A Variational Study of the Nucleon Wavefunction

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outline

- physical motivation
- variational method
- exploratory calculation
 - parameters explored
 - results
- summary

Physical Motivation

- spatial nucleon wavefunction (quarks)
- gluons
- diquarks
 - "good" (scalar) diquarks $(u C\gamma_5 d)$
 - "bad" (vector) diquarks $(u C \gamma_{\mu} d)$
 - motivated by one-gluon exchange
- flux tube model (c.f. Regge trajectories)



The Method

For a trial source

$$|\mathbf{J}\rangle = \mathbf{J} |\Omega\rangle$$

we calculate the (momentum-projected) twopoint function

$$\langle \overline{\mathbf{J}}(t) | \mathbf{J}(0) \rangle = \sum_{n} |\langle \overline{\mathbf{J}}(0) | n \rangle|^2 e^{-E_n t}$$

For large t, the ground state dominates the sum...

... so we fit the two-point function to a sum of exponentials:

$$C(t) \equiv \langle \overline{\mathbf{J}}(t) | \mathbf{J}(0) \rangle \stackrel{\cdot}{=} A_0 e^{-m_0 t} + A_1 e^{-m_1 t}$$

and estimate the (normalized) groundstate overlap for our trial source:

$$\frac{\left|\langle \bar{\mathbf{J}}(0)|0\rangle\right|^2}{\sum_n \left|\langle \bar{\mathbf{J}}(0)|n\rangle\right|^2} = \frac{A_0}{C(0)} \int_{-12}^{-7}$$



computational side note:

propagator correction at $t = t_{source}$

Using the "naive" Wilson propagator in the calculation of the two-point function is incorrect at t_{source} ! This can be seen by examining the transfer matrix formalism:

$$\langle \hat{O}_1 \hat{O}_2 \rangle = \frac{\int DU D\bar{q} Dq \, e^{-S(U,\bar{q},q)} O_1 O_2}{\int DU D\bar{q} Dq \, e^{-S(U,\bar{q},q)}} = \frac{\langle 0 | N\{\hat{T} \dots \hat{T} \hat{O}_1 \hat{T} \dots \hat{T} \hat{O}_2 \hat{T} \dots \hat{T}\} | 0 \rangle}{\langle 0 | \hat{T} \dots \hat{T} | 0 \rangle}$$

We want to find T (and appropriate coordinate transform) such that we recover the action exponential in the path integral.

Writing out the Wilson action, one finds that it is necessary to define the following normal ordering convention for the fermion field operators:

upper:	$N\{q^U \bar{q}^U\}$	=	$q^U \bar{q}^U$
lower:	$N\{q^L\bar{q}^L\}$	=	$\bar{q}^L q^L$

... where upper (lower) refers to the upper (lower) half spinor in the Dirac basis. The operators anticommute **except** when on the same timeslice, in which case we pick up a correction term from the anticommutator:

$$P_{correct}(x,y) = P_{naive}(x,y) - \frac{(1-\gamma_0)}{2}B^{-1}(\vec{x},\vec{y})\,\delta_{x_0,y_0}$$

where
$$B(\vec{x},\vec{y}) = \mathbb{I} - \kappa \sum_{j=1}^3 \left[U_j^{\dagger}(y)\delta_{x-y-\hat{j}} + U_j(x)\delta_{x-y+\hat{j}} \right]$$

Exploratory Calculation:

• Quenched Wilson, 16³x32 lattice

 $\beta = 6.0, \ \kappa = 0.1530 \ (m_{\pi} \approx 900 \ \text{MeV})$

- Why not use domain wall fermions?
 - nonlocal transfer matrix
 - oscillating terms (SS & JN, hep-lat: 0710.0425)

Exploratory Calculation:

Source interpolating field motivated by diquark model of nucleon:

 $J=(~U~C\gamma_5~D~)~U$

- Variational parameters:
 - spatial extent of quark sources (quark smearing)
 - number of dirac components (2 vs. 4)
 - gluon wavefunction (gauge field smearing)
 - relative size of quark, diquark
 - relative position of quark, diquark

Quark Smearing

• Gauge-invariant smearing:



• Fix
$$\alpha = 3.0$$
, vary N

• four-component spinors: variation with quark radius



two-component versus four-component spinors



• using only two components: $0.35 \rightarrow 0.50$!

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Apparently, we've been using a very poor trial wavefunction in the lower components. For any dirac spinor:

$$u = \left(\begin{array}{c} u^U \\ u^L \end{array}\right)$$

the dirac equation implies:

$$u^L = \frac{1}{E+m} (\mathbf{p} \cdot \sigma) u^U$$

So, if we put the upper components into something like an s-wave state, the lower components ought to be in a p-wave. But we've been setting: $u^L \approx u^U$, which is clearly a bad approximation to the true wavefunction. It's better to set $u^L = 0$ than to introduce so many excited state contaminants.

Add Gauge Smearing:

• Smear gauge links used in constructing the source:

$$\left(\mid \rightarrow (1-c) \mid + c \sum \left[\right]^{N} \right)^{N}$$

• Parameters used:

$$c = 0.26, N = 25$$









• Improves overlap: $0.50 \rightarrow 0.80$

Change Relative Size of Quark, Diquark:

- Perhaps quarks in diquark feel a different potential than single quark does....
- Use trial source:

$$\mathbf{J} = (\mathbf{U}_{s1} \mathbf{C} \gamma_5 \mathbf{D}_{s1}) \mathbf{U}_{s2}$$

 Want to find maximum overlap in two-dimensional parameter space





Try displacing quark:

- Motivated by flux tube model
- Use a source with quark, diquark separated: $J = (U_{s1}(x) C\gamma_5 D_{s1}(x)) U_{s2}(y)$
- Symmetrize sink:

 $C(t) = \sum_{\vec{r}} \sum_{\hat{j}} \langle (\bar{u}(\vec{r},t)C\gamma_5 \bar{d}(\vec{r},t))\bar{u}(\vec{r}+\ell\hat{j},t) (u(0,0)C\gamma_5 d(0,0))u(\ell\hat{x},0) \rangle$



upper components only



quark smearing, RMS



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 - 2-vs. 4-component:
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