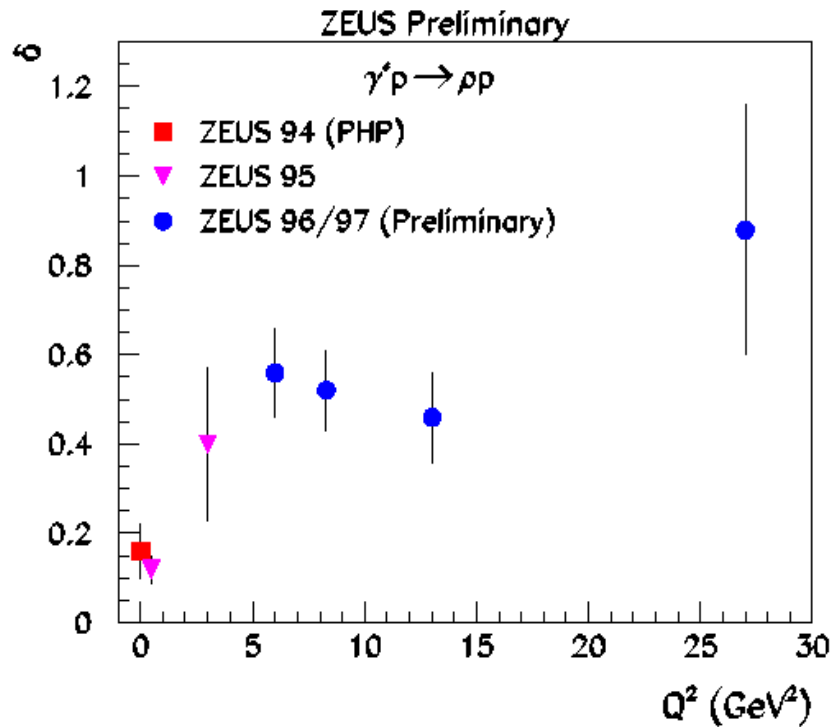


$\sigma \propto W^\delta$

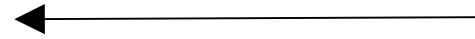
Still large errors: slope (δ) consistent with being constant as function of Q^2 .

What about the Q^2 dependence of ρ 's? \Rightarrow

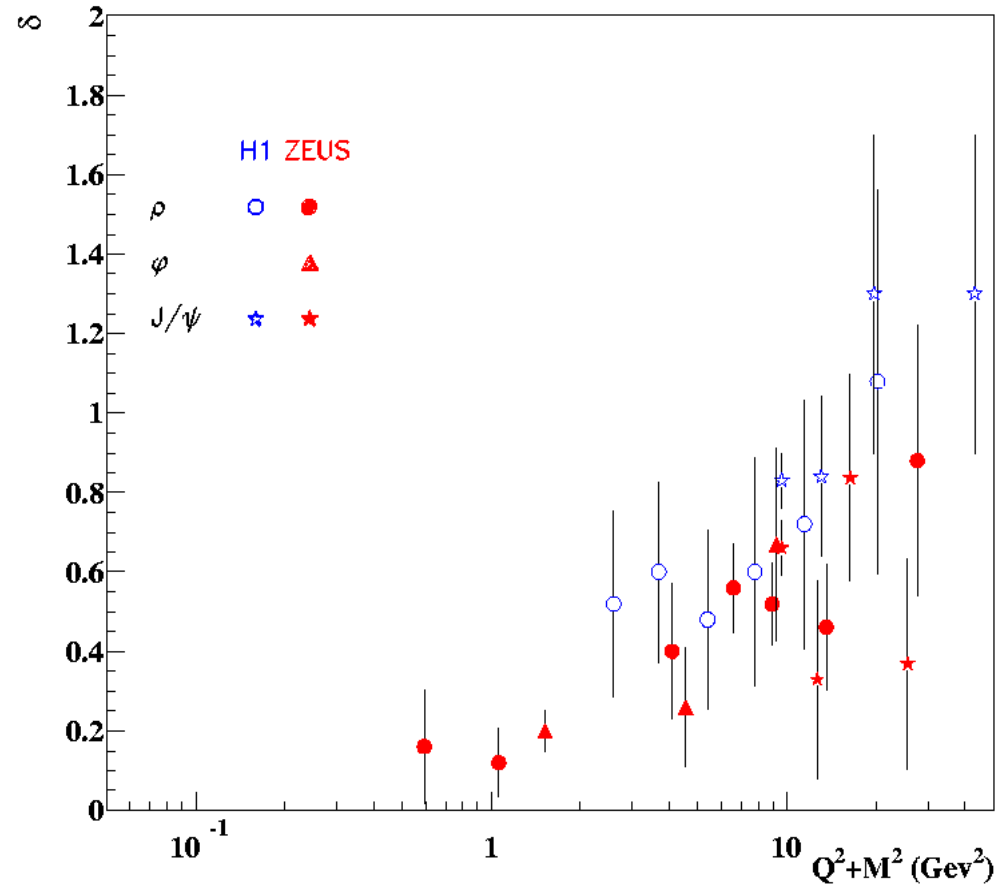




ρ mesons: increasing slope with Q^2 .

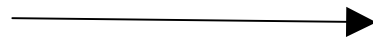


Elastic VM Production at HERA



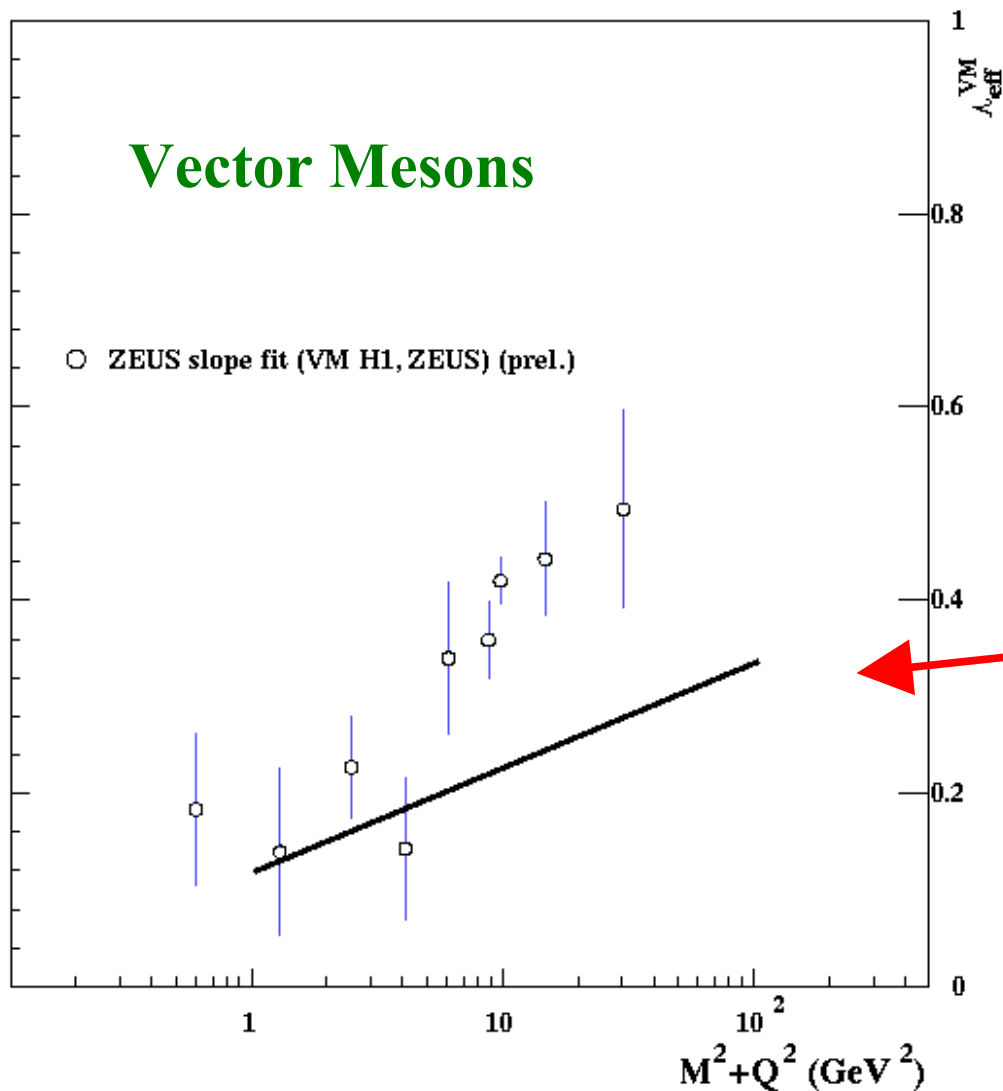
All vector mesons:

The rise, δ , scales with $Q^2 + M_{VM}^2$!

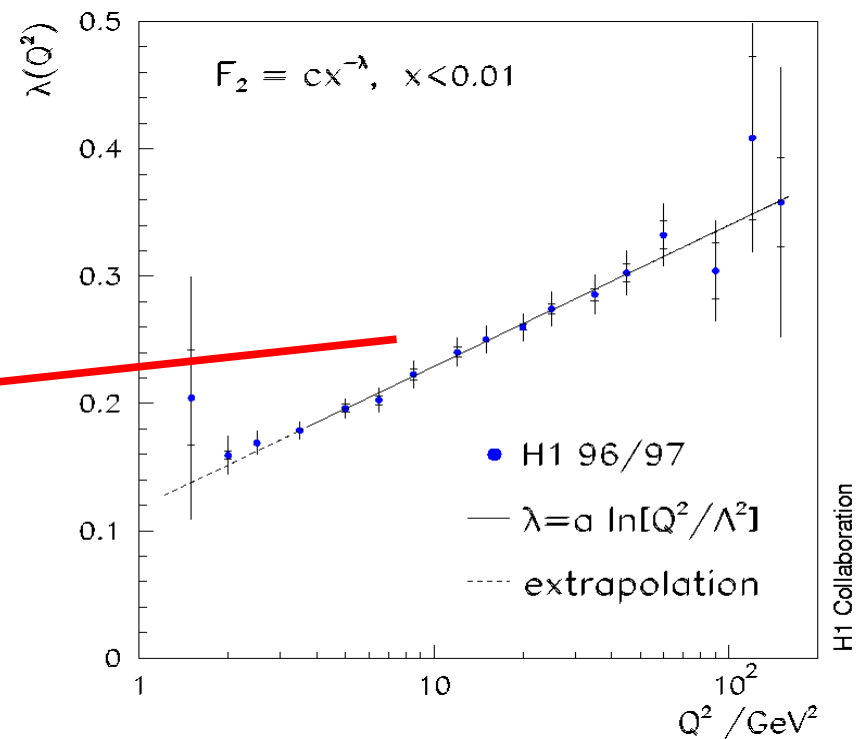


Now compare to F2 (i.e. inclusive DIS)

$$x^{-\lambda} \rightarrow W^{2\lambda}$$



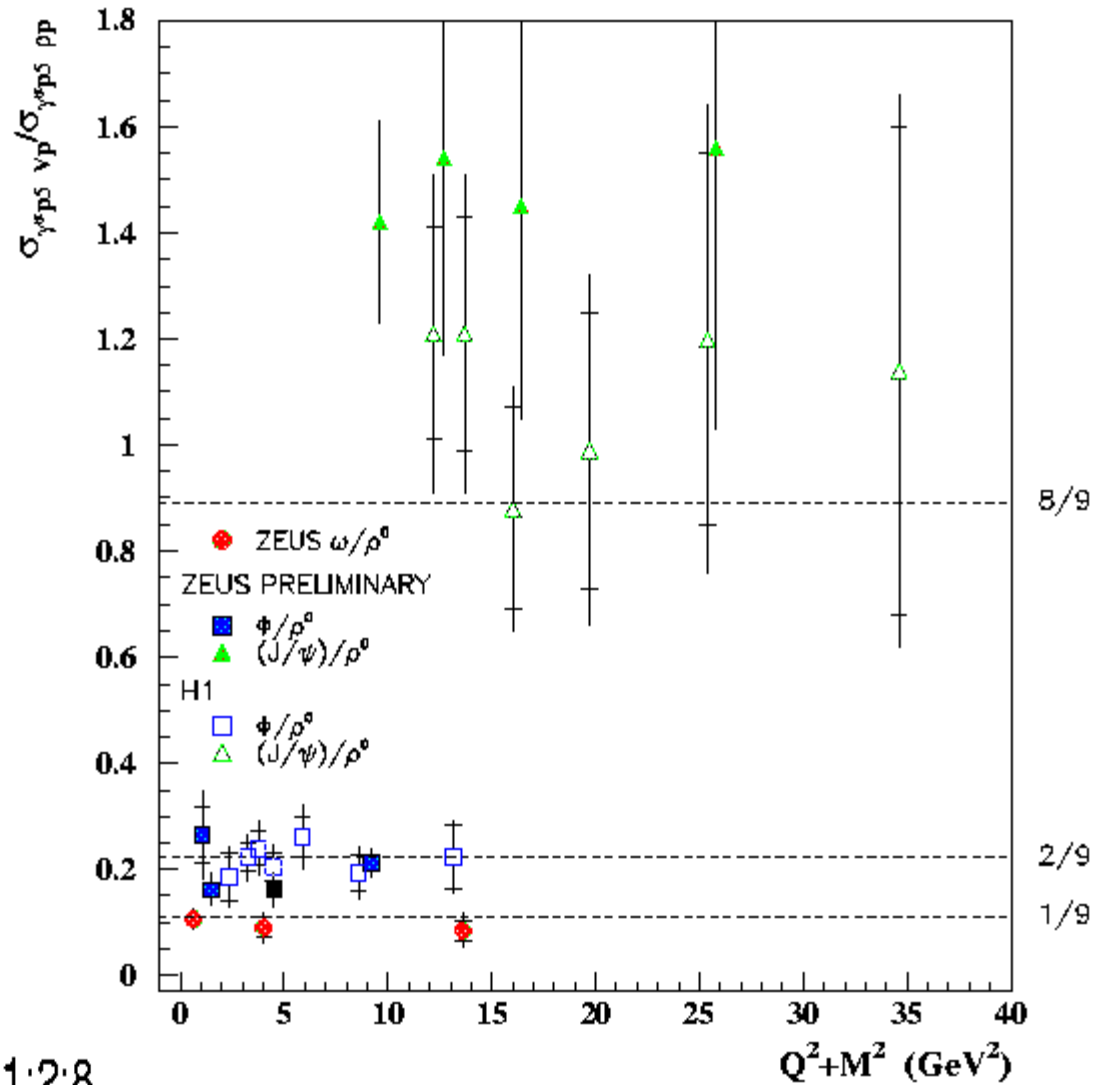
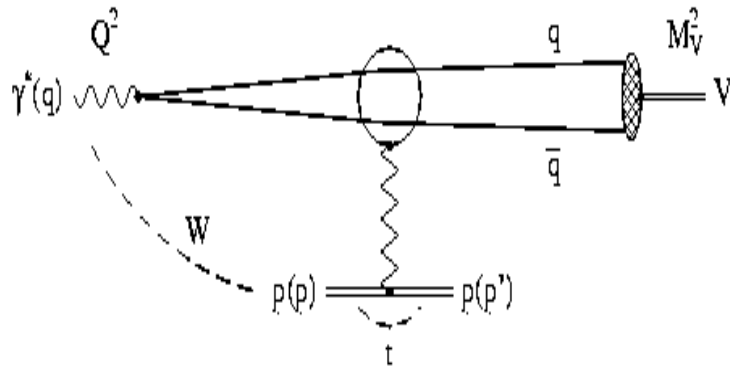
Recall F2 behavior



Elastic VM production at HERA

Naively: $\sigma_{VM} \sim f_V^2 \cdot \hat{\sigma}$

i.e. coupling times the (dipole?) cross section $\hat{\sigma}$?



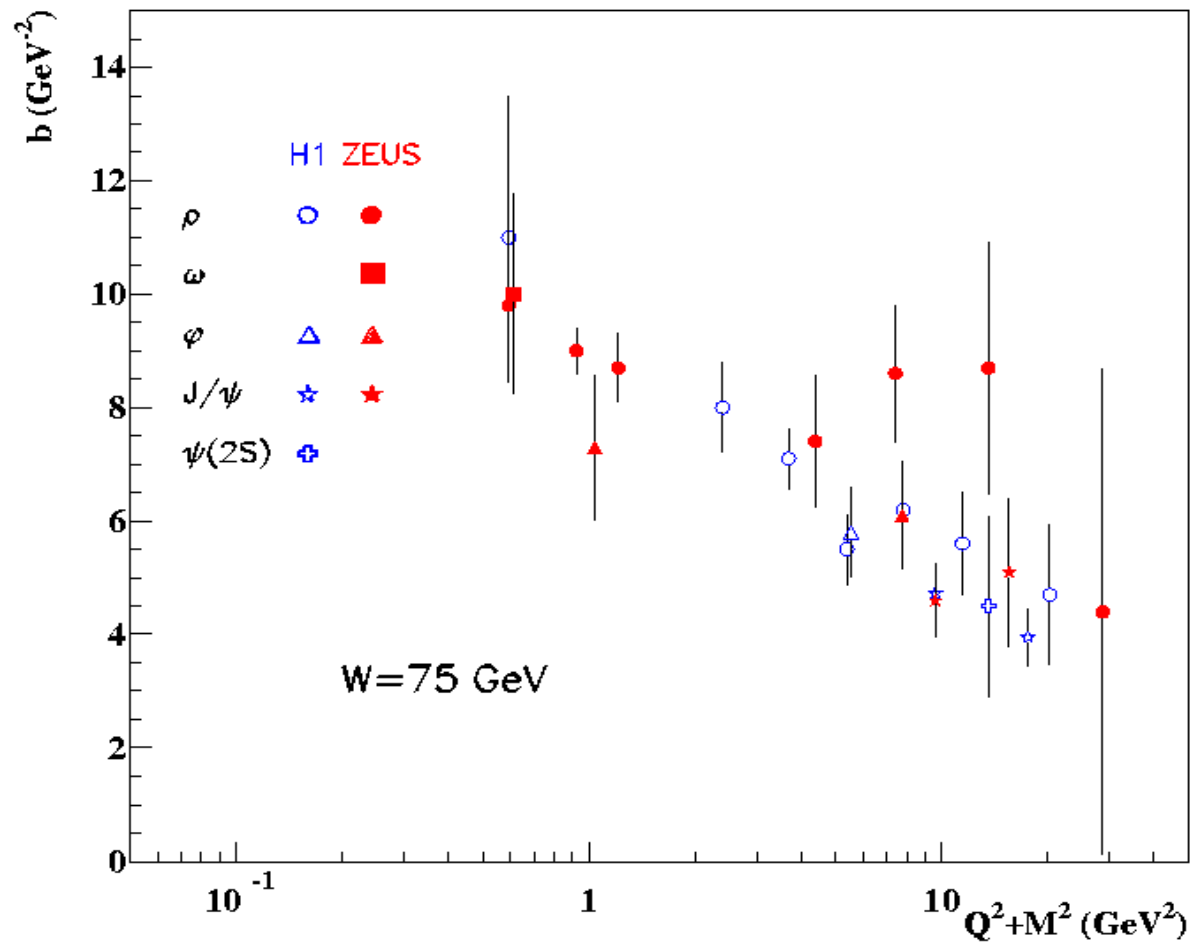
f_V^2 from charges of the quarks: $\rho:\omega:\phi:J/\psi \rightarrow 9:1:2:8$.



... and $\hat{\sigma}$ that depends on the dipole radius?

Look at the t slope: $\frac{d\sigma_{VM}}{d|t|} \propto e^{-b|t|}$

Elastic VM Production at HERA



Consistent with $Q^2 + M_{VM}^2$ giving the “size” of the interaction.



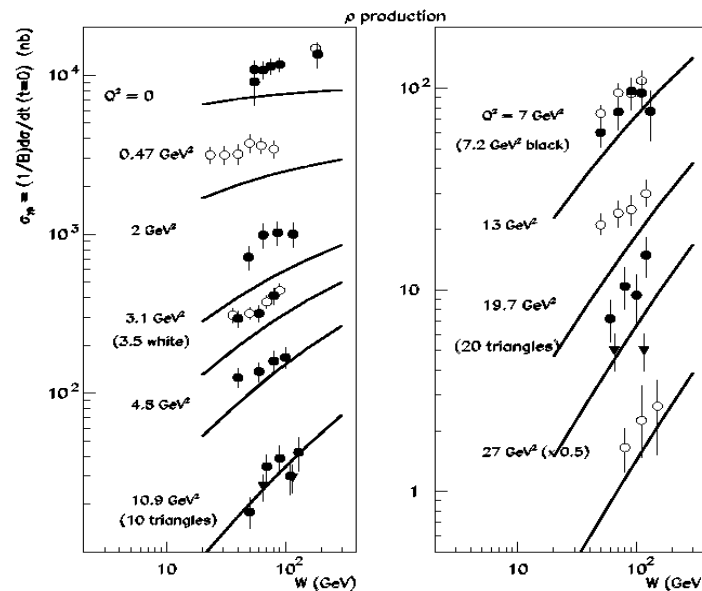
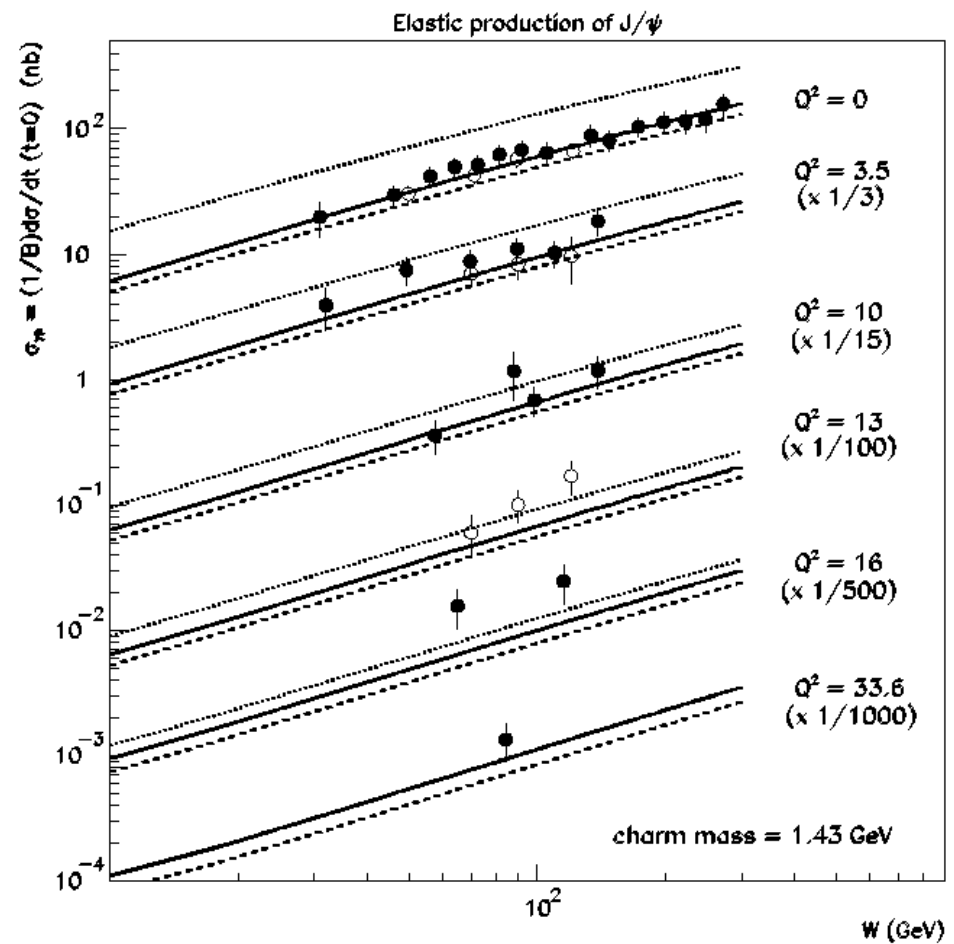
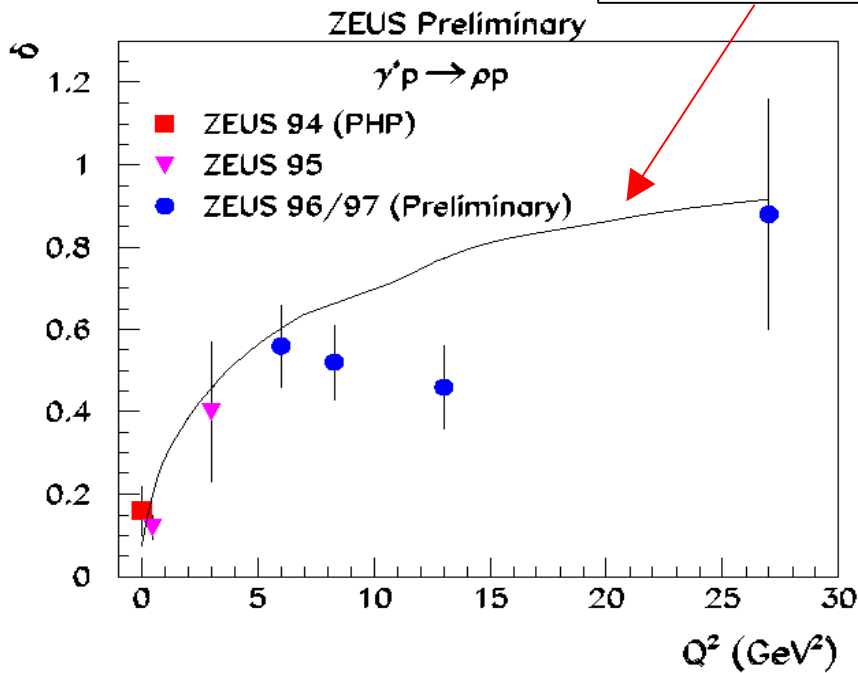
Dipole formulation of the VM cross section \Rightarrow

Amplitude: $A = \int dz \int d^2r \Psi_{VM} \hat{\sigma}_{dipole} \Psi_{\gamma}$

Ψ_{VM} : The vector meson wave function.

Ψ_{γ} : The photon wave function.

GB&W dipole



From paper
by Caldwell
and Soares



Conclusions and outlook:

Only a part of HERA results on proton structure was presented here.

The measurements of proton structure function F_2 at of 3% accuracy over 6 orders magnitude.

-leads to precise determinations of QCD parameters: parton distributions, strong coupling constant.

Looking more closely at small x :

Connection has to exist between small- x and diffraction.

And...



-A simple dipole model can qualitatively demonstrate the connection between DIS, diffraction, and VM production.

-May point the way to bridge the gap, in understanding, between DGLAP, and, low-x and low-Q² regions.

HERA detectors, will begin to take data again after the luminosity upgrade of the accelerator:

-Will concentrated on high-Q², high-x, physics.

Meanwhile..a large amount of HERA data on low-x and diffraction is available.

