

Hypernuclear Physics with Electromagnetic and hadronic Probes

Radhey Shyam, Saha Institute of Nuclear Physics, Kolkata, India

1. Introduction

A brief review and comparison of production reactions

2. Brief sketch of the theoretical model

3. Results, cross sections, spectroscopy

4. Conclusions

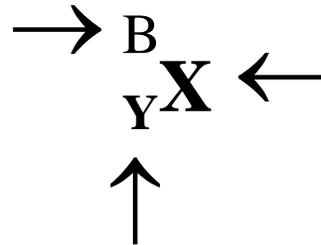
Seminar at Jlab, Oct. 27, 2008

Hypernuclei are systems where one nucleon is replaced by a hyperon

Λ Hypernuclei

${}^A_{\Lambda}Z$ is a bound state of Z protons ($A-Z-1$) neutrons and a Λ hyperon

Number of
Baryons
($N+Z+Y$)



Element
=
Total charge

Number of
hyperons

Double Hypernuclei, ${}^{10}_{\Lambda\Lambda}\text{Be}$, Σ Hypernuclei

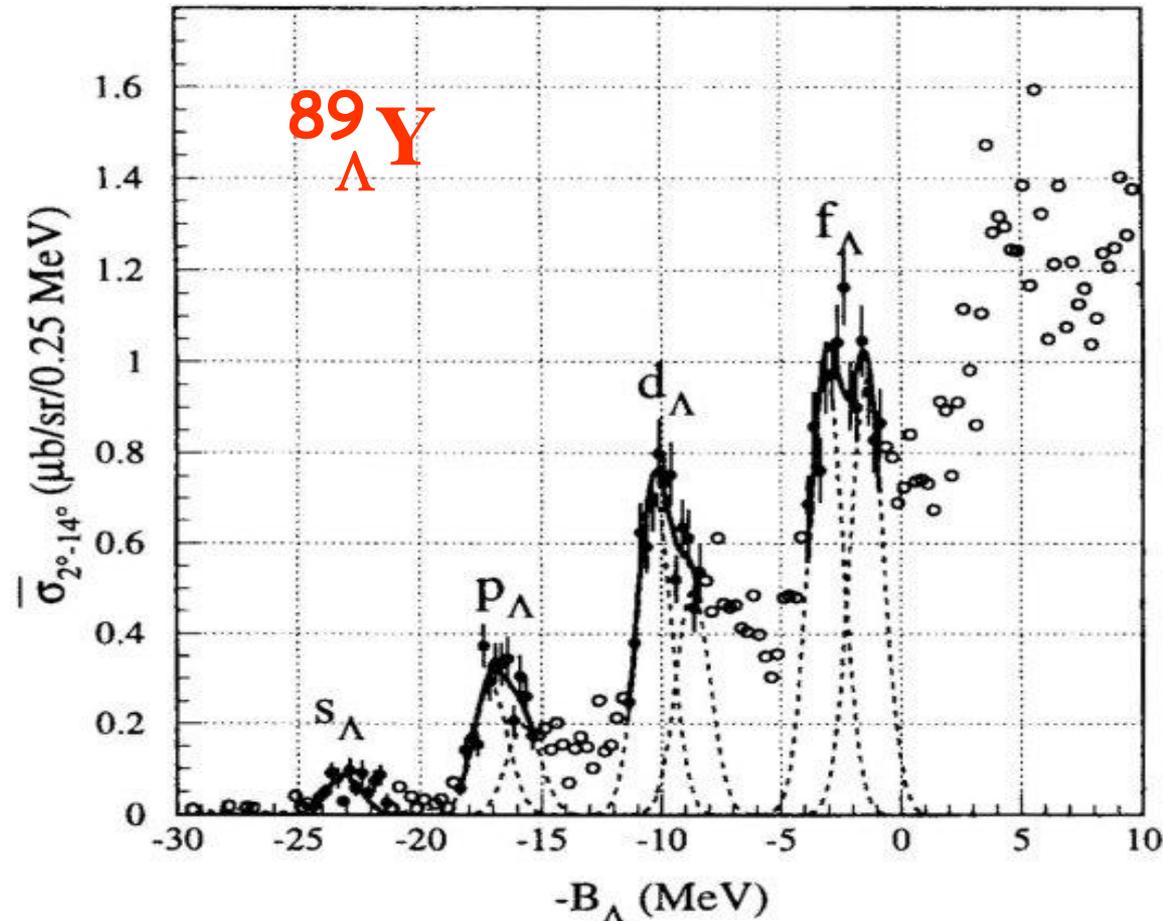
Why are Hypernuclei interesting!

New type of nuclear matter, new symmetries, New selection rules. First kind of flavored nuclei.

Hyperons are free from Pauli principle restrictions

Can occupy quantum states already filled up with nucleons

Good probe for deeply bound single particle states.

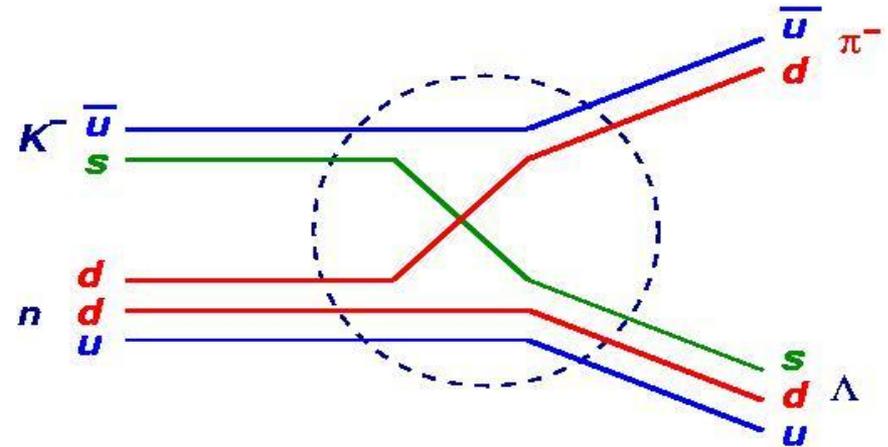


Production of Λ Hypernuclei

MESONIC PROBES

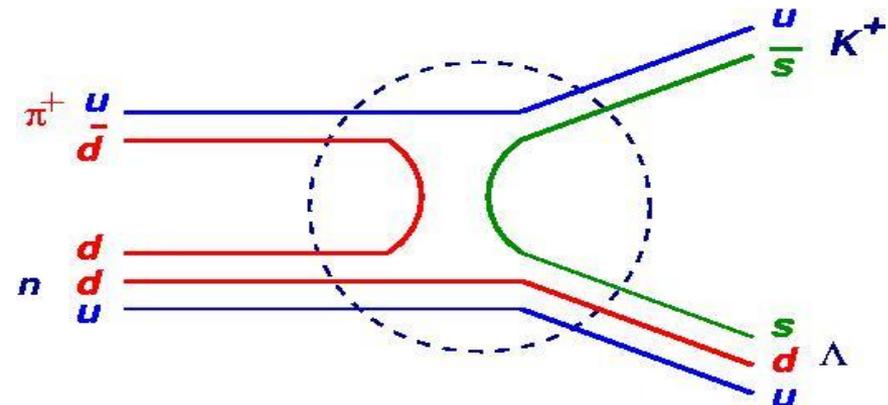
Prog. Part. Nucl. Phys. 57, 564 (2006)

(K^-, π^-) reaction



Strangeness exchange

(π^+, K^+) reaction



Associated strangeness production

Electromagnetic and baryonic Probes

(γ, K^+) reaction



H. Yamazaki et al., Phys. Rev. C 52, R1157 (1995)

($e, e' K^+$) reaction



L. Yuan et al. Phys. Rev. C 73, 044607 (2006)

M. Iodice et al, Phys. Rev. Lett. 99, 052501 (2007)

F. Cusanno et al., arXiv: 0810:3853

JLab

(p, K^+) reaction



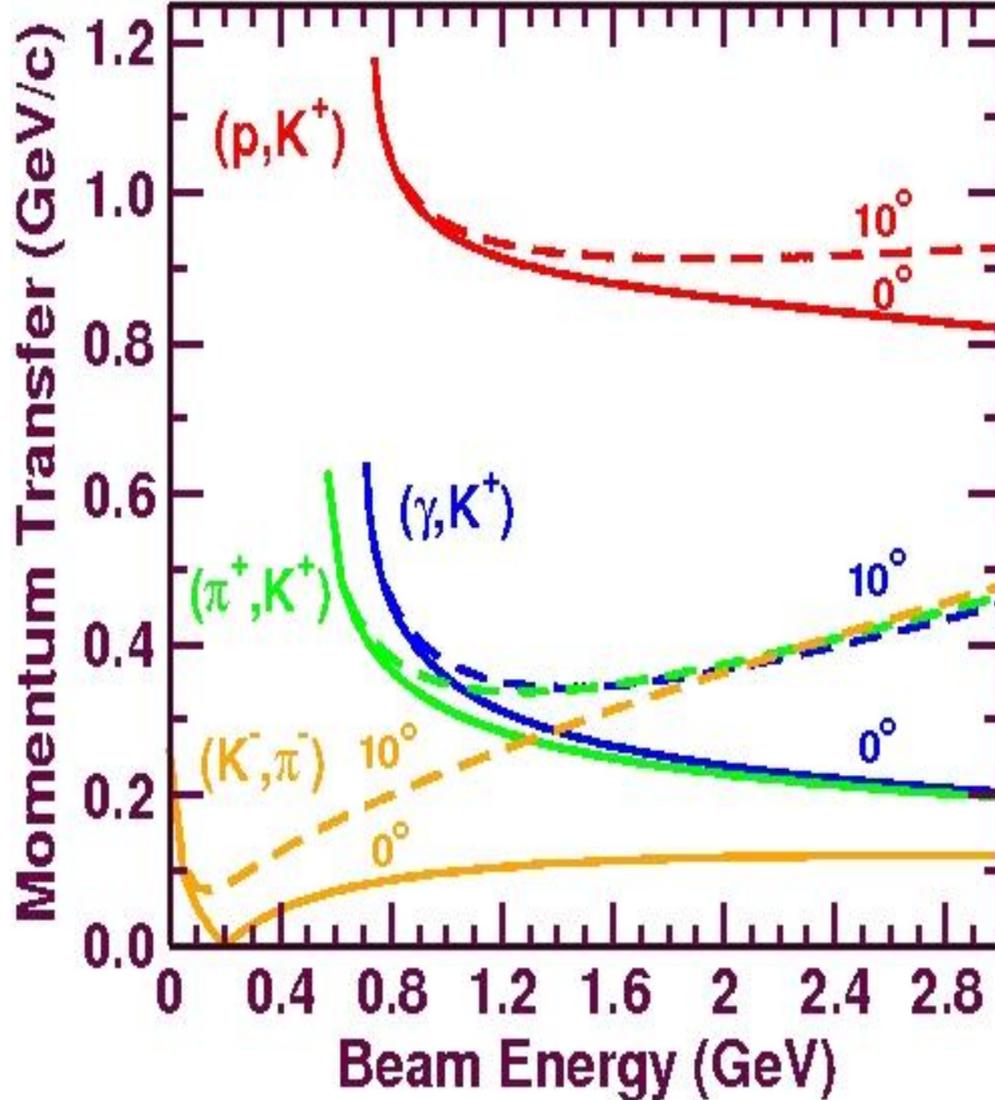
Heavy Ion reactions



GSI

^{12}C

KINEMATICS



(p, K⁺) reaction

Large Momentum transfers

(π⁺, K⁺) and (γ, K⁺) reactions

Momentum transfer > p_F

(K⁻, π⁻) reaction

Low momentum transfer

Excitation Energy Spectra

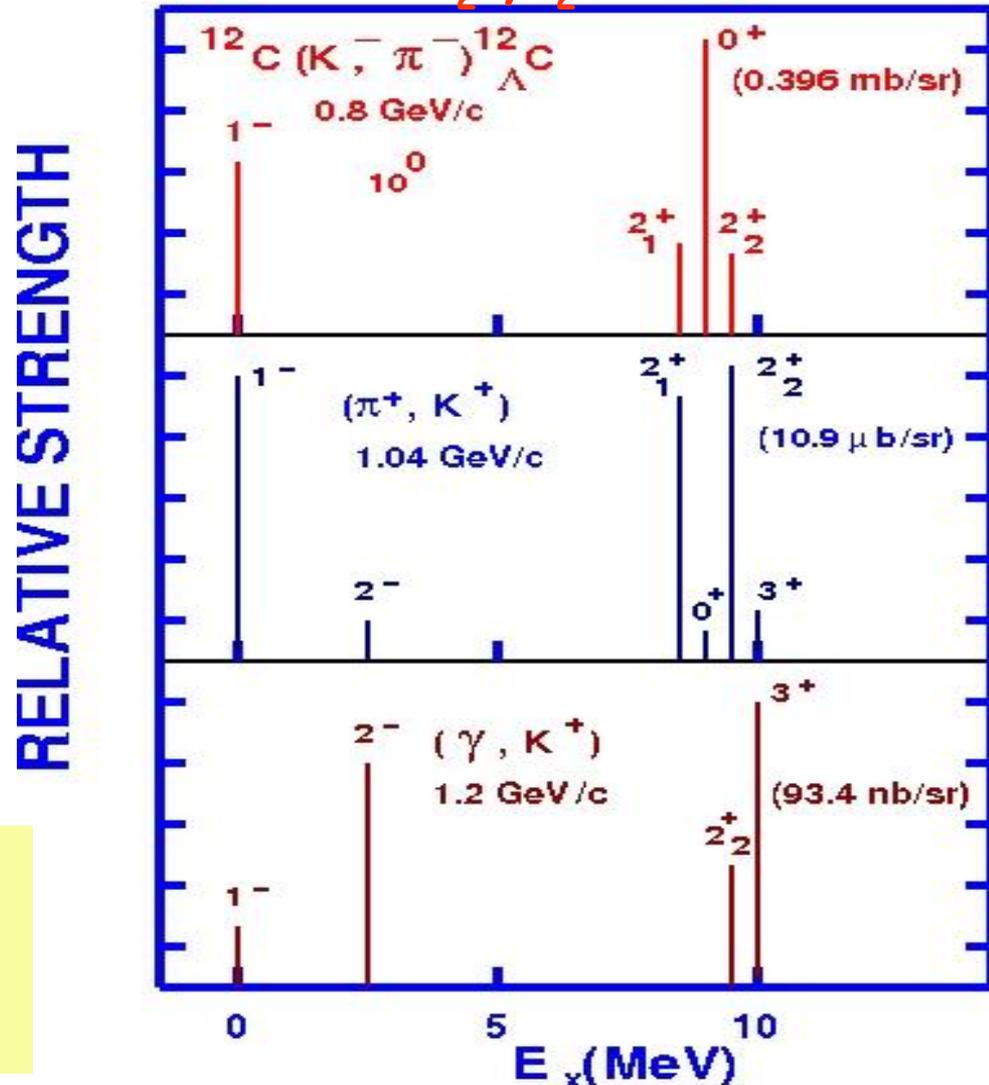
GS: $(p_{3/2}^{-1}, s_{1/2}^{\wedge}) 1^-, 2^-$ $(p_{3/2}^{-1}, p_{3/2}^{\wedge}) 1_1^+, 2_1^+, 3^+$ $(p_{3/2}^{-1}, p_{1/2}^{\wedge}) 1_2^+, 2_2^+$

(K^-, π^-) substitutional
 0^+ state dominates

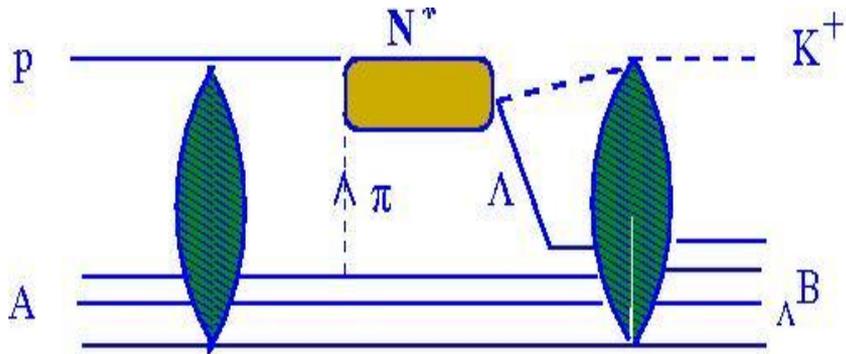
(π^+, K^+) nonsubstitutional
 1^- and 2^+ states dominate

(γ, K^+) unnatural parity
 states dominate

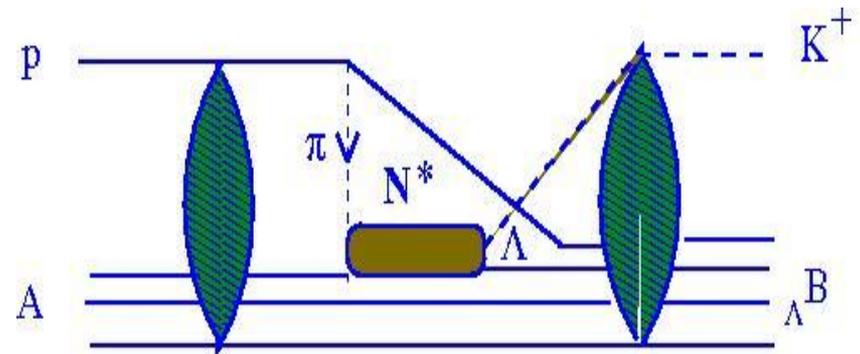
(γ, K^+) and $(e, e'K^+)$ reactions
 can also excite unnatural
 parity stretched states



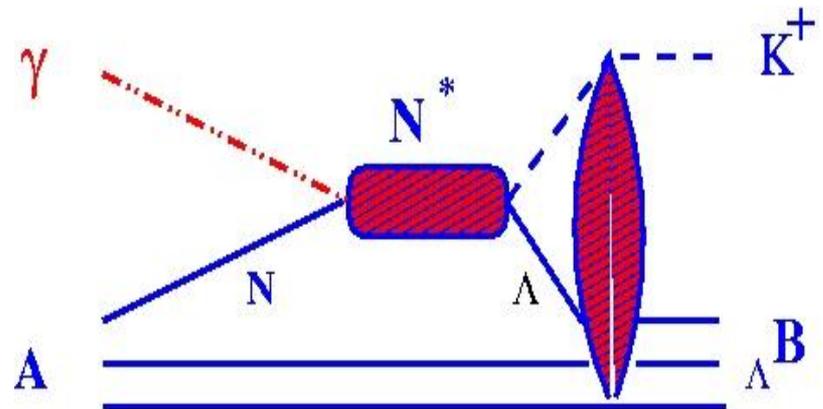
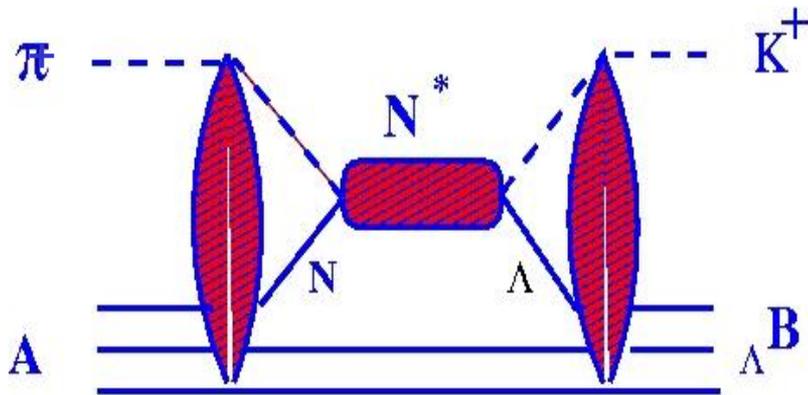
Production processes for various reactions



Target emission



Projectile emission



$N^*(1650)$, $N^*(1710)$, $N^*(1720)$ baryonic resonances.

A Covariant Description of $A(h\nu, K^+)_{\Lambda}B$ reaction

Effective Lagrangians at various vertices

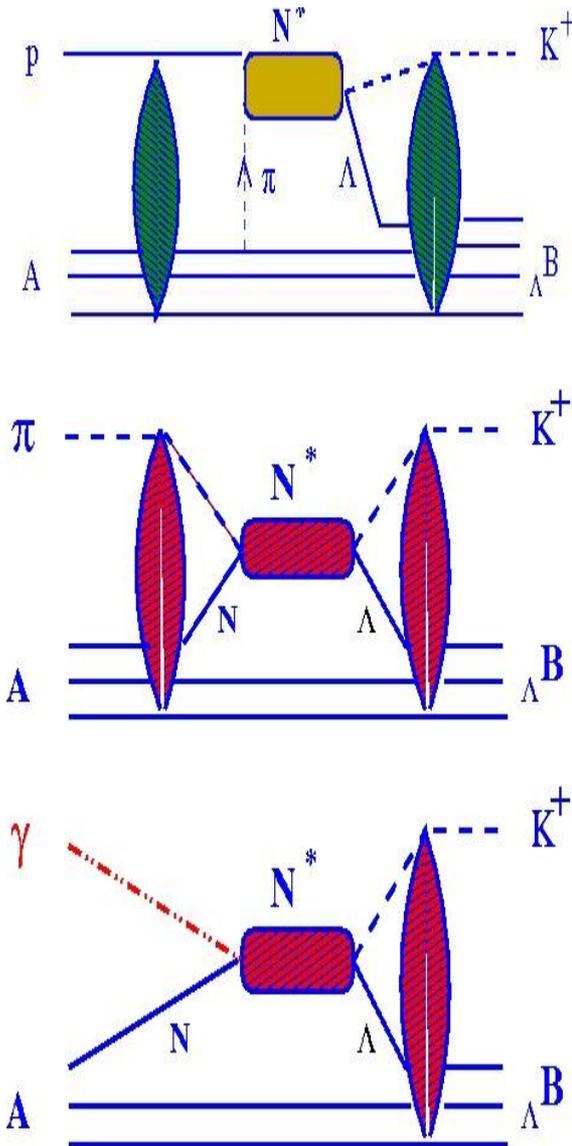
Coupling constants, form-factors

Bound state nucleon and hyperon spinors

Initial and final state interactions
(distorted waves).

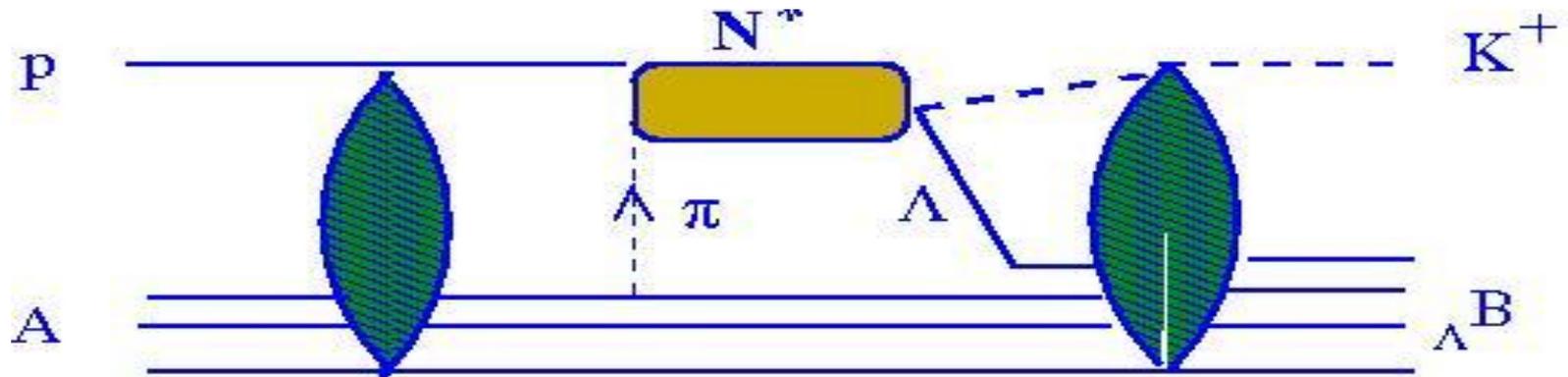
Medium modification of N^* (also of
intermediate mesons in proton induced
reactions) self energies.

All calculations in momentum space,
so nonlocalities are included.

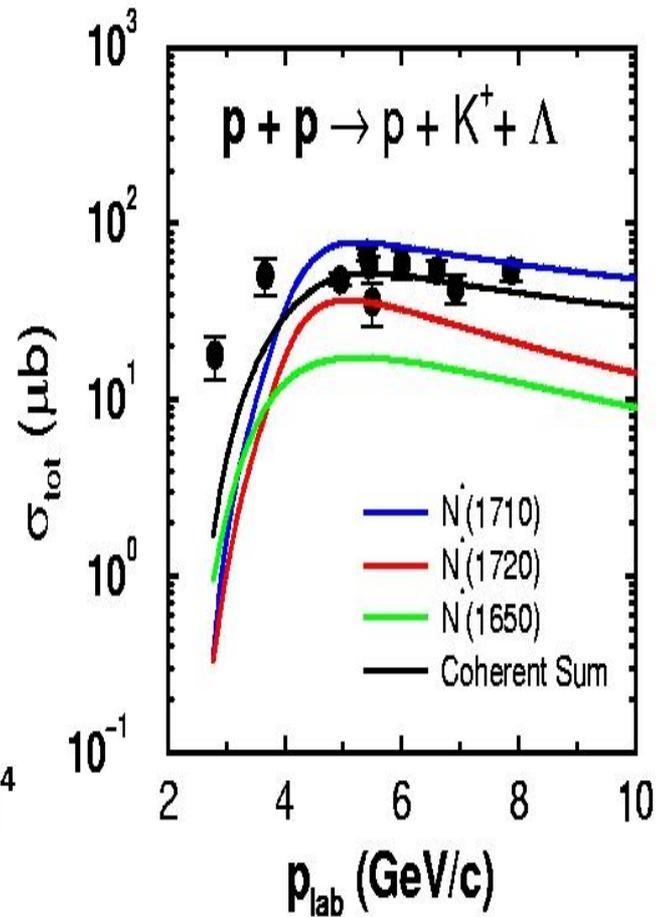
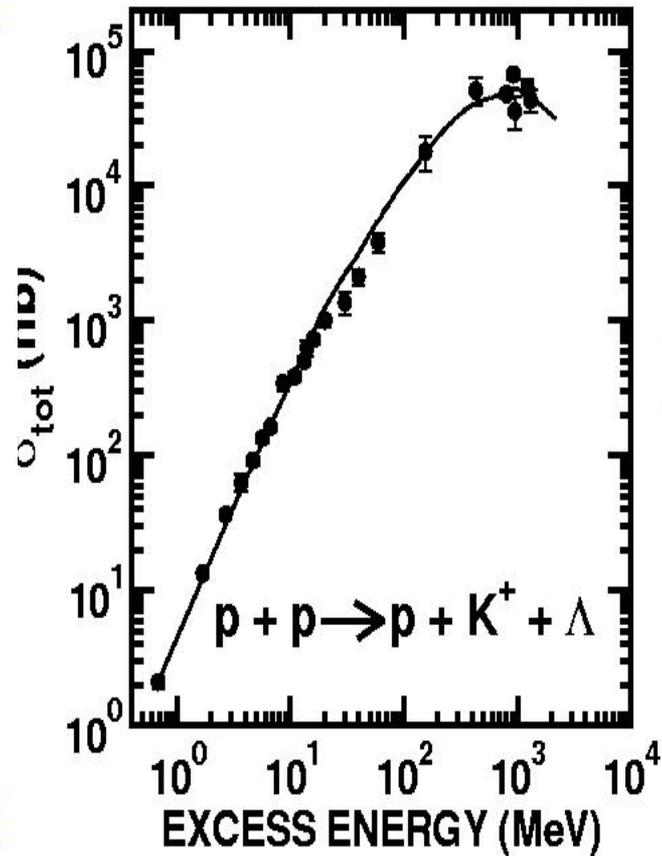
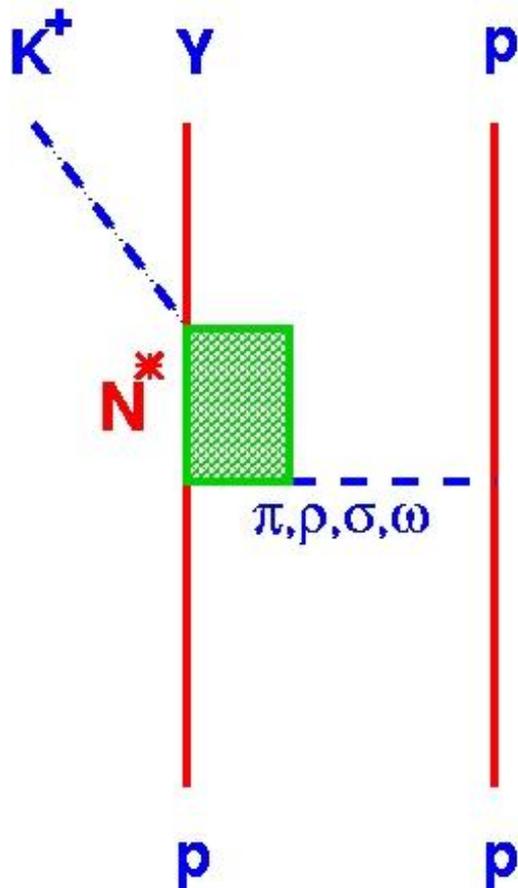


A typical amplitude

$$\begin{aligned}
 M_{2b}(N_{1/2}^*) &= C_{iso}^{2b} \left(\frac{g_{NN\pi}}{2m_N} \right) (g_{N_{1/2}^* N \pi}) (g_{N_{1/2}^* \Lambda K^+}) \bar{\psi}(p_2) \gamma_5 \gamma_\mu q^\mu \\
 &\times \psi(p_1) D_\pi(q) \bar{\psi}(p_\Lambda) \gamma_5 D_{N_{1/2}^*}(p_{N^*}) \gamma_5 \\
 &\times \Phi_K^{(-)*}(p'_K, p_K) \Psi_i^{(+)}(p'_i, p_i),
 \end{aligned}$$

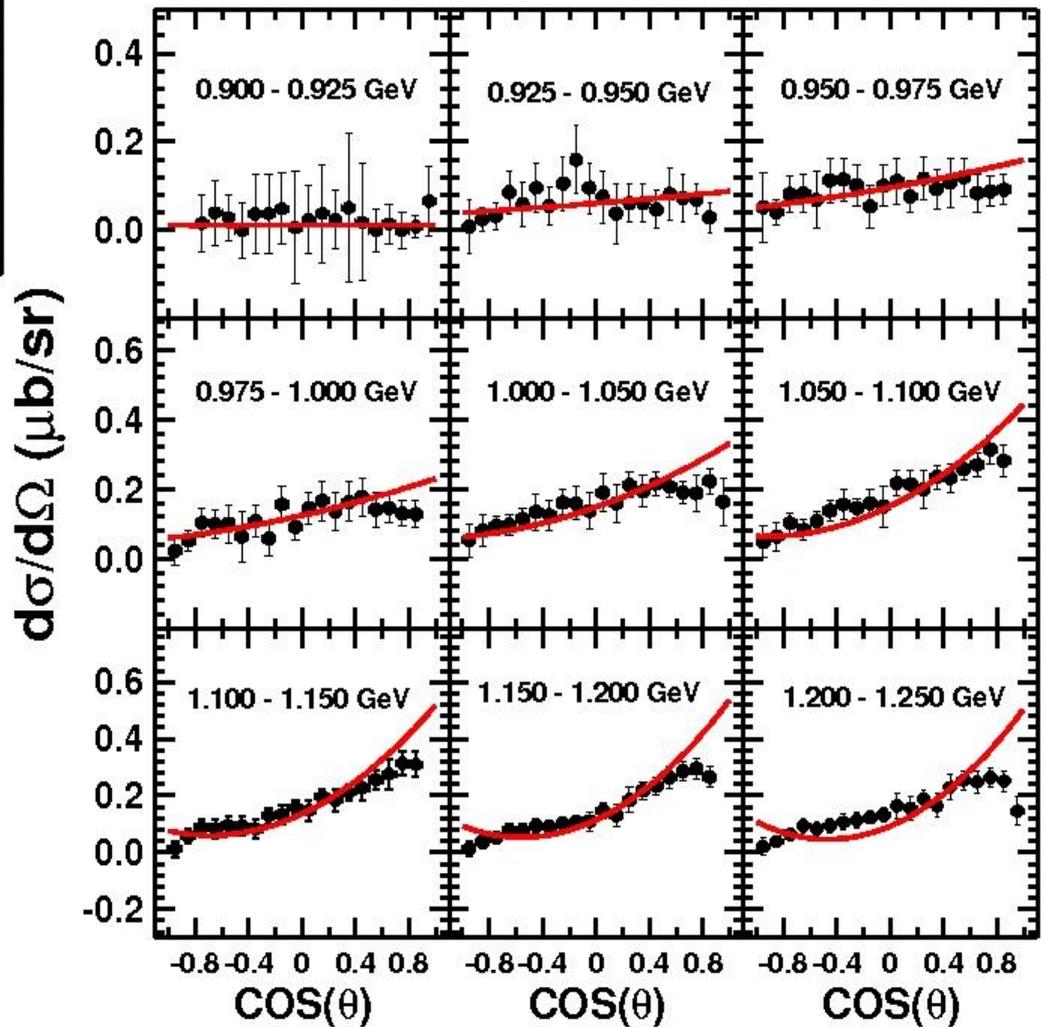
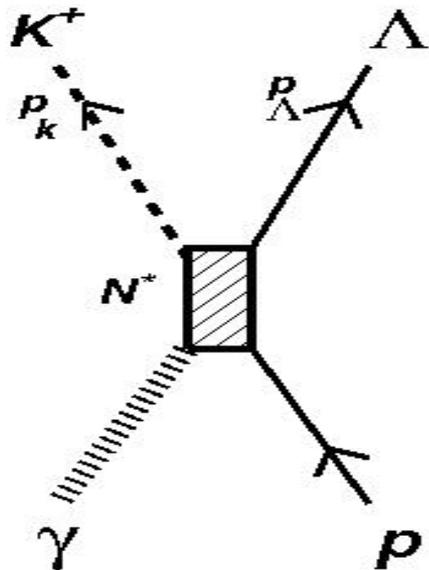
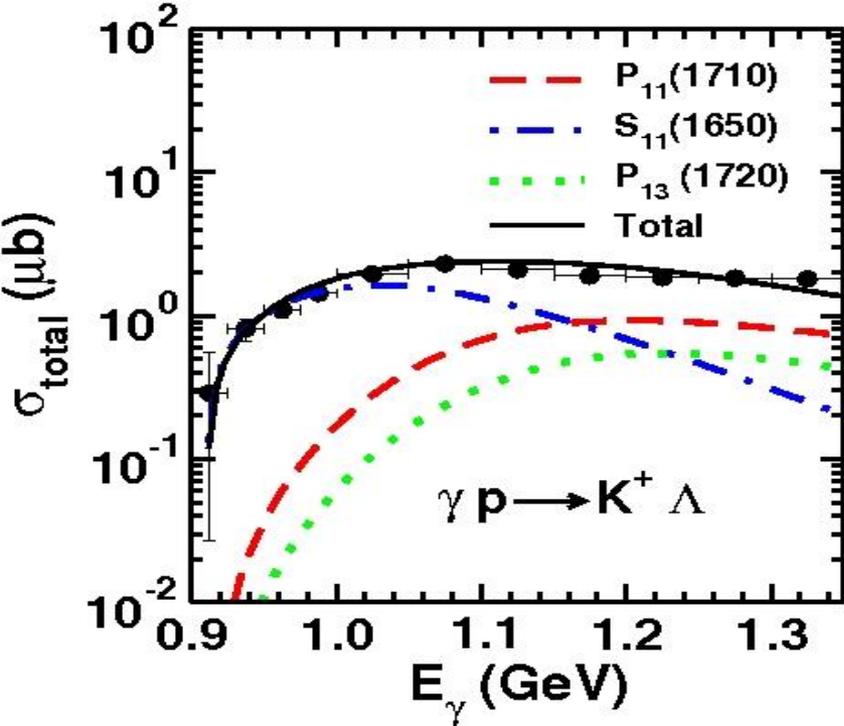


$pp \rightarrow p\Lambda K^+$ reaction



$p(\gamma, K^+) \Lambda$ reaction

RS, K. Tsushima, A.W. Thomas, In prep



Bound state spinors

A mean field approach

Momentum space Dirac Eq.

$$\not{p}\psi(p) = m_N\psi(p) + F(p),$$

$$F(p) = \delta(p_0 - E) \left[\int d^3p' V_s(-\mathbf{p}')\psi(\mathbf{p} + \mathbf{p}') \right. \\ \left. - \gamma_0 \int d^3p' V_v^0(-\mathbf{p}')\psi(\mathbf{p} + \mathbf{p}') \right].$$

$$\psi(p) = \delta(p_0 - E) \begin{pmatrix} f(k) \mathcal{Y}_{\ell 1/2j}^{m_j}(\hat{\mathbf{p}}) \\ -ig(k) \mathcal{Y}_{\ell' 1/2j}^{m_j}(\hat{\mathbf{p}}) \end{pmatrix},$$

$$F(p) = \delta(p_0 - E) \begin{pmatrix} \zeta(k) \mathcal{Y}_{\ell 1/2j}^{m_j}(\hat{\mathbf{p}}) \\ -i\zeta'(k) \mathcal{Y}_{\ell' 1/2j}^{m_j}(\hat{\mathbf{p}}) \end{pmatrix},$$

$^{16}\text{O} (\gamma, K^+) ^{16}_{\Lambda}\text{N}, ^{16}\text{O} (\gamma^*, K^+) ^{16}_{\Lambda}\text{N}$ Reactions

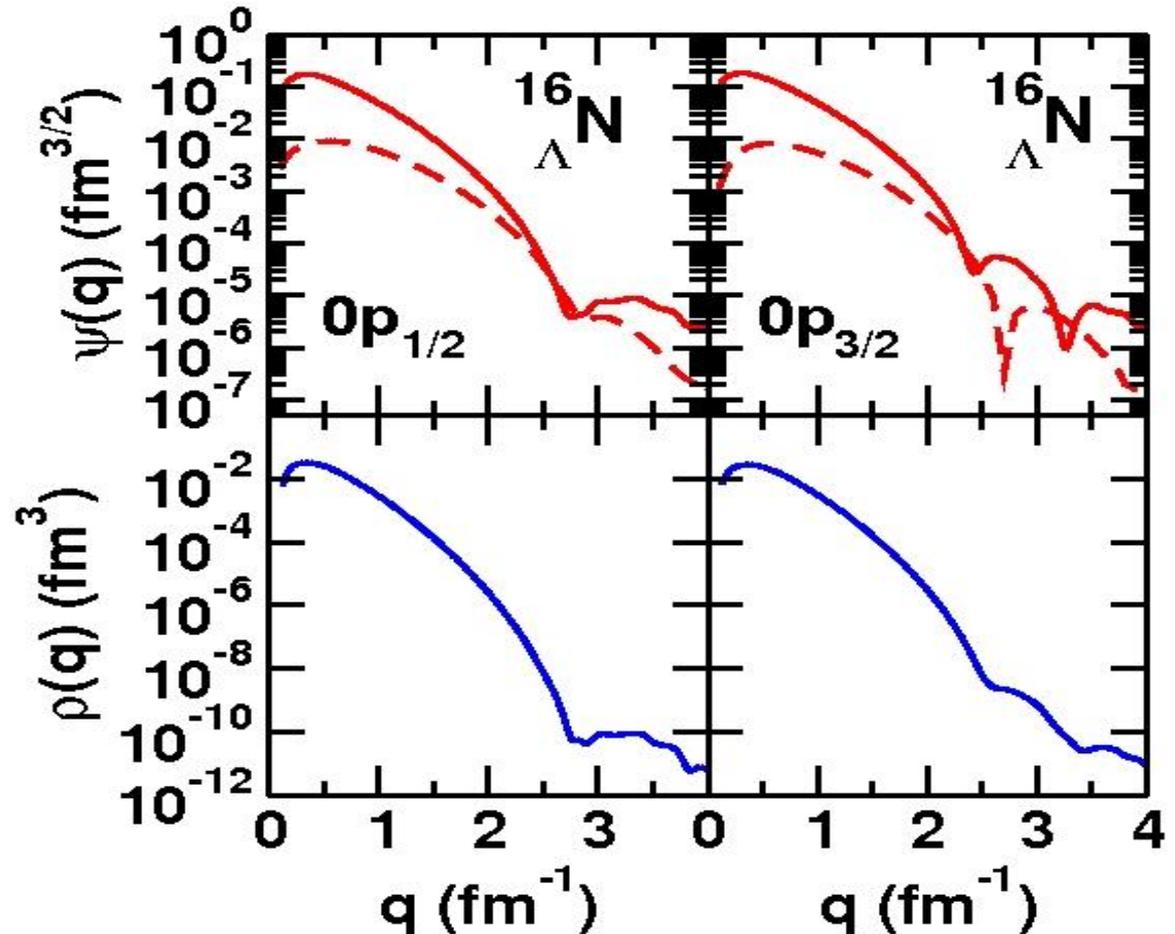
Hole states $B(p_{3/2}) - B(p_{1/2}) = 6.3$ MeV Jlab, arXiv: 0810.3852

$[p_{1/2}^{-1}, s^{\Lambda}], [p_{3/2}^{-1}, s^{\Lambda}], [p_{1/2}^{-1}, p^{\Lambda}], [p_{3/2}^{-1}, p^{\Lambda}]$ $^{16}_{\Lambda}\text{N}$ spectrum

Bound Hypernuclear spinors

Single particle model

It is only for
 $q \ll 1.5 \text{ fm}^{-1}$
 $|\lg(q)| \ll |f(q)|$



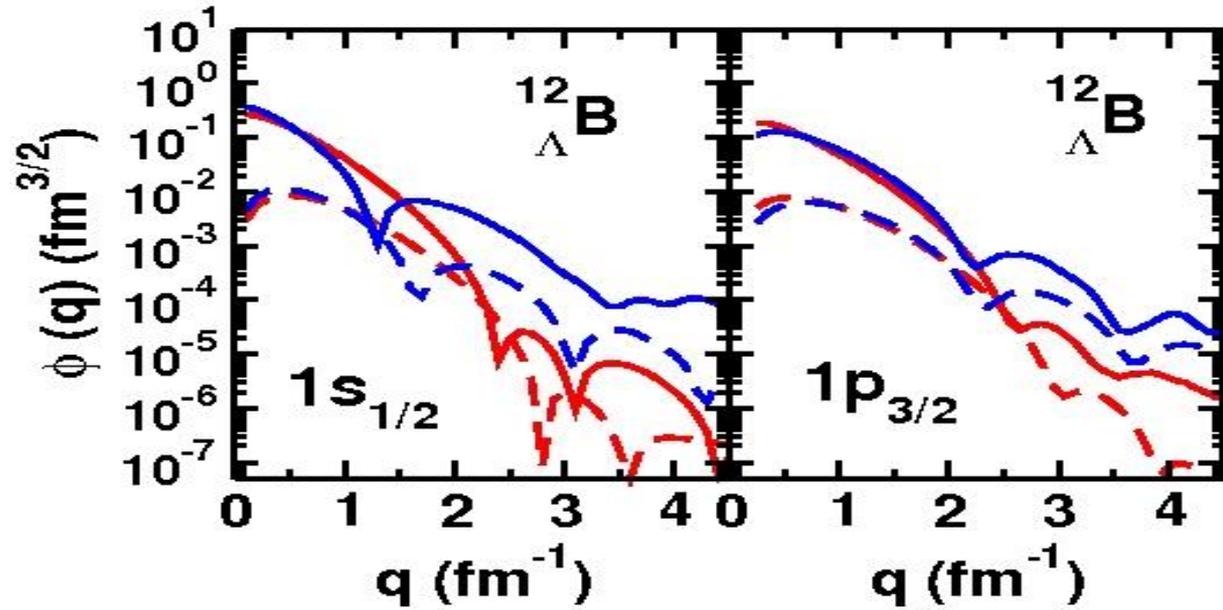
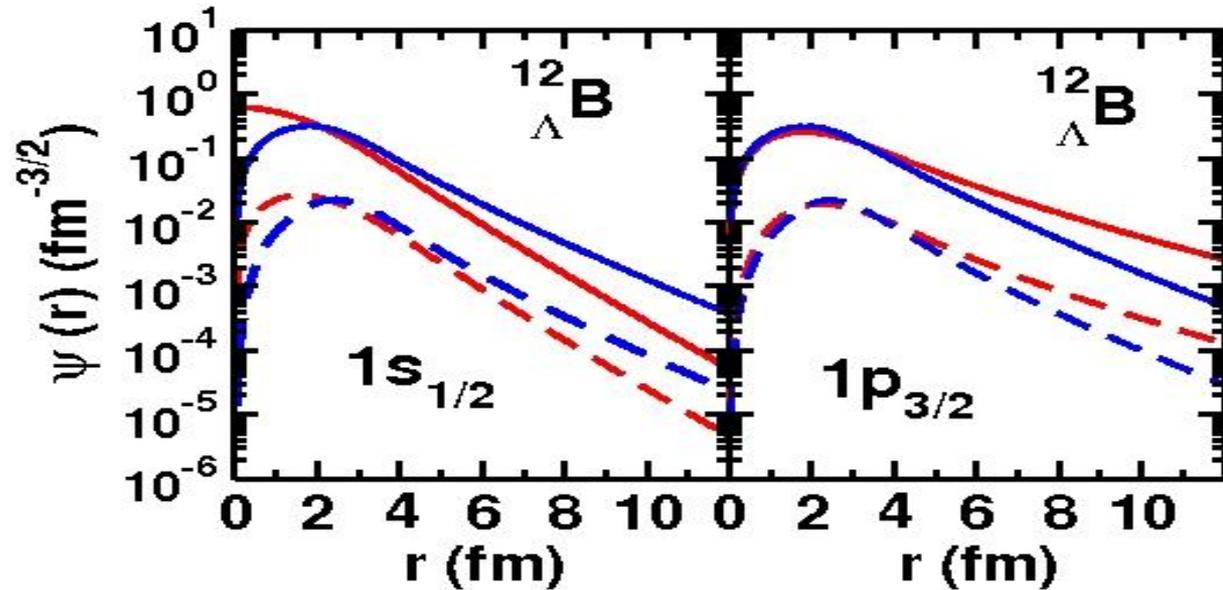
$^{12}\text{C} (\gamma, K^+) ^{12}_{\Lambda}\text{B}$, $^{12}\text{C} (\gamma^*, K^+) ^{12}_{\Lambda}\text{B}$ Reactions

PRL 99, 052501 (2007)

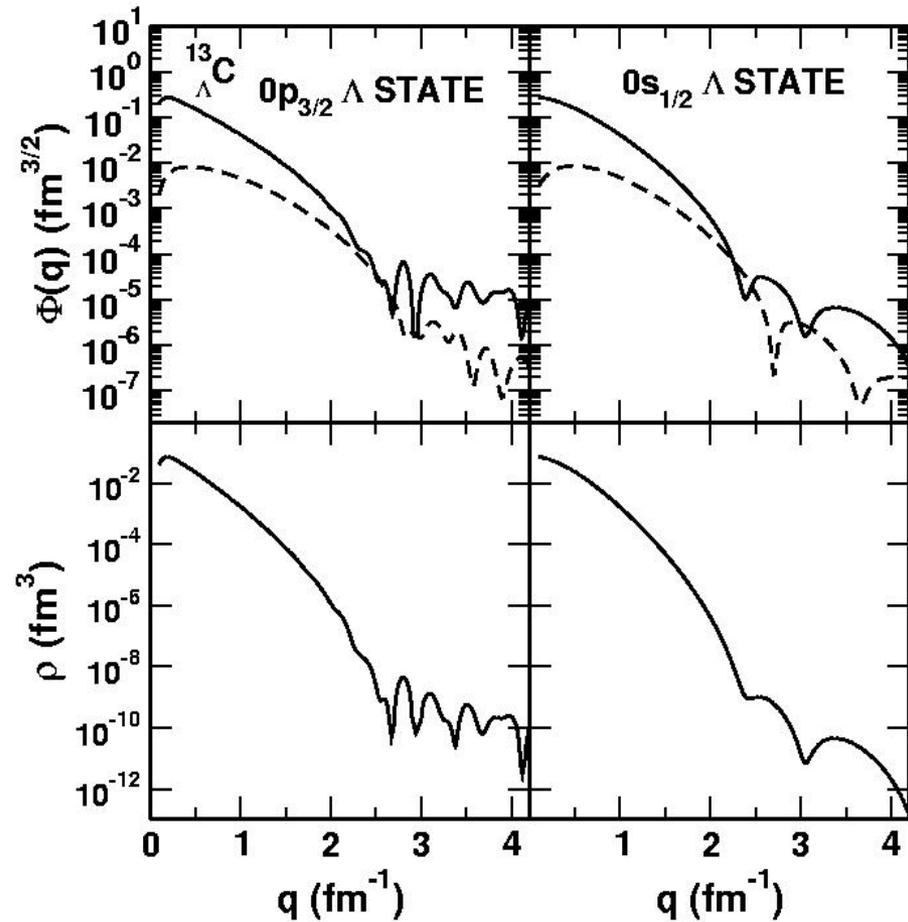
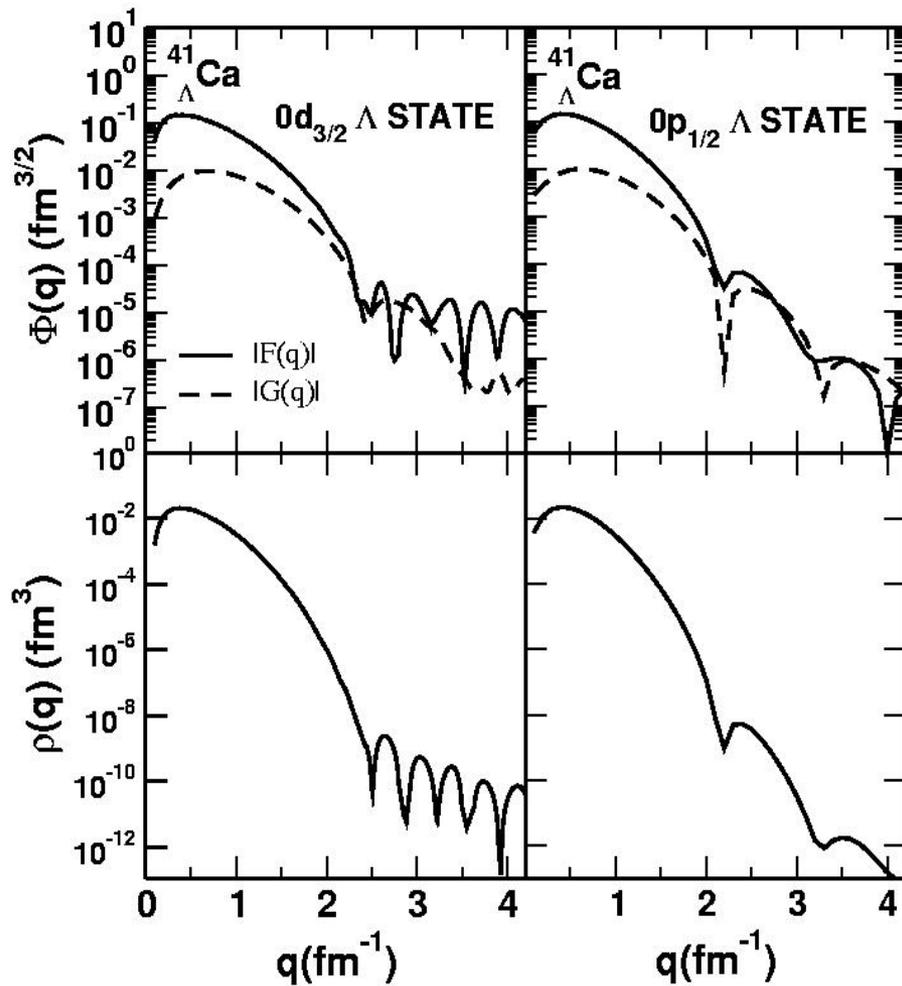
Quark meson
Coupling model

Saito, Tsushima, Thomas,
Prog. Nucl. Part. Phys.
58, 1-167 (2007).

Single Particle
model



Bound Hypernuclear wave spinors



In the region of the momentum transfer of interest, the lower component of the spinor is not negligible.

scattering states

$$\Phi_{pf}(p_K) = \delta(p_K^0 - E_K) \sum_{\ell m} (-)^{\ell} Y_{\ell m}(p_f) Y_{\ell m}^*(p_K) F_{\ell}(p_K)$$

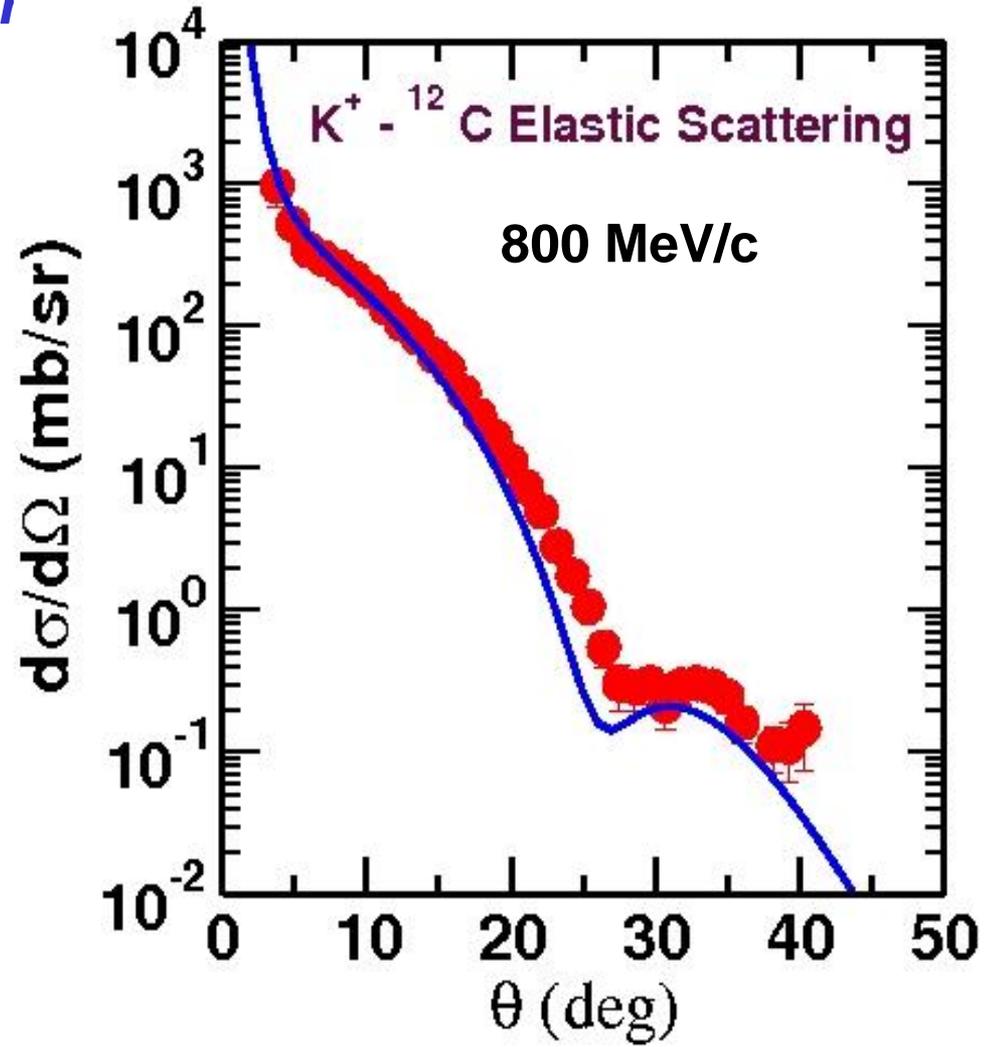
$$F_{\ell}(p_K) = \int j_{\ell}(p_K r) f_{\ell}(p_f, r) r^2 dr$$

Kaon optical potential

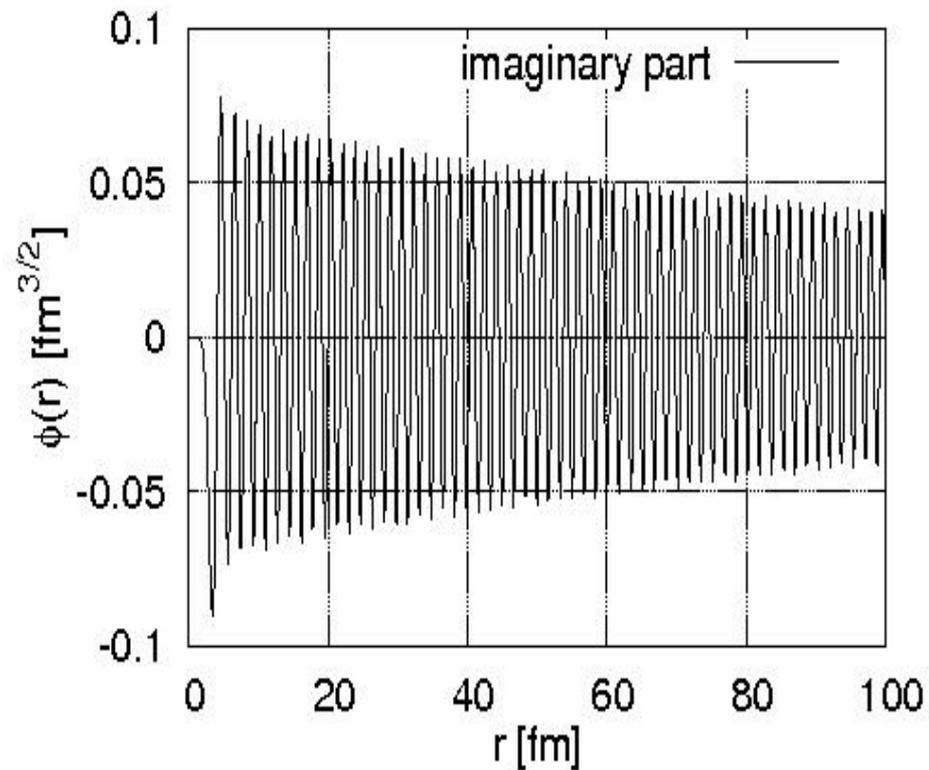
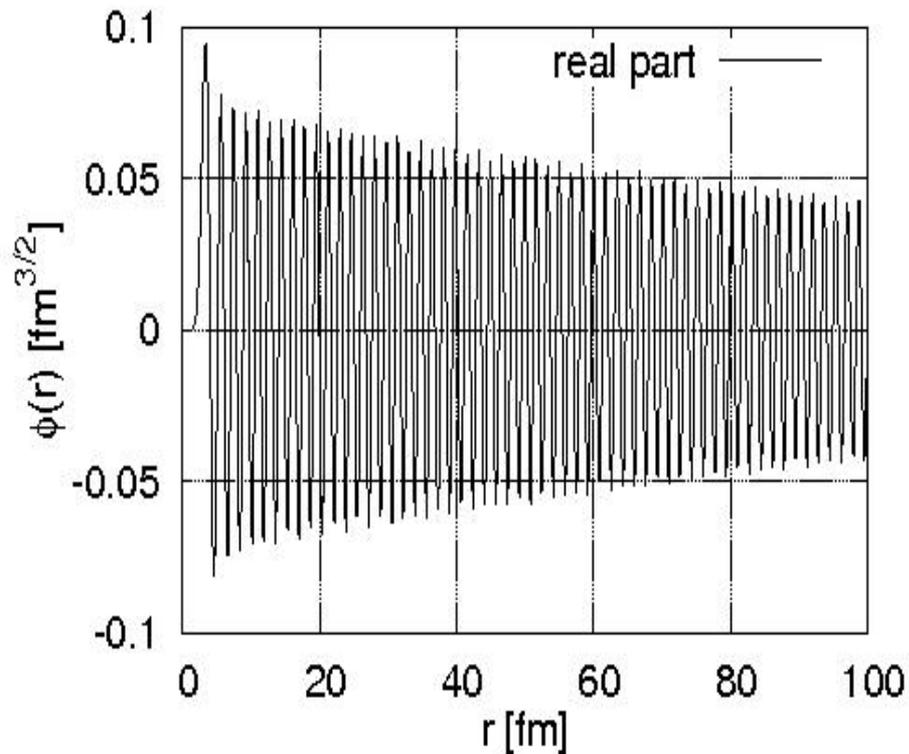
$$2EV_K(r) = -Ab_0 k^2 \rho(r) + Ab_1 \nabla \cdot \rho \nabla$$

b_0 and b_1 are parameters of the potential

This provides $f_{\ell}(p_f, r)$



$$F_e(p_K) = \int_0^{\infty} j_e(p_K r) f_e(p_f, r) r^2 dr$$



Scattering wave function in momentum space

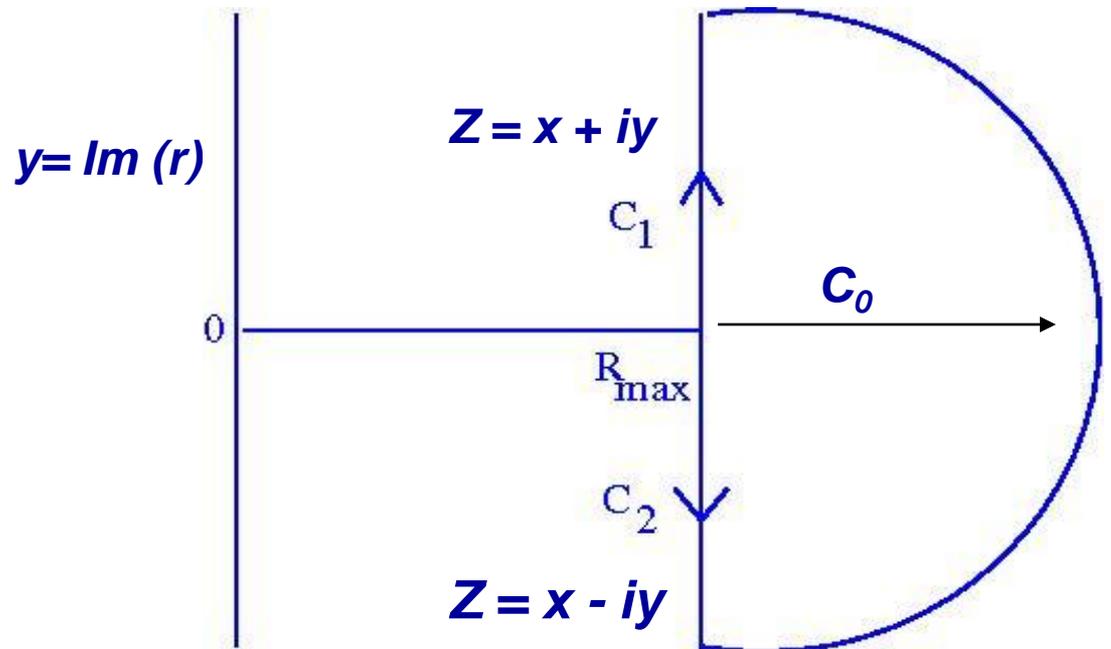
$$F_e(\mathbf{p}_K) = \int_0^\infty j_e(\mathbf{p}_K r) f_e(\mathbf{p}_f, r) r^2 dr$$

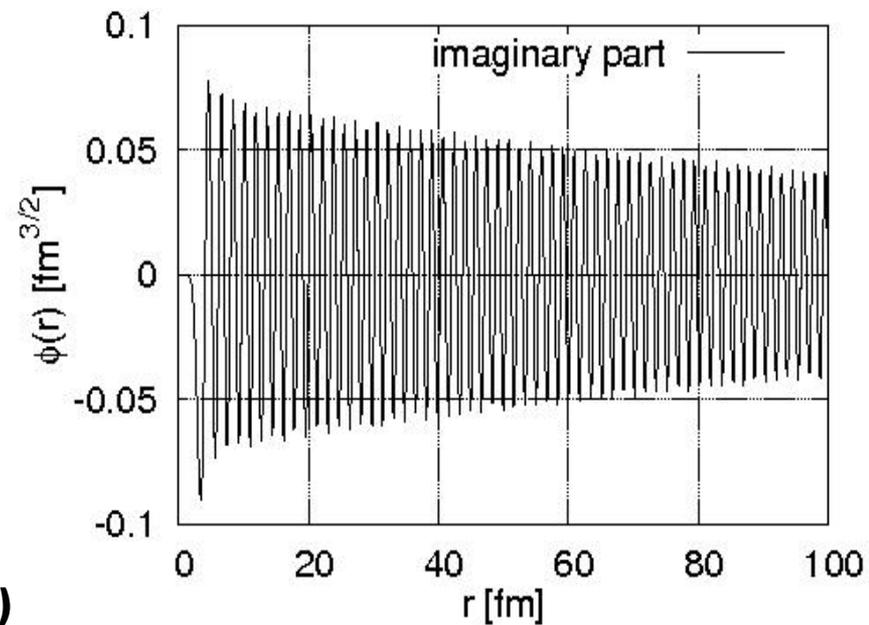
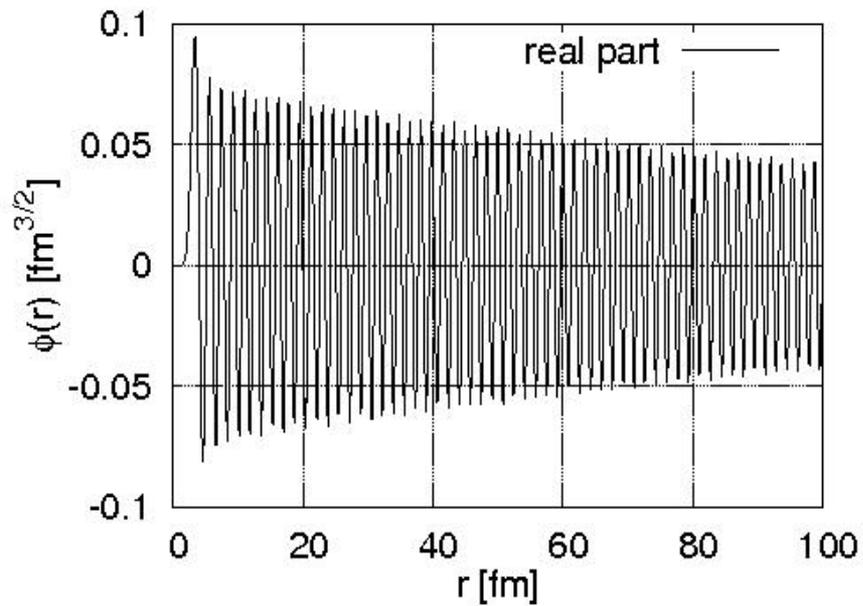
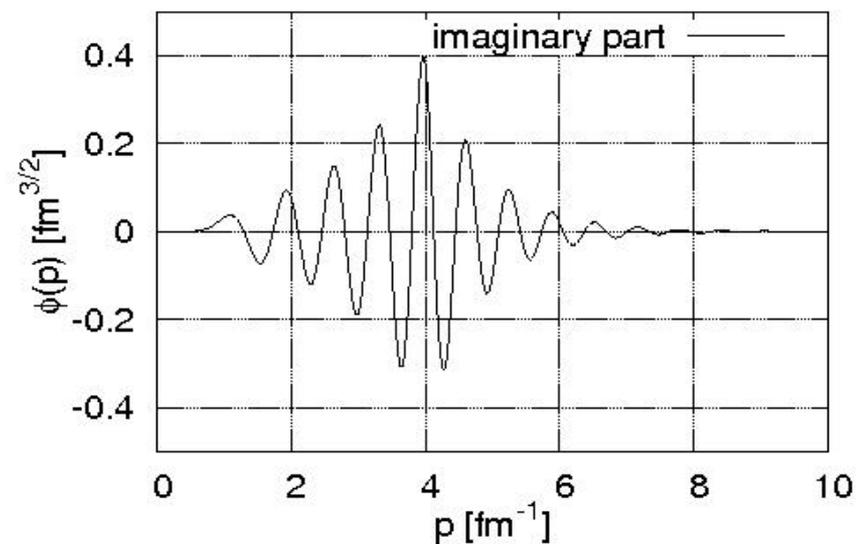
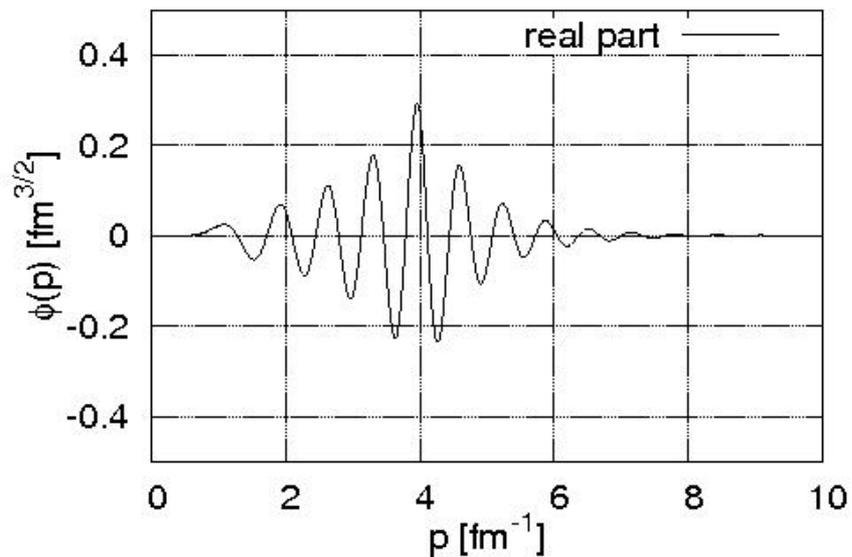
$$= \int_0^{R_{\max}} j_e(\mathbf{p}_K r) f_e(\mathbf{p}_f, r) r^2 dr + \int_{R_{\max}}^\infty j_e(\mathbf{p}_K r) f_e(\mathbf{p}_f, r) r^2 dr$$

$$h_\ell^1(\mathbf{p}_K r) + h_\ell^2(\mathbf{p}_K r)$$

$$H_\ell^-(\mathbf{p}_f r) - S_\ell H_\ell^+(\mathbf{p}_f r)$$

$$I_2 = \exp[\pm i(\mathbf{p}_K \mp \mathbf{p}_f)r]$$

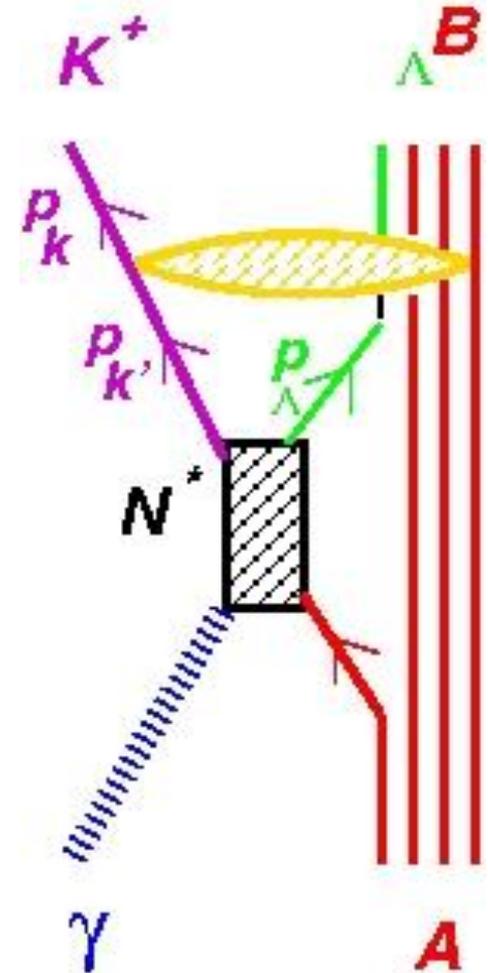


$f_\ell(p_f, r)$ $\ell = 15$  $F_\ell(p_\kappa)$ 

(γ, K^+) reaction on Nuclei

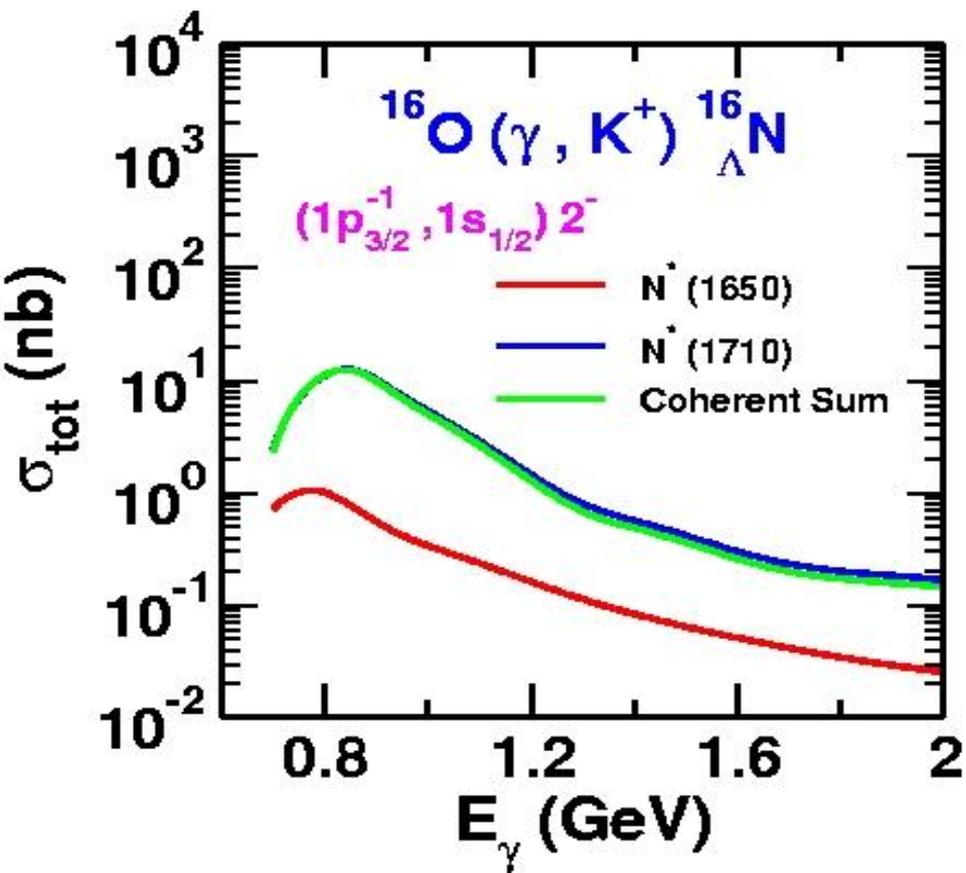
- K^+ is weakly absorbing so reaction occurs deep in the nuclear interior.
- *A proton is converted into a Λ , produces neutron rich hypernuclei.*
- Unnatural parity states strongly excited
- $\gamma p \rightarrow \Lambda K^+$ reaction well understood within an effective Lagrangian picture

Excitations of $N^*(1650)$, $N^*(1710)$, $N^*(1720)$ resonances.

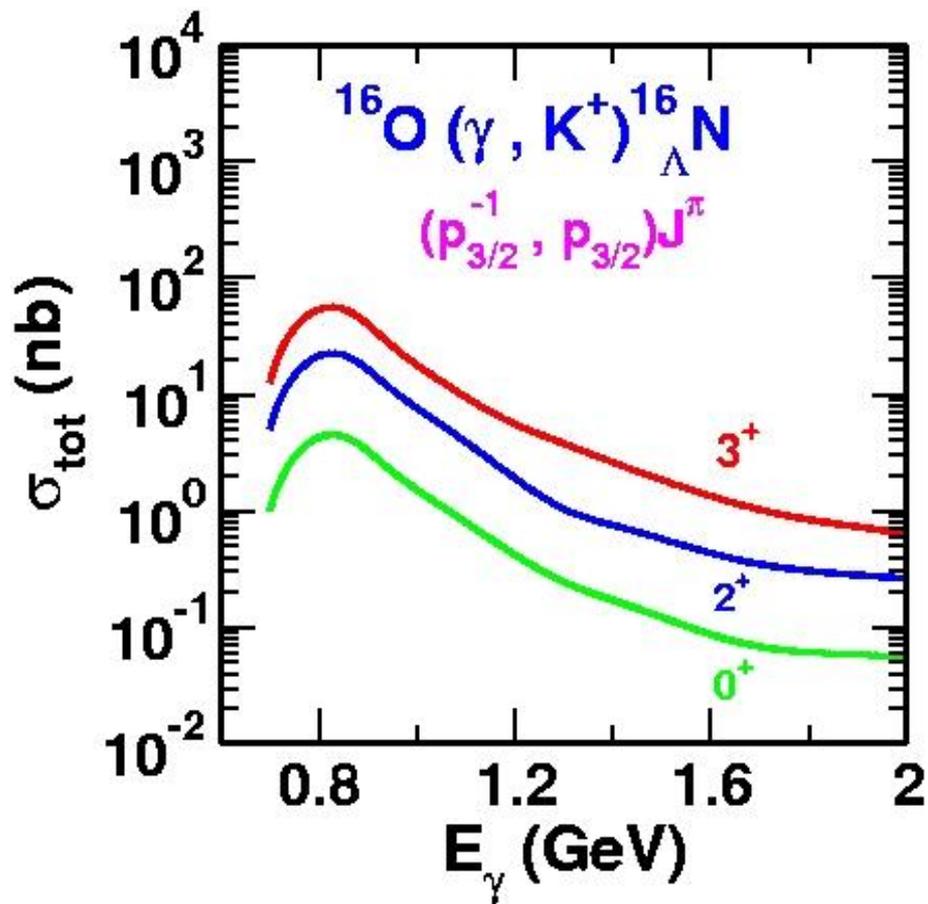


$^{16}\text{O} (\gamma, K^+) ^{16}\text{N}$ Reaction

Relative Contribution of Various resonances



Different J excitation

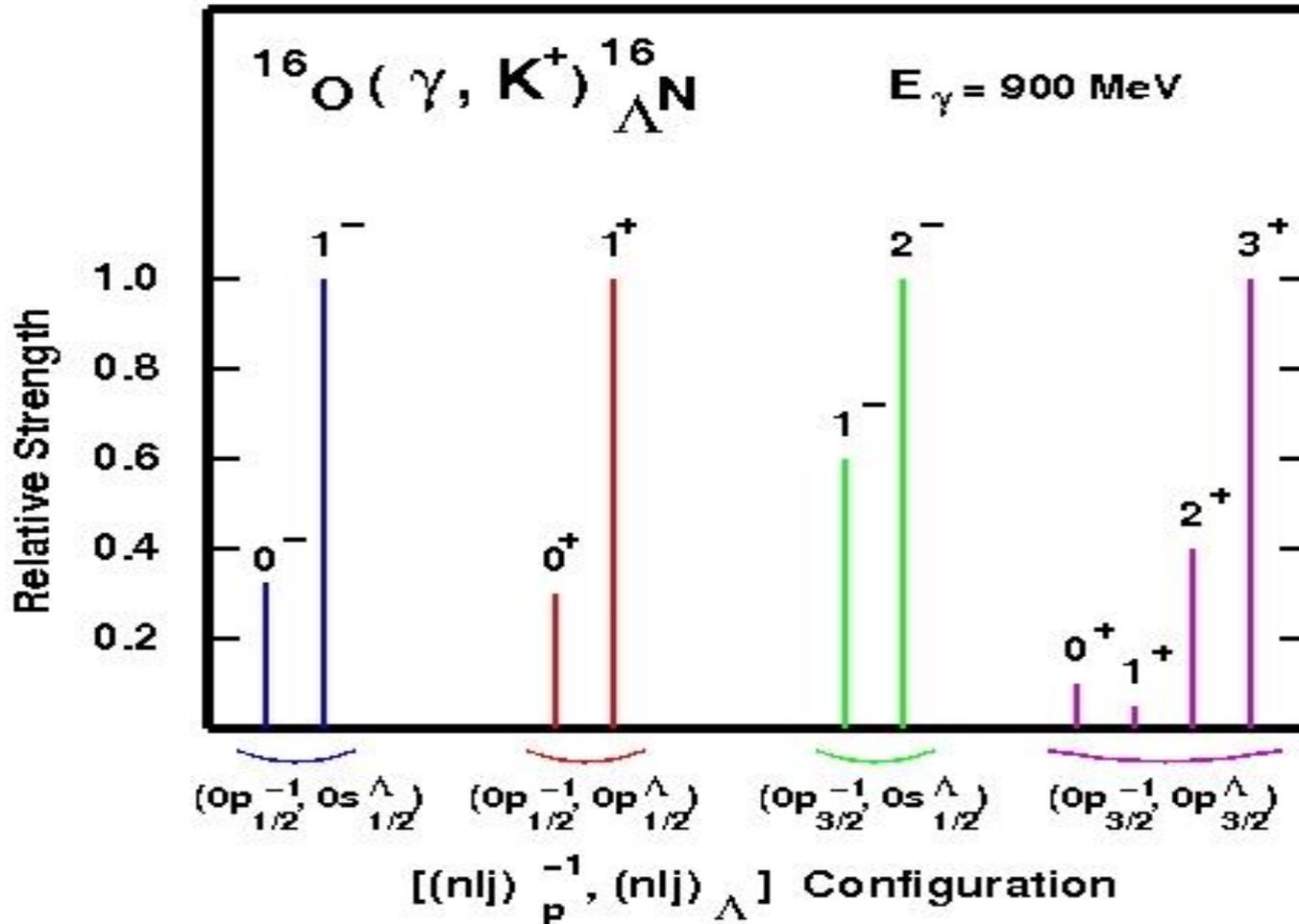


$N^*(1720)$ weaker by 3-4 orders of magnitude

Largest J state is dominant

Jlab results from (γ^*, K^+)

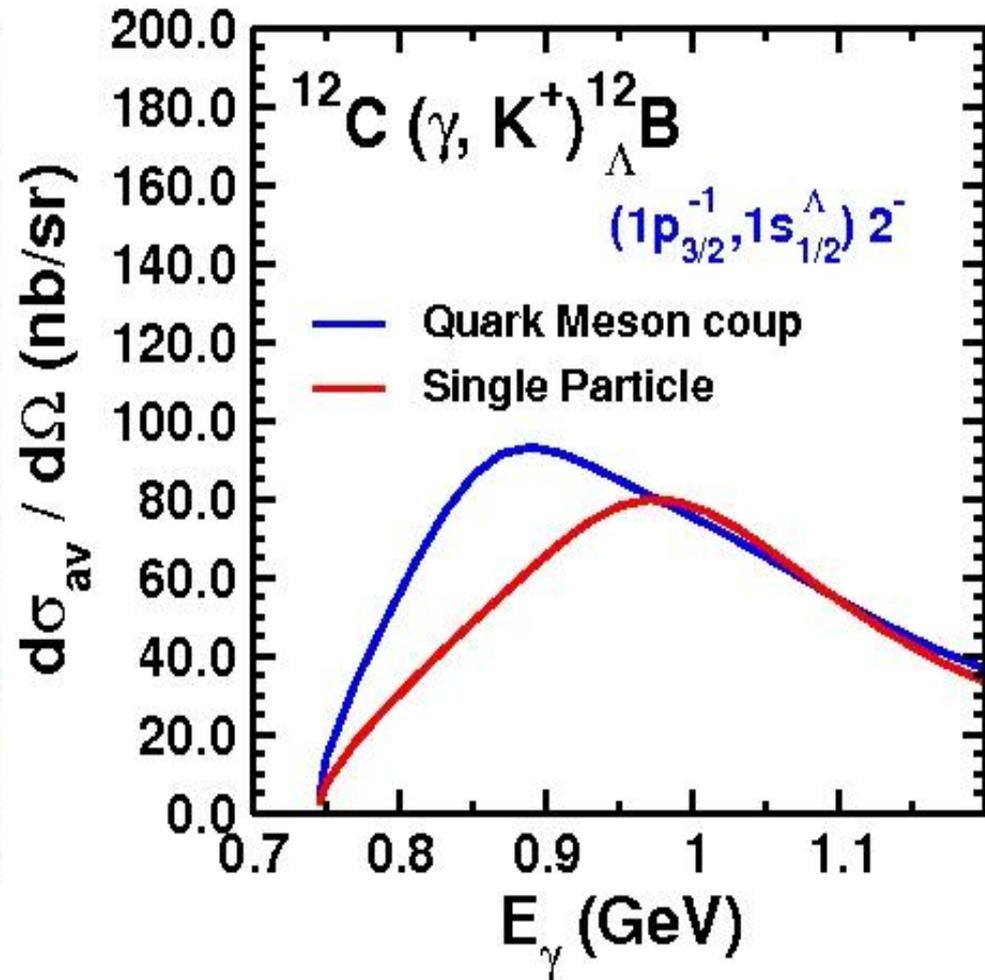
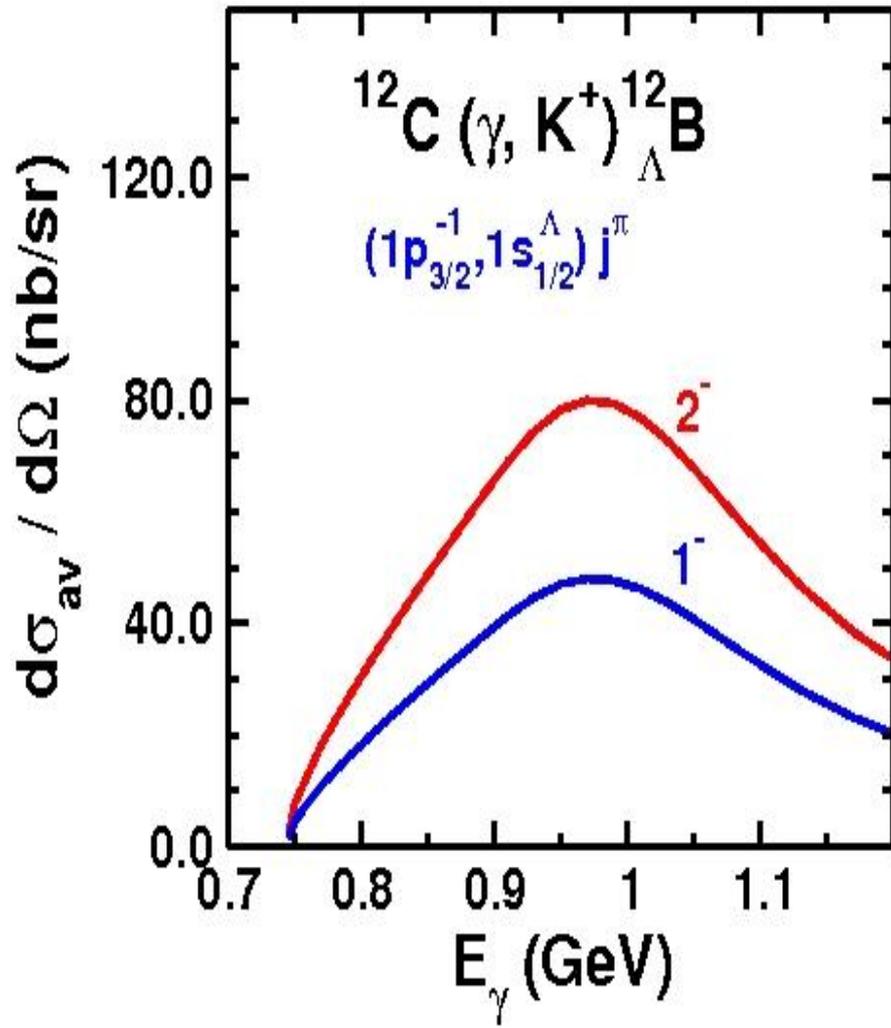
Excitation spectrum of 4 groups of (p^{-1}, Λ) states



In each group highest J is most strong.

Unnatural Parity states dominate.

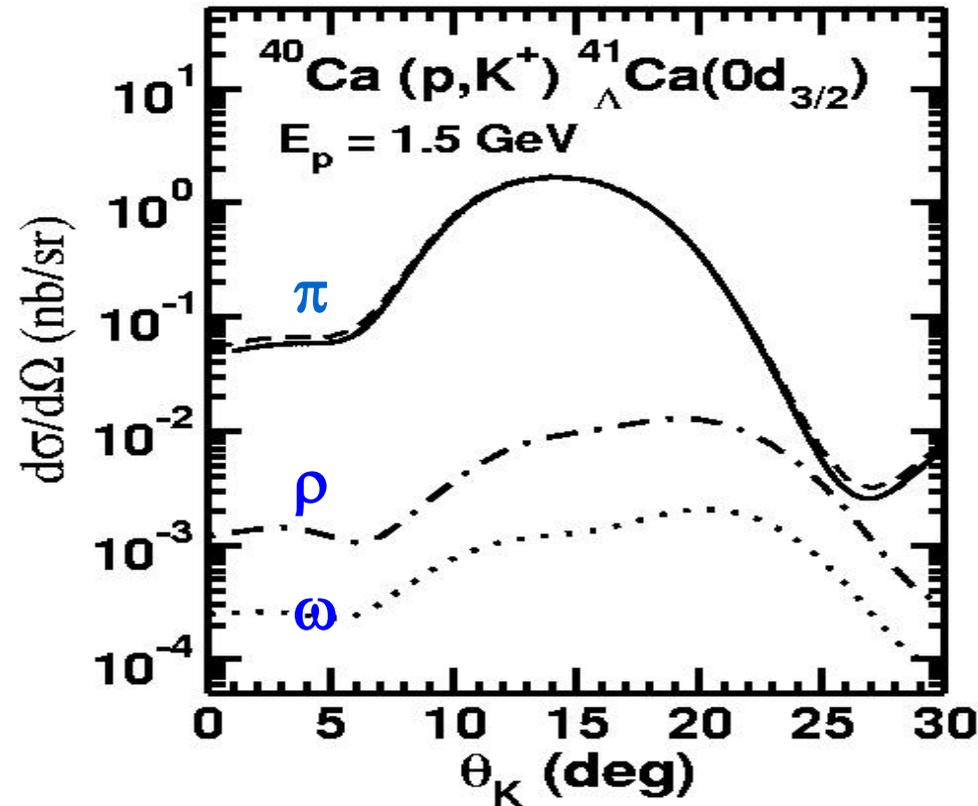
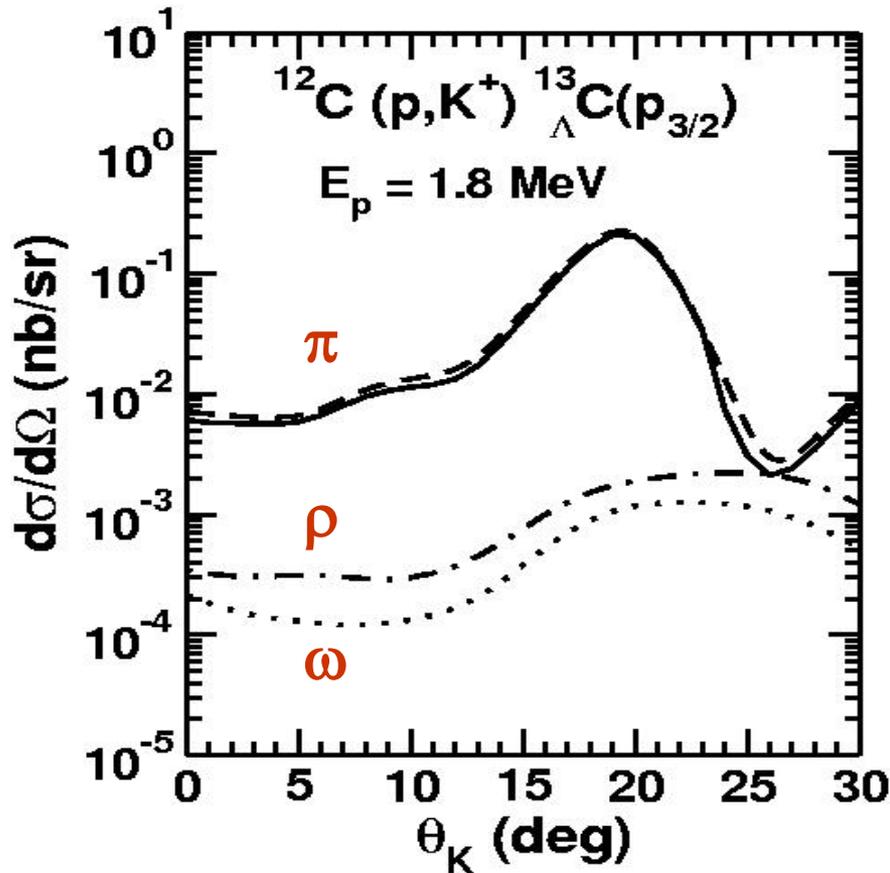
$^{12}\text{C} (\gamma, K^+) ^{12}_{\Lambda}\text{B}$ Reaction



RS, K. Tsushima and A.W. Thomas, In preparation

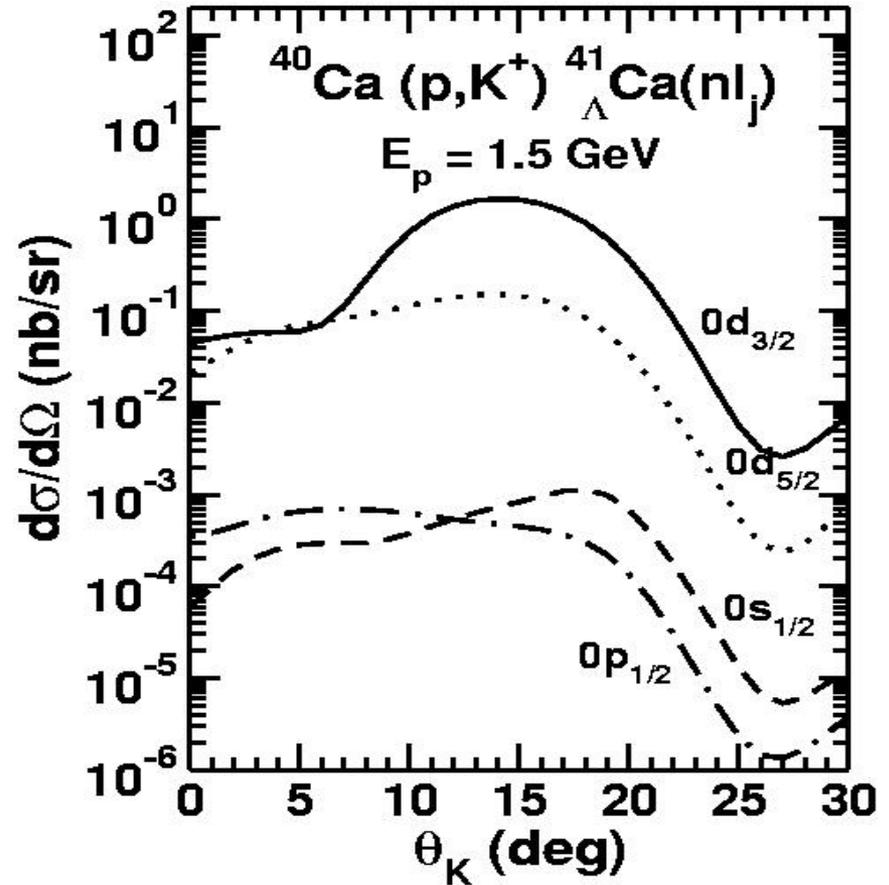
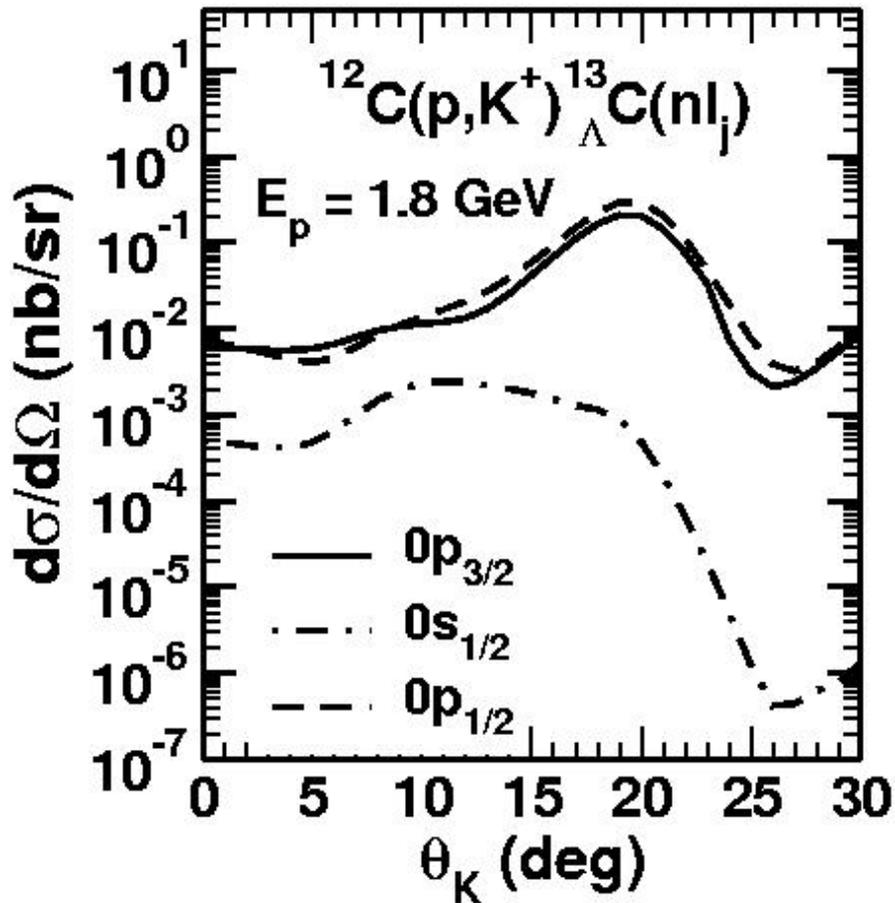
$A(p, K^+)_{\Lambda}B$

Contributions of Various Meson Exchange Processes



π exchange dominates, ρ and ω exchange more important at back angles due to large momentum transfers.

Binding energy selectivity



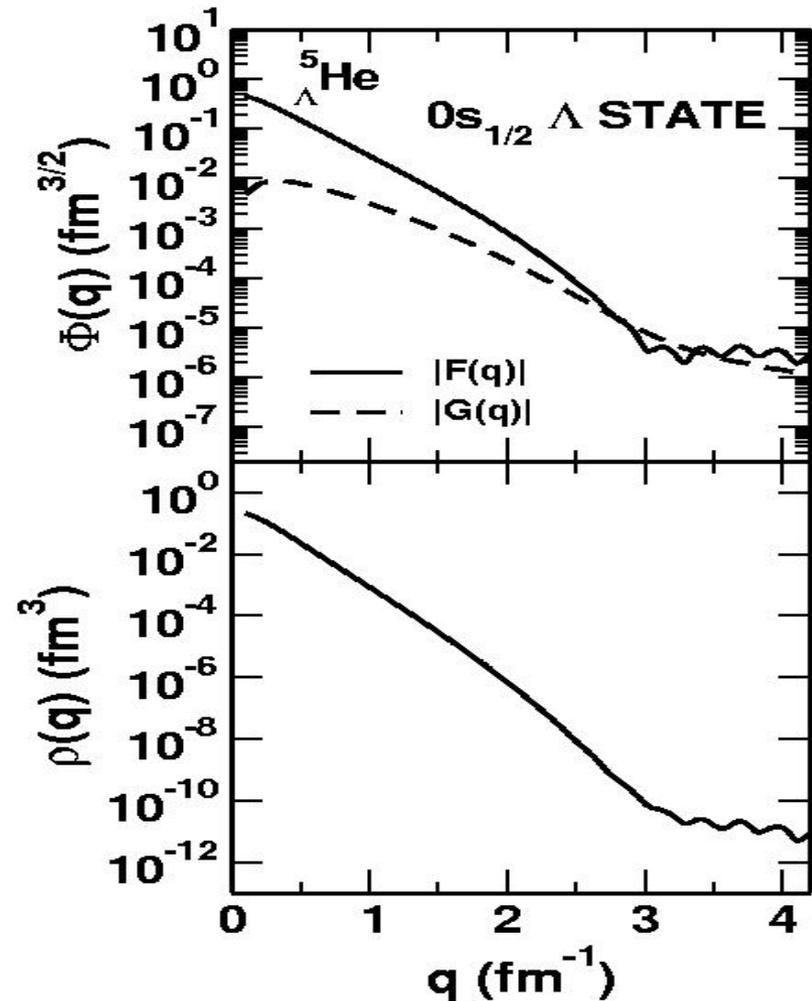
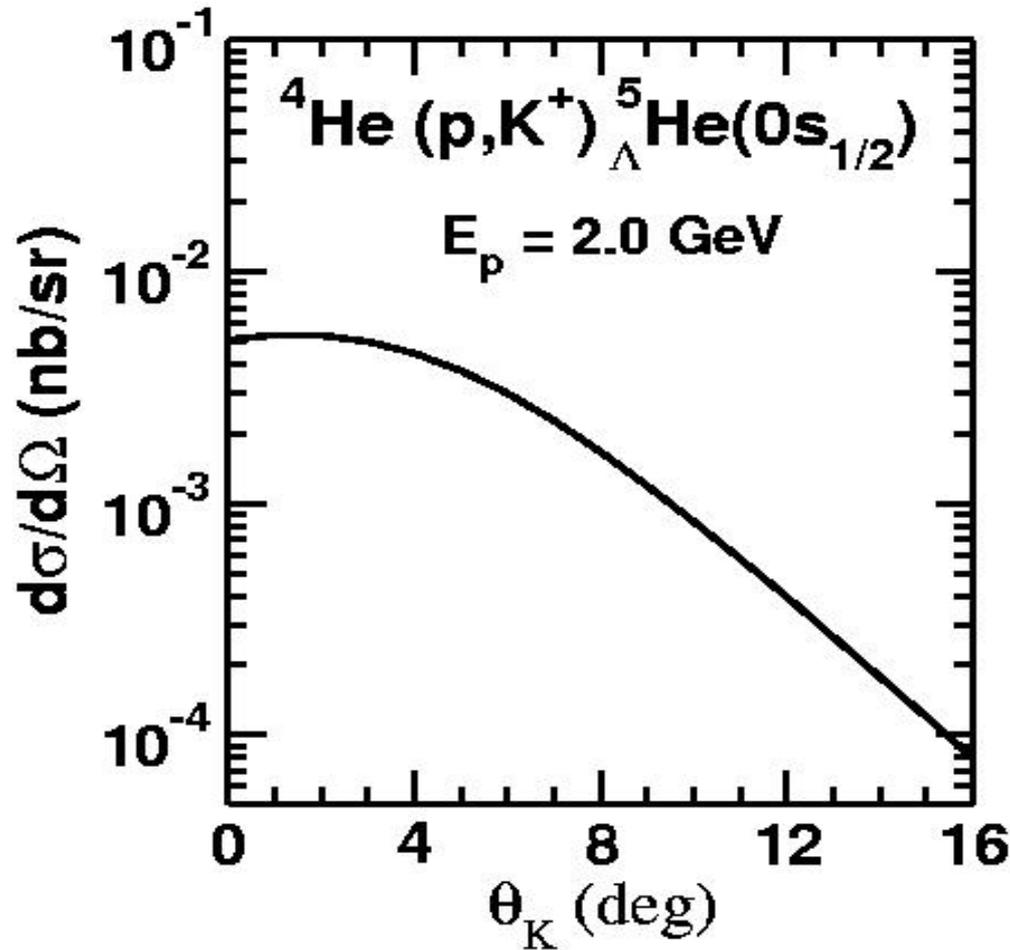
$B_{\Lambda} (^{13}_{\Lambda}\text{C})$

$0p_{1/2}$	0.708 MeV
$0p_{3/2}$	0.860
$0s_{1/2}$	11.690

$B_{\Lambda} (^{41}_{\Lambda}\text{Ca})$

$0d_{3/2}$	0.753
$0d_{5/2}$	1.544
$0p_{1/2}$	9.140
$0s_{1/2}$	17.882

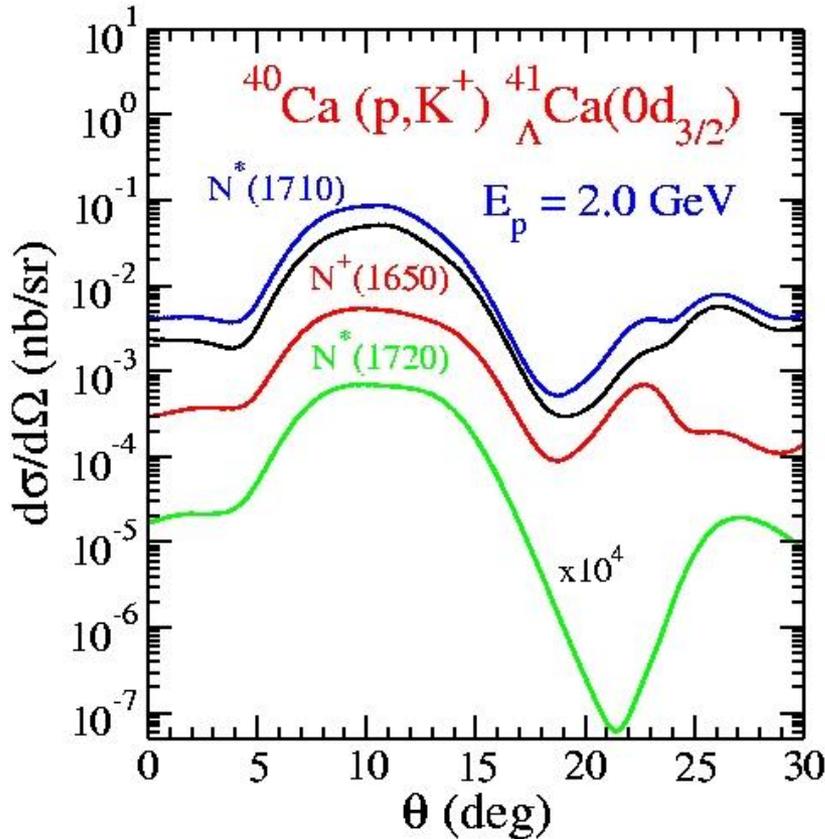
Results for ^4He target



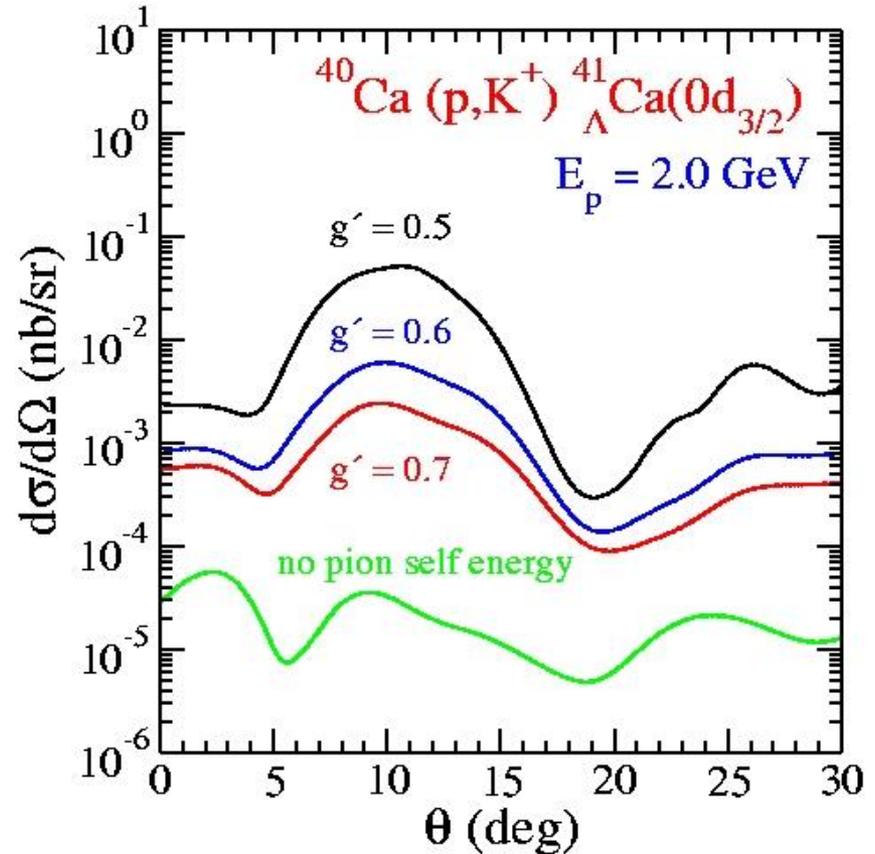
For momentum transfers of interest, Dirac spinors smoothly varying, and are devoid of structures.

$A(p, K^+)_{\Lambda} B$

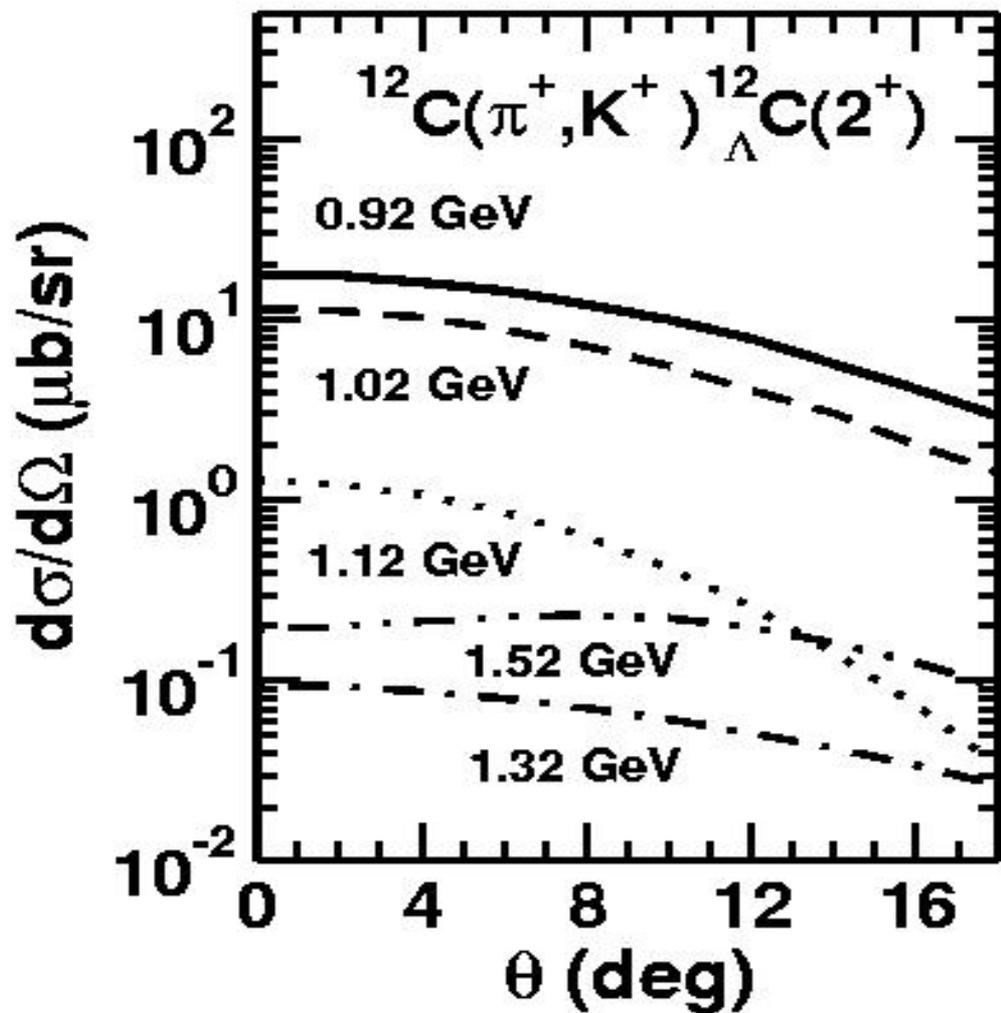
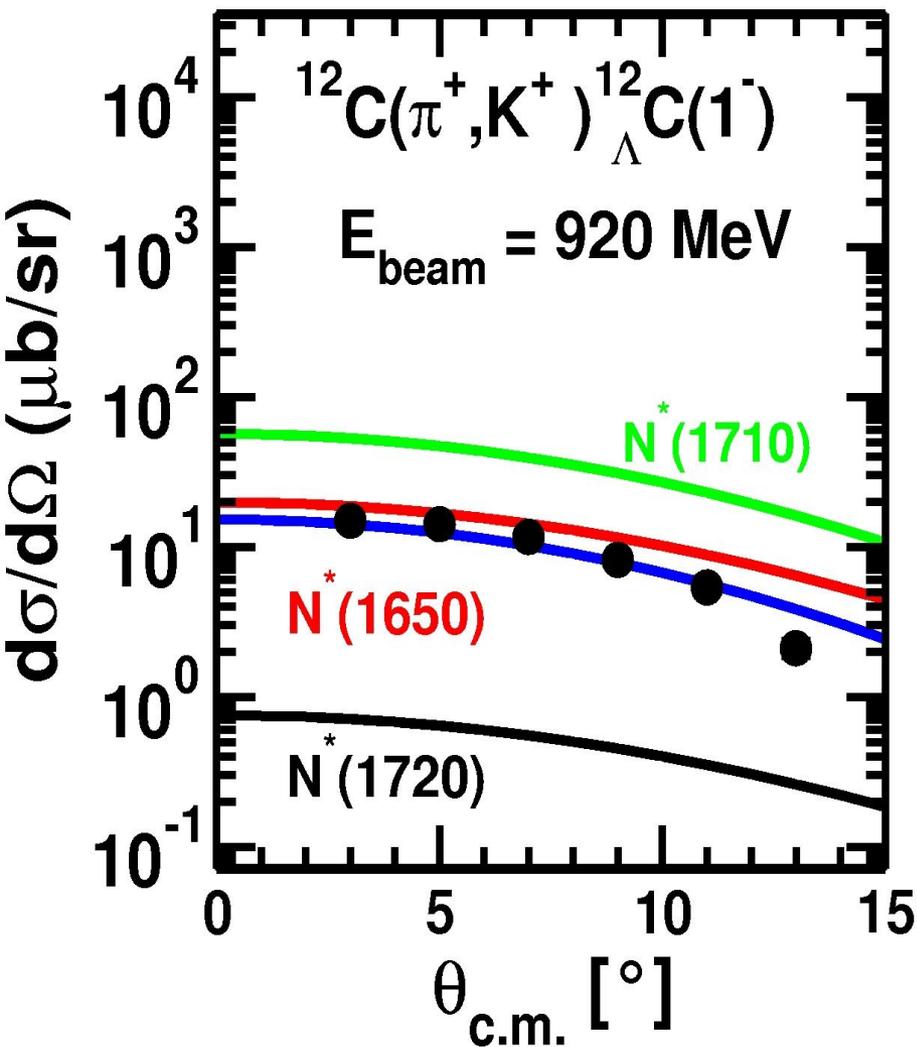
Relative Contribution of Various resonances



Role of pion self energy



(π^+, K^+) reaction



$N^*(1710)$ dominates in his case too.

S. Bender

SUMMARY AND OUTLOOK

$A(h\gamma, K^+)_{\Lambda}B$ reactions provide mutually complimentary information about the hypernuclear spectrum.

A fully covariant description of these reactions is desirable and is possible.

Hypernuclear states with largest angular momentum are dominant, typical of large momentum transfer reactions.

(γ, K^+) , (γ^*, K^+) strongly excites also the unnatural parity states.

Tighter constraints on the models of Λ -N interaction

In-medium Chiral Dynamics, orxiv:0709.2298 (Weise)

Quark Meson Coupling Model (Tsushima, A.W. Thomas, Saito, Guichon)

Coll: H. Lenske, U. Mosel, K. Tsushima and A.W. Thomas