

Efficient use of the Generalised Eigenvalue Problem

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in collaboration with

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Williamsburg, July 2008

The GEVP
Outline of a proof
Application to HQET

The GEVP

matrix of correlation functions on an infinite time lattice

$$C_{ij}(t) = \langle O_i(0)O_j(t) \rangle = \sum_{n=1}^{\infty} e^{-E_n t} \psi_{ni} \psi_{nj}, \quad i, j = 1, \dots, N$$

$$\psi_{ni} \equiv (\psi_n)_i = \langle n | \hat{O}_i | 0 \rangle = \psi_{ni}^* \quad E_n \leq E_{n+1}$$

the GEVP is

$$C(t) v_n(t, t_0) = \lambda_n(t, t_0) C(t_0) v_n(t, t_0), \quad n = 1, \dots, N \quad t > t_0,$$

Lüscher & Wolff showed that

$$E_n^{\text{eff}} = \frac{1}{a} \log \frac{\lambda_n(t, t_0)}{\lambda_n(t+a, t_0)} = E_n + \varepsilon_n(t, t_0)$$

$$\varepsilon_n(t, t_0) = O(e^{-\Delta E_n (t-t_0)}), \quad \Delta E_n = \left| \min_{m \neq n} E_m - E_n \right|.$$

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correction term

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- ▶ problematic when ΔE_n is small
- ▶ no gain for the ground state ??

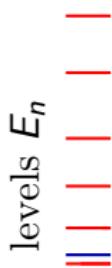


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- ▶ mostly: no statement about corrections
- ▶ now we say: ...



The GEVP: a simplified situation

with only N states

$$C_{ij}^{(0)}(t) = \sum_{n=1}^N e^{-E_n t} \psi_{ni} \psi_{nj}.$$

the dual (time-independent) vectors are defined by

$$(u_n, \psi_m) = \delta_{mn}, \quad m, n \leq N. \quad (u_n, \psi_m) \equiv \sum_{i=1}^N (u_n)_i \psi_{mi}$$

one then has

$$\begin{aligned} C^{(0)}(t) u_n &= e^{-E_n t} \psi_n, \\ C^{(0)}(t) u_n &= \lambda_n^{(0)}(t, t_0) C^{(0)}(t_0), \\ \lambda_n^{(0)}(t, t_0) &= e^{-E_n(t-t_0)}, \quad v_n(t, t_0) \propto u_n \end{aligned}$$

an orthogonality for all t

$$(u_m, C^{(0)}(t) u_n) = \delta_{mn} \rho_n(t), \quad \rho_n(t) = e^{-E_n t}.$$

The GEVP with N states

$$\hat{Q}_n = \sum_{i=1}^N (u_n)_i \hat{O}_i \equiv (\hat{O}, u_n),$$

create the eigenstates of the Hamilton operator

$$|n\rangle = \hat{Q}_n |0\rangle, \hat{H}|n\rangle = E_n |n\rangle.$$

So matrix elements are

$$p_{0n} = \langle 0 | \hat{P} | n \rangle = \langle 0 | \hat{P} \hat{Q}_n | 0 \rangle$$

generalization:

$$\begin{aligned} p_{0n} &= \langle P(t) O_j(0) \rangle (u_n)_j = \frac{\langle P(t) Q_n(0) \rangle}{\langle Q_n(t) Q_n(0) \rangle^{1/2}} e^{E_n t/2} \\ &= \frac{\langle P(t) O_j(0) \rangle v_n(t, t_0)_j}{(v_n(t, t_0), C(t) v_n(t, t_0))^{1/2}} \frac{\lambda_n(t_0 + t/2, t_0)}{\lambda_n(t_0 + t, t_0)} \end{aligned}$$

Corrections are due to the excited states

$$C_{ij}^{(1)}(t) = \sum_{n=N+1}^{\infty} e^{-E_n t} \psi_{ni} \psi_{nj}$$

Perturbation theory

[Ferenc Niedermayer & Peter Weisz, 1998, unpublished]

$$A v_n = \lambda_n B v_n, \quad A = A^{(0)} + \epsilon A^{(1)}, \quad B = B^{(0)} + \epsilon B^{(1)}.$$

We will later set

$$\begin{aligned} A^{(0)} &= C^{(0)}(t), & \epsilon A^{(1)} &= C^{(1)}(t), \\ B^{(0)} &= C^{(0)}(t_0), & \epsilon B^{(1)} &= C^{(1)}(t_0) \end{aligned}$$

$$(v_n^{(0)}, B^{(0)} v_m^{(0)}) = \rho_n \delta_{nm}.$$

$$\begin{aligned} \lambda_n &= \lambda_n^{(0)} + \epsilon \lambda_n^{(1)} + \epsilon^2 \lambda_n^{(2)} \dots \\ v_n &= v_n^{(0)} + \epsilon v_n^{(1)} + \epsilon^2 v_n^{(2)} \dots \end{aligned}$$

Perturbation theory

$$A^{(0)} v_n^{(1)} + A^{(1)} v_n^{(0)} = \lambda_n^{(0)} \left[B^{(0)} v_n^{(1)} + B^{(1)} v_n^{(0)} \right] + \lambda_n^{(1)} B^{(0)} v_n^{(0)},$$

$$A^{(0)} v_n^{(2)} + A^{(1)} v_n^{(1)} = \lambda_n^{(0)} \left[B^{(0)} v_n^{(2)} + B^{(1)} v_n^{(1)} \right] + \lambda_n^{(1)} \left[B^{(0)} v_n^{(1)} + B^{(1)} v_n^{(0)} \right] -$$

solve using orthogonality $(v_n^{(0)}, v_m^{(0)}) = \delta_{mn} \rho_n$

$$\lambda_n^{(1)} = \rho_n^{-1} \left(v_n^{(0)}, \Delta_n v_n^{(0)} \right), \quad \Delta_n \equiv A^{(1)} - \lambda_n^{(0)} B^{(1)}$$

$$v_n^{(1)} = \sum_{m \neq n} \alpha_{nm}^{(1)} \rho_m^{-1/2} v_m^{(0)}, \quad \alpha_{nm}^{(1)} = \rho_m^{-1/2} \frac{\left(v_m^{(0)}, \Delta_n v_n^{(0)} \right)}{\lambda_n^{(0)} - \lambda_m^{(0)}}$$

$$\lambda_n^{(2)} = \sum_{m \neq n} \rho_n^{-1} \rho_m^{-1} \frac{\left(v_m^{(0)}, \Delta_n v_n^{(0)} \right)^2}{\lambda_n^{(0)} - \lambda_m^{(0)}} - \rho_n^{-2} \left(v_n^{(0)}, \Delta_n v_n^{(0)} \right) \left(v_n^{(0)}, B^{(1)} v_n^{(0)} \right).$$

And a recursion formula for the higher order coefficients

Perturbation theory

insert specific case and use (for $m > n$)

$$\begin{aligned}(\lambda_n^{(0)} - \lambda_m^{(0)})^{-1} &= (\lambda_n^{(0)})^{-1} (1 - e^{-(E_m - E_n)(t - t_0)})^{-1} \\&= (\lambda_n^{(0)})^{-1} \sum_{k=0}^{\infty} e^{-k(E_m - E_n)(t - t_0)}\end{aligned}$$

we get

Perturbation theory

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we get

$$\begin{aligned} \varepsilon_n(t, t_0) &= O(e^{-\Delta E_{N+1,n} t}), \quad \Delta E_{m,n} = E_m - E_n, \\ \pi_n(t, t_0) &= O(e^{-\Delta E_{N+1,n} t_0}), \quad \text{at fixed } t - t_0 \\ \pi_1(t, t_0) &= O(e^{-\Delta E_{N+1,1} t_0} e^{-\Delta E_{2,1}(t-t_0)}) + O(e^{-\Delta E_{N+1,1} t}). \end{aligned}$$

to all orders in the convergent expansion

Effective theory to first order

$$C_{ij}(t) = C_{ij}^{\text{stat}}(t) + \omega C_{ij}^{1/\text{m}}(t) + O(\omega^2)$$

$$E_n^{\text{eff,stat}}(t, t_0) = \log \frac{\lambda_n^{\text{stat}}(t, t_0)}{\lambda_n^{\text{stat}}(t+1, t_0)} = E_n^{\text{stat}} + O(e^{-\Delta E_{N+1,n}^{\text{stat}} t}),$$

$$\begin{aligned} E_n^{\text{eff},1/\text{m}}(t, t_0) &= \frac{\lambda_n^{1/\text{m}}(t, t_0)}{\lambda_n^{\text{stat}}(t, t_0)} - \frac{\lambda_n^{1/\text{m}}(t+1, t_0)}{\lambda_n^{\text{stat}}(t+1, t_0)} \\ &= E_n^{1/\text{m}} + O(t e^{-\Delta E_{N+1,n}^{\text{stat}} t}). \end{aligned}$$

$$C^{\text{stat}}(t) v_n^{\text{stat}}(t, t_0) = \lambda_n^{\text{stat}}(t, t_0) C^{\text{stat}}(t_0) v_n^{\text{stat}}(t, t_0),$$

$$\lambda_n^{1/\text{m}}(t, t_0) = \left(v_n^{\text{stat}}(t, t_0), [C^{1/\text{m}}(t) - \lambda_n^{\text{stat}}(t, t_0) C^{1/\text{m}}(t_0)] v_n^{\text{stat}}(t, t_0) \right)$$

Numerics: only static GEVP needs to be solved

Demonstration in HQET

- ▶ $(1.5 \text{ fm})^3 \times 3 \text{ fm}$ with pbc (apbc for quarks), quenched
- ▶ $L/a = 16, L/a = 24$, i.e. 0.1 fm and 0.07 fm lattice spacing, with $\kappa = 0.133849, 0.1349798$ (strange quark)
- ▶ all-to-all propagators for the strange quark [Dublin]
50 (approximate) low modes and 2 noise fields
100 configs each
- ▶ 8 levels of Wuppertal smearing (Gaussian) [Güsken, Löw, Mütter, Patel, Schilling, S., 1989] after gauge-field APE smearing [Basak et al., 2006]
- ▶ Truncate to a $N \times N$ projecting with the N eigenvectors of $C(a)$ with the largest eigenvalues

[Niedermayer, Rufenacht, Wenger, 2000]

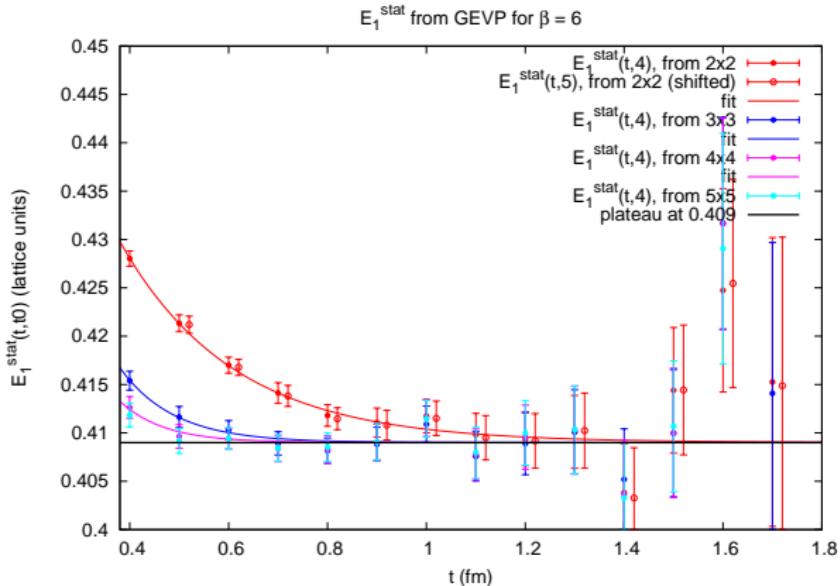
Demonstration in HQET

$$a = 0.1 \text{ fm}$$

$$aE_1^{\text{eff,stat}}(t, t_0)$$

curve:

$$E_1 + \alpha_N e^{-\Delta E_{N+1,1} t}$$



$\Delta E_{N+1,1}$ agree with plateaux of $E_{N+1}^{\text{eff,stat}}(t, t_0)$ for large N' and t .

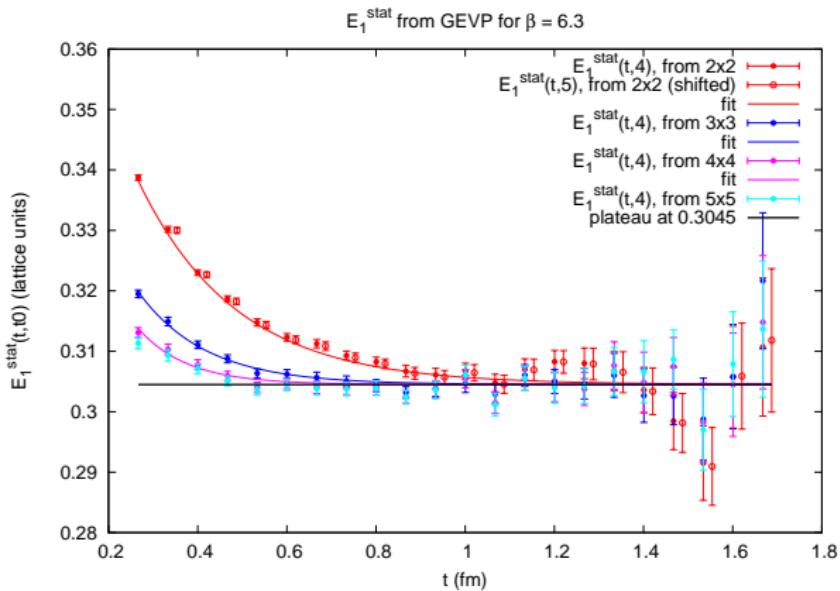
Demonstration in HQET

$$a = 0.07 \text{ fm}$$

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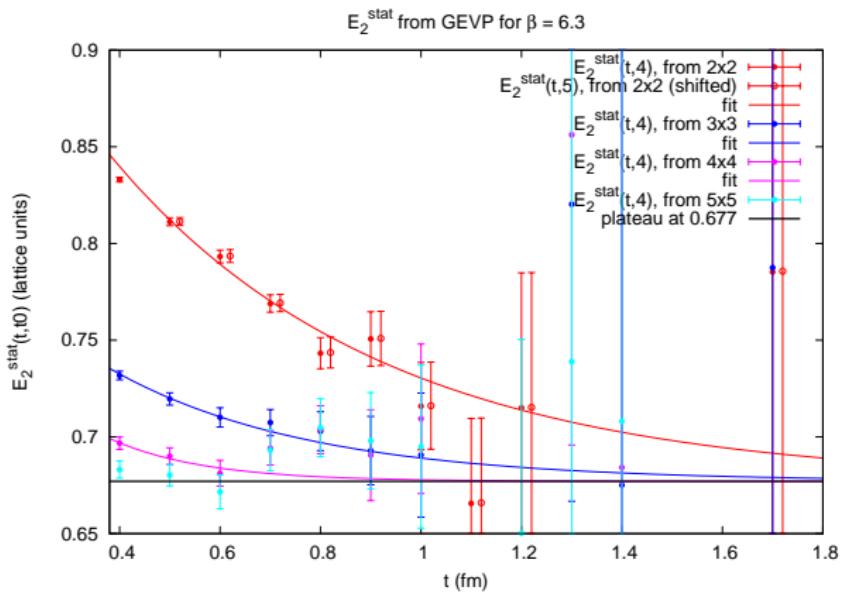
Demonstration in HQET

$$a = 0.1 \text{ fm}$$

$$aE_2^{\text{eff,stat}}(t, t_0)$$

curve:

$$E_2 + \alpha_N e^{-\Delta E_{N+1,2} t}$$



$\Delta E_{N+1,2}$ agree with plateaux of $E_{N+1}^{\text{eff,stat}}(t, t_0)$ for large N' and t .

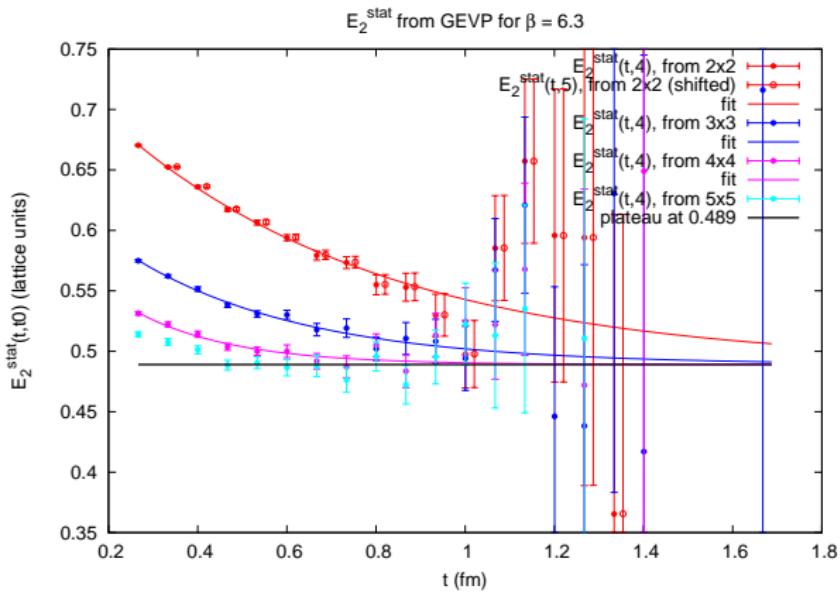
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Energy levels

how well do different determinations agree?

β	method	N	m	$r_0 \Delta E_{m,1}$
6.0219	plateau	6	3	2.3
	corr. to E_1	2	3	2.3
	corr. to E_2	2	3	2.3
6.2885	corr. to E_1	2	3	2.3
	corr. to E_2	2	3	2.2
	plateau	6	4	3.3
6.0219	corr. to E_1	3	4	4
	corr. to E_2	3	4	2.6
	corr. to E_1	3	4	3.3
6.2885	corr. to E_2	3	4	2.7
	plateau	6	5	4.2
	corr. to E_1	4	5	5
6.0219	corr. to E_2	4	5	4
	corr. to E_1	4	5	4.5
	corr. to E_2	4	5	3.7

Table: Estimates of energy differences. Errors are roughly 1 or 2 on the last digit.

Static B-meson decay constant

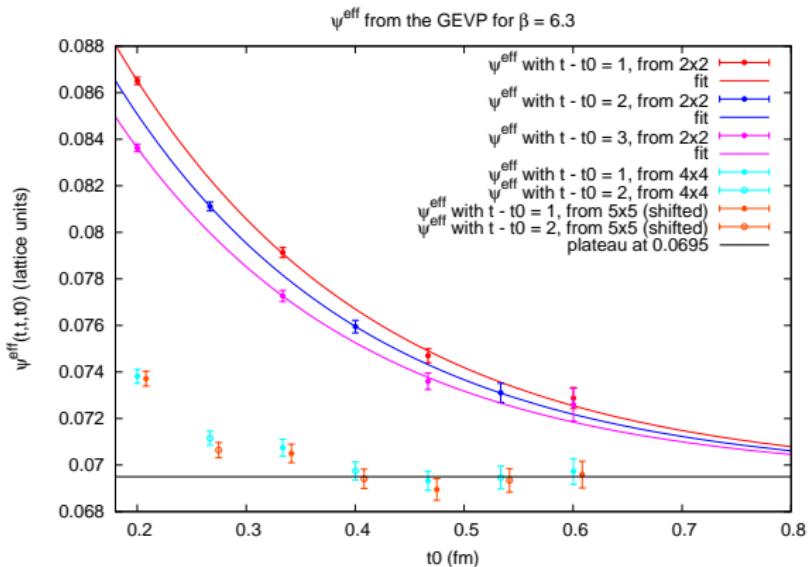
$$a = 0.07 \text{ fm}$$

$$a^{3/2} p_{01}^{\text{eff,stat}}(t, t_0)$$

(bare, unimproved)

curve:

$$F_{B_s}^{\text{stat}} + \alpha_N(t - t_0) e^{-\Delta E_{N+1,1} t_0}$$



$\Delta E_{N+1,1}$ agree with plateaux of $E_{N+1}^{\text{eff,stat}}(t, t_0)$ for large N' and t

systematical and statistical precision better than 1%

Conclusions

Corrections to $E_n^{\text{eff}}, p_{n0}^{\text{eff}}$ of GEVP understood

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- ▶ Matrix elements can be computed systematically, also excited state ME's
 - corrections $\exp(-(E_{N+1} - E_n) t_0)$
- ▶ everything is applicable also to HQET (and other EFT)

The end

thank you for your attention