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# Generalised Parton Distributions and Transversity from Lattice QCD

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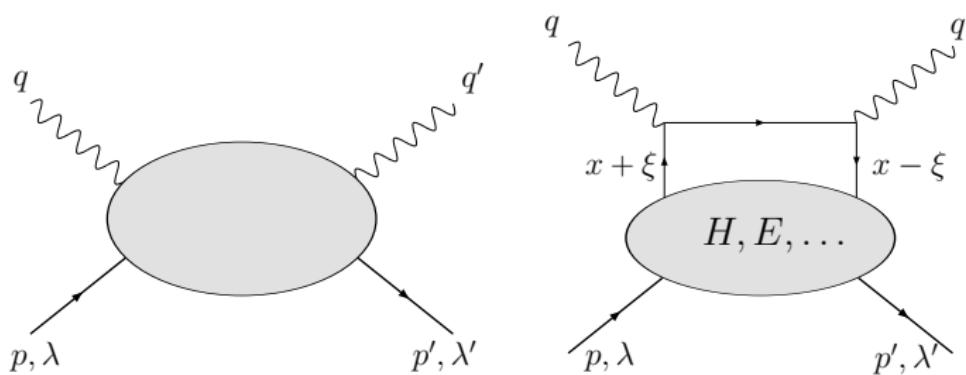


## Motivation

- Unique opportunity for lattice calculations to compliment experiment
- Assist in the full mapping of the parameter space spanned by GPDs

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$$\Delta = p' - p, \quad t = \Delta^2, \quad \bar{p} = \frac{p' + p}{2}, \quad \xi = -n \cdot \Delta, \quad n \cdot \bar{p} = 1$$

# Generalised Parton Distributions

- GPDs are defined through nucleon matrix elements of the light cone operators

$$\int \frac{d\lambda}{4\pi} e^{i\lambda x} \langle P' | \bar{\psi}(-\frac{\lambda}{2}n) \gamma^\mu \psi(\frac{\lambda}{2}n) | P \rangle = \\ \overline{U}(P') (\gamma^\mu H(x, \xi, t) + \frac{i\sigma^{\mu\nu}\Delta_\nu}{2m} E(x, \xi, t)) U(P)$$

$$\int \frac{d\lambda}{4\pi} e^{i\lambda x} \langle P' | \bar{\psi}(-\frac{\lambda}{2}n) \gamma_5 \gamma^\mu \psi(\frac{\lambda}{2}n) | P \rangle = \\ \overline{U}(P') (\gamma_5 \gamma^\mu \tilde{H}(x, \xi, t) + \frac{i\gamma_5 \Delta^\mu}{2m} \tilde{E}(x, \xi, t)) U(P)$$

$$\int \frac{d\lambda}{4\pi} e^{i\lambda x} \langle P' | \bar{\psi}(-\frac{\lambda}{2}n) \sigma^{\mu\nu} \psi(\frac{\lambda}{2}n) | P \rangle = \\ \overline{U}(P') \left( i\sigma^{\mu\nu} H_T(x, \xi, t) + \frac{\gamma^{[\mu} \Delta^{\nu]}}{2m} E_T(x, \xi, t) \right. \\ \left. + \frac{\overline{P}^{[\mu} \Delta^{\nu]}}{m^2} \tilde{H}_T(x, \xi, t) + \frac{\gamma^{[\mu} \overline{P}^{\nu]}}{m} \tilde{E}_T(x, \xi, t) \right) U(P)$$

# Generalised Parton Distributions

- The forward limit  $t = \Delta^2 \rightarrow 0$  reproduces the parton distributions,  $H(x, 0, 0) = q(x)$ ,  $\tilde{H}(x, 0, 0) = \Delta q(x)$  and  $H_T(x, 0, 0) = \delta q(x)$
- $\int dx \rightarrow$  Dirac, Pauli, axial, pseudo-scalar, tensor etc. form factors,  $\int dx H(x, \xi, t) = F_1(t)$ ,  $\int dx E = F_2(t)$ ,  
 $\int dx \tilde{H} = g_A(t)$ , ...
- FT  $\int d\Delta_\perp e^{-ib_\perp \cdot \Delta_\perp}$  GPDs  $H$ ,  $\tilde{H}$  and  $H_T$  ( $\xi = 0$ ) are coordinate space probability densities in the impact parameter  $b_\perp$
- $\int dx x E(x, 0, 0) = B_{20}(0)$ ,  $\rightarrow$  quark orbital angular momentum contribution to the nucleon spin,  $L^q = 1/2(\langle x \rangle + B_{20} - \Delta q)$ , where  $\langle x \rangle$  is the quark momentum fraction



# Generalised Form Factors

- Construct Mellin moments
- non-forward MEs of tower of local twist-2 operators

$$\mathcal{O}_q^{\{\mu_1 \dots \mu_n\}} = \bar{q} \gamma^{\{\mu_1} \overleftrightarrow{D}^{\mu_2} \dots \overleftrightarrow{D}^{\mu_n\}} q$$

→ Generalised Form Factors

$$\begin{aligned} \langle p', s' | \mathcal{O}^{\{\mu_1 \dots \mu_n\}}(\Delta) | p, s \rangle &= \\ &\bar{u}(p', s') \gamma^{\{\mu_1} u(p, s) \sum_{i=0}^{\frac{n-1}{2}} A_{qn,2i}(t) \Delta^{\mu_2} \dots \Delta^{\mu_{2i+1}} \bar{p}^{\mu_{2i+2}} \dots \bar{p}^{\mu_n\}} \\ &+ \bar{u}(p', s') \frac{i\sigma^{\{\mu_1 \nu} \Delta_\nu}{2m} u(p, s) \sum_{i=0}^{\frac{n-1}{2}} B_{qn,2i}(t) \Delta^{\mu_2} \dots \Delta^{\mu_{2i+1}} \bar{p}^{\mu_{2i+2}} \dots \bar{p}^{\mu_n\}} \\ &+ C_{qn}(t) \frac{1}{m} \bar{u}(p', s') u(p, s) \Delta^{\mu_1} \dots \Delta^{\mu_n} |_{n \text{ even}} \end{aligned}$$

- Similar for

$$\tilde{A}_{qn,2i}(t), \tilde{B}_{qn,2i}(t), A_{qn,2i}^T(t), B_{qn,2i}^T(t), \tilde{A}_{qn,2i}^T(t), \tilde{B}_{qn,2i}^T(t)$$



# Moments of GPDs

- GPDs are defined through their moments.

$$\int_{-1}^1 dx x^{n-1} H_q(x, \xi, t) = H_{qn}(\xi, t)$$

$$\int_{-1}^1 dx x^{n-1} E_q(x, \xi, t) = E_{qn}(\xi, t)$$

where  $\xi = -n \cdot \Delta$ ,  $n \cdot \bar{p} = 1$ ,  $\bar{p} = \frac{1}{2}(p' + p)$  and

$$H_{qn}(\xi, t) = \sum_{i=0}^{\frac{n-1}{2}} A_{qn,2i}(t)(-2\xi)^{2i} + C_{qn}(t)(-2\xi)^n|_{n \text{ even}}$$

$$E_{qn}(\xi, t) = \sum_{i=0}^{\frac{n-1}{2}} B_{qn,2i}(t)(-2\xi)^{2i} - C_{qn}(t)(-2\xi)^n|_{n \text{ even}}$$



## Moments of GPDs II

- zeroth moment,  $n = 1$

$$\int_{-1}^1 dx \, H_q(x, \xi, t) = F_1(t)$$

$$\int_{-1}^1 dx \, E_q(x, \xi, t) = F_2(t)$$

- first moment,  $n = 2$

$$\int_{-1}^1 dx x \, H_q(x, \xi, t) = A_{2,0}^q(t) + \xi^2 C_2^q(t)$$

$$\int_{-1}^1 dx x \, E_q(x, \xi, t) = B_{2,0}^q(t) - \xi^2 C_2^q(t)$$

# Impact Parameter GPDs

- Probabilistic interpretation of  $H(x, \xi, t)$  (and  $\tilde{H}(x, \xi, t)$ ,  $H_T(x, \xi, t)$ ) at  $\xi = 0$  due to M.Burkardt

$$q(x, \vec{b}_\perp) = \frac{1}{(2\pi)^2} \int d^2 \Delta_\perp e^{-i \vec{b}_\perp \cdot \Delta_\perp} H(x, 0, -\Delta_\perp^2)$$

- Also shows  $H(x, 0, t = -\Delta_\perp^2) \rightarrow t\text{-independent for } x \rightarrow 1$

$$\Rightarrow \lim_{x \rightarrow 1} q(x, \vec{b}_\perp) \propto \delta^2(\vec{b}_\perp)$$

## Ji's Angular Momentum Sum Rule

- Forward limit,  $t = \xi = 0$

$$A_{2,0}^q(0) = \langle x_q \rangle \equiv \int_{-1}^1 dx x (q_\uparrow(x) + q_\downarrow(x))$$

$$\frac{1}{2}(A_{2,0}^q(0) + B_{2,0}^q(0)) = J_q$$

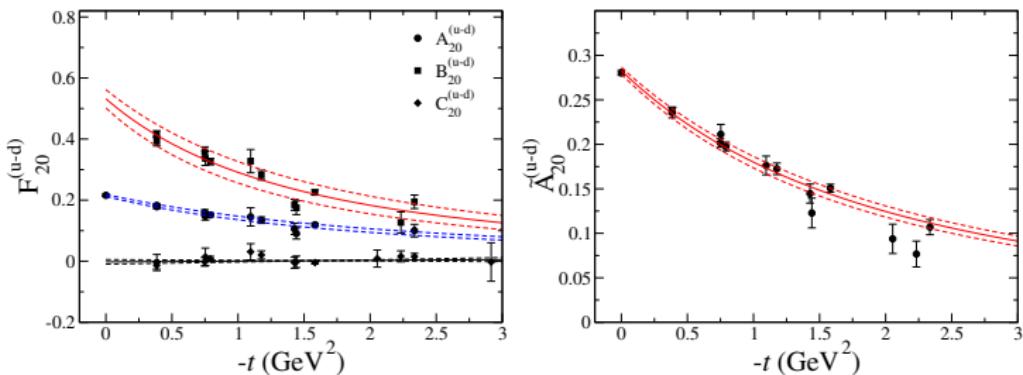
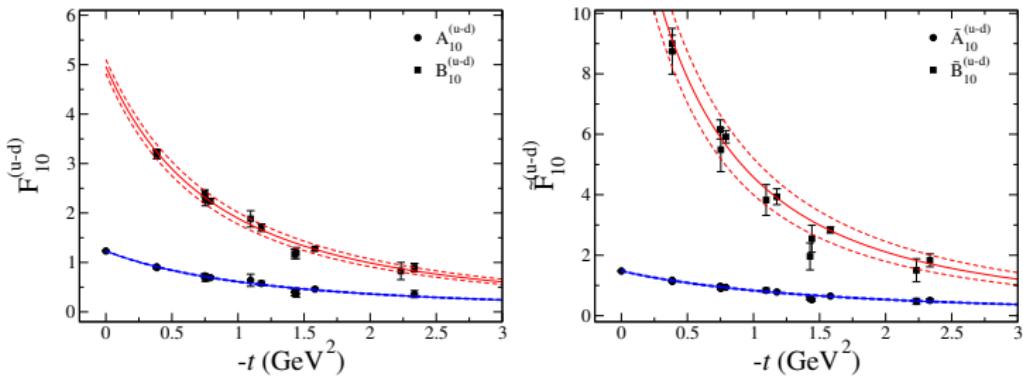
- $J_q = L_q + S_q$  - angular momentum of quark, q
  - $L_q$  - orbital angular momentum of q
  - $S_q$  - spin of q

$$S_q = \frac{1}{2} \Delta q \equiv \frac{1}{2} \int_{-1}^1 dx (q_{\uparrow}(x) - q_{\downarrow}(x))$$

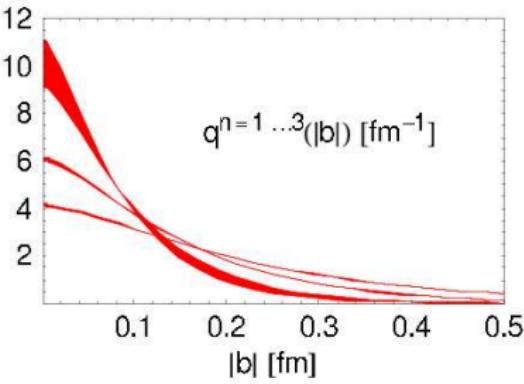
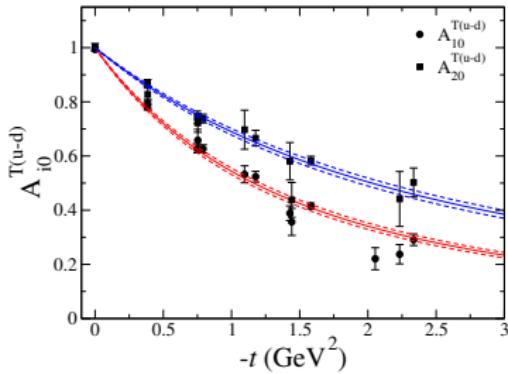
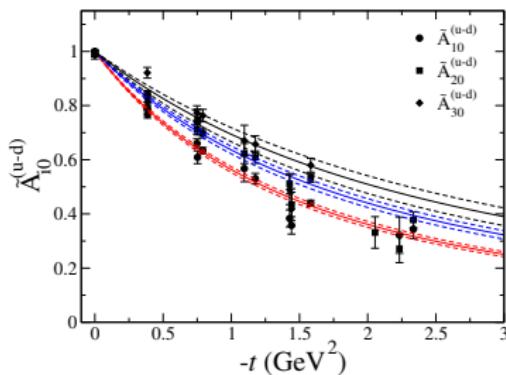
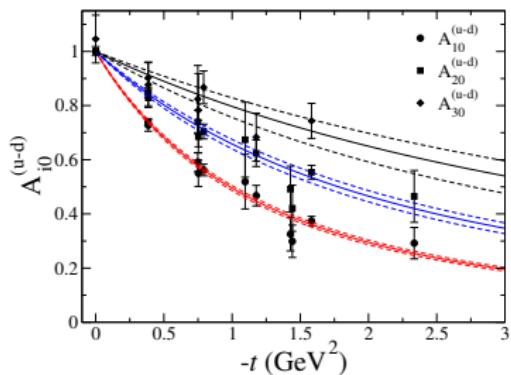
$$J = \sum_q J_q$$



# Generalised Form Factors, ( $m_\pi \approx 950\text{MeV}$ )



# Impact Parameter Space, ( $m_\pi \approx 950\text{MeV}$ )

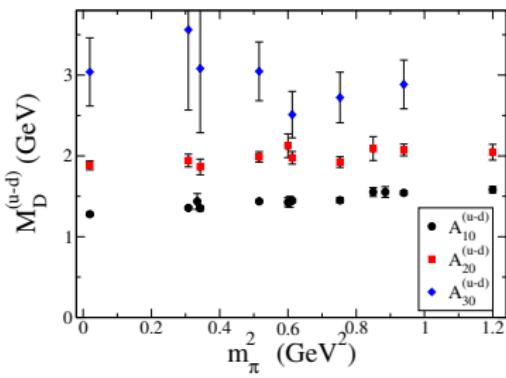
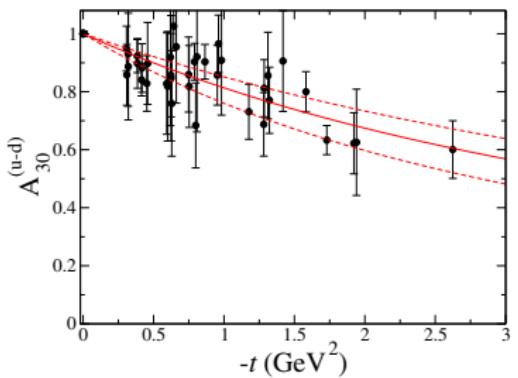
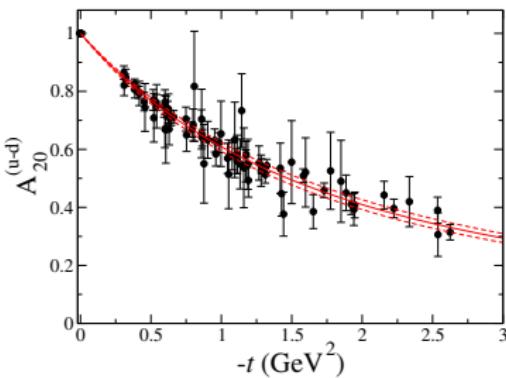
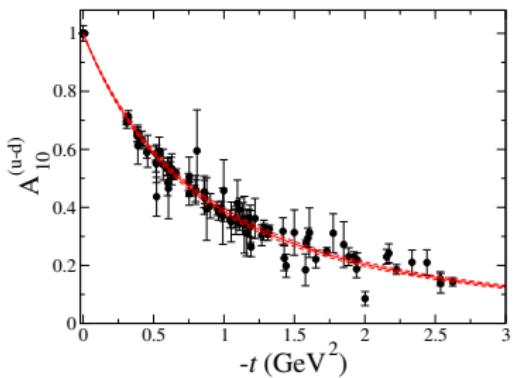


# Impact Parameter Space

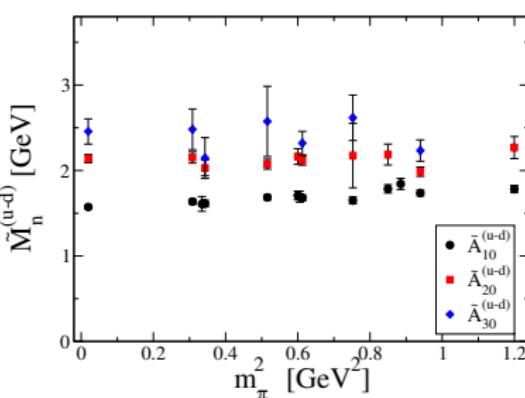
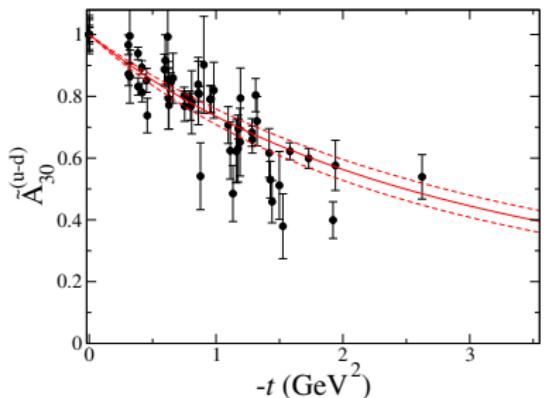
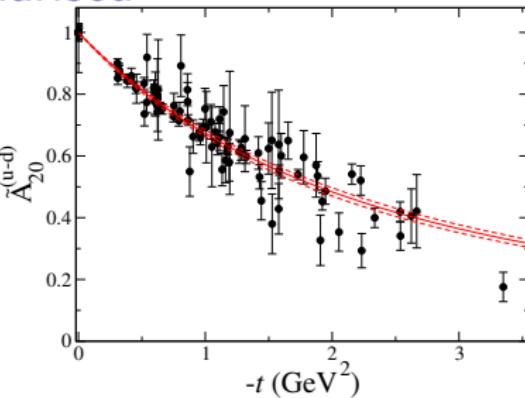
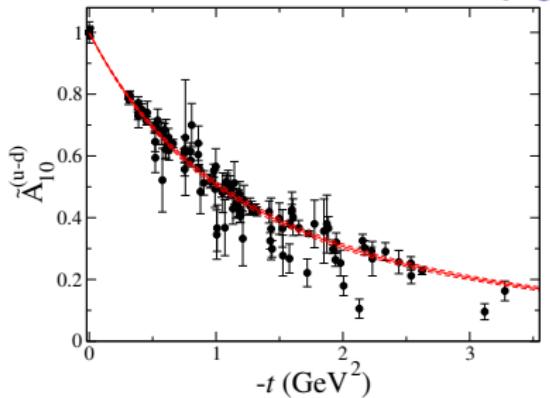
- Flattening of the GFFs  $\Rightarrow$   
quark distribution  $\rightarrow \delta^2(\vec{b}_\perp)$  as  $x \rightarrow 1$
- Form factors well described by a **dipole**
- Extrapolate dipole masses to chiral limit

$$A_{n0}(t) = \frac{A_{n0}(0)}{(1 - t/M_n^2)^2} \rightarrow \bar{A}_{n0}(t) = \frac{1}{(1 - t/(M_n^0 + \alpha_n m_\pi^2))^2}$$

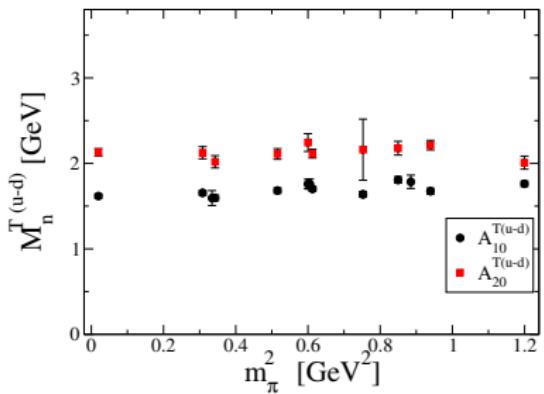
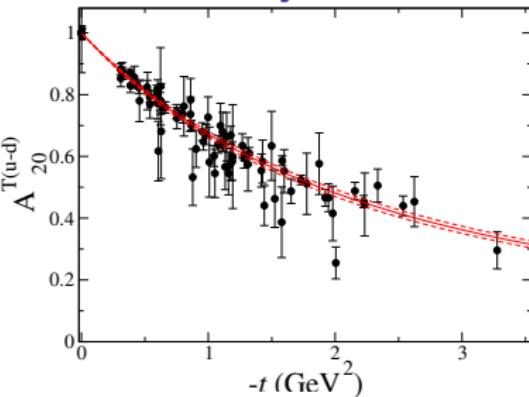
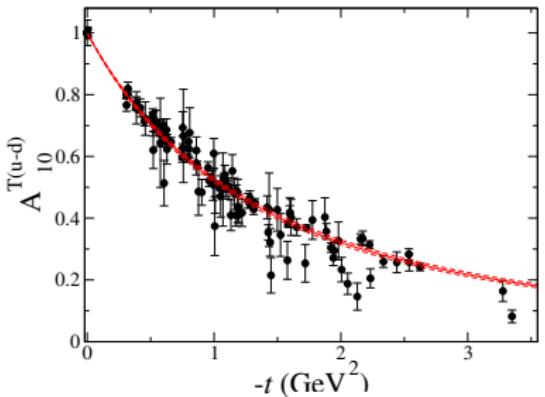
# Described By A Dipole



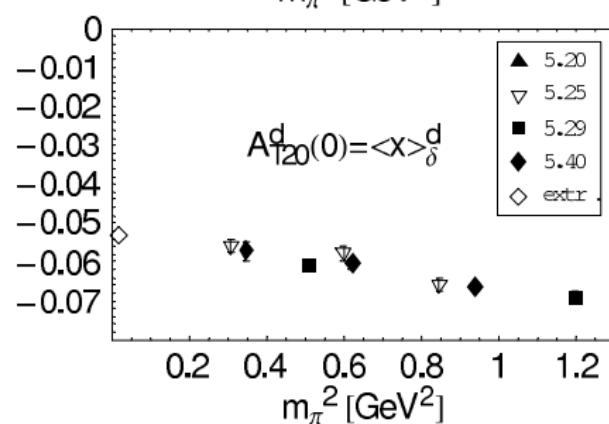
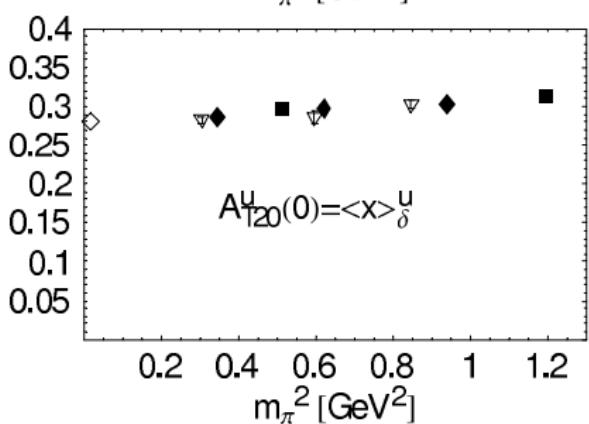
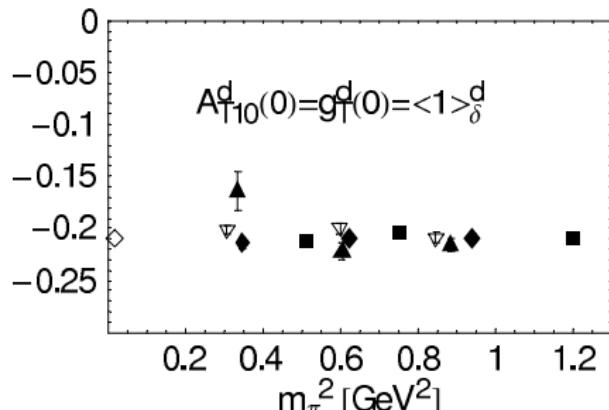
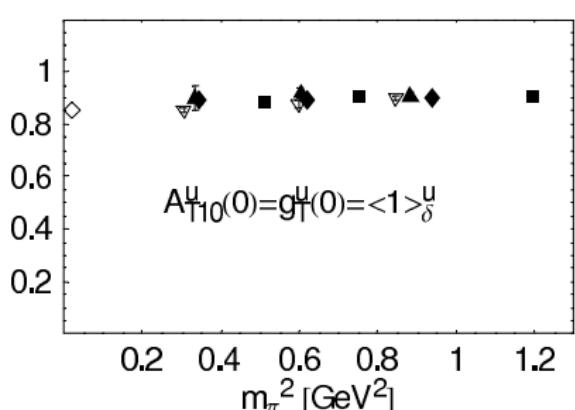
# Polarised



# Generalised Transversity



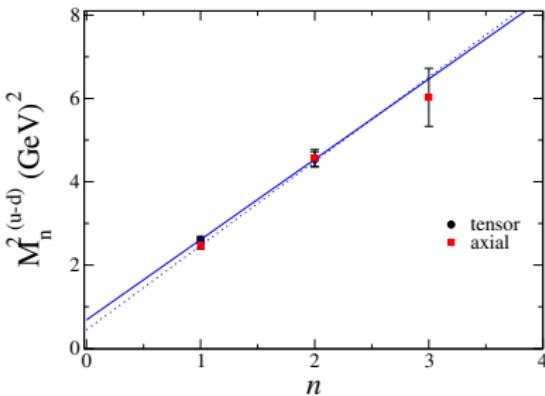
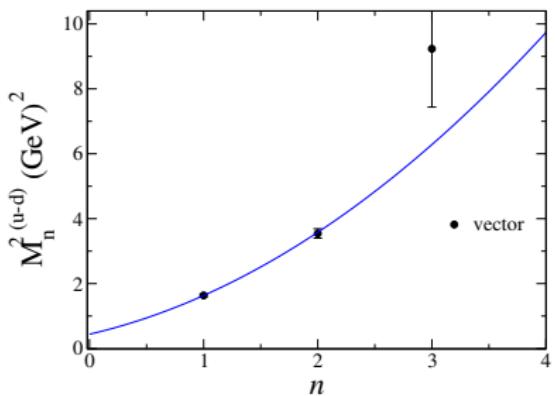
## Generalised Transversity, ( $\overline{\text{MS}}^4$ $\text{GeV}^2$ )



# Regge Trajectory?

Regge:  $\sqrt{M_n^2} = \alpha + n/\alpha'$

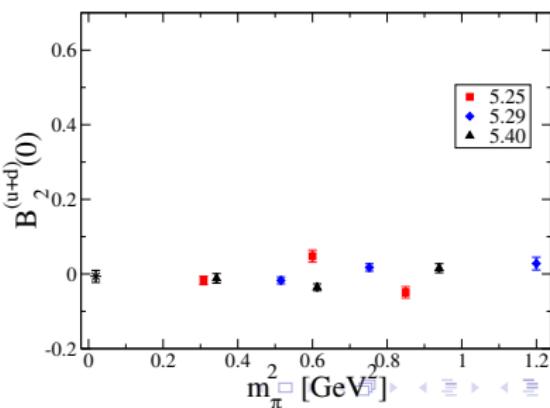
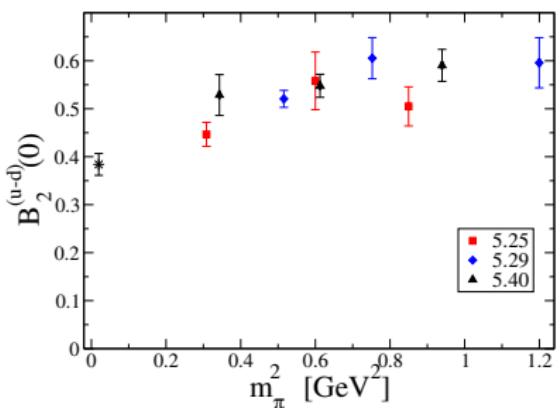
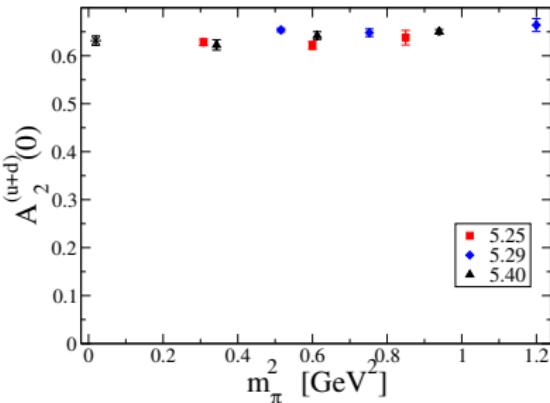
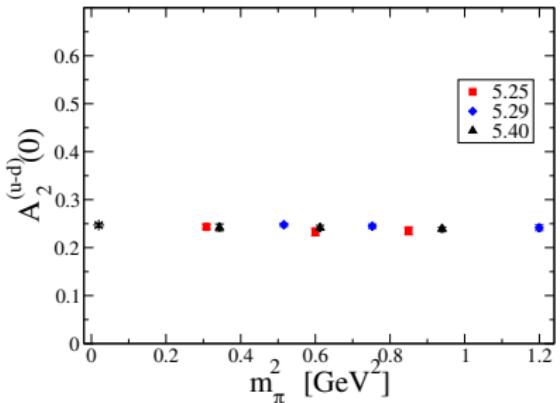
$$M_n^2 = \alpha + n/\alpha'$$



$$\int_{-1}^1 dx x^{n-1} H_q(x, 0, t) = \frac{\langle x_q^{n-1} \rangle}{(1 - t/M_n^2)^2},$$



# Angular Momentum $J^q = L^q + S^q = \frac{1}{2}(A_2^q + B_2^q)$ , ( $\overline{\text{MS}}4$ GeV $^2$ )



# Conclusions and Outlook

- Lattice provides a useful tool for investigating (moments of) GPDs
  - “Flattening” of GFFs  $A_{n0}(t)$  for increasing  $n$
  - $\lim_{x \rightarrow 1} q(x, \vec{b}_\perp) \propto \delta^2(\vec{b}_\perp)$
- Complete Current Analysis
  - Increase statistics for 2<sup>nd</sup> moment
  - Renormalisation
  - Finite volume effects
  - Chiral extrapolation
  - Partially Quenched
- Quark contribution to nucleon spin and angular momentum
- Compute  $H_q(x, 0, t)$  and  $\tilde{H}_q(x, 0, t)$  via inverse Mellin Transform