Axial coupling and momentum fraction of the nucleon with twisted mass fermions

Rémi Baron CEA Saclay

17 July 2008, Lattice 2008



Collaboration

- C. Alexandrou, T. Korzec, G. Koutsou
- V. Drach, P. Harraud, Z. Liu, V. Morenas, J. Carbonell, M. Brinet, P. Guichon, O. Pene
- C. Urbach, D. Renner



3-point functions

Results



Long term motivation : compute the nucleon observables with twisted mass

- Form factors (we start with g_A)
- Parton distributions PDF (< x >)
- Generalized parton distributions GPD (nothing yet)

Intense experimental program at JLab and CERN

How to compute

Form of the matrix element we are interested in

$$egin{array}{rcl} R &=& \displaystyle{\sum_{ec{y},ec{z}} e^{iec{
ho}\cdotec{y}}e^{iec{
ho}^\prime\cdotec{z}}\left\langle 0|J_{\gamma}(y)O(z)ar{J}_{
ho}(0)|0
ight
angle} \end{array}$$

 $\propto \langle N(p,s)|O(z)|N(p',s') \rangle$ if J is a nucleon interpolating field



• Look for a plateau on t_z with $t_y \gg t_z \gg 0$

Generalized source

$$R = \sum_{ec y,ec z} e^{iec p\cdot ec y} e^{iec p^ec \cdot ec z} ig\langle 0|J_\gamma(y)O(z) ar J_
ho(0)|0ig
angle$$

• After Wick contraction R has the form :

$$R = B^{\dagger} \gamma_5 \Lambda S_u$$

- With B solution of Dirac equation DB = Σ_G where Σ_G is a combination of quarks propagators.
- Σ_G is called generalized or sequential source
- New inversions are necessary
- To create Σ_G
 - Fix time of sink (ie ty, 12 in our case)
 - Fix momentum at sink. Operator momentum inserted when doing contractions
 - Choose how to project states

Contractions



- Green arrows : 2 propagators used in sequential source
- Red part : generalized propagator
- Blue part : normal propagator

Just need to combine correctly Dirac, color, and space indices.

Quark disconnected diagrams



- Evaluation of disconnected diagram numerically demanding (all to all propagator)
- Consider only non singlet observables

Renormalization issues

- We will use the RI-MOM scheme to renormalize non-perturbatively the bare quantities
- Z₄₄ computed for some quark masses, but not yet extrapolated to chiral limit (Z.Liu, V.Morenas)
- Hypercubic artefacts reduced using arXiv :0705.3523 technique
- Ideally (no mixing) multiplicative renormalization using scheme ${\cal S}$ at scale μ

$$\mathcal{O}^{\mathcal{S}}(\mu) = \textit{Z}^{\mathcal{O}}_{\mathcal{S}}(\mu)\mathcal{O}_{\textit{bare}}$$

• Some renormalization constants have already been computed by other members of the collaboration $(Z_A = 0.76(1) \text{ from arXiv :0710.0975})$

- g_A is one of the first check point for nucleon structure calculation
- Well known experimentally from neutron beta decay $(g_A \simeq 1.269)$
- Can be extracted with $O(z) = A_{\mu}^{u-d} = \bar{u}\gamma_{\mu}\gamma_5 u \bar{d}\gamma_{\mu}\gamma_5 d$
- We will need to use chiral perturbation theory, to check for finite size effects, for cutoff effects

< X >

- < x > = momentum fraction carried by the quarks
- Twist-2 operators

$$\mathcal{O}^{\{\mu_1\cdots\mu_n\}} = (\frac{i}{2})^{n-1} G_{ff'} \bar{\psi}_f \gamma^{\{\mu_1} \stackrel{\leftrightarrow}{D}^{\mu_2} \cdots \stackrel{\leftrightarrow}{D}^{\mu_n\}} \psi_{f'} - traces,$$

 $\{\cdots\}$ means symmetrization on the Lorentz indices.

· In our calculation, we use the operator

$$\mathcal{O}_{44}(x) = \frac{1}{2}\overline{u}(x)[\gamma_4 \stackrel{\leftrightarrow}{D}_4 - \frac{1}{3}\sum_{k=1}^3 \gamma_k \stackrel{\leftrightarrow}{D}_k]u(x),$$

where $D_{\mu} = \frac{1}{2} (\nabla_{\mu} + \nabla^{*}_{\mu}).$

< X >

 Then the bare lowest moment of the quark distribution functions (x) is given by

$$\langle x \rangle = rac{1}{m_N^2} \langle N, \vec{0} | \mathcal{O}_{44} | N, \vec{0} \rangle = rac{1}{m_N} rac{C_{44}(t)}{C_N(t_y)} \quad (0 \ll t \ll t_y).$$

Here

$$egin{aligned} \mathcal{C}_{44}(t) &= \sum_{ec{y},ec{z}} \langle J(t_y,ec{y}) \mathcal{O}_{44}(t,ec{z}) ar{J}(0,0)
angle, \ \mathcal{C}_{N}(t_y) &= \sum_{ec{y}} \langle J(t_y,ec{y}) ar{J}(0,0)
angle, \end{aligned}$$

and J(x) is the interpolating field for the nucleon.



3-point functions

Results

Ensembles

 $\beta = 3.9, a = 0.0855(6)$ fm from f_{π} .

| $\pmb{a}\mu$ | $m_{\pi}({ m GeV})$ | $L^3	imes T$ | $m_{\pi}L$ | Nmeas |
|--------------|---------------------|-----------------|------------|-------|
| 0.0100 | 0.483(1) | $24^3 	imes 48$ | 5 | 173 |
| 0.0085 | 0.446(1) | $24^3 	imes 48$ | 4.7 | 161 |
| 0.0064 | 0.389(1) | $24^3 	imes 48$ | 4.1 | 131 |
| 0.0040 | 0.312(2) | $24^3 	imes 48$ | 3.2 | 392 |

APE and Gaussian smearings were used

Vector current

• Preliminary (10 configurations only)



• Close to the value found by the collaboration $(Z_V = 0.6104(2)(3))$

< X >



g_A



Summary

- First results for *g_A* and *< x >*
- Error bars will be strongly reduced in the near future
- Planned work
 - Higher moments of GPDs
 - Renormalization for more complicated operators
 - Partially twisted boundary conditions for low transfer
 - Disconnected diagrams