# Phenomenological extraction of Transversity from COMPASS SIDIS and Belle e<sup>+</sup>e<sup>-</sup> data

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QCD Evolution 2014
Santa Fe



### with the logitudinally polarised 160 GeV $\mu^+$ beam and

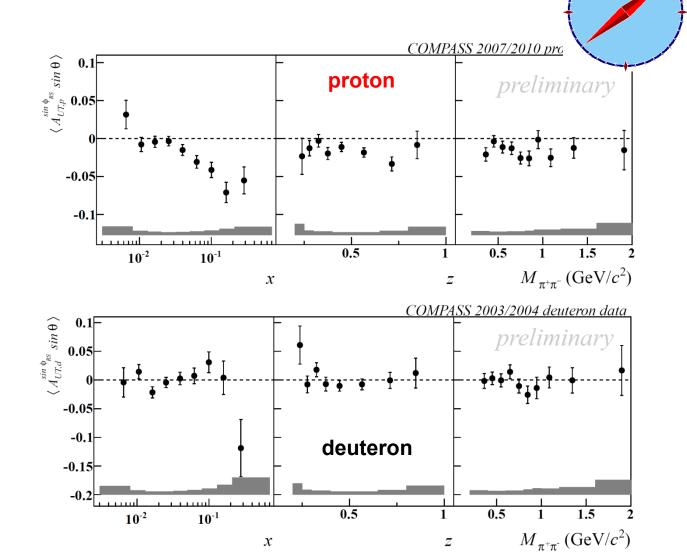
- unpolarised d target
  - h<sup>±</sup> multiplicities vs p<sub>t</sub><sup>2</sup>
  - h + azimuthal asymmetries
- longitudinally polarised d target
  - h<sup>±</sup> azimuthal asymmetries
- transversely polarised d and p target
  - Collins and Sivers asymmetries for h<sup>±</sup>, π<sup>±</sup>, K<sup>±</sup>, K<sup>0</sup>
  - the other 6 transverse spin asymmetries for h<sup>±</sup>
  - dihadron asymmetry for h<sup>+</sup>h<sup>-</sup>,  $\pi$ <sup>+</sup> $\pi$ <sup>-</sup>, K<sup>+</sup>K<sup>-</sup>,  $\pi$ <sup>+</sup>K<sup>-</sup>,  $\pi$ <sup>-</sup>K<sup>+</sup>

and other results will come soon

multidimensional analysis of TSA

some examples

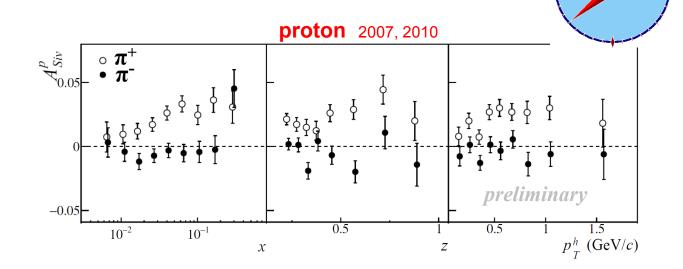
dihadron asymmetries  $\pi^+\pi^-$ 

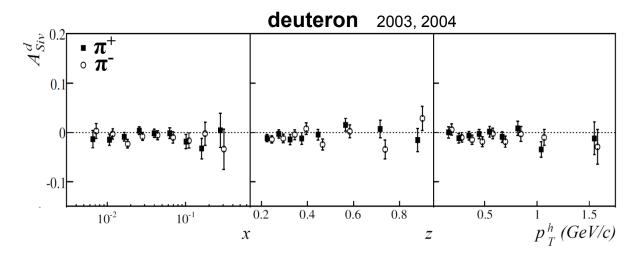


COMPAS

some examples

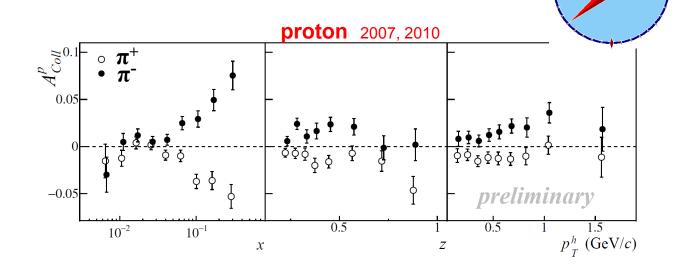
Sivers asymmetry



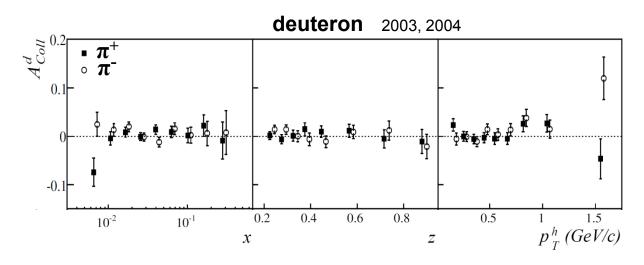


some examples

Collins asymmetry



**COMPA** 

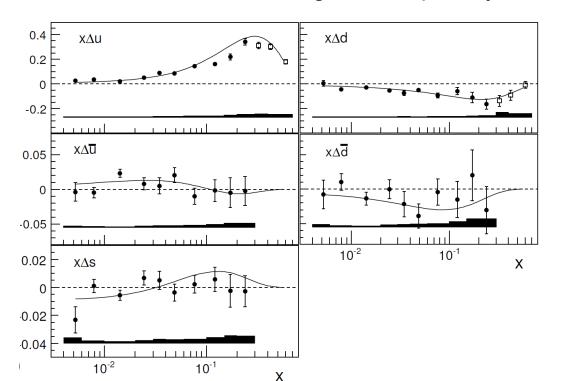


# COMI

measurements on p and d at the same beam momentum

with the same x binning  $\rightarrow$  same < x >,  $< Q^2 >$  in each bin this allows to combine p and d asymmetries to perform flavour separation and point-to-point extraction of PDFs

as done in the case of longitudinal spin asymmetries using the DSS FFs



COMPASS Collaboration, PLB 693 (2010) 227



#### new:

COMPASS result on transversity measured in each x bin from pion-pair asymmetry on p and d using results of the Pavia group analysis for the FFs

C. Braun, DIS2014 (PhD Thesis)

#### method:

COMPASS results for  $A_d^{2\pi}$  and  $A_p^{2\pi}$  as function of x

use the same coefficients evaluated by A. Bacchetta et al. from Belle data [JHEP1303 (2013)119]

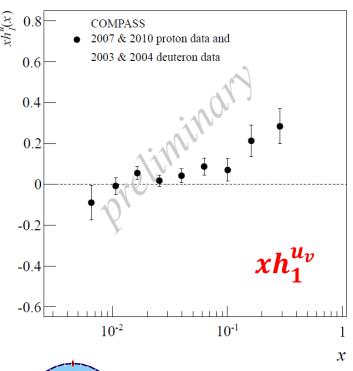
numerical values for  $4xh_1^{u_v} - xh_1^{d_v}$  and  $xh_1^{u_v} + xh_1^{d_v}$  in each x bin

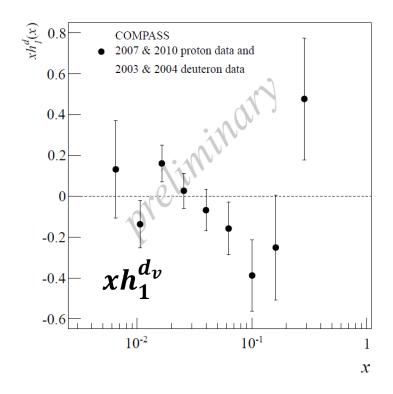


numerical values for  $xh_1^{u_v}$  and  $xh_1^{d_v}$  in each x bin

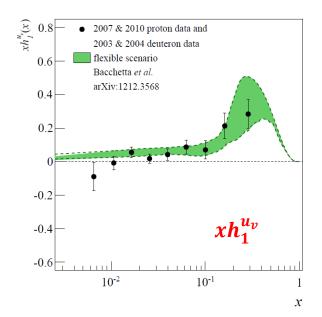


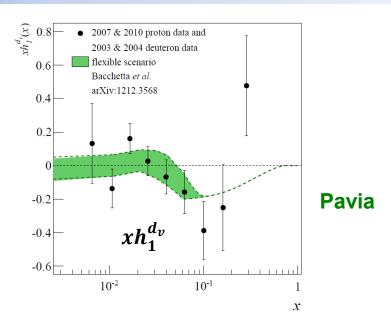
### results:





COMPASS DIS2014

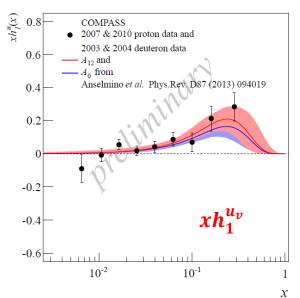


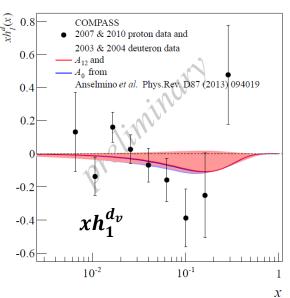


### **Torino**



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**this work** [F. B., Anna Martin, Vincenzo Barone]

COMPASS results for  $A_d^{2h}$  and  $A_p^{2h}$  as function of x



directly use the Belle data following Bacchetta et al., PRL107(2011)012001

numerical values for  $4xh_1^{u_v} - xh_1^{d_v}$  and  $xh_1^{u_v} + xh_1^{d_v}$ in each x bin



numerical values for  $xh_1^{u_v}$  and  $xh_1^{d_v}$  in each x bin

and follow a similar procedure for the Collins asymmetries, without using parametrisations for Collins FFs and transversity PDFs

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### results obtained using

Belle results for pion and pion-pair asymmetries

PRL 107(2011)072004, PRD78(2008)032011 / 86(2012)039905

- COMPASS results on
  - p and d dihadron asymmetries vs x (integrated over z, M)
  - -p and d Collins asymmetry vs x (integrated over z,  $p_T$ )

h<sup>+</sup> and h<sup>-</sup> assuming that all hadrons are pions

unpolarised PDFs and FFs parametrisations

- PDFs: CTEQ5D

-FFs: DSS LO

# dihadron asymmetries

$$A^{2h} = \frac{R}{M_{2h}} \frac{\sum_{q,\bar{q}} e_q^2 x h_1^q H_q^{\angle}}{\sum_{q,\bar{q}} e_q^2 x f_1^q D_q} = \frac{\sum_{q,\bar{q}} e_q^2 x h_1^q H_q}{\sum_{q,\bar{q}} e_q^2 x f_1^q D_q} \qquad \text{measured as function} \\ H_q(z,M_{2h}) = \sin\theta_q \cdot R/M_{2h} \cdot H_q^{\angle}(z,M_{2h})$$

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### with reasonable assumptions on the FFs

$$H_q = -H_{\bar{q}}, \ H_u = -H_d, \ H_s = H_c = 0$$
  
 $D_u = D_d = D_{\bar{u}} = D_{\bar{d}}, \ D_s = D_{\bar{s}}, \ D_c = D_{\bar{c}}$   $D_s \simeq D_c \simeq 0.5 D_u$ 

Bacchetta et al. PRL107(2011)012001

$$D_s \simeq D_c \simeq 0.5 D_u$$

### neglecting s and c quark contributions, and integrating over z, M:

$$A_p^{2h}(x) \simeq \frac{4xh_1^{u_v} - xh_1^{d_v}}{4xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{d_v}}{xf_1^{*u} + xf_1^{*d}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{xf_1^{*u} + xf_1^{*u}} < H_u > A_d^{2h}(x) \simeq \frac{3}{5} \frac{xh_1^{u_v} + xh_1^{u_v}}{x$$

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### dihadron asymmetry - Belle data

$$a_{12} = \frac{s^2}{1 + c^2} \frac{\sum_q e_q^2 H_q H_{\bar{q}}}{\sum_q e_q^2 D_q D_{\bar{q}}} \qquad \text{measured as function of } z_1, z_2, M_1, M_2$$
 
$$s^2 = \sin^2 \theta, \ c^2 = \cos^2 \theta$$
 
$$H_q(z, M_{2h}) = \sin \theta_q \cdot R / M_{2h} \cdot H_q^{\angle}(z, M_{2h})$$

with the previous assumptions on the FFs and  $D_c$  fixed in order to reproduce the charm yield, the fully integrated  $a_{12}$  asymmetry given by Belle is

$$a_{12}^{I} \simeq -\frac{5}{8} \frac{s^2}{1+c^2} \left(\frac{\langle H_u \rangle}{\langle D_u \rangle}\right)^2$$

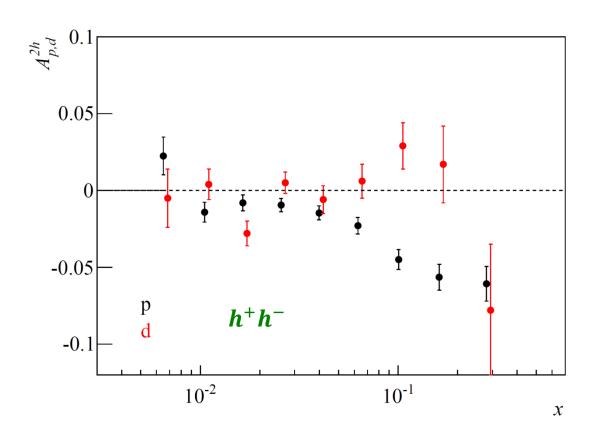
$$a_{12}^{I} = -0.0196 \pm 0.0002 \pm 0.0022$$

and the anaysing power is

$$< a_P > = \frac{< H_u >}{< D_u >} = 0.203$$

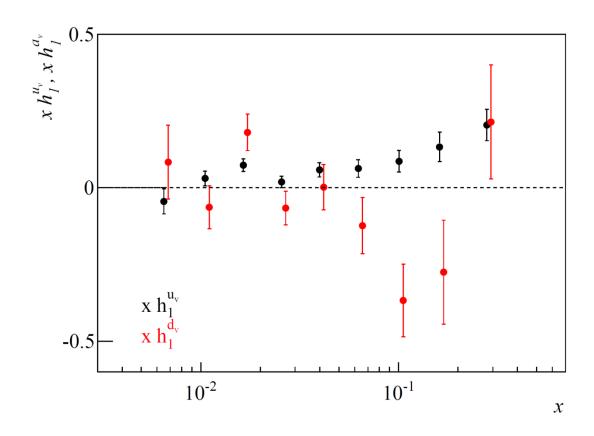
here we have used this value at all COMPASS Q<sup>2</sup>, neglecting evolution

effect evaluated by the Pavia group: ~ -8%



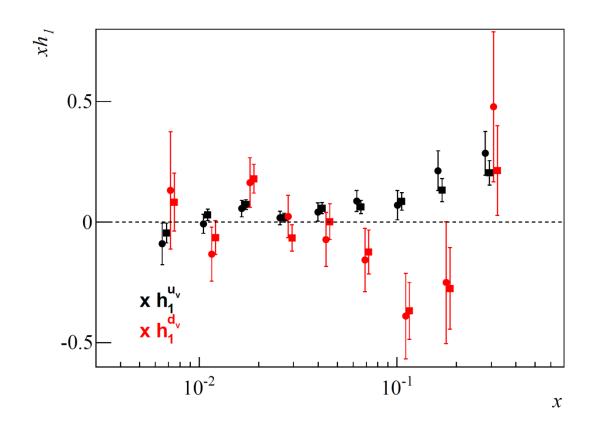
# dihadron asymmetry - transversity

### present result



### dihadron asymmetry - transversity

present result 
$$u = d$$
  $h^+h^-$  compared with DIS2014  $u = d$   $u = d$ 



# **Collins asymmetries**

### interlude

# some thoughts on the interplay between the Collins and the dihadron FFs

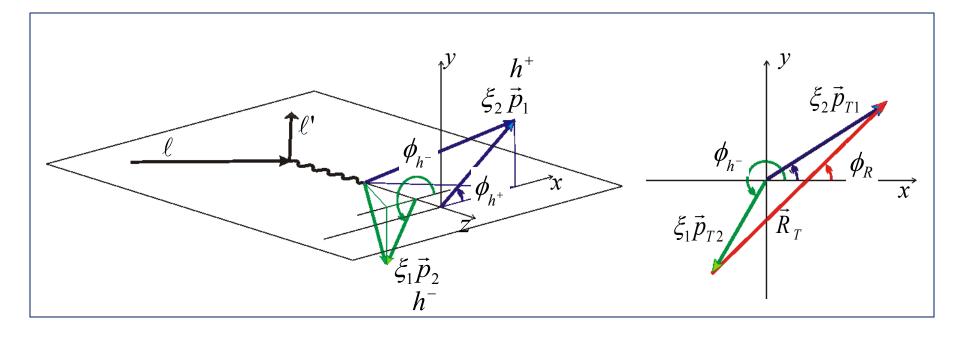
see FB, Structure of Nucleons and Nuclei, Como,10-14 June 2013, FB, Dspin-13, Dubna, October 8-12 2013, C. Adolph et al. [COMPASS Collaboration], hep-ex/1401.7873

### interlude - dihadron and Collins asymmetries

$$N_{2h}^{\pm}(\phi_{RS}) = N_{2h}^{0} \cdot \{1 \pm A \cdot \sin \phi_{RS} \}$$

$$\phi_{RS} = \phi_R + \phi_S - \pi$$

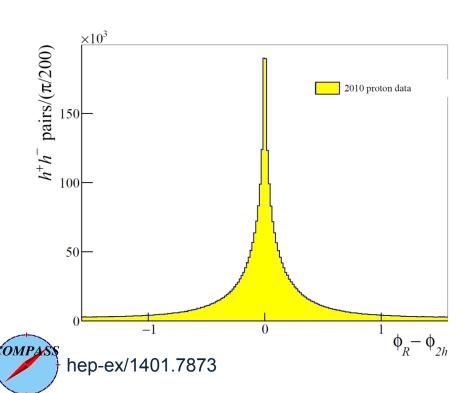


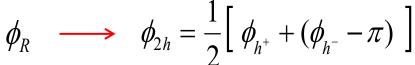


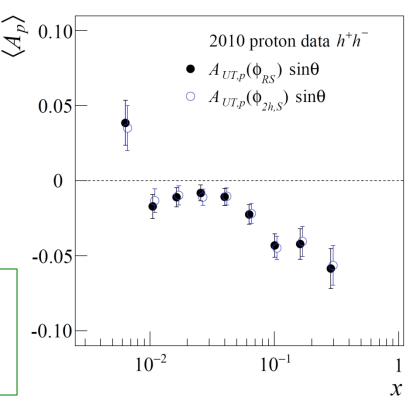
$$\phi_R \longrightarrow \phi_{2h} = \frac{1}{2} \left[ \phi_{h^+} + (\phi_{h^-} - \pi) \right]$$

 $mod \; 2\pi$ 

### interlude - dihadron and Collins asymmetries







... Due to local compensation of transverse momentum, the one-particle Collins effect generates a two-particle effect, and viceversa (X. Artru, arXiv:hep-ph/0207309)

# **Collins asymmetries**

in this case, from the Belle data one has to calculate the analysing power

$$A_P = \frac{\langle H^{fav} \rangle}{\langle D^{fav} \rangle}$$

### we have used the asymmetry (corrected for charm contribution)

$$A_{12}^{UL}(z_1, z_2) = \frac{\langle s^2 \rangle}{\langle 1 + c^2 \rangle} \left[ P_U(z_1, z_2) - P_L(z_1, z_2) \right]$$

integrated over  $M_1, M_2$ 

where

$$P_{U}(z_{1}, z_{2}) = \frac{\sum_{q} e_{q}^{2} [H_{1q}^{+}(z_{1}) H_{1\bar{q}}^{-}(z_{2}) + H_{1q}^{-}(z_{1}) H_{1\bar{q}}^{+}(z_{2})]}{\sum_{q} e_{q}^{2} [D_{1q}^{+}(z_{1}) D_{1\bar{q}}^{-}(z_{2}) + D_{1q}^{-}(z_{1}) D_{1\bar{q}}^{+}(z_{2})]}$$

$$P_L(z_1, z_2) = \frac{\sum_q e_q^2 [H_{1q}^+(z_1) H_{1\bar{q}}^+(z_2) + H_{1q}^-(z_1) H_{1\bar{q}}^-(z_2)]}{\sum_q e_q^2 [D_{1q}^+(z_1) D_{1\bar{q}}^+(z_2) + D_{1q}^-(z_1) D_{1\bar{q}}^-(z_2)]}$$

$$H_{1q}^{\pm} = H_{1(q \to \pi^{\pm})}^{\perp (1/2)}, \quad D_{1q}^{\pm} = D_{1(q \to \pi^{\pm})}$$

Efremov et al., PRD73 (2006) Bacchetta et al., PLB659 (2008) Anselmino et al., PRD75 (2007) Seidl et al., PRD78 (2008)

### for the FFs we have made the assumptions

ignoring the c and s quark contributions, in the case  $z_1 = z_2 = z$  it is

$$A_{12}^{UL}(z) = \frac{\langle s^2 \rangle}{\langle 1 + c^2 \rangle} \left[ \frac{H_1^{fav}(z)}{D_1^{fav}(z)} \right]^2 B(z)$$

where 
$$B(z) = \frac{b(z)[1 + a^2(z)] - [1 + b^2(z)]a(z)}{b(z)[1 + b^2(z)]}$$
  $a(z) = \frac{H_1^{dis}(z)}{H_1^{fav}(z)}$   $b(z) = \frac{D_1^{dis}(z)}{D_1^{fav}(z)}$ 

not so simple as in the 2h case →

### we have done 2 alternative assumptions

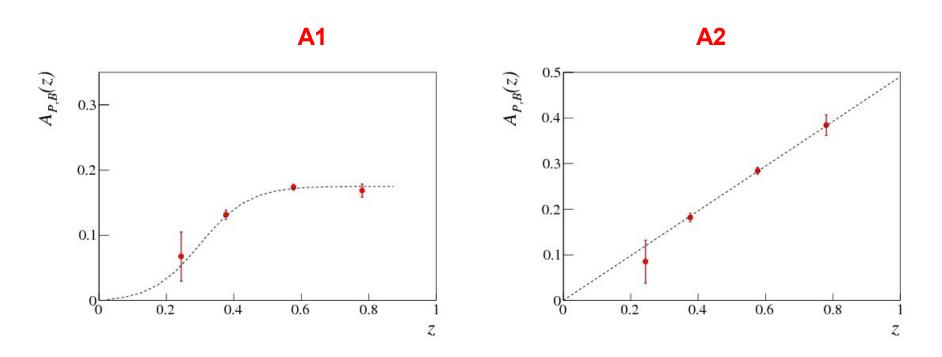
**A1** 
$$H_1^{fav}(z) = -H_1^{dis}(z)$$
 i.e.  $a(z) = -1$ 

**A2** 
$$\frac{H_1^{fav}(z)}{D_1^{fav}(z)} = -\frac{H_1^{dis}(z)}{D_1^{dis}(z)} \quad i.e. \quad a(z) = -b(z)$$

both in agreement with the considerations on the "interplay between the Collins and the dihadron FFs"

and already used / suggested / found as a result of global fits

- these assumptions allow to evaluate  $\frac{H_1^{fav}(z)}{D_1^{fav}(z)}$  in the four z bins
- the values are then fitted with a function of z



to obtain the analysing power the functions are integrated over z

### finally:

**A1** 
$$\frac{\langle H_1^{fav} \rangle}{\langle D_1^{fav} \rangle} \sim 0.10$$
 **A2**  $\frac{\langle H_1^{fav} \rangle}{\langle D_1^{fav} \rangle} \sim 0.18$ 

if the evolution of  $H_1^{fav}$  is negligible;

if the evolution of  $H_1^{fav}$  is the same as that of  $D_1^{fav}$  the analysing powers decrease by  $\sim 10\%$ 

# Collins asymmetry – COMPASS data

$$A_{Coll}^{\pm}(x,z) = \frac{\sum_{q,\bar{q}} e_q^2 x h_1^q(x) \otimes H_{1q}^{\pm}(z)}{\sum_{q,\bar{q}} e_q^2 x f_1^q(x) \otimes D_{1q}^{\pm}(z)}$$

### "gaussian ansatz":

$$A_{Coll}^{\pm}(x,z) = C_G \cdot \frac{\sum_{q,\bar{q}} e_q^2 x h_1^q(x) H_{1q}^{\pm}(z)}{\sum_{q,\bar{q}} e_q^2 x f_1^q(x) D_{1q}^{\pm}(z)}$$

$$C_G = \frac{1}{\sqrt{1 + z^2 < p_{h_1}^2 > / < p_{H_1}^2 > }}$$

Efremov et al., PRD73 (2006)

### we have assumed

- $C_G = 1$
- the previous relations among the FFs
- the s and c quark contributions to be negligible

# Collins asymmetry – COMPASS data

### the measured asymmetries as function of x can be written as

$$A_{Coll,p}^{+} = \frac{\langle H_{1}^{fav} \rangle 4(xh_{1}^{u} + \alpha xh_{1}^{\bar{u}}) + (\alpha xh_{1}^{d} + xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{d}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + \alpha xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{d} + xh_{1}^{\bar{u}})}{\langle D_{1}^{fav} \rangle} \\ A_{Coll,p}^{-} = \frac{\langle H_{1}^{fav} \rangle 4(\alpha xh_{1}^{u} + xh_{1}^{\bar{u}}) + (xh_{1}^{u} + xh_{1}^{\bar{u}})}{\langle D_{1}^{\bar{u}} \rangle} \\ A_{Col$$

$$\alpha = \frac{\langle H_1^{dis} \rangle}{\langle H_1^{fav} \rangle} = \begin{bmatrix} -1 & A1 \\ \langle D_1^{dis} \rangle \\ \overline{\langle D_1^{fav} \rangle} & A2 \end{bmatrix}$$

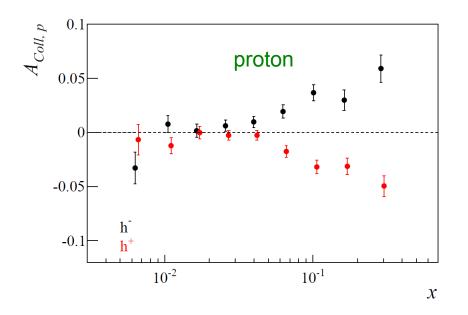
$$\begin{array}{lll} A_{Coll,d}^{+} & = & \frac{< H_{1}^{fav}>}{< D_{1}^{fav}>} \frac{(xh_{1}^{u}+xh_{1}^{d})(4+\alpha)+(xh_{1}^{\bar{u}}+xh_{1}^{\bar{d}})(1+4\alpha)}{d_{d}^{+}} & \\ A_{Coll,d}^{-} & = & \frac{< H_{1}^{fav}>}{< D_{1}^{fav}>} \frac{(xh_{1}^{u}+xh_{1}^{d})(4\alpha+1)+(xh_{1}^{\bar{u}}+xh_{1}^{\bar{d}})(4+\alpha)}{d_{\bar{d}}^{-}} & \\ \end{array} \begin{array}{l} \text{corresponding quantities with unpole PDFs and FFs} \\ \text{PDFs and FFs} & \\ \end{array}$$

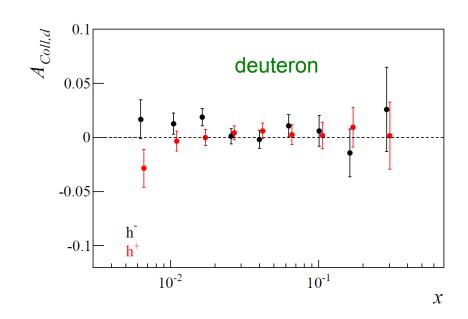
$$A_{Coll,d}^{-} = \frac{\langle H_1^{fav} \rangle}{\langle D_1^{fav} \rangle} \frac{(xh_1^u + xh_1^d)(4\alpha + 1) + (xh_1^{\bar{u}} + xh_1^{\bar{d}})(4 + \alpha)}{d_d^{-}}$$

from Belle

neglecting qbar transversity, in each x bin, the only unknowns are the u and d quark transversity PDFs

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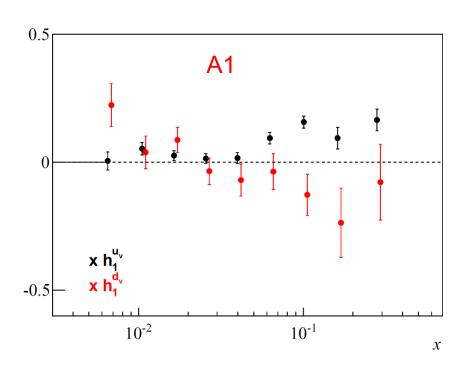


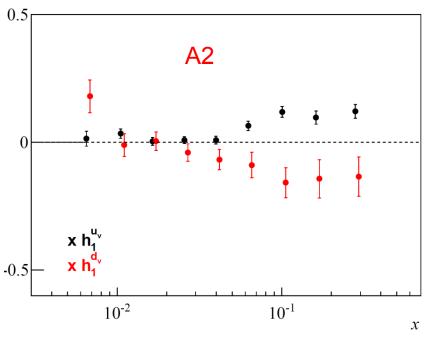


un-identified hadrons

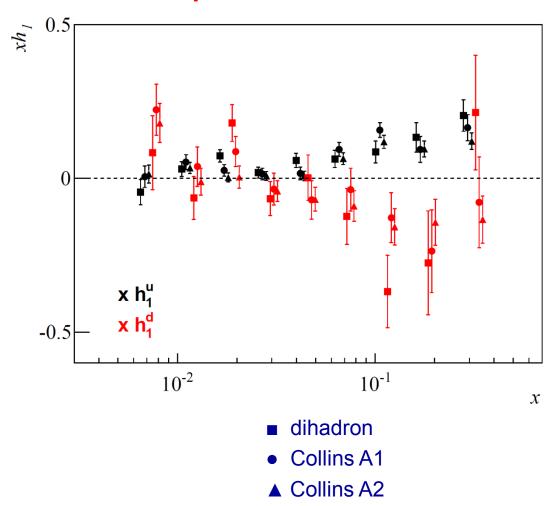
# **Collins asymmetry – transversity**

### results

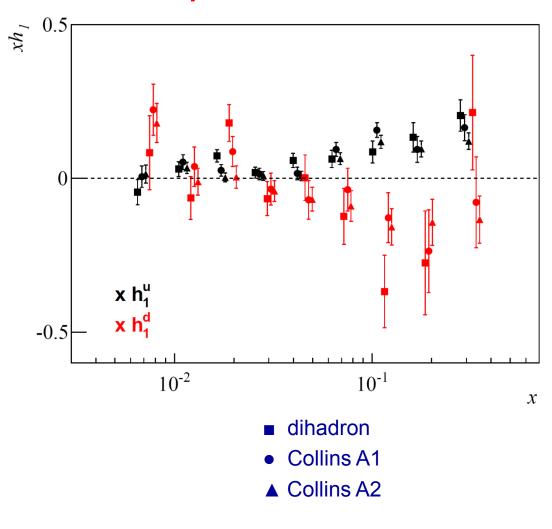




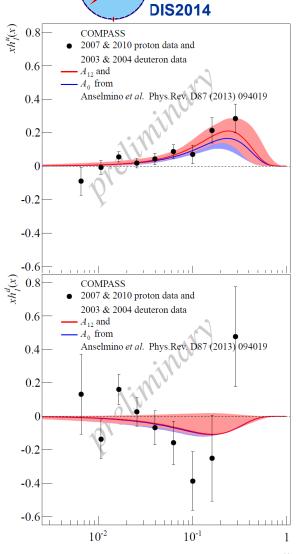
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### summary

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- the physics is simple

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and Q<sup>2</sup> evolution is needed!