

Spectroscopy with dynamical Chirally Improved quarks

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July 14 2008



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Outline

1 Chirally Improved fermions

- Chirally Improved Dirac operator
- Simulation details
- Scale setting
- Variational method and interpolator basis

2 Quenched spectroscopy

- Excited mesons with derivative sources
- Ghost contributions to the isovector scalar

3 Spectroscopy with dynamical sea quarks

- Results for baryons (preliminary)
- Results for mesons (preliminary)

Chirally Improved Dirac operator

General ansatz for fermion action:

$$D_{mn} = \sum_{\alpha=1}^{16} \Gamma_\alpha \sum_{p \in \mathcal{P}_{m,n}^\alpha} c_p^\alpha \prod_{l \in p} U_l \delta_{n,m+p}$$

- Insert ansatz into Ginsparg-Wilson-equation
- Truncate the length of the contributions and compare the coefficients
- \Rightarrow Set of algebraic equations (solved by norm minimization)

The diagram shows the expansion of the Wilson term in the Ginsparg-Wilson equation. It starts with a red dashed box labeled "Wilson" containing terms $s_1 + s_2$ (with arrows pointing up and down), s_3 (with arrows pointing left and right), and s_4 (a 2x2 grid of arrows). This is followed by a plus sign and a series of terms involving γ_μ , v_i , and t_i . The first term is $+\gamma_\mu(v_1 - +)$, followed by $v_2 - +$, $v_3 - +$, and so on. Below this, there is a bracketed term $+ \gamma_\mu \gamma_\nu (t_1 - +)$, followed by $\gamma_\mu \gamma_\nu \gamma_p (a_1 \dots)$ and $\gamma_s (p_1 \dots)$.

(Gattringer, PRD63(2001)114501)

Simulation Details

- Lüscher-Weisz gauge action
- Two flavors of Chirally Improved fermions
- Stout smearing
- Hybrid Monte Carlo simulation

Earlier results (2005/2006): $12^3 \times 24$ (Lang/Majumdar/Ortner)

Presently: Three ensembles for $16^3 \times 32$

ensemble	β_{LW}	m_0	HMC time	$a[\text{fm}]^*$	$m_\pi[\text{MeV}]^*$
A	4.70	-0.05	492	0.1507(17)	526(7)
B	4.65	-0.06	993	0.1500(11)	469(4)
C	4.58	-0.077	996	0.1440(11)	318(5)

*)using $r_{0,\text{exp}} = 0.48 \text{ fm}$

Parameters: Lattice spacing

Lattice spacing determined from potential
(for hypercubic smeared configurations).

Potential fit:

$$V_L(r) = A + \frac{B}{r} + \sigma r + c_3 \Delta V(r)$$

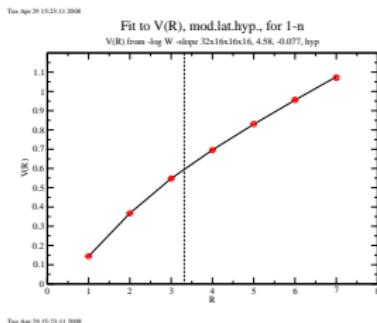
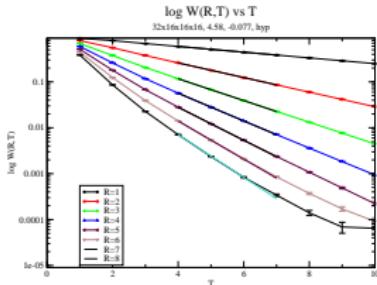
with $\Delta V(r) \equiv \left[\frac{1}{r} \right] - \frac{1}{r}$.

(Perturbative lattice Coulomb potential:

$$\left[\frac{1}{r} \right]$$

gives the Sommer parameter

$$\sqrt{\frac{1.65 + B}{\sigma}} = \frac{r_{0,\text{exp}}}{a}$$



Fit to $\ln W(r, t)$ in the range
 $4 \leq t \leq 7$ (upper plot) and
to the potential in the range
 $1 \leq r \leq 7$ (lower plot).

Variational method (C.Michael; Lüscher and Wolff)

- Matrix of correlators projected to fixed momentum (will assume 0)

$$C(t)_{ij} = \sum_n \langle 0 | O_i | n \rangle \langle n | O_j^\dagger | 0 \rangle$$

- Solve the generalized eigenvalue problem:

$$\begin{aligned} C(t) \vec{v}_k &= \lambda_k(t) C(t_0) \vec{v}_k \\ \lambda_k(t) &\propto e^{-tM_k} \left(1 + \mathcal{O}\left(e^{-t\Delta M_k}\right) \right) \end{aligned}$$

- At large time separation: only a single mass in each eigenvalue.
- Eigenvectors can serve as a fingerprint.

Quark sources

- Jacobi smeared quark sources, e.g., $u_s \equiv (S u)_x$

$$S = M S_0 \quad \text{with} \quad M = \sum_{n=0}^N \kappa^n H^n$$

$$H(\vec{n}, \vec{m}) = \sum_{j=1}^3 \left(U_j(\vec{n}, 0) \delta(\vec{n} + \hat{j}, \vec{m}) + U_j(\vec{n} - \hat{j}, 0)^\dagger \delta(\vec{n} - \hat{j}, \vec{m}) \right).$$

- Fewer quark propagators
- Combination allows nodes in the interpolating operators
- Derivative quark sources W_{d_i} :

$$D_i(\vec{x}, \vec{y}) = U_i(\vec{x}, 0) \delta(\vec{x} + \hat{i}, \vec{y}) - U_i(\vec{x} - \hat{i}, 0)^\dagger \delta(\vec{x} - \hat{i}, \vec{y}),$$
$$W_{d_i} = D_i S_w.$$



Mesons from quenched ensemble

Mesons with derivative sources: Pion channel

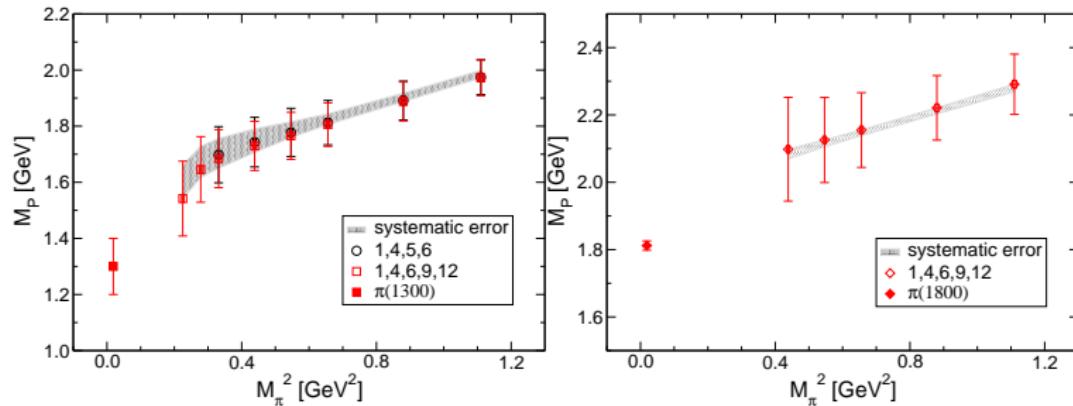


Figure: 1st and 2nd excitation of π

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD(2008)

Mesons from quenched ensemble

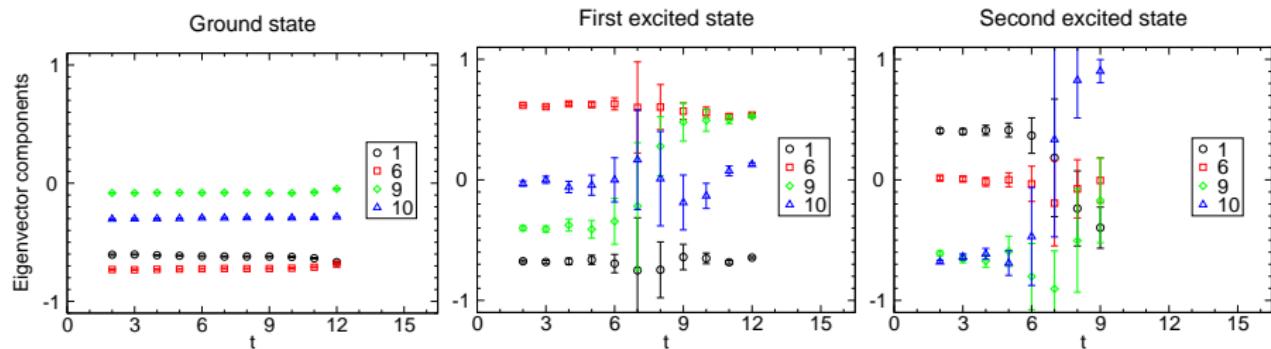


Figure: Eigenvector components for ground state and lowest excitations

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD(2008)

Mesons from quenched ensemble

Mesons with derivative sources: Isovector Scalar 1 0^{++}

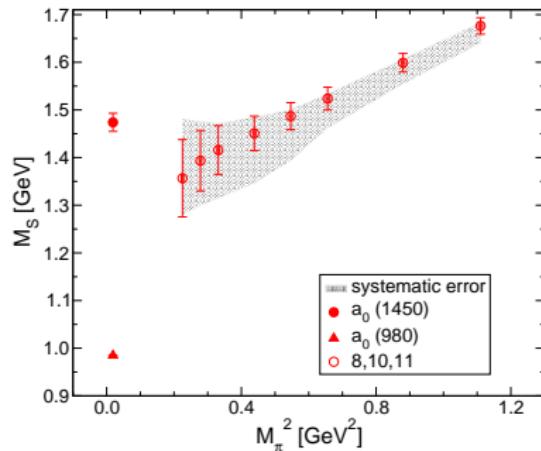
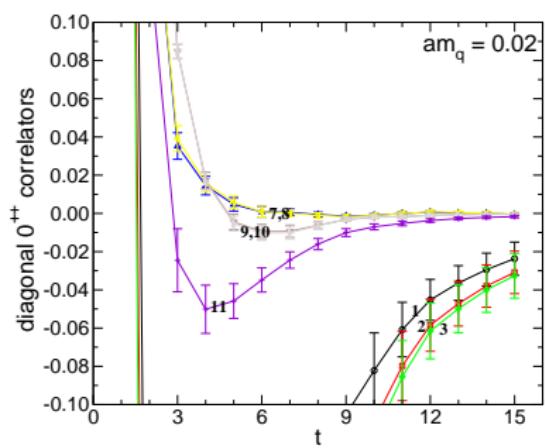


Figure: Diagonal correlators and ground state in the 0^{++} channel

Gattringer et al., ArXiv:0802.2020 [hep-lat], PRD(2008)

Run parameters

- Lattice size $16^3 \times 32 \sim (2.4 \text{ fm})^3 \times 4.8 \text{ fm}$
- We analyze every fifth configuration and shift the sources for consecutive configurations.
- Currently only preliminary data
- We also show results with larger valence quark masses: “partially quenched”.

set	β_{LW}	m_0	# conf.s	a [fm]*	m_{AWI} [MeV]*	m_π [MeV]*	$a m_\pi L$
A	4.70	-0.05	100/100	0.1507(17)	43.0(4)	526(7)	6.4
B	4.65	-0.06	100/200	0.1500(11)	35.1(2)	469(4)	5.8
C	4.58	-0.077	100/200	0.1440(11)	15.0(4)	318(5)	3.8

Nucleon - positive parity

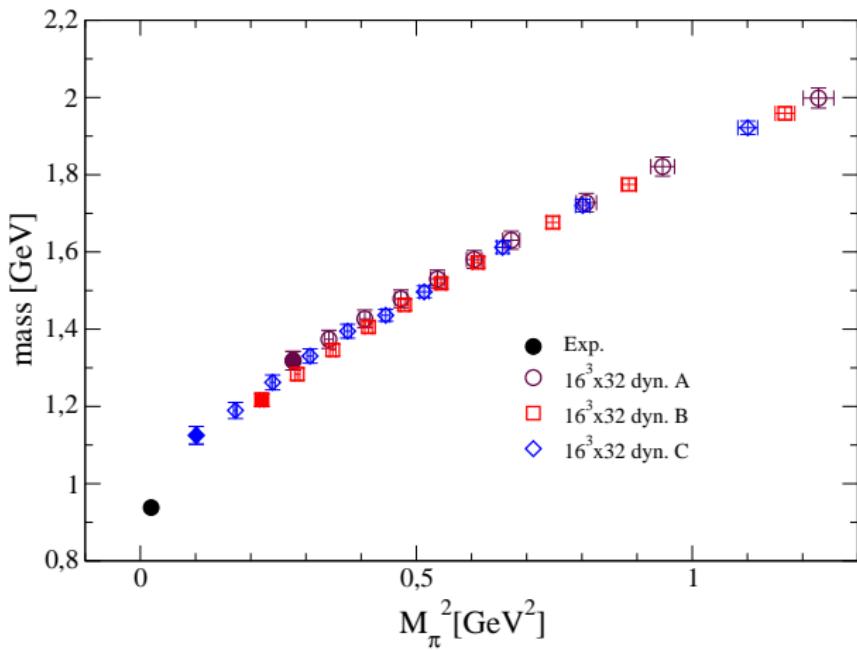


Figure: Mass of the nucleon ground state

Nucleon - negative parity

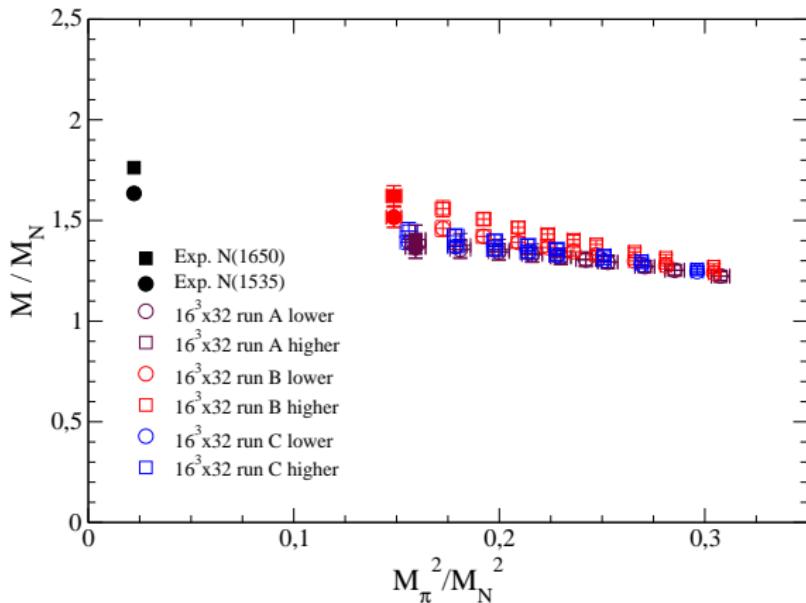
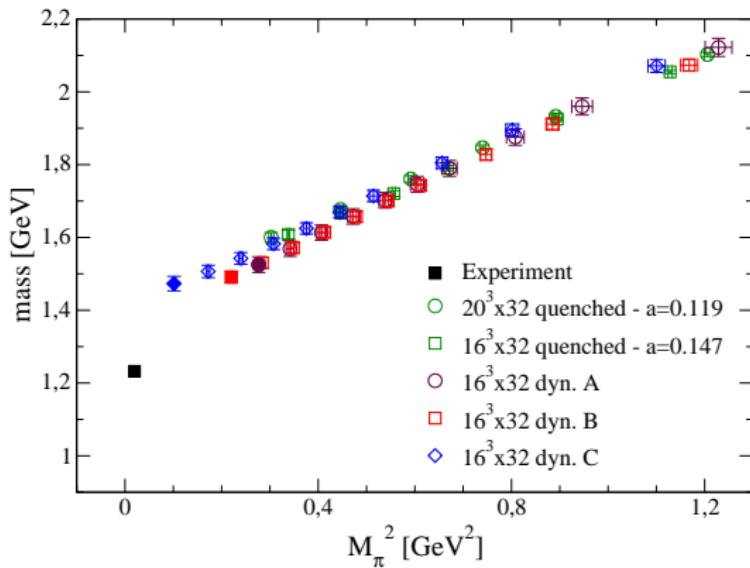


Figure: Mass of the negative parity ground state and 1st excited state

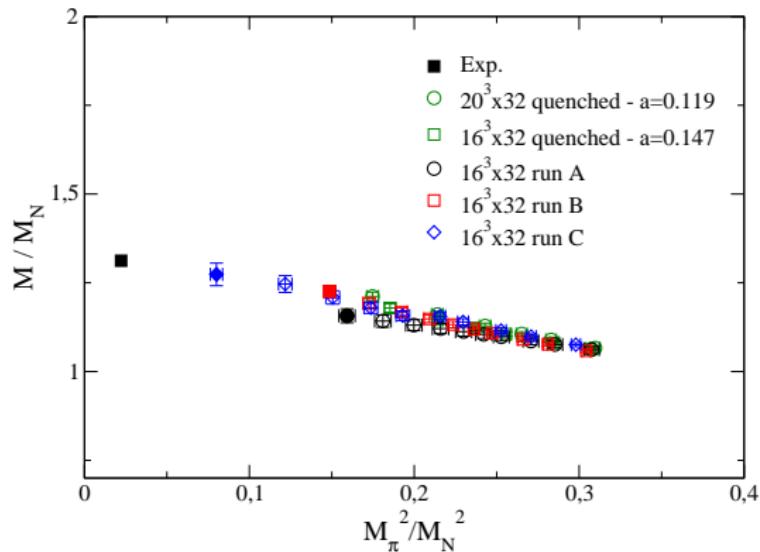
No reasonable signal for run C at small valence quark masses.

Delta (pos. parity, ground state)



..too high at small quark masses...

Delta (pos. parity)



High masses in N and Δ^{++} cancel...

rho

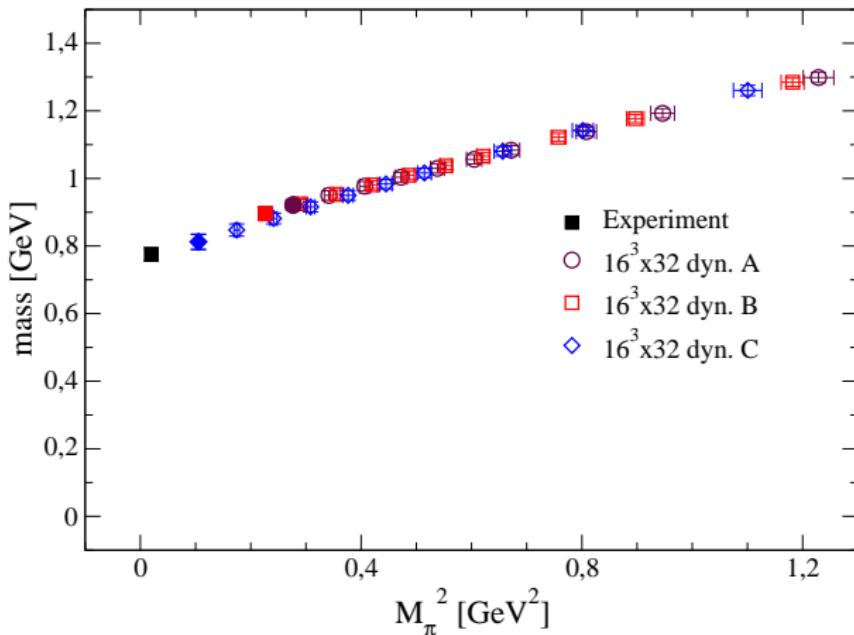


Figure: Ground state in the 1^{--} channel ($\rho(772)$)

Scalar meson a_0

Quenched: Extrapolates to $a_0(1450)$

Dynamical: Unclear signal; operator dependent 10% variation.

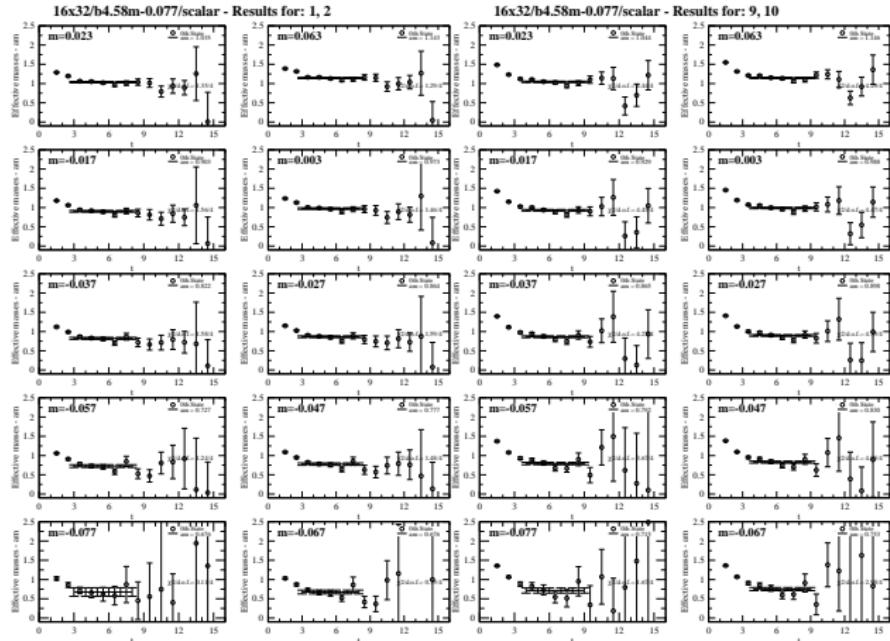


Figure: Example: ground state signals for two combinations, ensemble C.

Scalar meson a_0

Mass of a_0 (interpolators 9, 10)

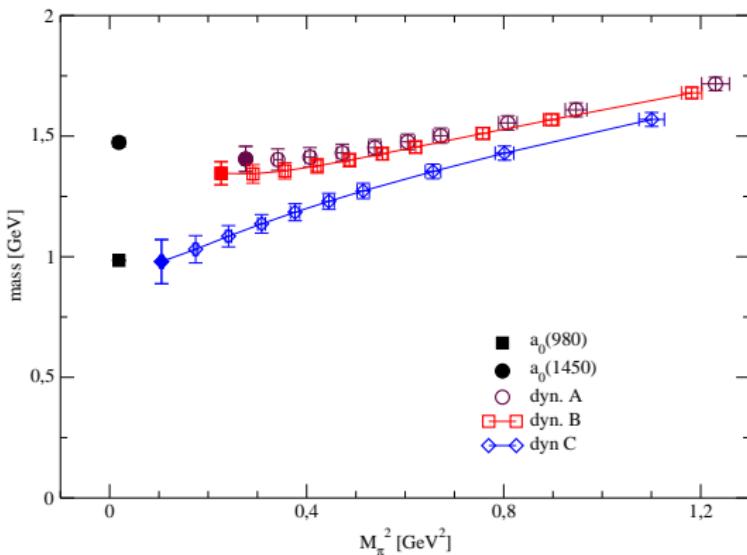


Figure: a_0 from one possible set of interpolators

Do we see a level crossing here ($\pi\eta$ -channel?)

Summary

- Progress report on spectroscopy with dynamical CI fermions
 - Good signals for (most) ground state mesons and baryons
 - Demonstrated power of the method for excited states (quenched)
 - So far: unclear signal for excited states (full QCD)
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- Poster by [Christian Hagen](#) (Regensburg) on heavy-light mesons
 - Talk by [Sasa Prelovsek](#) on tetraquarks with quenched CI fermions

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