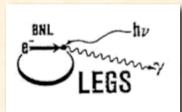


# RESULTS FROM POLARIZED EXPERIMENTS FROM LEGS AND GRAAL



#### Annalisa D'ANGELO

University of Rome "Tor Vergata" and INFN Rome Tor Vergata



## Outline

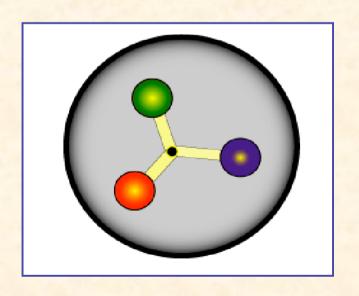
- Introduction: missing resonances and hadronic degrees of freedom
- The Legs and Graal Experimental set-ups.
- Results for  $\Sigma$  beam polarization asymmetries at Graal:

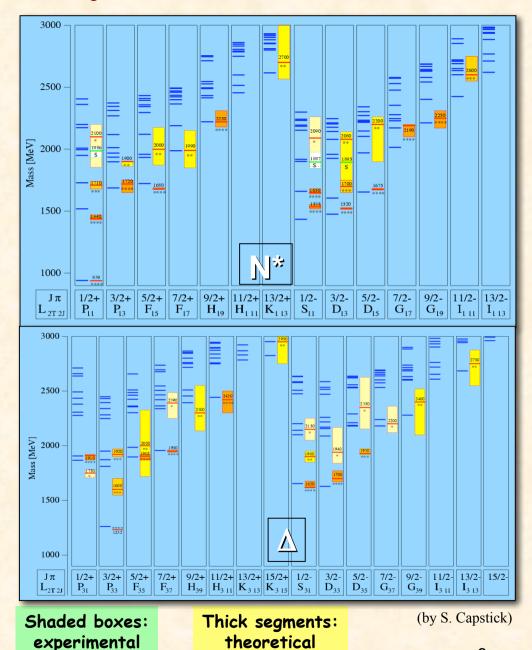
$$\begin{cases} \vec{\gamma} + n \to \pi^0 + n \\ \vec{\gamma} + n \to \pi^- + p \\ \vec{\gamma} + N \to \omega + N \end{cases}$$

- Results for  $O_x$  and  $O_y$  double polarization asymmetries for  $k^+\Lambda$  photoproduction on the proton at Graal.
- Results on E and G double polarization asymmetries at Legs:  $\begin{cases} \vec{\gamma} + \vec{H}D \rightarrow \pi^0 p \\ \vec{\gamma} + \vec{H}D \rightarrow \pi^+ n \end{cases}$
- Conclusions

#### QCD-inspired Constituent Quark Models

- Chiral symmetry breaking of the QCD Lagrangian generates Constituent Q with effective masses confirmed by LQCD and DSE calculations.
- •Asymmetry of the baryon wave function is guaranteed by color, but color degrees of freedom are integrated out and play no dynamical role.
- States classified by isospin, parity and spin within each oscillator band.





predictions

results

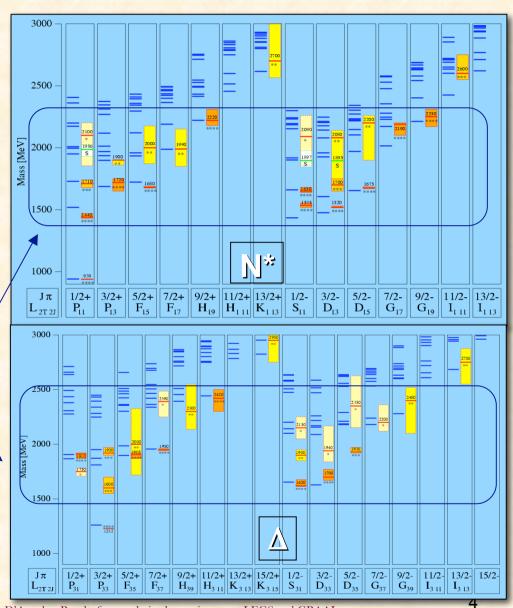
Annalisa ]

ized ex

3

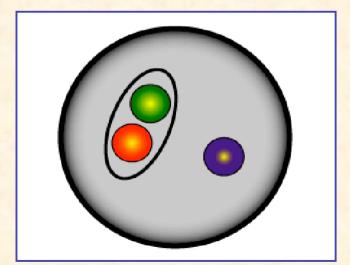
#### QCD-inspired Constituent Quark Models

- Chiral symmetry breaking of the QCD Lagrangian generates Constituent Q with effective masses confirmed by LQCD and DSE calculations.
- •Asymmetry of the baryon wave function is guaranteed by color, but color degrees of freedom are integrated out and play no dynamical role.
- States classified by isospin, parity and spin within each oscillator band.
- only lowest few in each band seen (in  $\pi N$ ) with  $4 \bigstar$  or  $3 \bigstar$  status
- $g(\pi N)$  couplings predicted to decrease rapidly with mass in each oscillator band
- higher levels predicted to have larger couplings to  $K\Lambda$ ,  $K\Sigma$ ,  $\pi\pi N$ , ...



#### QCD-inspired di- Quark Models

- 2 quarks in nucleon assumed to be quasibound in a color isotriplet; diquark-quark is a net color isosinglet.
- all possible internal di-quark excitations ⇔ full spectrum of CQM
- internal di-quark excitations are frozen out (spin 0; isospin 0)  $\Leftrightarrow$  large reduction in the number of degrees of freedom  $\Leftrightarrow$  predicts less N\* states than seen in  $\pi$ N



	N*	Status	$SU(6) \otimes U(3)$	Parity	Δ*	Status	$SU(6) \otimes U(3)$	Parity
Ī	P <sub>13</sub> (938)	****	(56, 0+)	+	$P_{33}(1232)$	****	(56,0+)	+
Ì	S <sub>11</sub> (1535)	****	$(70, 1^{-})$	-	S <sub>31</sub> (1620)	****	$(70, 1^{-})$	-
	$S_{11}(1650)$	****	$(70, 1^{-})$	-	$D_{13}(1700)$	***	$(70, 1^{-})$	-
	$D_{13}(1520)$	****	$(70, 1^{-})$	-				
	$D_{13}(1700)$	***	$(70, 1^{-})$	-				
	$D_{15}(1675)$	****	$(70, 1^{-})$	-				
	$P_{11}(1520)$	****	$(56,0^+)$	+				
					$P_{31}(1875)$	****	$(56, 2^+)$	+
	$P_{11}(1710)$	***	$(70,0^+)$	+	$P_{31}(1835)$		$(70,0^+)$	+
	$P_{11}(1880)$		$(70, 2^+)$	+				
	$P_{11}(1975)$		$(20, 1^+)$	+				
					$P_{33}(1600)$	***	$(56,0^+)$	+
	$P_{13}(1720)$	****	$(56, 2^+)$	+	$P_{33}(1920)$	***	$(56, 2^+)$	+
	$P_{13}(1870)$	*	$(70,0^+)$	+				
	<i>P</i> <sub>13</sub> (1910)		$(70, 2^+)$	+	$P_{33}(1985)$		$(70, 2^+)$	+
	$P_{13}(1950)$		$(70, 2^+)$	+				
	$P_{13}(2030)$		$(20, 1^+)$	+				
	$F_{15}(1680)$	****	$(56, 2^+)$	+	$F_{35}(1905)$	****	$(56, 2^+)$	+
	$F_{15}(2000)$	**	$(70, 2^+)$	+	$F_{35}(2000)$	**	$(70, 2^+)$	+
	$F_{15}(1995)$		$(70, 2^+)$	+				
	$F_{17}(1990)$	**	$(70, 2^+)$	+	$F_{37}(1950)$	****	(56, 2+)	+

the challenge:  $\Leftrightarrow$  unravel the N\* spectrum

### Experimental Requirements

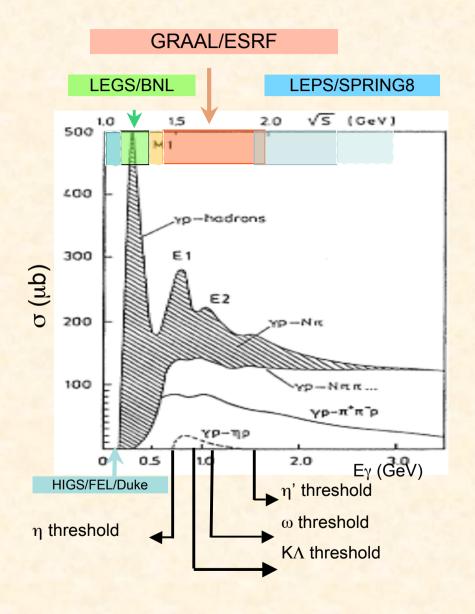
- ☐ Tagged and polarized photon beam
- ☐ Large acceptance detector
- H and D polarized targets

Both Legs and Graal experiment were constructed to meet all above requirements

in the energy ranges:

$$E_{\gamma} = (180 - 450) \text{MeV}$$
 and  $E_{\gamma} = (500 - 1500) \text{MeV}$ 

#### Polarized photon beams: Compton backscattering



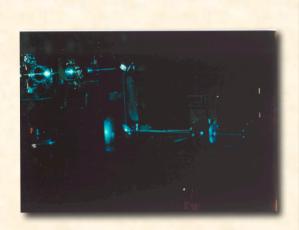
- •Hiys  $\rightarrow$  below  $\pi$  threshold
- ·Legs  $\rightarrow \Delta_{33}(1232)$  resonance region

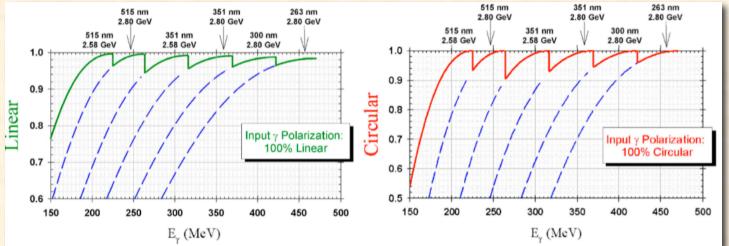
•Graal  $\rightarrow$  E $_{\gamma}$  = .6-1.5 GeV / W=1.4-1.9 GeV Region of the second and third baryon resonances  $\eta$ , K,  $\omega$ , thresholds

•Leps  $\rightarrow$  E $_{\gamma}$  = 1.5-2.5 GeV  $\eta' \phi$  thresholds

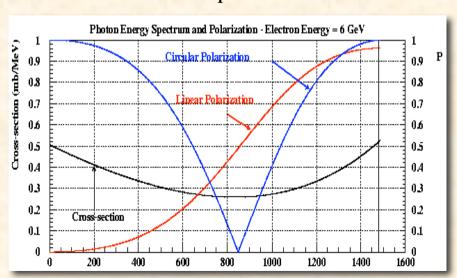
#### Polarized photon beams: Compton backscattering and Bremsstrahlung

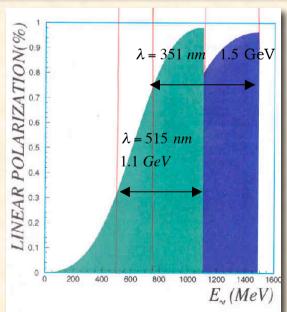
#### LEGS beam polarization



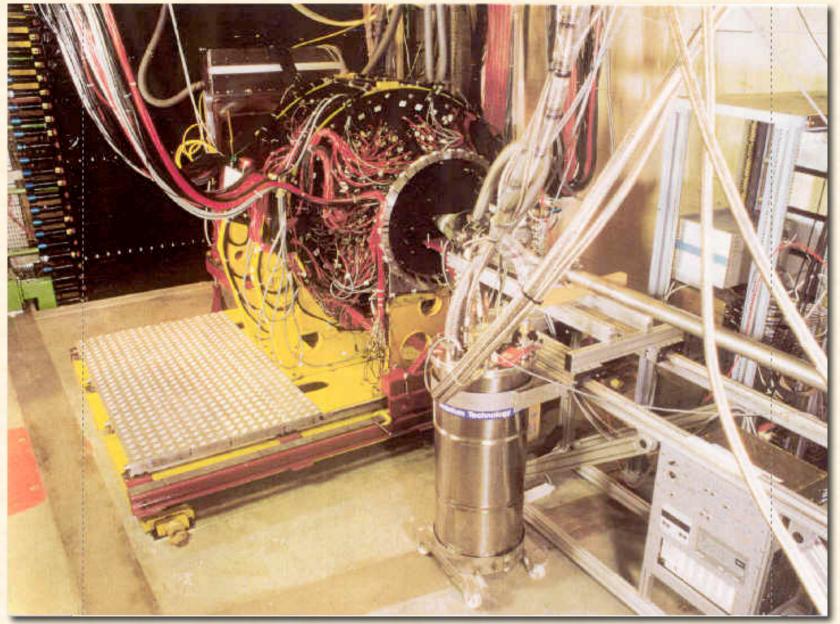


#### GRAAL beam polarization





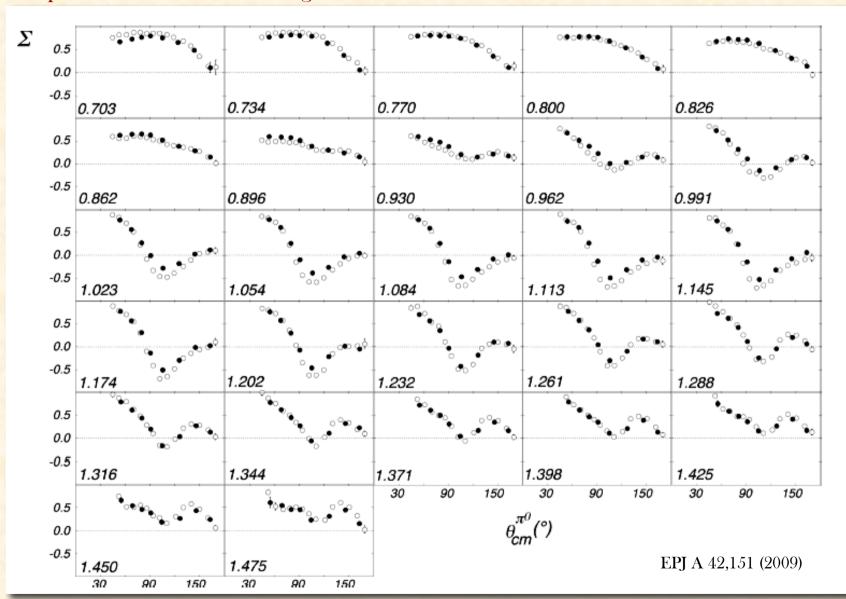
#### The Graal detector: Lagranye



Large Acceptance Graal Apparatus for Nuclear y Experiments

## Σ measurements at Graal on proton and deuteron targets

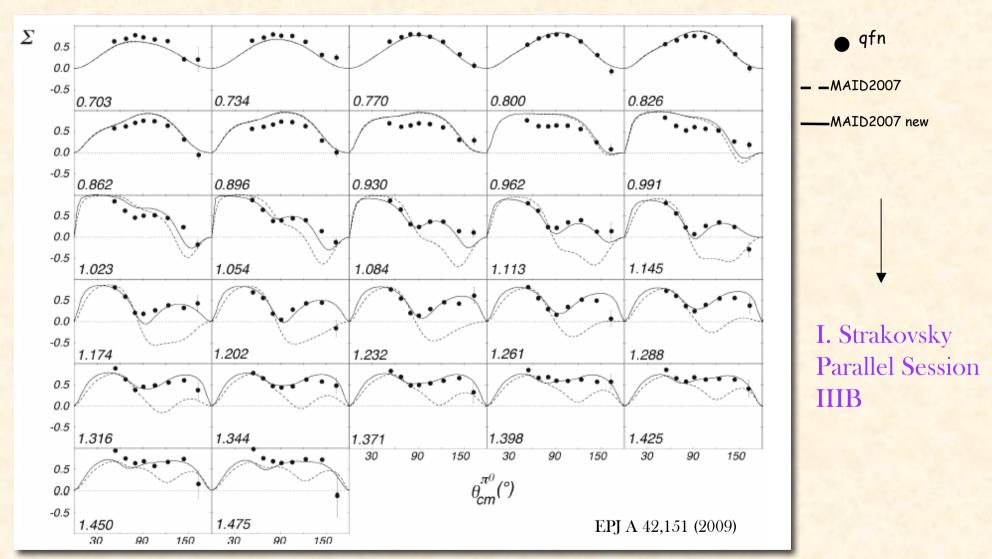
$$\bigcirc \quad \vec{\gamma} + p \rightarrow \pi^0 + p \qquad \bullet \quad \vec{\gamma} + p(+n) \rightarrow \pi^0 + p(+n)$$



Very nice agreement between free and quasi-free results on the proton

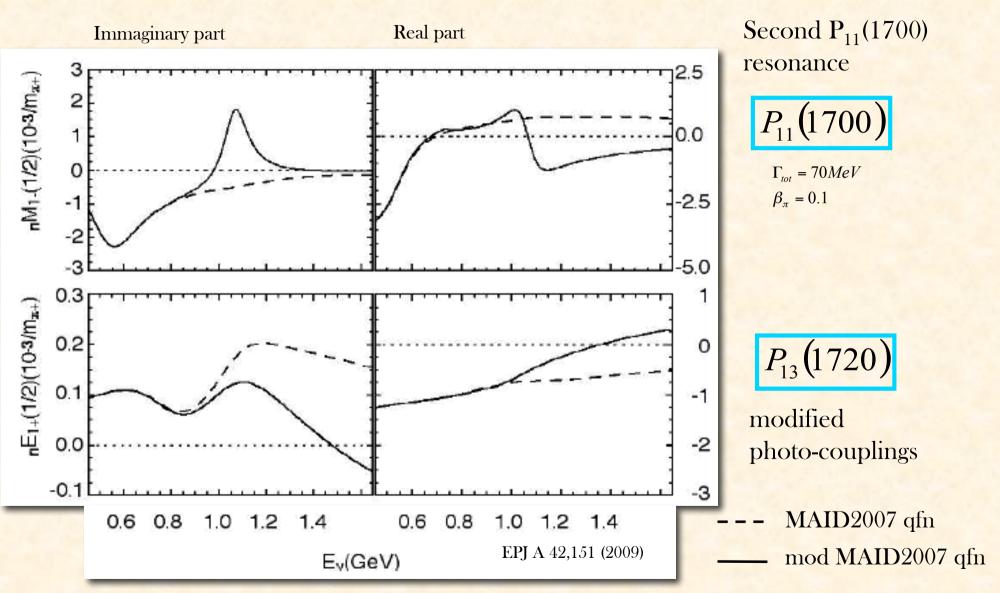
# Σ measurements at GRAAL deuteron target

$$\overrightarrow{\gamma} + n(+p) \rightarrow \pi^0 + n + (p)$$

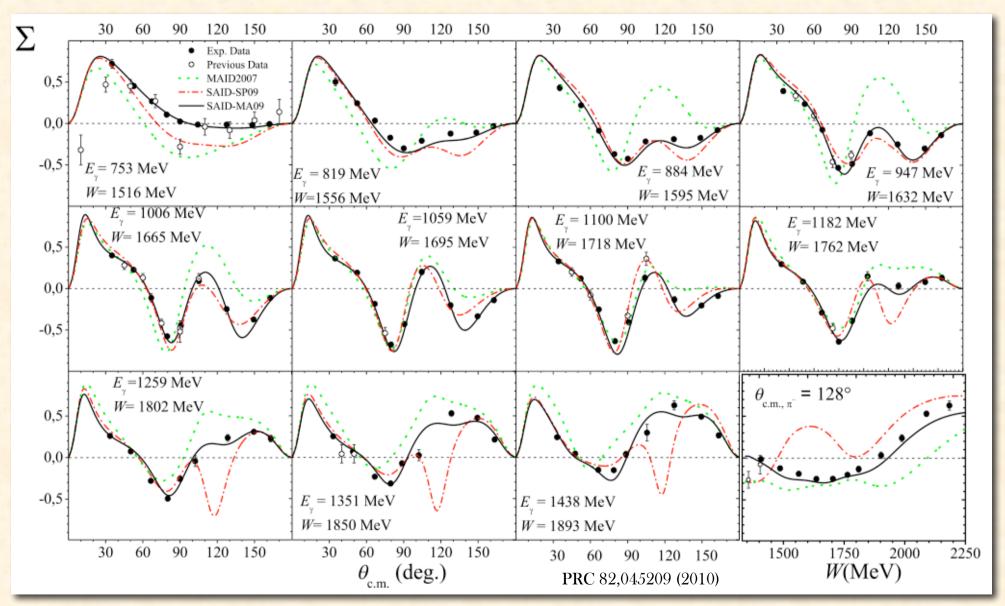


We may assume that results from quasi-free neutrons may represent the free neutron response → final state interactions and re-scattering are negligible)

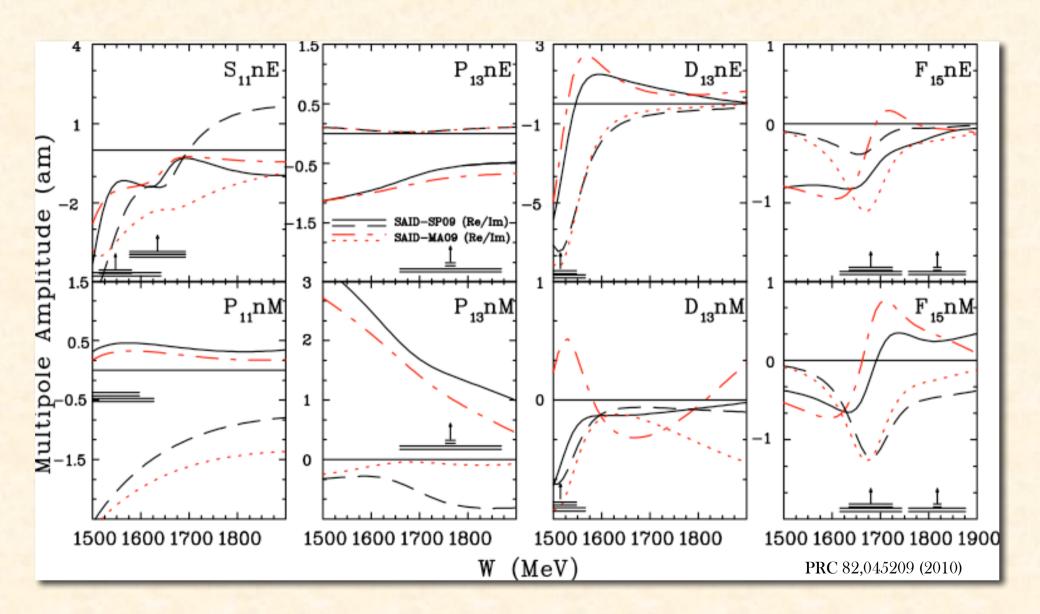
# $\Sigma$ for $\pi^0$ photoproduction on qfn Multipole extraction in MAID2007



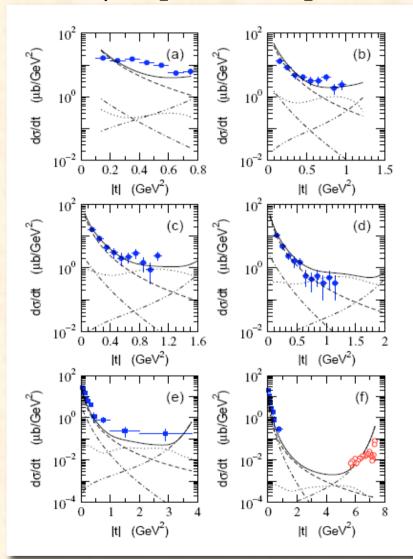
#### $\Sigma$ results on $\vec{\gamma} + n (+p) \rightarrow \pi^- + p (+p)$ at GRAAL



#### Multipole modifications due to $\Sigma$ results on $\vec{\gamma} + n \ (+p) \rightarrow \pi^- + p \ (+p)$ at GRAAL



### $\vec{\gamma} + p \rightarrow \omega + p$ : Differential Cross-Section



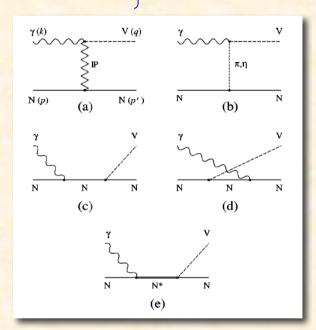
Oh,Titov,Lee PRC63 (2001) 025201

Low *t* diffractive behavior:

Vector Dominance Model (1960), J.J.Sakurai

- → Pomeron exchange
- $\rightarrow \pi^0/\eta$  exchange

t-channel



Eγ=(a) 1.23GeV

- (b) 1.45GeV
- (c) 1.68GeV
- (d) 1.92GeV
- (e) 2.80GeV
- (f) 4.70 GeV

Large t behavior : s- and u-channel contributions  $\rightarrow$  intermediate resonant states (N\*).

----

pseudo-scalar meson exchange

----

Pomeron exchange

\_..\_.

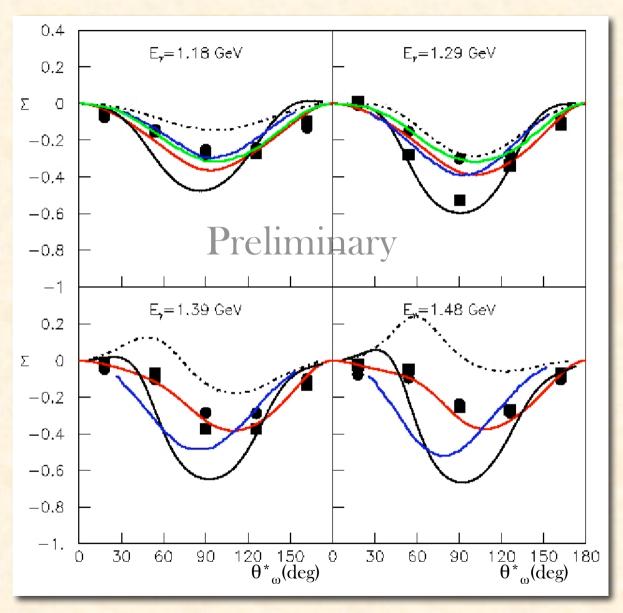
direct and crossed nucleon terms

.....

N\* excitation

### $\Sigma$ results on $\vec{\gamma} + p \rightarrow \omega + p$ at GRAAL:

$$\omega \to \pi^0 \gamma$$
 and  $\omega \to \pi^+ \pi^- \pi^0$ 



O. Zhao s and u-channel including  $P_{13}(1720)$ PRC63(2001)025203

Bonn-Gatchina dominant  $P_{13}(1720)$ Eur. Phys.J.A 25(2005)427

Giessen model PRC71(2005)055206

Oh, Titov and Lee PRC66 (2002)015204

M. Paris PRC79 (2009) 025208

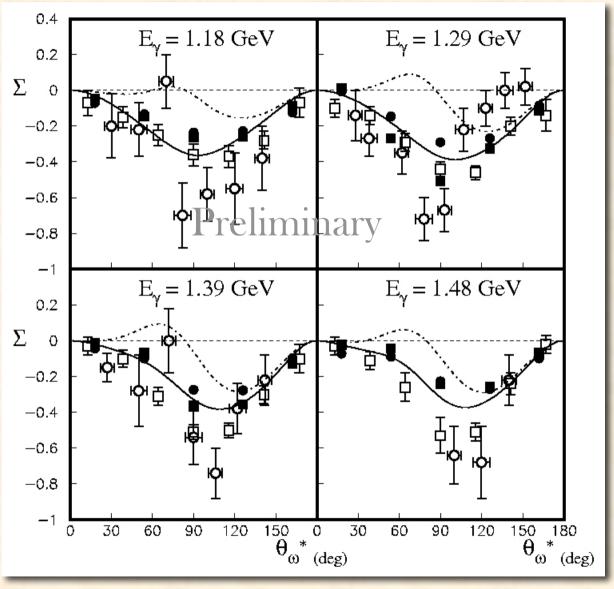
• 
$$\omega \rightarrow \pi^0 \gamma$$

$$\bullet \quad \omega \to \pi^0 \gamma$$

$$\bullet \quad \omega \to \pi^+ \pi^- \pi^0$$

$$\Sigma$$
 results on  $\vec{\gamma} + p \rightarrow \omega + p$  at GRAAL:

$$\omega \to \pi^0 \gamma$$
 and  $\omega \to \pi^+ \pi^- \pi^0$ 



• Graal 
$$\omega \to \pi^0 \gamma$$

Graal 
$$\omega \rightarrow \pi^+ \pi^- \pi^0$$

$$\circ$$
 Bonn  $\omega \to \pi^0 \gamma$ 

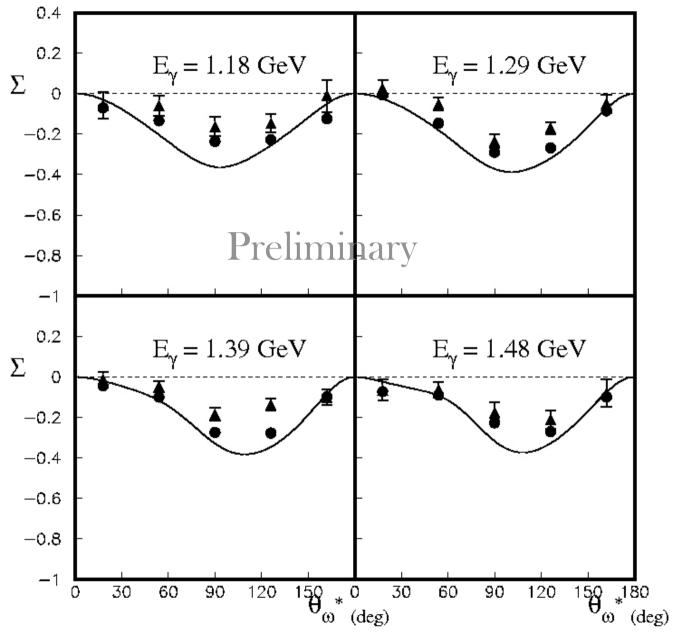
$$_{\sqcap}$$
 PRL96(06)  $\omega 
ightharpoonup \pi^{+}\pi^{-}\pi^{0}$ 

Zhao model

s and u-channel
including P<sub>13</sub>(1720)

s and u-channel no P<sub>13</sub>(1720)

### $\Sigma$ results on $\vec{\gamma} + p \rightarrow \omega + p$ and $\vec{\gamma} + p (+n) \rightarrow \omega + p (+n)$ at GRAAL

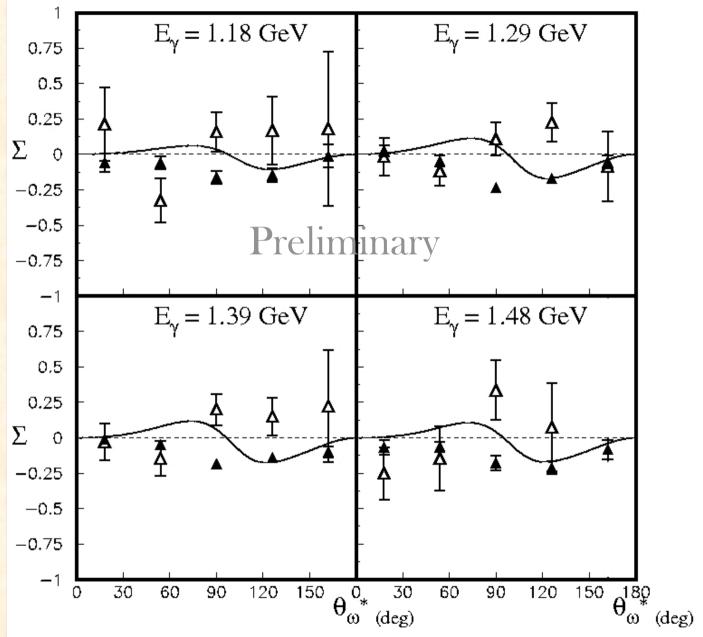


Zhao model
\_\_\_\_\_s and u-channel
including P<sub>13</sub>(1720)

- $\omega \rightarrow \pi^0 \gamma$  free-proton
- $\omega \to \pi^0 \gamma$ A Quasi-free-proton

18

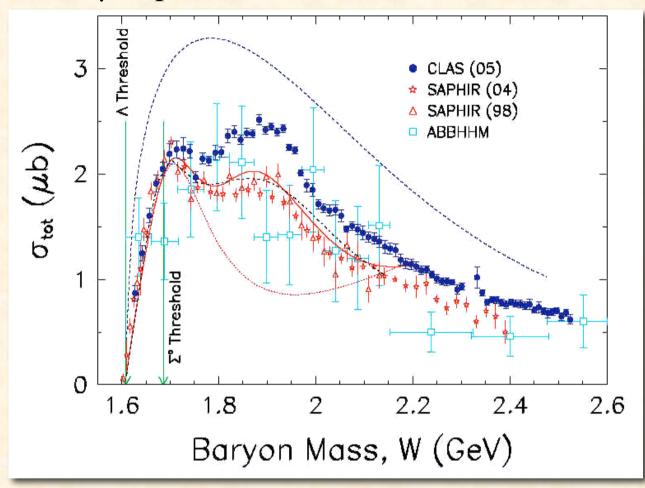
### $\Sigma$ results on $\vec{\gamma} + n (+p) \rightarrow \omega + n (+p)$ at GRAAL



\_\_\_\_ Zhao model

 $\omega \to \pi^0 \gamma$   $\Delta \text{ quasi-free- neutron}$ 

### $\vec{\gamma} + p \rightarrow k^+ + \Lambda$ : Total Cross-Section



Cross section data show a structure at W=1900 MeV.

Coupled-channel analysis finds that  $S_{11}(1650)$ ,  $P_{11}(1710)$  and  $P_{13}(1720)$  have the most significant decay widths in the  $k+\Lambda$  channel.

Isobar model requires the inclusion of a "missing"  $D_{13}(1895)$  resonance to reproduce the cross section data.

**– – – – Regge model calculation** 

...... KAON-Maid without D<sub>13</sub>(1895)

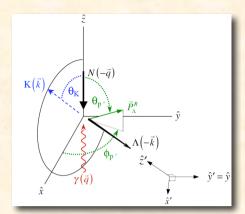
KAON-Maid with  $D_{13}$ (1895)

— · — · - Saclay dynamical coupled channel

 $S_{11}$  (1800) and  $P_{13}$  (1900) also seem to play a role

#### Polarization observables in

$$\vec{\gamma} + p \rightarrow k^+ + \vec{\Lambda}$$

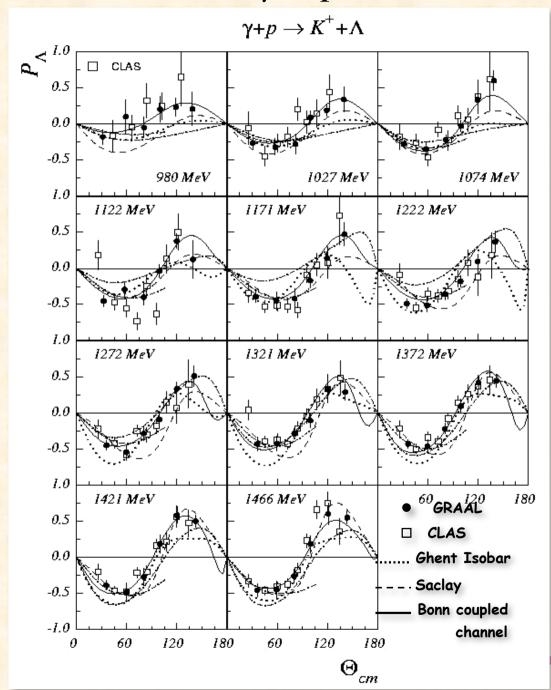


#### Weak Λ decay is self-analyzing

	Photon bean	Target			Recoil			
					000 - 700 - 7000 - 7000 - 7000 - 7000 - 7	<i>x'</i>	<i>y</i> '	<b>z</b> '
			X	У	Z	Amp Camp Camp Camp Camp Camp Camp Camp Ca		
	unpolarized	$\sigma_0$	2   168   168   168   168   168   166	$egin{array}{cccccccccccccccccccccccccccccccccccc$	the land and had been land as		$oldsymbol{P}$	
	linearly P <sub>γ</sub>	Σ	H	- <b>P</b>	-G	$O_x$ ,	<b>-T</b>	$O_z$ ,
	circular P <sub>γ</sub>	THE VIEW VIEW VIEW VIEW VIEW VIEW VIEW VIE	F		- <b>E</b>	-C <sub>x</sub> ,		$C_z$ ,

## PA in: $\vec{\gamma} + p \rightarrow k^+ + \vec{\Lambda}$

at Graal



A.Lleres et al., EPJ A 31, 79-93 (2007)

$$W(\cos \theta_p) = \frac{1}{2} \left( 1 + \alpha |\vec{P}_{\Lambda}| \cos \theta_p \right)$$

$$P_{\Lambda} = \frac{2}{\alpha} \frac{N_{(\cos \theta_p > 0)} - N_{(\cos \theta_p < 0)}}{N_{(\cos \theta_p > 0)} + N_{(\cos \theta_p < 0)}}$$

$$\alpha = 0.642 \pm 0.013$$

#### From $\Sigma$ and P measurements:

· Saclay Model:

$$S_{11}(1700) P_{13}(1800) D_{13}(1850)$$

· Ghent Isobar Model:

$$D_{13}(1900)$$

· Reggeized Model:

$$P_{13}(1900) D_{13}(1900)$$

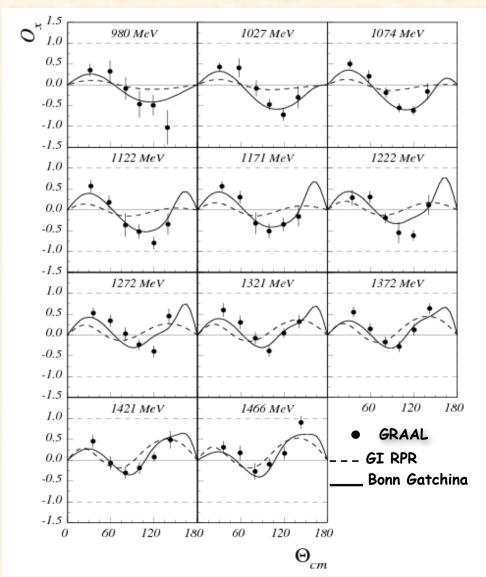
· Bonn Coupled Channel Model:

$$D_{13}(1875)$$

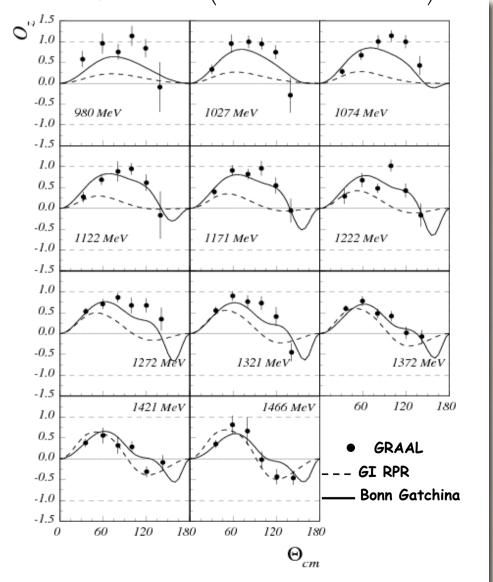
#### Double Polarization Observables in K+Λ Photoproduction

A.Lleres et al., EPJ A 39, 149-161 (2009)

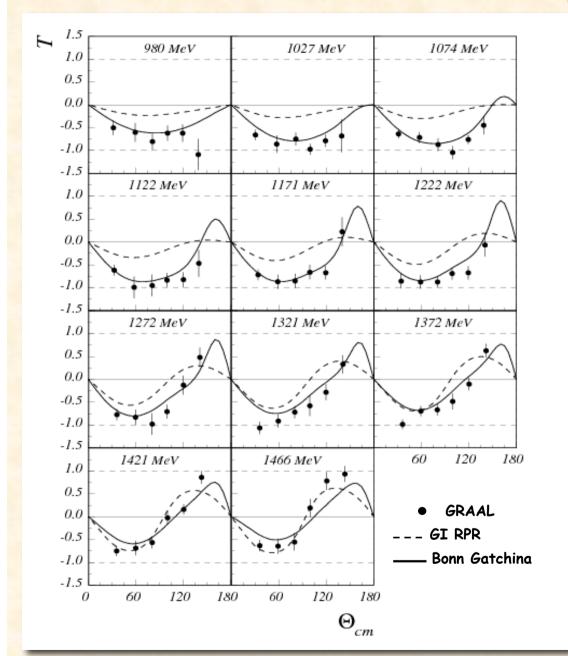
$$\frac{2N_{+}^{x'}}{N_{+}^{x'} + N_{-}^{x'}} = \left(1 + \alpha \frac{2P_{\gamma}O_{x}}{\pi} \cos \theta_{p}^{x'}\right)$$



$$\frac{2N_{+}^{z'}}{N_{+}^{z'} + N_{-}^{z'}} = \left(1 + \alpha \frac{2P_{\gamma}O_{z}}{\pi} \cos \theta_{p}^{z'}\right)$$



#### T in K<sup>+</sup>Λ Photoproduction



A.Lleres et al., EPJ A 39, 149-161 (2009)

$$\frac{2N_{+}^{y'}}{N_{+}^{y'} + N_{-}^{y'}} = \left(1 + \frac{2P_{\gamma}\Sigma}{\pi}\right) \left(\frac{1 + \alpha \frac{P\pi + 2P_{\gamma}T}{\pi + 2P_{\gamma}\Sigma} \cos\theta_{p}^{y'}}{1 + \alpha P \cos\theta_{p}^{y'}}\right)$$

From  $O_x$ ,  $O_z$  and T results:

• Ghent Isobar RPR Model:

$$S_{11}(1650)$$
  $P_{11}(1710)$   $P_{13}(1720)$ 

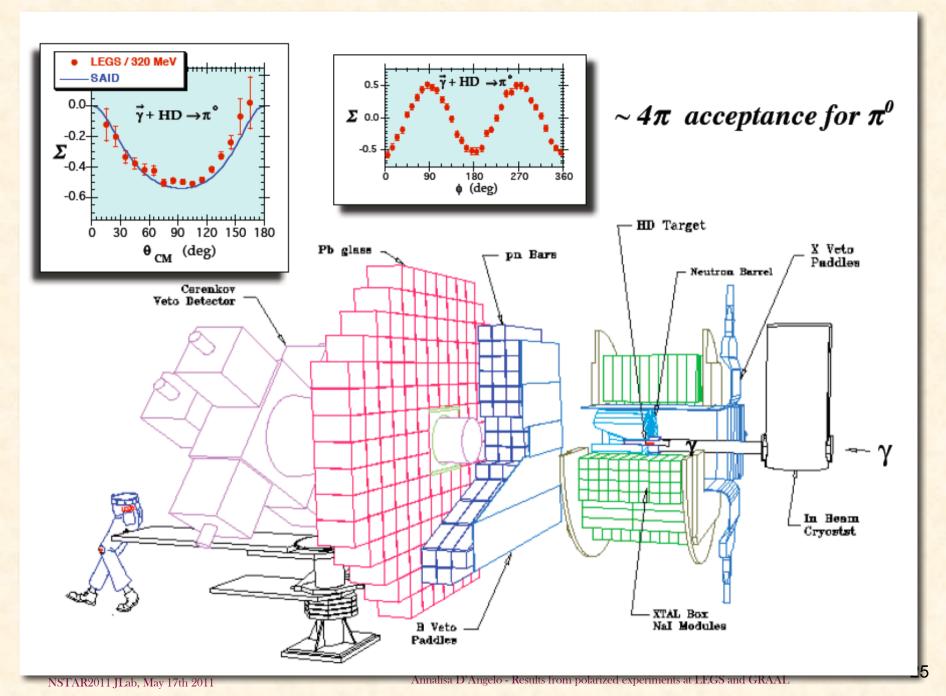
$$P_{13}(1900)$$
  $D_{13}(1900)$ 

• Bonn Gatchina Model:

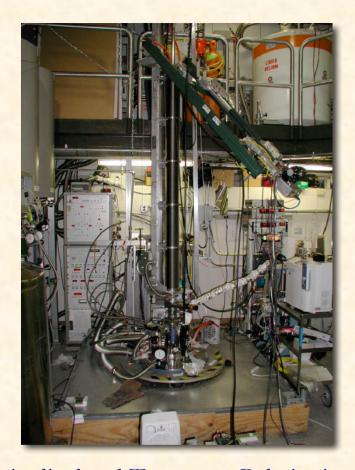
$$S_{11}(1535)S_{11}(1650)P_{13}(1720)P_{11}(1840)$$

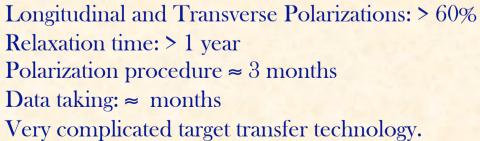
$$P_{13}(1900)$$

#### LEGS Spin ASYmmetry Array (SASY)

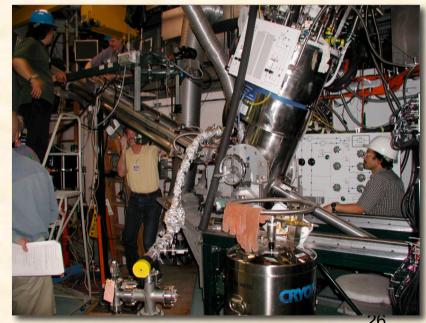


#### Polarized targets: frozen spin HD target at LEGS









#### Polarized targets: frozen spin HD target at LEGS

Very clean signal/background separation

	TARGET			
PHOTON BEAM				
		×	У	z
unpolarized	$\sigma_{\scriptscriptstyle 0}$		Т	
linearly P,	Σ	I	-P	-G
circular P <sub>y</sub>		F		-E

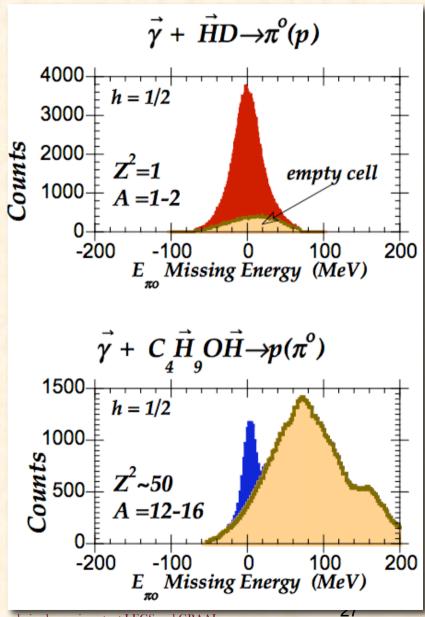
Longitudinal and Transverse Polarizations: > 60%

Relaxation time: > 1 year

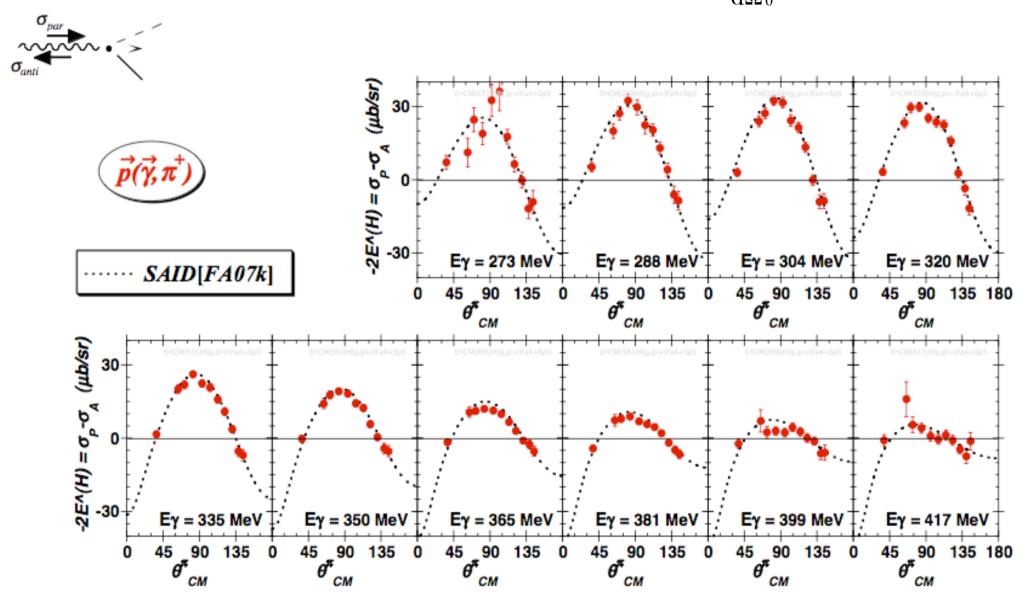
Polarization procedure ≈ 3 months

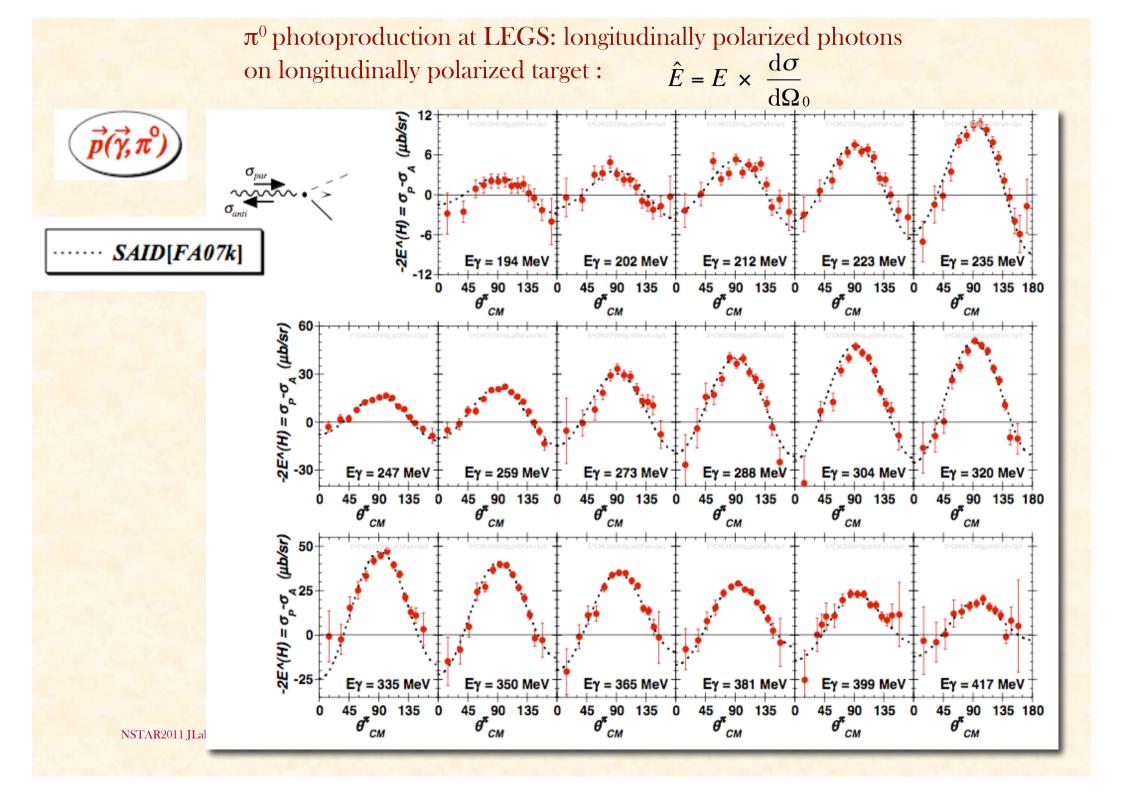
Data taking: ≈ months

Very complicated target transfer technology.



 $\pi^{\theta}$  photoproduction at LEGS: longitudinally polarized photons on longitudinally polarized target :  $\hat{E} = E \times \frac{\mathrm{d}\sigma}{\mathrm{d}\Omega_0}$ 





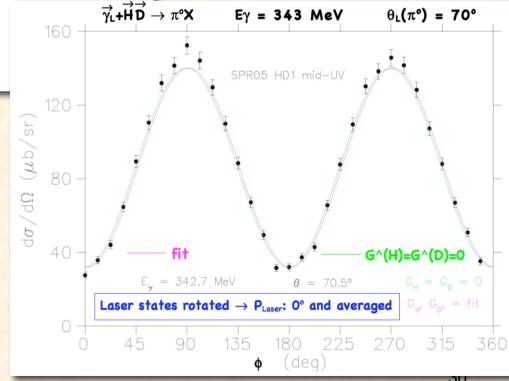
# Extraction of observable G linearly polarized photons on longitudinally polarized targets

$$d\sigma = d\sigma_o(HD) + P_{\gamma}^L \cdot \left[ \hat{\Sigma}(HD) + \frac{1}{\sqrt{2}} P_D^T \cdot T_{20}^L(D) \right] \cdot \cos 2\phi$$

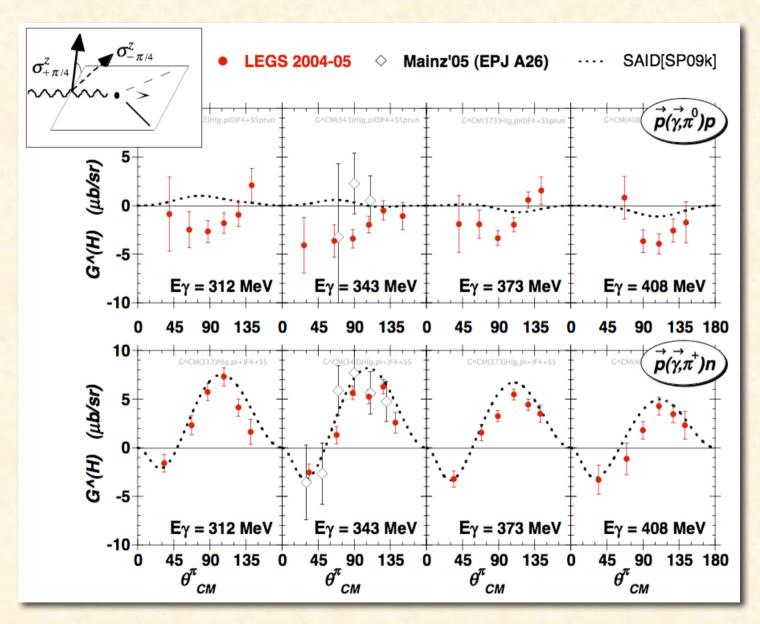
$$+ P_{\gamma}^L \cdot \left[ P_H \cdot \hat{G}(H) + P_D^V \cdot \hat{G}(D) \right] \cdot \sin 2\phi$$

$$- P_{\gamma}^C \cdot \left[ P_H \cdot \hat{E}(H) + P_D^V \cdot \hat{E}(D) \right] + \frac{1}{\sqrt{2}} P_D^T \cdot T_{20}^0(D)$$

$$\uparrow_{L} + \overrightarrow{HD} \rightarrow 0$$



#### G asymmetry from $\pi^+$ and $\pi^0$ photoproduction on the proton at LEGS



$$\hat{G} = G \times \frac{d\sigma}{d\Omega_0}$$

Surprise: opposite sign and one order of magnitude larger than expected.

Under investigation.

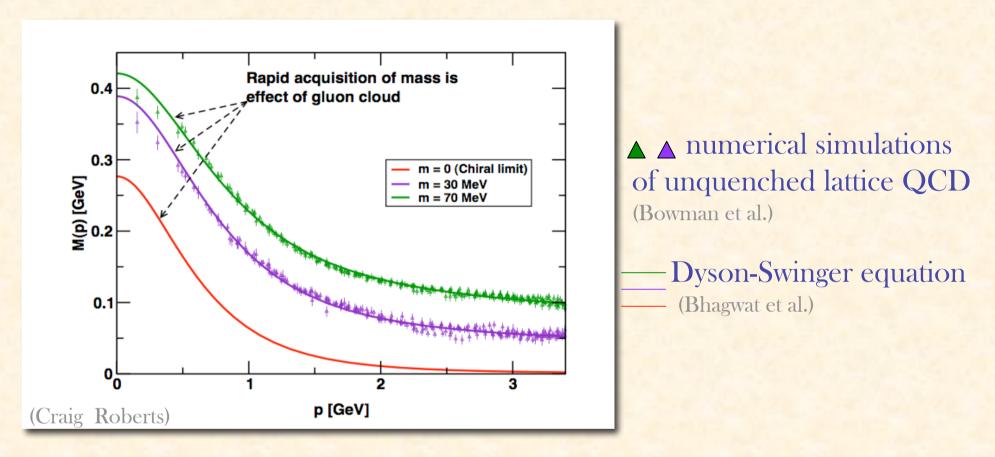
D-wave component under  $P_{33}(1232)$  larger than expected.

#### Conclusions

- $\Sigma$  asymmetry for  $\pi^0$  and  $\pi^-$  production on quasi-free neutrons provided new challenging constraints on  $P_{11}(1700)$  and  $P_{13}(1720)$  properties.
- $\Sigma$  asymmetry for  $\omega$  photoproduction on the nucleon is a benchmark prediction for most existing models sensitive to  $P_{13}(1720)$  resonance.
- Double polarization observables in  $k^+\Lambda$  photoproduction are mostly consistent with Bonn-Gatchina CC-PWA predictions the role of the "missing"  $D_{13}$ (1900) is still uncertain.
- First results on G double polarization observable at LEGS suggest a strong D-wave component in the  $\Delta$  resonance.
- The next step is performing complete experiments.

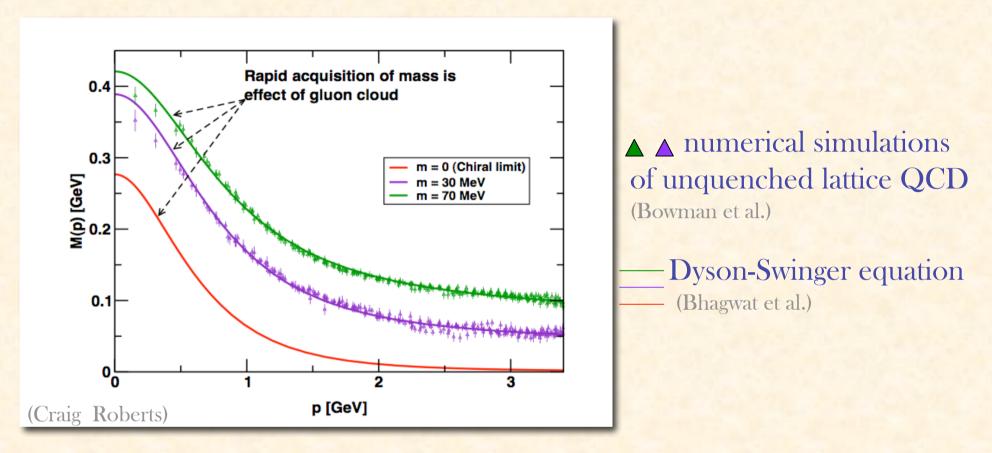
# Backup slides

#### Hadron Models: connection between constituent and current quarks



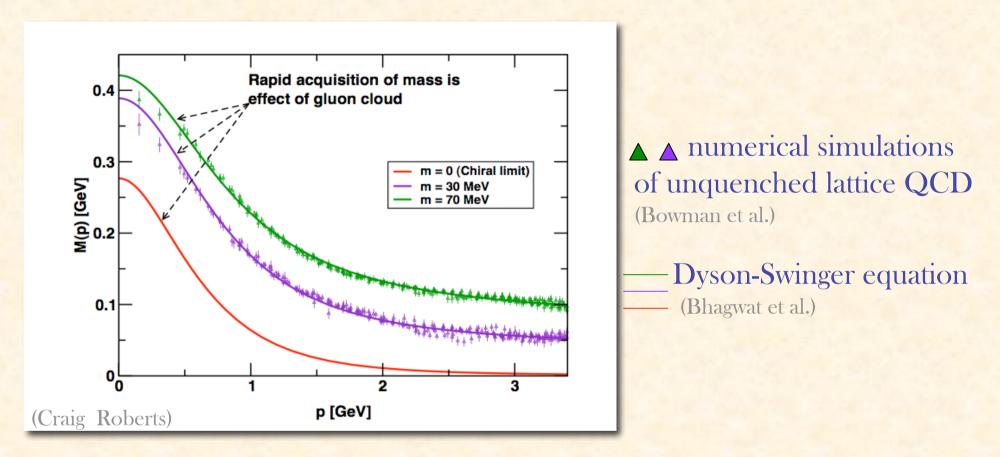
Current-quarks of perturbative QCD evolve into constituent quarks at low momentum \_\_\_\_\_\_ the constituent quark mass arises from low momentum gluons attaching them selves to current quarks.

#### Hadron Models: connection between constituent and current quarks



This effect is a dynamical chiral symmetry breaking (DCSB): a non-perturbative QCD effect that occurs also at the chiral limit generates mass from nothing

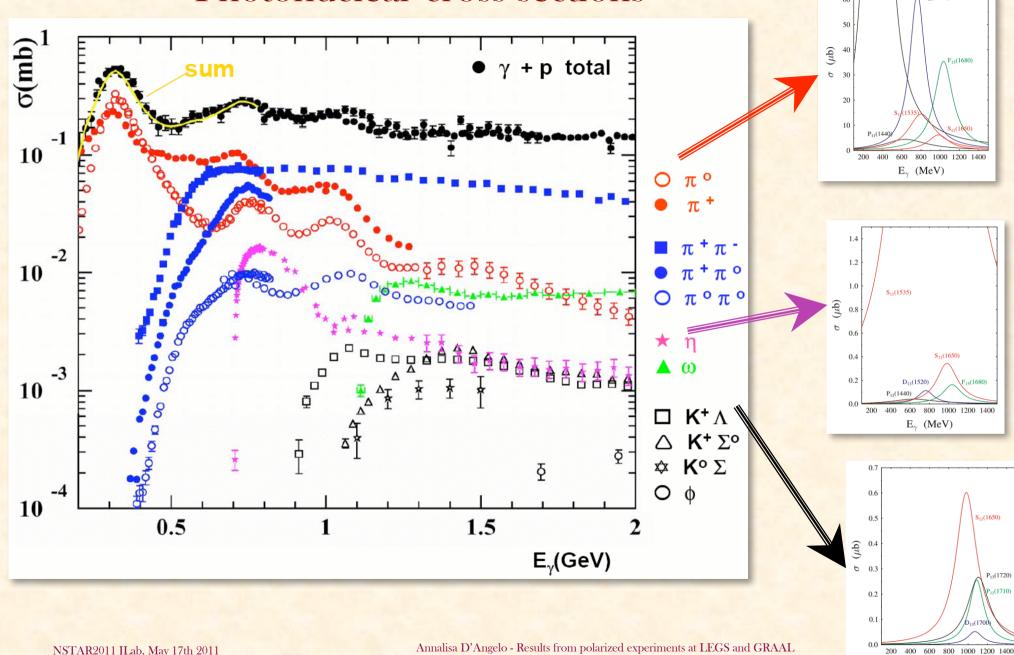
#### Hadron Models: connection between constituent and current quarks



The interaction that describes color-singlet mesons also generates axial-vector isotriplet quark-quark correlations with <u>significant attraction</u>:

$$m[ud]_0 = 0.74 - 0.82 \text{ GeV}$$
  
 $m[ud]_1 = m[uu]_1 = m[dd]_1 = 0.95 - 1.02 \text{ GeV}$  di-Quarks





P33(1232)

 $E_{\gamma}$  (MeV)

## From the Experiment to Theory

Experiment cross section, spin observables

 $\sigma, d\sigma/d\Omega, \Sigma, P, T$ 

(beam-target) E,F,G,H,

(beam-recoil)  $C_x, C_z, O_x, O_z$ ,

(target-recoil)  $L_x, L_z, T_x, T_z$ 

Amplitude analysis

→multipole amplitudes

→phase shifts

Reaction Theory
dynamical frameworks

Theory LQCD, quark models, QCD sum rules,

Coupled channels: resonance parameter extraction

(by Eugene Pasyuk)

## From the Experiment to Theory

Idealized path to search for  $N^*$ ,  $\Delta^*$  states via meson photo-production:

(1) determine the production amplitude from experiment

search for resonant structure: Argand circles, phase motion speed plots, etc.

Experiment cross section, spin observables

 $\sigma$ ,d $\sigma$ /d $\Omega$ ,Σ,P,T (beam-target) **E,F,G,H,** (beam-recoil) **C**<sub>x</sub>,**C**<sub>z</sub>,**O**<sub>x</sub>,**O**<sub>z</sub>, (target-recoil) **L**<sub>x</sub>,**L**<sub>z</sub>, **T**<sub>x</sub>,**T**<sub>z</sub>,

components

determine resonant γN\* and decay couplings; contact with LQCD, DSE, Hadron models

(A. Sandorfi et al.)

<u>Theory</u> LQCD, quark models, QCD sum rules, Coupled channels: resonance parameter extraction

(by Eugene Pasyuk)

Reaction Theory dynamical frameworks

## From the Experiment to Theory

Idealized path to search for  $N^*$ ,  $\Delta^*$  states via meson photo-production:

(1) determine the production amplitude from experiment

search for resonant structure: Argand circles, phase motion speed plots, etc.

Never been done after 50 years of experiments

(2) separate resonance and background components

determine resonant γN\* and decay couplings; contact with LQCD, DSE, Hadron models

(A. Sandorfi et al.)

Without exp Amplitudes models have conjectured resonances and adjusted couplings to compare with limited data