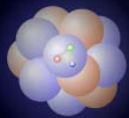


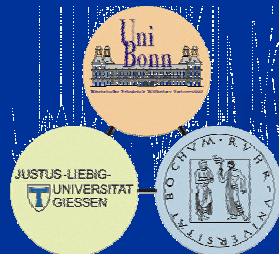
Hadrons in Nuclei

Predictions and Observables

Ulrich Mosel



Institut für
Theoretische Physik



Why study in-medium hadrons?

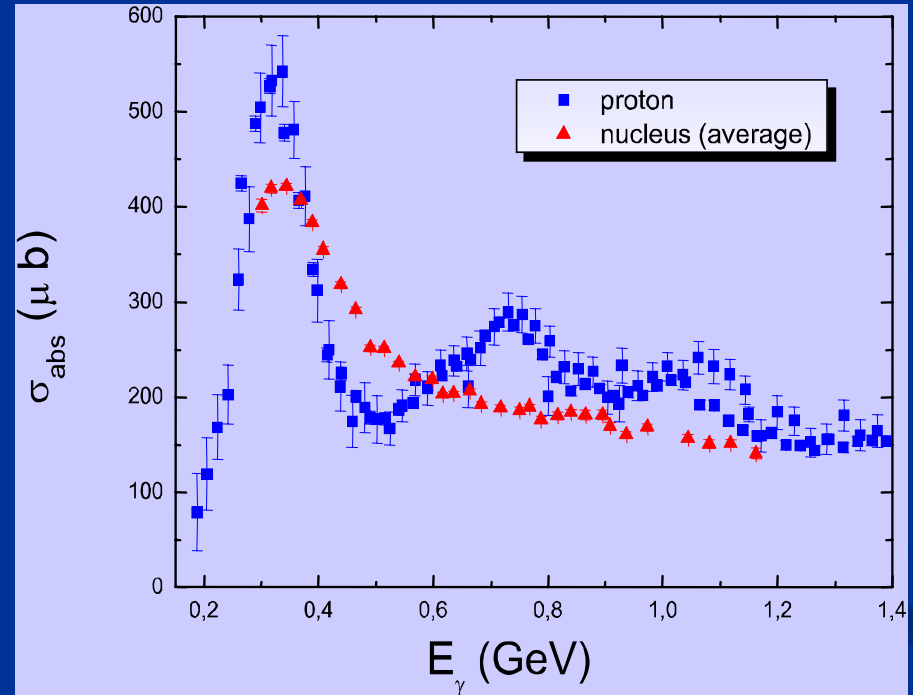
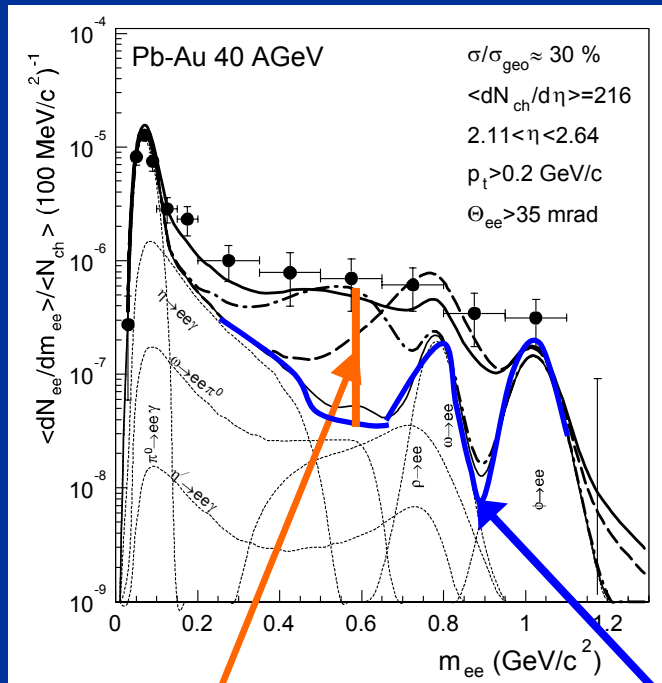
- Nucleon-meson physics in medium
 - ➔ Link to Nucleon-Resonance Properties
- Density may restore symmetries of QCD
 - Drop of condensates
 - Degeneracy of chiral partners
- Nucleus as a ‚microdetector‘:
 - access to production- and form.-times in quark-fragmentation, color transparency
- In-medium properties may signal exotic states of nuclear matter (e.g.: QGP)
 - ➔ need baseline effects in normal nuclear matter

In-medium changes: experiment 2000

Evidence for QGP at Cern
Invariant ($e+e^-$) mass spectrum

Total photoabsorption cross section

explained by spectral change of ρ meson in dense, (hot) matter



CERES effect

Cocktail plot of free sources
PN12 2004



In-medium changes: experiment 2003

- evidence for QGP at RHIC:
Jet Quenching
- Related Photonuclear Effect at HERMES:
Jet Quenching

Equilibrium vs. Nonequilibrium

or: why it is better to work with microscopic probes on nuclei

- URHIC signal sums over very different stages of reaction:
 - Start: highly non-equilibrated, very high density
 - End: equilibrated, high temperature, low density
 - In *thermal* equilibrium all infos about interactions are lost
- Theories assume equilibrium for calculating hadronic in-medium properties
- Photonuclear reactions much closer to (cold) equilibrium, low but constant density throughout

Observables

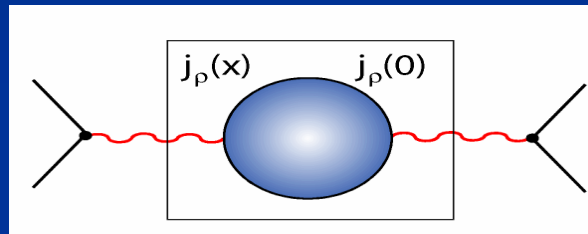
- Experimental data:
incoherent photo- and electroproduction of hadrons on nuclei from 100 MeV (MAMI, ELSA) over few GeV (JLAB) to ~ 20 GeV (HERMES)
- Experiment:
 - weak ISI $\rightarrow \gamma$ best
 - FSI
 - Hadronic, e.g. $\phi \rightarrow K^+K^-$ difficult
 - Semihadronic, e.g. $\omega \rightarrow \pi^0 \gamma$ possible
 - Electromagnetic, e.g. $\omega \rightarrow e^-e^-$ best

Observables

- FSI important:
 - *attenuation* of primary particle
can be treated by Glauber (standard)
 - *Side feeding* from other channels,
e.g. $\gamma N \rightarrow \pi N$, $\pi N' \rightarrow K \Lambda$
must be treated by Coupled Channels (BUU)
Method

Vector Mesons: Hadronic Tensor

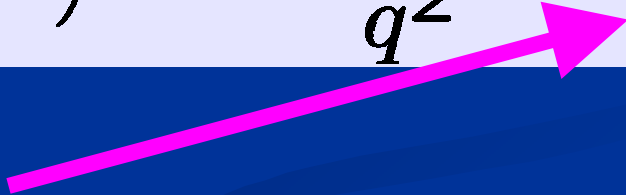
- determines electromagnetic coupling to hadrons:



$$\begin{aligned}\Pi^{\mu\nu} &\sim \int d^4x e^{iqx} \langle 0|T [j^\mu(x)j^\nu(0)] |0\rangle \\ &= (q^2 g^{\mu\nu} - q^\mu q^\nu) \Pi(q^2)\end{aligned}$$

Vector Mesons: Hadronic Tensor

Total cross section for hadron production:

$$\sigma(e^+e^- \rightarrow \text{hadrons}) = -\frac{4\pi\alpha}{q^2} \Im \Pi(q^2)$$


Experimentally known for free hadrons
contains spectral information about hadrons

Vector Mesons: Hadronic Tensor

Vector Meson Dominance

Photon \cong Vector meson ($J^{\pi}=1^{-}$)

$$\begin{aligned}\Pi^{\mu\nu}(q) &\sim \int d^4x e^{iqx} \langle 0 | T [j^\mu(x) j^\nu(0)] | 0 \rangle \\ &= \int d^4x e^{iqx} \langle 0 | T [\rho^\mu(x) \rho^\nu(0)] | 0 \rangle \frac{m_\rho^2}{g_\rho^2} = \frac{m_\rho^2}{g_\rho^2} D_\rho^{\mu\nu}(q)\end{aligned}$$

ρ meson propagator

Spectral Function

$$\mathcal{A}(\omega, \vec{q}) = -\frac{1}{\pi} \Im D(\omega, \vec{q}) \sim \Im \Pi(\omega, \vec{q})$$

- Two ways:
 - Determine constraints on in-medium Π from QCD sum rules
 - Calculate in-medium Π in hadronic model

QCD Sum Rule

only ,clean‘, but indirect connection between hadronic and quark world

Compare spectral function in time-like region with OPE of current-correlator for space-like distances

$$\frac{Q^2}{\pi} \int_0^\infty ds \frac{\Im \Pi(s)}{s(s+Q^2)} = -\frac{1}{8\pi^2} \left(1 + \frac{\alpha_s}{\pi}\right) \ln \frac{Q^2}{\Lambda^2} + \frac{m_q \langle \bar{q}q \rangle}{Q^4} + \frac{1}{24} \frac{\langle \frac{\alpha_s}{\pi} G^2 \rangle}{Q^4} + \frac{\langle (\bar{q}q)^2 \rangle}{Q^6} + \dots$$

- Lhs dominated by soft scale $\sim m_\rho$
- Rhs separates hard scale $\sim Q^2$ from soft scale (condensates)

- Model condensate (ρ)
- Parametrize $\Im \Pi$ in terms of few parameters, to be extracted from sum rule

$$\Im \Pi(s) = \pi F \frac{S(s)}{s} \Theta(s_0 - s) + \frac{1}{8\pi} \left(1 + \frac{\alpha_s}{\pi} \right) \Theta(s - s_0)$$

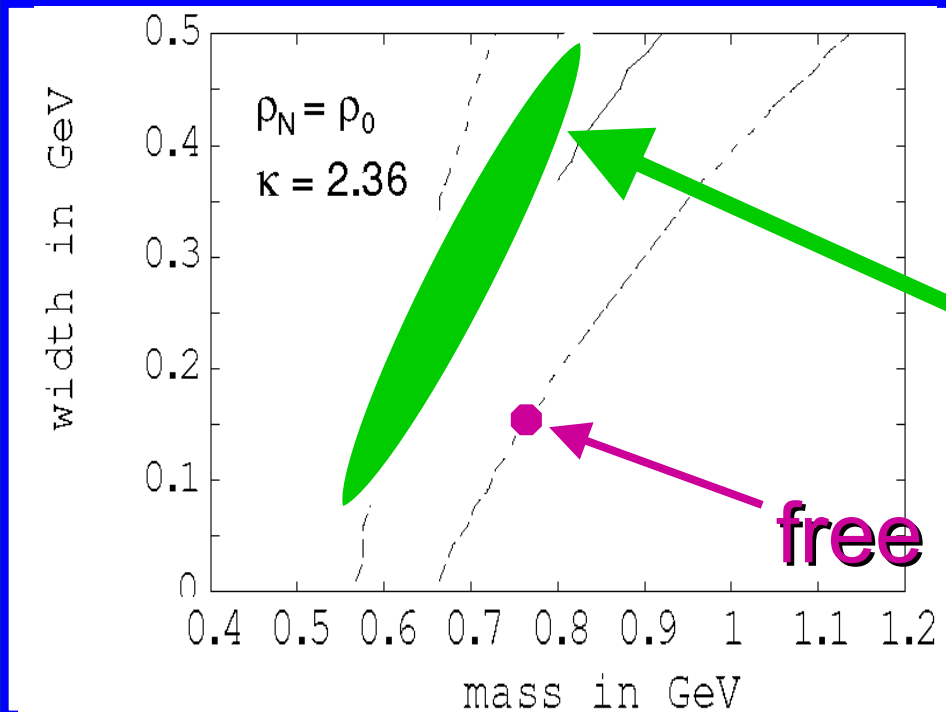
- Hatsuda-Lee (1991):

$$S(s) = \delta(s - m_\rho^2) \implies m_\rho$$

- Leupold-Peters-Mosel (1998):

$$S(s) = \frac{1}{\pi} \frac{\sqrt{s} \Gamma_\rho(s)}{(s - m_\rho^2)^2 + s \Gamma_\rho(s)^2} \implies (m_\rho, \Gamma_\rho)$$

ρ spectral function in medium



QCDSR-allowed (Γ, m)
at saturation density

free ρ meson

QCD Sum Rules
provide constraints
for, but do not fix in-
medium hadron props

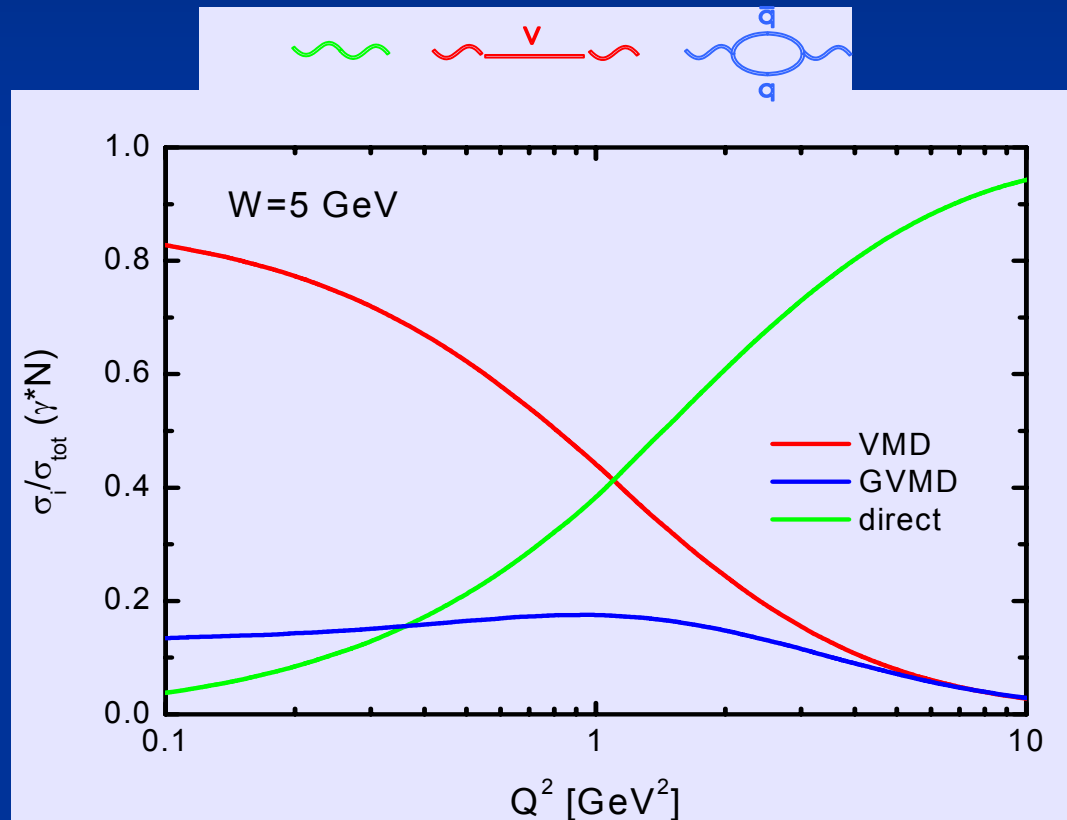
Leupold et al, Phys.Rev.C58:2939-2957,1998

Theoretical Method

- Entrance Channel:
 - Quantum Coherence: Shadowing
- Primary Production
 - Resonance Decay or String-Fragmentation (PYTHIA)
- Exit Channel:
 - Incoherent Final State Interactions (Absorption + Scattering + Side-Feeding), Propagation with self-energies and interactions through resonances or fragmentation

Theoretical Method: Coherence in Entrance Channel

■ Hadronic Structure of the Photon

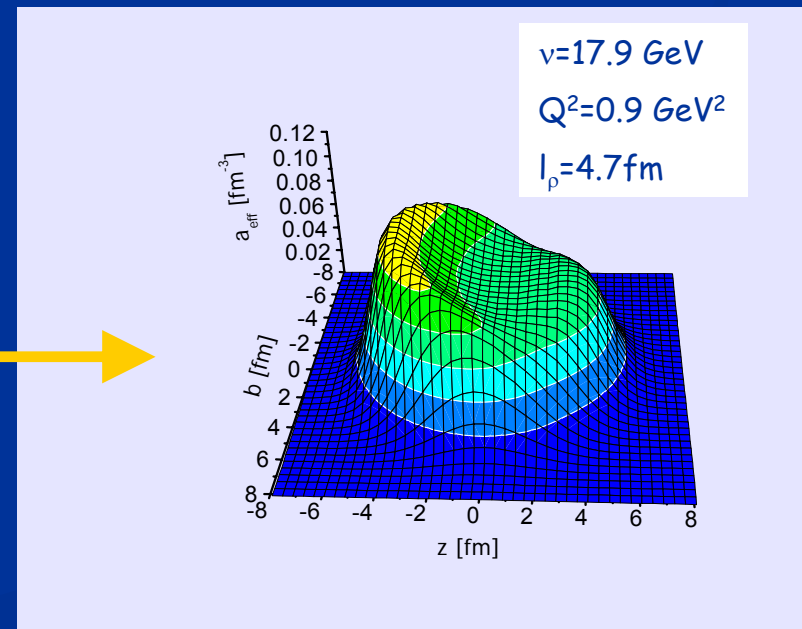


Theoretical Method: Coherence in Entrance Channel

- Coherence Length: $l_\rho = \frac{2v}{Q^2 + m_V^2}$
 - Distance that the photon travels as V meson
- Treat by Glauber

Vector Meson Components in
the Nucleus Shadowed

γ^*



Theoretical Method: BUU CC transport model

- Same Code for Same Physics
 - Photonuclear reactions $\gamma + A$
 - Hadronic reactions $\pi, \rho + A$
 - Heavy-ion reactions $A + A$
- Resonance and Continuum Region treated
 - Resonance Decays from data
 - Continuum Decays from String Fragmentation

Theoretical Method: BUU CC Transport Model

$$\left(\frac{\partial}{\partial t} + (\nabla_{\vec{p}} H) \nabla_{\vec{r}} - (\nabla_{\vec{r}} H) \nabla_{\vec{p}} \right) f_i(\vec{r}, \vec{p}, t) = I_{\text{coll}} [f_1, \dots, f_i, \dots, f_M]$$

set of BUU equations coupled via I_{coll} and mean field

f_i : phase space density

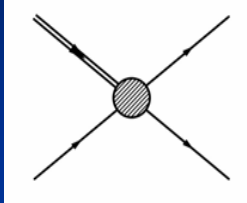
H : Hamilton function

$$H = \sqrt{(\mu + U_s)^2 + \vec{p}^2}$$

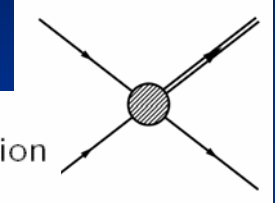
collision integral accounts for changes in f_i due to 2 particle collisions: creation, annihilation, elastic scattering (Pauli blocking for fermions)

1. products of $\gamma^* A$ reaction need not be created in primary $\gamma^* N$ reaction
2. In-medium changes can be modelled in H (selfenergies) and in I_{coll} (reaction rates, form. times, prehadron cross sections)
3. Experimental acceptance can be simulated event by event

Nucleon-Spectral Functions



$$a(\omega, p) = \frac{\Gamma(\omega, p)}{(\omega - \frac{p^2}{2m_N} - \text{Re}\Sigma(\omega, p))^2 + \frac{1}{4}\Gamma^2(\omega, p)}$$



Assume **constant Scattering Amplitude** = Short-Range Correlation

$$\Sigma^>(\omega, p) = -4i \frac{|\mathcal{M}|^2}{(2\pi)^6} \int d\omega_3 \int d\omega_2 \int dp_3 p_3^2 \int dp_2 p_2^2 \frac{d\cos\vartheta_2}{p_{\text{tot}} p_3} \int dp_4 p_4 \\ \times a(\omega_2, p_2) f(\omega_2, p_2) a(\omega_3, p_3) (1 - f(\omega_3, p_3)) a(\omega_4, p_4) (1 - f(\omega_4, p_4))$$

$$\Sigma^<(\omega, p) = 4i \frac{|\mathcal{M}|^2}{(2\pi)^6} \int d\omega_3 \int d\omega_2 \int dp_3 p_3^2 \int dp_2 p_2^2 \frac{d\cos\vartheta_2}{p_{\text{tot}} p_3} \int dp_4 p_4 \\ \times a(\omega_2, p_2) (1 - f(\omega_2, p_2)) a(\omega_3, p_3) f(\omega_3, p_3) a(\omega_4, p_4) f(\omega_4, p_4)$$

with $p_{\text{tot}} = |\vec{p} + \vec{p}_2|$

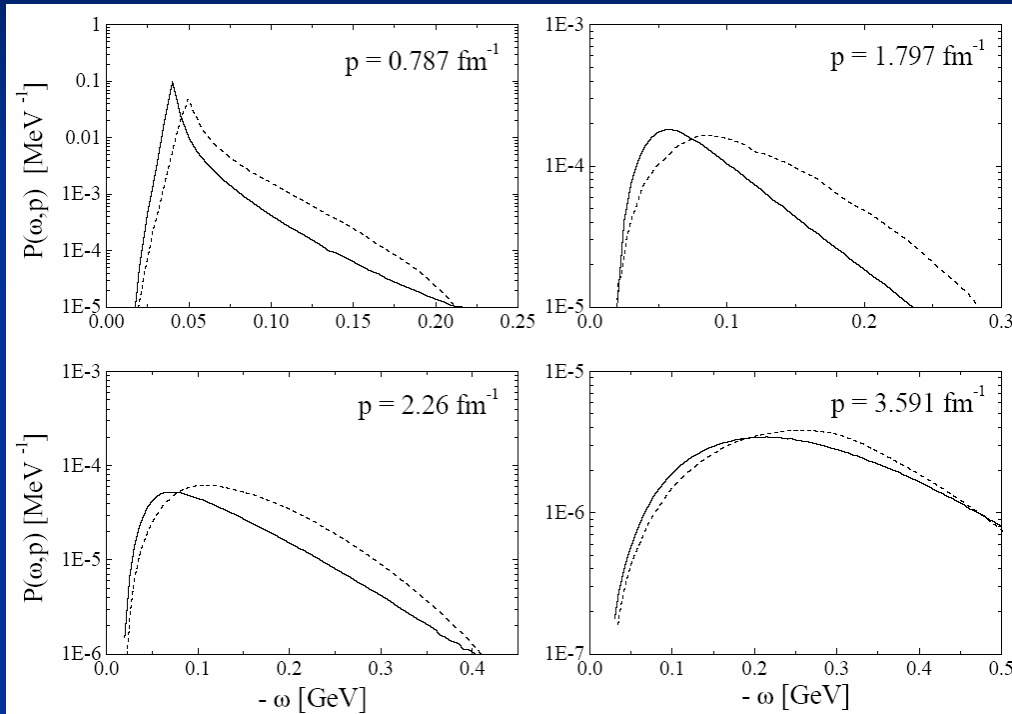
Selfconsistency Problem

$$\Gamma \rightarrow a \rightarrow \Sigma^{\langle \rangle} \rightarrow \Gamma$$

Collisional Width

$$\Gamma(\omega, p) = 2 \text{Im} \Sigma(\omega, p) = i(\Sigma^>(\omega, p) - \Sigma^<(\omega, p))$$

Nucleon-Spectral Functions

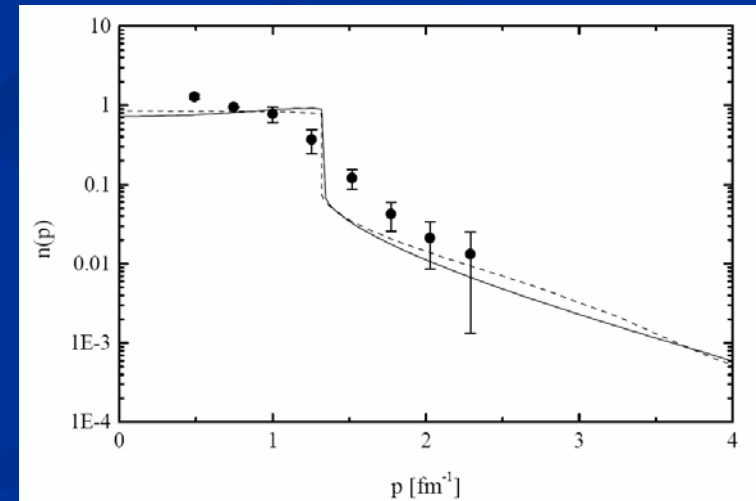


--- Benhar et al

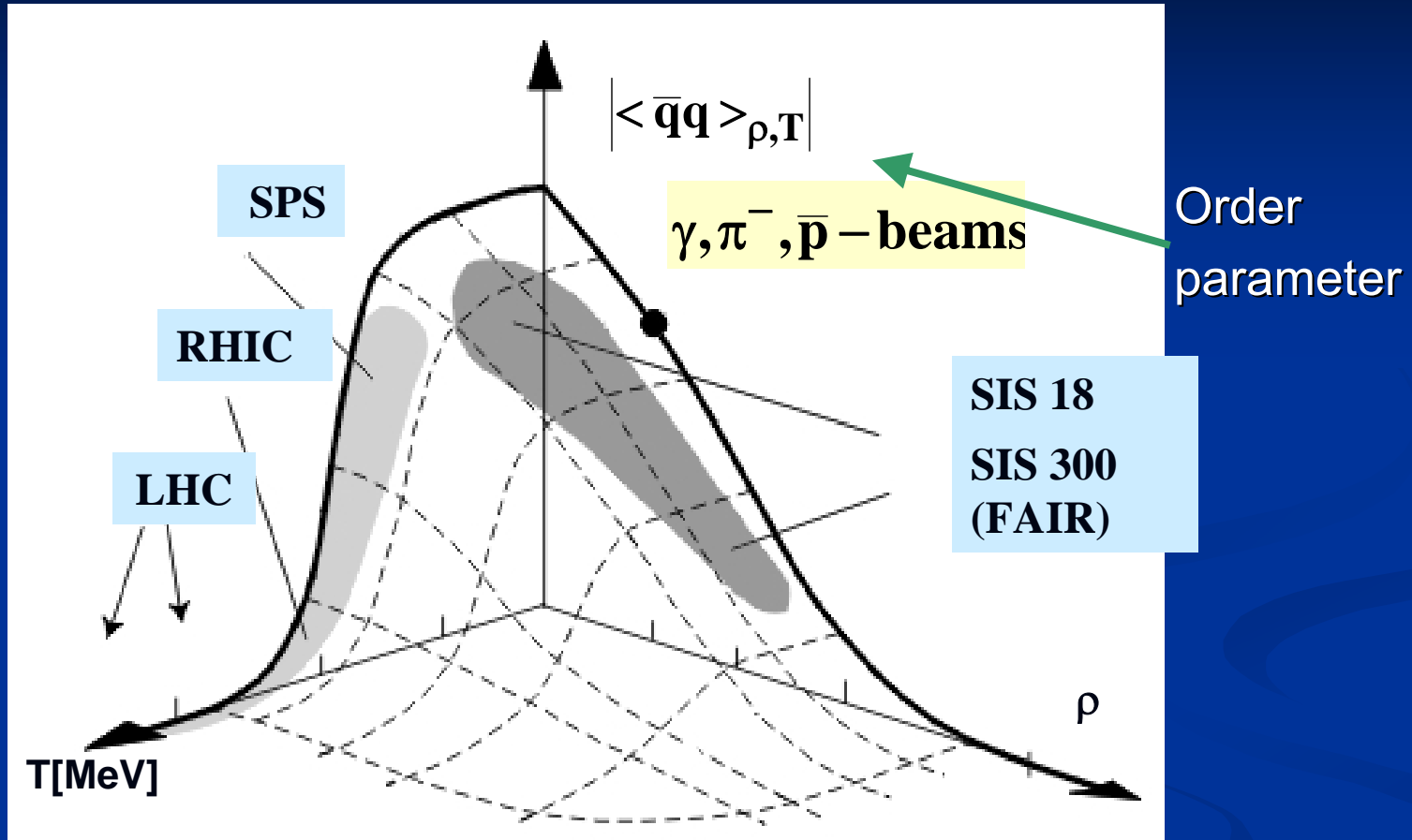
Data: Sick et al

Lehr et al, Nucl.Phys.A703:393,2002

Spectral Functions determined by coll. width
 → Phase space



Chiral Symmetry Restoration

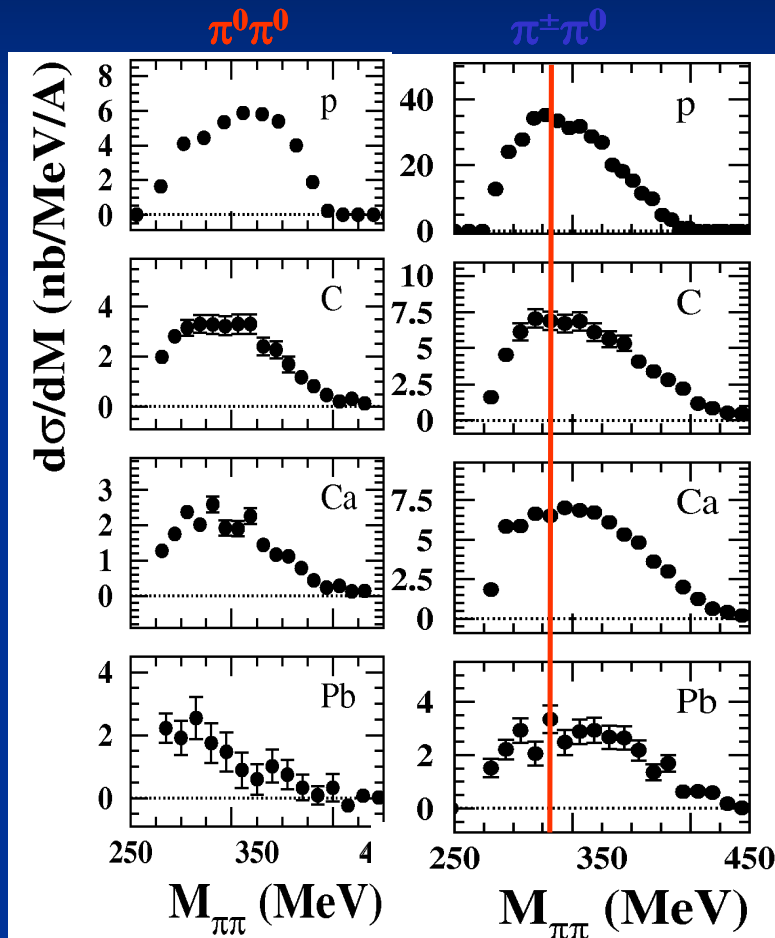


Klimt et al, 1990

Connection to Observables??

Chiral Symmetry Restoration

$2\pi^0$ Production on Nuclei



■ Expected:
 σ - π degenerate in chiral limit
 \rightarrow shift of σ strength
 to lower masses

TAPS data

$E_\gamma = 400 - 500$ MeV

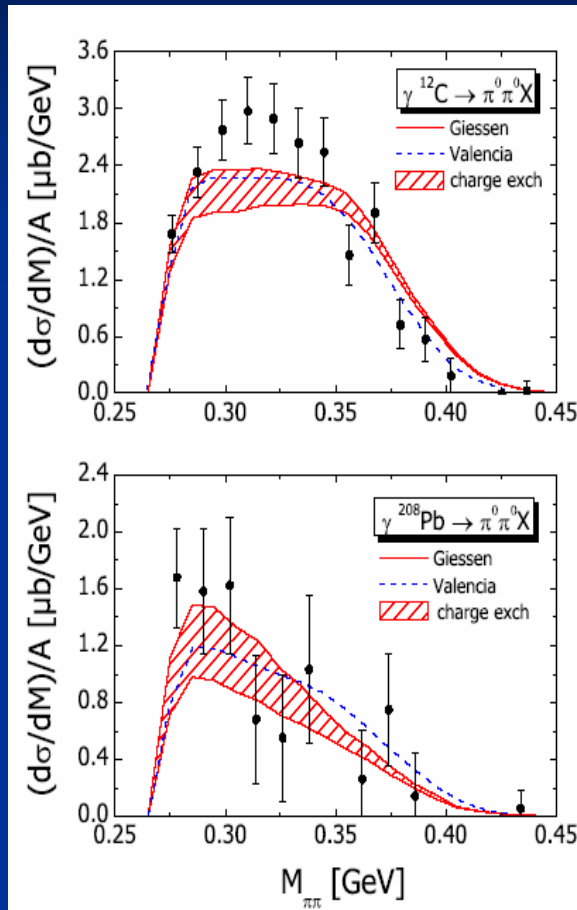
PN12 2004



Institut für
Theoretische Physik

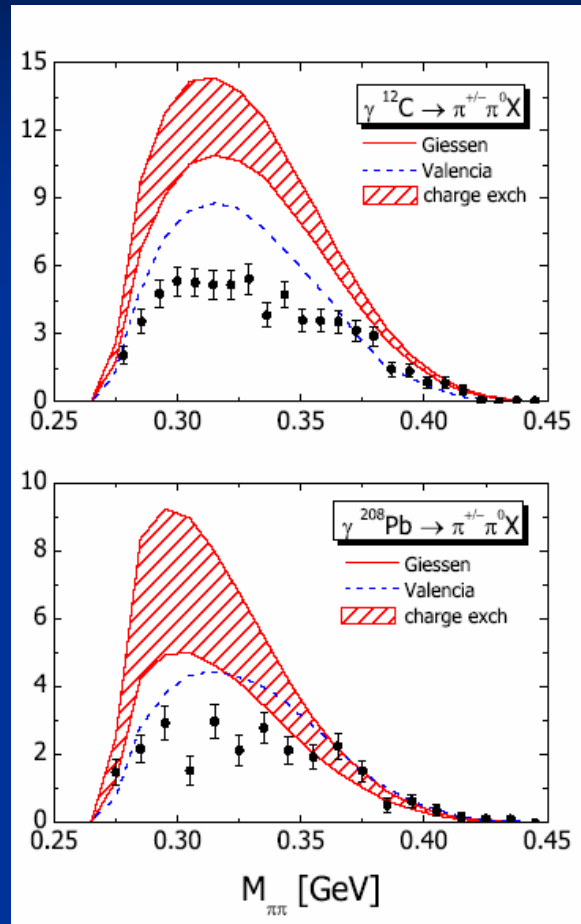
JUSTUS-LIEBIG-
UNIVERSITÄT
GIESSEN

$2\pi^0$ Production on Nuclei



TAPS Data

Calculations without
chiral restoration
explain $2\pi^0$ data!



P. Muehlich et al, Nucl.Phys.A703:393-408,2002

Problem with
 $\pi^0\pi^{+/-}$ channel

Need:
Single (γ, π^\pm)
data

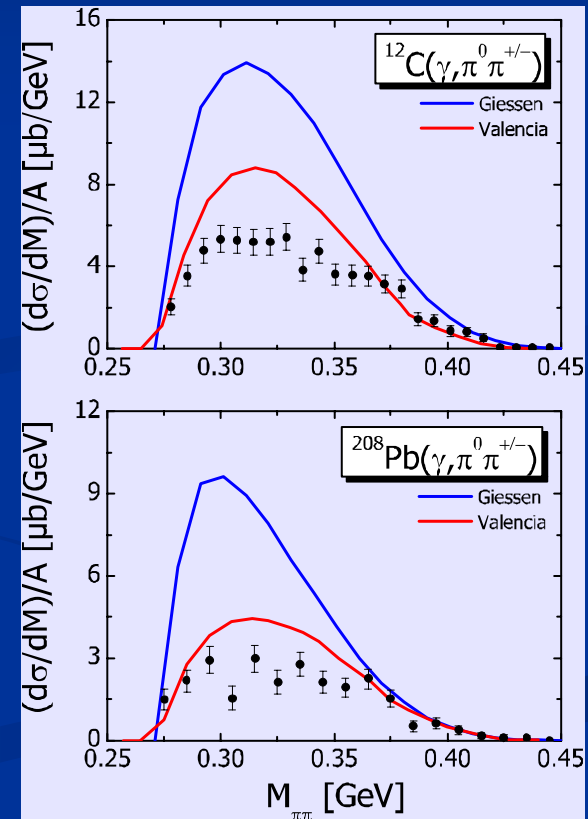
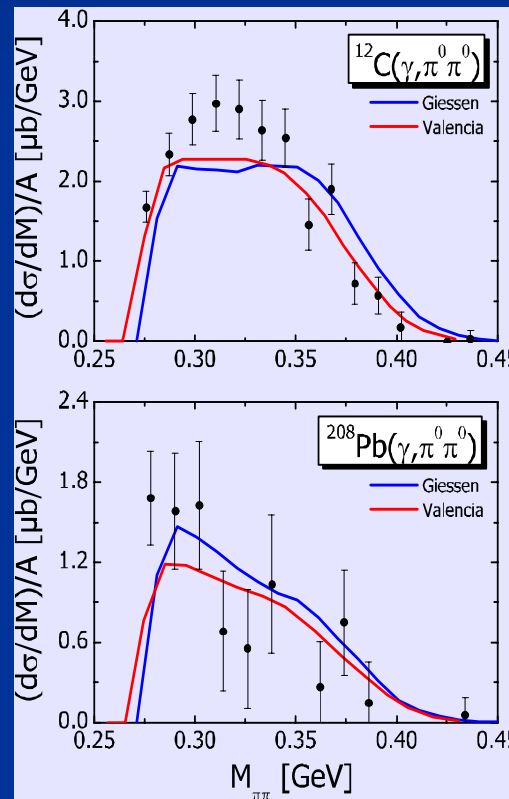
$2\pi^0$ Production on Nuclei

■ Chiral symmetry restoration??

Evidence
for lowering
of scalar
strength in
nuclei (??)

Giessen:
only FSI
Valencia:

$\pi\pi$ correlation



Need: Hadronic Model

- Propagator in Medium

$$D(\omega, \vec{q}) = \frac{1}{q^2 - m^2 - \Pi_{\text{vac}} - \Pi_{\text{med}}}$$

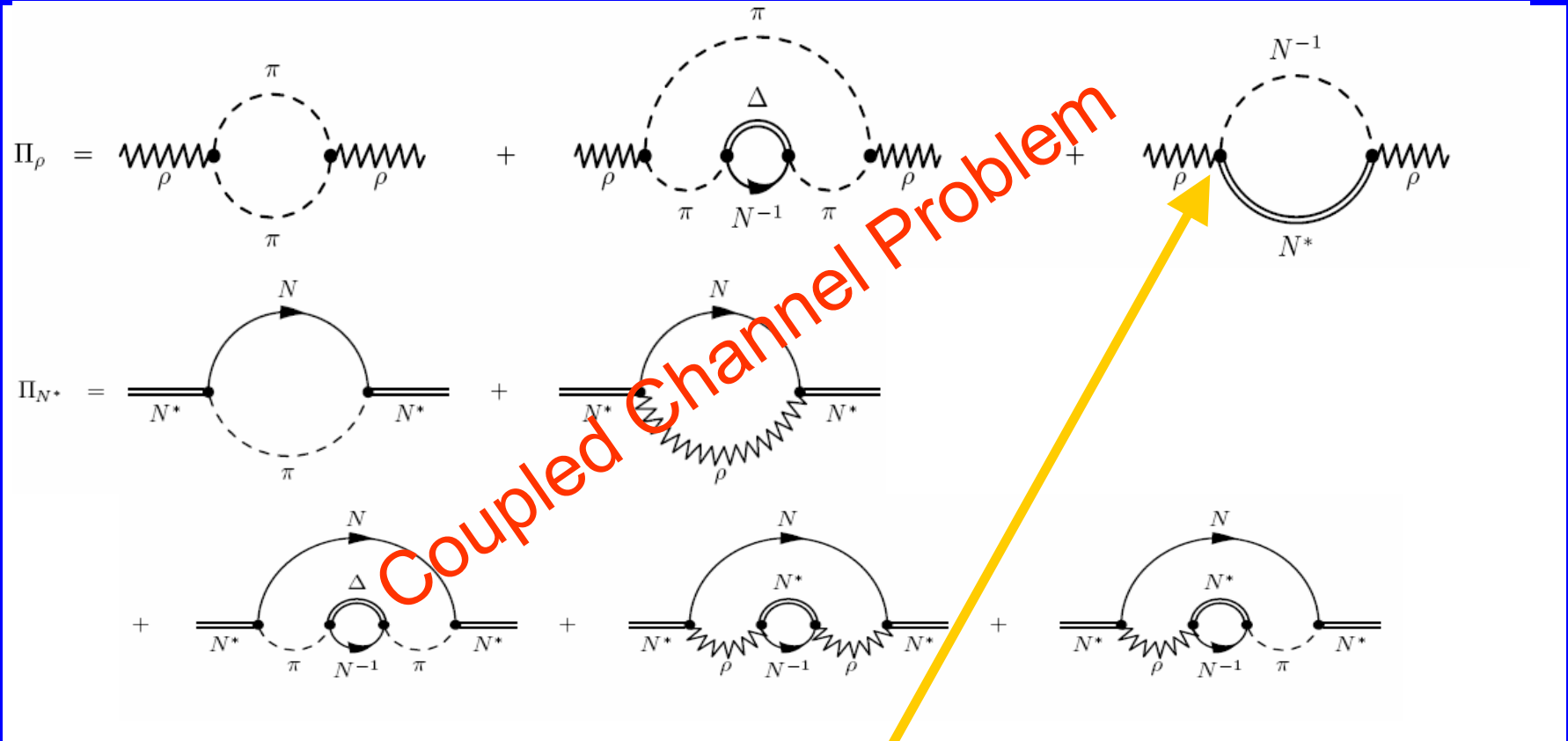
- Low density theorem: $\Pi_{\text{med}} = t \rho$

with (Optical Theorem): $\Im t \sim \sigma_{\text{tot}}$

- Mass shift, Collisional broadening:

$$\Delta m^2 = \Re \Pi_{\text{med}} \quad \Delta \Gamma \sim \Im \Pi_{\text{med}} \rightarrow v \sigma_{\text{tot}} \rho$$

ρ and N^* selfenergy in medium



Relevant Resonances:

$$D_{13}(1520), P_{13}(1720)$$

■ PDG:

- $D_{13} : \Gamma_{N\rho} \approx 20 \%$ **subthreshold !**
- $P_{13} : \Gamma_{N\rho} \approx 80 \%$

■ D_{13} subthreshold \rightarrow large coupling to $N\rho$

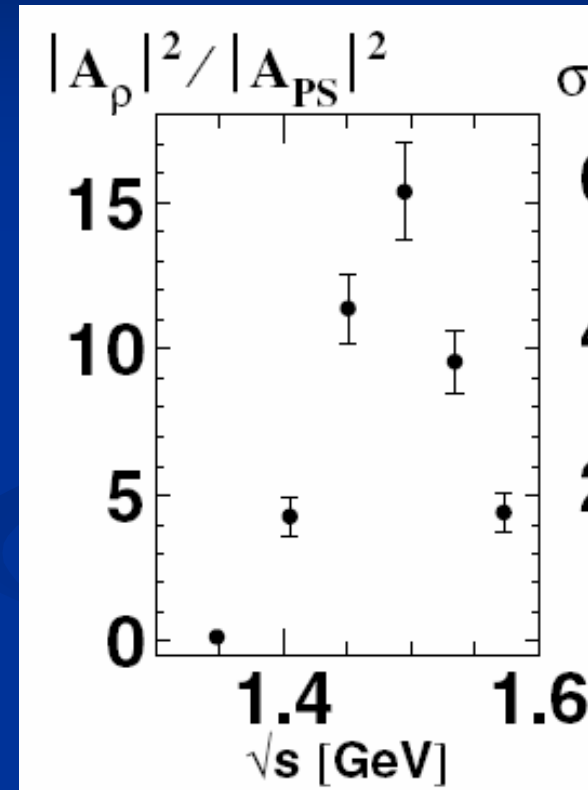
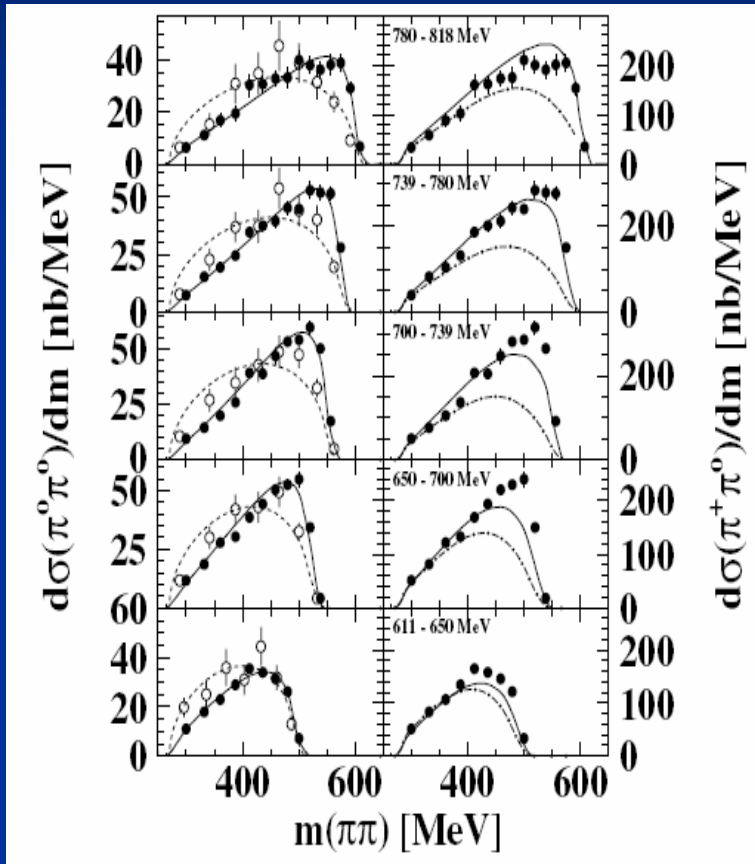
But: so far not directly been measured!

Only indirect infos from 2π invariant masses

$N(1520) \rightarrow N \rho$

Experiments: Daphne, TAPS at MAMI

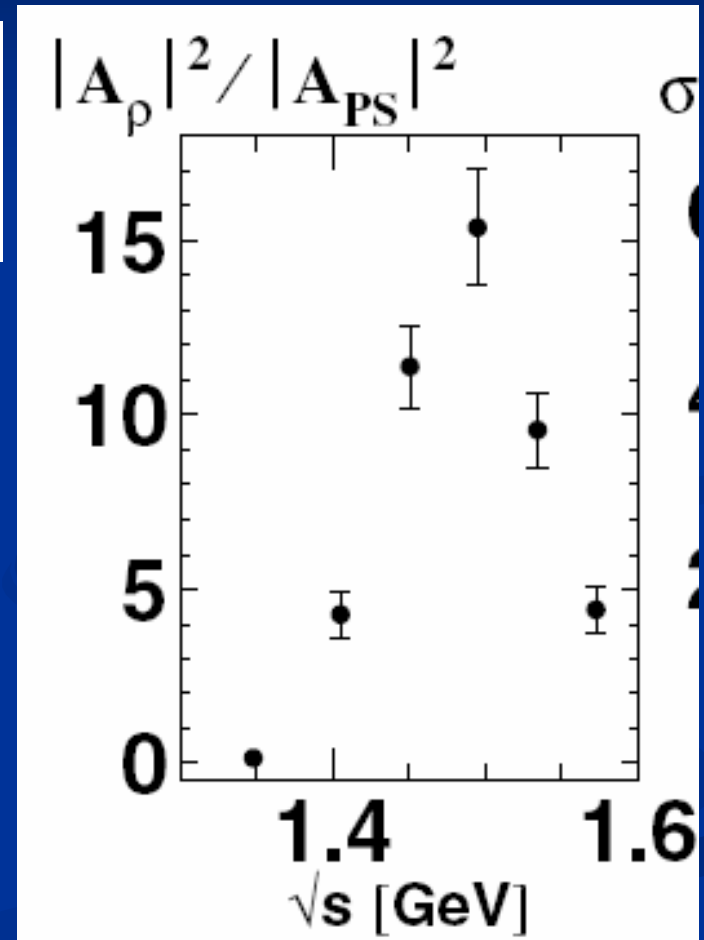
ρ - strength



PWA identification of ρ -meson still missing

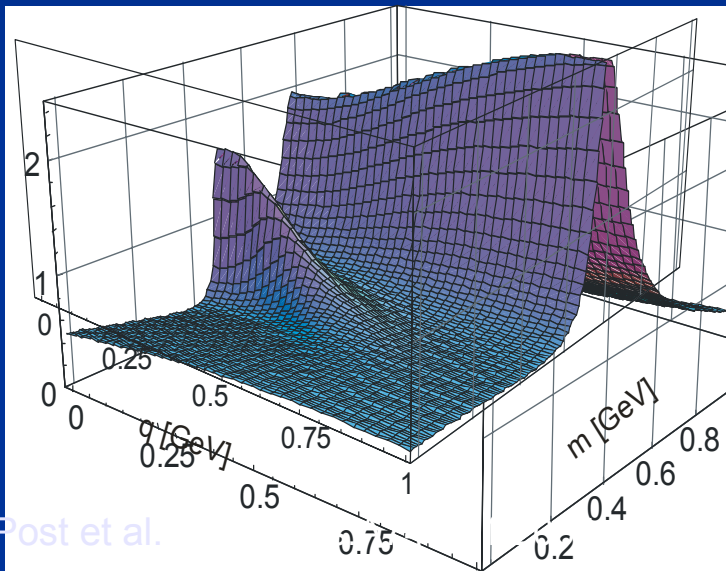
$N(1520) \rightarrow N \rho$

$$\frac{d\sigma}{dm} \sim |a(\sqrt{s}) + b(\sqrt{s})p_{\pi}(m_{\pi\pi})| \cdot |D_{\rho}(m_{\pi\pi})|^2 P_{\sqrt{s} \rightarrow \pi\pi N},$$



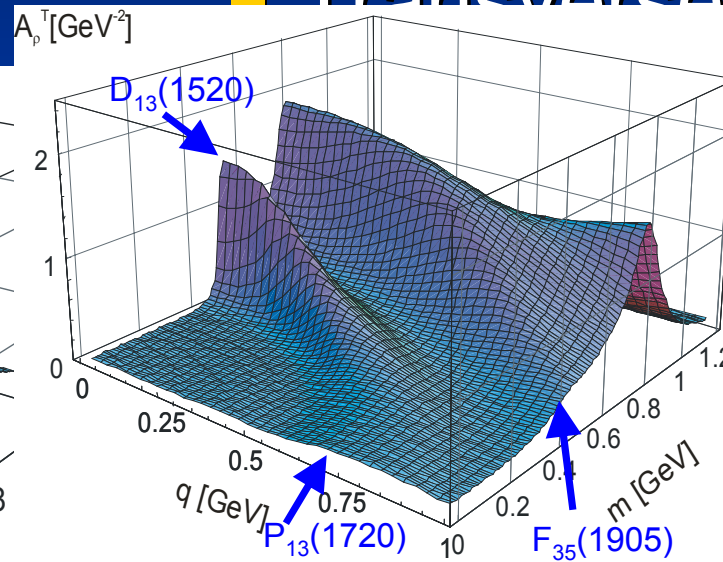
Rho meson in matter

A_L [GeV⁻²] ■ Longitudinal



Post et al.

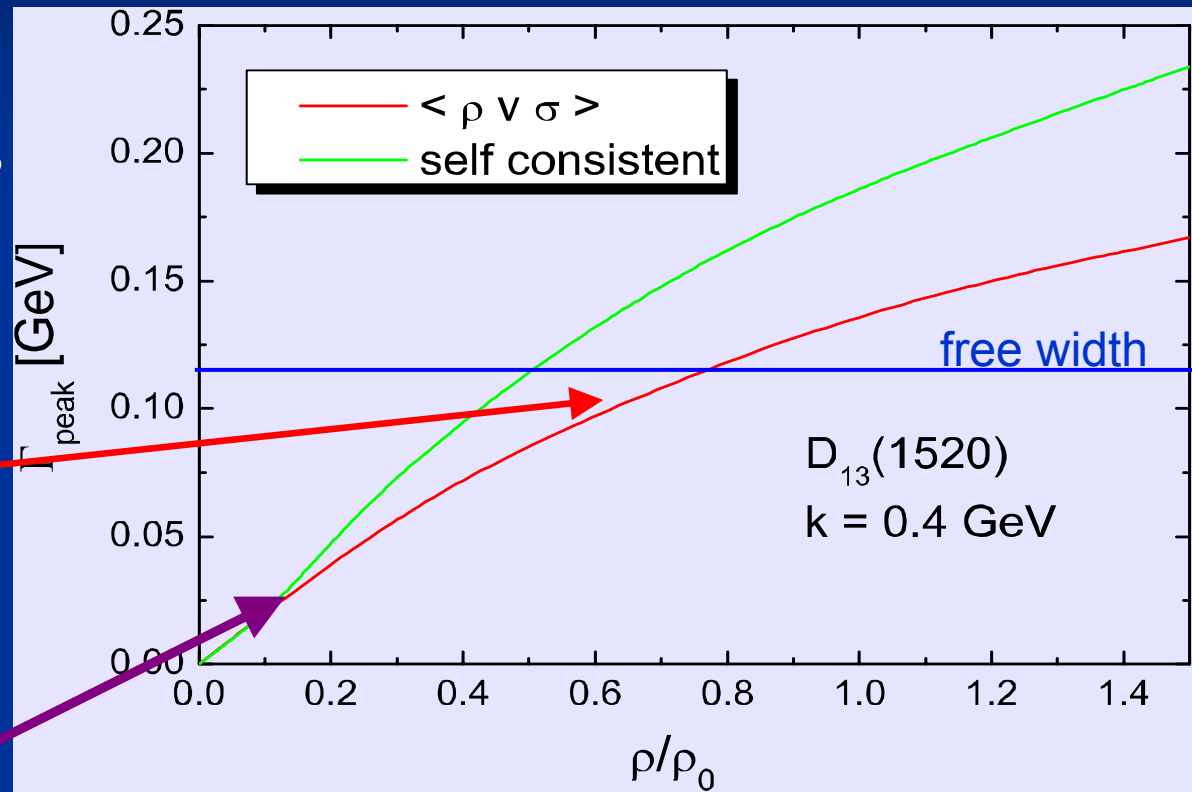
■ Transverse



$D_{13}(1520)$ and 2π (ρ) strongly mixed

In-medium Resonance D13(1520)

Collision width:
N(1520) doubles
width!



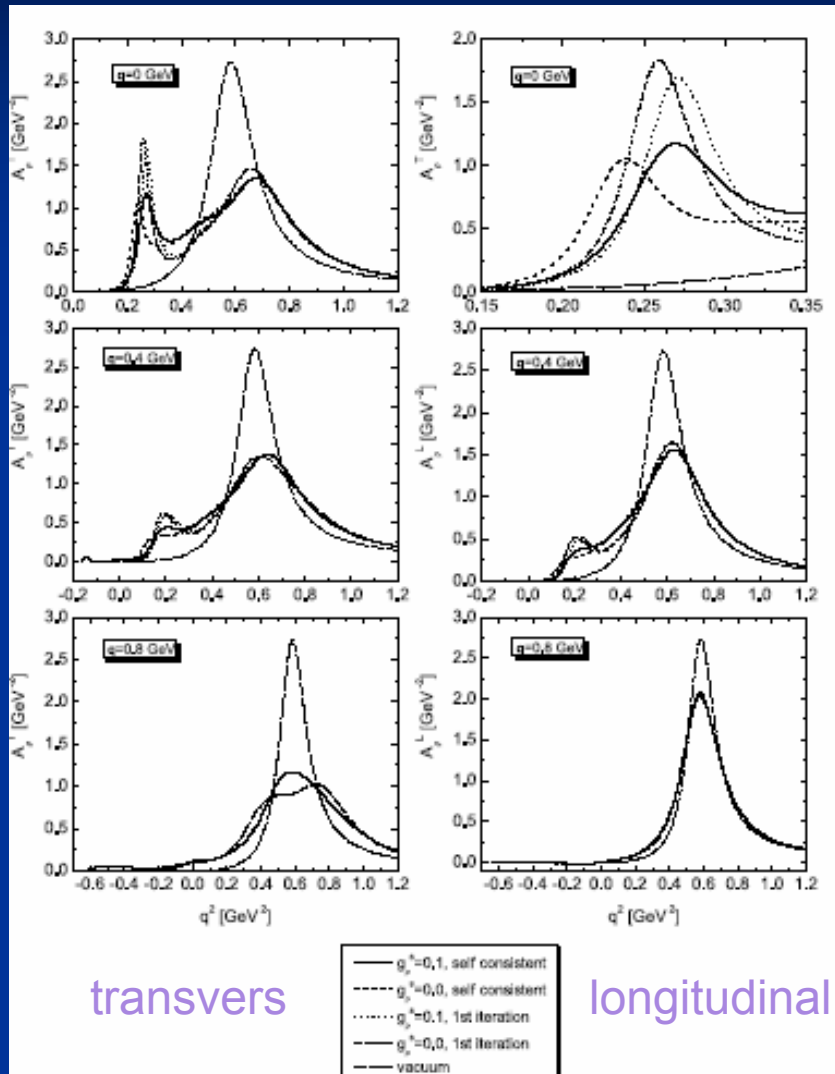
t- ρ approximation

fails already at very low density

Solves photoabsorption puzzle?

Post et al, Nucl.Phys.A741:81,2004

In-medium Resonance N(1520)

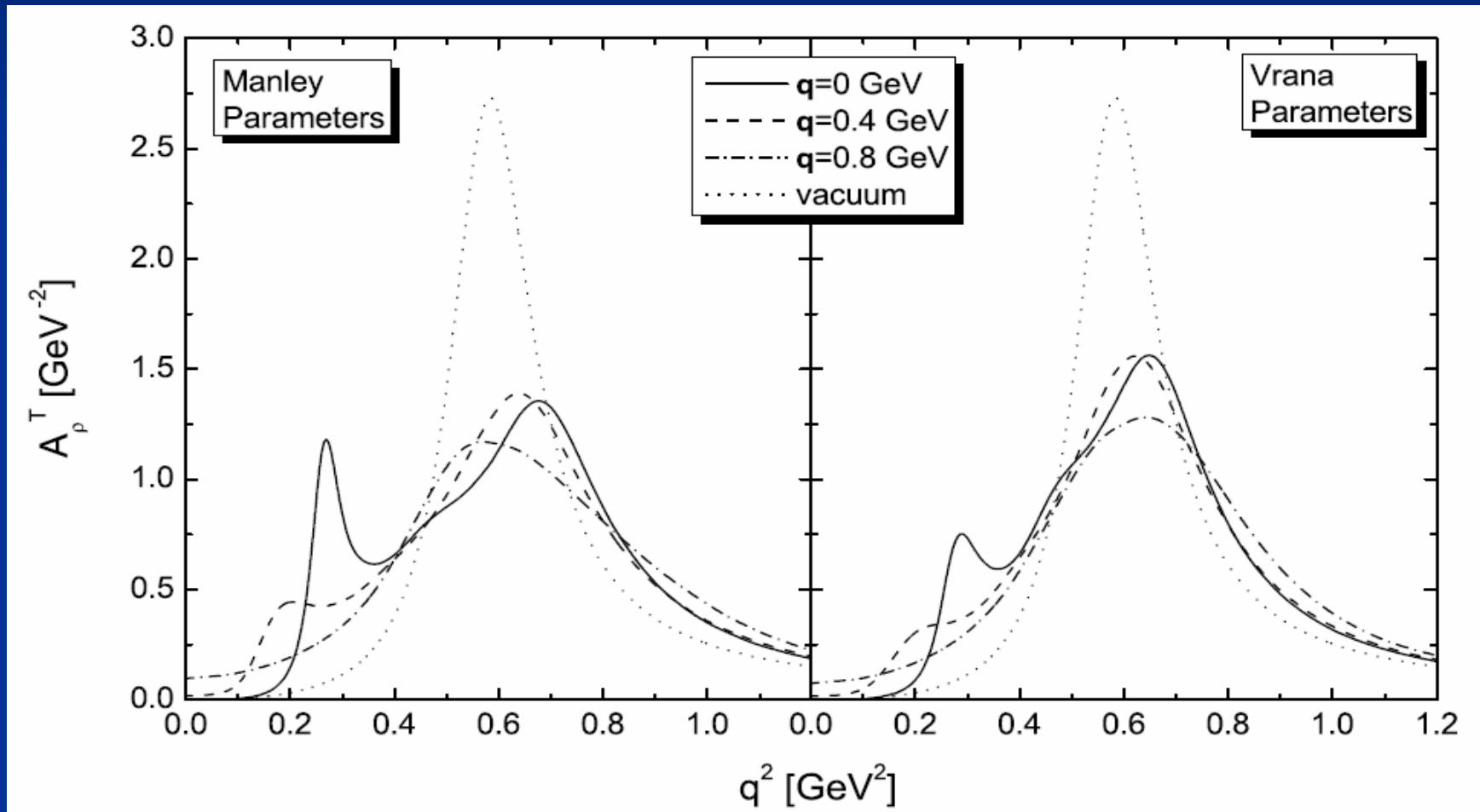


Double bump from s-wave coupling to D13(1520)

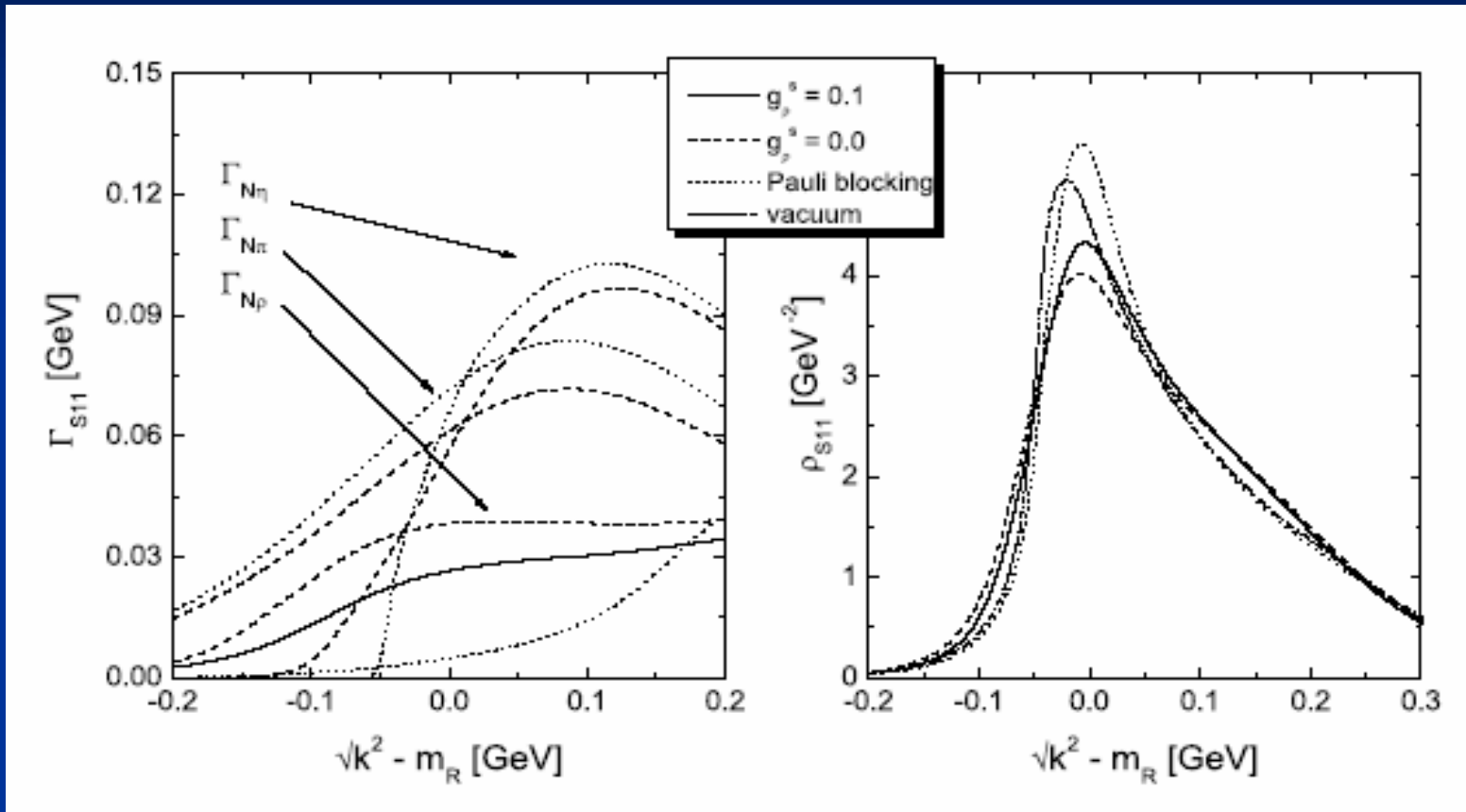
D13 influence reduced, p-wave coupling to P13(1720) and F35(1905)

Transverse broadened by coupling to higher resonances
 Longitudinal does not couple to p-waves, only weakly to s-waves at large q

In-medium Resonance D13(1520) parameter-sensitivity



In-medium Resonance S11(1535)



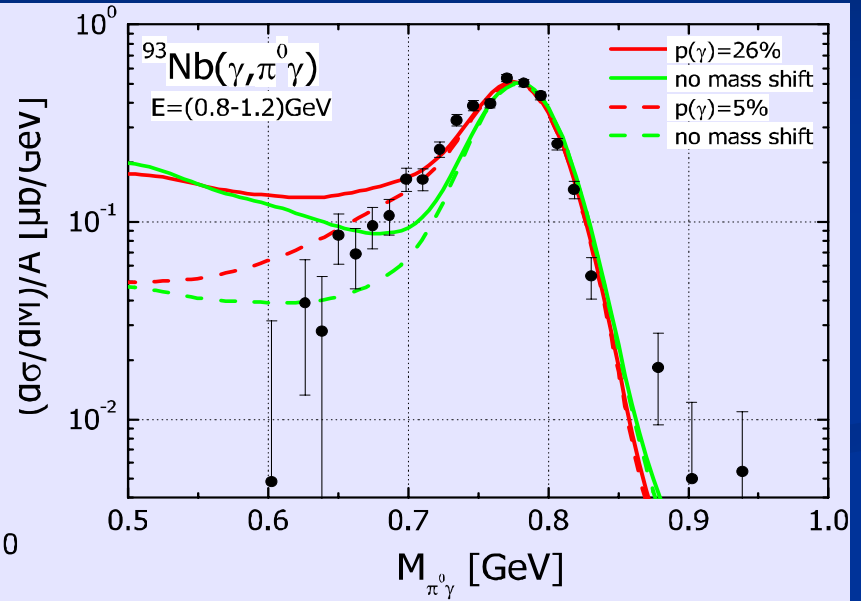
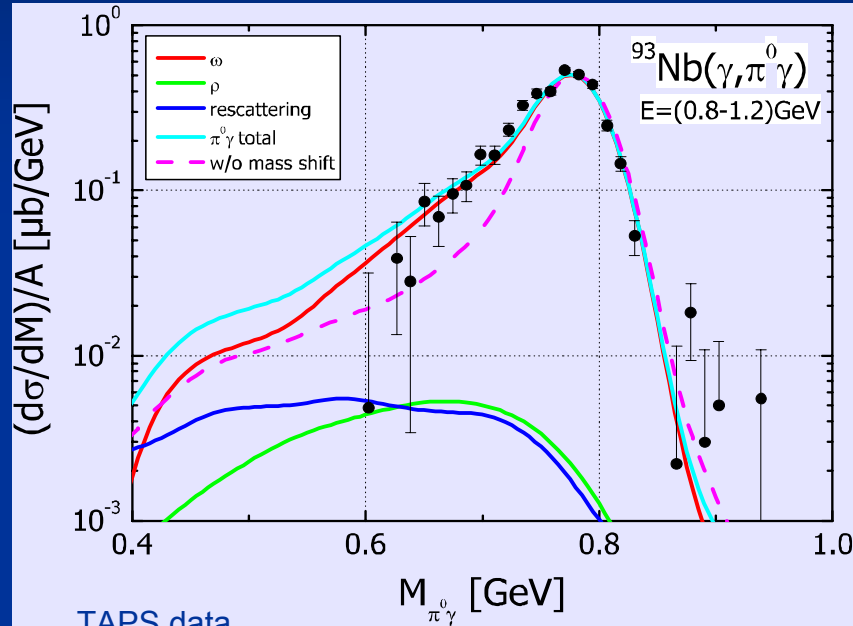
S11(1535) only little modified

Post et al, Nucl.Phys.A741:81,2004

Largest changes from Pauli block and mod. ρ spectral function

Omega in Medium

■ $\omega \rightarrow \gamma \pi^0$ (TAPS@ELSA)



$$M_\omega = M_\omega^0 (1 - 0.16 \rho/\rho_0)$$

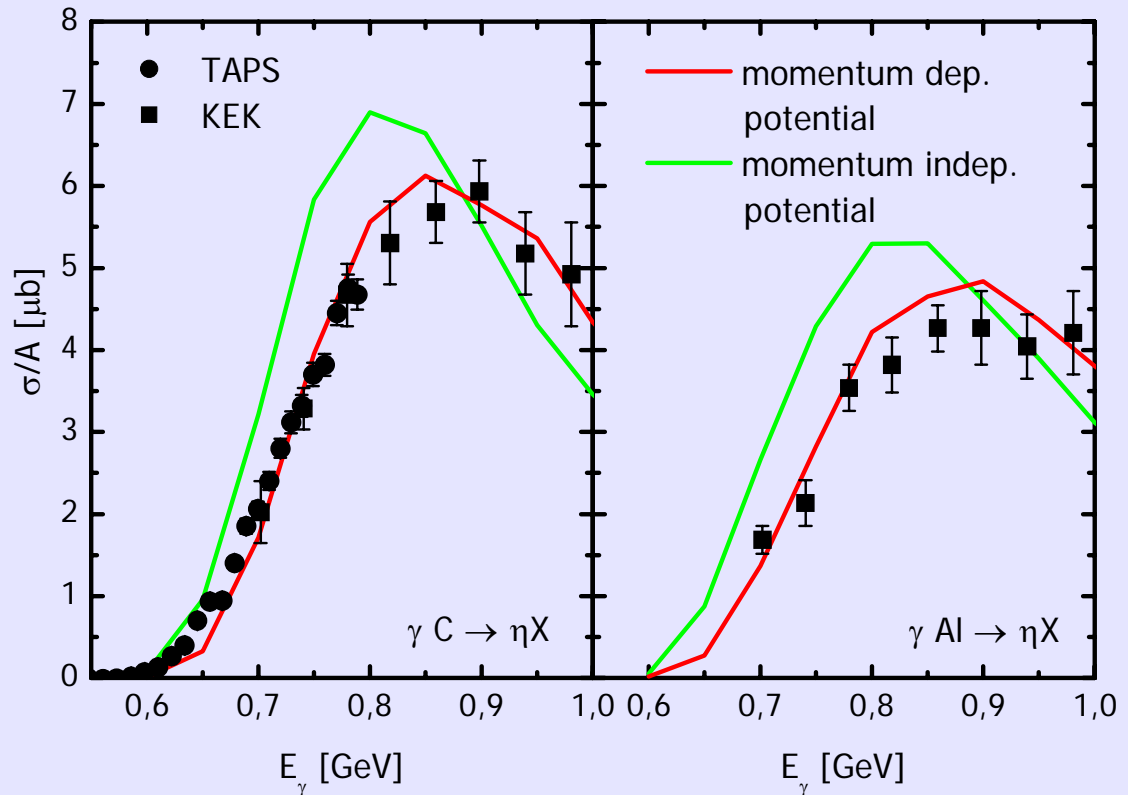
$$\Gamma_\omega \approx 40 \text{ MeV at } \rho_0 \text{ and } p_\omega = 0$$

$$M \text{ independent of } p_\omega$$

Background from
 $2\pi^0 \rightarrow 4\gamma$,
 one escapes with $p(\gamma)$

η Photo-Production on nuclei

- $S_{11}(1535)$: small in-medium change for $p=0$, sizeable momentum dependence of self-energy.



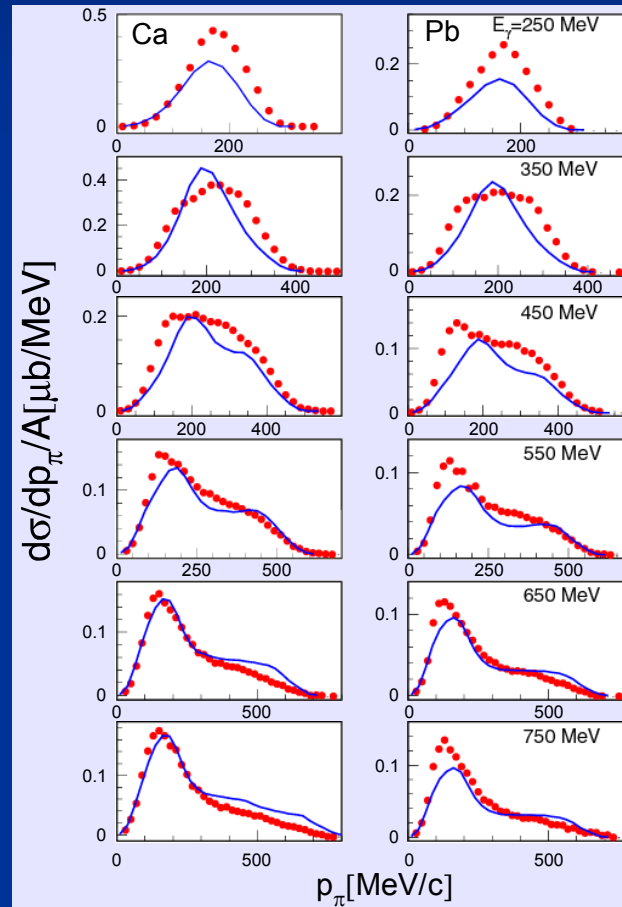
J. Lehr et al, Phys.Rev.C68:044601,2003

Photo-pion production on nuclei

TAPS Data on π^0 production (Krusche et al)

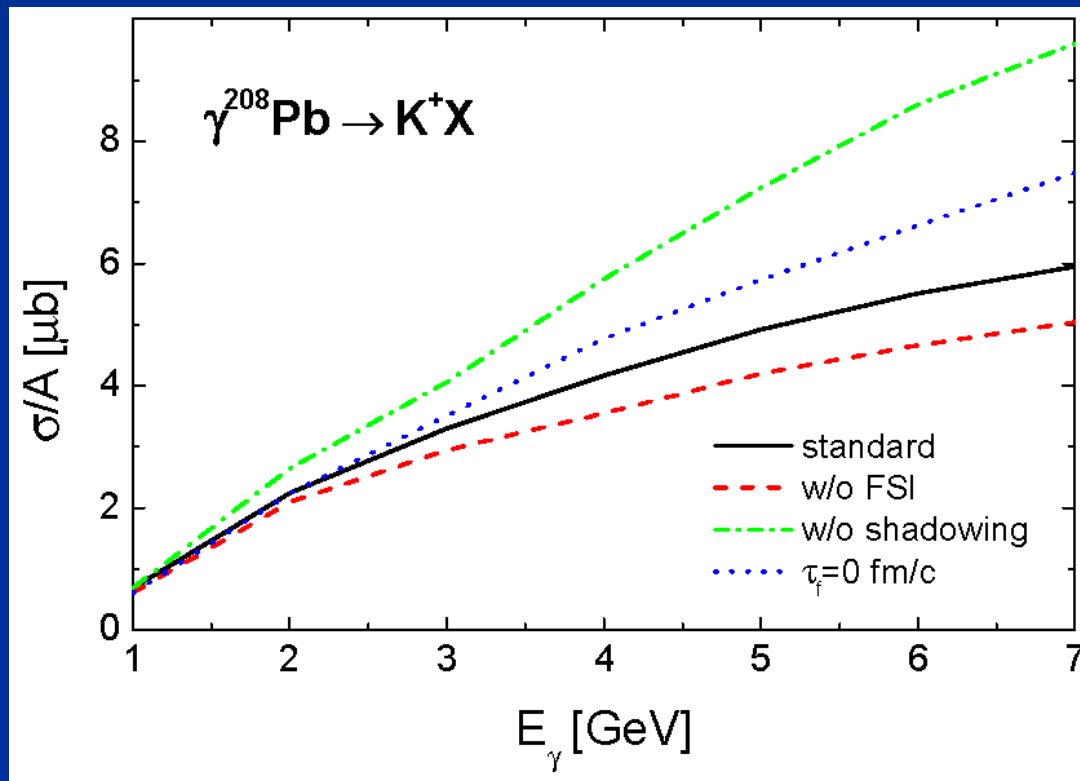
Theory: 
Lehr et al

From:
Krusche et al, nucl-ex/0406002



Coupled Channel Effects

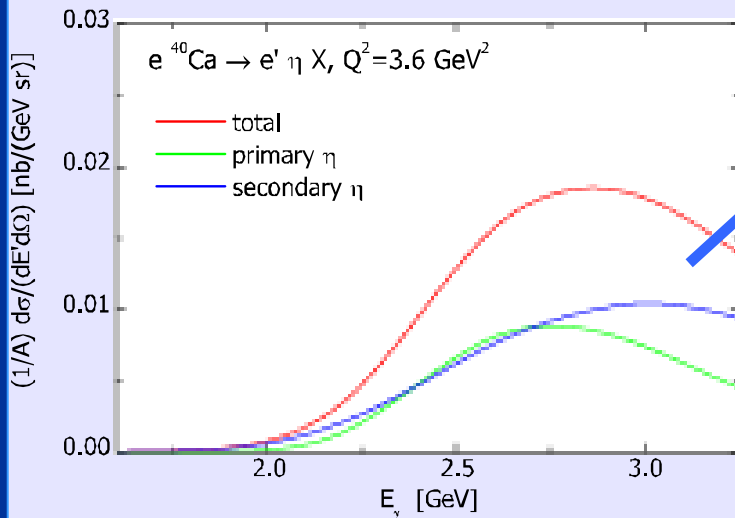
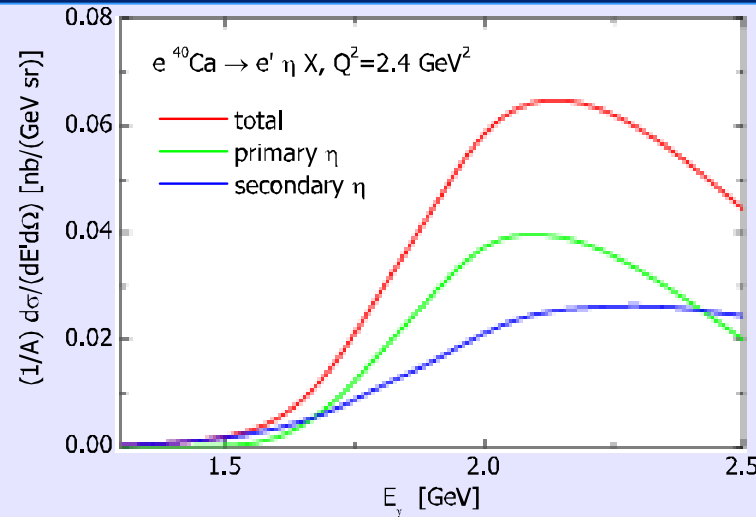
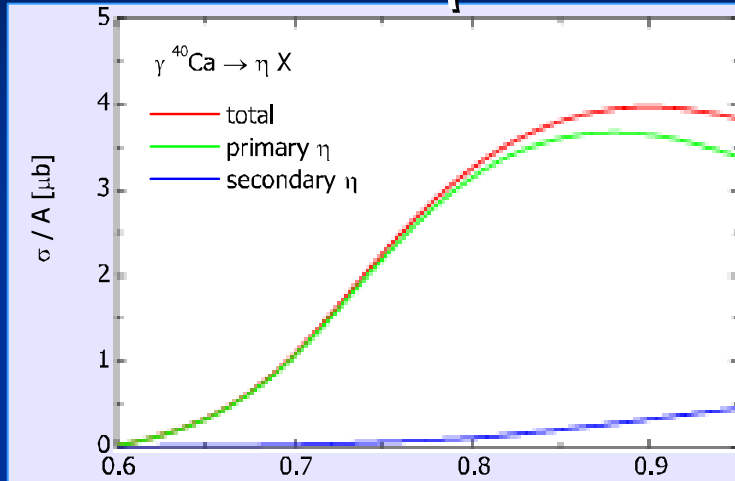
- Final state interactions can **increase** cross section



Falter et al

Coupled Channel Effects

η Electroproduction



Secondaries

Lehr et al,
PRC68:044603,2003

$\gamma N \rightarrow \pi N, \pi N \rightarrow \eta N$

**Dominate
Coupled channel
treatment necessary !**

ϕ photoproduction (Spring 8)

- Predicted:

- $\Delta m \approx -30 \text{ MeV}$

- $\Gamma_{\text{coll}} \approx 24 \text{ MeV}$

- Problem:

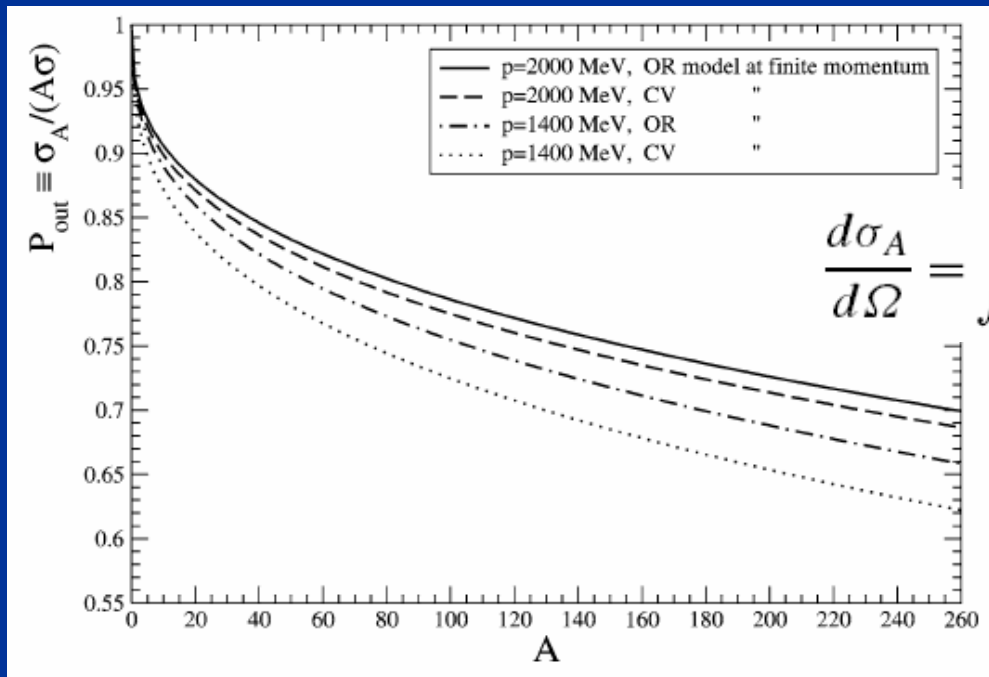
- Strong FSI on kaons ($\phi \rightarrow K^+K^-$)

- Coulomb effects on kaons

- In-medium self-energies of kaons

ϕ photoproduction (Spring 8)

- Integral test: $\Im(\Pi)$ through absorption length (Glauber) \rightarrow A-dependence

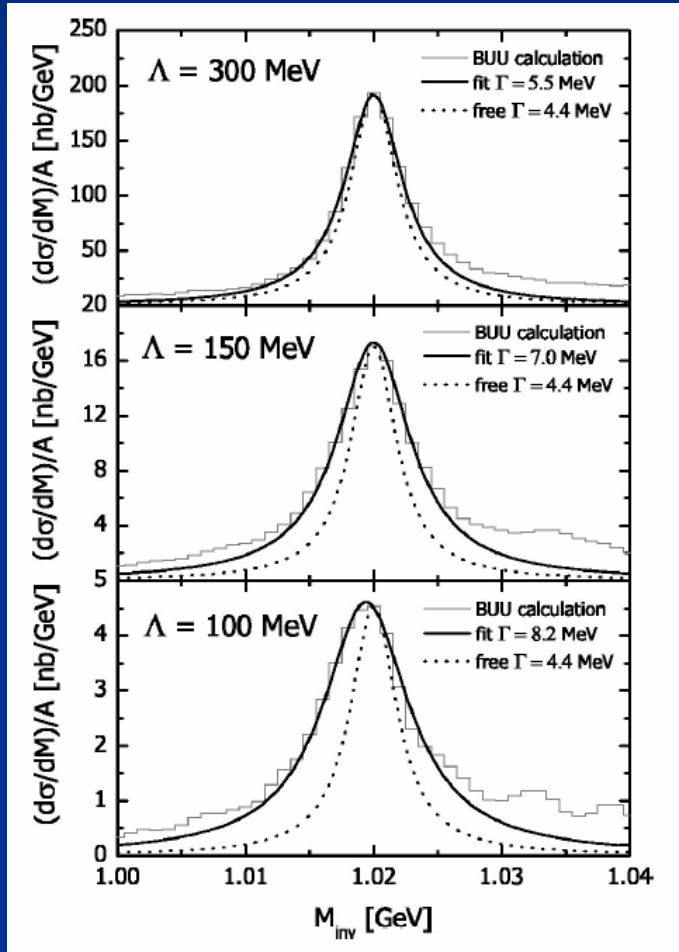


Cabrera et al, Poster

$$\frac{d\sigma_A}{d\Omega} = \int d^3\vec{r} \rho(r) \frac{d\sigma}{d\Omega} e^{-\int_0^\infty dl \frac{-1}{p} \text{Im} \Pi(p, \rho(r'))}$$

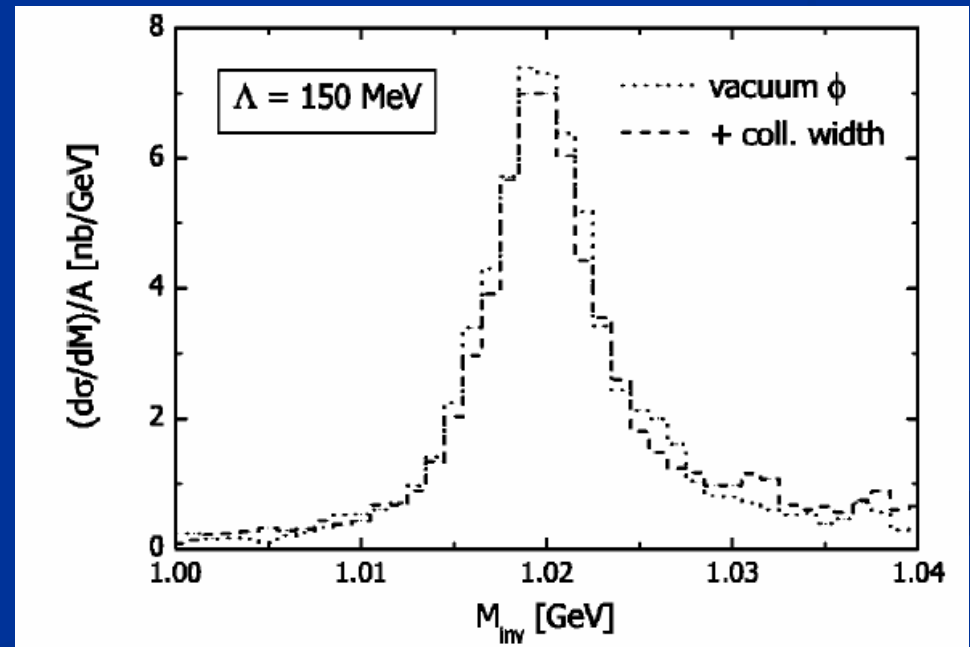
ϕ photoproduction (Spring 8)

Muehlich et al, Phys.Rev.C67:024605,2003



Collisional broadening $\approx \times 2$

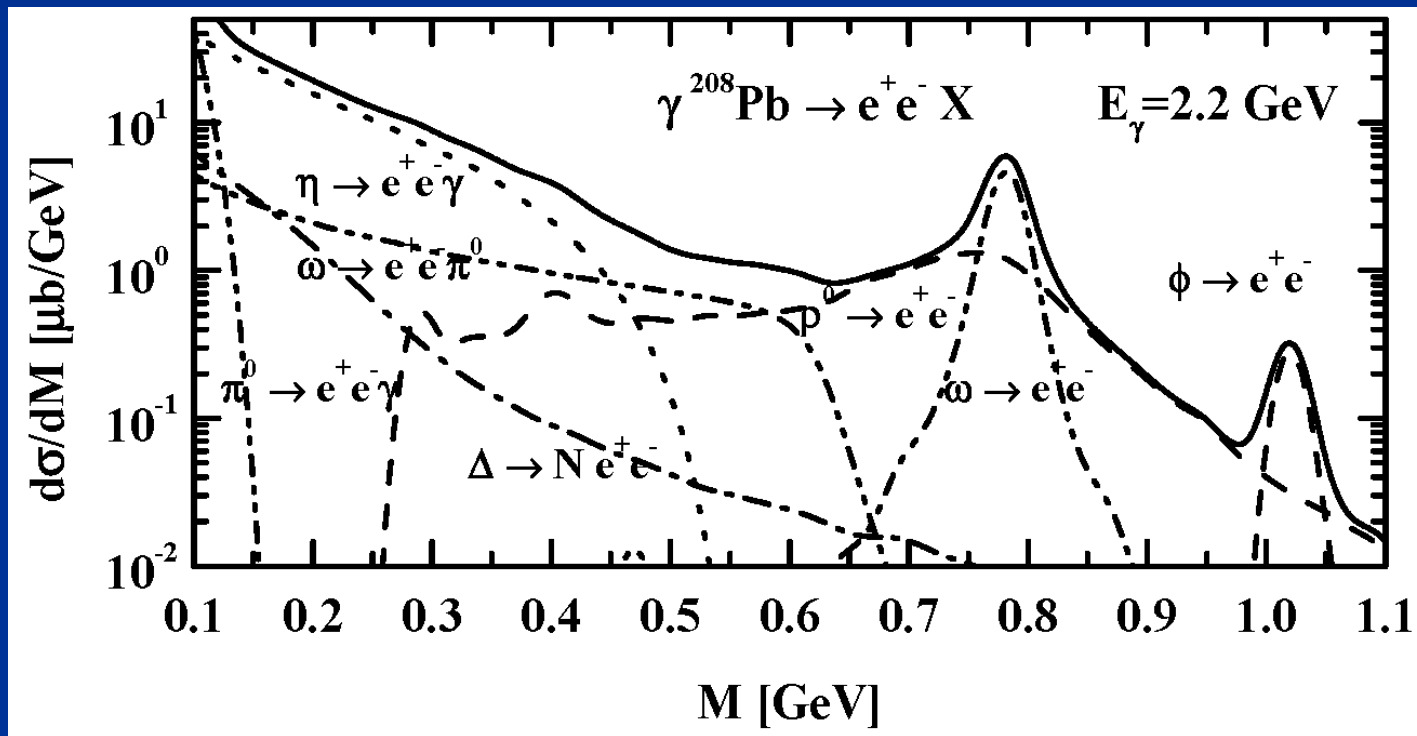
But: Coulomb kills it all



Dileptons: light from the interior

- JLab Experiment E-01-112, g-7

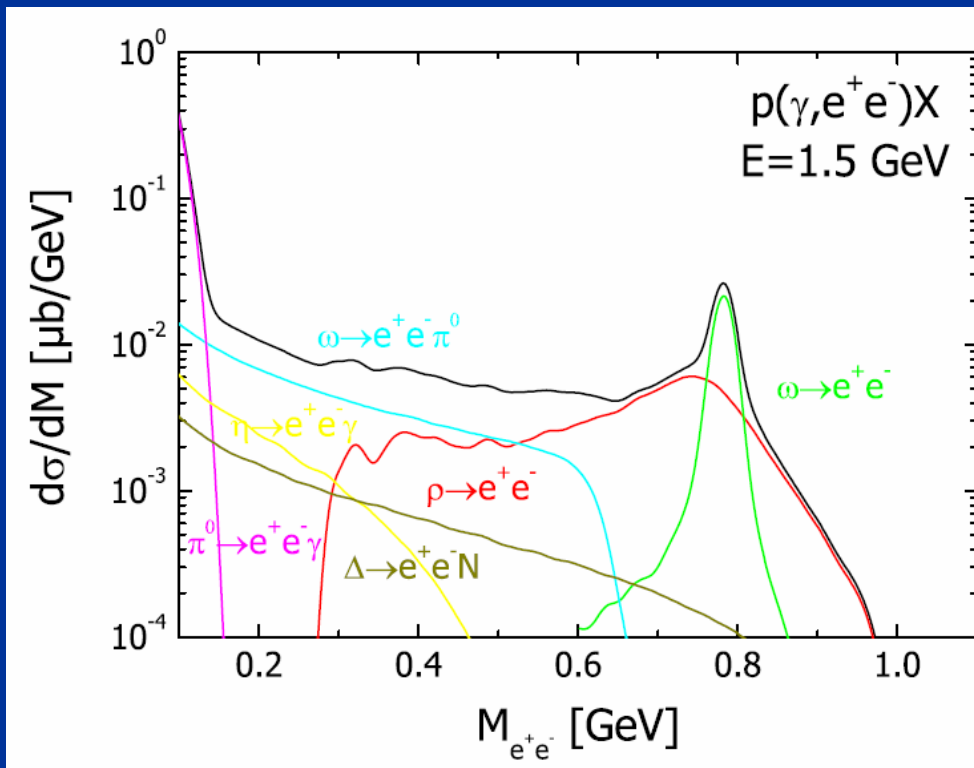
(Djalali, Weygand, Tur et al.)



Same sources as in URHICs!

Dileptons: light from the interior

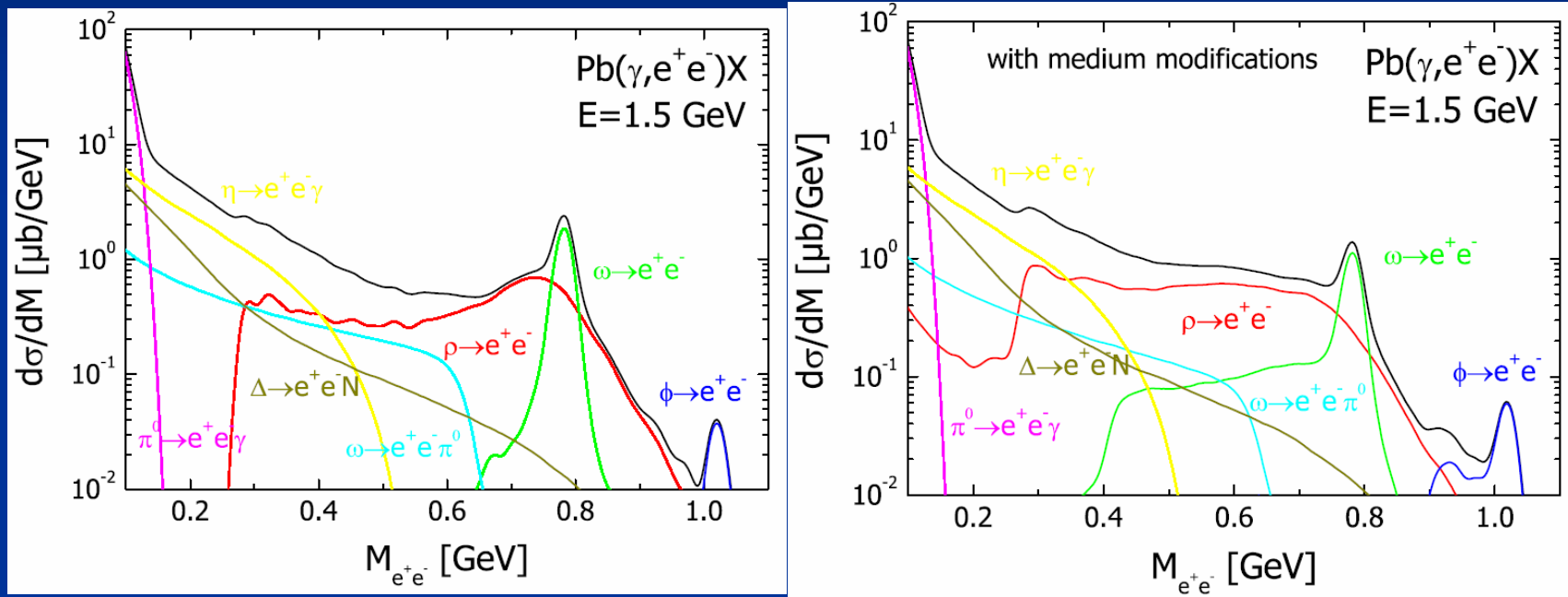
- JLab Experiment E-01-112, g-7
(Djalali, Weygand, Tur et al.)



$$\gamma + p \rightarrow p + e^+e^-$$

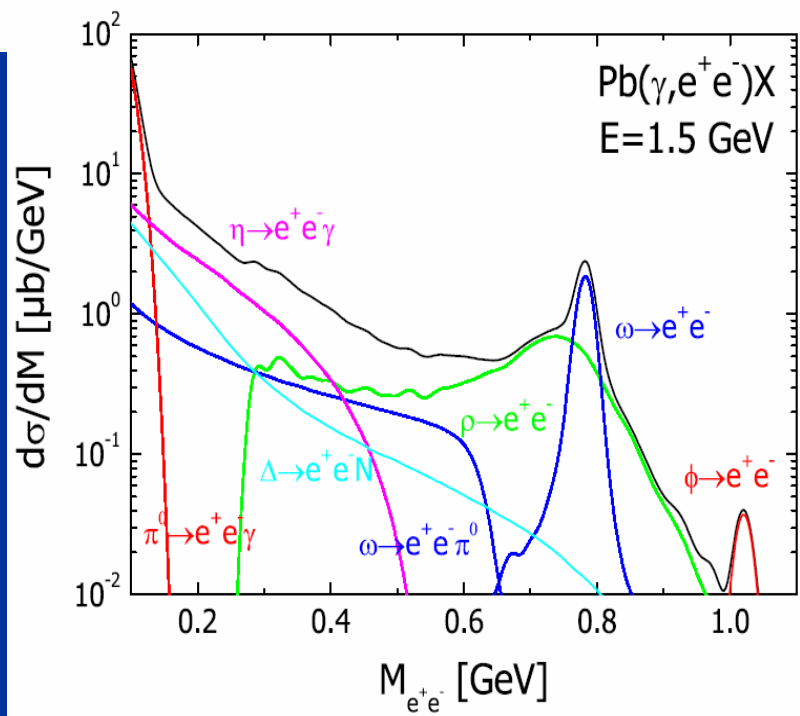
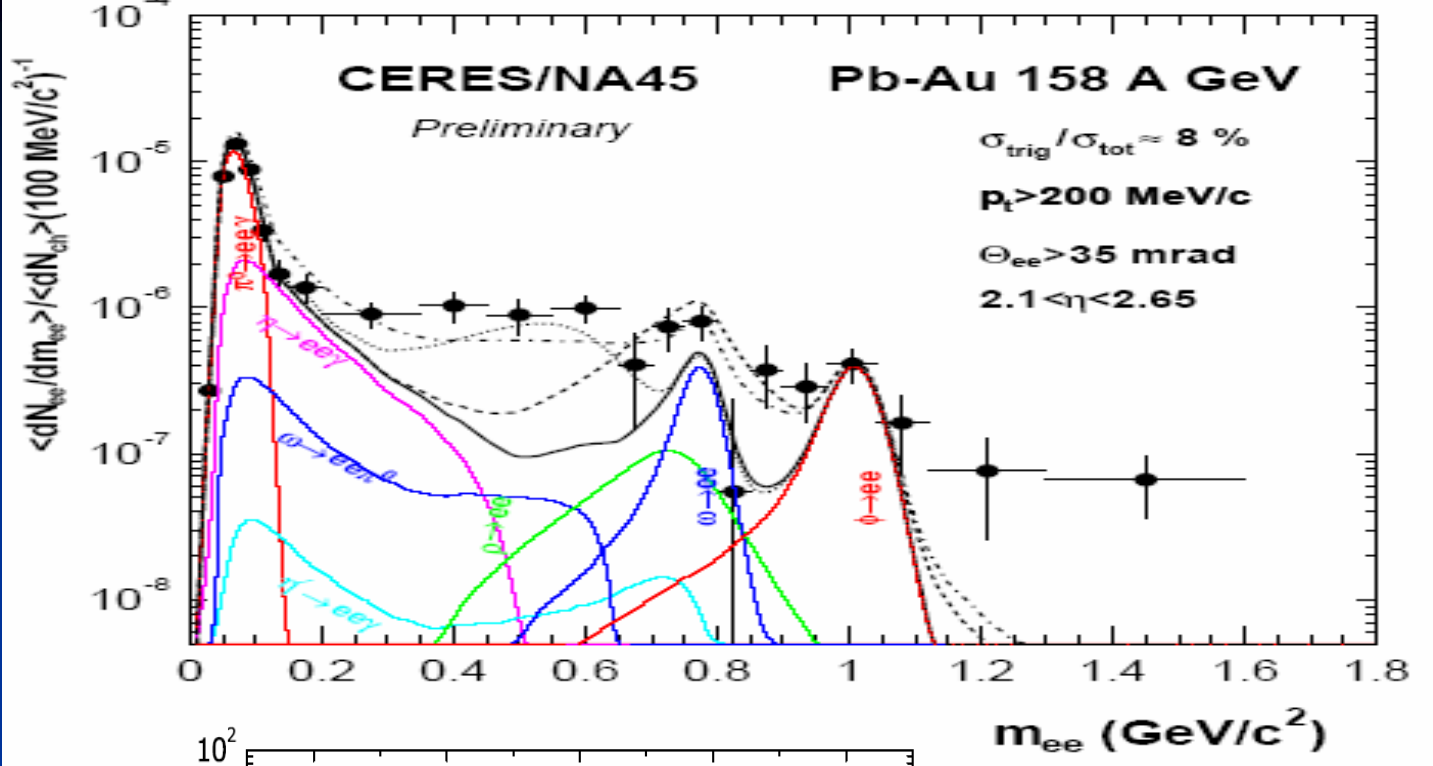
gives access to
em formfactor in
time-like region

JLab Experiment E-01-112, g-7 (Djalali, Weygand, Tur et al.)



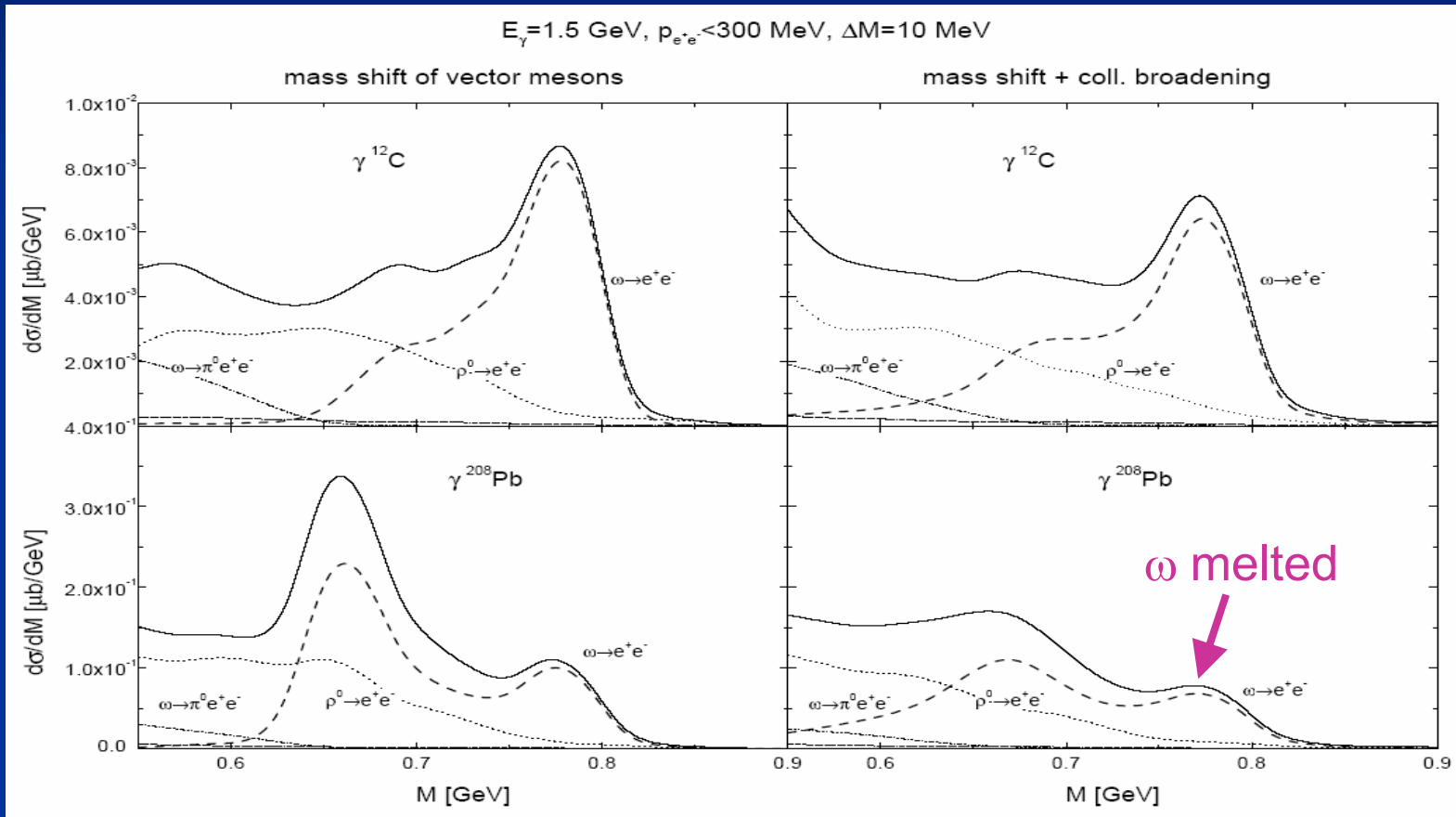
M. Effenberger et al, Phys.Rev.C60:044614,1999

Same sources as in URHICs!



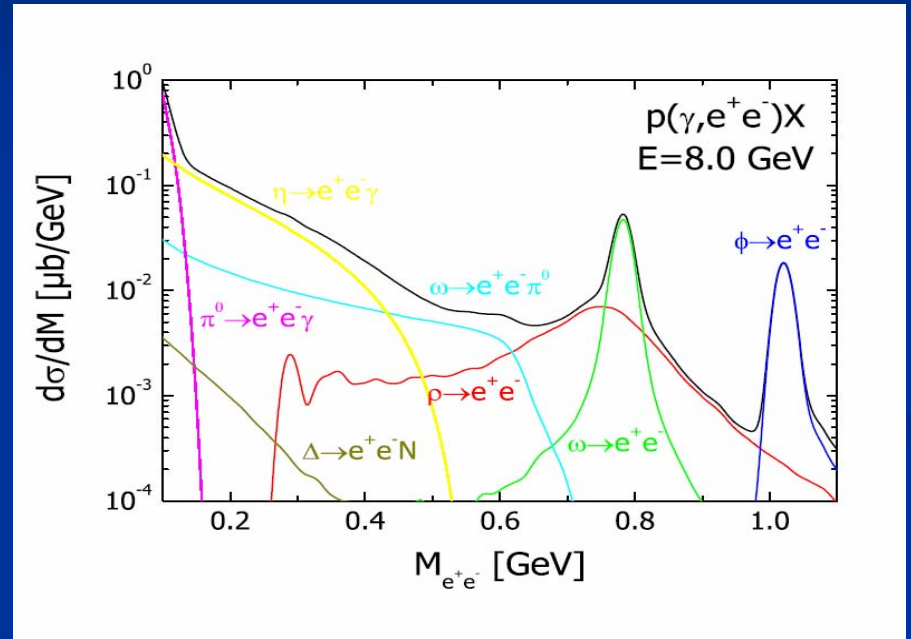
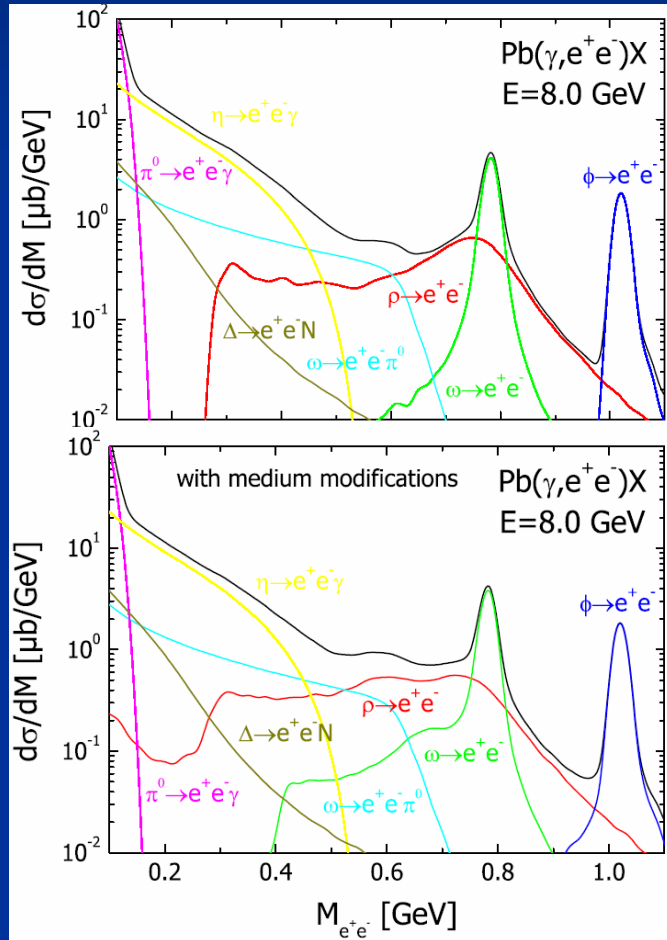
Same Sources !

Dileptons: light from the interior



ω peak fades away with A with cut on low momenta

Dileptons at 8 GeV

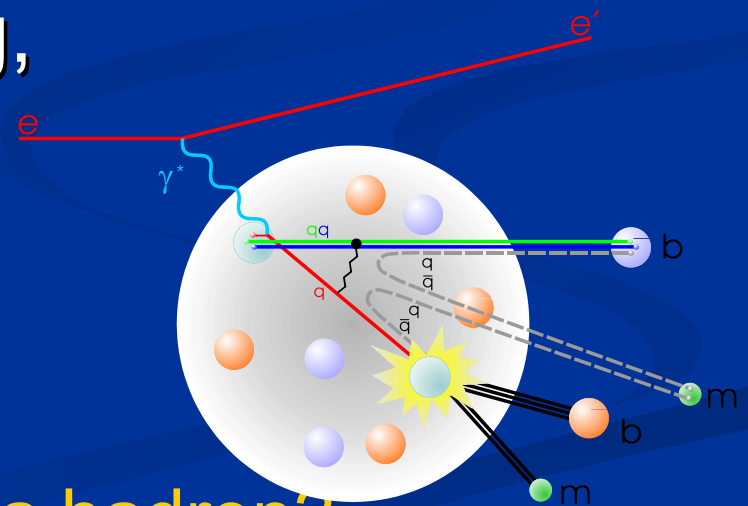
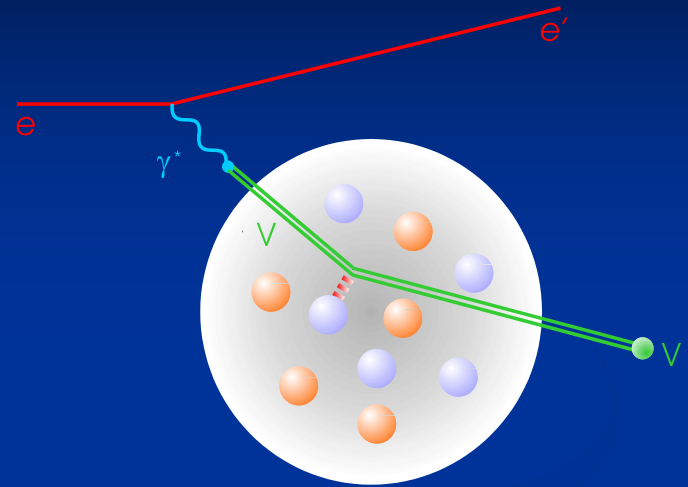


P. Muehllich, 2004

ϕ meson enhanced

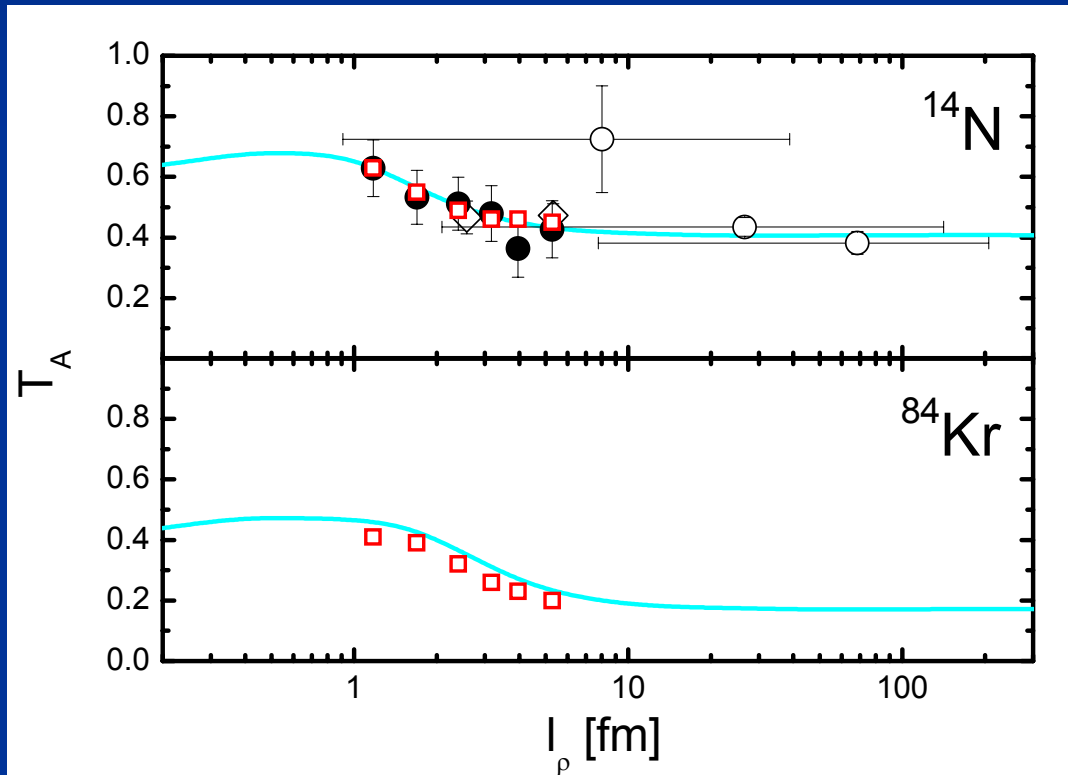
High Energy γ Production Processes

- Diffractive VMD-Event
 - Main contribution to exclusive ρ^0 -production
- Deep inelastic scattering, Jets



How long does it take to form a hadron?

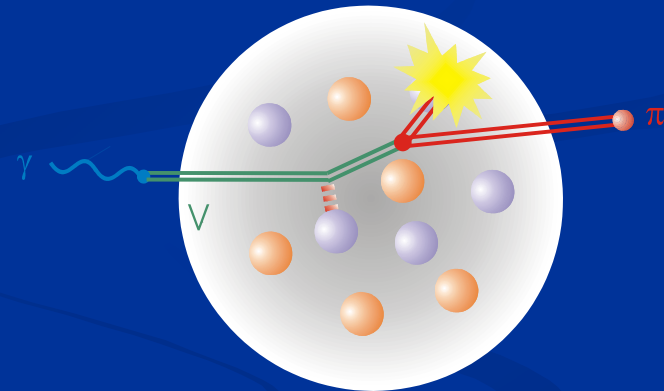
Color Transparency (?)



■ Glauber:

$$\text{FSI} \sim \exp \left[-\sigma_{\rho N}^{\text{inel}} \int_z^\infty dz' n(\vec{b}, z') \right]$$

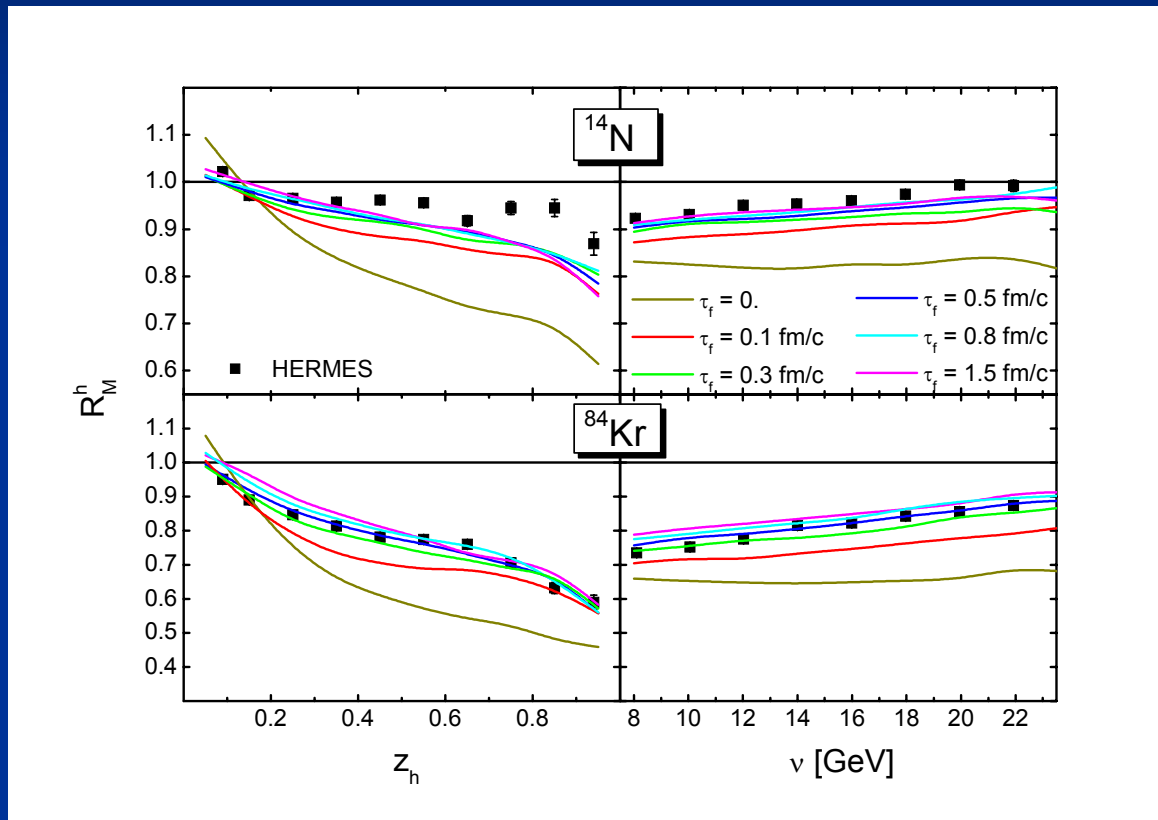
■ BUU:



T. Falter et al, Phys.Rev.C67:054606,2003

In-medium fragmentation (HERMES)

Charged hadron ratios (to p)



Yields info on
hadronization times
→
nucleus as micro-
detector

T. Falter et al,

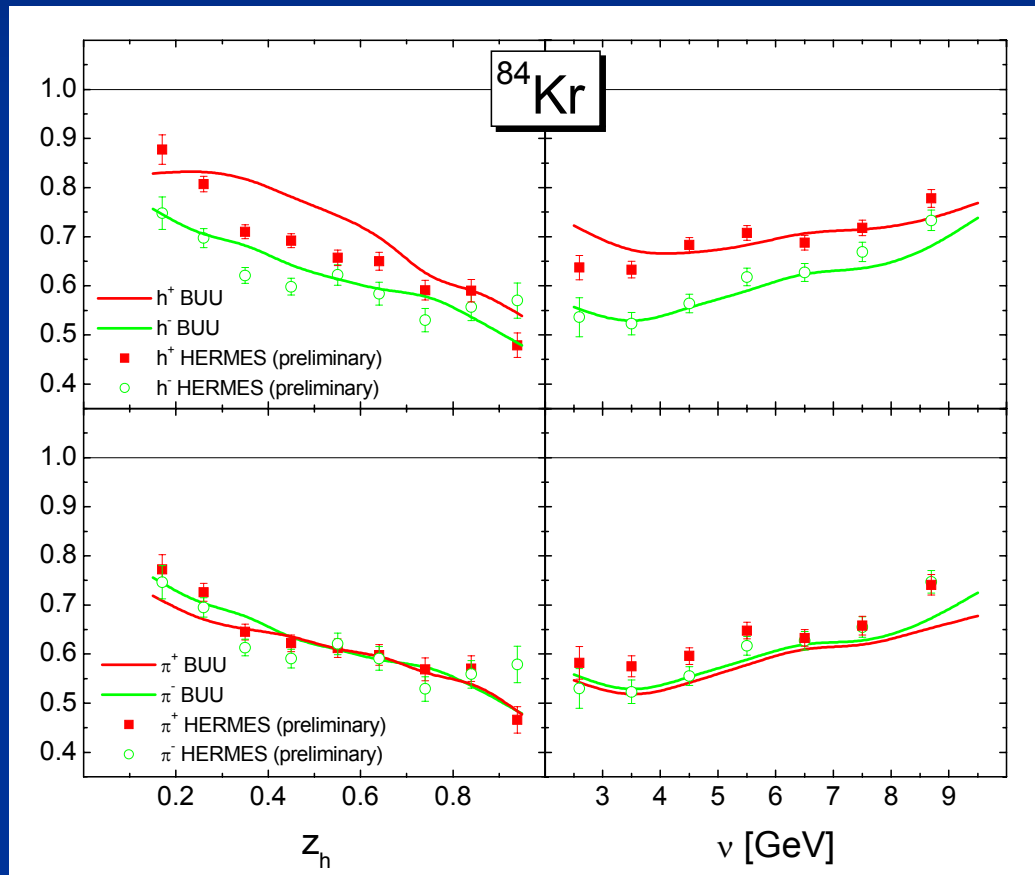
Phys.Lett.B594:61-68,2004

Fast (pre)hadrons see nucleus

Formation time
 $\tau_f > 0.3$ fm/c

■ HERMES @ 12 GeV ($\tau_f = 0.5 \text{ fm}/c$)

Model works also at lower energies



- Jefferson Lab
($\tau_f = 0.5 \text{ fm}/c$)

- CLAS detector
larger geometrical
acceptance

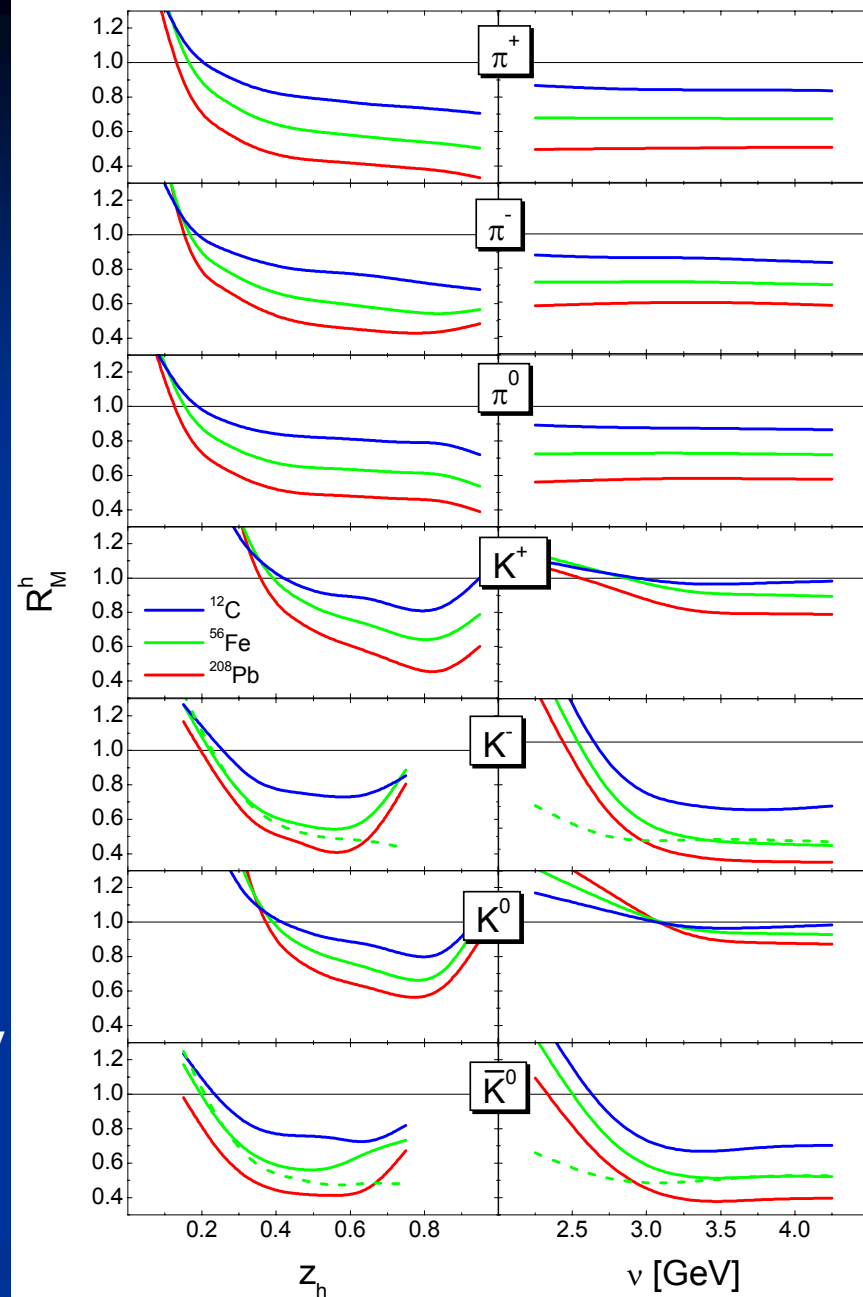
- detects more secondary
particles from FSI

- CEBAF

lower energy $E_e = 5 \text{ GeV}$
 $v_{\text{max}} = 4.25 \text{ GeV}$

- strong effect of
Fermi-motion

T. Falter, PhD thesis, Giessen, 2004

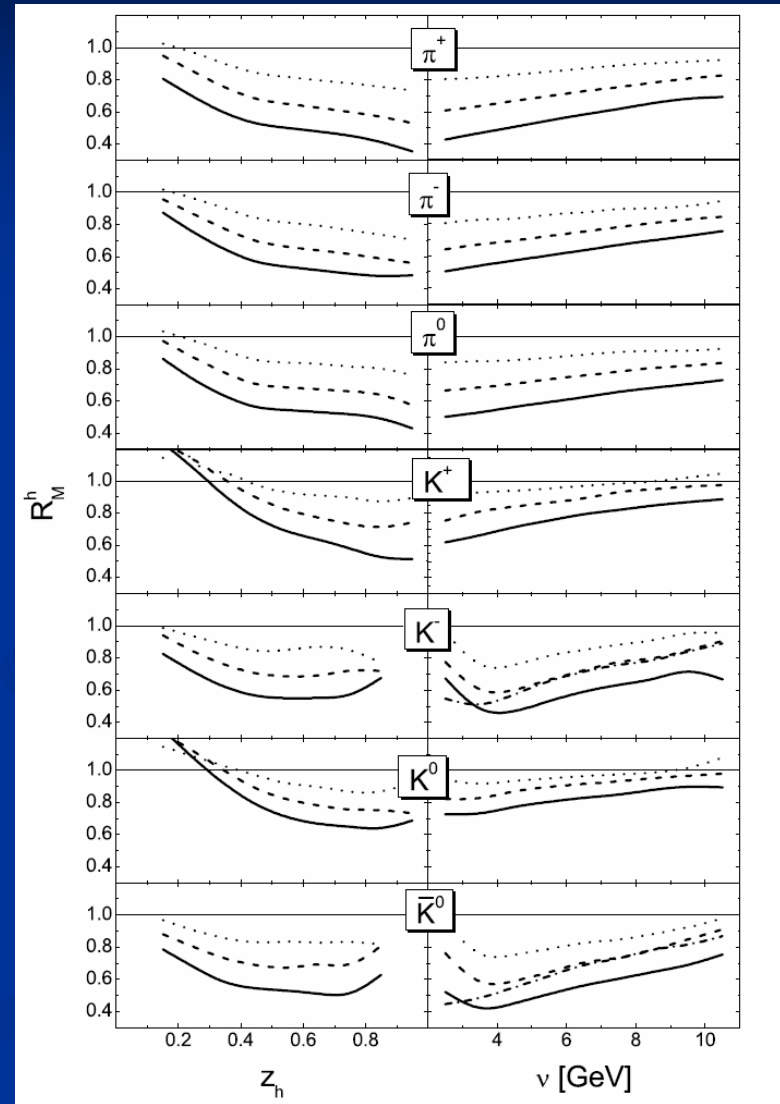
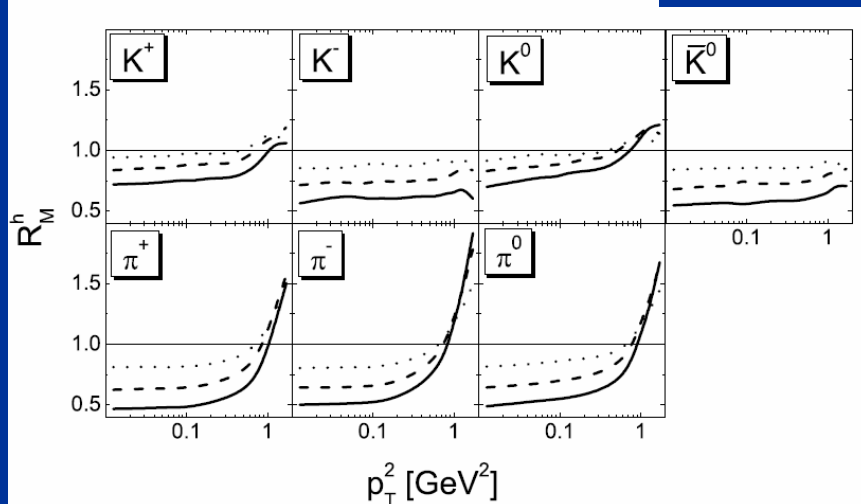
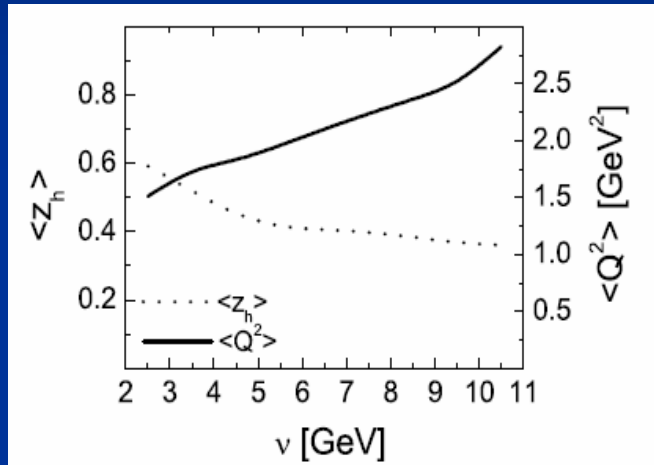


Jlab at 12 GeV

CLAS acceptance modelled

T. Falter, PhD thesis, Giessen, 2004

C
Fe
Pb

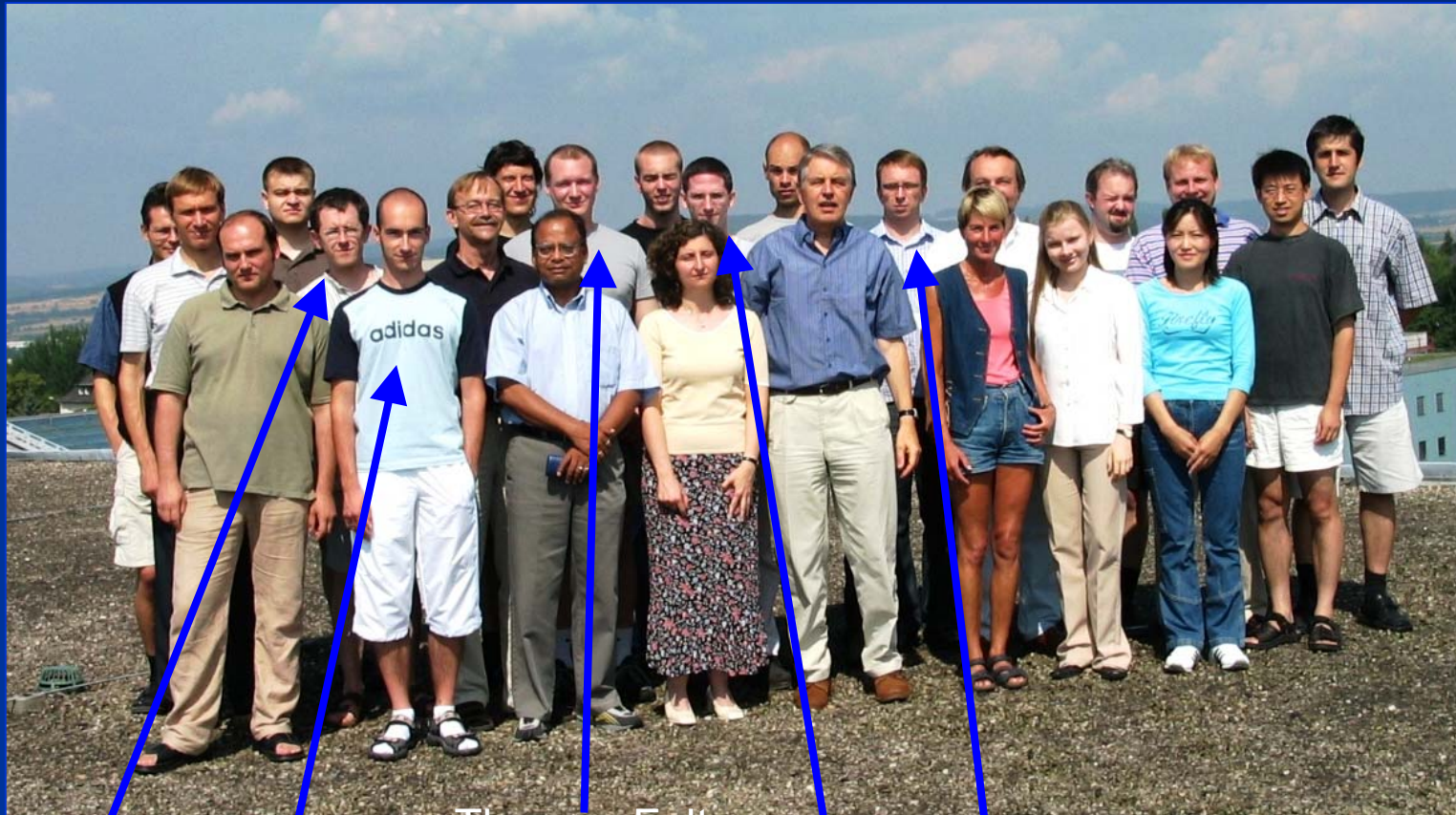


Summary

- Method combines coherence in entrance channel (shadowing) with coupled channel (incoherent) transport in exit channel
- Reliable tool for interpretation of many-body processes: same physics (and code!) for photonuclear and heavy-ion reactions. Allows off-shell transport of broad resonances and implementation of experimental acceptance

Summary

- Theoretical methods to calculate equilibrium in-medium properties quite advanced: selfconsistency possible, but still formidable many channel problem
- Chiral symmetry restoration in nuclei suggestive, but hard to pin down
- Reactions with microscopic probes (p , π , γ) provide baseline for ,exotic' phenomena in URHICs
- At high energies access to formation times and CT



Stefan Leupold

Pascal Mühlich

Thomas Falter

Jürgen Lehr

Marcus Post

