



Hadron Structure Approach to the Running Coupling at Low Energy

Aurore Courtoy
IFPA, Universite de Liege

Workshop on Confinement Physics
March 12-15, 2012, Jefferson Lab

Outline

talk in collaboration with S. Liuti

- Strong Coupling Constant
 - Perturbative determination
 - Non-perturbative approaches
- Hadron Structure Phenomenology
 - Final State Interaction and Parton Distribution Functions
 - Parton-Hadron Duality
- Nonperturbative QCD coupling from Phenomenology PRELIMINARY

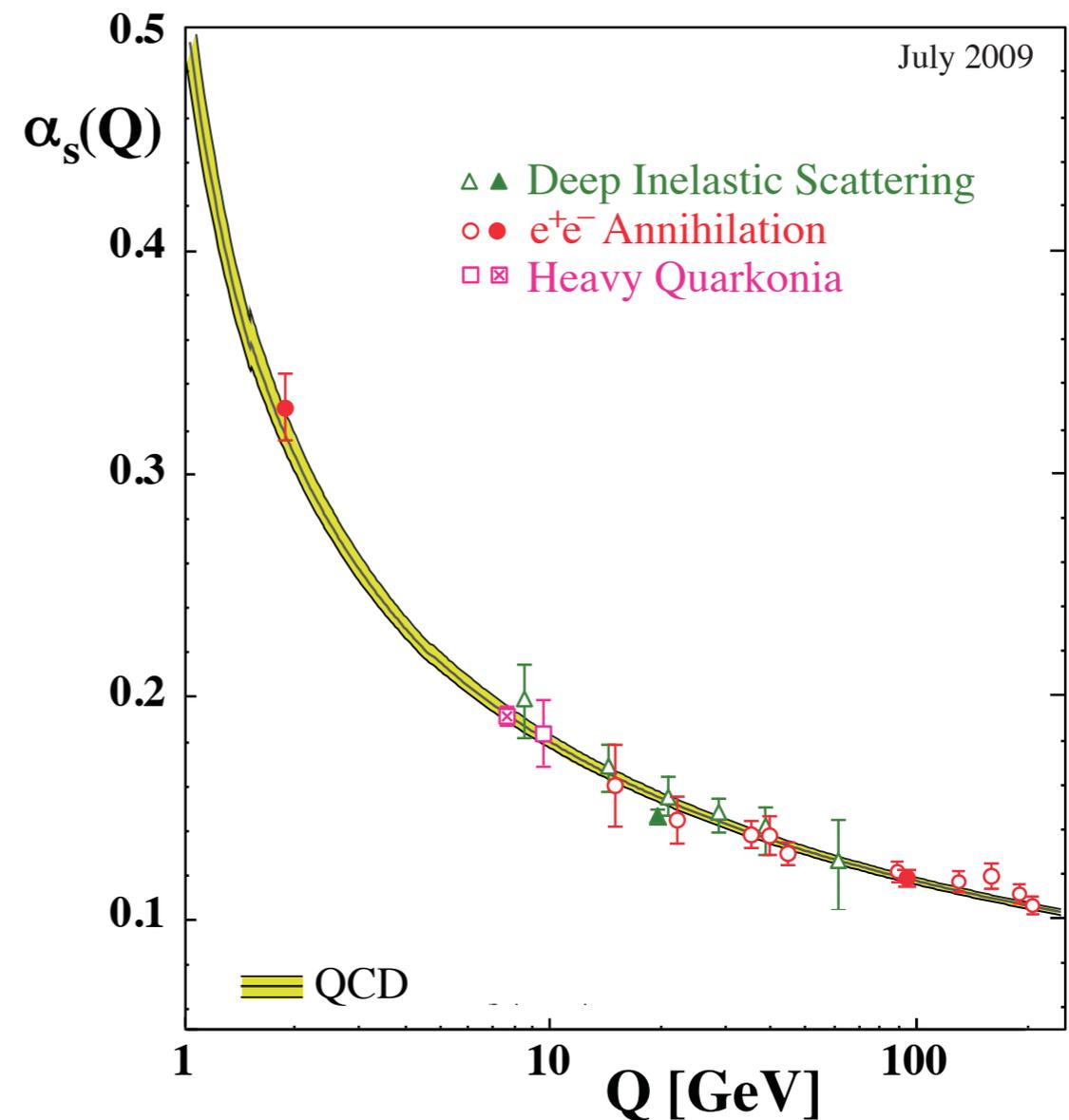
QCD Coupling Constant in pQCD

- QCD with massless quarks
 - ➔ no scale parameters
- RGE introduces a momentum scale Λ
 - ➔ interaction strength =1
- Renormalization scheme dependence of Λ
- World data average (2009)

$$\alpha_s(M_{Z^0}) = 0.1184 \pm 0.0007$$

that corresponds to

$$\Lambda_{\overline{MS}}^{(5)} = (213 \pm 9) \text{ MeV}$$



QCD Running Coupling Constant

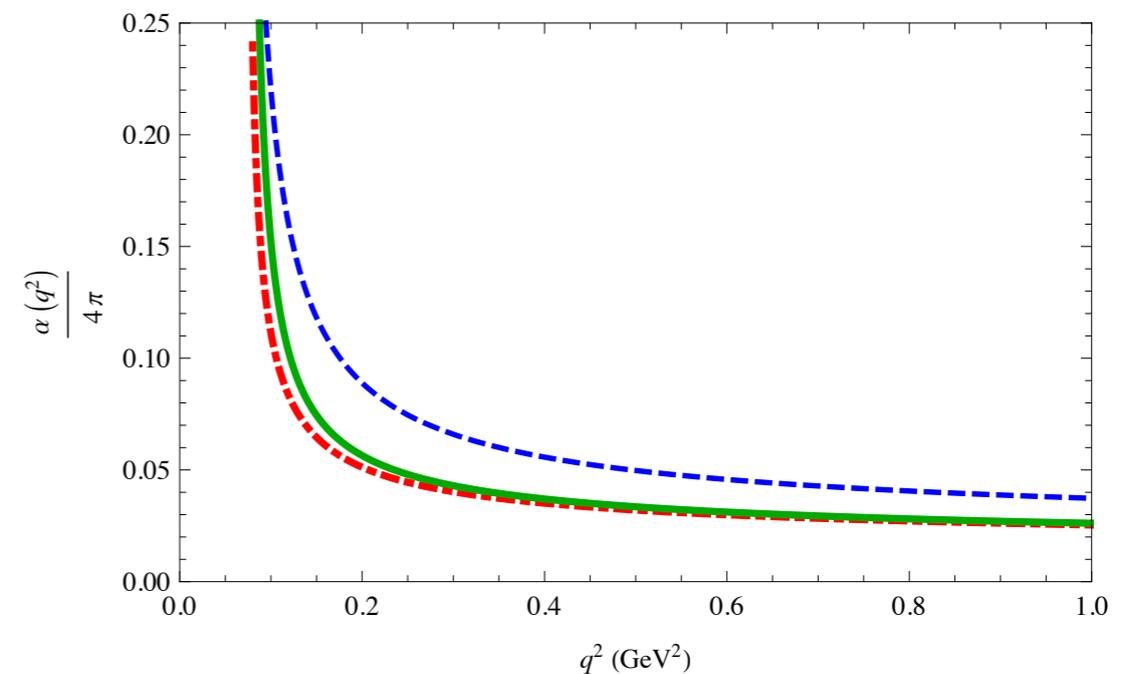
$$\frac{d a(Q^2)}{d(\ln Q^2)} = \beta_{N^m LO}(\alpha) = \sum_{k=0}^m a^{k+2} \beta_k$$

$\overline{\text{MS}}$ scheme
 $a = \alpha_s / 4\pi$

LO exact perturbative solution $\Lambda=250$ MeV

NLO exact perturbative solution $\Lambda=250$ MeV

NNLO exact perturbative solution $\Lambda=250$ MeV



QCD predicts the shape of the running coupling constant, not its value

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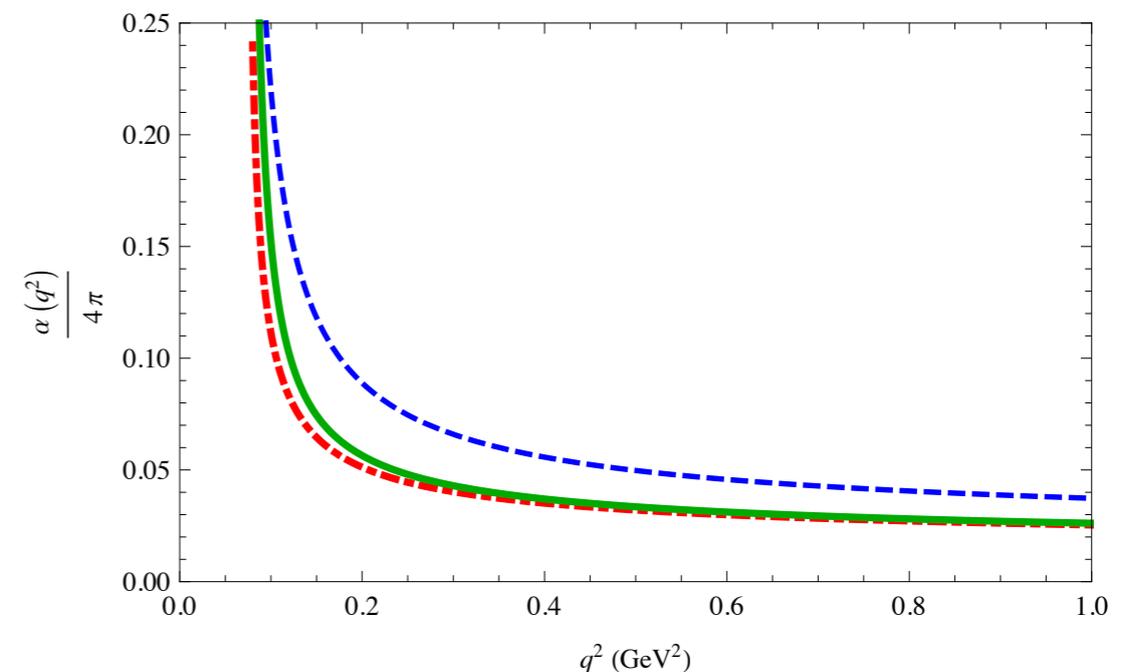
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Intermediate energy?

Perturbative to nonperturbative transition?

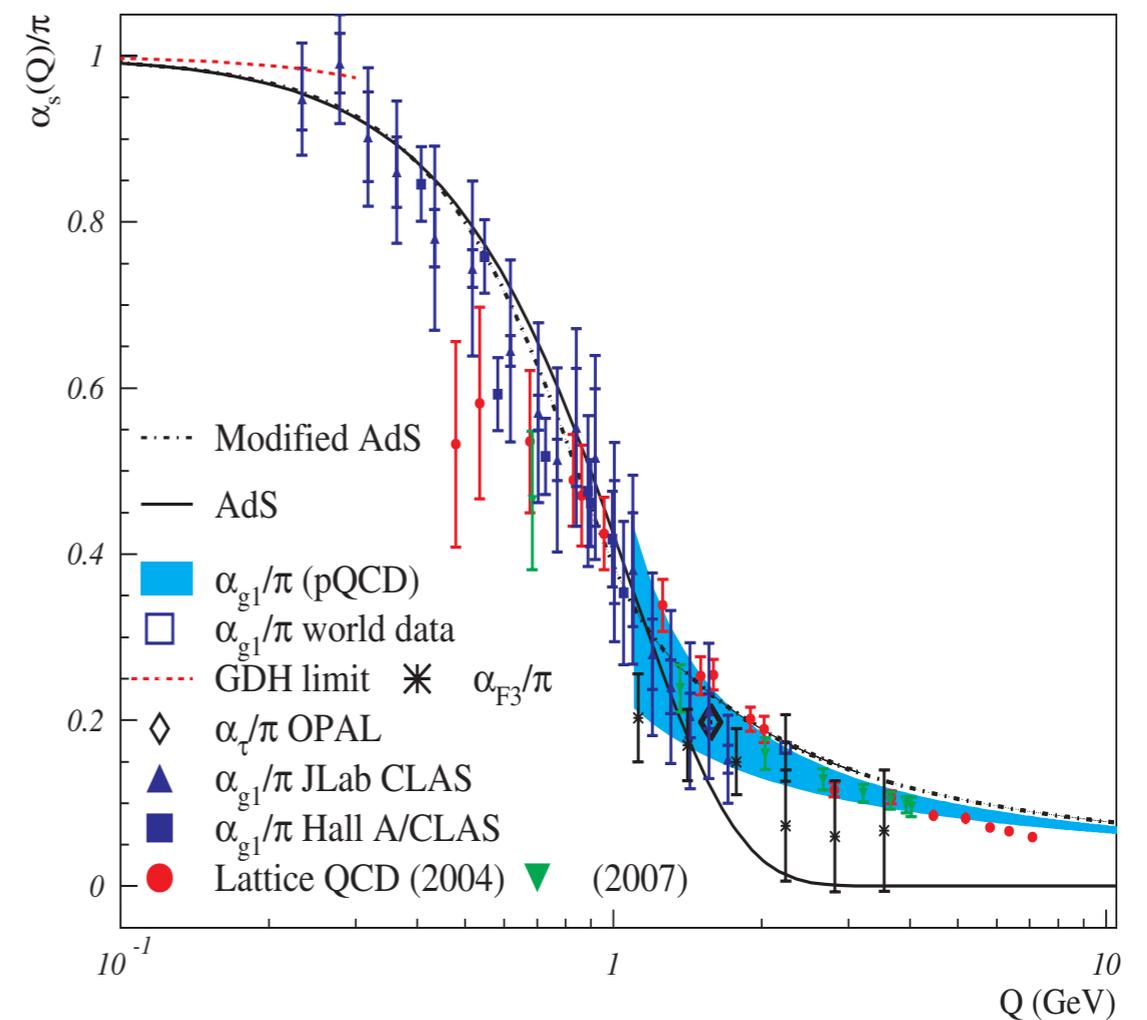
Effective Charges

The non-perturbative approach:

- Importance of finite couplings
- Taming the Landau pole

The non-perturbative extraction:

- Effective couplings from phenomenology
 - Dimensional transmutation (RG-improved)
- ➔ from RS dependence to Observable dependence (à la Grunberg)



[Brodsky et al., Phys.Rev.D81]
[Deur et al., Phys.Lett.B60]

Nonperturbative Gluon Propagator

Solving the Schwinger-Dyson eqs ...

$$\Delta^{-1}(Q^2) = Q^2 + m^2(Q^2)$$

J. M. Cornwall, Phys. Rev. D26, 1453 (1982)

A. C. Aguilar and J. Papavassiliou, JHEP0612, 012 (2006)

$$m^2(Q^2) = m_0^2 \left[\ln \left(\frac{Q^2 + \rho m_0^2}{\Lambda^2} \right) / \ln \left(\frac{\rho m_0^2}{\Lambda^2} \right) \right]^{-1-\gamma}$$

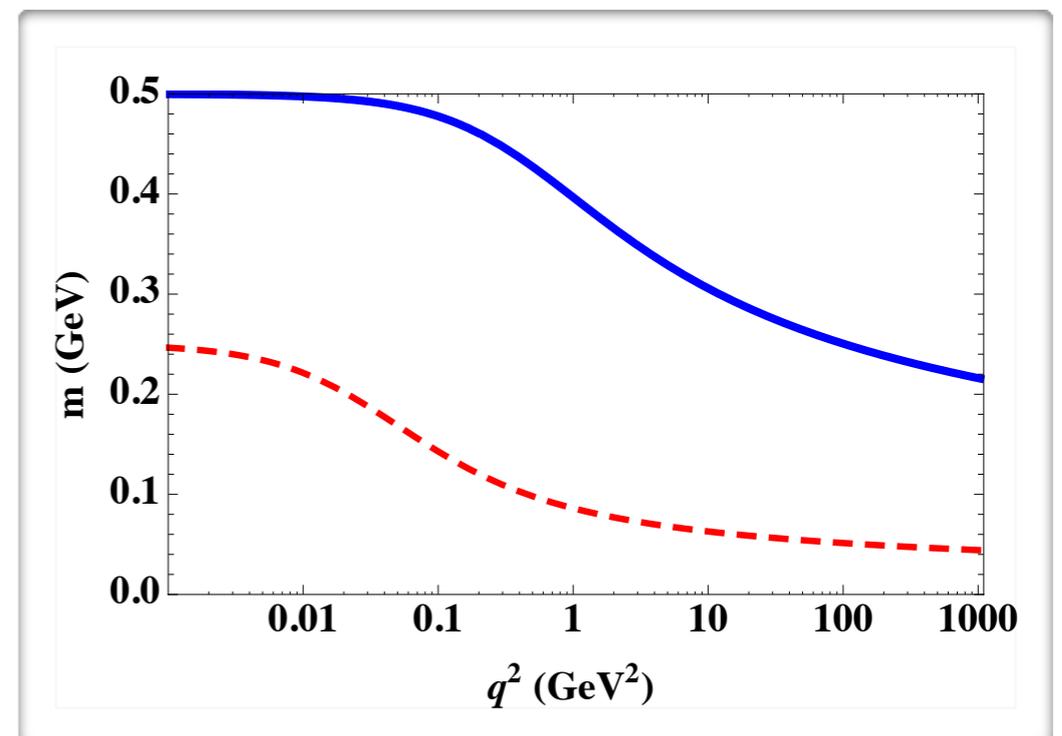
Gluon Mass as IR Regulator

- **effective gluon mass**
phenomenological estimates

$$m_0 \sim \Lambda - 2\Lambda$$

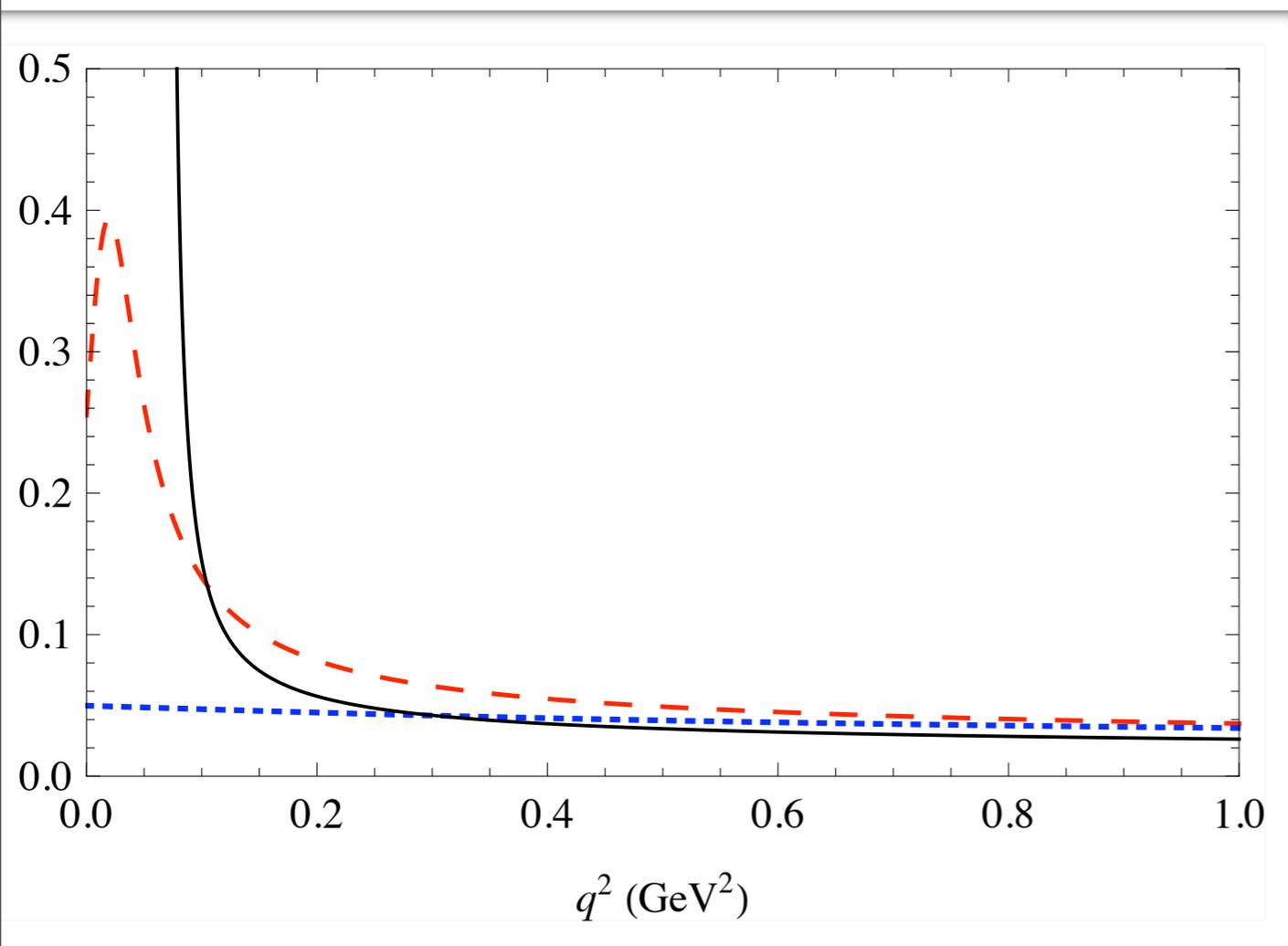
- **Solution free of Landau pole**
- **Freezes in the IR**

Low mass scenario
High mass scenario



NP Momentum-dependence of the Coupling Constant

$$\frac{\alpha_{\text{NP}}(Q^2)}{4\pi} = \left[\beta_0 \ln \left(\frac{Q^2 + \rho m^2(Q^2)}{\Lambda^2} \right) \right]^{-1}$$



L0 perturbative evolution
 $\Lambda=250$ MeV ; \overline{MS} scheme

Low mass scenario NP coupling constant
 $m_0=250$ MeV ; $\Lambda=250$ MeV ; $\rho=1.5$

High mass scenario NP coupling constant
 $m_0=500$ MeV ; $\Lambda=250$ MeV ; $\rho=2$.

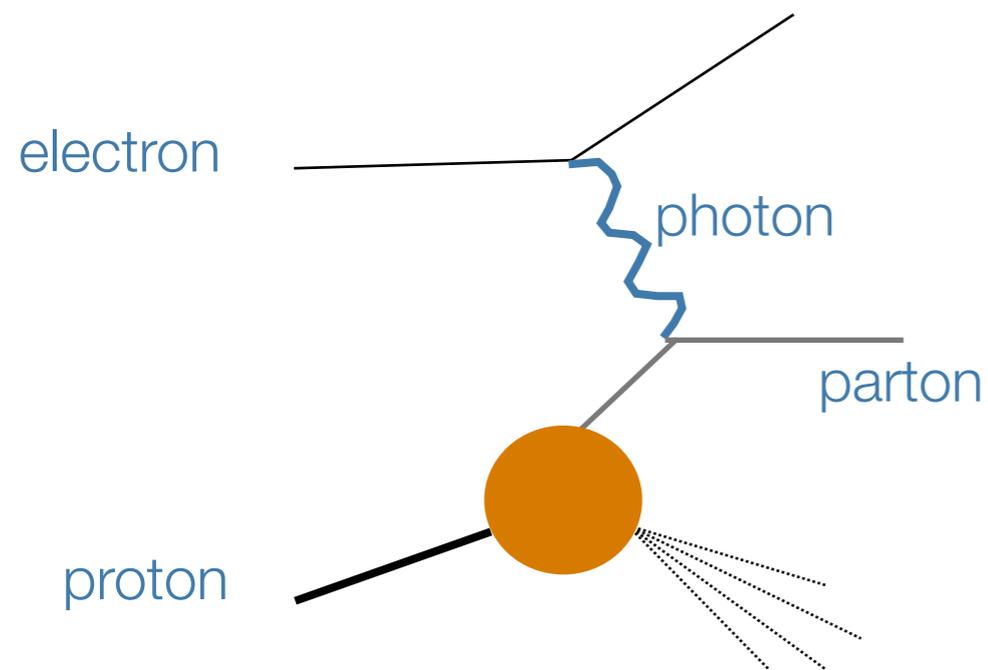
Hadron Structure Phenomenology

Final State Interaction and Parton Distribution Functions

Hard Probes and Factorization

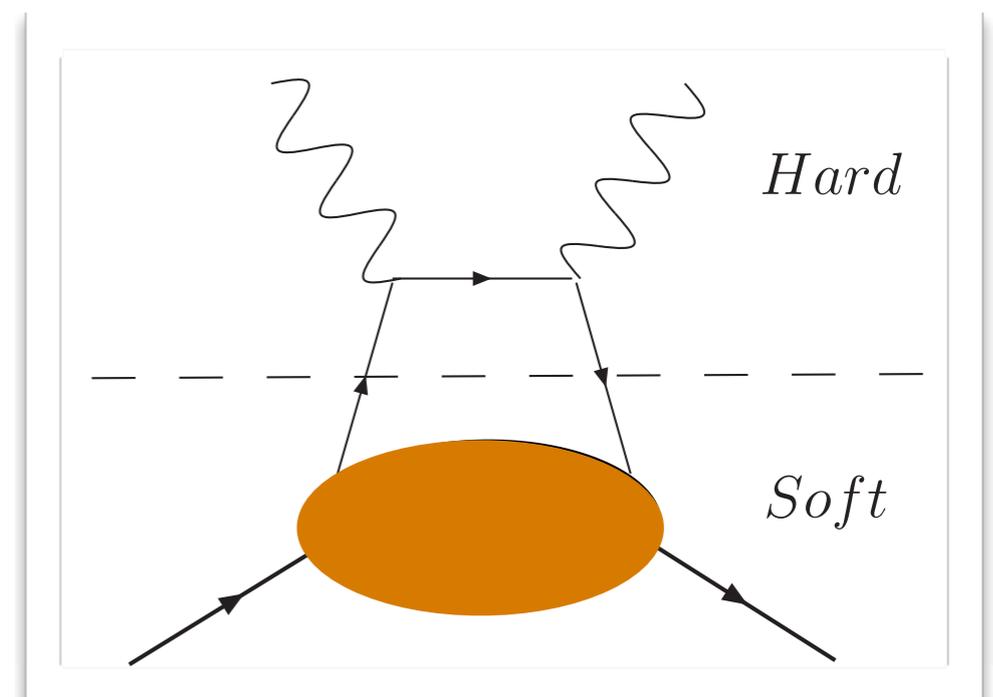
Small size configuration \Rightarrow Hard Probes \Rightarrow Hard processes

Deep Inelastic Scattering



Hadronic tensor \Rightarrow

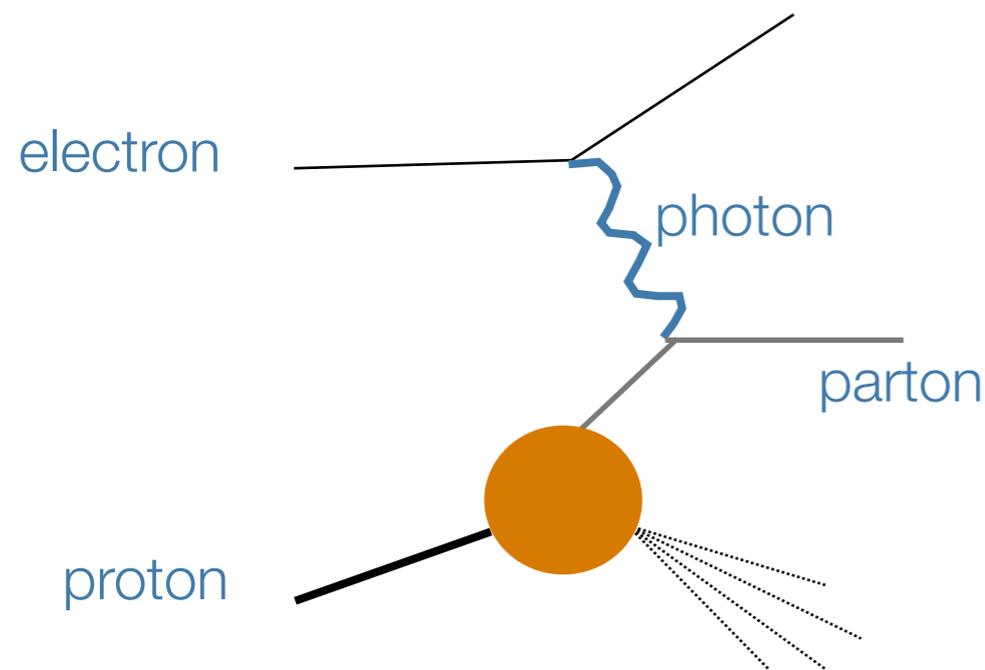
Parton Model
High energy photon Q^2
Fast-moving proton
Bjorken scaling



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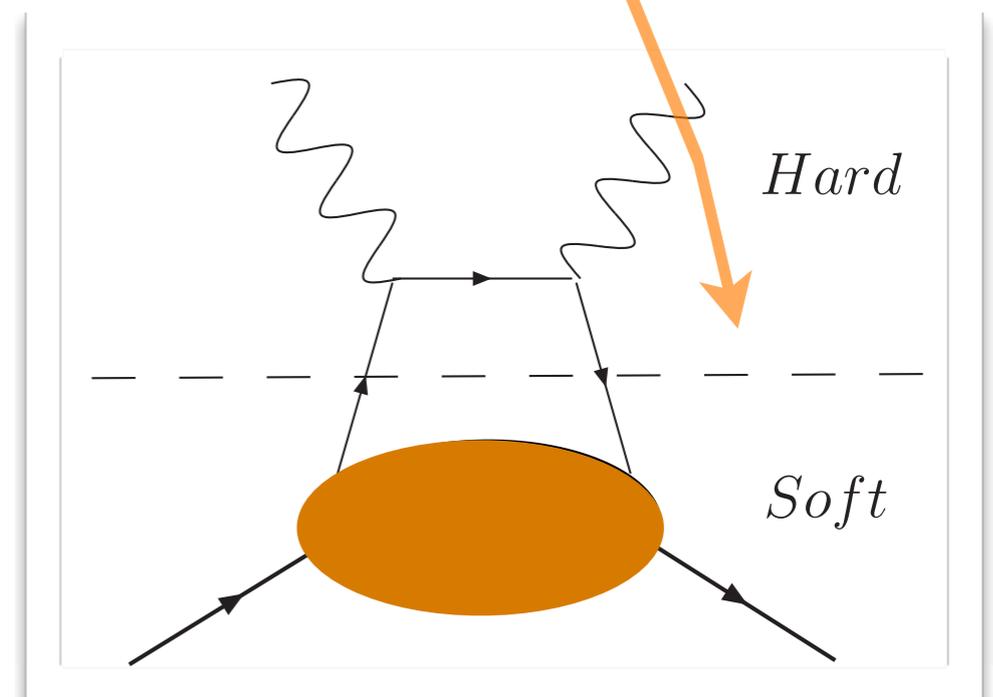
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Factorization
& factorization scale



Transverse Momentum Dependent PDFs

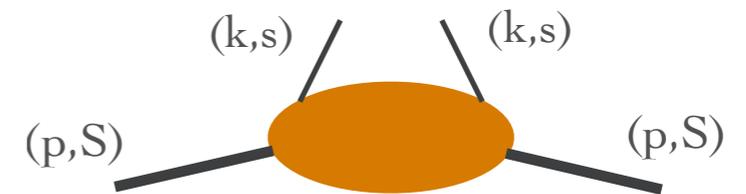
Hadronic matrix elements to $f(x, k_T)$



Number of independent structure functions



Number of Lorentz scalars +hermiticity+parity invariance+Time-reversal invariance



- Relaxing Time-reversal invariance \Rightarrow naive T-odd functions

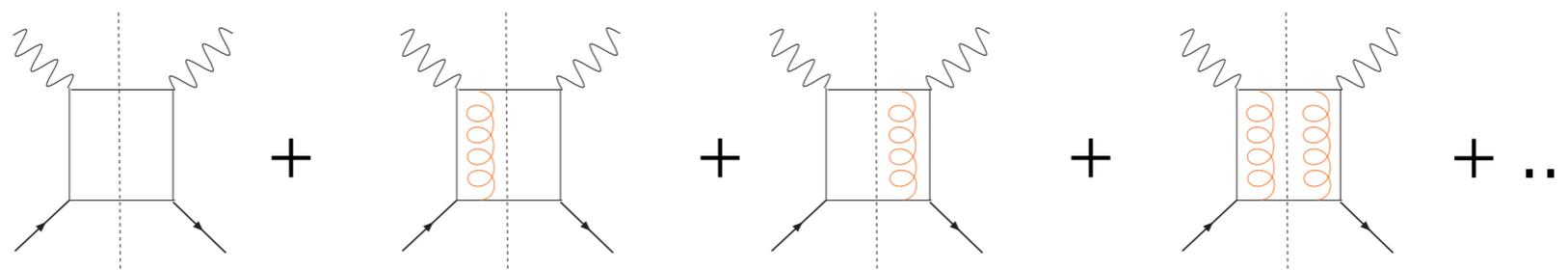
Sivers & Boer-Mulders functions

Sivers, Phys.Rev.D41
Boer & Mulders, Phys.Rev.D57

- Existence of Final State Interactions (FSI) at leading-order

Brodsky, Hwang & Schmidt, Phys.Lett.B530

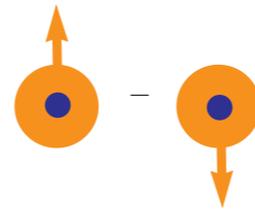
- Importance of the gauge link



T-odd TMDs

The Sivers function $f_{1T}^{\perp Q}(x, k_T)$

⇒ Distribution of **unpolarized quarks** inside a **transversely polarized proton**



The Boer-Mulders functions $h_1^{\perp Q}(x, k_T)$

⇒ Distribution of **transversely polarized quarks** inside a **unpolarized proton**



- Matrix element of low twist operator

$$f_{1T}^{\perp q}(x, k_T) = -\frac{M}{2k_x} \int \frac{d\xi^- d^2\vec{\xi}_T}{(2\pi)^3} e^{-i(xp^+ \xi^- - \vec{k}_T \cdot \vec{\xi}_T)}$$

$$\times \frac{1}{2} \sum_{S_y = -1, 1} S_y \langle PS_y | \bar{\psi}_q(\xi^-, \vec{\xi}_T) \mathcal{L}_{\vec{\xi}_T}^\dagger(\infty, \xi^-) \gamma^+ \mathcal{L}_0(\infty, 0) \psi_q(0, 0) | PS_y \rangle + \text{h.c.}$$

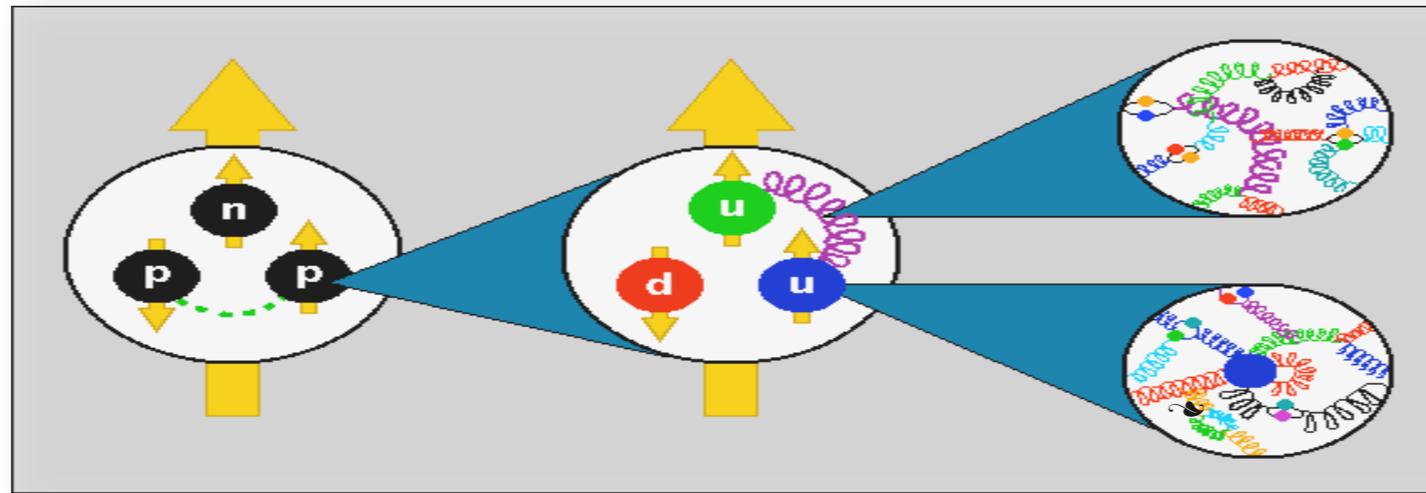
- Importance of gauge link

$$\mathcal{L}_{\vec{\xi}_T}(\infty, \xi^-) = \mathcal{P} \exp \left(-ig \int_{\xi^-}^{\infty} A^+(\eta^-, \vec{\xi}_T) d\eta^- \right)$$

- holds in covariant gauges
- process dependent

Hadronic Models

Hadron \Leftrightarrow Constituent quarks \Leftrightarrow Current quarks



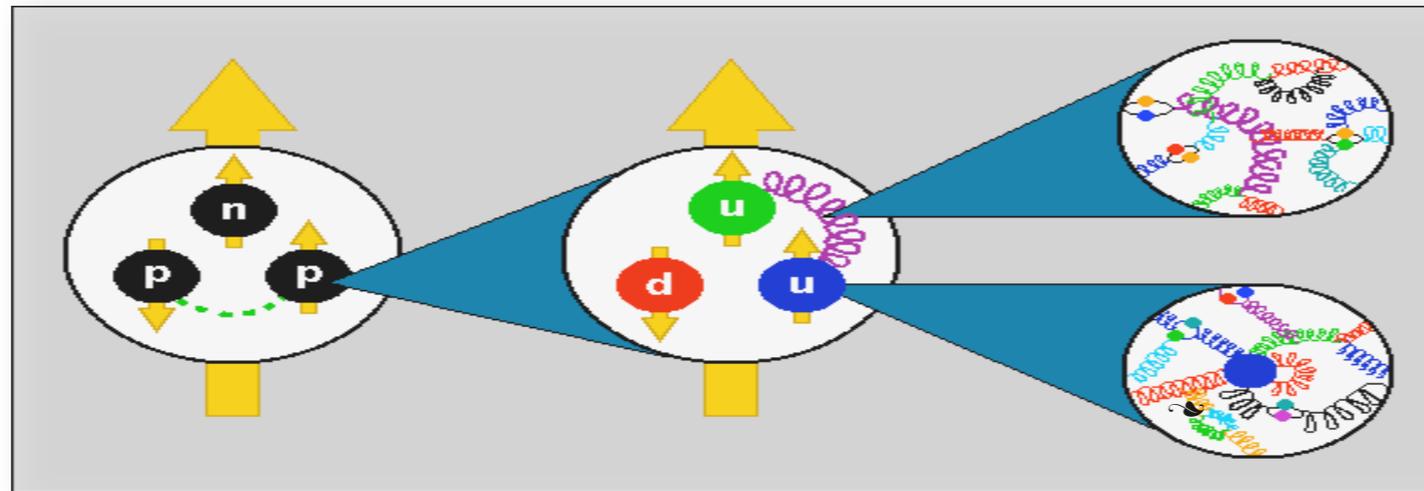
Nonperturbative vs. Perturbative QCD

Models of Hadron Structure

Renormalization Group Eqs.

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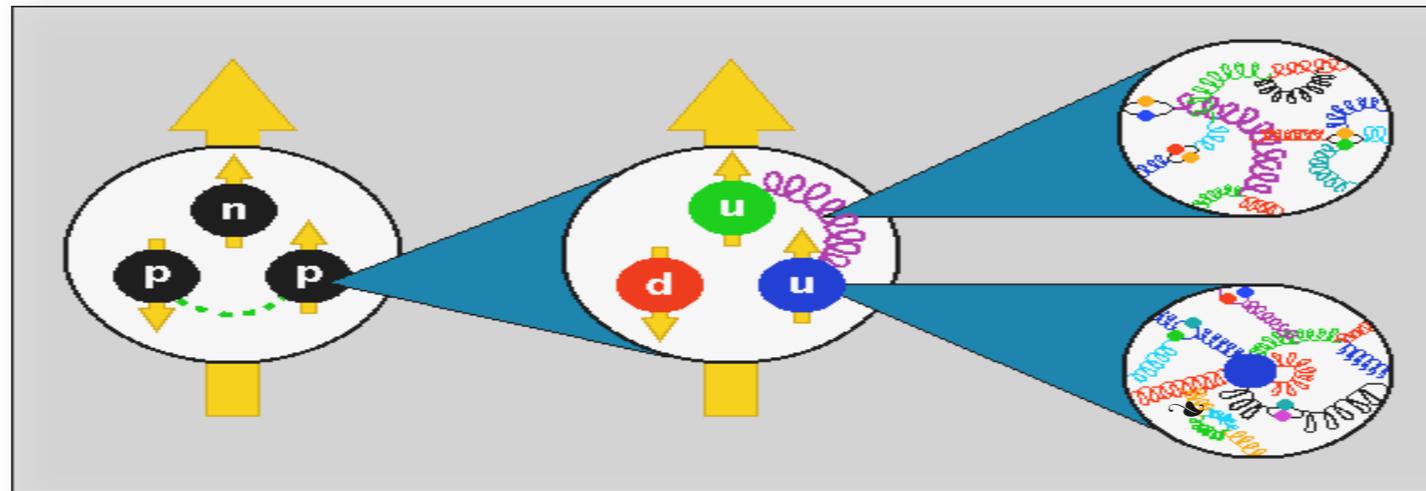
Renormalization Group Eqs.

Observable

- calculated in hadronic model
- at scale μ_0
- switch on QCD evolution

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Hadronic Scale

from collinear PDFs, e.g. CTEQ, GRV,...

**We use RGE and one *first principle* based assumption.
Then we set scenarios ...**

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Say there exists a scale at which there is no sea and no gluon, then

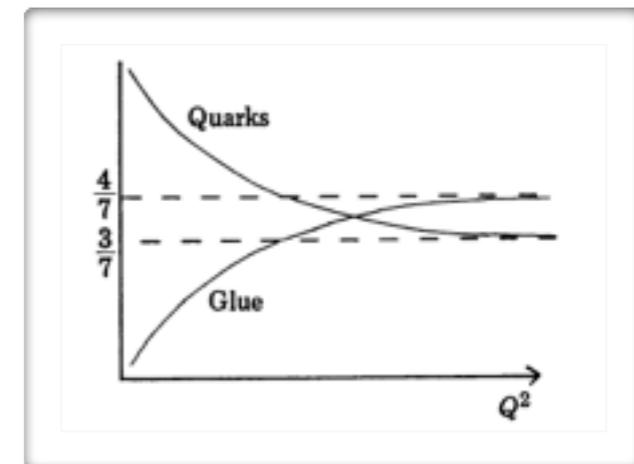
$$\langle (u_v + d_v) (\mu_0^2) \rangle_{n=2} = 1$$

QCD evolution introduces gluons and sea quarks:

$$\langle (u_v + d_v) (Q^2 = 10 \text{ GeV}^2) \rangle_{n=2} = 0.36$$



DATA= PDFs parameterization



R.G.Roberts
"The Structure of the Proton"

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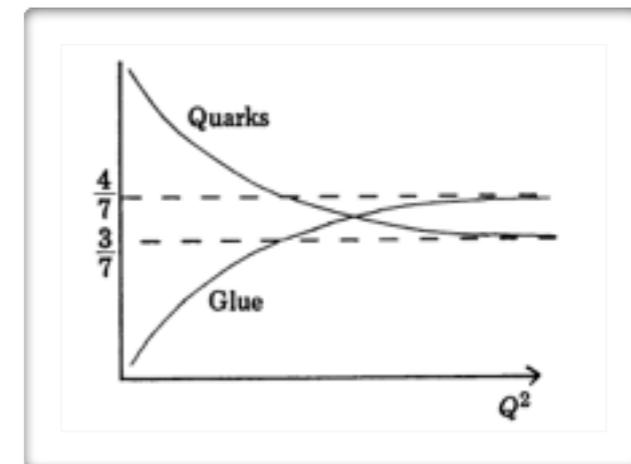
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Evolve in energy until 2nd moment=1
Find $\mu_0^2 \sim 0.1 \text{ GeV}^2 + \Delta\mu_0^2$

Perturbative vs. NP 'evolution': Fixing the hadronic scale

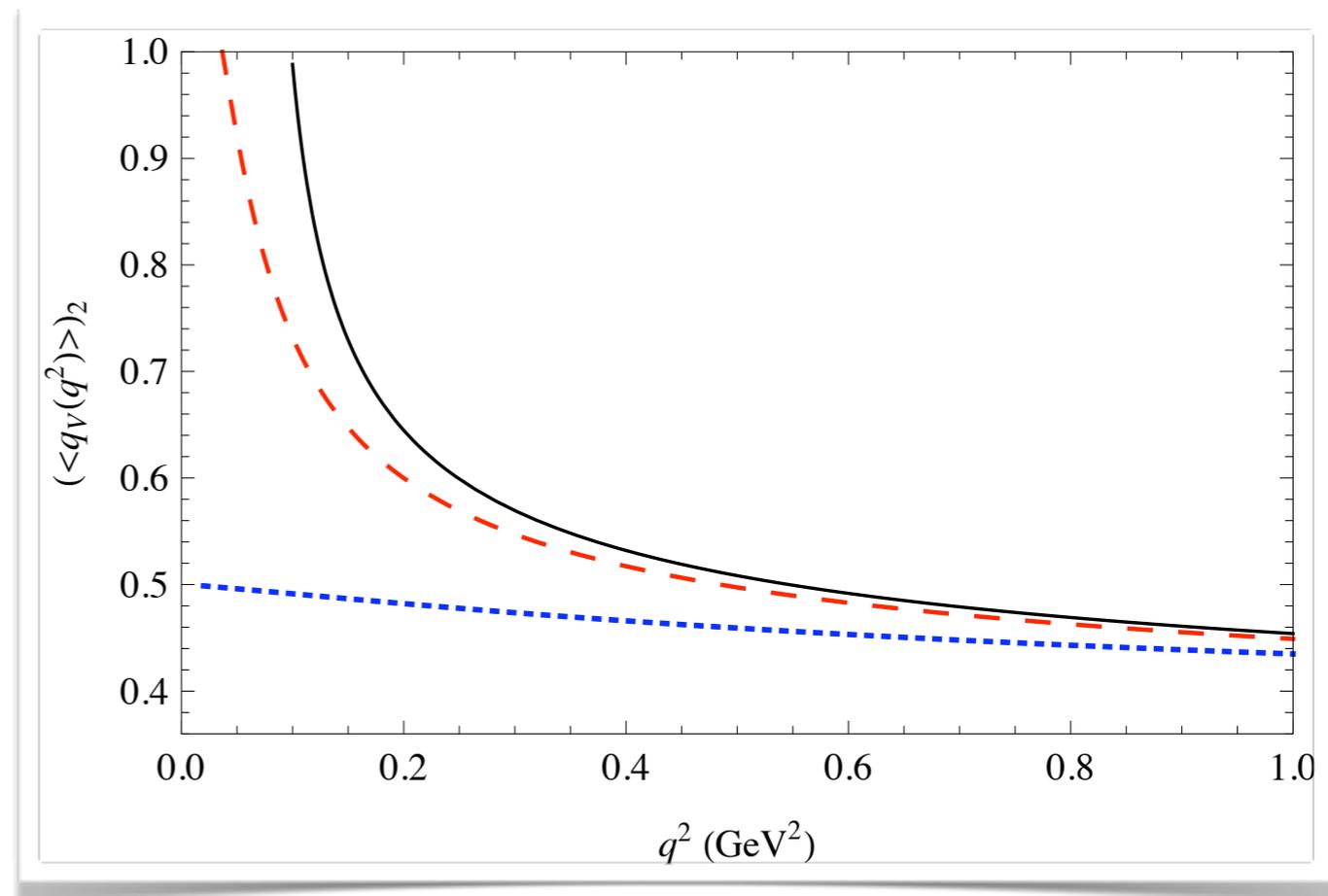
2nd moment of f_1

$$\langle q_v(Q^2) \rangle_n = \langle q_v(\mu_0^2) \rangle_n \left(\frac{\alpha(Q^2)}{\alpha(\mu_0^2)} \right)^{d_{NS}^n}$$

L0 perturbative evolution
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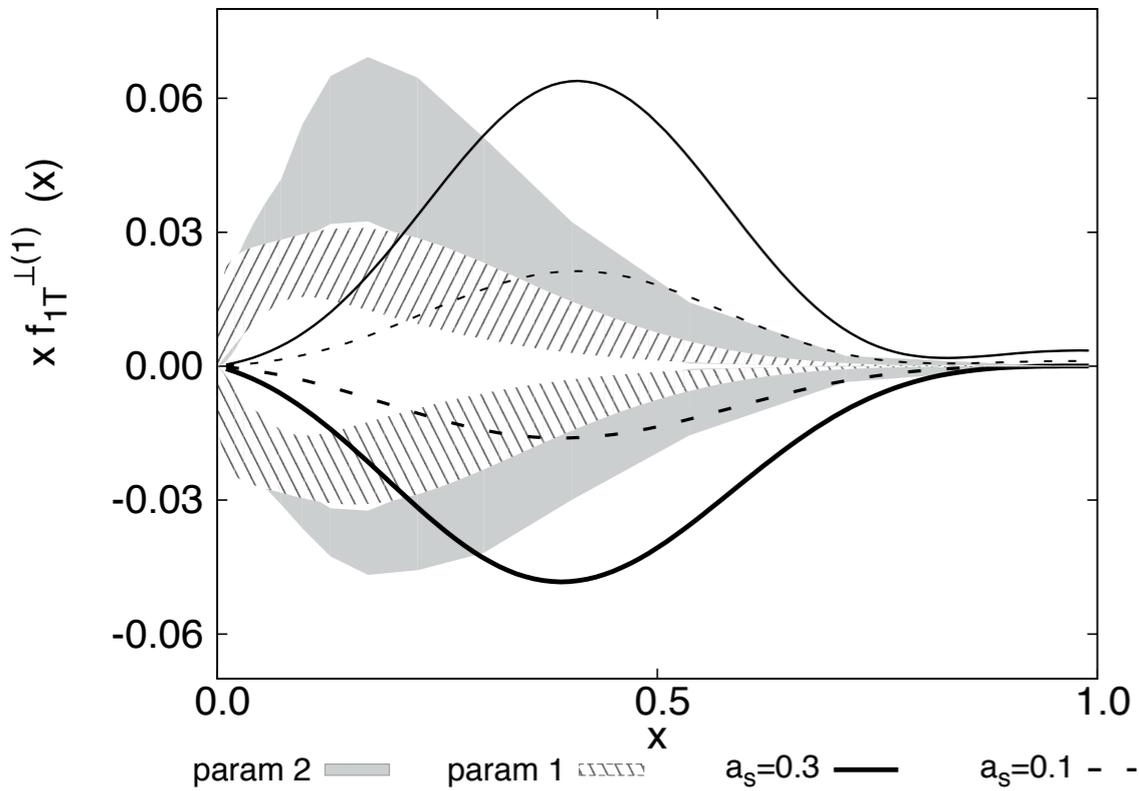
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Sivers & Boer-Mulders functions

[A.C., Scopetta & Vento, Eur.Phys.J. A47]



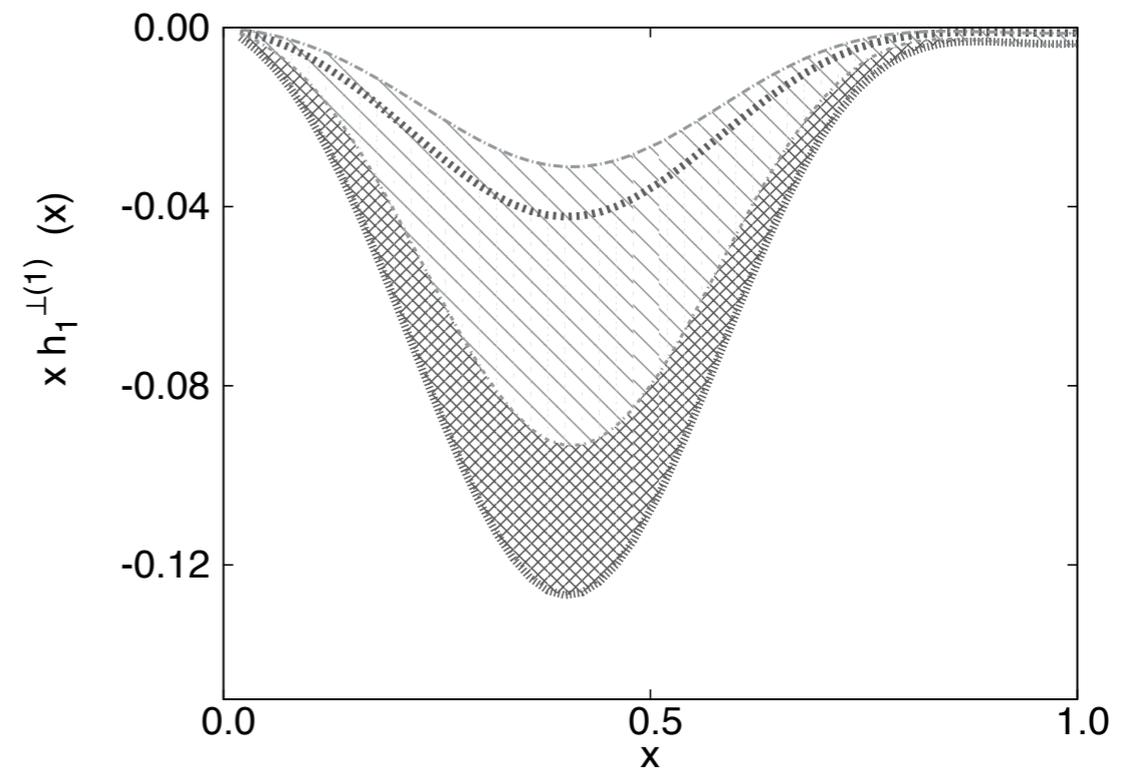
$$0.1 < \frac{\alpha_s(\mu_0^2)}{4\pi} < 0.3$$

dashed $\frac{\alpha_s(\mu_0^2)}{4\pi} = 0.1$

solid $\frac{\alpha_s(\mu_0^2)}{4\pi} = 0.3$

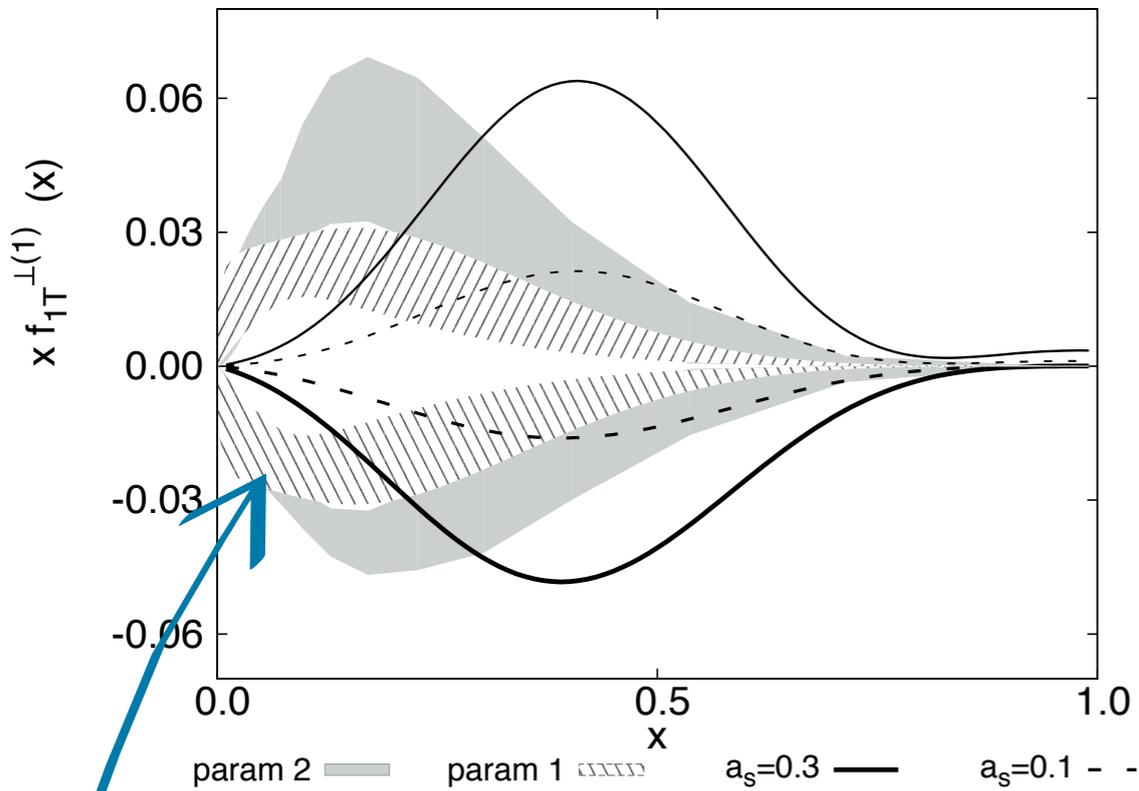
Bag Model:

rescaling/defining error
of f_{1T^\perp} & h_{1^\perp}



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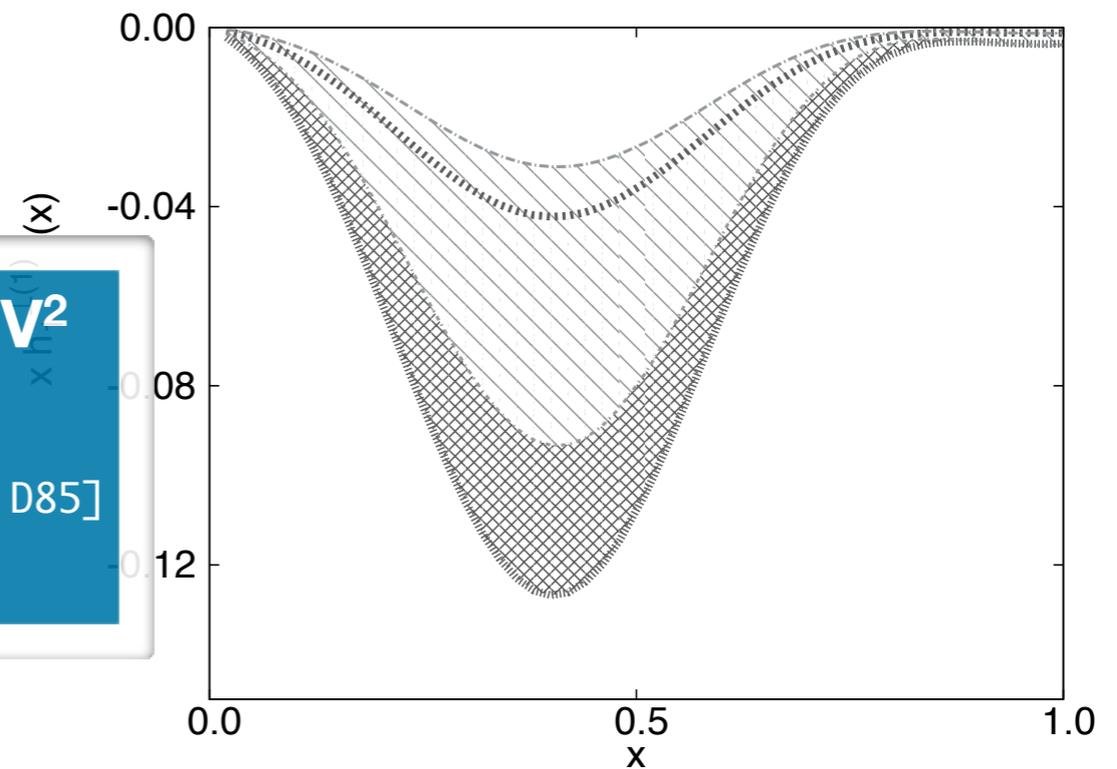
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Phenomenological extractions at $Q^2=2.5\text{GeV}^2$

→ Need for QCD formalism for T-odd TMDs

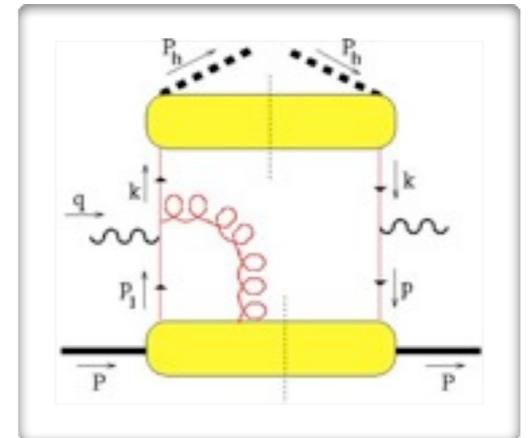
[Aybat, Collins, Qiu & Rogers, Phys.Rev. D85]

→ Additional source of error

Work in progress for T-odd TMDs

- Ambiguity Sivers function and Qiu-Sterman function

- Model dependent definition of the FSI and of the proton



- TMD evolution: Coupled CSS and RGE \rightarrow two scales ! [Aybat et al., PRD85]

- Definition of momentum regions

[Bacchetta et al., JHEP 0808]

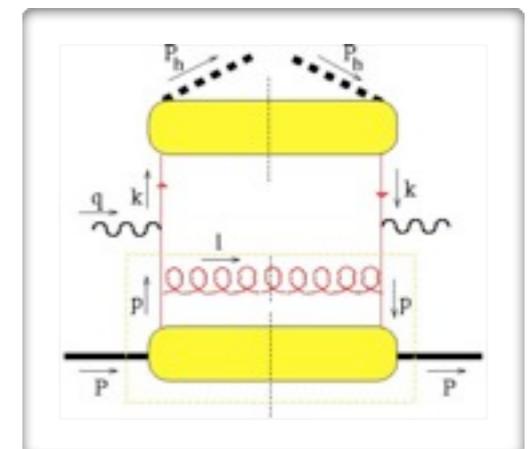
- Redefinition of both scale for model calculations (in collaboration with T. Rogers)

- Correspondance effective coupling from the soft blob with pQCD

- [Brodsky et al., Phys.Rev.D81] *À la Grunberg?* [Phys. Rev. D29]

- Commensurate Scale Relations

[Brodsky & Lu, Phys. Rev. D251]



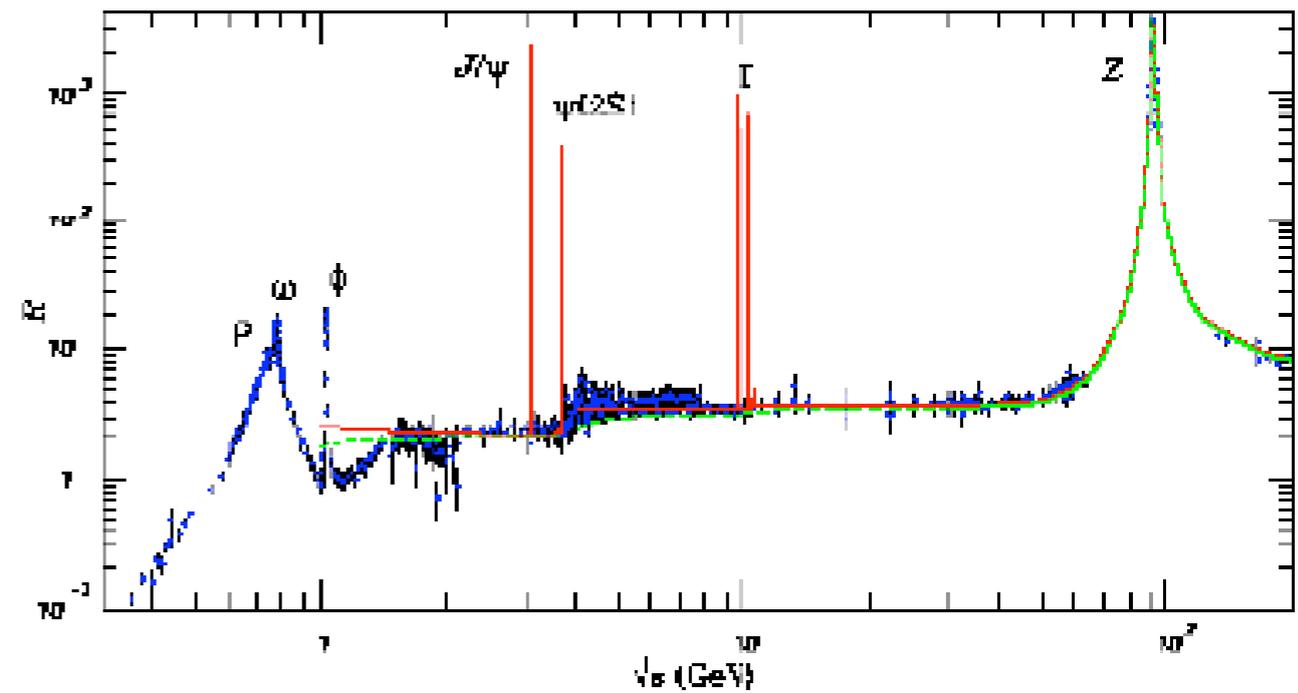
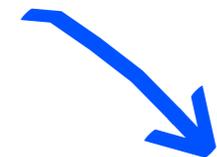
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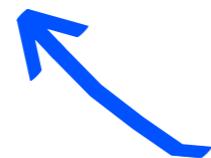
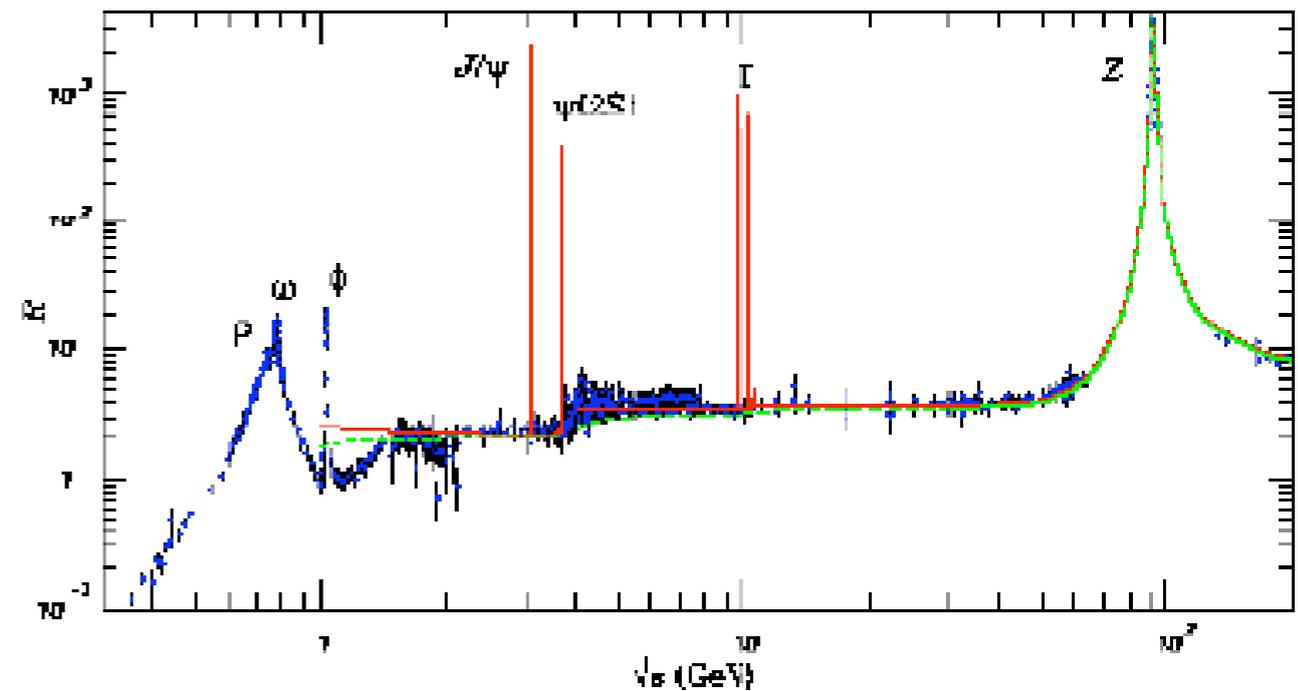
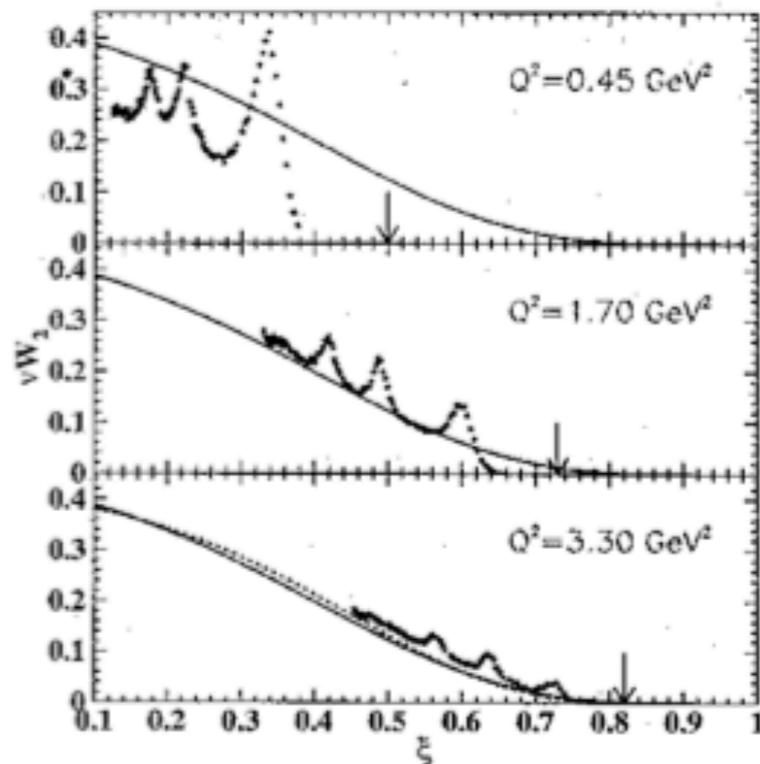
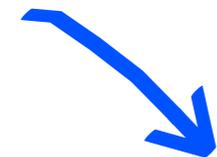


Complementarity between Parton and Hadron descriptions of observable

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Structure functions

Resonance region \Leftrightarrow Scaling region

$x_{Bj} > 0.5$, Q^2 multi-GeV region $\Rightarrow W^2 \leq 5 \text{ GeV}^2$

[Bloom & Gilman, Phys.Rev.Lett.25]

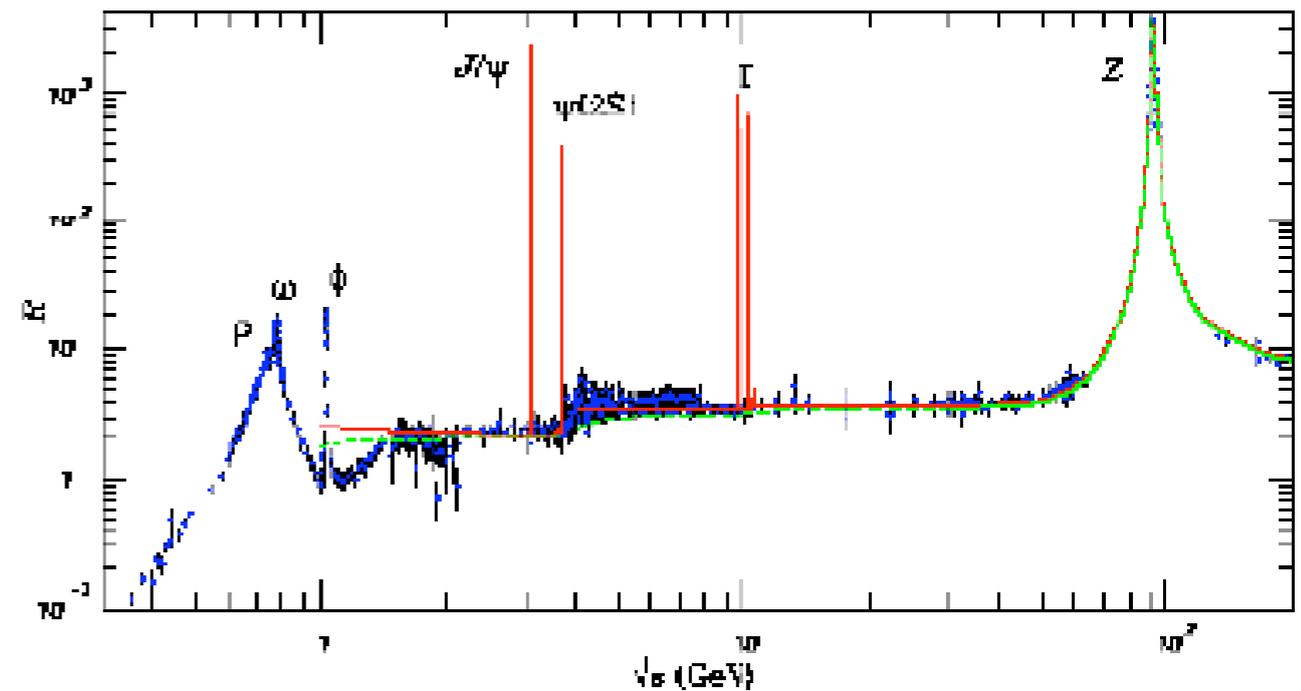
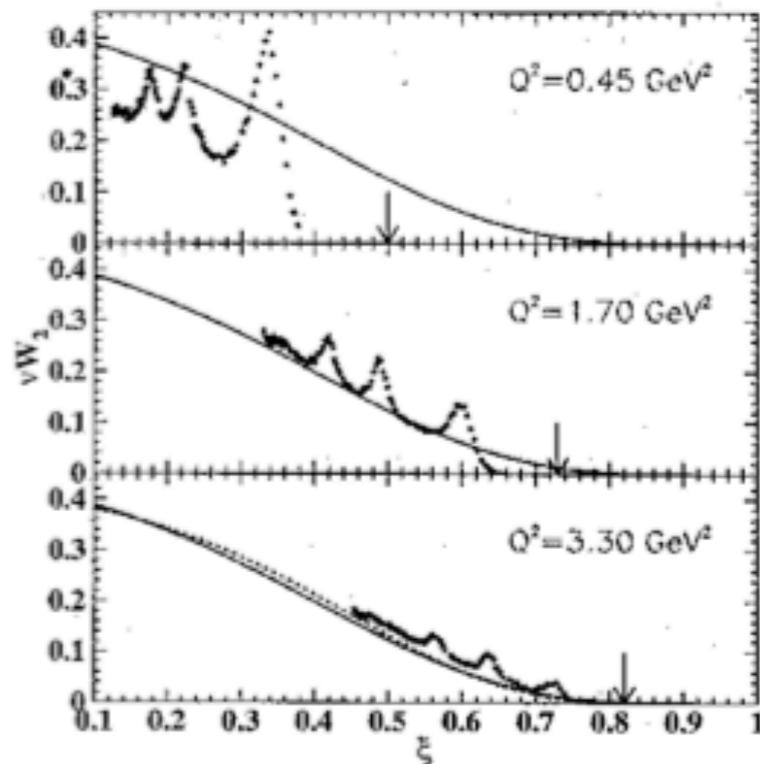
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Talk by W. Melnitchouk



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Two Complementary Approaches to Structure Functions

experiment

$$I^{res}(Q^2) = \int_{x_m}^{x_M} F_2^{Res}(x, Q^2) dx$$
$$I^{DIS}(Q^2) = \int_{x_m}^{x_M} F_2^{DIS}(x, Q^2) dx$$

perturbative QCD

$$x_M \div x_m \Leftrightarrow W_m^2 \div W_M^2 \Rightarrow 1 \div 4 \text{ GeV}^2$$

- **Nonperturbative models analysis**
- **Perturbative analysis**

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Start with NLO PDF and then ...

- Target Mass Corrections (TMC)
- Large-x Resummation (LxR)
- Higher-order in pQCD
- Higher-Twists

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[A. De Rujula, et al., Phys. Lett. B64]

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- Higher-order in pQCD \longrightarrow
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Ok

pQCD

?

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Large-x Resummation

Text Book, e.g. *Cornerstones of QCD*, M. Pennington.

- Large invariants: $\Lambda^2 \ll W^2 \ll Q^2$
- Argument for α_s is ω^2 , mass square of final state of γ^* parton collision

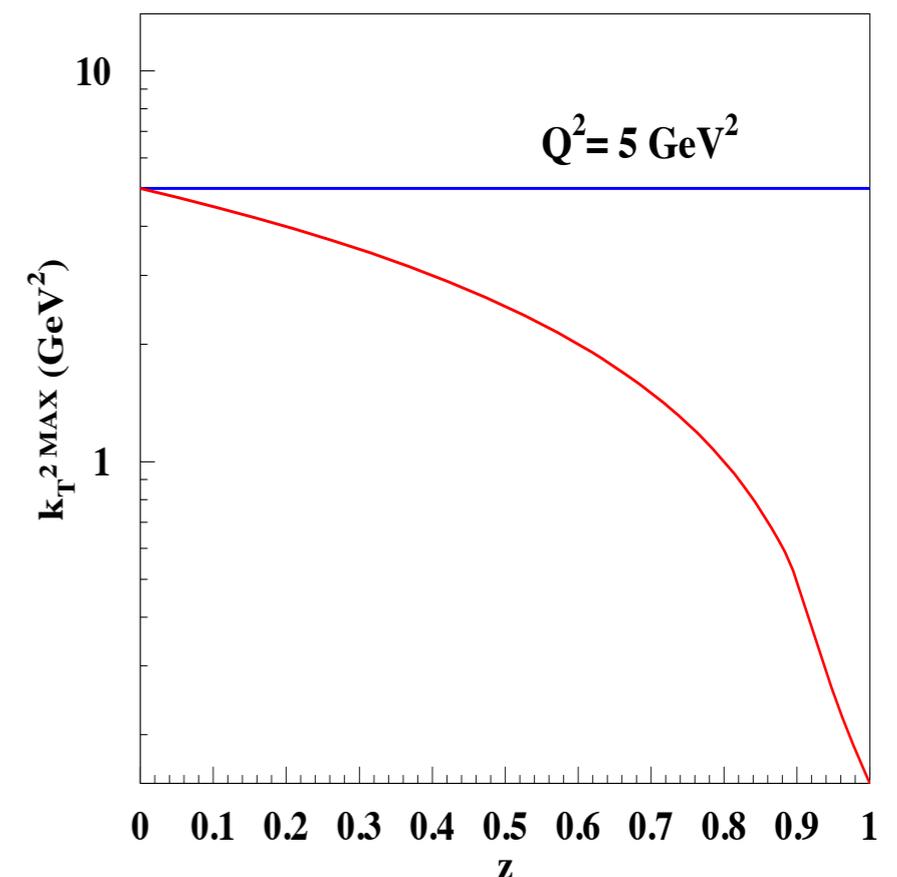
$$\omega^2 = \frac{Q^2}{z} (1-z)$$

Without LxR, upper limit = Q^2

$$q(x, Q^2) = \int_x^1 \frac{dz}{z} \int_{\mu^2}^{Q^2 \frac{1-z}{4z}} dk_T^2 \alpha_S(k_T^2) P_{qq}(z) q\left(\frac{x}{z}, k_T^2\right)$$

The structure functions become

$$F_2^{NS}(x, Q^2) = \sum_q \int_x^1 dz \frac{\alpha_s\left(\frac{Q^2(1-z)}{4z}\right)}{2\pi} C_{NS}(z) q_{NS}\left(\frac{x}{z}, Q^2\right)$$



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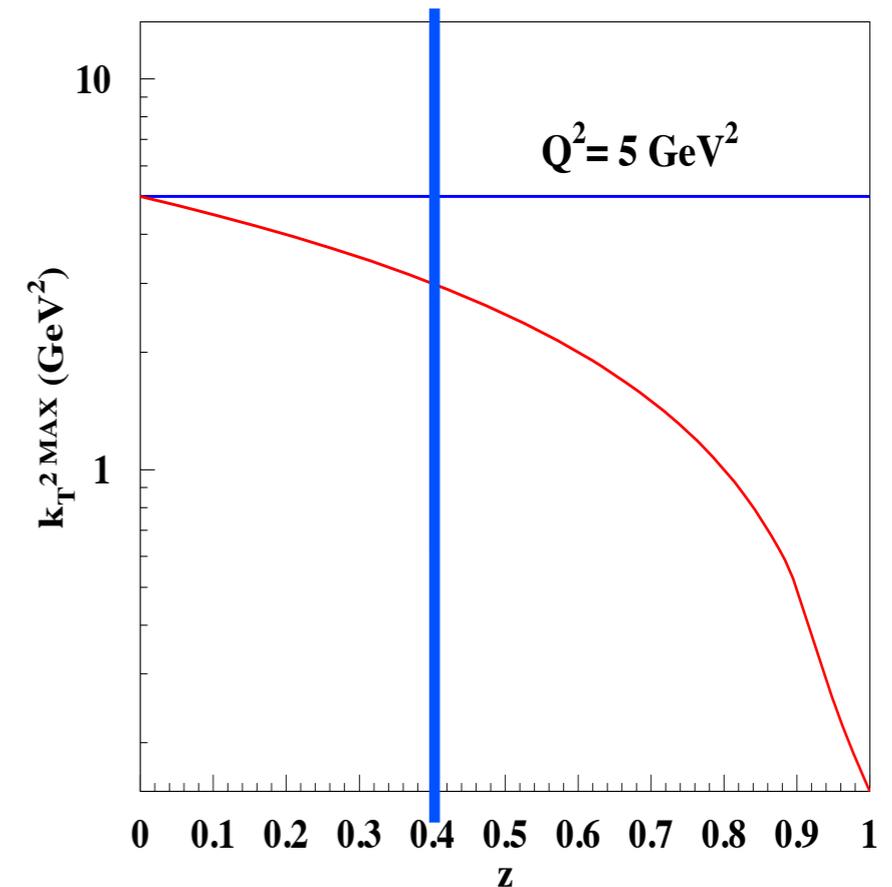
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$$q(x, Q^2) = \int_x^1 \frac{dz}{z} \int_{\mu^2}^{Q^2 \frac{1-z}{4z}} dk_T^2 \alpha_S(k_T^2) P_{qq}(z) q\left(\frac{x}{z}, k_T^2\right)$$

The structure functions become

$$F_2^{NS}(x, Q^2) = \sum_q \int_x^1 dz \frac{\alpha_s\left(\frac{Q^2(1-z)}{4z}\right)}{2\pi} C_{NS}(z) q_{NS}\left(\frac{x}{z}, Q^2\right)$$

x-values



Large-x Resummation

Text Book, e.g. *Cornerstones of QCD*, M. Pennington.

- Large invariants: $\Lambda^2 \ll W^2 \ll Q^2$
- Argument for α_s is ω^2 , mass square of final state of γ^* parton collision

$$\omega^2 = \frac{Q^2}{z} (1-z)$$

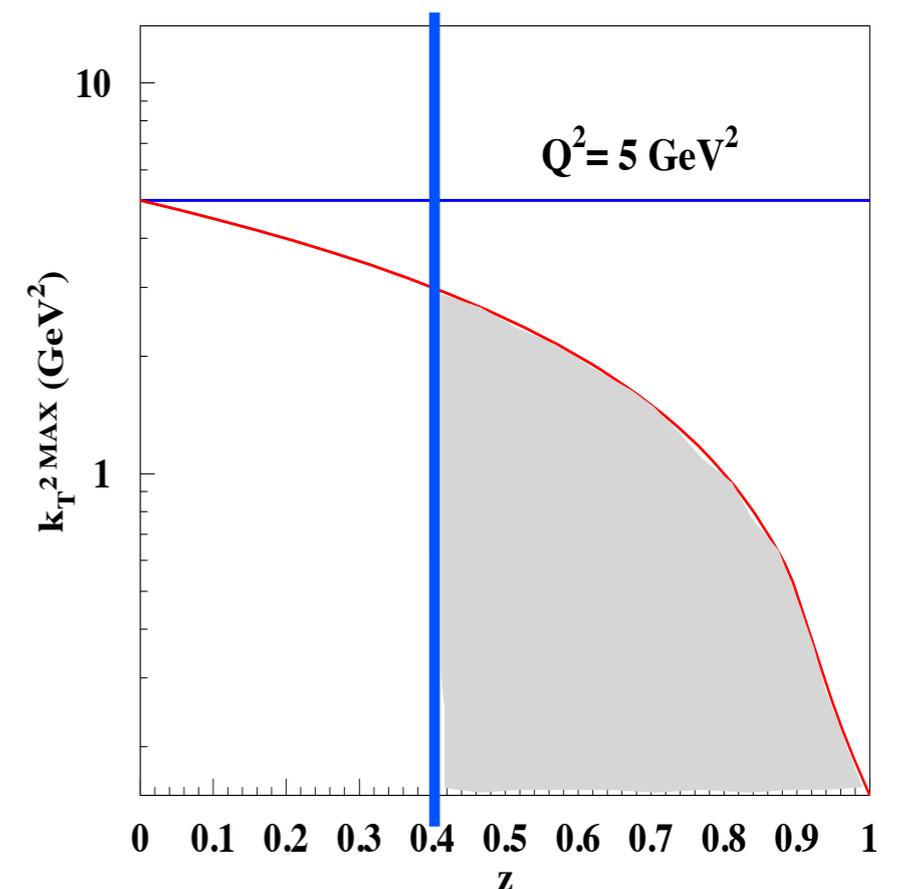
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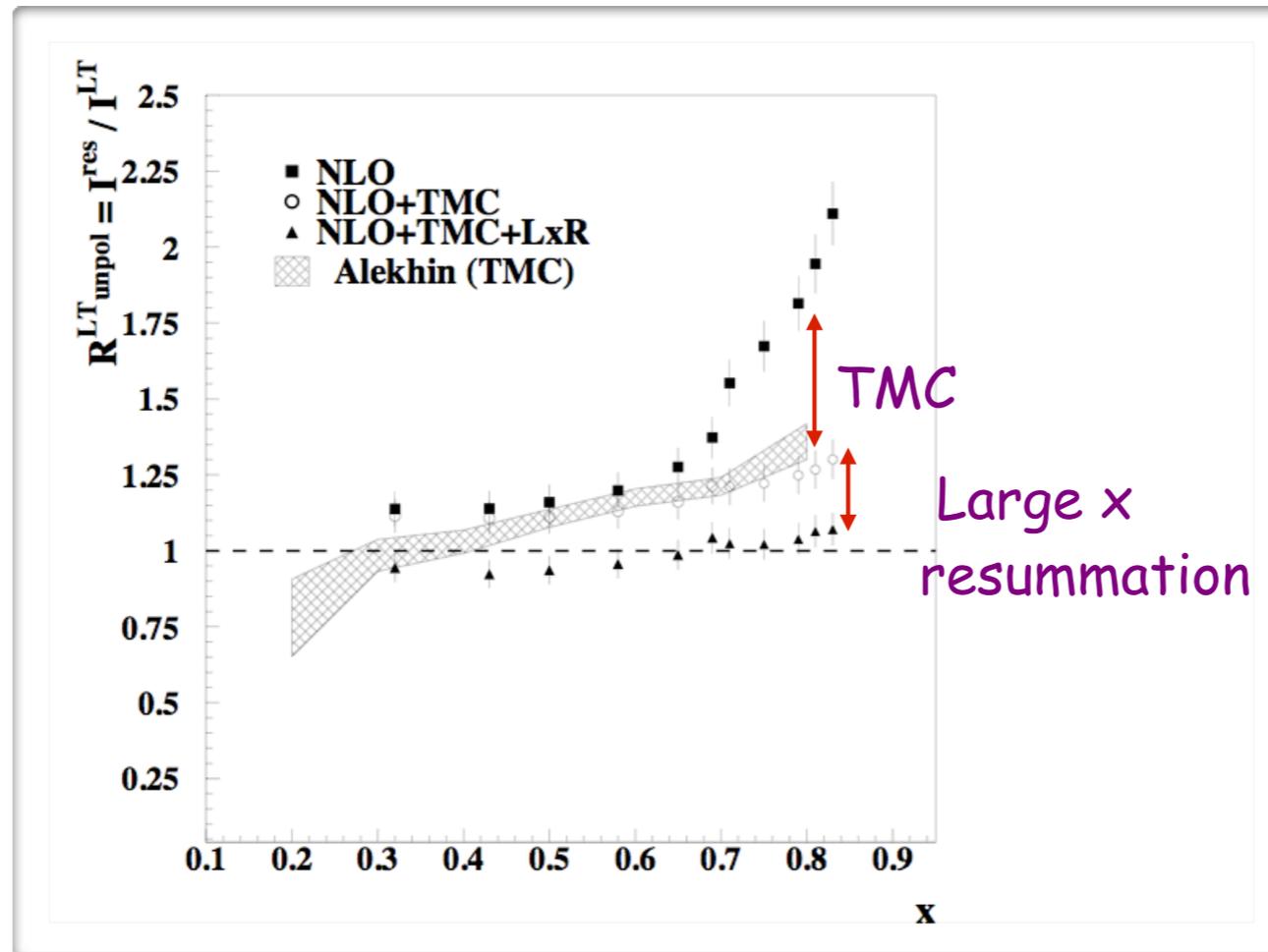
x-values



Size of Nonperturbative Contributions

[Niculescu et al., PRD60]

[Bianchi, Fantoni & Liuti, PRD69]

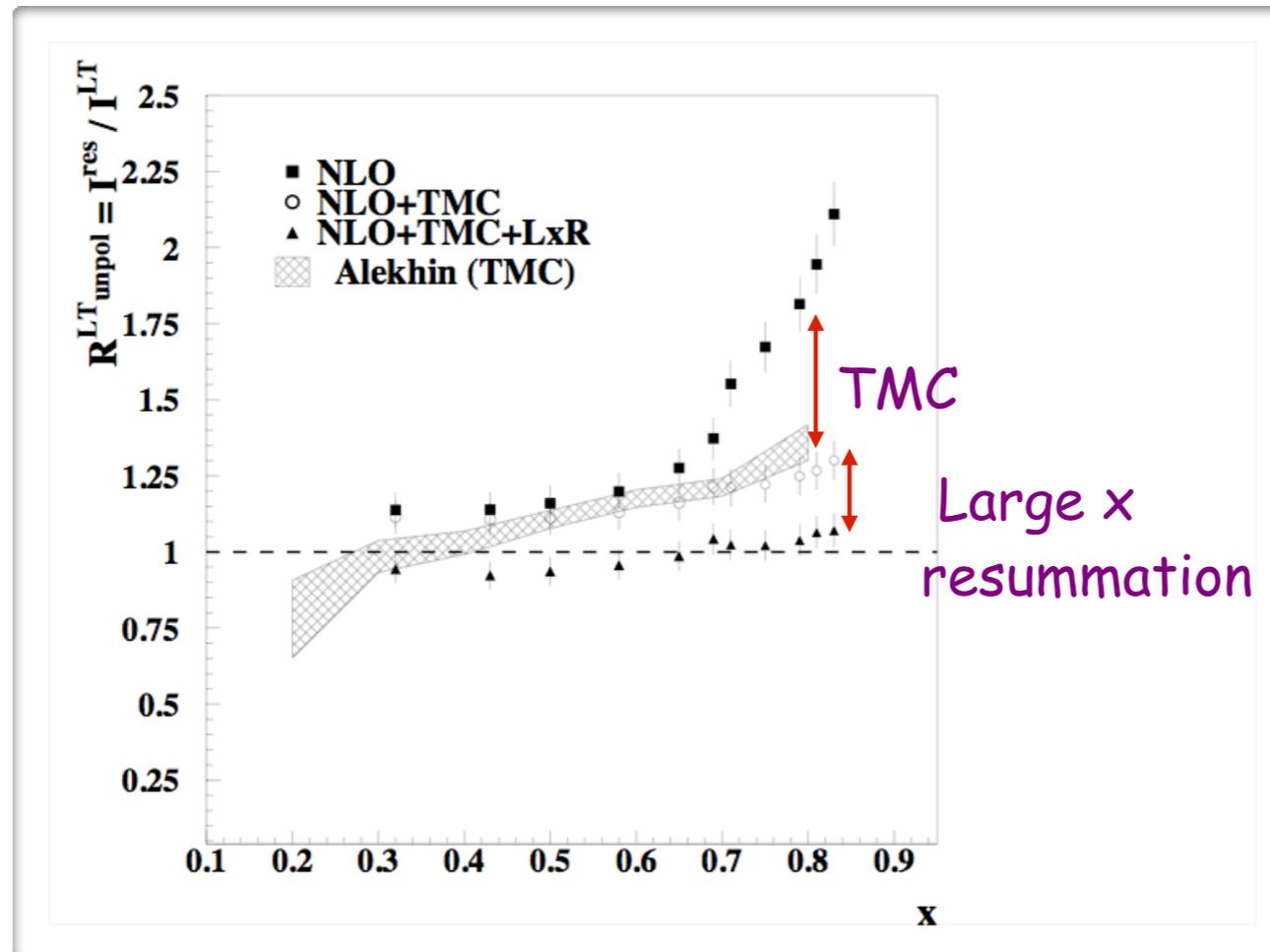


$$R \equiv \frac{I^{res}(Q^2)}{I^{DIS}(Q^2)} \Leftrightarrow \text{Duality fulfilled if } R=1$$

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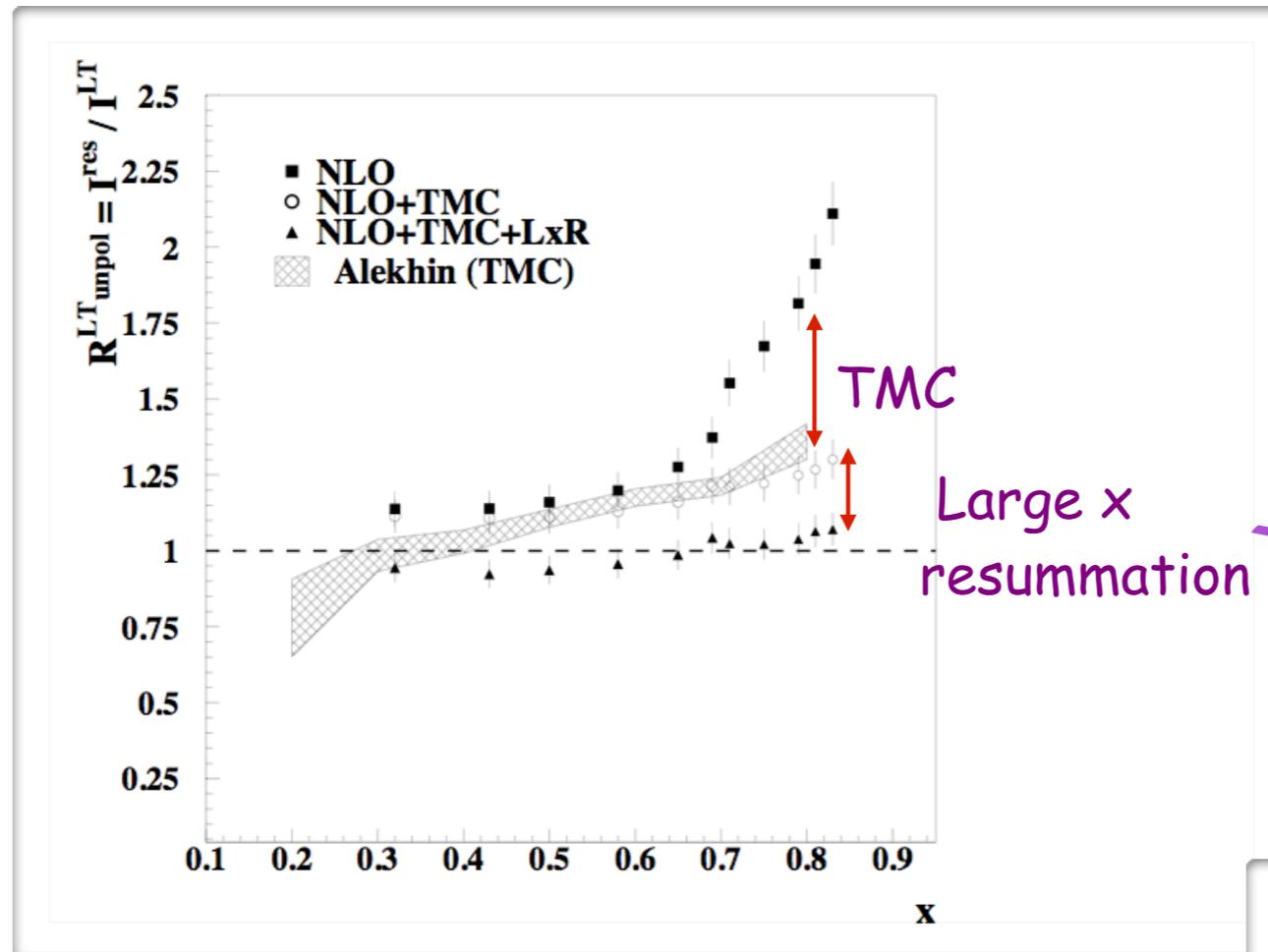
\Leftrightarrow Duality fulfilled if $R=1$

\Leftarrow LxR sensitive to α_s

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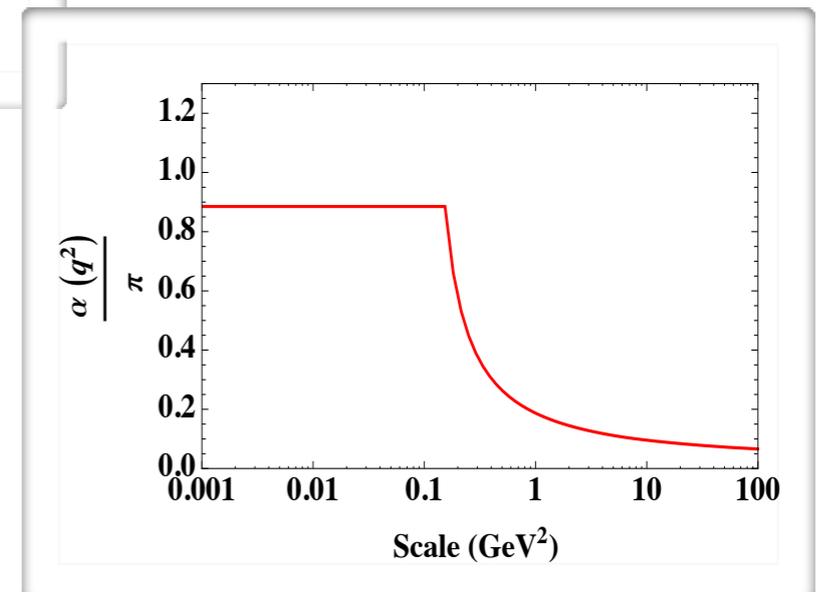
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⇔

Duality fulfilled if $R=1$

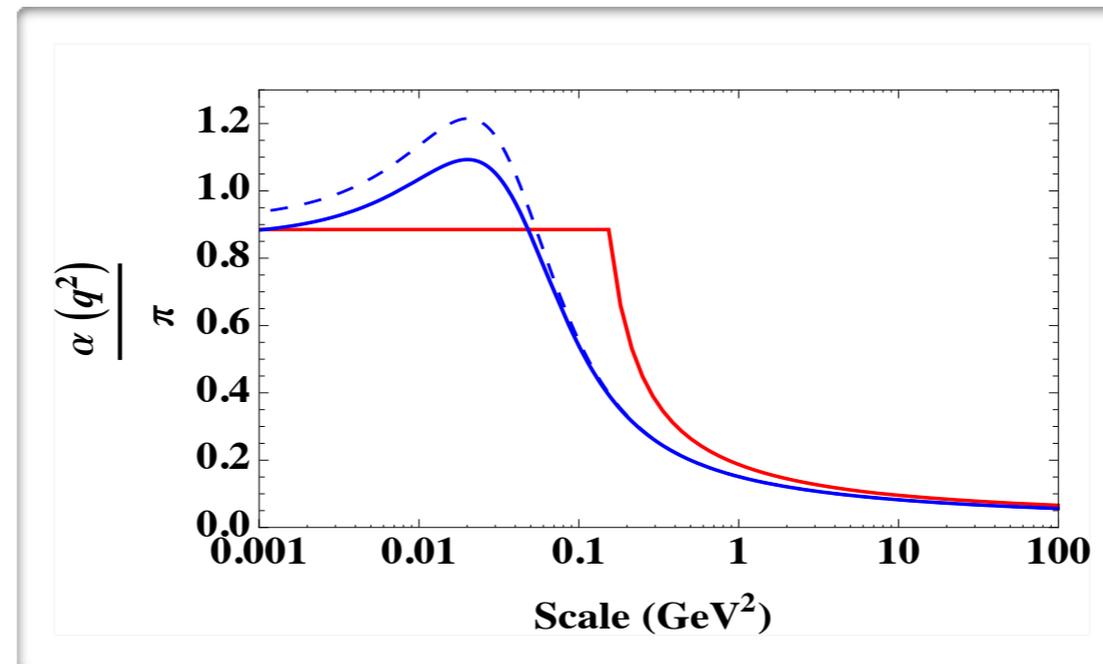
⇐

LxR sensitive to α_s



Nonperturbative Coupling Constant & LxR

How we go further : Nonperturbative Coupling Constant from DSE



Cornwall α_s^{NP}

3-4 free parameters

(up to physical constrains)

- Nonperturbative effects gathered in effective coupling α_s^{NP}
- Use of NP running coupling that scales to LO pQCD result
- Include in LxR
- **Parameterization of the realization of duality**
- Understand Higher-Twists ?
- Go to NNLO ?

Work in progress with S. Liuti

Nonperturbative QCD coupling from Phenomenology

Joint analysis: Chen, Courtoy, Deur, Liuti & Vento

Work in progress about α_s at low energy

- **Nonperturbative to perturbative transition**

- Final States Interactions and pQCD
- Errorbands to measurements (even if error on “model dependence” is immeasurable)

- **Perturbative to nonperturbative transition**

- Realization of duality & parametrization via α_s^{NP}
- New data for F_2 in the resonance region at JLab

- **How to relate the coupling constant?**

- Commensurate Scale Relations?
- RG-improved perturbation theory?

[Brodsky & Lu, Phys. Rev. D251]

[Grunberg, Phys. Rev. D29]

Extraction of α_s at low energy

- Polarized scattering from both proton and neutron

Deur et al. Phys.Lett. B650 (2007) 244-248

Natale, PoS QCD-TNT09 (2009) 031

Bjorken Sum Rule from JLab & GDH Sum Rule at $Q^2=0$ GeV²

- Deep Inelastic Scattering (DIS) at large Bjorken-x & parton-hadron duality

Liuti, [arXiv:1101.5303 [hep-ph]].

- Semi-Inclusive DIS & Extraction of T-odd TMDs from SSAs

A.C., Vento & Scopetta, Eur. Phys. J. A47, 49 (2011)

Joint analysis: Chen, Courtoy, Deur, Liuti & Vento