

Linking Nuclear Observables through Quark-Hadron Duality

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Outline

1. Quark-hadron (“Bloom-Gilman”) duality
2. Local duality
 - elastic duality
 - medium modifications
3. Truncated moments

Quark-Hadron ("Bloom-Gilman") duality

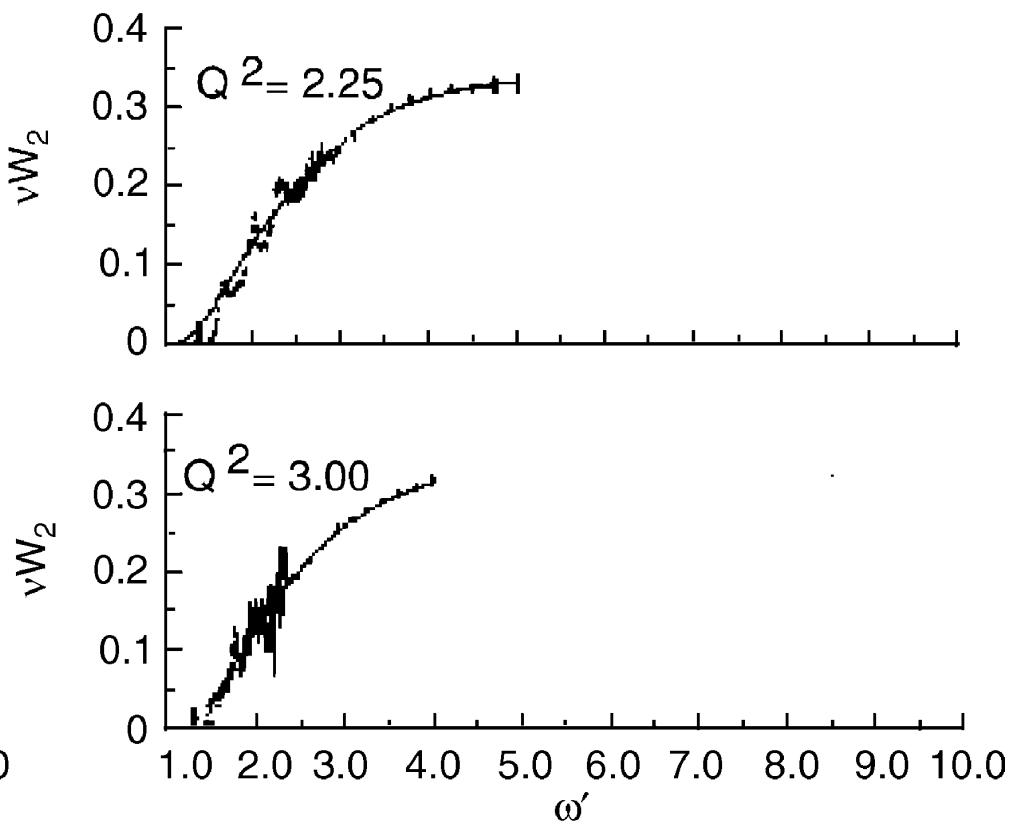
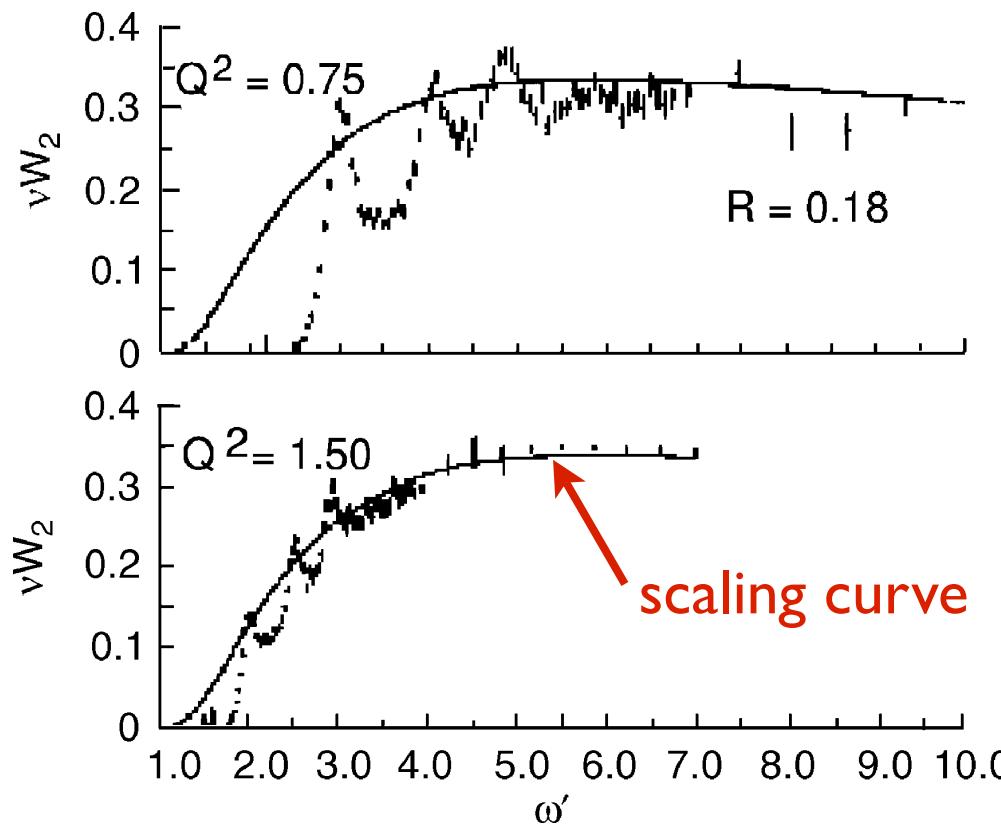
Quark-hadron duality

Complementarity between *quark* and
hadron descriptions of observables

$$\sum_{\text{hadrons}} = \sum_{\text{quarks}}$$

Can use either set of complete basis states
to describe all physical phenomena

Electron scattering



Bloom, Gilman, Phys. Rev. Lett. 85 (1970) 1185

→ resonance – scaling duality in
proton $\nu W_2 = F_2$ structure function

Bloom-Gilman duality

Average over (strongly Q^2 dependent) resonances
 $\approx Q^2$ independent scaling function

Finite energy sum rule for eN scattering

$$\frac{2M}{Q^2} \int_0^{\nu_m} d\nu \nu W_2(\nu, Q^2) = \int_1^{\omega'_m} d\omega' \nu W_2(\omega')$$

measured structure function
(function of ν and Q^2)

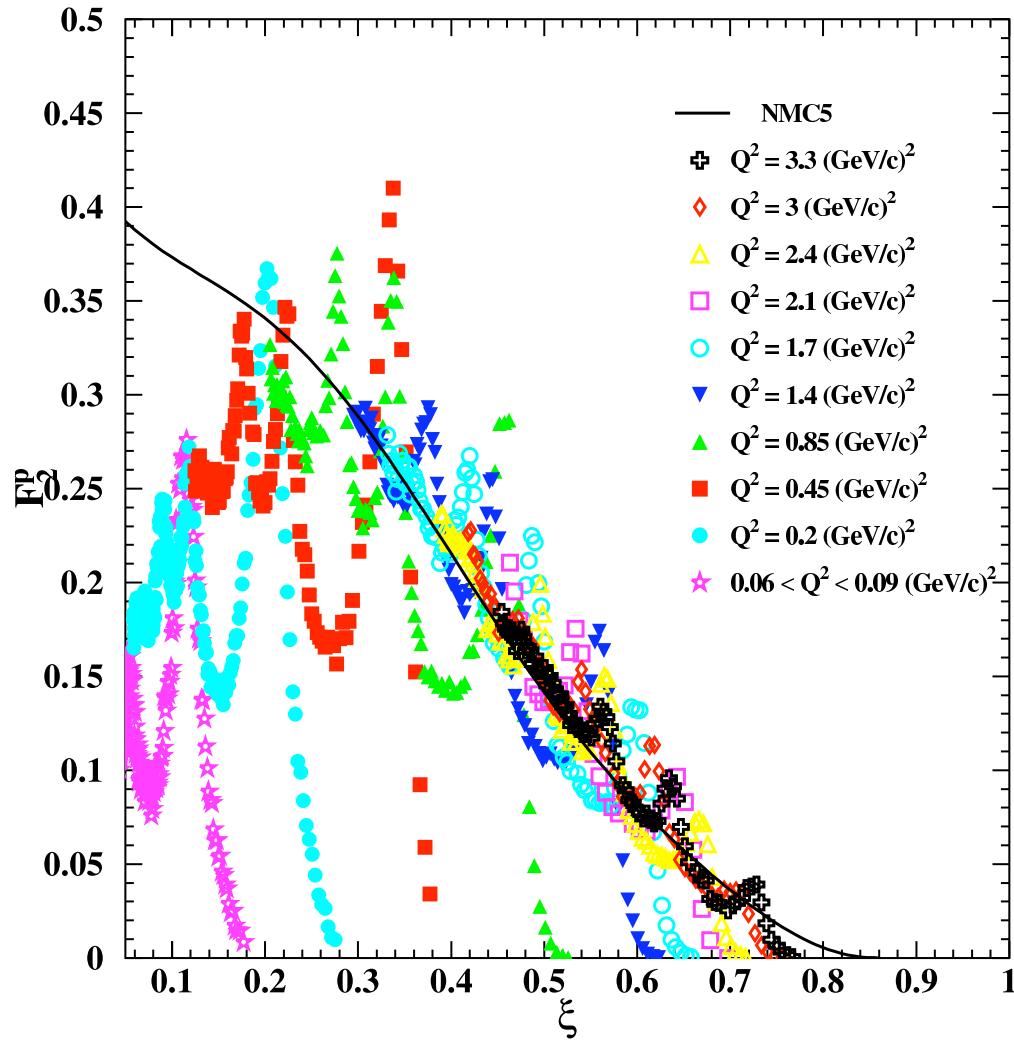
“hadrons”

$$\omega' = \frac{1}{x} + \frac{M^2}{Q^2}$$

scaling function
(function of ω' only)

“quarks”

Bloom-Gilman duality

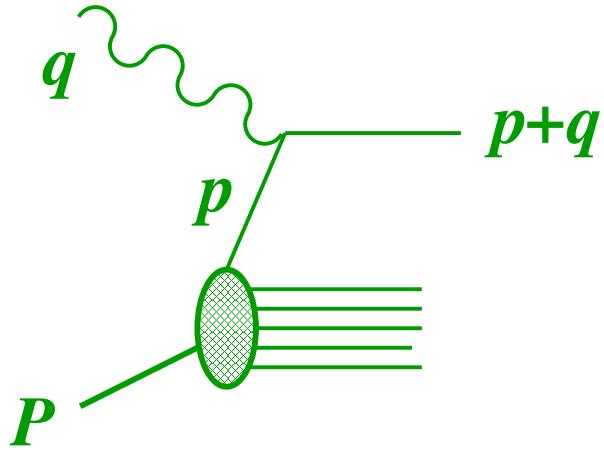


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Niculescu *et al.*, Phys. Rev. Lett. 85 (2000) 1182

Average over
(strongly Q^2 dependent)
resonances
 $\approx Q^2$ independent
scaling function

Scaling variables



$$(p + q)^2 = m_q^2 \quad \left\{ \begin{array}{l} m_q = 0 \\ p_T = 0 \end{array} \right.$$

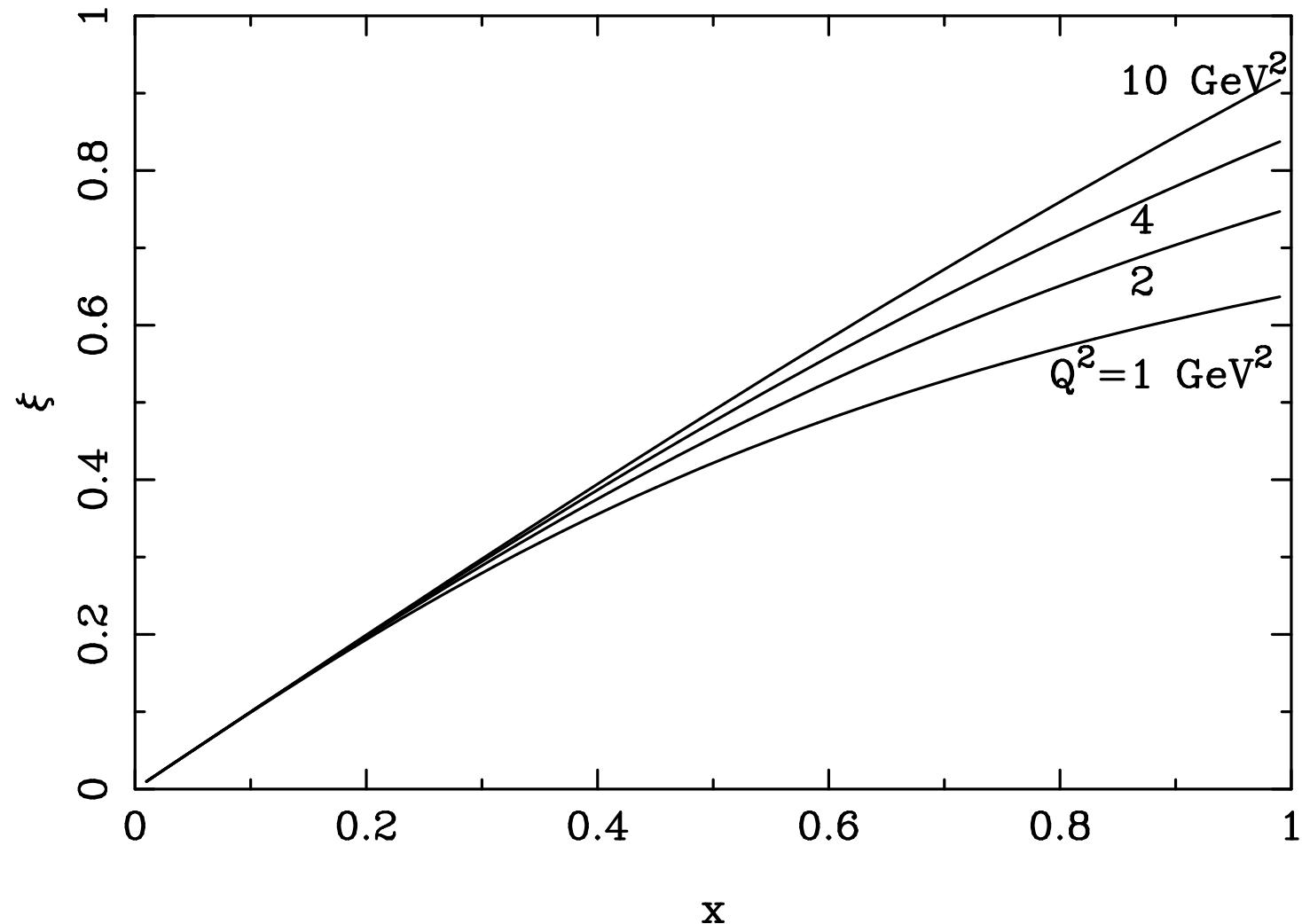
light-cone fraction of target's momentum carried by parton

$$\xi = \frac{p^+}{P^+} = \frac{p^0 + p^z}{M}$$

$$\rightarrow \xi = \frac{2x}{1 + \sqrt{1 + 4x^2 M^2 / Q^2}} \rightarrow x \text{ as } Q^2 \rightarrow \infty$$

Nachtmann scaling variable

Scaling variables



Duality in QCD

Operator product expansion

→ expand moments of structure functions
in powers of $1/Q^2$

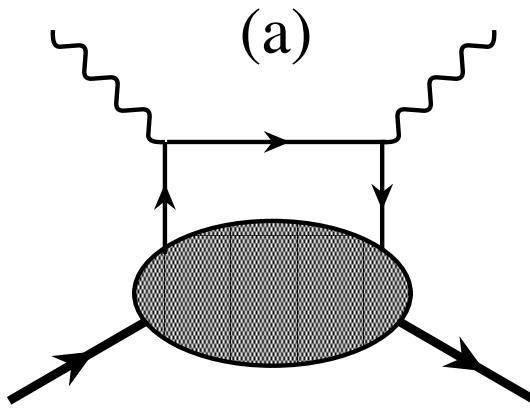
$$\begin{aligned} M_n(Q^2) &= \int_0^1 dx \ x^{n-2} \ F_2(x, Q^2) \\ &= A_n^{(2)} + \frac{A_n^{(4)}}{Q^2} + \frac{A_n^{(6)}}{Q^4} + \dots \end{aligned}$$



matrix elements of operators with
specific “twist” τ

$\tau = \text{dimension} - \text{spin}$

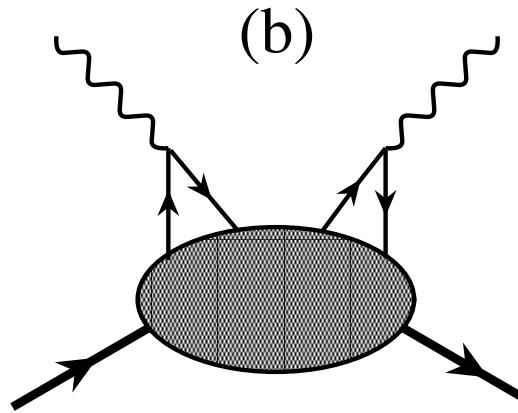
Higher twists



$$\tau = 2$$

single quark
scattering

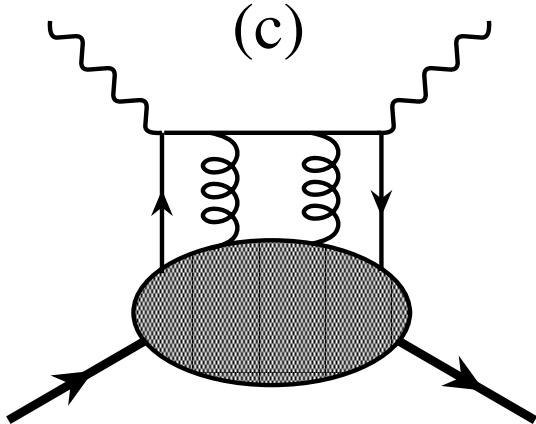
e.g. $\bar{\psi} \gamma_\mu \psi$



$$\tau > 2$$

qq and qg
correlations

e.g. $\bar{\psi} \gamma_\mu \psi \bar{\psi} \gamma_\nu \psi$
or $\bar{\psi} \tilde{G}_{\mu\nu}\gamma^\nu \psi$



Duality in QCD

Operator product expansion

→ expand moments of structure functions
in powers of $1/Q^2$

$$\begin{aligned} M_n(Q^2) &= \int_0^1 dx \ x^{n-2} \ F_2(x, Q^2) \\ &= A_n^{(2)} + \frac{A_n^{(4)}}{Q^2} + \frac{A_n^{(6)}}{Q^4} + \dots \end{aligned}$$

If moment \approx independent of Q^2

→ higher twist terms $A_n^{(\tau>2)}$ small

Duality in QCD

Operator product expansion

→ expand moments of structure functions
in powers of $1/Q^2$

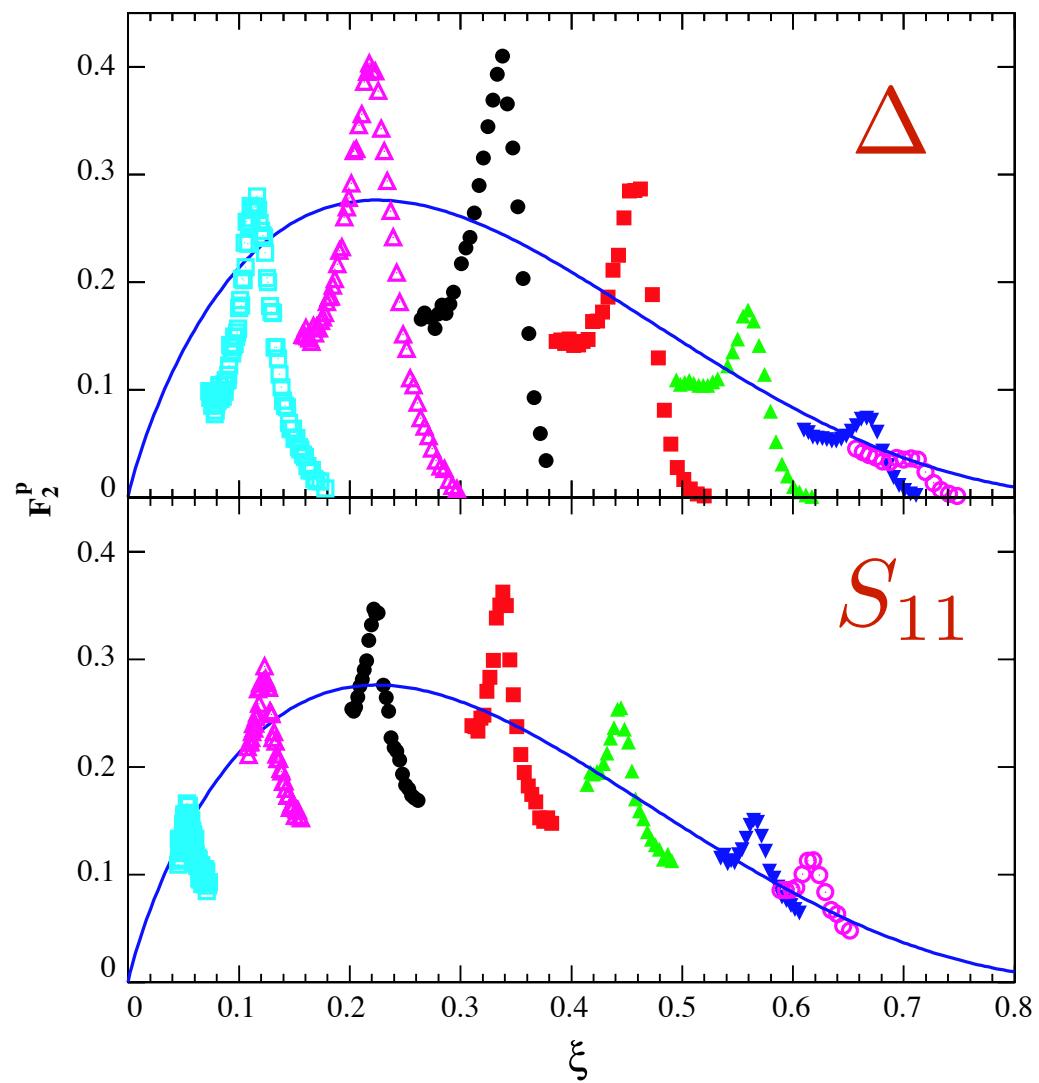
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Duality \iff suppression of higher twists

*de Rujula, Georgi, Politzer,
Ann. Phys. 103 (1975) 315*

Local Duality

Local Bloom-Gilman duality



Local Bloom-Gilman duality

■ contribution of (narrow) resonance R to structure function

$$F_2^{(R)} \approx 2M\nu \left(G_R(Q^2)\right)^2 \delta(W^2 - M_R^2)$$

if $G_R(Q^2) \sim (1/Q^2)^n$ then for $Q^2 \gg M_R^2$

$$F_2^{(R)} \approx (1 - x_R)^{2n-1} \quad \text{"Drell-Yan-West relation"}$$

with

$$x_R = \frac{Q^2}{Q^2 + M_R^2 - M^2}$$

→ as $Q^2 \rightarrow \infty$, $x_R \rightarrow 1$

resonances move to larger x

Local elastic duality

- extreme case of local duality for elastic peak

→ elastic contribution to structure function

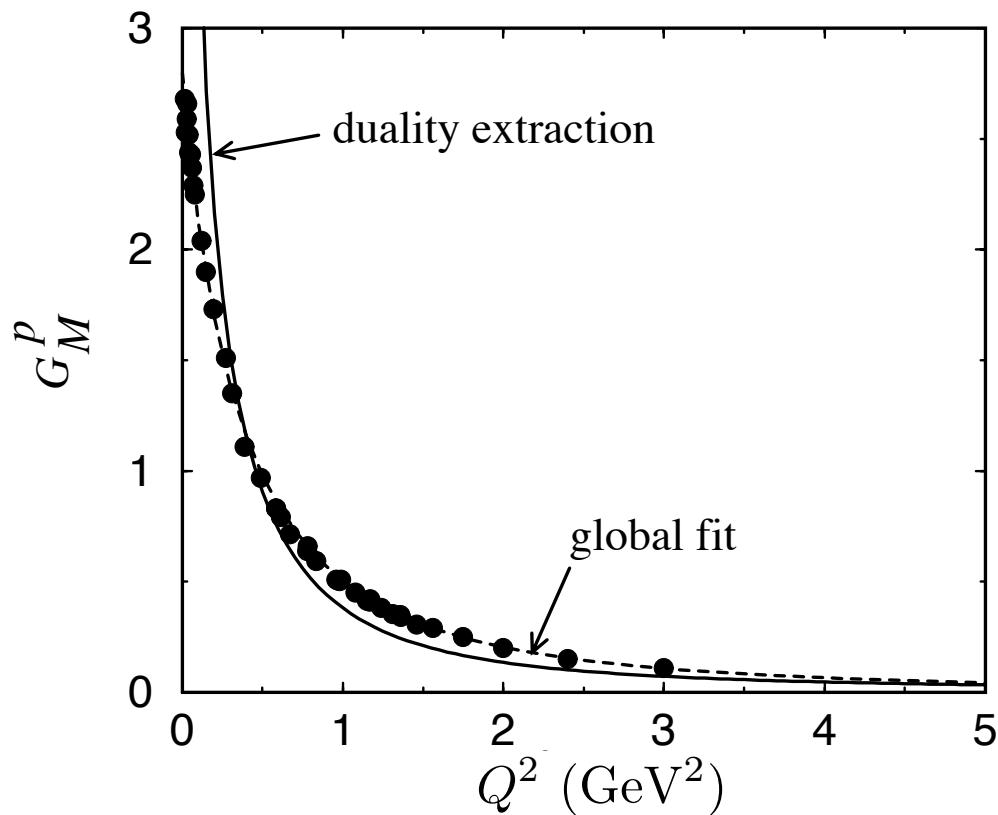
$$F_2^{(el)} = \frac{2M\tau}{1+\tau} (G_E^2 + \tau G_M^2)^2 \delta(\nu - Q^2/2M) \quad \tau = \frac{Q^2}{4M^2}$$

- hypothesis: area under elastic peak same as integral of scaling structure function below threshold

$$\int_1^{\delta\omega'} d\omega' F_2^{\text{LT}}(\omega') = \frac{2M}{Q^2} \int d\nu F_2^{(el)}(\nu, Q^2) \quad \omega' = \frac{2M\nu + M^2}{Q^2}$$
$$= \frac{G_E^2 + \tau G_M^2}{1 + \tau} \quad \begin{matrix} \text{Bloom-Gilman} \\ \text{scaling variable} \end{matrix}$$

Local elastic duality

- extract magnetic form factor from integral of F_2



→ good to $\sim 30\%$ for $Q^2 \sim \text{few GeV}^2$

Local elastic duality

- conversely, differentiate local duality relation w.r.t. Q^2 to obtain structure function at threshold

$$F_2(x = x_{\text{th}}) = \beta \left[\frac{G_M^2 - G_E^2}{2M^2(1 + \tau)^2} + \frac{2}{1 + \tau} \left(\frac{dG_E^2}{dQ^2} + \tau \frac{dG_M^2}{dQ^2} \right) \right]$$

where

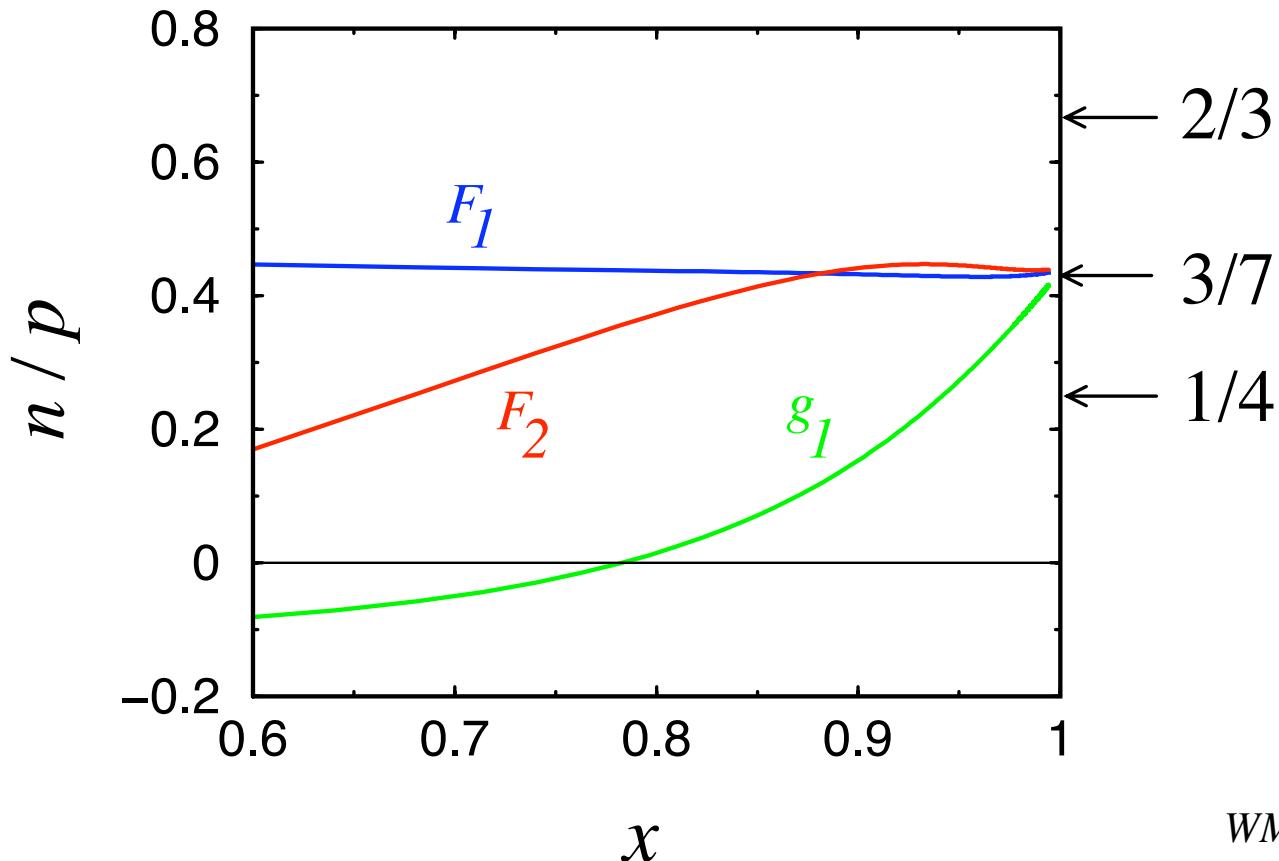
$$\beta = \frac{(Q^4/M^2)(\xi_0^2/\xi^3)(2 - \xi/x)}{2\xi_0 - 4}$$

$$\xi_0 = \xi(x = 1)$$

→ structure functions at large x from form factors !

Local elastic duality

- neutron to proton structure function ratios



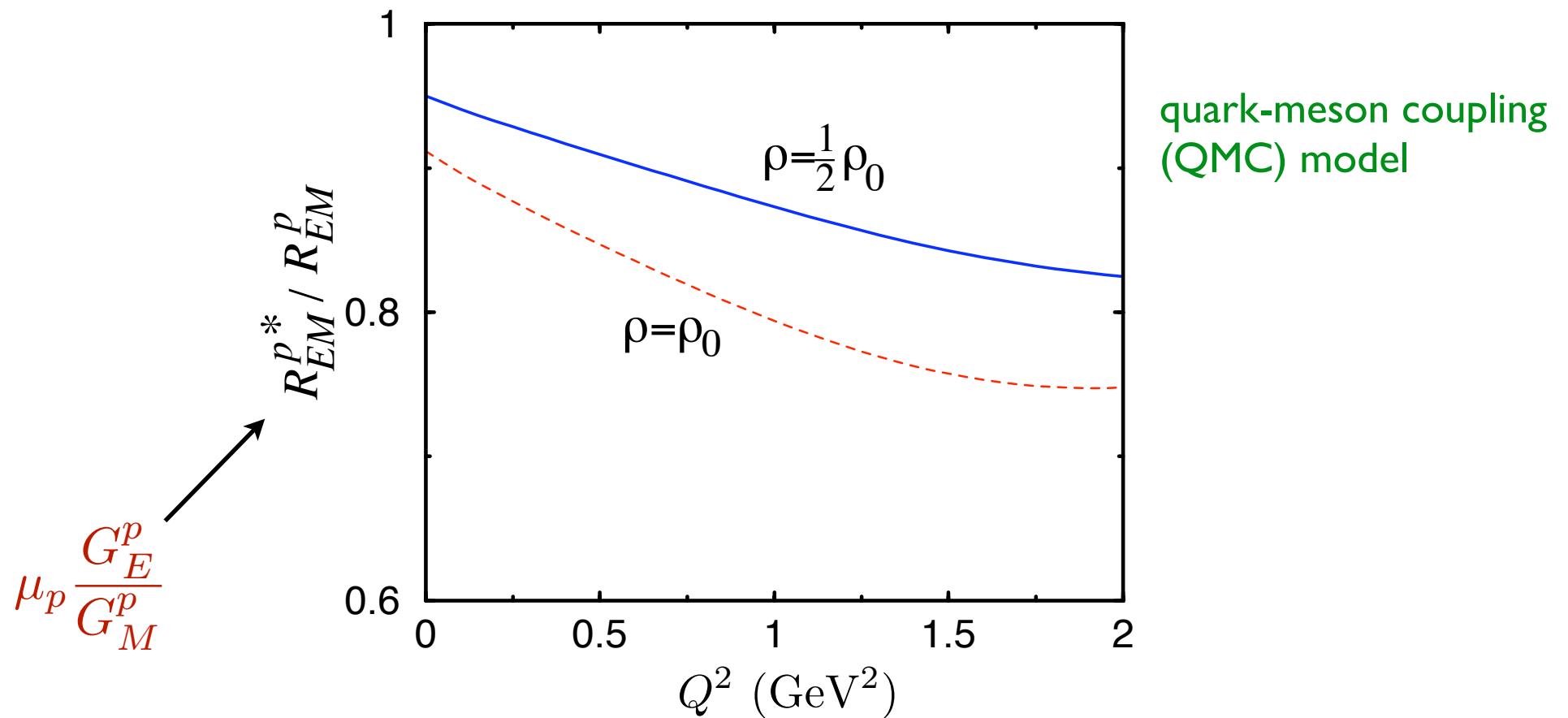
WM, PRL 86 (2001) 35

→ testable predictions for $x \rightarrow 1$ behavior

Local Duality & Nuclear Modifications

WM, Tsushima, Thomas

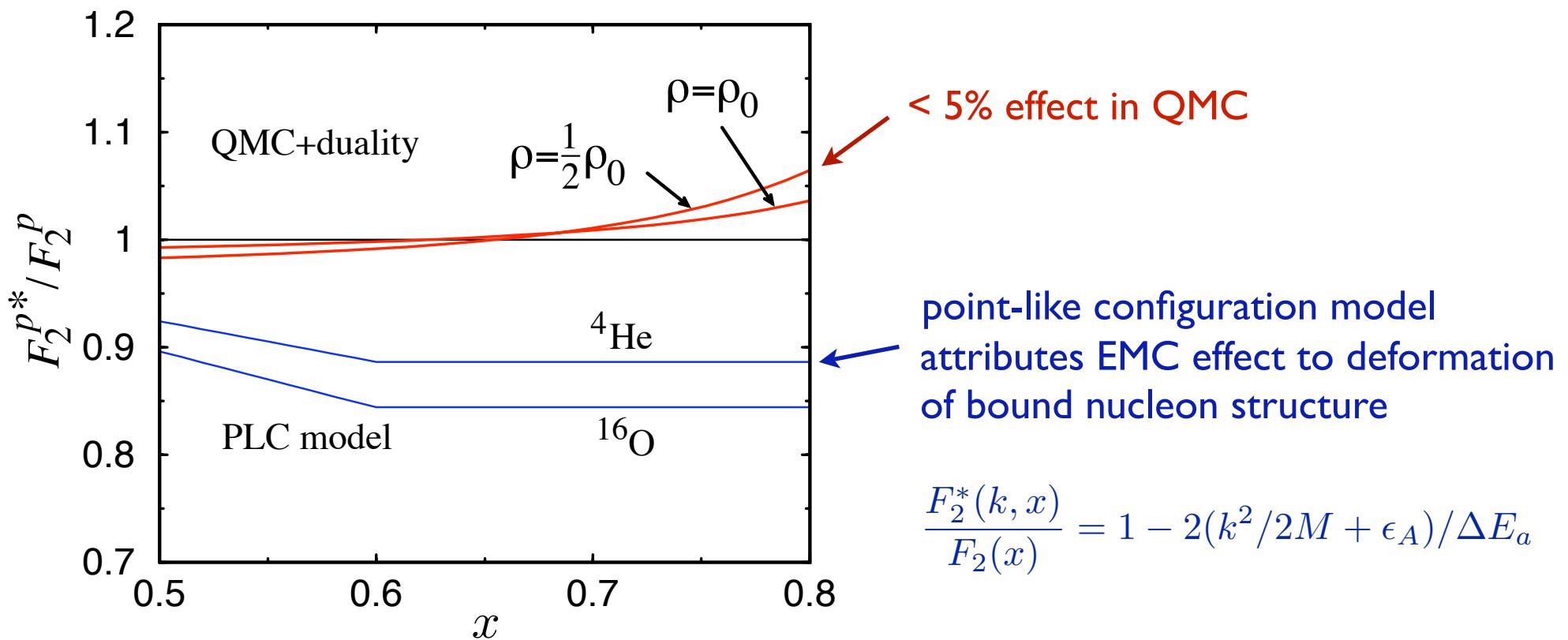
- can recent ${}^4\text{He}$ ($e,e'p$) data be interpreted in terms of medium modified form factors ?
- use local duality to relate medium modified *form factors* to medium modified *structure functions* (EMC effect)



■ medium modified structure functions

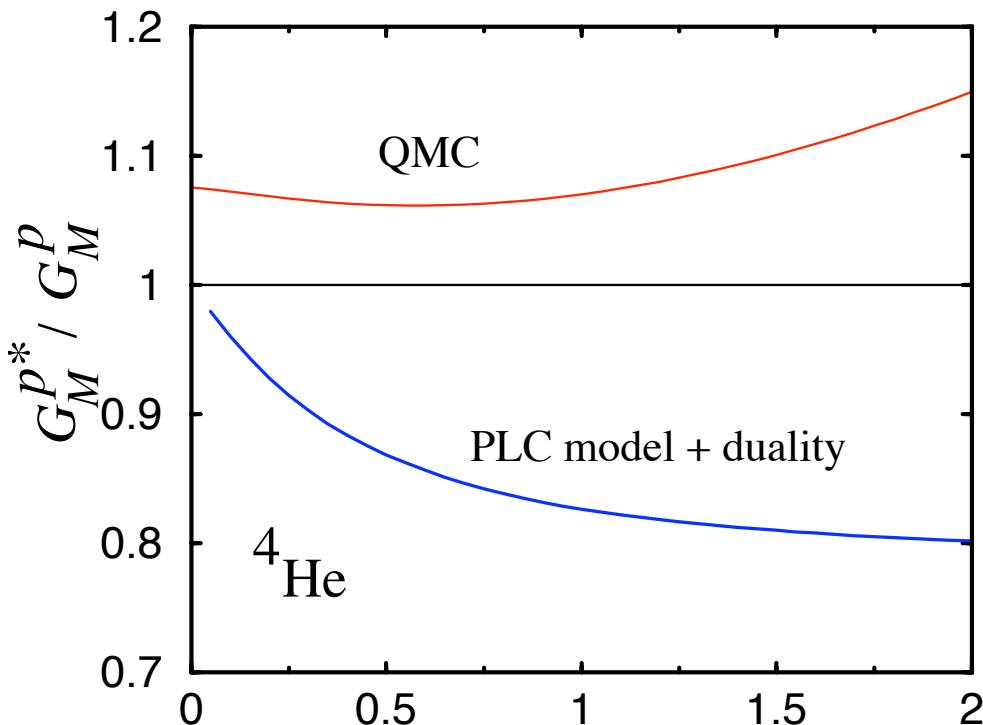
$$\frac{F_2^{p*}}{F_2^p} \approx \frac{dG_M^{p*2}/dQ^2}{dG_M^{p2}/dQ^2} \quad \text{large } Q^2$$

note: threshold for bound nucleon at $x_{\text{th}}^* = \left(\frac{m_\pi(2M + m_\pi^2) + Q^2}{m_\pi(2(M^* + V) + m_\pi) + Q^2} \right) x_{\text{th}}$



- conversely, change in form factor of bound nucleon implied by change in structure function in medium

$$[G_M^p(Q^2)]^2 \approx \frac{2 - \xi_0}{\xi_0^2} \frac{(1 + \tau)}{(1/\mu_p^2 + \tau)} \int_{\xi_{\text{th}}}^1 d\xi \ F_2^p(\xi)$$



predicts 20% suppression
 in magnetic form factor
 → enhancement of PT ratio
 → contrary to ${}^4\text{He}$ data

Truncated Moments

Psaker, Christy, Keppel, WM (2007)

Truncated moments

- complete moments can be studied in QCD via twist expansion
 - Bloom-Gilman duality has a precise meaning
(*i.e.*, duality violation = higher twists)
- for “local” duality, difficult to make rigorous connection with QCD
 - *e.g.* need prescription for how to average over resonances
- truncated moments allow study of restricted regions in x (or W) within QCD in well-defined, systematic way

$$\overline{M}_n(\Delta x, Q^2) = \int_{\Delta x} dx \ x^{n-2} \ F_2(x, Q^2)$$

Truncated moments

- truncated moments obey DGLAP-like evolution equations, similar to PDFs

$$\frac{d\overline{M}_n(\Delta x, Q^2)}{d \log Q^2} = \frac{\alpha_s}{2\pi} \left(P'_{(n)} \otimes \overline{M}_n \right) (\Delta x, Q^2)$$

where modified splitting function is

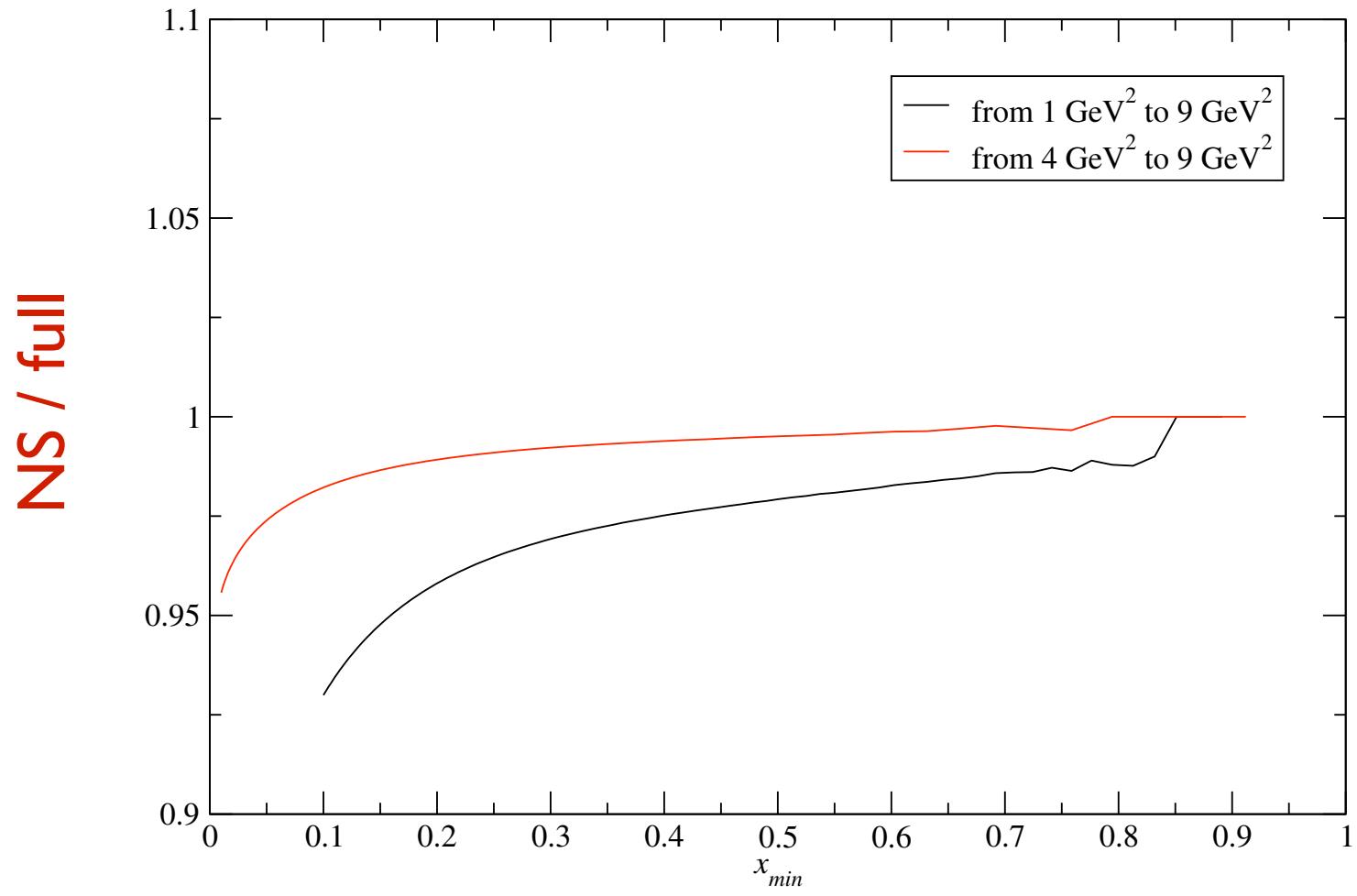
$$P'_{(n)}(z, \alpha_s) = z^n \ P_{NS,S}(z, \alpha_s)$$

- can follow evolution of specific resonance (region) with Q^2 in pQCD framework!
- suitable when complete moments not available

Truncated moments

- truncated moment evolution equations exist for singlet (S) and nonsinglet (NS) separately
- for analysis of data, do not know much of experimental structure function is NS and how much is S
 - for lowest ($n=2$) truncated moment, assumption that total \approx NS is good to few % for $x_{\min} > 0.2$
 - for higher moments, small- x region is further suppressed, so that NS is a very good approximation to total

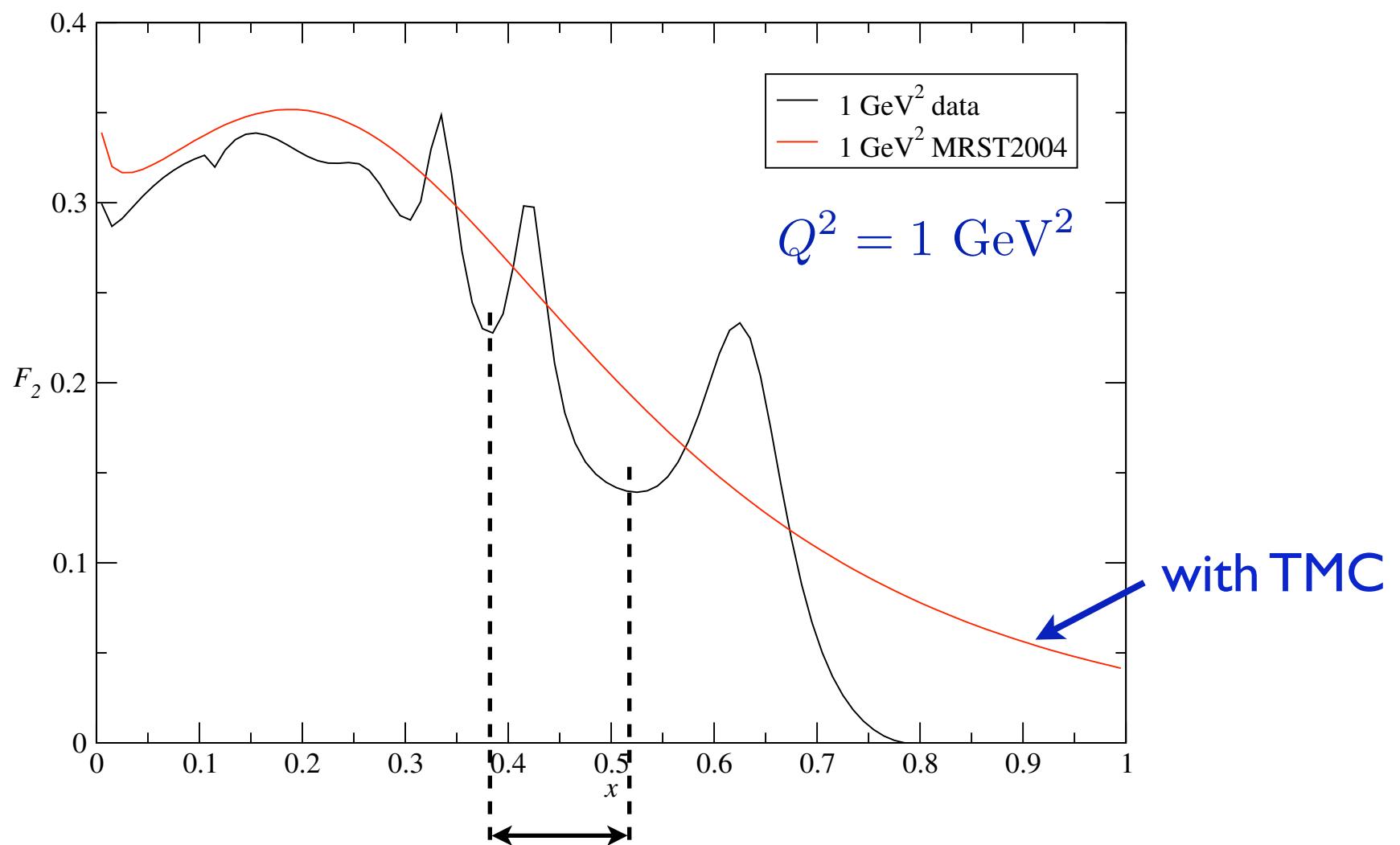
$n = 2$ truncated moment of F_2^p



$$\Delta x = [x_{min}, 1]$$

Psaker et al. (2007)

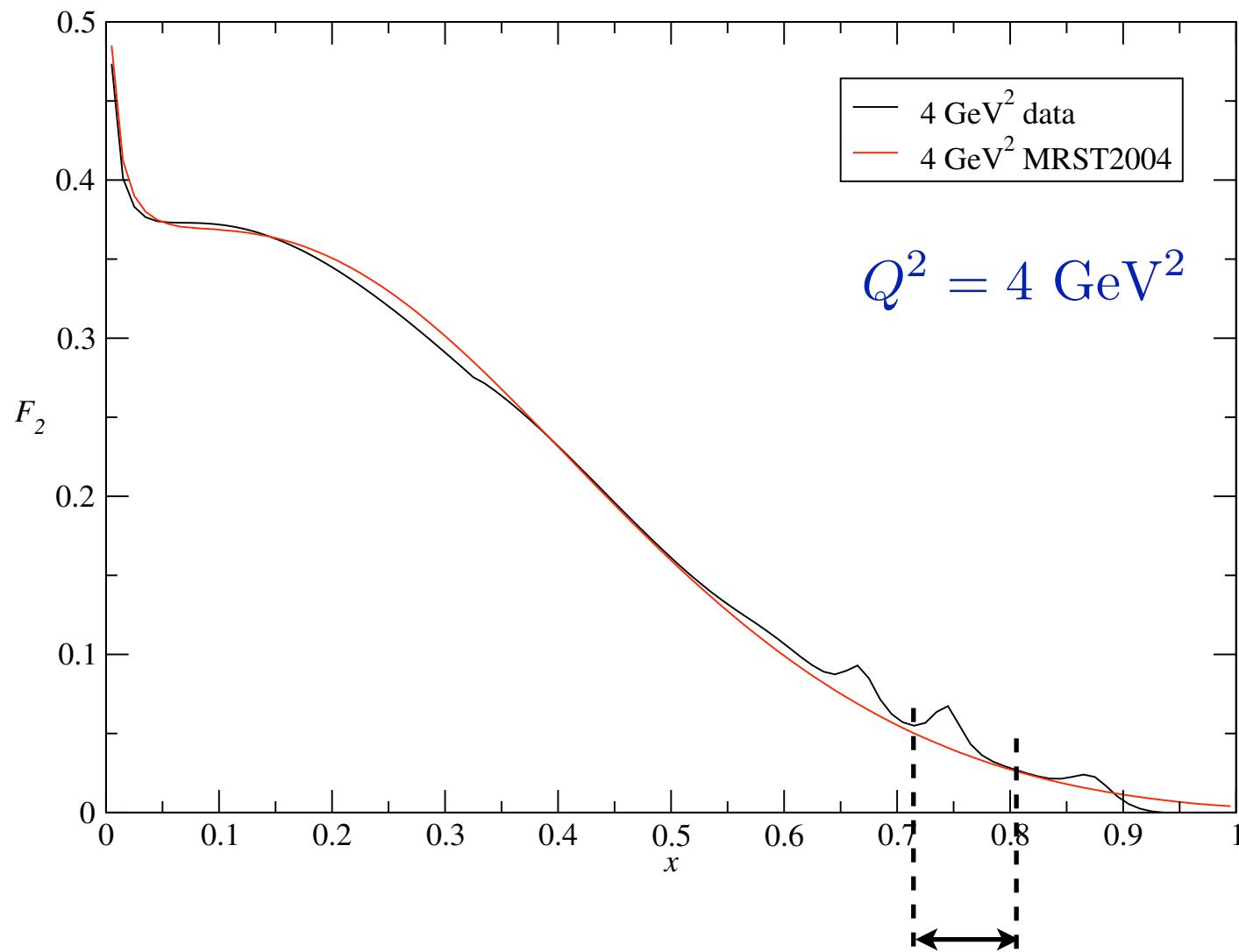
Parameterization of F_2^p data



how much of this region
is leading twist ?

Psaker et al. (2007)

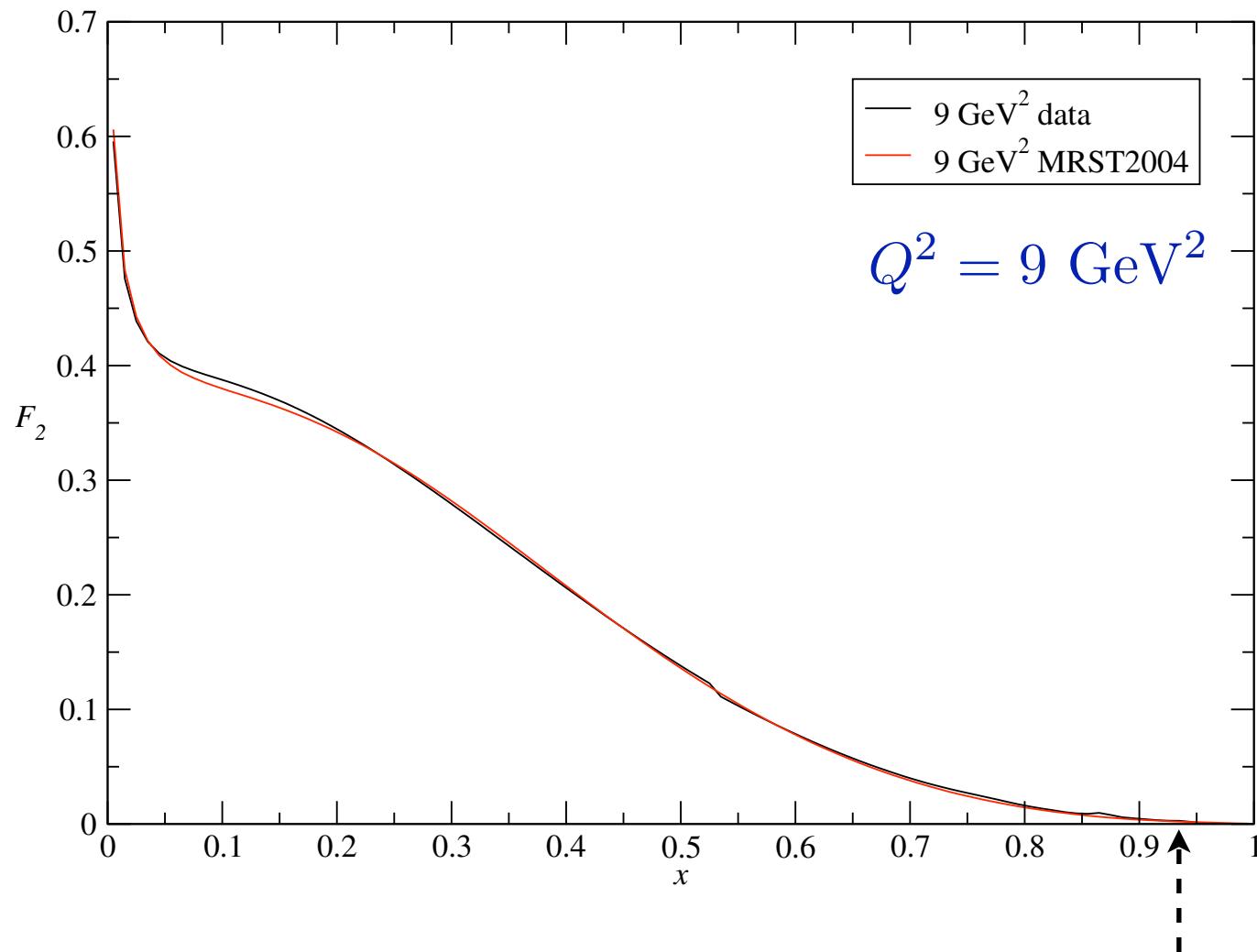
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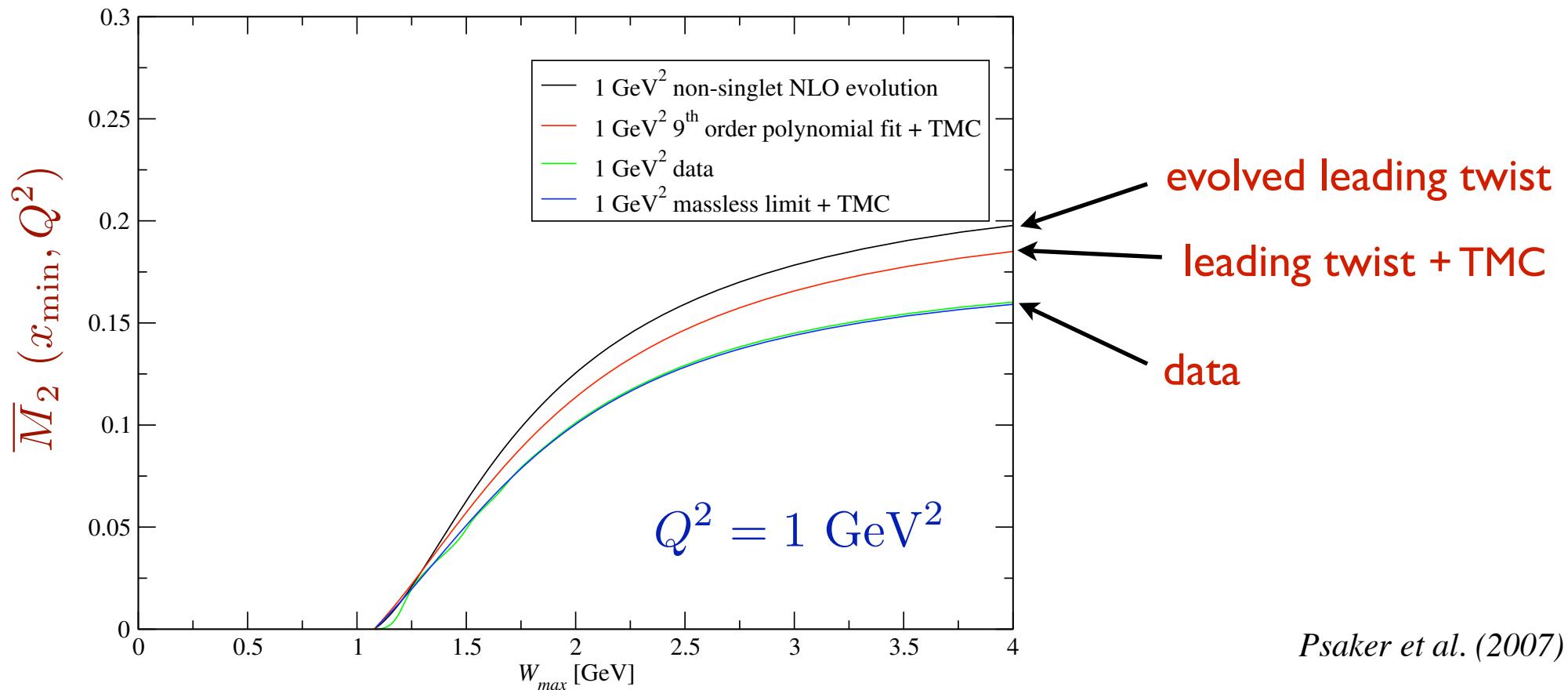


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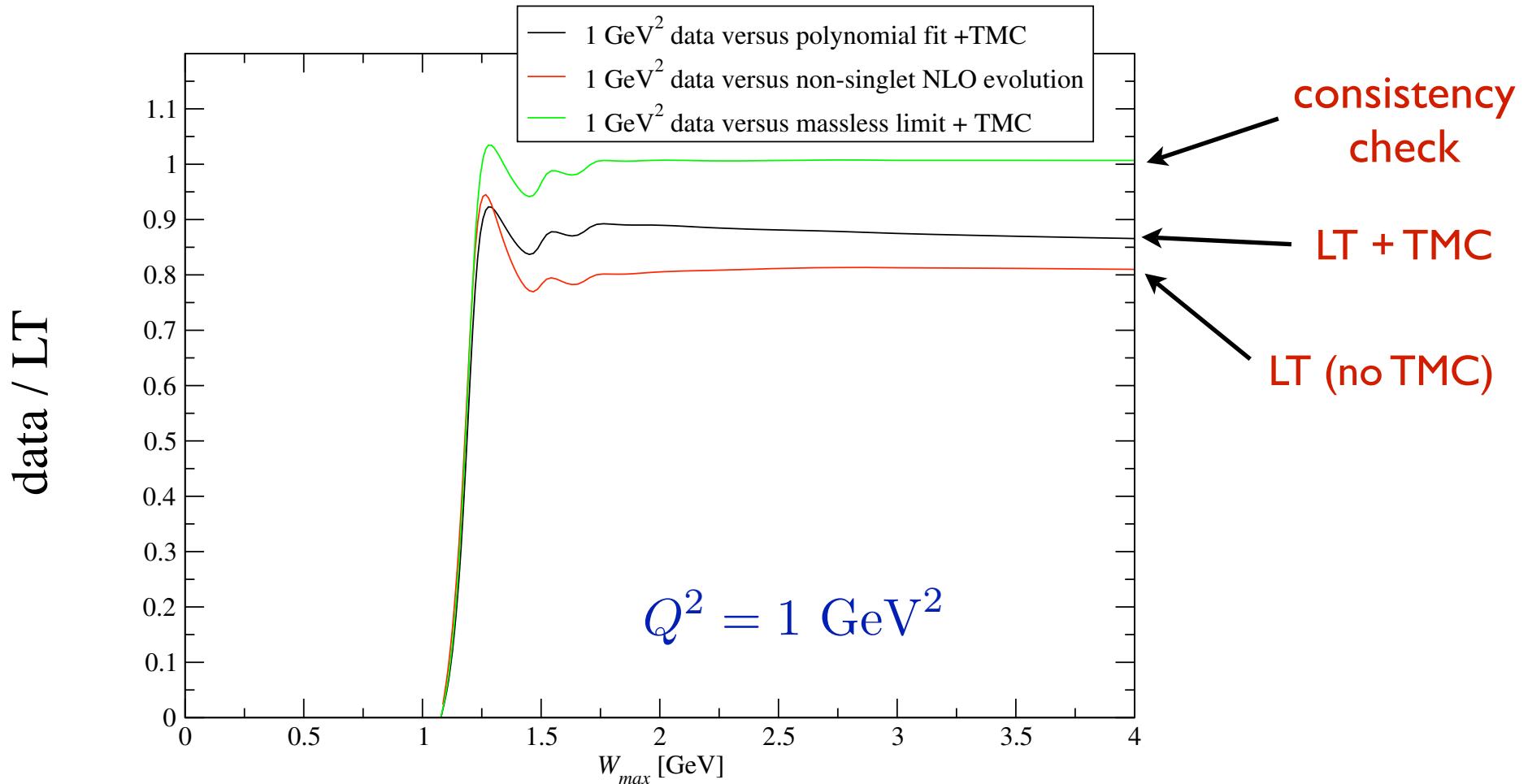
Analysis of Hall C data

- assume data at highest Q^2 ($Q^2 = 9 \text{ GeV}^2$) is entirely leading twist
- evolve (as NS) fit to data at $Q^2 = 9 \text{ GeV}^2$ down to lower Q^2
 - apply TMC, and compare with data at lower Q^2



Analysis of Hall C data

ratio of data to leading twist



Analysis of Hall C data

- consider individual resonance regions:

$$W_{\text{thr}}^2 < W^2 < 1.9 \text{ GeV}^2 \quad \text{“}\Delta(1232)\text{”}$$

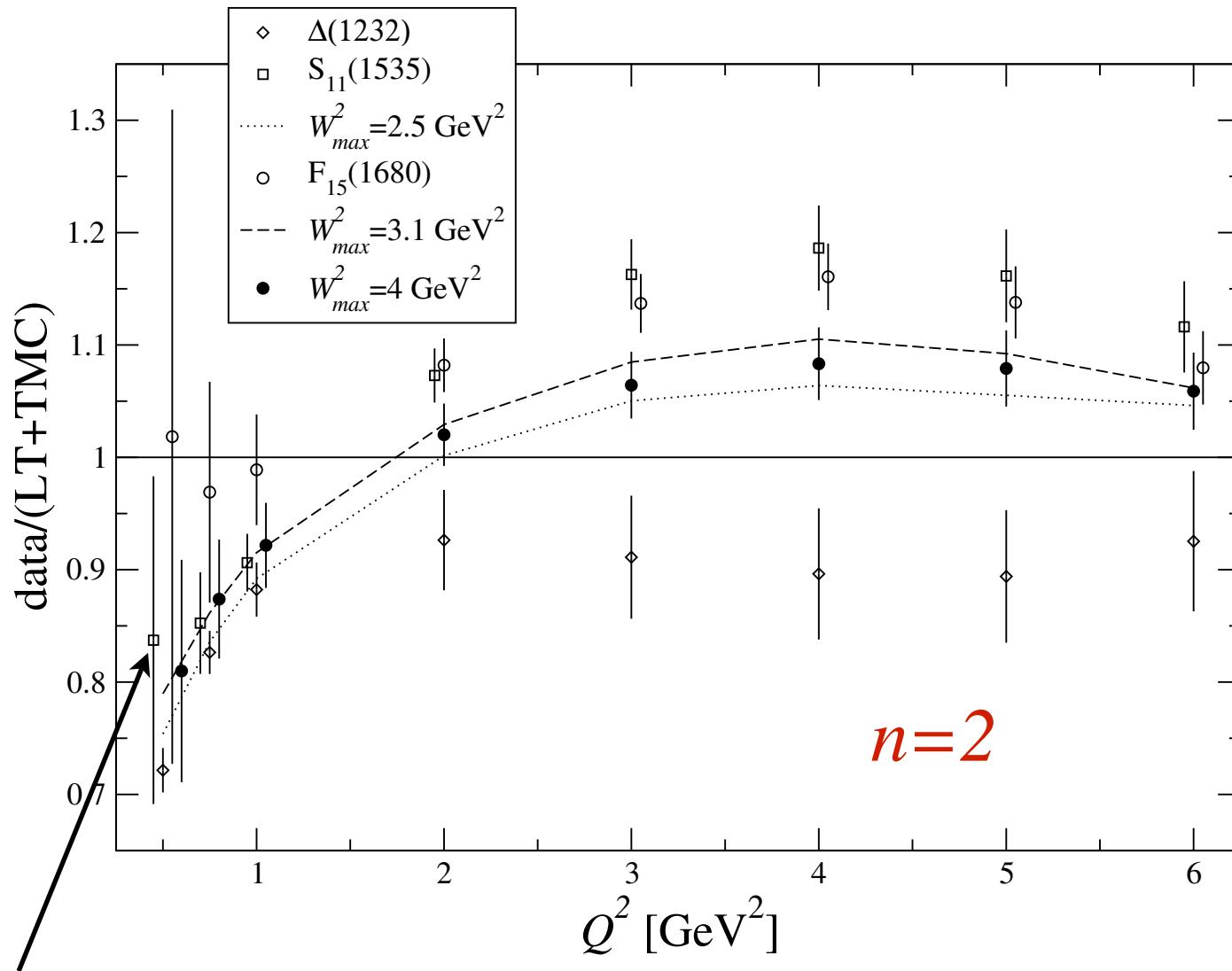
$$1.9 < W^2 < 2.5 \text{ GeV}^2 \quad \text{“}\textit{S}_{11}(1535)\text{”}$$

$$2.5 < W^2 < 3.1 \text{ GeV}^2 \quad \text{“}\textit{F}_{15}(1680)\text{”}$$

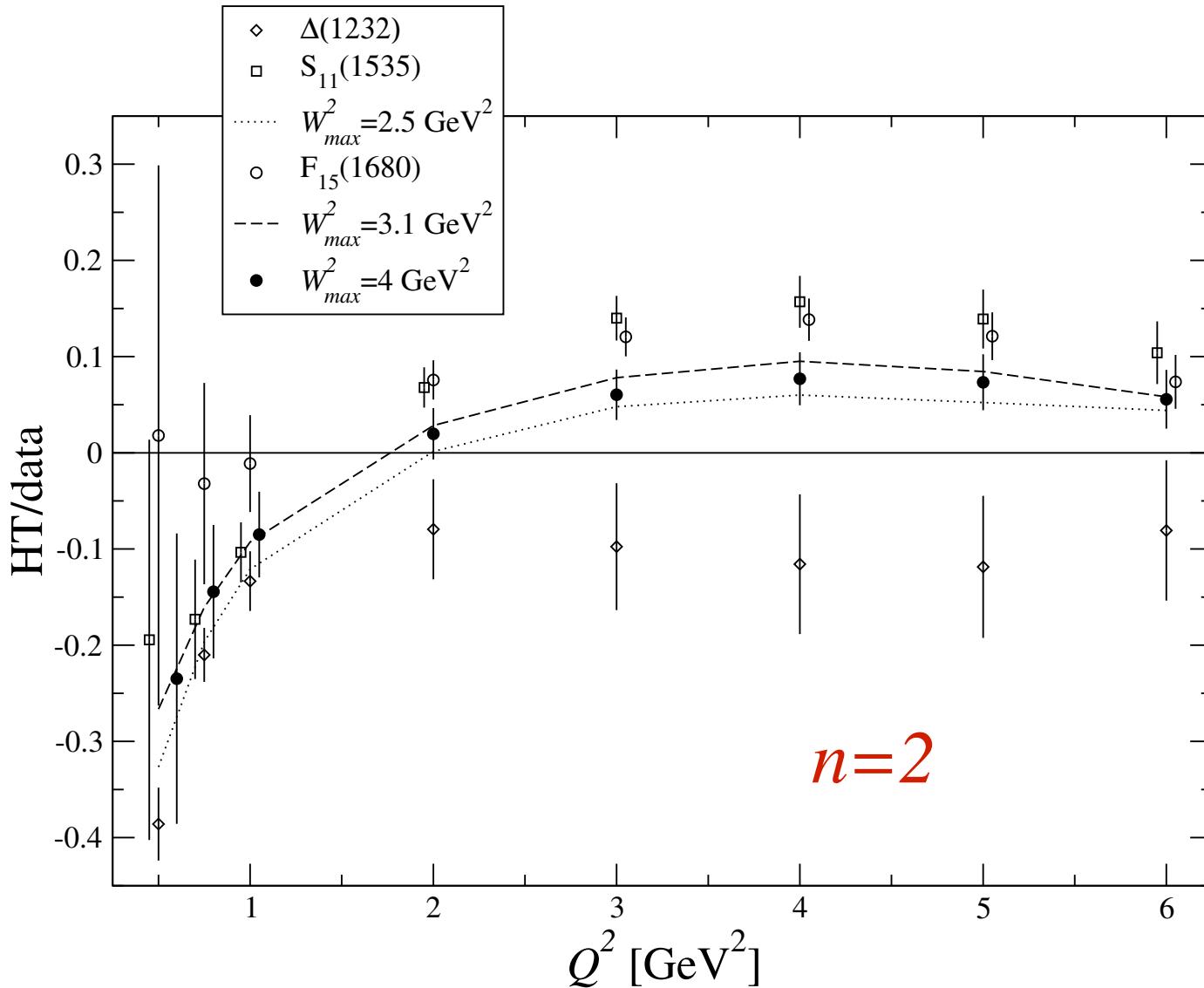
as well as total resonance region:

$$W^2 < 4 \text{ GeV}^2$$

Analysis of Hall C data

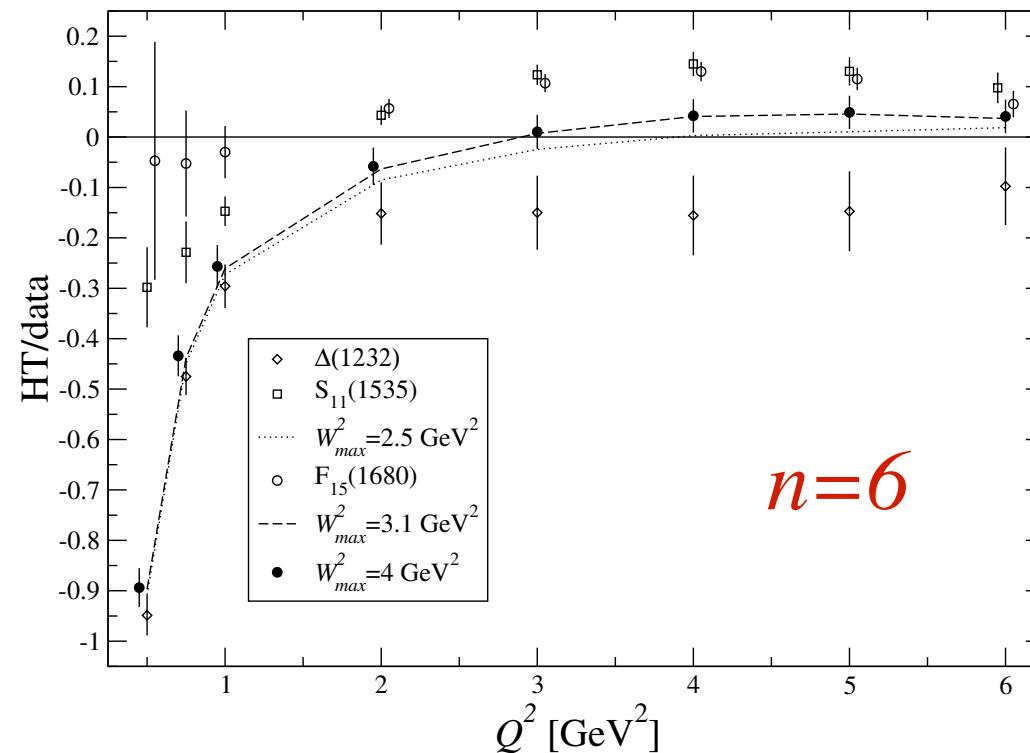
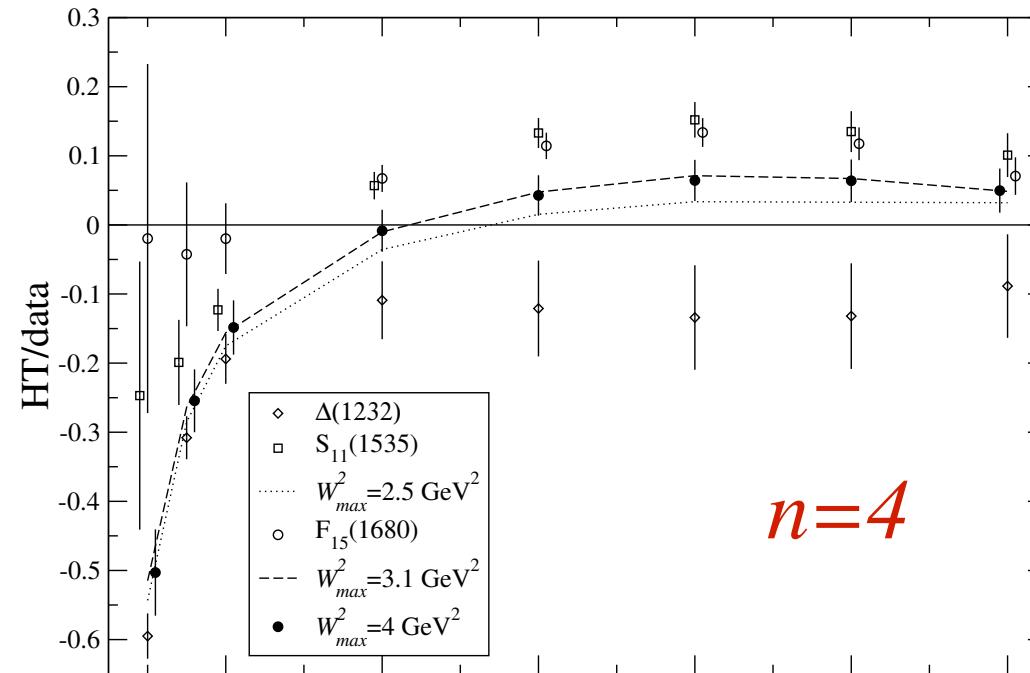


method breaks down for
low x (high W) at low Q^2



→ higher twists < 10% for $Q^2 > 1 \text{ GeV}^2$

■ higher moments



Summary

- Remarkable confirmation of quark-hadron duality in structure functions
 - higher twists “small” down to low Q^2 ($\sim 1 \text{ GeV}^2$)
- Local (elastic) duality
 - constraints on nuclear EMC effect and medium modified form factors
- Truncated moments
 - firm foundation for study of local duality in QCD
 - method can be applied to nuclear cross sections, relating nuclear structure functions to transition form factors