Vector meson electromagnetic form factors

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Lattice 2008, College of William & Mary



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Form factors

- Internal structure of hadrons \rightarrow (generalised) form factors
- low energy quantities \rightarrow lattice
- nucleon, pion, ρ meson
- heavy pions, ρ meson stable
- representative for spin 1 particle

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Form factors

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interaction hadron - e.m. current

$$egin{aligned} &\langle m{p}',m{s}'\,|m{J}^lpha|\,m{p},m{s}
angle = \ &\left(2\sqrt{E_
ho(m{p}')E_
ho(m{p})}
ight)^{-1}\epsilon_ au'^{\,*}(m{p}',m{s}')J^{ aulpha\sigma}(m{p}',m{p})\epsilon_\sigma(m{p},m{s}) \end{aligned}$$

for spin one particle parametrised by three form factors

$$J^{\tau\alpha\sigma}(p',p) = -G_{1}\left(Q^{2}\right)g^{\tau\sigma}\left(p^{\alpha}+p'^{\alpha}\right)$$
$$-G_{2}\left(Q^{2}\right)\left(g^{\alpha\sigma}q^{\tau}-g^{\alpha\tau}q^{\sigma}\right)$$
$$+G_{3}\left(Q^{2}\right)\left(q^{\sigma}q^{\tau}\frac{p^{\alpha}+p'^{\alpha}}{2m_{\rho}^{2}}\right)$$
$$= -g^{2} - (p'-\rho)^{2}$$

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Form factors

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Interesting static quantities

- charge radius $\langle r^2 \rangle = -6 \frac{\partial G_C}{\partial (Q^2)} |_{Q^2=0}$
- magnetic moment $\mu_M = \frac{e}{2m_o}G_M(0)$, aka *g* factor
- quadrupole moment $\mu_Q = \frac{e}{m_a^2} G_Q(0)$

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What to expect

- Samsonov et al: $\mu_M = 1.8(3)$ (QCD sum rules in ext. fields)JHEP 0312:061,2003
- Bhagwat *et al*: $\langle r^2 \rangle = 0.54 \text{ fm}^2, \mu_M = 2.01, \mu_Q = -0.41 \text{ fm}^2$ (Dyson-Schwinger eqs.) Phys.Rev.C77:025203,2008
- Aliev et al:G_M/G_C > 2 (light cone sum rules; don't work at small Q²)_{Phys.Rev.D70:094007,2004}
- Choi *et al*: μ_M = 1.92, μ_Q = -0.43 fm² (Light front quark model)_{Phys.Rev.D70:053015,2004}
- Alexandrou *et al:* ρ is non-spherical; (density-density correlators, lattice)_{Phys.Rev.D66:094503,2002}
- Hedditch *et al*: $\langle r^2 \rangle \sim 0.6 \text{ fm}^2$, $\mu_M \sim 2.3$, $\mu_Q \sim -0.005 \text{ fm}^2$ (quenched lattice simulation, standard 3pt technique)_{Phys.Rev.D75:094504,2007}

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Lattice method

compute three point functions involving $\langle p', s' | J^{\alpha} | p, s \rangle$ (and two point functions)

system of equations $R^{lpha}_{\mu
u}(p,p') = \sum_i c_i G_i$ for each Q^2

solve numerically (χ^2 minimisation) $\rightsquigarrow G_i(Q^2)$



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Lattice matrix elements

can be extracted from three point functions



transfer matrix formalism; 0 $<< \tau << t$

$$\lim_{t\to\infty} G_{\mu\nu}(t,\vec{p}) = -\frac{e^{-E_{\rho}(\vec{p})\,t}}{2E_{\rho}(\vec{p})}\lambda_{\rho}(\vec{p})\bar{\lambda}_{\rho}(\vec{p})\left(g_{\mu\nu} - \frac{p_{\mu}p_{\nu}}{m_{\rho}^2}\right)$$

$$\begin{split} \lim_{\substack{\tau \to \infty \\ t - \tau \to \infty}} G^{\alpha}_{\mu\nu}(t,\tau,\vec{p}',\vec{p}) &= \frac{e^{-E_{\rho}(\vec{p}')(t-\tau)}e^{-E_{\rho}(\vec{p})\tau}}{4E_{\rho}(\vec{p}')E_{\rho}(\vec{p})}\lambda_{\rho}(\vec{p}')\bar{\lambda}_{\rho}(\vec{p}) \\ &\times \left(g_{\mu\tau} - \frac{p'_{\mu}p'_{\tau}}{m_{\rho}^2}\right)J^{\tau\alpha\sigma}\left(g_{\sigma\nu} - \frac{p_{\sigma}p_{\nu}}{m_{\rho}^2}\right) \end{split}$$

$$\begin{aligned} \lambda \text{-overlap of interpolating operator } \chi_{\mu}^{\,\prime} &= d\gamma_{\mu} u \text{ with } \rho \\ \left\langle \Omega \left| \chi_{\mu}(0) \right| \rho(\vec{p}, s) \right\rangle &= \sqrt{2E_{\rho}(\vec{p})}^{-1} \lambda_{\rho}(\vec{p})\epsilon_{\mu}(\rho, s) \end{aligned}$$

sum over polarisations using transversality condition $\sum_s \epsilon_\mu(p,s) \epsilon^*_\nu(p,s) = -g_{\mu\nu} + p_\mu p_\nu / m_\rho^2$

 $-g_{\mu\nu} + \rho_{\mu}\rho_{\nu} / m_{\rho}^{2}$ Technische Universität München
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Ratios

$$R^{\alpha}_{\mu\nu}(\tau,\vec{p}',\vec{p}) = \frac{G^{\alpha}_{\mu\nu}(t,\tau,\vec{p}',\vec{p})}{G_{\mu\mu}(t,\vec{p}')} \sqrt{\frac{G_{\nu\nu}(t-\tau,\vec{p})G_{\mu\mu}(\tau,\vec{p}')G_{\mu\mu}(t,\vec{p}')}{G_{\nu\nu}(\tau,\vec{p})G_{\mu\mu}(t-\tau,\vec{p}')G_{\nu\nu}(t,\vec{p})}}$$

is independent of $\boldsymbol{\tau}$

 $(\mu, \nu = 1...3)$ t fixed;potential problems with $\sqrt{}$, argument can be negative

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Details of the lattice calculation

- QCDSF-UKQCD configurations
- 2 dynamical flavours of Wilson fermions
- non-perturbatively improved Dirac operator $i/4c_{SW}ag^2\bar{\psi}(x)\sigma_{\mu\nu}F_{\mu\nu}(x)\psi(x), c_{SW}(g)$ (ALPHA coll.)
- (Jacobi) smeared sources and sinks
- local vector current \rightsquigarrow renormalisation $Z_V = 1/G_1^{\text{unren}}(0)$
- compute for 3 values of \vec{p}' and 17 values of \vec{p} and all polarisation combinations
- no disconnected contribution (G^{disc}(U) = -G^{disc}(U^{*})), both have equal weight

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Volume	β	κ	m_{π} [MeV]	a [fm]
16 ³ 32	5.29	0.13500	929(2)	0.089
16 ³ 32	5.29	0.13550	784(3)	0.084
24 ³ 48	5.29	0.13590	591(2)	0.080
24 ³ 48	5.29	0.13620	406(6)	0.077

two-point function example



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Ratio example

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Example for form factor fits

 $\beta = 5.29$ $\kappa = .13620$ Vol $24^{3}48$ $m_{\pi} = 406 \text{MeV}$



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Charge radii



Results

Charge radii



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Magnetic moment



Magnetic moment



 g_{ρ}

Results

Quadrupole moment



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Quadrupole moment



- first unquenched direct computation of the vector meson e.m.form factors
- still preliminary
- Charge radii
 - slightly larger than found by Hedditch et al (larger Q² range)
 - growing towards smaller m_q
- g-factor
 - $m \circ ~ \sim$ 2; close to quark model expectation
 - o chiral curvature?
- quadrupole moment
 - $\bullet~\sim$ 0 at large pion masses
 - decreasing quark mass: increasingly negative → oblate shape

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• next: axial/tensor form factors; GFF

When does the ρ decay?



Another ratio example

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