

“Ask not what the extra
dimension can do for you...”

(or the successes and challenges of AdS/QCD)

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Lattice2008

Overview

1. What is AdS/QCD?
2. Top-down vs. bottom-up paths.
3. Results/challenges facing the approach.

AdS/QCD

Old Myth: G.T. at large N simplifies  a weakly coupled string theory.

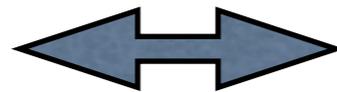
AdS/CFT: An additional simplification occurs for a large hierarchy of anomalous dimension of operators:

a few have $\Delta_O \sim 1$, but most have $\Delta_O \gg 1$.

AdS/CFT duality:

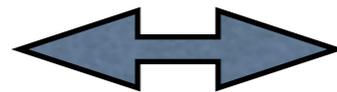
G.T. is dual to an extra-dimensional (ED) theory:

Local primary operators
of G.T., $O(x)$.



Fields propagating in ED
curved space $\phi(x, z)$.

Δ_O



m_ϕ

Holographic Dictionary

G.T.

$$\bar{q}q(x)$$

$$J_\mu(x) = \bar{q}\gamma_\mu q(x)$$

$$T_{\mu\nu}(x)$$

$$\bar{q}\gamma_{\mu_1} D_{\mu_2} \cdots D_{\mu_n} q(x)$$

ED dual

$$\phi(x, z)$$

$$A_\mu(x, z)$$

$$h_{\mu\nu}(x, z)$$

$$\phi_{\mu_1\mu_2\dots\mu_n}(x, z)$$

$$\Delta_0 \gg 1 \quad \longrightarrow \quad m_\phi \sim \Delta_0 \gg 1$$

Most ED dual fields decouple. Only left with duals to currents:

$$J_\mu, T_{\mu\nu} \text{ (+ superpartners)} \quad \longleftrightarrow \quad A_\mu, h_{\mu\nu} \text{ (+ super)}$$

(The string theory simplifies, becoming a local field theory)

$$g_{ED} \sim \frac{1}{\sqrt{N_c}}$$

At first: ED duals only for conformal theories were

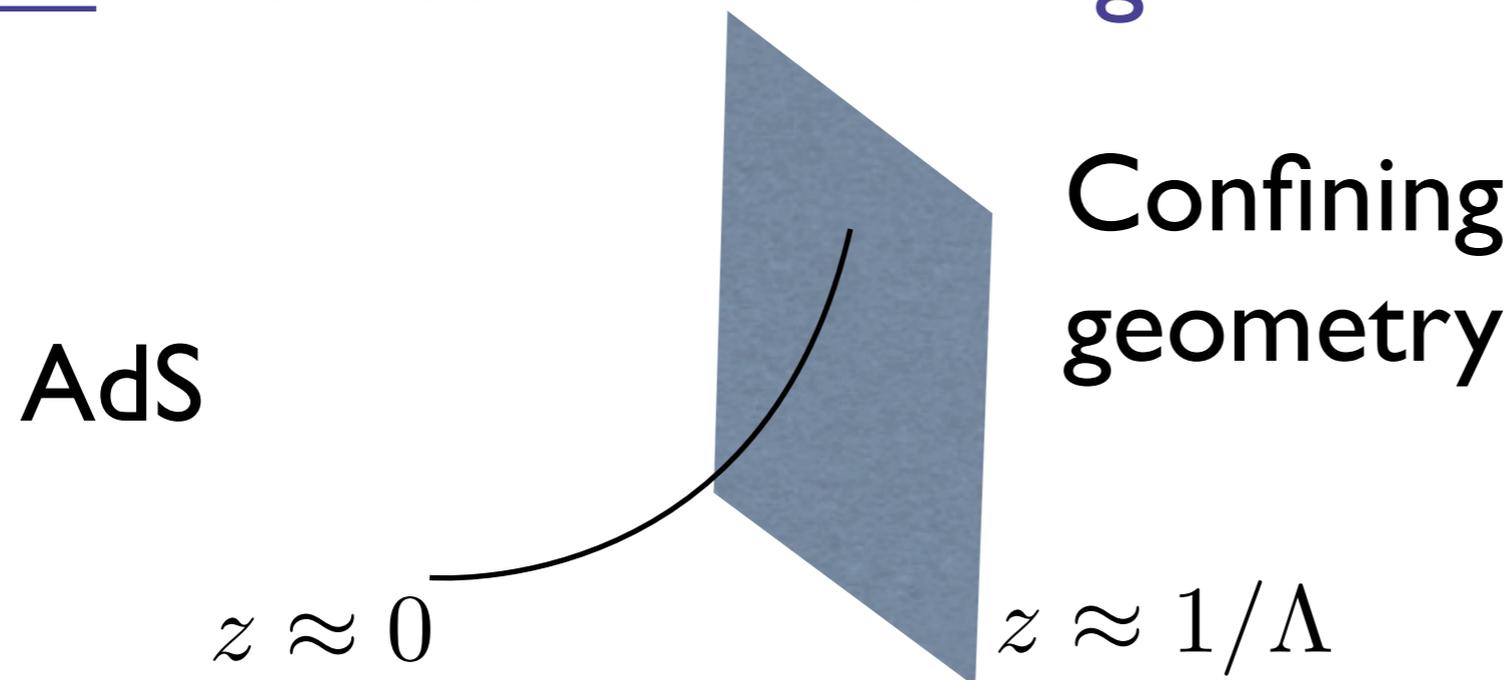
found:

$$ds^2 = \frac{1}{z^2} (dx_\mu dx^\mu - dz^2)$$

Dilatation: $x^\mu \rightarrow \lambda x^\mu, z \rightarrow \lambda z$

UV: $z \ll 1,$ IR: $z \gg 1$

Then: Non conformal backgrounds found.



Towards QCD

Top-down approach: Study of particular field theories (mostly SUSY) with known ED duals in an attempt to extract universal/qualitative features.

Bottom-up approach: An attempt to build ED models for certain limited sectors of QCD that make quantitative predictions.

Top-down approach

- Studies of χSB , spectra, DIS, etc in theories with large hierarchies in anomalous dimensions of operators.
- QFT dynamics at finite temperature and chemical potential: Phase diagrams, transport properties, and response to external E & B fields:
 1. RHIC physics: small shear viscosity consistent with strong coupling (qualitatively similar to strongly coupled known duals).
 2. Condensed matter physics: theories which reflect some properties of superconductors.

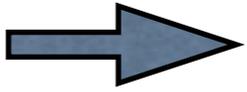
- Challenges for top-down:
 1. Calculations are under control at large N .
 2. All field theories described have a large hierarchy in the dimensions of operators.
 3. It is not yet clear how to connect the calculations to QCD data at a more quantitative level.

FT with simple ED duals are
Probably rare

1. Generically, a strongly coupled CFT has $\Delta_O \sim 1$.
2. Even for a CFT with a marginal deformation, there is no guarantee that $\Delta_O \sim (\text{def. parameter})$.
3. FT with weakly coupled duals have correlation functions with certain polarization structures suppressed (not due to any symmetry).

QCD as a test for ED models

w/ Erlich, Schwartz, Son, and Stephanov.

QCD: $\Delta_O \sim O(1)$ (no hierarchy). 

ED dual contains infinitely many fields.

Conjecture: Lightest states are well described by the lowest dimension operators. 

Might only need to consider a few ED fields.

Evidence for the conjecture:

1. True for G.T. with known ED duals.
2. Suggested by success of ED models for lightest QCD resonances.
3. True in 2D QCD at large N .
4. LATTICE?

Modeling the lightest QCD states

1. Include only lowest dimension operators.
2. Use QCD OPE's as a guide for constructing the 5D metric.
3. Model confinement with a hard wall + b.c.

$$SU(3)_L \times SU(3)_R \rightarrow SU(3)_V$$

QCD

$$\bar{q}q(x)$$

$$J_\mu^V(x), J_\mu^A(x)$$

$$T_{\mu\nu}(x)$$

$$\bar{q}\gamma_{\{\mu_1}D_{\mu_2}\dots D_{\mu_n}\}q(x), G_{\nu\{\mu_1}^aD_{\mu_2}\dots D_{\mu_{n-1}}G_{\mu_n}^{a\nu}\}(x)$$

5D dual

$$\phi(x, z)$$

$$V_\mu(x, z), A_\mu(x, z)$$

$$h_{\mu\nu}(x, z)$$

$$\phi_{\{\mu_1\mu_2\dots\mu_n\}}(x, z)$$

$$\int_x e^{iqx} \langle J_\mu^a(x) J_\nu^b(0) \rangle = \delta^{ab} (q_\mu q_\nu - q^2 g_{\mu\nu}) \Pi(Q^2)$$

$$\Pi_V(Q^2) = -\frac{1}{8\pi^2} \left(1 + \frac{\alpha_s}{\pi}\right) \ln Q^2 + \frac{1}{24\pi} \cdot \frac{\langle \alpha_s G^2 \rangle}{Q^4} - \frac{m_q \langle q\bar{q} \rangle}{Q^4} - \frac{14}{9} \pi \frac{\langle \alpha_s \bar{q}q \rangle^2}{Q^6} + \dots$$

$$S = \int d^5x \left\{ -\frac{1}{4g_5^2 z} \left(1 + \frac{c_0}{\text{Log}(\Lambda z)} + c_{4V}^{ab} z^4 + c_{6V}^{ab} z^6\right) F_V^a F_V^b \right. \\ \left. - \frac{1}{4g_5^2 z} \left(1 + \frac{c_0}{\text{Log}(\Lambda z)} + c_{4A}^{ab} z^4 + c_{6A}^{ab} z^6\right) F_A^a F_A^b + \frac{v(z)^{ab}}{2z^3} D\pi^a D\pi^b \right\}.$$

$$v(z)_{ab} = (M_q)_{ab} z + \langle \bar{q}q \rangle \delta_{ab} z^3$$

LATTICE - Help find condensates at large N.

QCD-lite Model

Also
Da Rold & Pomarol

$$c_{nV}, c_{nA} = 0.$$

$$z_{IR} = 1/(346\text{MeV}), \langle \bar{q}q \rangle = (308\text{MeV})^3, m_q = 2.3\text{MeV}, m_s = 35\text{MeV}.$$

Observable	Measured (MeV)	Model (MeV)	Width (MeV)
m_π	139.6	141	-
m_ρ	775.8	832	146
m_{a_1}	1230±40	1220	250-400
f_π	92.4	84	-
$F_\rho^{1/2}$	345±8	353	-
$F_{a_1}^{1/2}$	433±13	440	-
m_{K^*}	892	897	51
m_ϕ	1020	994	4
m_{K_1}	1272±7	1290	90±2
m_K	498	411	-
f_k	113	117	-
m_{f_2}	1275	1236	185
m_{ω_3}	1667±4	1656	168±10
m_{f_4}	2025±8	2058	225±18
m_{ρ_5}	2330±35	2448	400±100
m_{f_6}	2465±50	2829	255±40
m_η	548	520	-
$m_{\eta'}$	958	867	-
$\Gamma(f_2 \rightarrow 2\gamma)$	2.6±.24 KeV	2.71 KeV	-

Additional Progress

1. Baryons as 5D solitons (similar to skyrmions).
Pomarol & Wulzer
2. Anomaly physics from 5D Chern-Simmons term.
3. Modification of IR geometry away from the hard wall to reproduce Regge physics: $M_J^2 \sim J$.

2D QCD at large N (w/ T. Okui)

The PDF of the n-th excited meson is given by

$$f_q^n(x) = |\phi_n(x)|^2$$

Solutions for highly excited modes are :

$$\phi_n(x) \simeq \sqrt{2} \cos[n\pi x] \quad , \quad m_n^2 \simeq \pi^2 \Lambda^2 n .$$

In terms of the primary operators, O_k , $\Delta_k \sim k$:

$$\phi_n(x) = \sum_k \langle 0 | O_k | n \rangle P_{k-1}(2x-1)$$

The lightest states have the biggest overlap with the lowest dimension operators:

$$\langle 0 | O_k | n \rangle \ll 1, \quad k \gtrsim n$$

Parton- x and the extra dimensional coordinate are related by a transform:

't Hooft Eqn  EOM for a scalar in ED.

Conclusions

1. AdS/QCD captures non-perturbative aspects of QCD both quantitatively and qualitatively.
2. Bottom-up models offer a connection between low-energy parameters (of lightest resonances) and the UV theory (via the OPE).
3. These models are a useful tool for the LHC, where we might access only a few resonances. Similar to Randall-Sundrum models.
4. They work better than they should.
5. LATTICE may help clarify the reach and limitations of the approach.