

# Summary of the workshop Hadrons in nuclear matter

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## Outline

### 1. Color transparency at intermediate energies

#### *Objectives:*

*Study of interplay of perturbative and nonperturbative QCD in hard processes at intermediate energies*

**Space-time evolution of small quark wave packages**

#### Tools:

#### *Nucleon form factor*

◇ Rescattering kinematics in  $e + D \rightarrow e + p + N$

◇ Transparency in  $e + A \rightarrow e + p + A'$

#### *Two body hard processes*

◇  $\gamma^* + p \rightarrow \rho + p \implies \gamma^* + D \rightarrow \rho + D$

◇ Large angle  $\gamma + n \rightarrow \pi^- + p \implies \gamma + {}^2H \rightarrow \pi^- p + p$

## 2. Quark-gluon structure of short-range correlations in nuclei

*Objective:* "Discovery of the fundamental nature of nuclear matter"

Tools:

◇ Tagged structure functions for DIS scattering off the deuteron:  $e + D \rightarrow \text{backward proton } (\Delta - \text{isobar}) + X$  to probe the origin of the EMC effect and nonnucleonic degrees of freedom in nuclei and measure  $F_{2n}(x, Q^2)/F_{2p}(x, Q^2)$  for  $x \sim 1$

◇ Study of correlation of backward nucleon ( $\Delta$ -isobar) and forward nucleon production to study two and three nucleon short-range correlations in nuclei.

◇ Superfast quarks in nuclei via DIS at  $x > 1$ , and  $Q^2 > 10 \text{ GeV}^2$ .

# Color coherent phenomena:

Color transparency, ...

\* Three components:

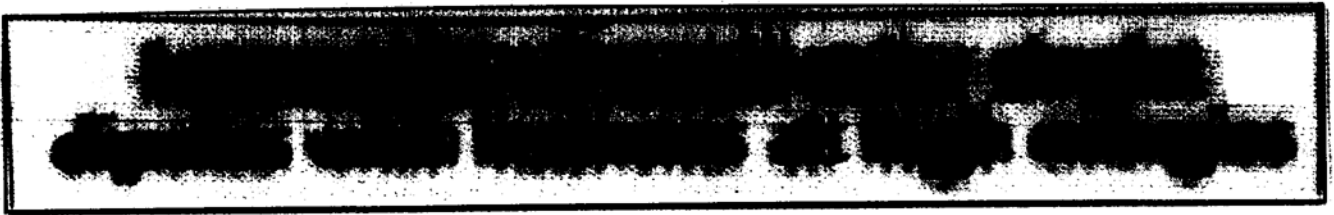
I  $\diamond$  Small color dipoles interact with small cross section

II  $\diamond$  Point-like (PLC) small size configurations are produced in hard exclusive processes

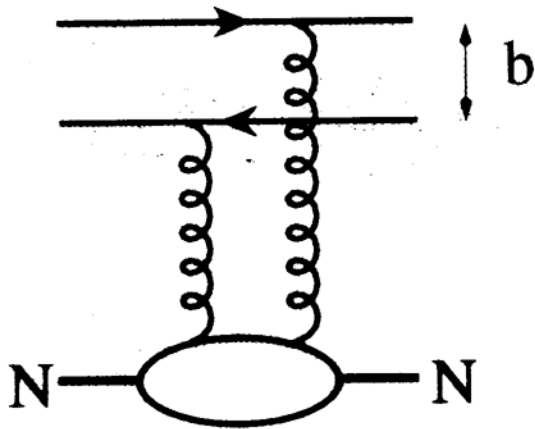
III  $\diamond$  Produced PLC can survive distances comparable to internucleon distances

High Energies: Only I & II

are necessary; III is always satisfied.



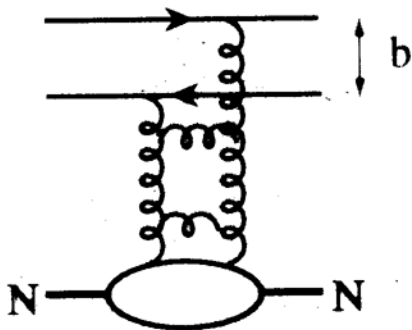
Gauge invariance for a small dipole-hadron interaction →



Two gluon exchange model

$$\sigma = C b^2 \quad (\text{F.Low 75})$$

C does not depend on  $E_{inc}$



pQCD in the leading log b approximation

(Baym, Blattel, FS, 93)

$$\sigma^{inel}(b, E_{inc}) = \frac{\pi^2}{3} b^2 \alpha_s \times G_N(x, \frac{\lambda}{b^2})$$

Qualitative difference from QED: cross section rapidly increases with  $W$  - a fingerprint of small size dipole interaction in DGLAP kinematics. ( $\lambda(x = 10^{-3}, Q^2 = 10 \text{ GeV}^2 \approx 9)$ )

legitimately calculated in pQCD due to color screening for interaction of small color singlets:

◇  $\pi + T \rightarrow 2 \text{ jets} + T$  Frankfurt, Miller, S. 93

◇  $\gamma_L^* N \rightarrow V(\rho, J/\Psi, \rho'..) + N$

Brodsky, Frankfurt, Gunion, Mueller, S., 94

◇  $\gamma_L^* N \rightarrow \text{Meson}(\pi, K, \eta, )$  [Few meson system] + Baryon,  
Collins, FS 96; M. Polyakov 98  $2\pi$  for finite  $x$

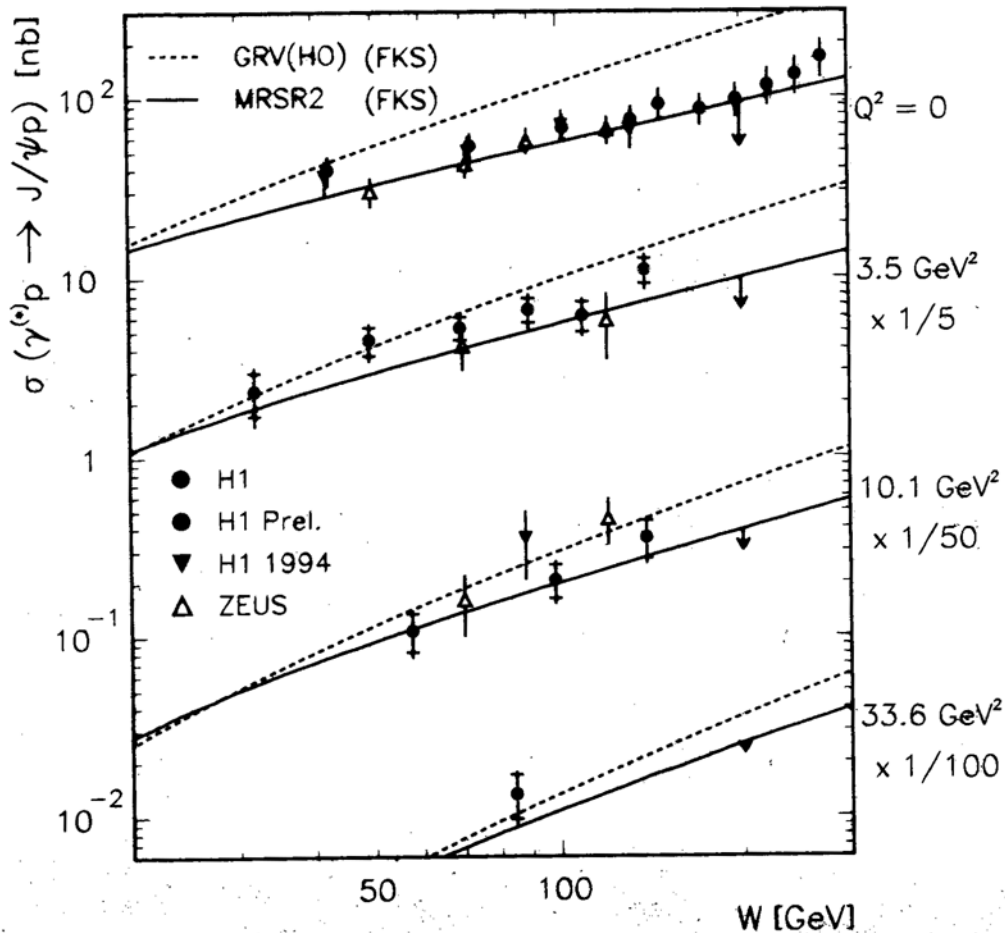
◇  $\gamma_{LP}^* \rightarrow \text{forward } N + \pi, \gamma_{LP}^* \rightarrow \text{forward } \Lambda + K^+$   
 $\gamma_{LP}^* \rightarrow \text{forward } \bar{p} + NN,$  FS & Polyakov 98

◇  $\gamma^* + N \rightarrow \gamma + N$  Bartels & Loewe 82; Dittes, Muller,  
Robaschik, Geyer, Horejse 88; Ji 96-97, Radyushkin 96-98,  
Freund & FS 97, Freund Collins 98

◇  $\gamma^* + \gamma \rightarrow \pi\pi$  M. Diehl, T. Gousset, B. Pire 98

Extensive data on  $VW$  production from HERA support dominance of the  $pQCD$  dynamics. Numerical calculations including finite  $b$  effects in  $\psi_V(b)$  explain key elements of high  $Q^2$  data. The most important ones are:

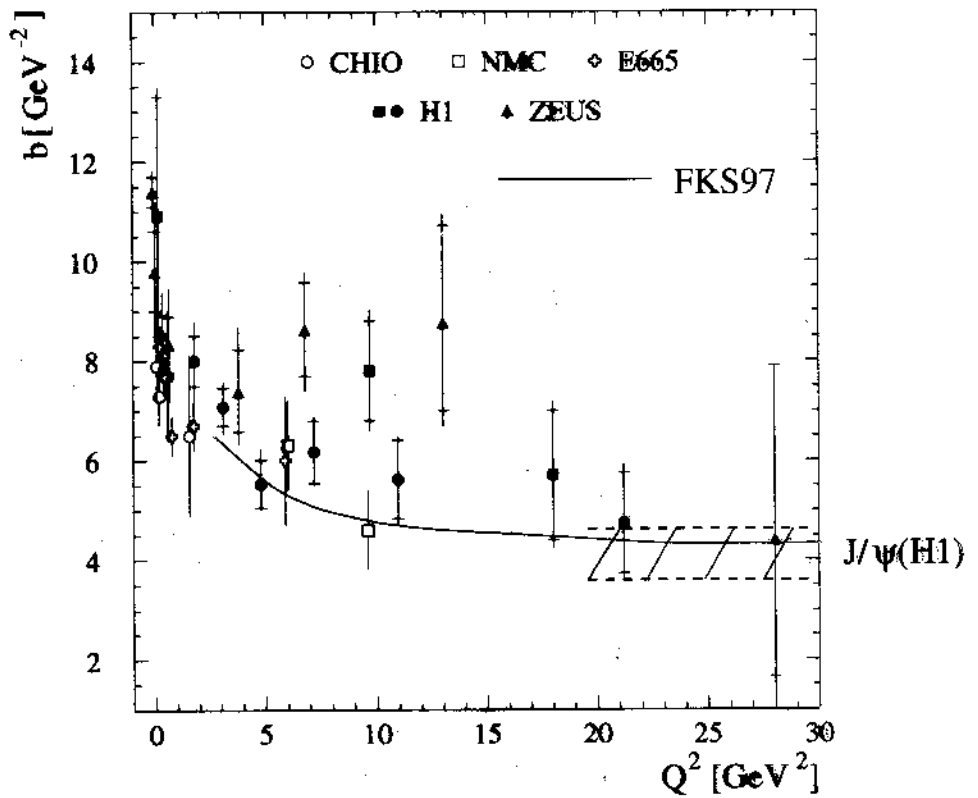
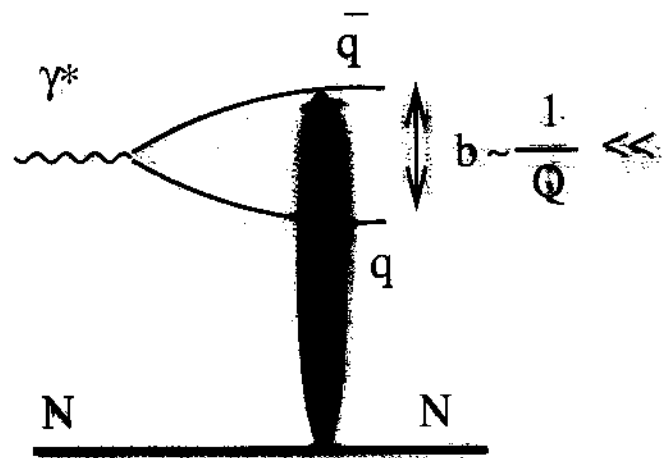
(i) Energy dependence of  $J/\psi$  production; absolute cross section of  $J/\psi, \Upsilon$  production.



(ii) Absolute cross section of  $\rho$  production at  $Q^2 \sim 20-30 \text{ GeV}^2$  and its energy dependence at  $Q^2 \sim 20 \text{ GeV}^2$ . Explanation of the data at lower  $Q^2$  is more sensitive to the higher twist effects, and uncertainties of the low  $Q^2$  gluon densities.

(iii) Convergence of the t slopes

$B$  ( $\sigma = A \exp(Bt)$ ) of  $\rho$ -meson production at large  $Q^2$  and  $J/\psi$  production (Brodsky et al 94)



⇒ Small size  $q\bar{q}$  Fock components are present in light mesons.

⇒ At the transverse separations  $b \leq 0.3 \text{ fm}$  pQCD reasonably describes "small  $q\bar{q}$  dipole" - nucleon interactions.

Color transparency effects for scattering off nuclei are observed at FNAL

(1)  $\gamma + A \rightarrow \psi + A$

(2)  $\gamma^* A \rightarrow p + A'$  30 effect at  $Q^2 \sim 10^2$

(3)  $\pi + A \rightarrow \text{jet}(k_+) + \text{jet}(-k_+) + A$  ( $k_+ > 1.5 \text{ GeV}$ )

suggested by Bartch, Brodsky, Goldhaber, Gunion & pQCD analysis F, Miller, S 93

E 971 A-dependence

Data	$A^{1.61 \pm 0.08}$
BB66	$A^{1/3}$
FMS	$A^{1.52 + E}$
Soft Data	$A^{2/3}$

a factor  $\sim 7$  difference.

$z = \frac{E_{\text{jet}}}{E_\pi}$

dependence:  $\psi_\pi^2(z) \propto (z \cdot (1-z))^2$  [BB66, FMS] consistent with data



⇒ Firm base for using CT

to study QCD at intermediate energies

New features:

- \* Study of CT for "3q" no analogy naively with QED ( $e^+e^-$ )
- \*\* Possibility to study space-time evolution of small wave packet.

At what  $q\bar{q}$  ( $qq$ ) separations chiral symmetry breaking becomes important? Since  $\sigma_{pQCD}(b)$  smaller at low energies ( $\propto b(x,b)$  at larger  $b$ )

- sharper transition

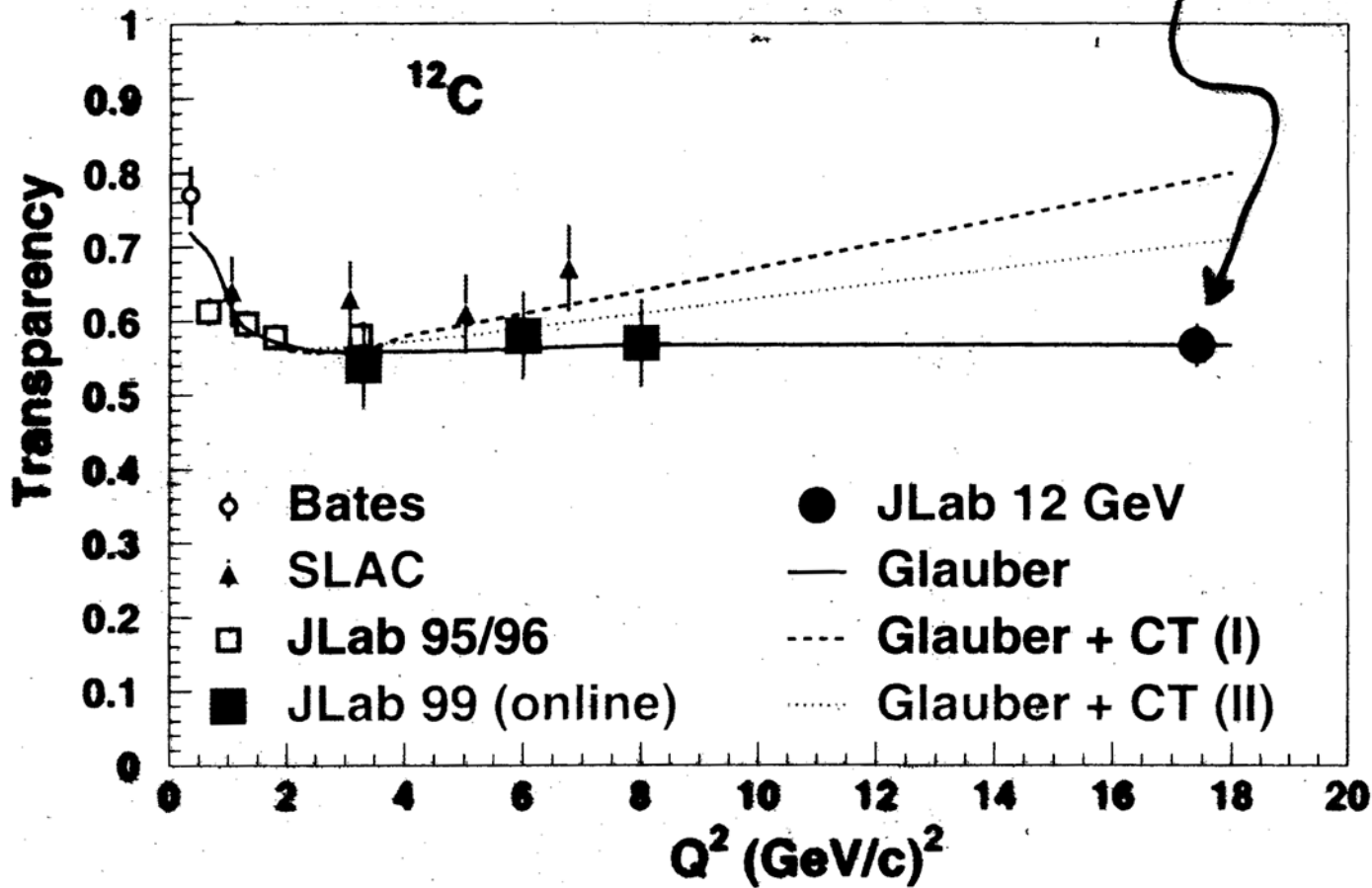
⇒ Need to scan interaction of ejectile as a function of distance from production point

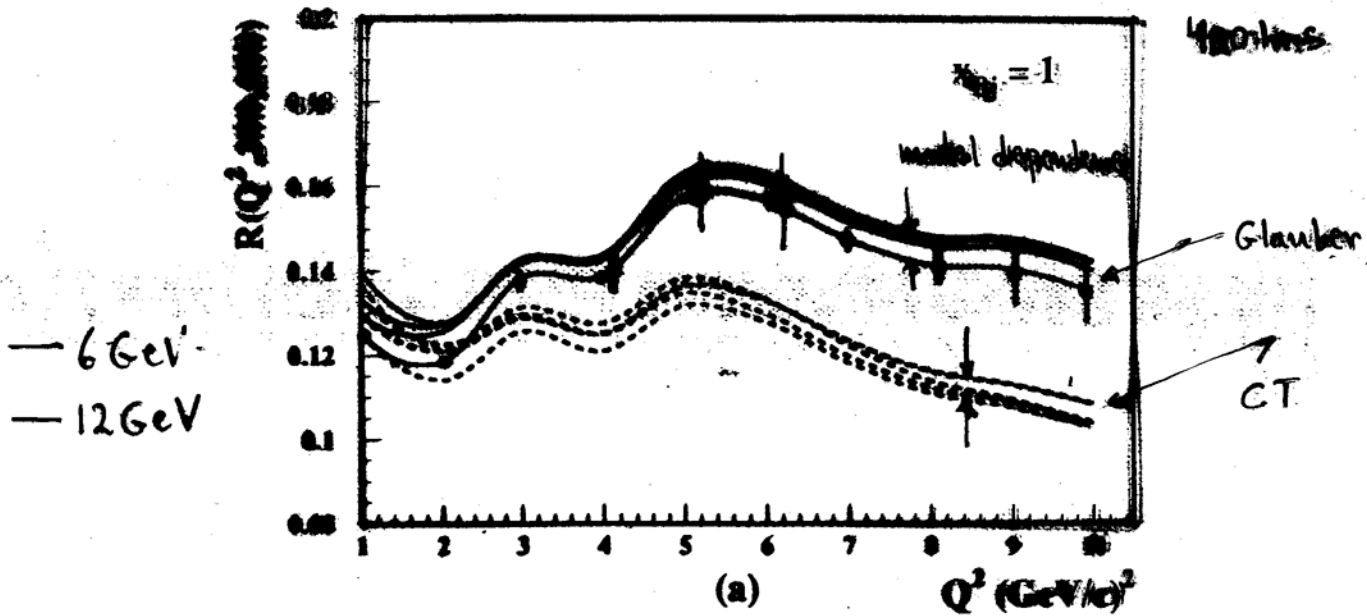
$$E = 12 \text{ GeV}$$

$$q = 10.2 \text{ GeV}/c$$

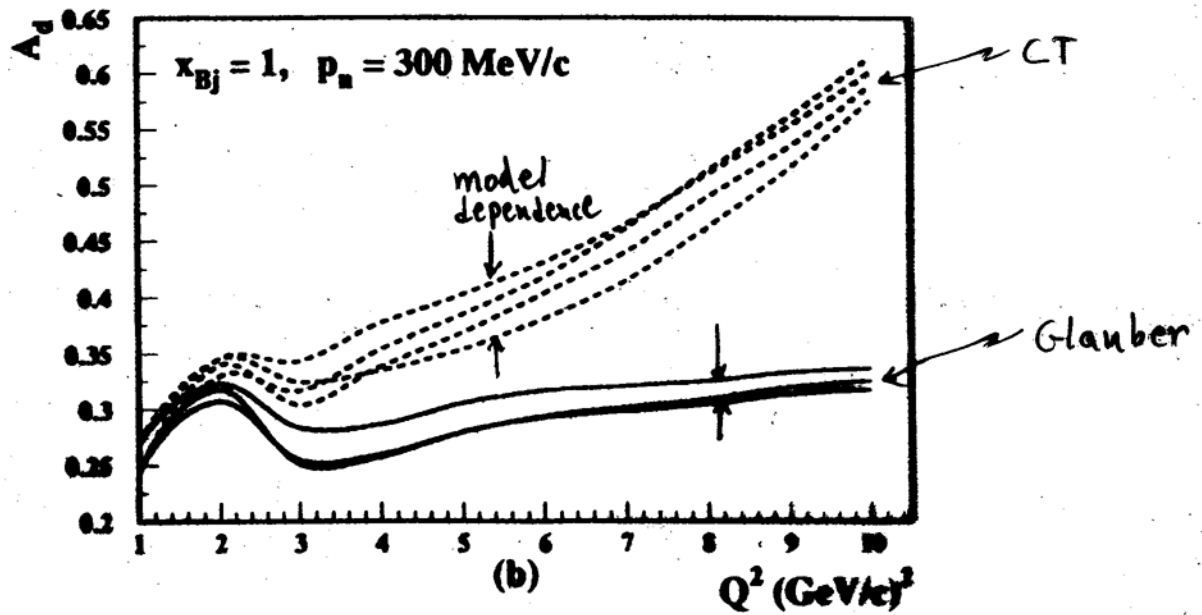
$$Q^2 \approx 17.5 (\text{GeV}/c)^2$$

$\sim \rho_0$  has





— 6 GeV  
 - - 12 GeV



$$A_d = \frac{T_{20}}{\left( \frac{d\sigma^{unp}}{dE_e d\Omega_e d^2p_T} \right)}$$

if one has tensor polarisation

$$R = \frac{J(p_n \approx 300 \text{ MeV/c})}{\sigma(p_n \approx 200 \text{ MeV/c})}$$

## Two complementary methods

Survival - increases with  $Q^2$  (Brodsky & Mueller 88)

$A(e, e'p)$  (distances  $l \geq 2 \text{ fm}$ )

Note:  $Q^2 \geq 26 \text{ GeV}^2$  necessary to avoid  $Q^2$ -dependence of quenching

Reinteraction - decreases with  $Q^2$  (Egiyan, Miller, Sargsian, FS 95)

$eD \rightarrow e p n$  (distances  $l \sim 1 \text{ fm}$ )

R. Ent:  $(e, e'p)$  with SHMS  $Q^2 = 17.5 \text{ GeV}^2$   
feasible  $\square$

K. Griffioen  $eD \rightarrow epn$   $Q^2 = 10 \text{ GeV}^2$   
easy with CLAS  $\square$

Note: BNL E850 confirm transparency

$\geq 2$  Glauber for  $pN \geq 6.6 \text{ GeV}/c$ , hence for these momenta expansion does not mask effect  $\Rightarrow$  if no CT in  $e, e'p$  nucleon

form factor for  $Q^2 \sim 10 \text{ GeV}^2$  due to average size.

# Vector meson production

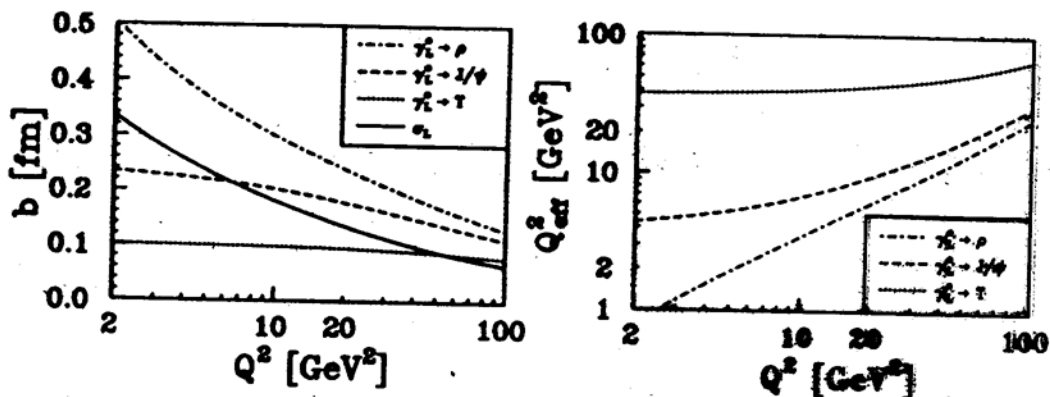
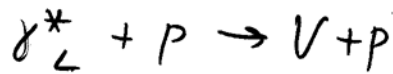


Figure 5. The dependence of average  $b$  and effective  $Q^2$  of  $Q^2$  for production of vector mesons [Köpf, FS 96]

Transverse size maybe  $\ll 2r_p$

even for  $Q^2 = 2 \text{ GeV}^2$  due to structure of wave function of  $\gamma_L^*$ .

How to check?

$$\sigma^* d \rightarrow \rho d'$$

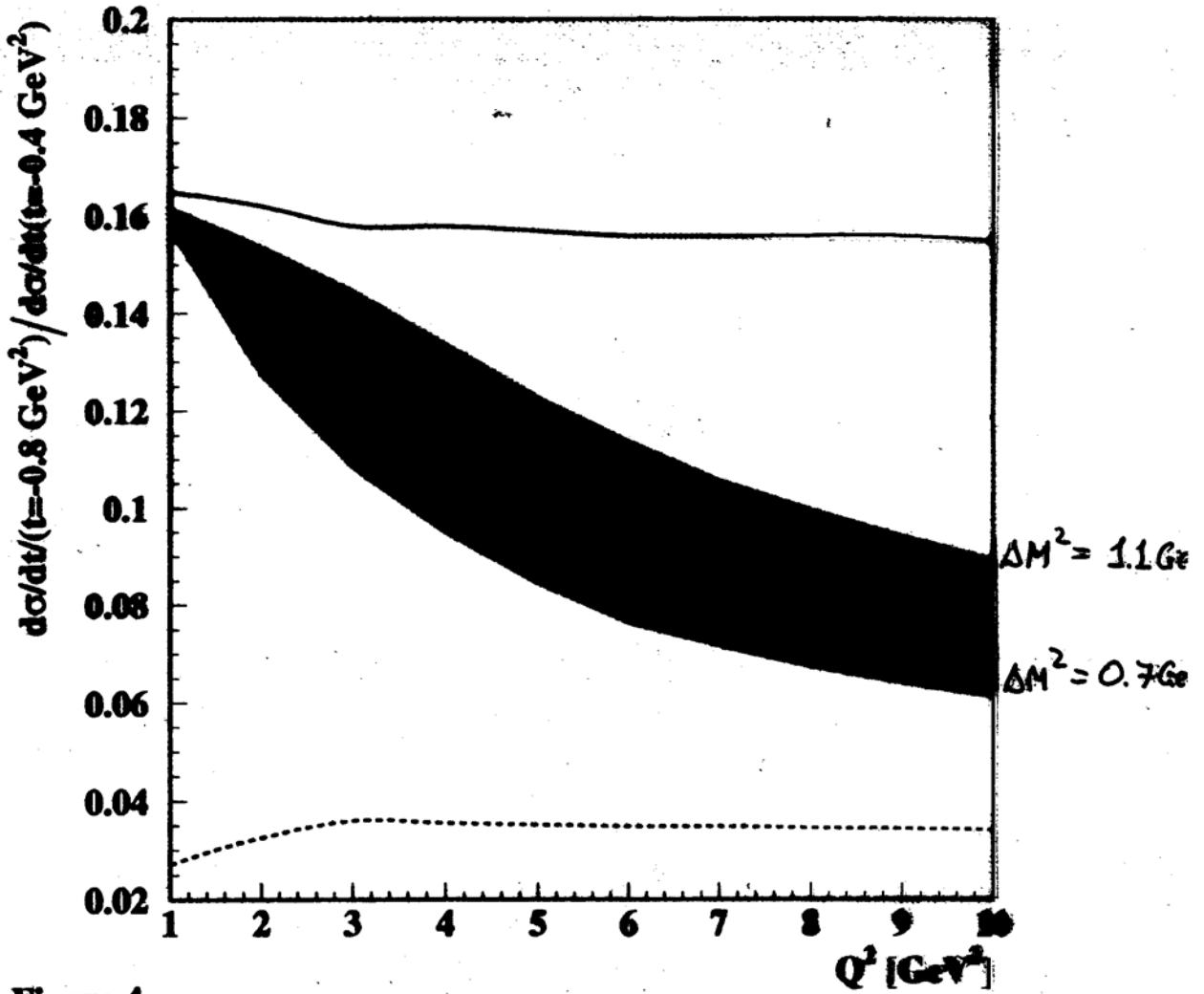
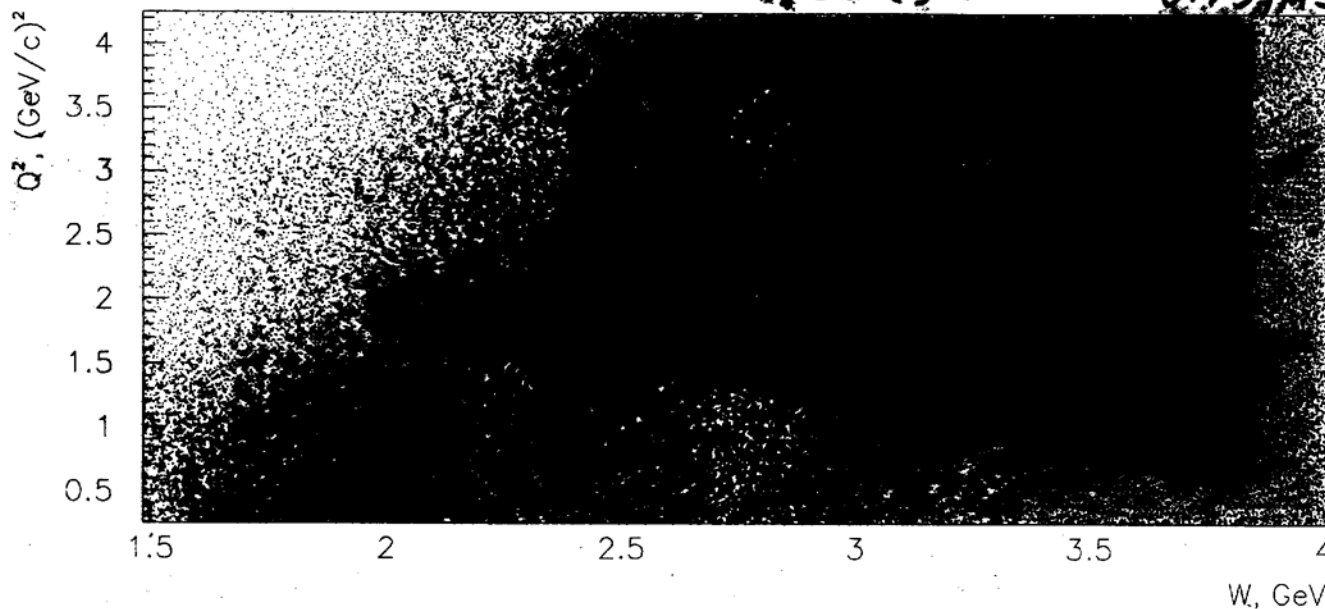


Figure 4

ed=ed $\rho^0, \varphi$  at  $E_e = 11.5$  GeV (cuts for  $Z_T = -100$ .cm,  $I = \pm 3375$ A)

$l_h = 0.5$  fm

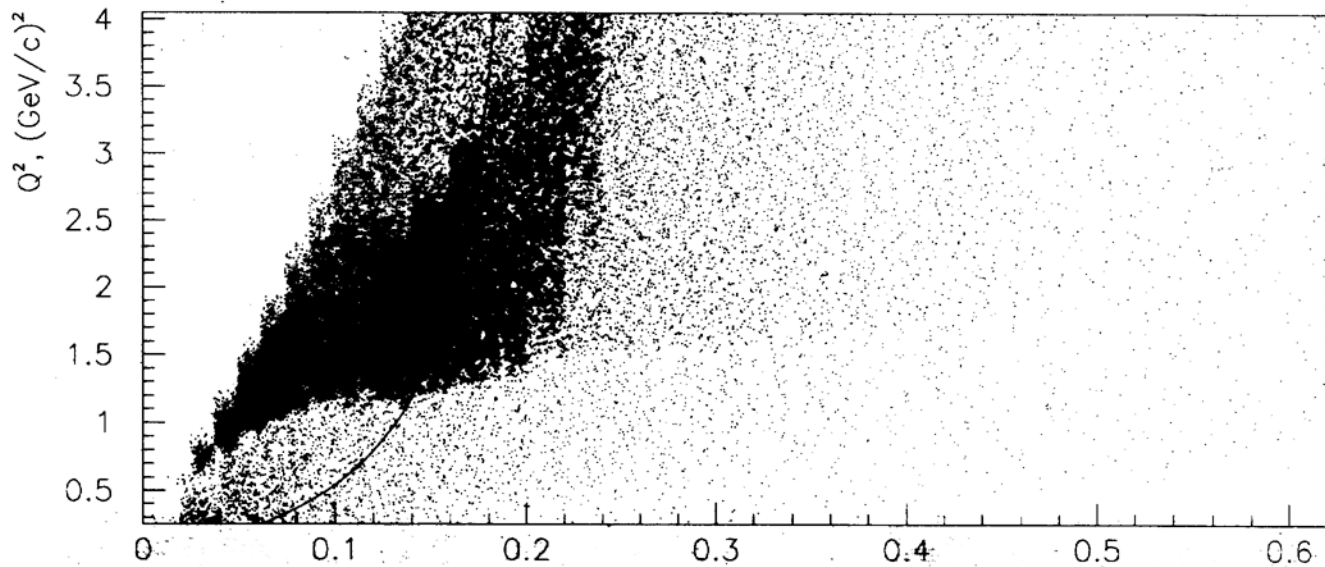
$0.75$  fm  $\leq l_h \leq 1.25$  fm



$l_h = 1$  fm

$l_h = 1.5$  fm  
+ 3375 A

+ 3375 A



**Problem:** Coherence length at Jlab is rather small – need to use the lightest nuclei

The simplest option - study double scattering for  $\gamma^* + d \rightarrow \rho + d$

Look for disappearance of double scattering for  $|t| \geq 0.4 \text{GeV}^2$

Sargsian talk: significant effect for  $Q^2 \geq 3 \text{GeV}^2$

Stepanyan's talk - CLAS has good acceptance, counting rates are reasonable for  $Q^2 \leq 3 \text{GeV}^2$

*Further studies are planned*

**In  $\gamma + n \rightarrow \pi^- + p$  at large  $t$ , the  $s$  dependence consistent with quark counting rules. H.Gao suggested to study  $\gamma + A \rightarrow \pi^- + p + (A - 1)$ ,  $\gamma + d \rightarrow \pi^- p + \text{recoil } p$  to check CT in these processes - This interesting option needs further studies of the missing mass resolution.**



Short-range correlations in nuclei.

Expectation: Large probability

$\sim 20-25\%$  for  $K > K_F$  &  $A=200$

$\sim 50\%$  kinetic energy

Dominant contribution: two-nucleon correlations.

— drops of cold superdense nuclear matter

Quark-gluon structure? Maybe rather

complicated: hints from EMC effect  $\frac{\bar{q}_A}{\bar{q}_N} < 1$

Nuclear forces: phenomenological potentials work  
but microscopic dynamics not clear

meson models of nuclear forces gross problem

for  $r_{NN} \leq 1.2$  fm.

Are quark & gluon degrees of freedom

like quark interchange necessary?

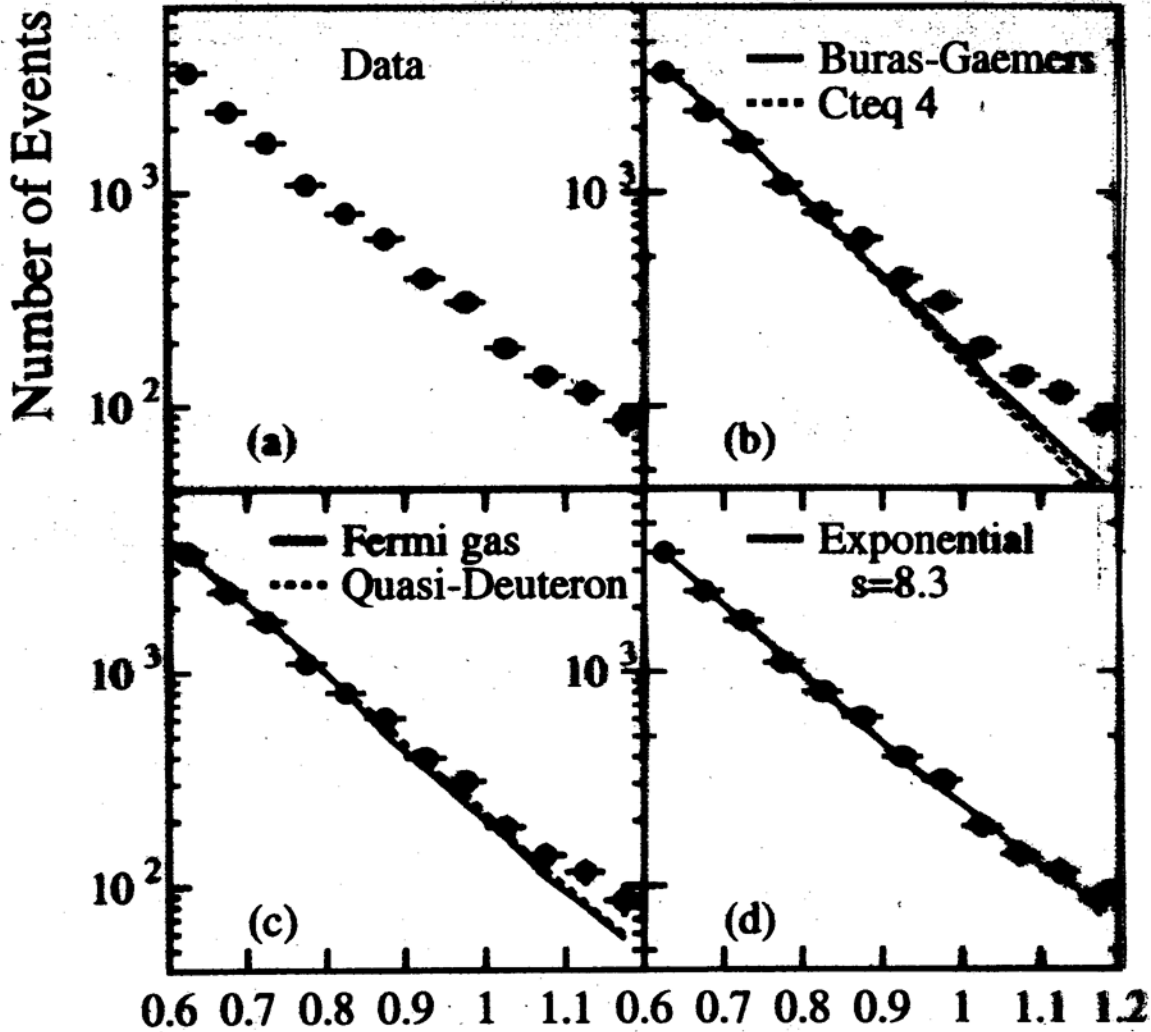
Current best evidence

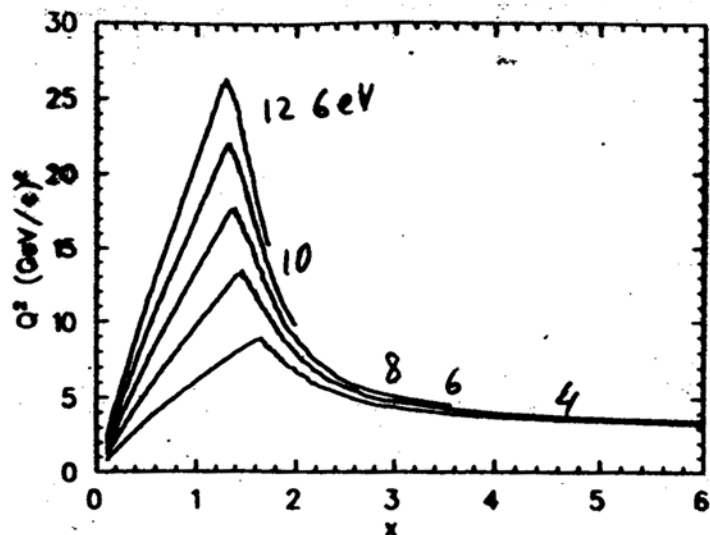
$(e, e')$  at  $x > 1$ ,  $Q^2 \geq 26 \text{ eV}^2$  (D. Day talk)

& new CCFR  $\nu \text{Fe} \rightarrow \mu + X$  data  
in DIS for  $x > 1$ .

JLab with HMS at  $12 \text{ GeV}$  can reach  
 $x \sim 1.5$  and cover  $Q^2 = 10 - 20 \text{ GeV}^2$

to prove scaling





**FIG. 2.** Kinematic coverage  $x$  vs  $Q^2$  for inclusive scattering from nuclear targets. Lower curve is with a 4 GeV CEBAF beam and in increasing order 6, 8, 10, and 12 GeV beam energies.

Consistent numbers for the probability  
of  $a_2$  2-nucleon correlations

	$a_2$	$\nu/A + N$	$p+A \rightarrow$ backward p, $\pi$ , X
	(e, e')		
$^3\text{He}$	$1.7 \pm 0.3$		
$^4\text{He}$	$3.3 \pm 0.5$		$\approx 4$
$^{12}\text{C}$	$5.0 \pm 0.5$		5-6
$^{20}\text{Ne}$		$6 \pm 2$	
$^{27}\text{Al}$	$5.3 \pm 0.6$		
$^{56}\text{Fe}$	$5.2 \pm 0.9$		
$^{197}\text{Au}$	$4.8 \pm 0.7$		

Inf. nuclear matter  $\approx 5$  (Fantoni, et al)

$a_2 \approx 5$  corresponds to

$\approx (25 \pm 5)\%$  nucleons above

Fermi surface.

# Recent developments

BNL E850 observed strong correlation

in  $pA \rightarrow p + (A-1)$   
 $\downarrow$   
 backward nucleon (neutron)

between removal of fast forward  
 proton and emission of backward  
 neutron (Note pn correlations  
 at  $k \sim 500 \text{ MeV}/c \Rightarrow$  pp correlations)

From  $\int P_A^N(\alpha, p_t) d p_t \propto \exp(-b\alpha) (\alpha > 1)$

$$F_{2A}(x, Q^2) \propto \frac{1}{x^3} \exp(-b\alpha) \approx \exp(-1)$$

$B \approx b+1$  (modulus EMC effect correction)

$b \approx 7-7.5 \rightarrow B \sim 8-8.5$  FS 79

CCFR data  $B \approx 8.2$  at  $x=1.2-1.4$   
 $Q^2 \sim 16$  12.6 eV scaling region

# **Active Target**

- **30 cm of 10 Atm Deuterium Gas**
  - 100 nA →  $L=10^{34}$
- **Sweeping Magnetic Field to Bottle Up Mollers**
- **Need Good Timing to Reduce Randoms - 10 ns**
- **GEM - Gas Electron Multiplier**
  - Fast
  - Low Density
  - Developed at CERN
  - Good Resolution

Fundamental question: how different  
quark distribution in bound and free nucleon.

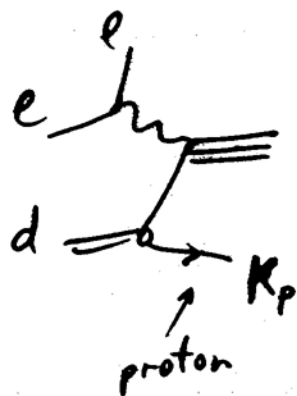
Simplest case: deuteron

Need to know the answer to measure

$$F_{2n}(x, Q^2) / F_{2p}(x, Q^2) \text{ at } x \sim 1$$

In  $F_{2D}(x, Q^2)$  large corrections due to  
the EMC effect FS 85

⇒ Two directions: @ tagging for small  
momenta  $k_p$



& do Chew & Low  
extrapolation to the pole

$$\tilde{m}^2 - m^2 = (p_d - k_p)^2 - m^2 \approx -(k^2 + m_n^2)$$

R. Ent Active target: tagging  
down to  $p_N = 80 \text{ MeV}/c$  !!



① Tagging ~~for~~ ~~for~~  $\geq 200$  MeV

"gold plated test of  
the models of the EMC  
effect"

FS 85"

Doable with CLAS K. Griffioen ~~talk~~

$$\sigma \propto \psi_D^2(\alpha) F_{2n}^{\text{bound}}\left(\frac{x}{z-\alpha}, k\right)$$

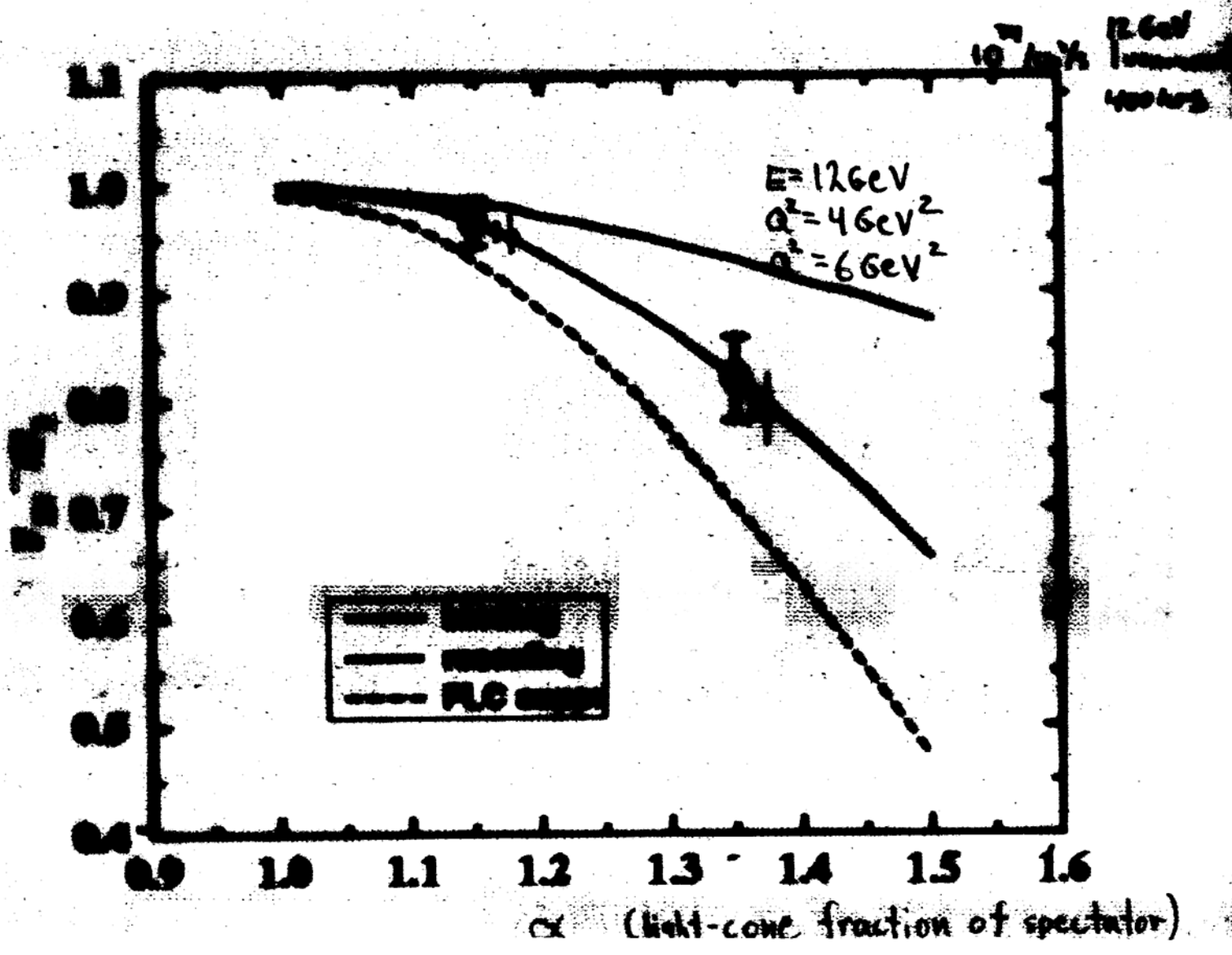
$$\frac{F_{2n}^{\text{bound}}\left(\frac{x}{z-\alpha}, k\right)}{F_{2n}\left(\frac{x}{z-\alpha}\right)} \neq 1$$

$$\alpha = \frac{2}{m_D} \left( E_N - \frac{\vec{p}_N \cdot \vec{q}}{|\vec{q}|} \right)$$

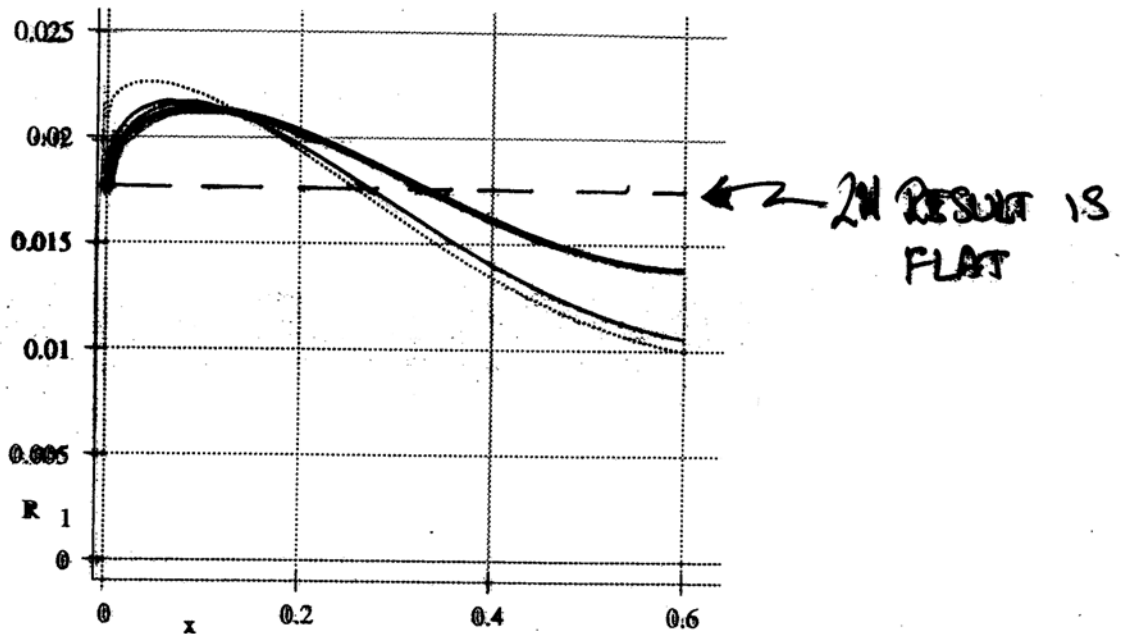
$\alpha > 1$   
nucleon flies backward

6 GeV  $Q^2 = 2 \text{ GeV}^2$  } expected data  
 $x \sim 0.5$  - for one  $Q^2$  interval

$Q^2 = 5 \text{ GeV}^2$  } calculations  
 $x \sim 0.6$



Large effects in 6q model C. Carlson



## Putative $R_1$

- Simple old quark distributions for  $F_{2n}$
- LS for 6q: heavy is their A, normal B, dotted C.
- $\alpha = 1.4$

for fixed

$$E = E_{in}$$

**Will it work at CEBAF?**

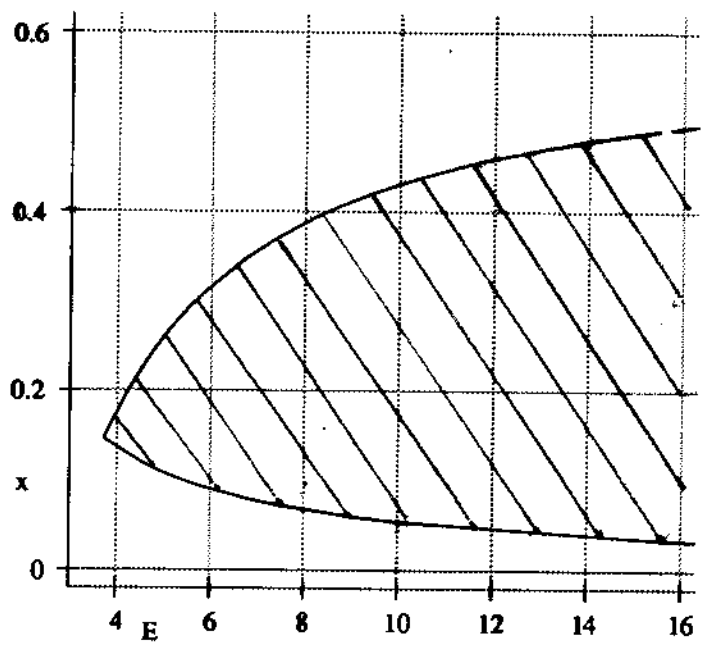
**Yes — if we are in the scaling region. Means**

♠  $Q^2 > 1 \text{ GeV}^2$ , setting lower limit to  $x$   $\left[ x = \frac{Q^2}{2Mv} \right]$

♣  $W > 2 \text{ GeV}$ , setting upper limit to  $x$   $\left[ \frac{1}{x} = 1 + \frac{W^2 - M^2}{Q^2} \right]$

( $W$  is  $\gamma$ -nucleon c.m. energy)

**Gives “scaling window” for CEBAF, say for  $\alpha = 1.4$ ,**

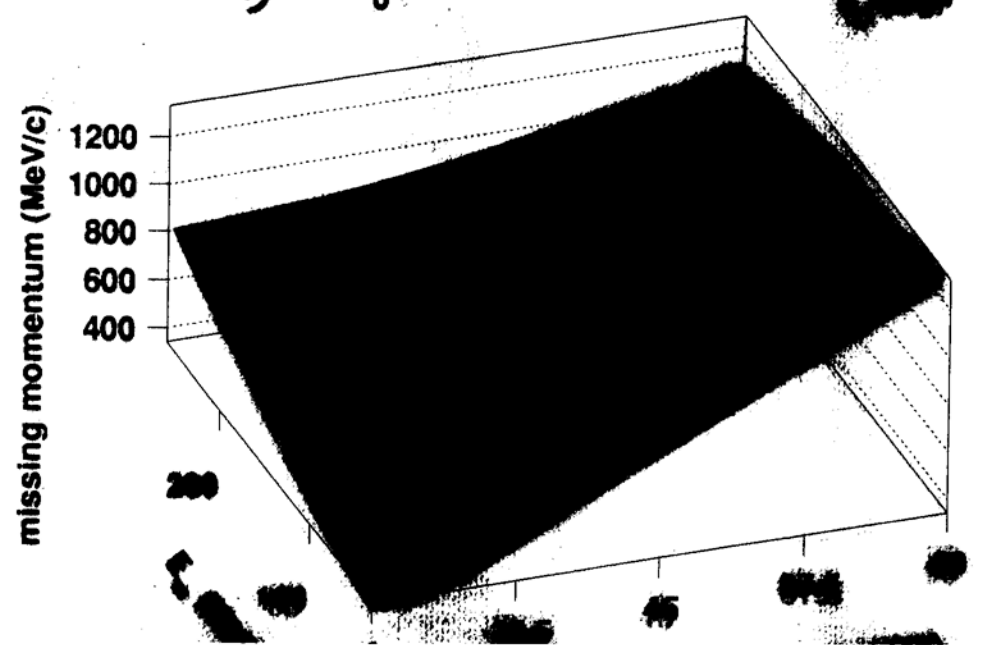
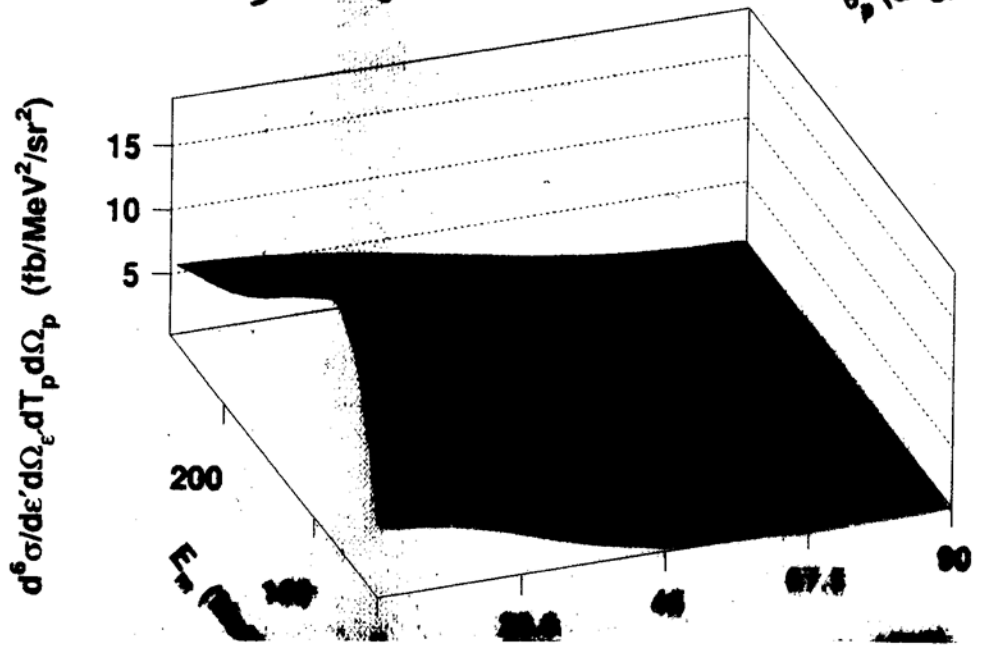
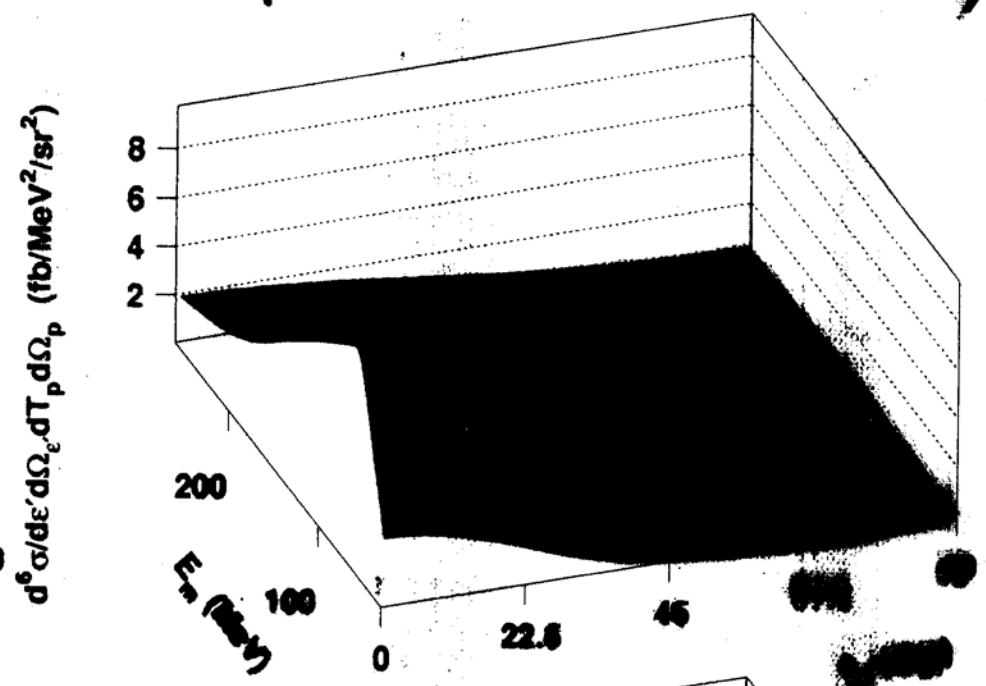
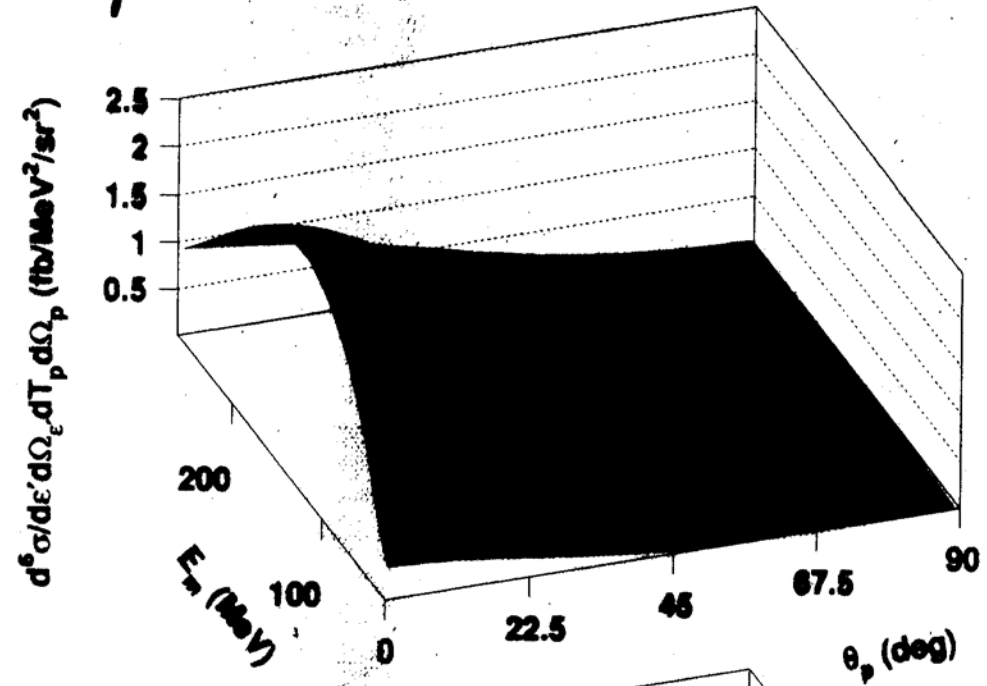


$E$  = energy of incoming electron beam.

Want a bit more than 4 GeV.

$(e, e'p)$   $x \approx 2$   $Q^2 \approx 1.1 (\text{GeV})^2$

$(E = 2.5 \text{ GeV}, \theta_e = \dots)$



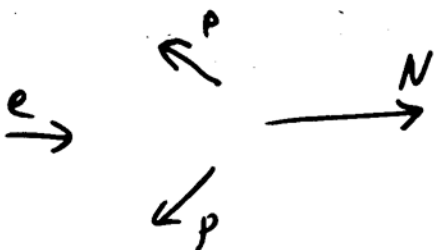
Other important directions:

- ② looking for backward  $\Delta$ 's  
- feasible in CLAS (L. Weinstein)

⇒ Sensitive to pre-mordial  $\Delta$ 's,  
and to knotted  $6q$  states.

- ⊗ Study of  $pp$ ,  $pn$ ,  $nn$  short-range correlations via correlation of backward and forward production was pretty difficult at  $E_e \leq 16\text{eV}$  (Jan Ryckebusch talk)

- ⊗  $3N$  correlations



## Conclusions

\* Decisive tests of dominance of pQCD vs soft physics in nucleon form factors at  $Q^2 = 10 - 15 \text{ GeV}^2$

\*\* Transition from study of bulk properties of short-range correlations to the study of parton structure of correlations, determining the dynamic origin of the EMC effect.