

*Proton Form Factor Ratio, G_E^P/G_M^P
From
Double Spin Asymmetries*

Spin
Asymmetries of the
Nucleon
Experiment
(E07-003)

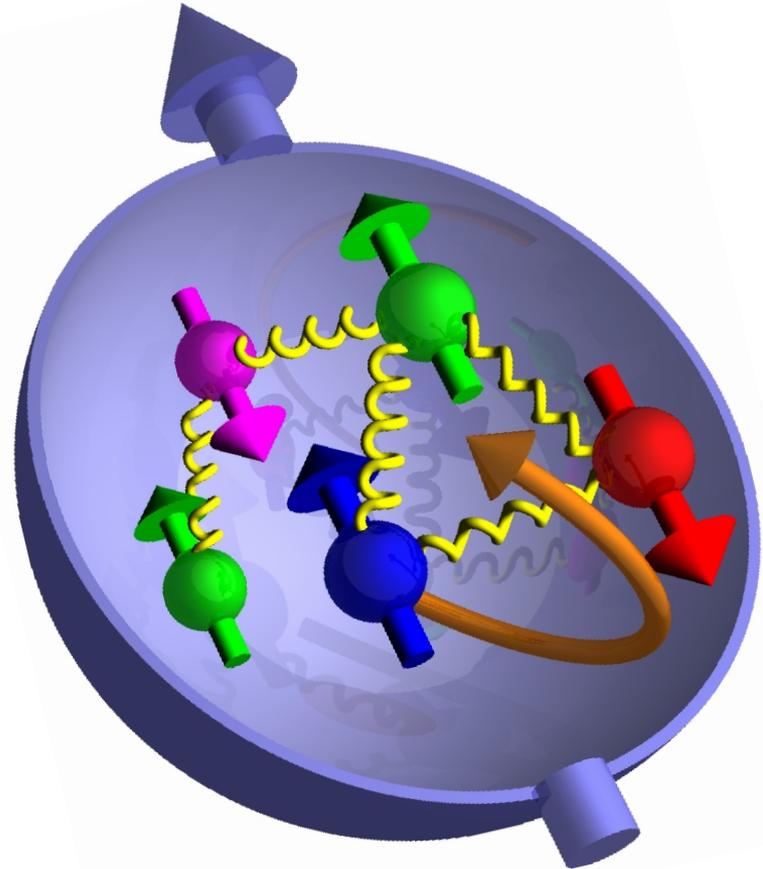


Jefferson Lab
● Thomas Jefferson National Accelerator Facility

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Pizza seminar, October 17, 2012

Outline

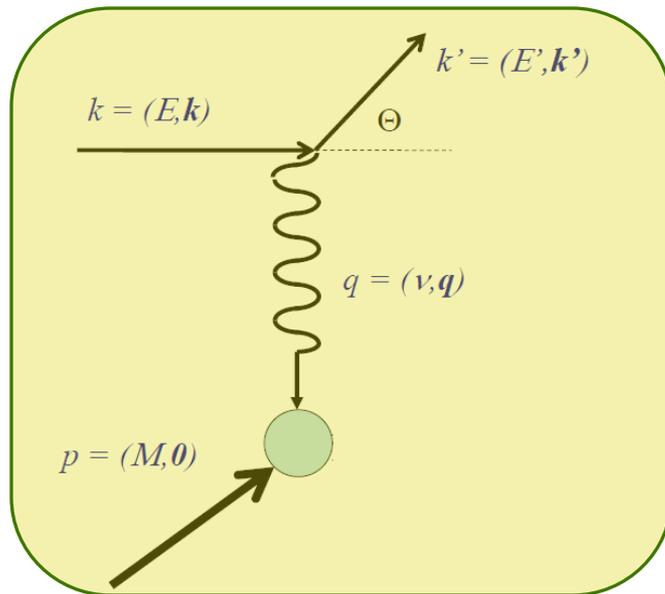
- Introduction
- Physics Motivation
- Experiment Setup
 - BETA Detector
 - HMS Detector
 - Polarized Target
- Elastic Kinematic
- Data Analysis &
MC/SIMC Simulation
- Conclusion



Introduction

Nucleon Elastic Form Factors

- Defined in context of single-photon exchange.
- Describe how much the nucleus deviates from a point like particle.
- Describe the internal structure of the nucleons.
- Provide the information on the spatial distribution of electric charge (by electric form factor, G_E) and magnetic moment (by magnetic form factor, G_M) within the proton.
- Can be determined from elastic electron-proton scattering.
- They are functions of the four-momentum transfer squared, Q^2



The four-momentum transfer squared,

$$Q^2 = -q^2 = 4EE' \sin^2\left(\frac{\Theta}{2}\right)$$

$$E - E' = \frac{Q^2}{2M}$$

General definition of the nucleon form factor is

$$\langle N(P') | J_{EM}^\mu(0) | N(P) \rangle = \bar{u}(P') \left[\gamma^\mu F_1^N(Q^2) + i\sigma^{\mu\nu} \frac{q_\nu}{2M} F_2^N(Q^2) \right] u(P)$$

Sachs Form Factors $G_E = F_1 - \tau F_2$; $G_M = F_1 + F_2$; $\tau = \frac{Q^2}{4M^2}$

F_1 – non-spin flip (Dirac Form Factor) describe the charge distribution

F_2 – spin flip (Pauli form factor) describe the magnetic moment distribution

At low $|q^2|$

$$G_E(q^2) \approx G_E(\vec{q}^2) = \int e^{i\vec{q}\cdot\vec{r}} \rho(\vec{r}) d^3\vec{r}$$

$$G_M(q^2) \approx G_M(\vec{q}^2) = \int e^{i\vec{q}\cdot\vec{r}} \mu(\vec{r}) d^3\vec{r}$$

Fourier transforms of the charge, $\rho(r)$
and magnetic moment, $\mu(r)$ distributions
in Breit Frame

At $q^2 = 0$

$$G_E(0) = \int \rho(\vec{r}) d^3\vec{r} = 1$$

$$G_M(0) = \int \mu(\vec{r}) d^3\vec{r} = \mu_p = +2.79$$

$$\mu \frac{G_E^p}{G_M^p} = 1$$

Form Factor Ratio Measurements

1. Rosenbluth separation method.

- Measured the electron - unpolarized proton elastic scattering cross section at fixed Q^2 by varying the scattering angle, θ_e .
- Strongly sensitive to the radiative corrections.

$$\frac{d\sigma}{d\Omega} = \underbrace{\frac{\alpha^2 E' \cos^2 \frac{\theta_e}{2}}{4(1+\tau)E^3 \sin^4 \frac{\theta_e}{2}}}_{\sigma_{Mott} / (1+\tau)} \left[G_E^2 + \frac{\tau}{\epsilon} G_M^2 \right]$$

$$Q^2 = 2EE'(1 - \cos\theta_e)$$

$$\tau = \frac{Q^2}{4M^2}$$

$$\epsilon = \left[1 + 2(1+\tau) \tan^2 \frac{\theta_e}{2} \right]^{-1}$$

$$\frac{d\sigma}{d\Omega} \cdot \frac{\epsilon(1+\tau)}{\sigma_{Mott}} = G_E^2 \epsilon + \tau G_M^2$$

$$Y = m X + C$$

The gradient = G_E^2 , The Intercept = τG_M^2 ,

E - Incoming electron energy

E' - Outgoing electron energy

θ_e - Outgoing electron's scattering angle

M - Proton mass

2. Polarization Transfer Technique.

- Measured the recoil proton polarization from the elastic scattering of polarized electron-unpolarized proton.
- Insensitive to absolute polarization, analyzing power.
- Less sensitive to radiative correction.

$$\frac{G_E}{G_M} = -\frac{P_T}{P_L} \frac{(E + E') \tan\left(\frac{\theta_e}{2}\right)}{2M_p}$$

E - Incoming going electron energy

E' - Out going electron energy

θ_e - Outgoing electron's scattering angle

M_p - Proton mass

$$P_L = M_p^{-1} (E + E') \sqrt{\tau(1 + \tau)} G_M^2 \tan^2(\theta_e / 2) \longrightarrow \text{Polarization along } q$$

$$P_T = 2\sqrt{\tau(1 + \tau)} G_E G_M \tan(\theta_e / 2) \longrightarrow \text{Polarization perpendicular to } q$$

(in the scattering plane)

$$P_N = 0 \longrightarrow \text{Polarization normal to scattering plane.}$$

3. Double-Spin Asymmetry.

- Measured the cross section asymmetry between + and – electron helicity states in elastic scattering of a polarized electron on a polarized proton.
- The systematic errors are different when compared to either the Rosenbluth technique or the polarization transfer technique.
- The sensitivity to the form factor ratio is the same as the Polarization Transfer Technique.

$$A_p = \frac{-br \sin \theta^* \cos \phi^* - a \cos \theta^*}{r^2 + c}$$

$$\frac{G_E}{G_M} = -\frac{b}{2A_p} \sin \theta^* \cos \phi^* + \sqrt{\frac{b^2}{4A_p^2} \sin^2 \theta^* \cos^2 \phi^* - \frac{a}{A_p} \cos \theta^* - c}$$

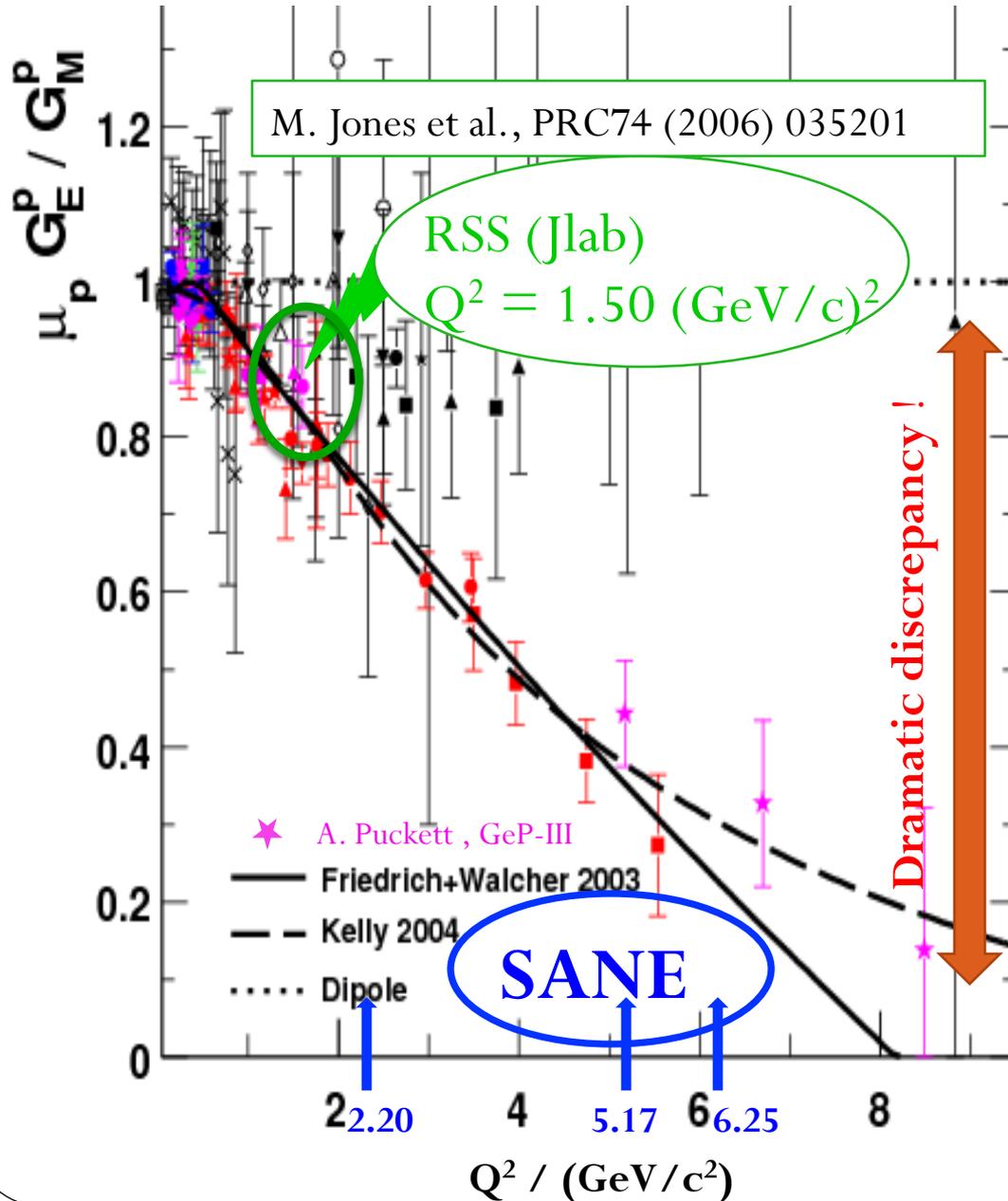
Here, $r = G_E / G_M$

$a, b, c =$ kinematic factors

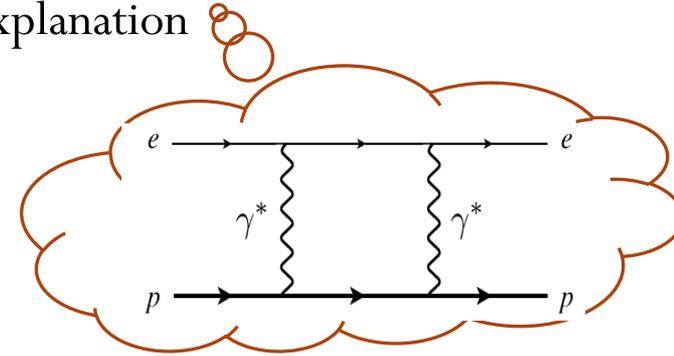
$\theta^*, \phi^* =$ pol. and azi. Angles between \vec{Q} and \vec{S}

$A_p =$ The beam - target asymmetry

Physics Motivation



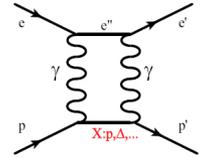
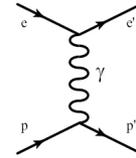
- Dramatic discrepancy between Rosenbluth and recoil polarization technique.
- Multi-photon exchange considered the best candidate for the explanation



- **Double-Spin Asymmetry** is an Independent Technique to verify the discrepancy

Two-Photon Exchange

- Both Rosenbluth method and the polarization transfer technique account for radiative correction, but neither consider two photon exchange.



- Contribution of the TPE amplitude has calculated theoretically and, has an ϵ dependence that has the same sign as the G_E contribution to the cross section and is large enough to effect the extracted value of G_E .

Therefore, the extracted G_E/G_M for the Rosenbluth technique is reduced.

- The effect of TPE amplitude on the polarization components is small, though the size of the contribution change with ϵ
- The size of the TPE would measure by taking the ϵ dependence of the ratio of cross sections, R for elastic electron-proton scattering to positron-proton scattering at a fixed Q^2 and measuring the deviation from 1.

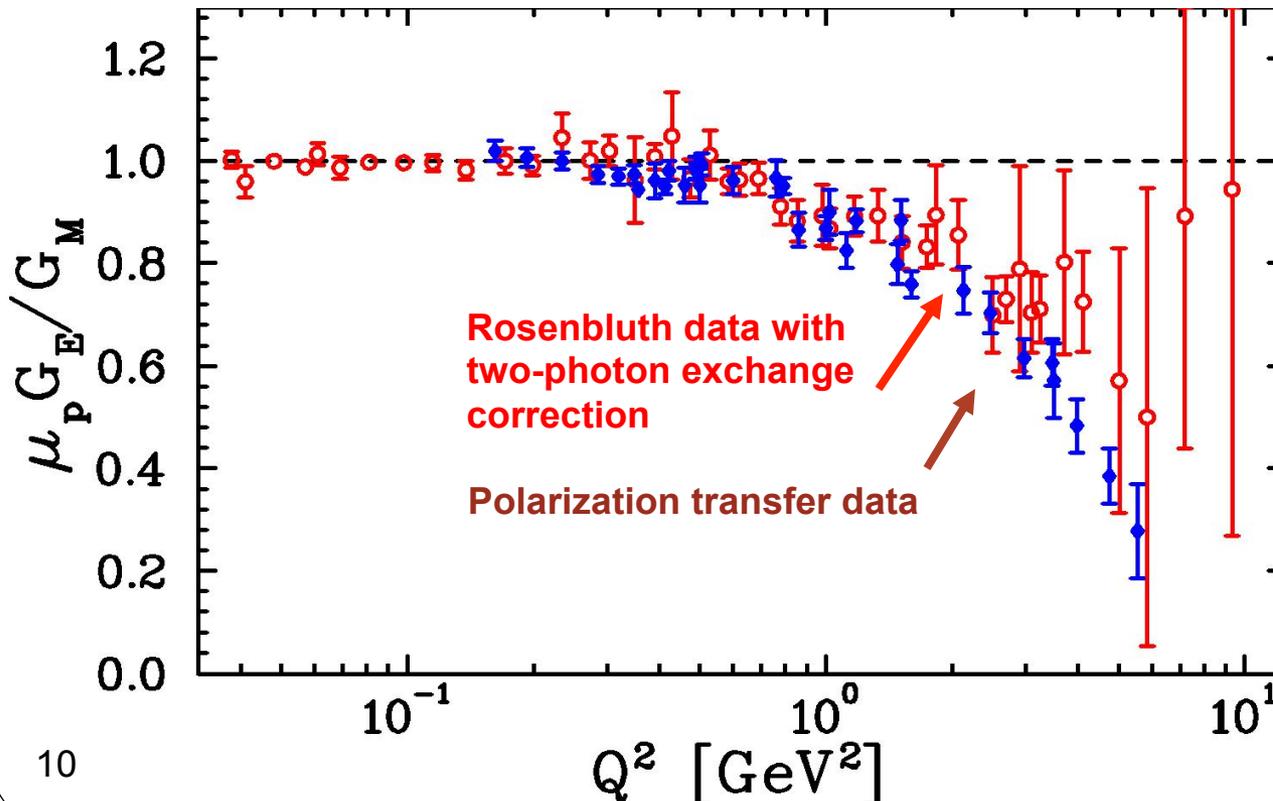
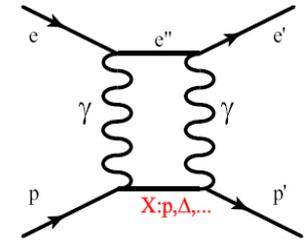
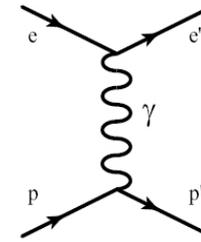
$$R = \frac{\sigma_{e^+}}{\sigma_{e^-}} = \frac{(A_{1\gamma} + A_{2\gamma})^2}{(A_{1\gamma} - A_{2\gamma})^2} \approx 1 + 4 \operatorname{Re} \left(\frac{A_{2\gamma}}{A_{1\gamma}} \right)$$

Two-Photon Exchange: Exp. Evidence

Two-photon exchange theoretically suggested

TPE can explain form factor discrepancy

J. Arrington, W. Melnitchouk, J.A. Tjon,
Phys. Rev. C 76 (2007) 035205



Asymmetry measurements

$$\sigma = \sigma_0 + P_E P_T \Delta\sigma$$



$$\sigma_{++} = \sigma_0 + P_E P_T \Delta\sigma$$

$$\sigma_{+-} = \sigma_0 - P_E P_T \Delta\sigma$$

$$\frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = P_E P_T \cdot \frac{\Delta\sigma}{\sigma_0} = \frac{N_+ - N_-}{N_+ + N_-} = A_r$$

$$\frac{A_r}{P_E P_T} = \frac{\Delta\sigma}{\sigma_0} = A_p$$

Hence,

A_p , known as the physics asymmetry is the relative scattering cross section correction due to the spin.

A_r is the raw asymmetry

σ - Scattering cross section

σ_0 - Scattering cross section at unpolarized target

σ_B - Scattering cross section from background

$\Delta\sigma$ - σ correction due to the spin

P_E - Beam polarization

P_T - Target polarization

f - Dilution factor

With background....

$$\sigma_{++} = \sigma_0 + P_E P_T \Delta\sigma + \sigma_B$$

$$\sigma_{+-} = \sigma_0 - P_E P_T \Delta\sigma + \sigma_B$$

$$A_r = P_E P_T \cdot \frac{\Delta\sigma}{(\sigma_0 + \sigma_B)}$$

$$A_r = P_E P_T \cdot \frac{\Delta\sigma}{\sigma_0} \cdot \frac{\sigma_0}{(\sigma_0 + \sigma_B)} f$$

$$A_p = \frac{A_r}{f P_E P_T}$$

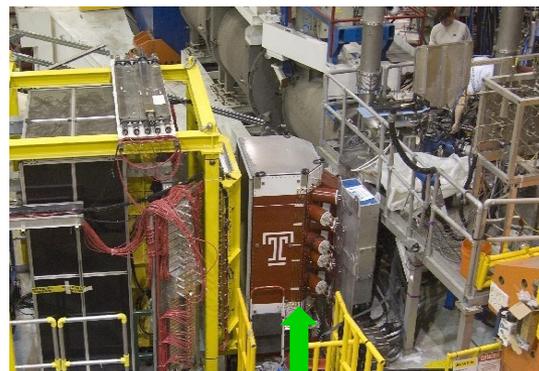
Experiment Setup



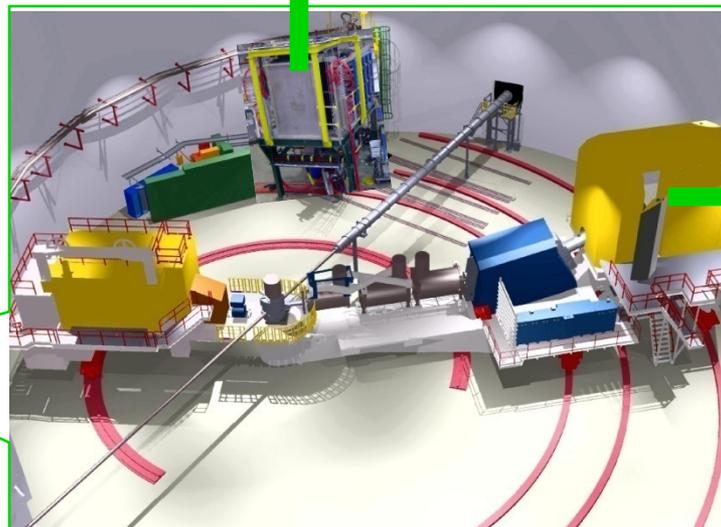
Hall C at
Jefferson Lab

Elastic ($e, e'p$) scattering from the polarized NH_3 target using a longitudinally polarized electron beam

(Data collected from Jan – March, 2009)



- BETA for coincidence electron detection
- Central scattering angle : 40°
- Over 200 msr solid angle coverage

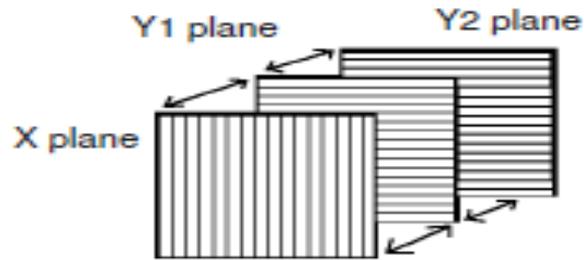


- HMS for the scattered proton detection
- Central angles are 22.3° and 22.0°
- Solid angle ~ 10 msr

Big Electron Telescope Array – BETA

Forward Tracker

- 3 planes of Bicron Scintillator provide early particle tracking

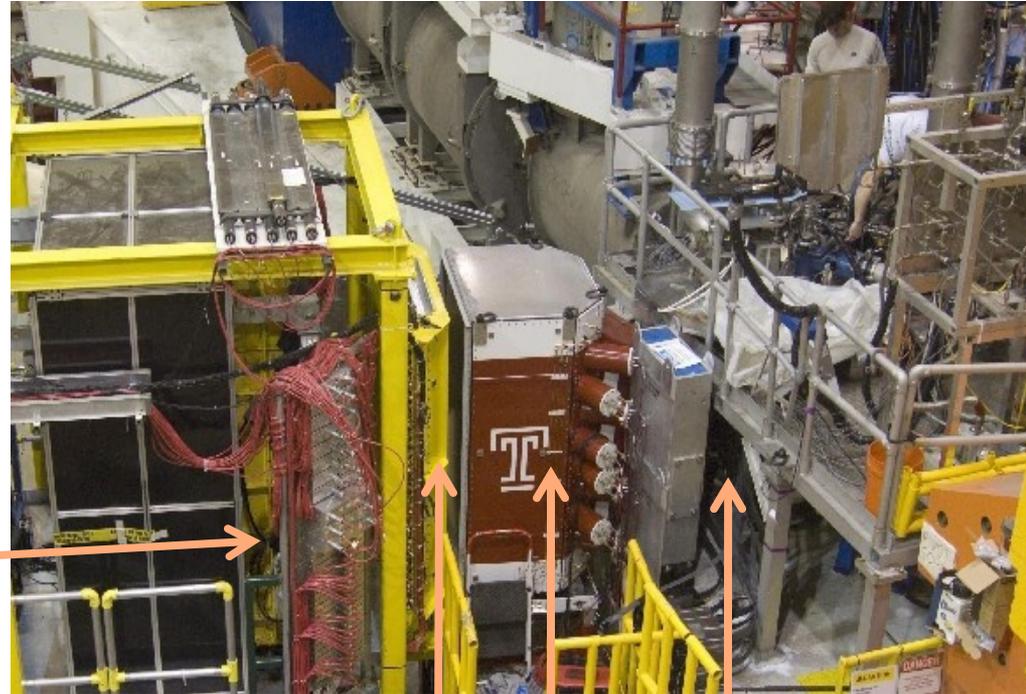


Cerenkov

- N₂ gas cerenkov
- Provides particle ID
- 8 mirrors and 8 PMTs

Lucite Hodoscope

- 28 bars of 6cm wide Lucite
- Bars oriented horizontally for Y tracking
- PMTs on either side of bar provides X resolution



BigCal

Lucite Hodoscope

Tracker

Cerenkov

Big Cal (GEP III Collaboration)

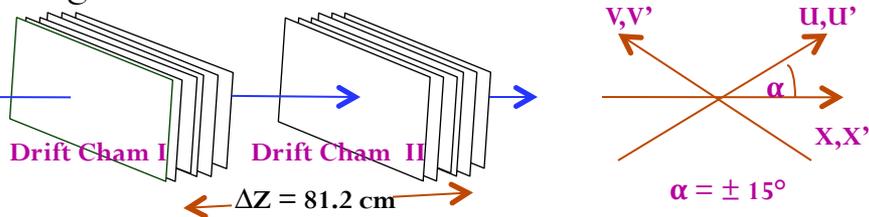
Lead glass calorimeter

- 1744 blocks aprx. 4cm x 4cm
- energy and position measurement

High Momentum Spectrometer – HMS

Drift Chambers

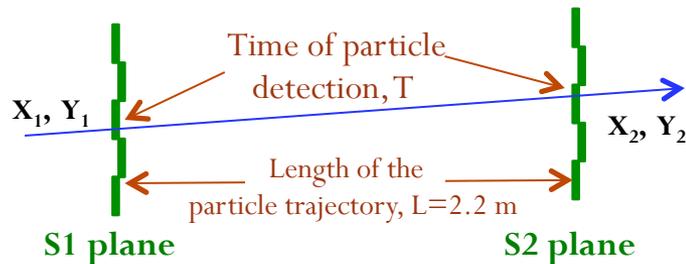
- Each plane has a set of alternating field and sense wires Filled with an equal parts Argon-Methane mixture



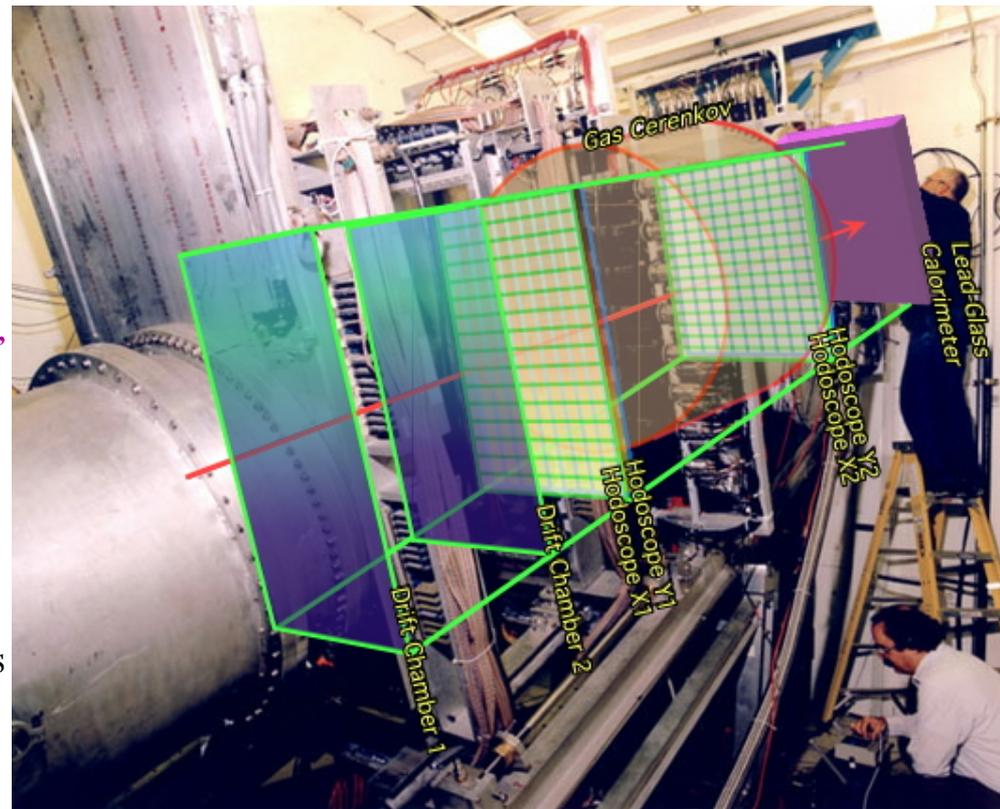
- Track particle trajectory by multiple planes.
- χ^2 fitting to determine a straight trajectory.

Hodoscopes

- Each plane contains 10 to 16 Scintillator paddles with PMTs on both ends
- Each Paddle is 1.0 cm thick and 8.0 cm wide



- Fast position determination & triggering
- Time of Flight (TOF) = $T_2 - T_1$ determines β
($\beta = L/c \times \text{TOF}$)



Gas Cerenkov

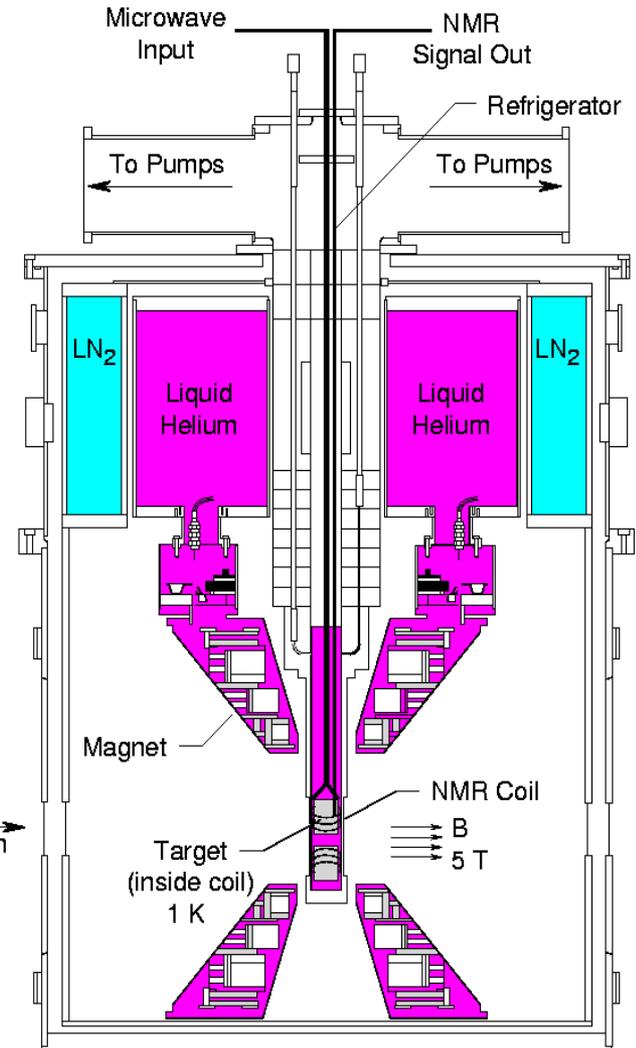
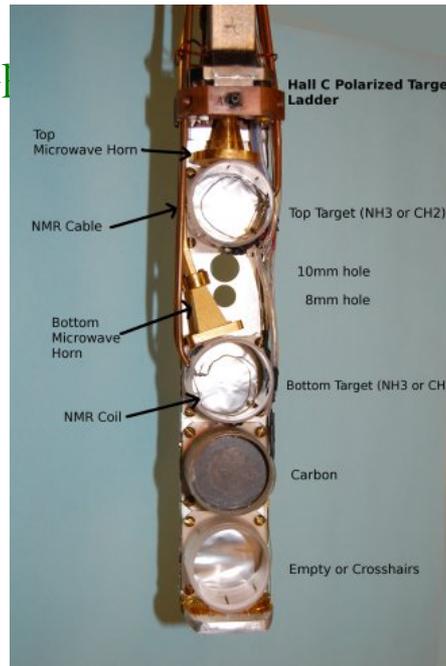
- Two mirrors (top & bottom) connected to two PMTs
- Used as a Particle ID

Lead Glass Calorimeter

- 4 layers of 10 cm x 10cm x 70cm blocks stacked 13 high.
- Used as a Particle ID

Polarized Target

- C, CH₂ and NH₃
- Dynamic Nuclear Polarization (DNP) polarized the protons in the NH₃ target up to 90% at
1 K Temperature
5 T Magnetic Field
- Temperature is maintained by immersing the entire target in the liquid He bath
- Used microwaves to excite spin flip transitions (55 GHz - 165 GHz)
- Polarization measured using NMR coils
- To maintain reasonable target polarization, the beam current,
 - limited to 100 nA
 - Was uniformly rastered.



The Polarized Target Assembly

Goal Of The SANE

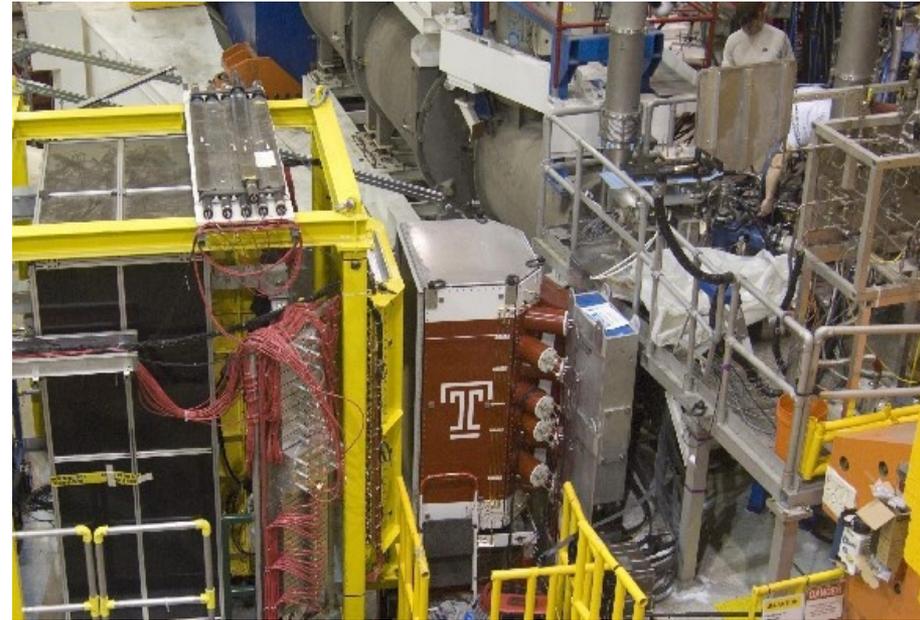
- SANE is a single arm inclusive scattering experiment. Used
 - **Big Electron Telescope Array** – **BETA** In single arm mode
 - **High Momentum Spectrometer** – **HMS** in both single arm and coincidence mode

Physics from BETA:

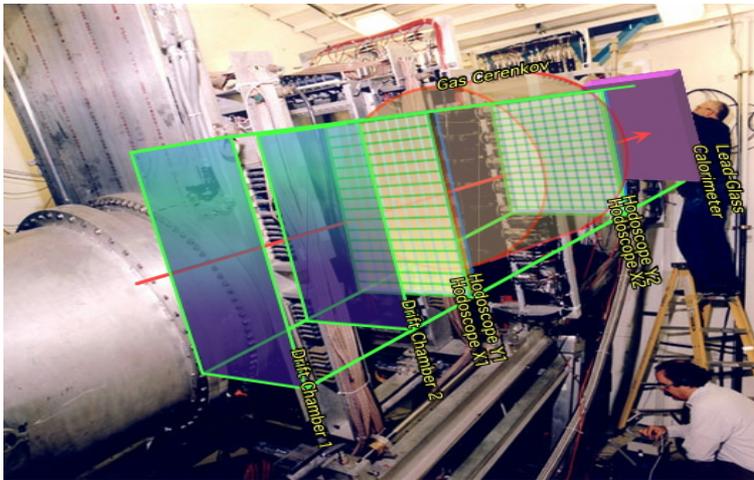
- Measure proton spin structure function $g_2(X, Q^2)$ and spin asymmetry $A_1(X, Q^2)$ at four-momentum transfer $2.5 < Q^2 < 6.5 \text{ GeV}^2$ and $0.3 < X < 0.8$

by measuring anti-parallel and near-perpendicular spin asymmetries.

- Study twist -3 effects (d2 matrix element) and moments of g_2 and g_1
- Comparison with Lattice QCD, QCD sum rule
- Explore “High” X_B region: A_1 at $X_B \sim 1$



What HMS use for



Packing Fraction determination.

- It will detect electrons with momenta from 1 to around 5 GeV/c
- Use the ratio of data/MC yields of C and C+He.

- In single arm mode HMS can be used to measure accurate pair symmetric backgrounds from $\gamma \rightarrow e^+e^-$ pair production.

HMS will detect positrons up to 2.2 GeV/c.

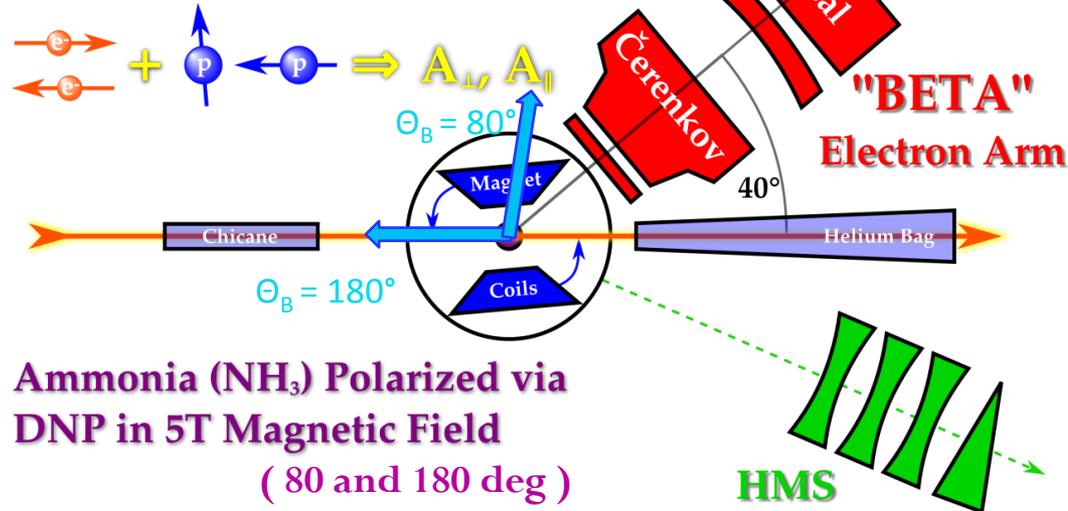
Asymmetries

- Inclusive Asymmetries;
 Q^2 of 0.8, 1.3 and 1.8 (GeV²)
- Elastic Asymmetries at magnetic field of 80⁰ and hence the ratio of form factors, G_E^p/G_M^p
 - From single arm data at $Q^2 = 2.2$ GeV²
 - From coincidence data at $Q^2 = 5.17$ and 6.25 GeV²

Polarized Target Magnetic Field

Polarized Electron Beam: 4.7, 5.9 GeV

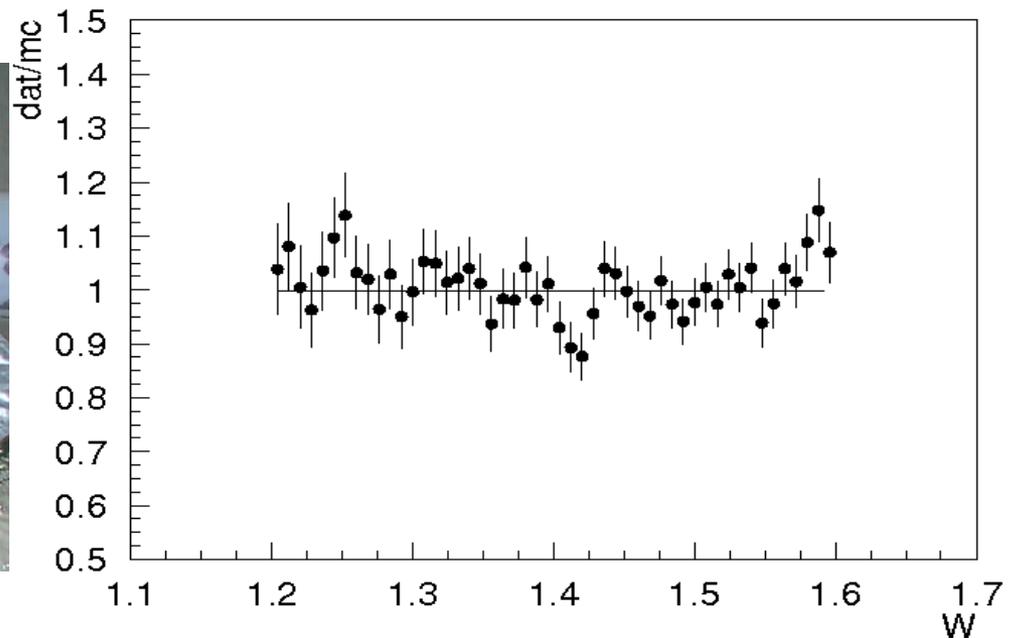
Polarized Proton Target: $\sim \perp, \parallel$



- Used only perpendicular magnetic field configuration for the elastic data
 - Average target polarization is $\sim 70\%$
 - Average beam polarization is $\sim 73\%$

Packing Fraction.

- Packing Fraction is the actual amount of target material used.
- Determined by taking the ratio of data to MC as a function of W .
- Need to determine the packing fractions for each of the NH_3 loads used during the data taking.



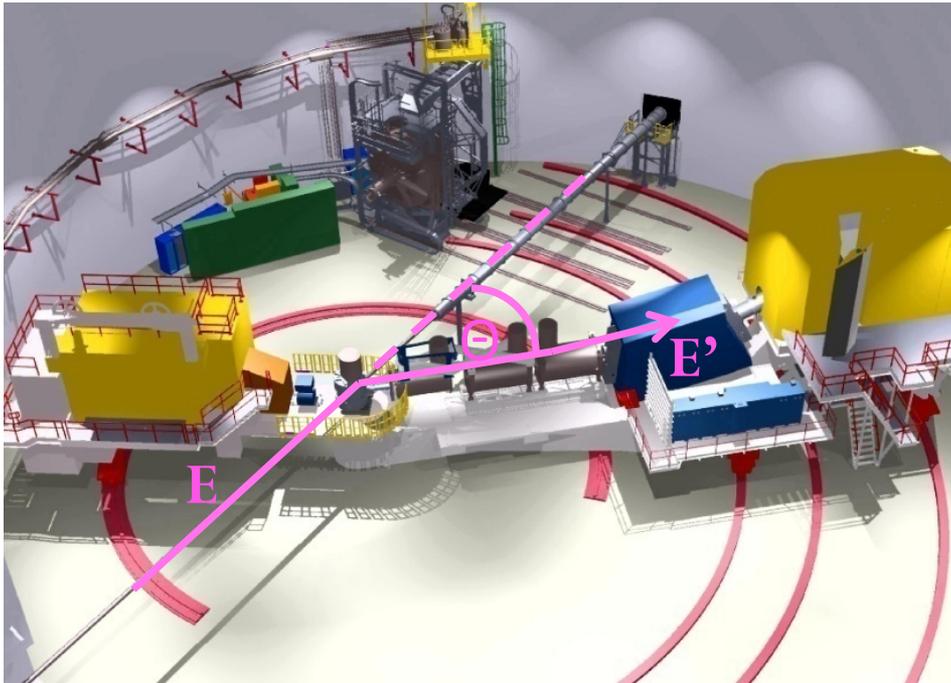
Elastic Kinematics

(From HMS Spectrometer)

Spectrometer mode	Coincidence	Coincidence	Single Arm
HMS Detects	Proton	Proton	Electron
E Beam GeV	4.72	5.89	5.89
P_{HMS} GeV/c	3.58	4.17	4.40
Θ_{HMS} (Deg)	22.30	22.00	15.40
Q^2 (GeV/c) ²	5.17	6.26	2.20
Total Hours (h)	~40 (~44 runs)	~155 (~135 runs)	~12 (~15 runs)
Elastic Events	~113	~1200	-

Data Analysis

Electrons in HMS



$$\vec{e}^- \vec{p} \longrightarrow e^- p$$

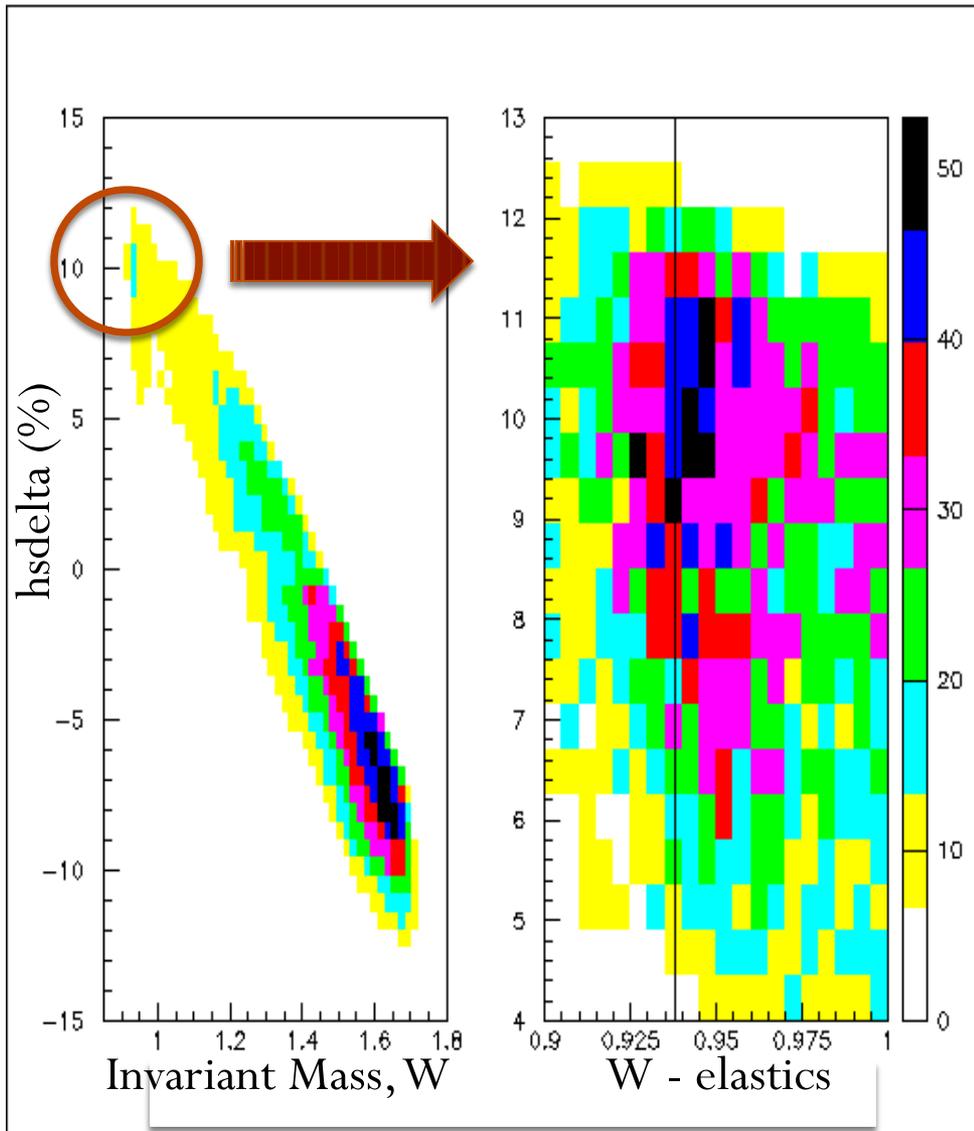
By knowing,
the incoming beam energy, E ,
scattered electron energy, E'
and
the scattered electron angle, θ

$$Q^2 = 4EE' \sin^2\left(\frac{\theta}{2}\right)$$



$$W^2 = M^2 - Q^2 + 2M(E - E')$$

- Momentum Acceptance



$$hsdelta = \left(\frac{P - P_c}{P_c} \right) = \frac{\delta p}{p}$$

P - Measured momentum in HMS

P_c - HMS central momentum

The elastic data are outside of the usual delta cut $\pm 8\%$

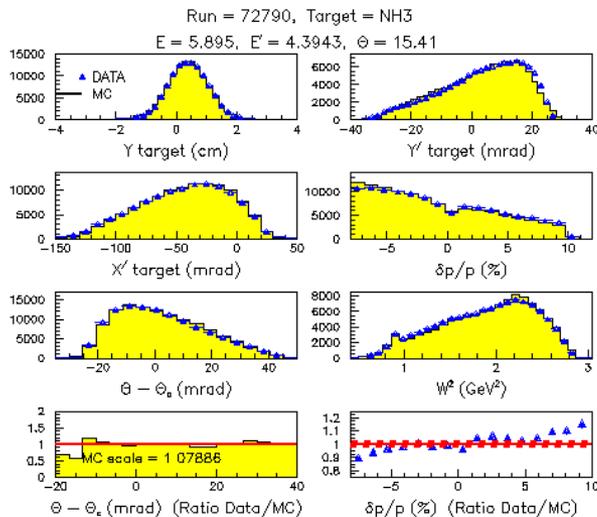
Because HMS reconstruction matrix elements work fine up to 10



Use $-8\% < hsdelta < 10\%$

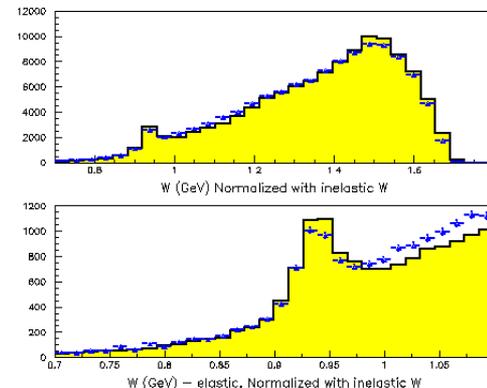
MC with NH3

- Generated N, H and He separately.
- Added Al come from target end caps and 4K shields as well.
- Calculated the MC scale factor using the data/MC luminosity ratio for each target type.
- Added all targets together by weighting the above MC scale factors.
- Used 60% packing fraction.
- Adjust acceptance edges in Ytar and yptar from adjusting the horizontal beam position.
- Adjust the vertical beam position to bring the W peak to 0.938 GeV



$$srast_x = -0.40 \text{ cm}$$

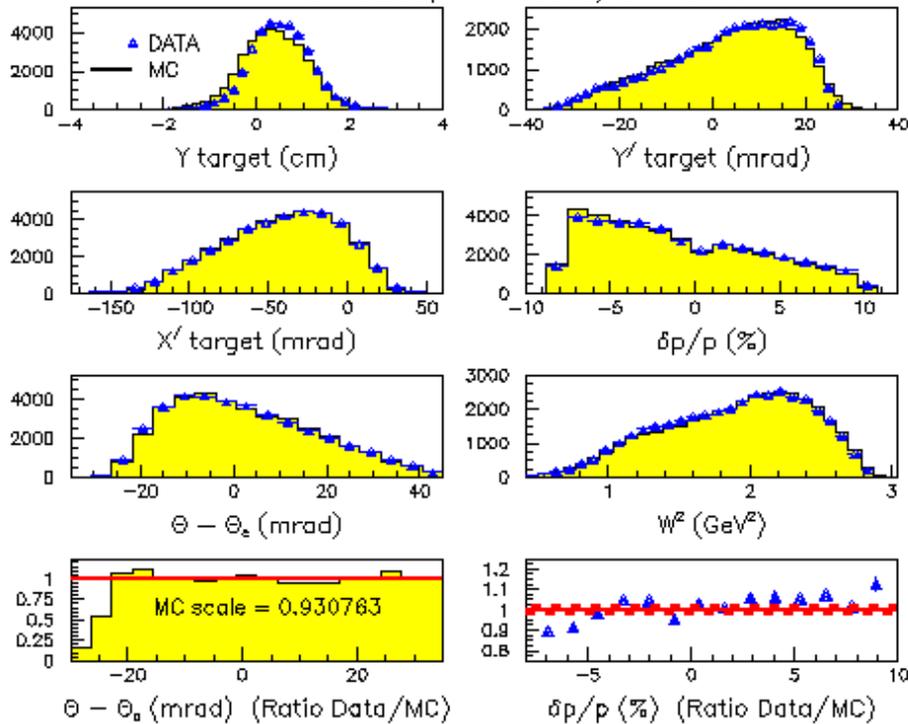
$$srast_y = 0.10 \text{ cm}$$



MC for C run

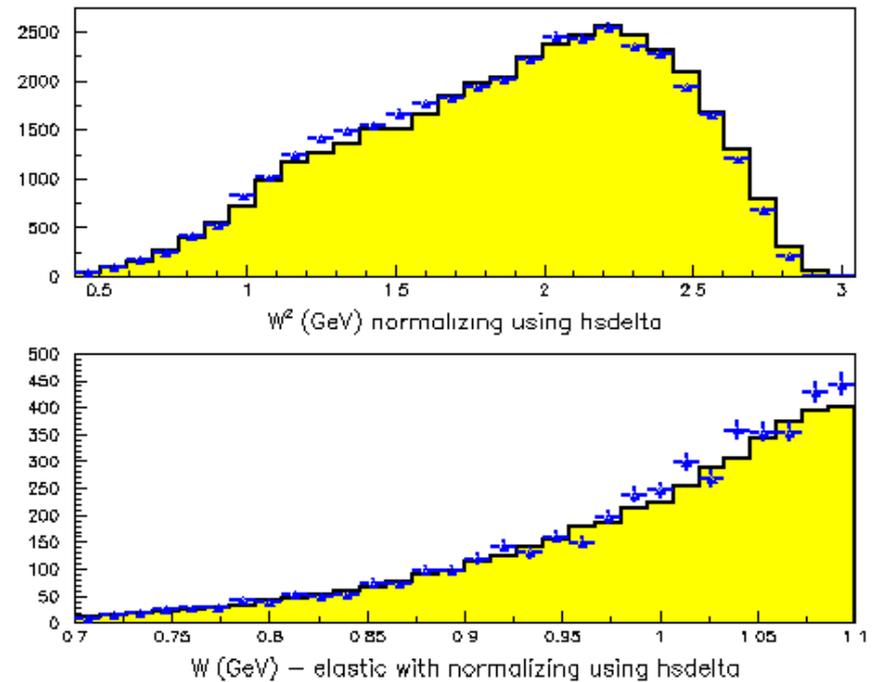
Run = 72782, Target = C

$E = 5.895$, $E' = 4.3943$, $\Theta = 15.41$



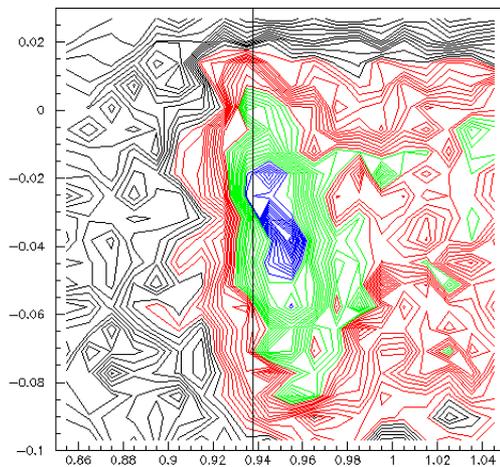
Srast x offset = -0.4 cm

Srast y offset = 0.1 cm



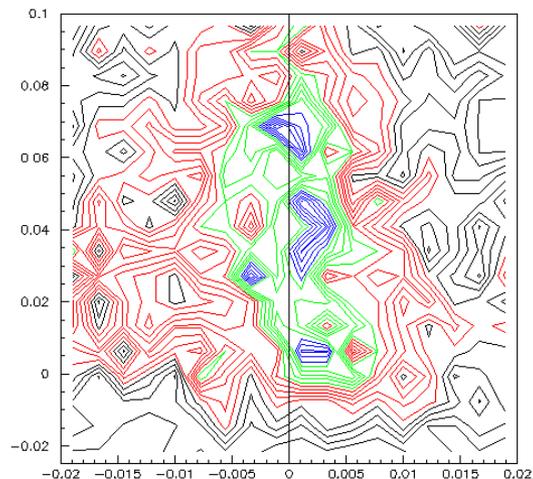
Perp. target magnetic field make some correlations....

In Single Arm electron data



Xptar vs W

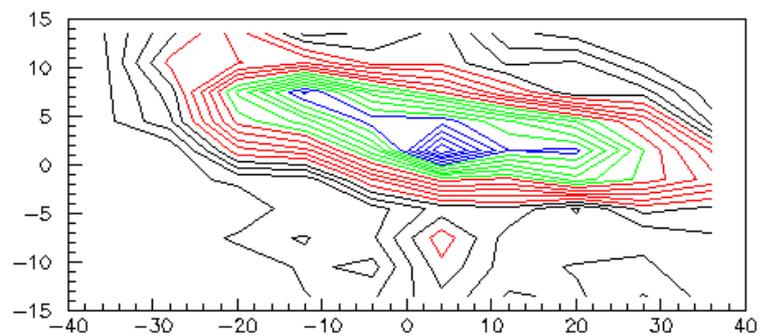
In COIN HMS data



Xptar vs dpel_hms

- Introduced an 'azimuthal angle correction' which correct the target magnetic field in vertical direction in terms of the azimuthal angle. (First make the same correlations on MC/SIMC by applying the correction only for the forward direction and then use the correction on data)
- Different corrections for different detector angles.

In COIN BETA data



Y_HMS-y_clust vs y_clust

Extract the electrons

- Used only Electron selection cuts.

of Cerenkov photoelectrons > 2

$$E_{sh}/E' > 0.7$$

$$\left(\frac{P - P_c}{P_c} \right) < 10 \text{ and } \left(\frac{P - P_c}{P_c} \right) > -8$$

- Cerenkov cut

- Calorimeter cut

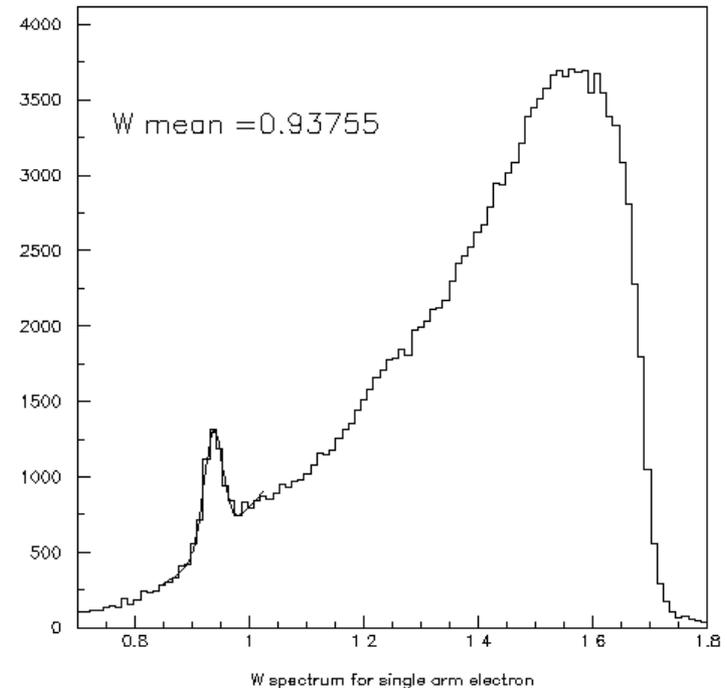
- HMS Momentum Acceptance cut

Here,

P/E' - Detected electron momentum/
energy at HMS

P_c - Central momentum of HMS

E_{sh} - Total measured shower energy
of a chosen electron track by
HMS Calorimeter



Extracted the Asymmetries

The raw asymmetry, A_r

$$A_r = \frac{N^+ - N^-}{N^+ + N^-}$$

$$\Delta A_r = \frac{2\sqrt{N^+} \sqrt{N^-}}{(N^+ + N^-)\sqrt{(N^+ + N^-)}}$$

$N^+ / N^- =$ Charge and life time normalized counts for the +/- helicities

$\Delta A_r =$ Error on the raw asymmetry

$P_B P_T =$ Beam and Target polarization

$N_c =$ A correction term to eliminates the contribution from quasi-elastic ^{15}N scattering under the elastic peak

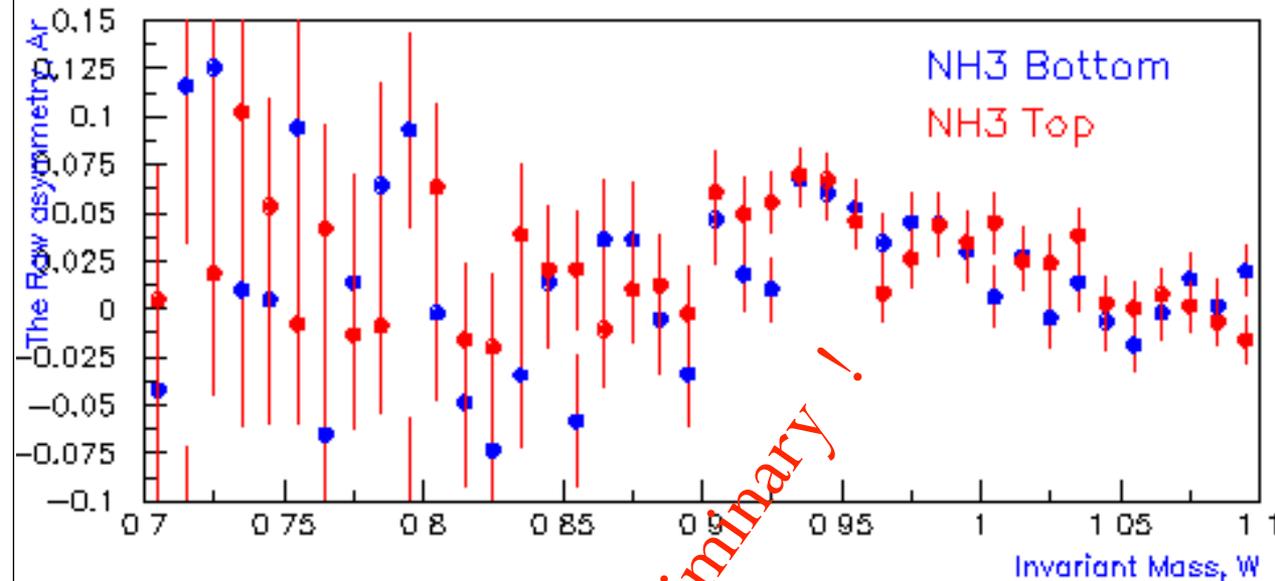
The Asymmetries

Need

dilution factor, f
in order to determine the
physics asymmetry,

$$A_p = \frac{A_r}{f P_B P_T} + N_c$$

and G_E^p / G_M^p
(at $Q^2 = 2.2 \text{ (GeV/c)}^2$)



Preliminary!

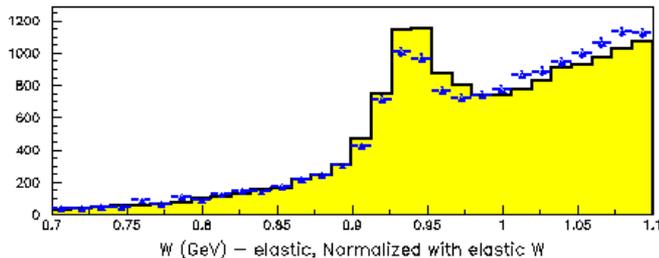
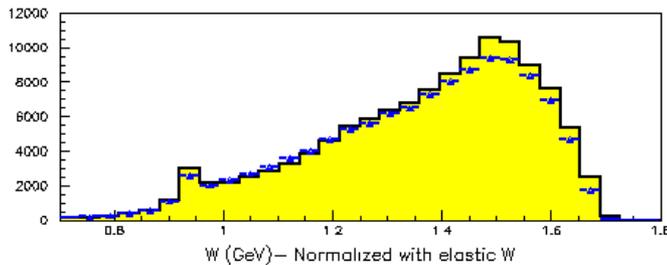
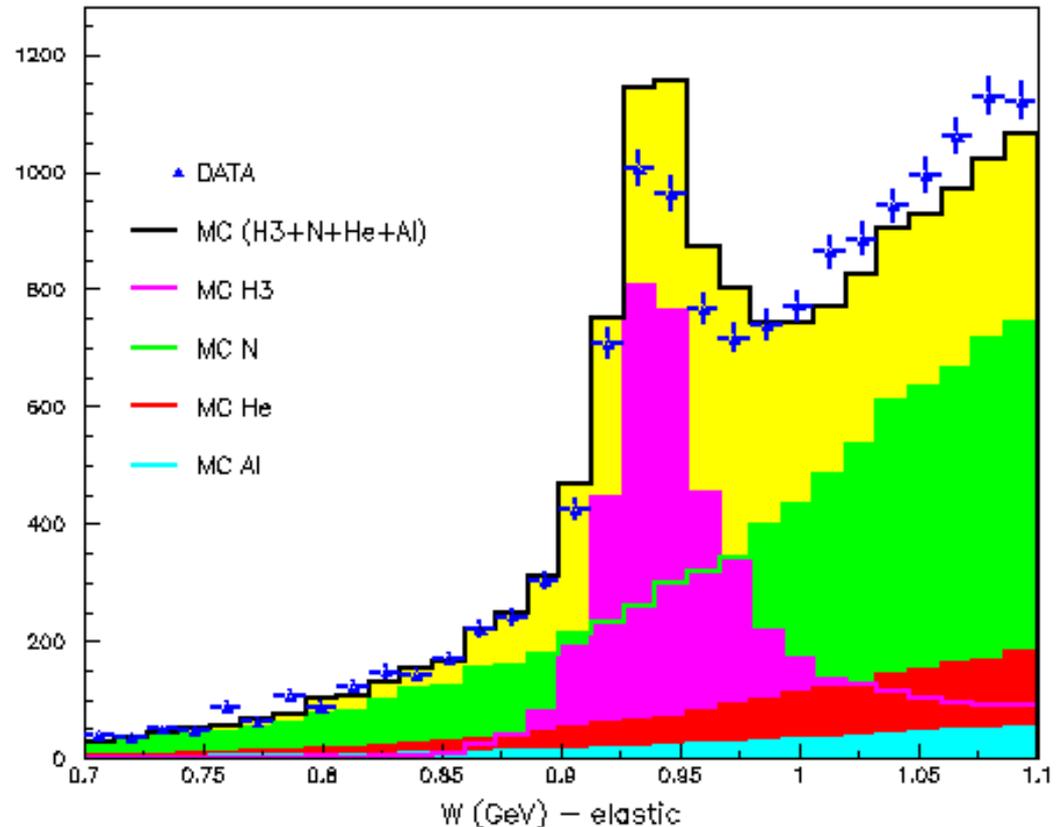
Determination of the Dilution Factor

What is the Dilution Factor ?

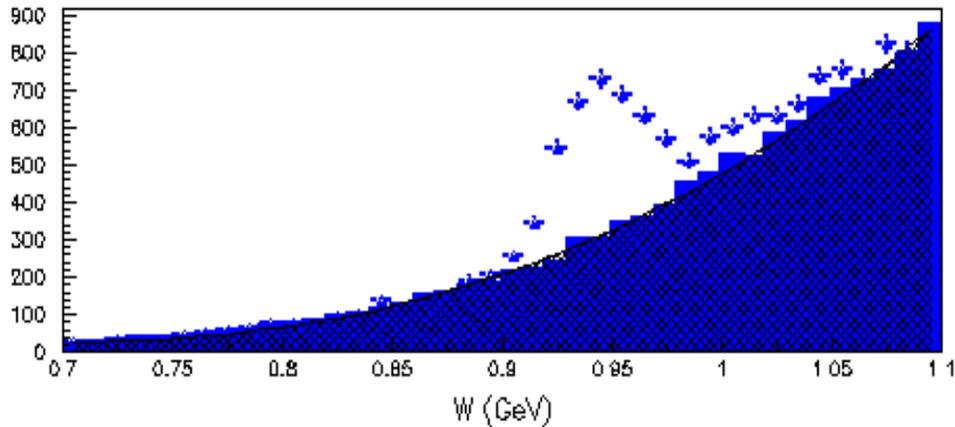
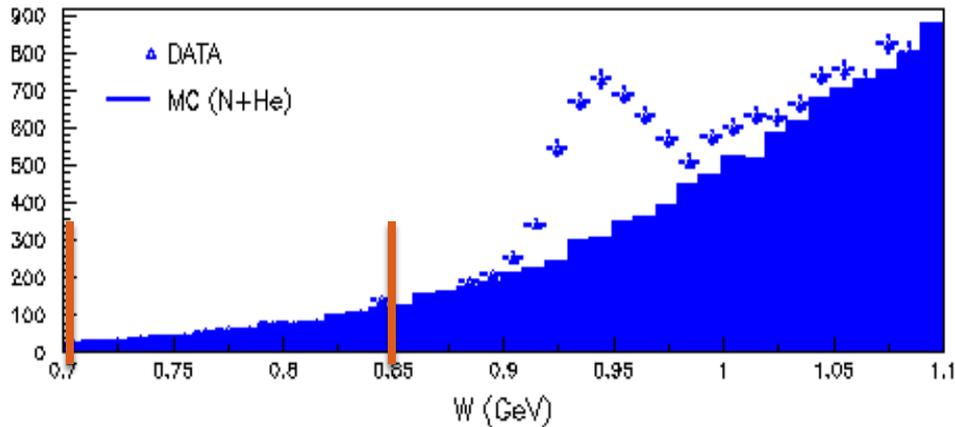
The dilution factor is the ratio of the yield from scattering off free protons (protons from H in NH₃) to that from the entire target (protons from N, H, He and Al)

Dilution Factor,

$$F = \frac{Yield_{Data} - Yield_{MC(N+He)}}{Yield_{Data}}$$



- MC Background contributions (Only He+N+Al)



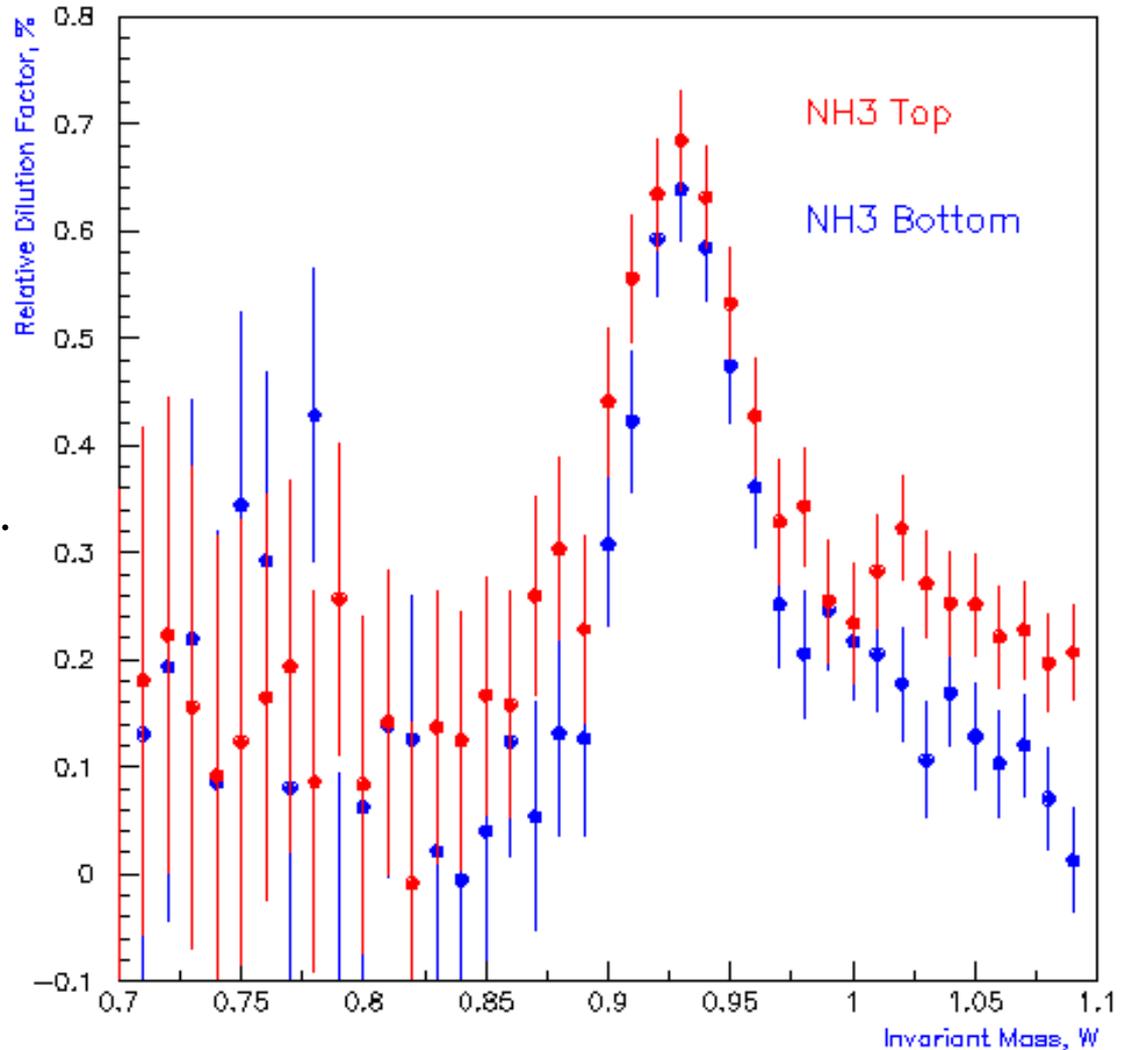
- Calculate the ratio of $\text{Yield}_{\text{Data}} / \text{Yield}_{\text{MC}}$ for the W region $0.7 < W < 0.85$ and MC is normalized with this new scaling factor.
- Used the polynomial fit to N+ He+Al in MC and
- Subtract the fit function from data

▪ The relative Dilution Factor (**Preliminary**)

Dilution Factor,

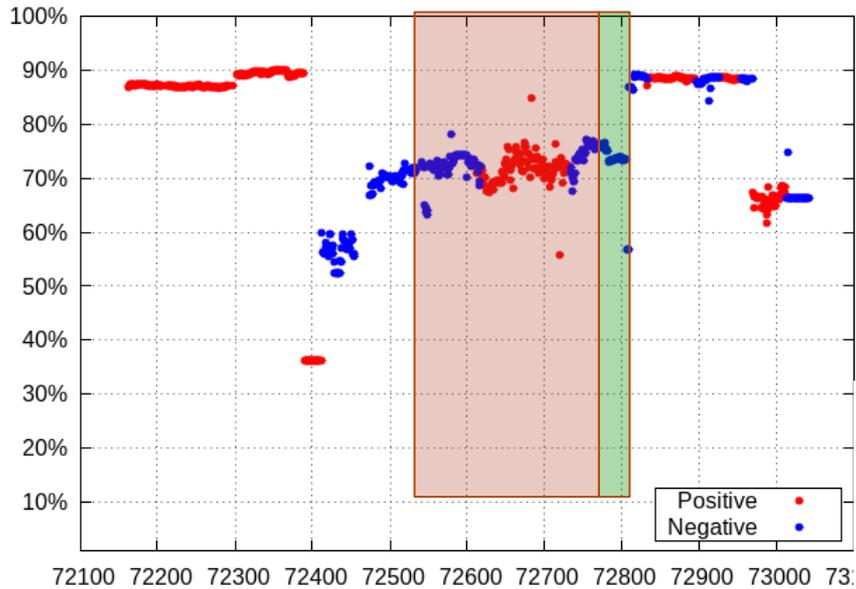
$$F = \frac{Yield_{Data} - Yield_{MC(N+He)}}{Yield_{Data}}$$

- We have taken data using both NH3 targets, called NH3 top and NH3 bottom.
- NH3 crystals are not uniformly filled in each targets which arise two different packing fractions and hence two different dilution factors.



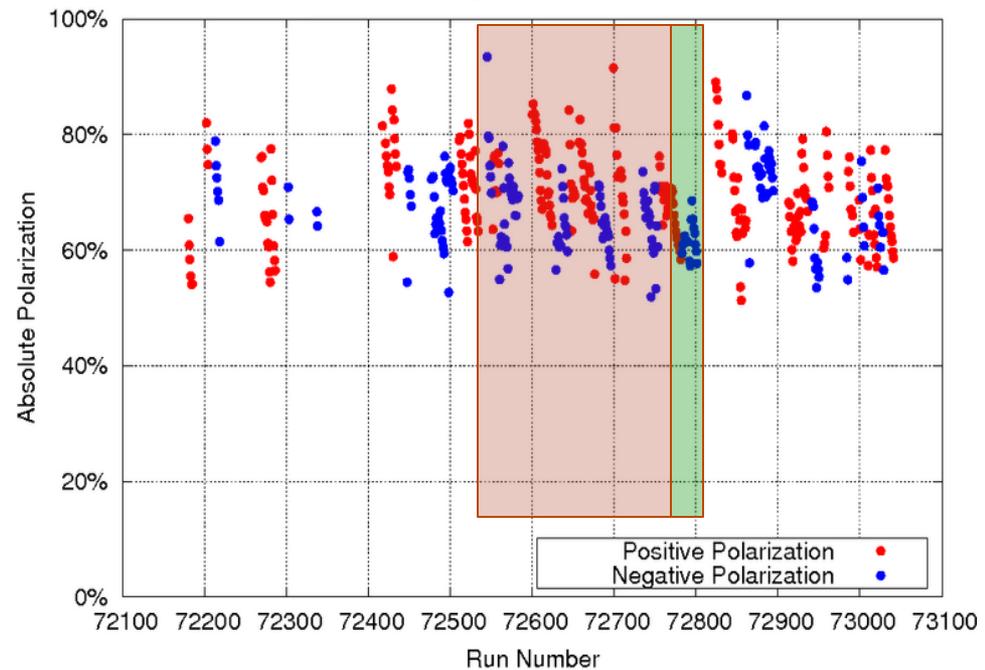
Beam / Target Polarizations

SANE Beam Polarization Per Run

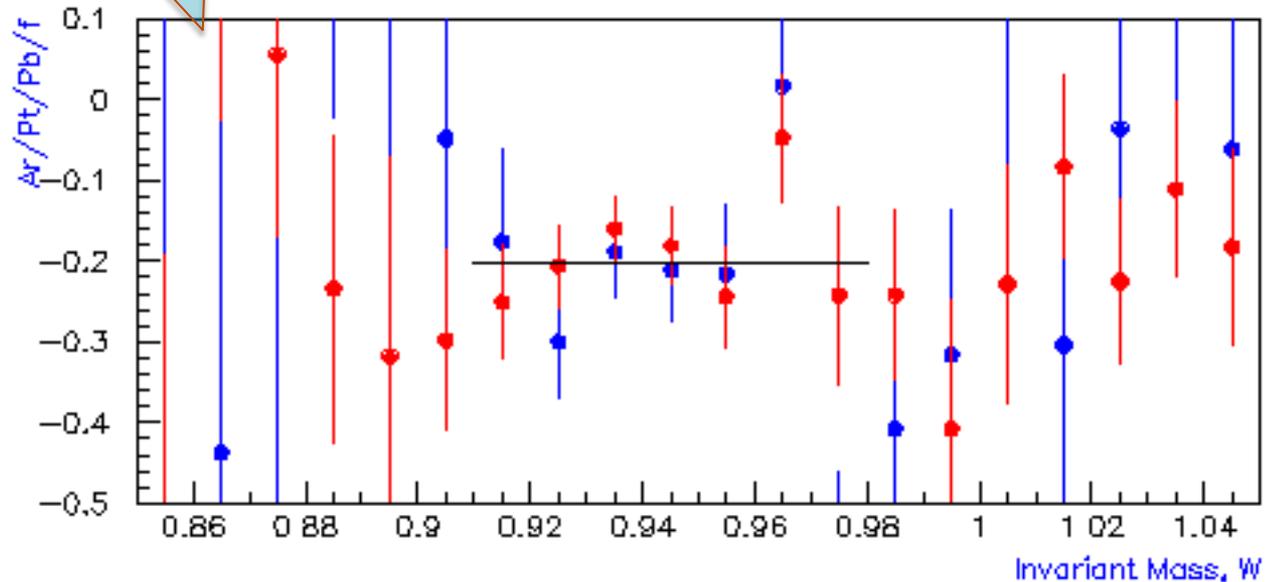
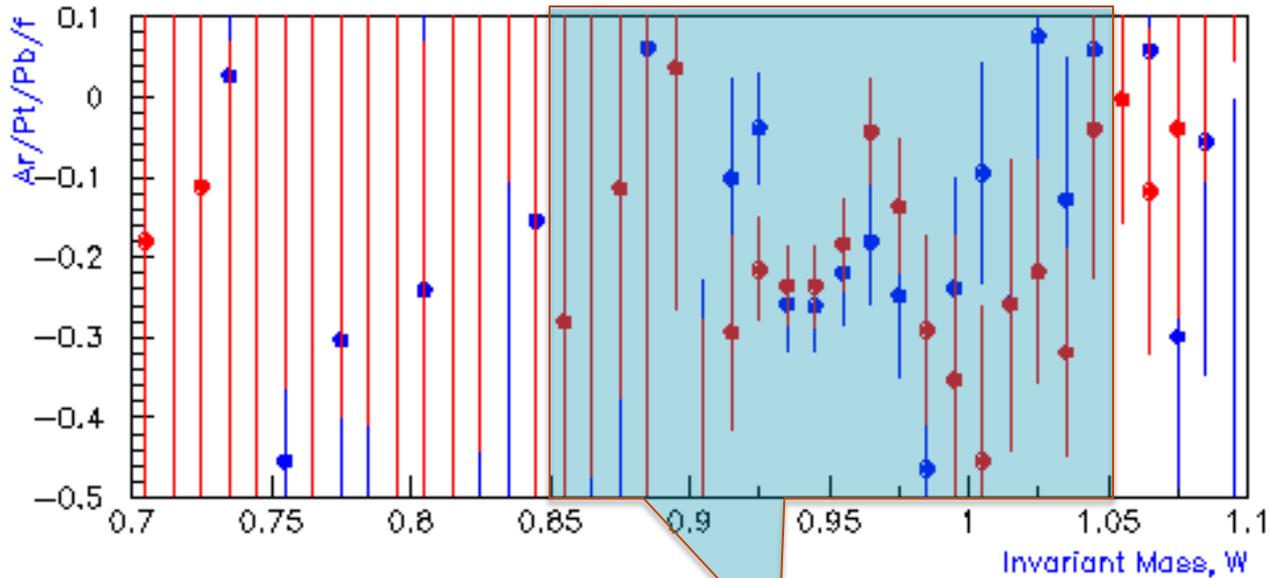


COIN data
Single arm electron data

Absolute Target Polarization for All SANE Runs



The Physics Asymmetry (Preliminary)



A_{phy}	Error A_{phy}
-0.201	0.0174

- The beam - target asymmetry, A_p

$$A_p = \frac{-br \sin \theta^* \cos \phi^* - a \cos \theta^*}{r^2 + c}$$

$$\frac{G_E}{G_M} = -\frac{b}{2A_p} \sin \theta^* \cos \phi^* + \sqrt{\frac{b^2}{4A_p^2} \sin^2 \theta^* \cos^2 \phi^* - \frac{a}{A_p} \cos \theta^* - c}$$

Using the experiment data at

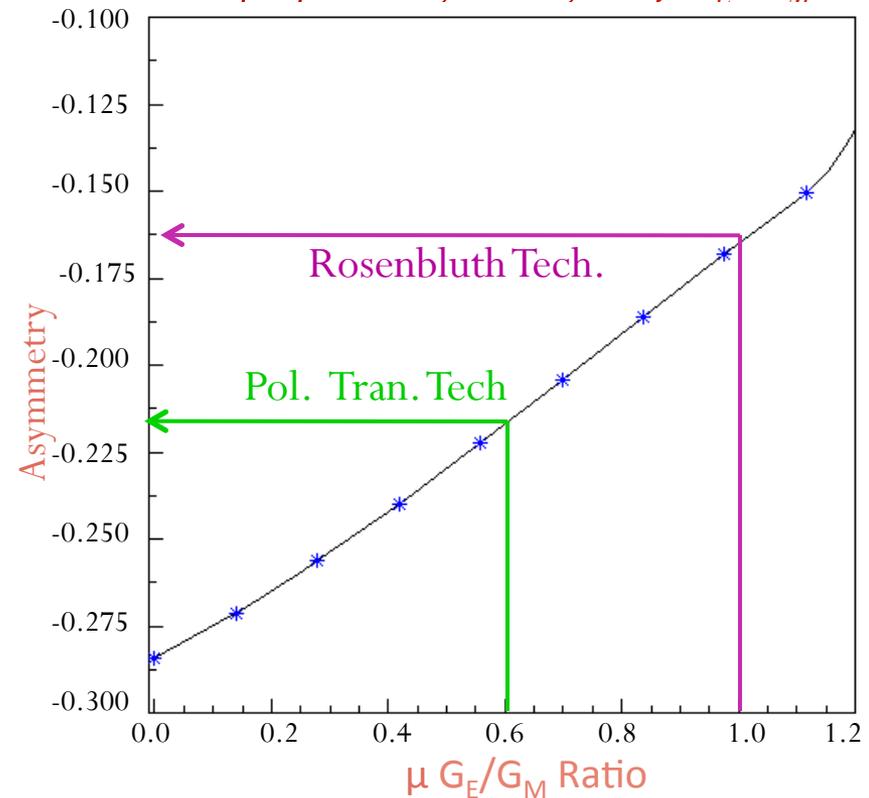
$$Q^2 = 2.2 \text{ (GeV/c)}^2$$

$$\theta^* \approx 34.55^\circ \text{ and } \phi^* = 180^\circ$$

From the HMS kinematics, $r^2 \ll c$

$$A_p = \frac{-b \sin \theta^* \cos \phi^* r}{c} - \frac{a \cos \theta^*}{c}$$

The projected asymmetry vs $\mu G_E/G_M$



Using the experiment data at $Q^2=2.2 \text{ (GeV/c)}^2$ and by knowing the $A_p=-0.201$,

$$r = \left(\frac{G_E}{G_M} \right) = 0.2416$$

$$\mu r = \mu \left(\frac{G_E}{G_M} \right) = 0.674$$

Where, μ – Magnetic Moment of the Proton=2.79

■ Error propagation from the experiment

$$A_p = \frac{-b \sin \theta^* \cos \phi^* r}{c} - \frac{a \cos \theta^*}{c}$$

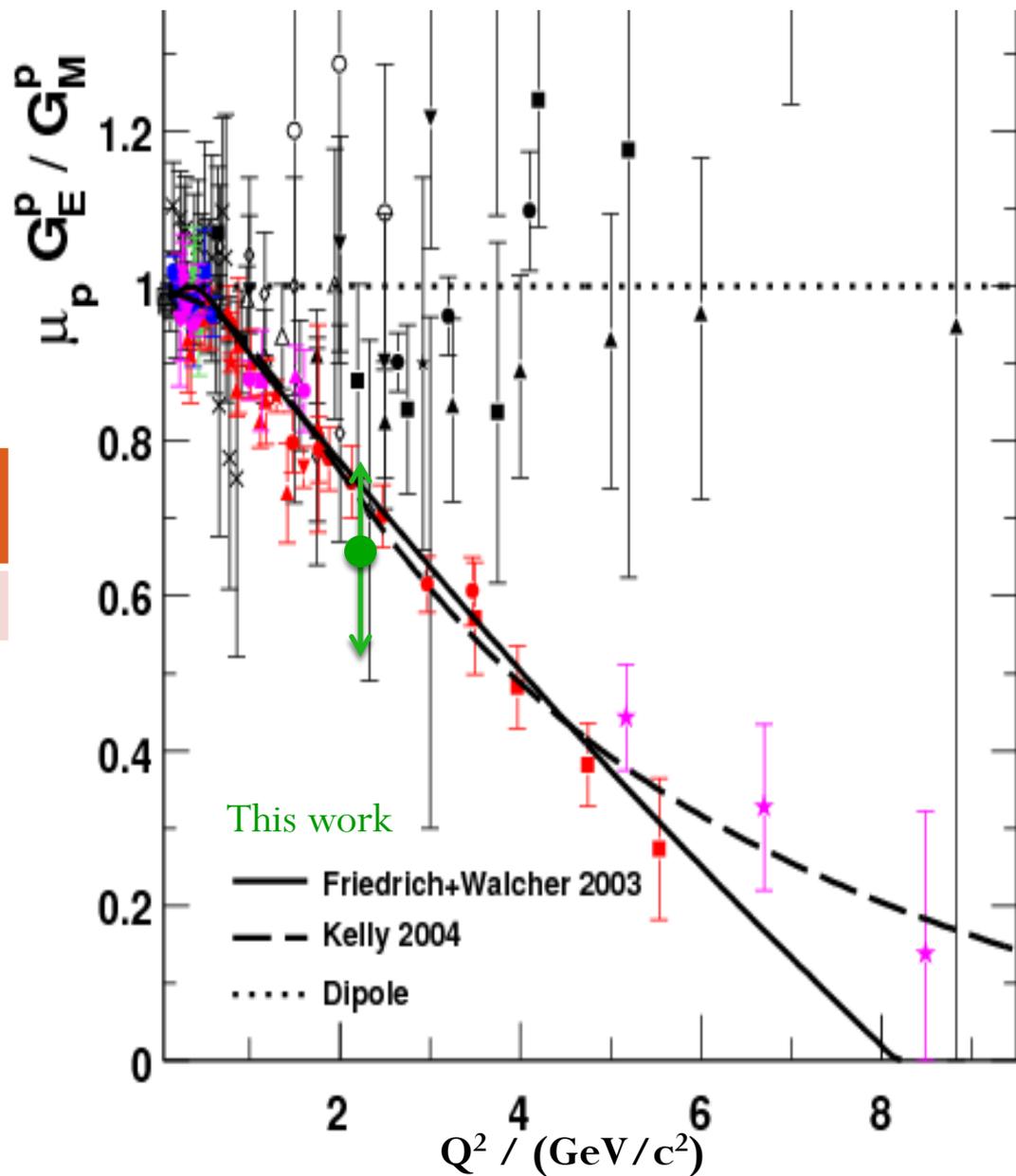
$$\Delta r = \Delta \left(\frac{G_E}{G_M} \right) = \left| \frac{c}{b \sin \theta^* \cos \phi^*} \right| \Delta A_p$$

By knowing the $\Delta A_p=0.017$,

$$\Delta(\mu r) = \Delta \left(\mu \frac{G_E}{G_M} \right) = 0.13$$

Preliminary

$\mu G_E/G_M$	$\Delta(\mu G_E/G_M)$
0.674	0.13



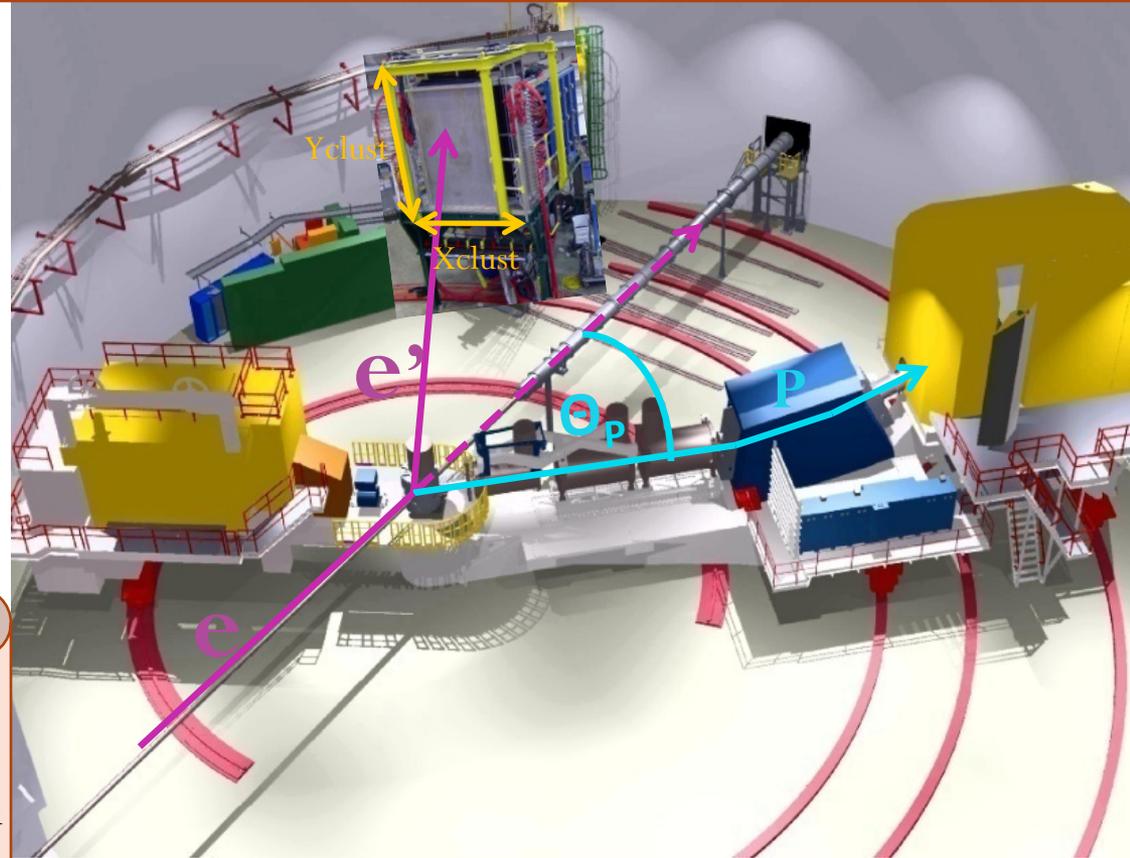
Coincidence Data (Electrons in BETA and Protons in HMS)

Definitions :

X/Y_{clust} - Measured X/Y positions on the BigCal

- X = horizontal / in-plane coordinate
- Y = vertical / out-of-plane coordinate

E_{clust} - Measured electron energy at the BigCal



By knowing
the energy of the polarized electron
beam, E_B
and
the scattered proton angle, Θ_p



We can predict the

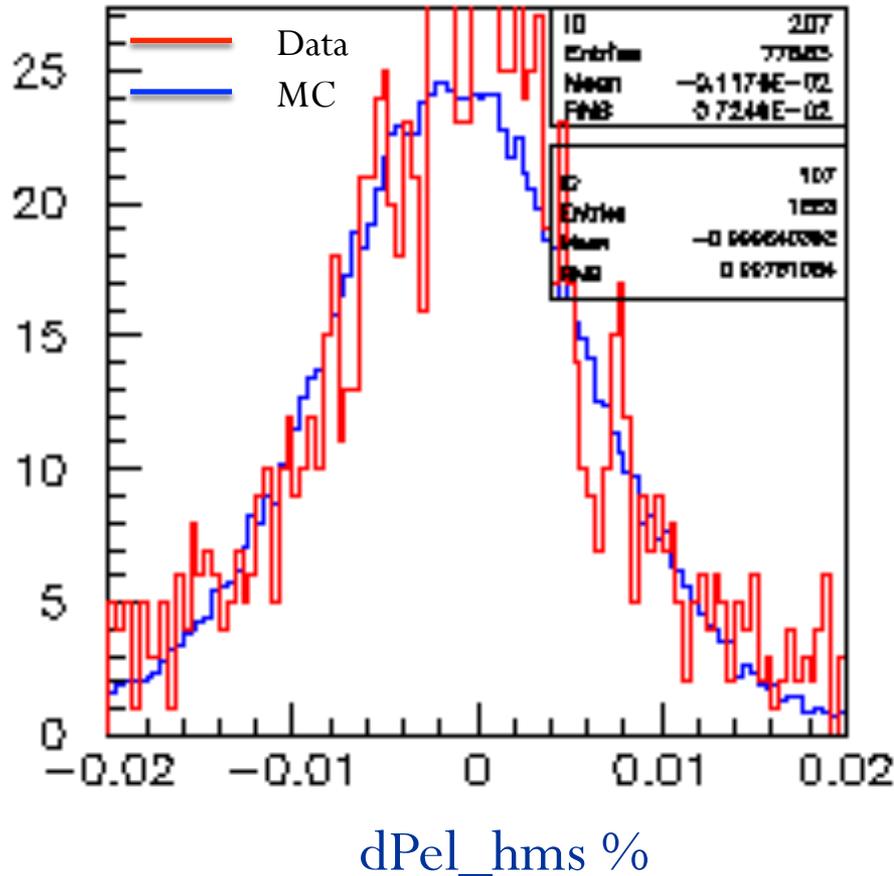
- X/Y coordinates - X_{HMS} , Y_{HMS} and (Target Magnetic Field Corrected)
- The Energy - E_{HMS} of the coincidence electron on the BigCal

Elastic Kinematics

(From HMS Spectrometer)

Spectrometer mode	Coincidence	Coincidence	Single Arm
HMS Detects	Proton	Proton	Electron
E Beam GeV	4.72	5.89	5.89
P_{HMS} GeV/c	3.58	4.17	4.40
Θ_{HMS} (Deg)	22.30	22.00	15.40
Q^2 (GeV/c) ²	5.17	6.26	2.20
Total Hours (h)	~40 (~44 runs)	~155 (~135 runs)	~12 (~15 runs)
e-p Events	~113	~1200	-

Fractional momentum difference



$$dPel_hms = \frac{P_{HMS} - P_{Cal}}{P_{cent}}$$

$$P_{Cal} = \sqrt{v^2 + 2Mv}$$

$$v = \frac{Q^2}{2M}$$

$$Q^2 = \frac{4M^2 E^2 \cos^2 \theta}{M^2 + 2ME + E^2 \sin^2 \theta}$$

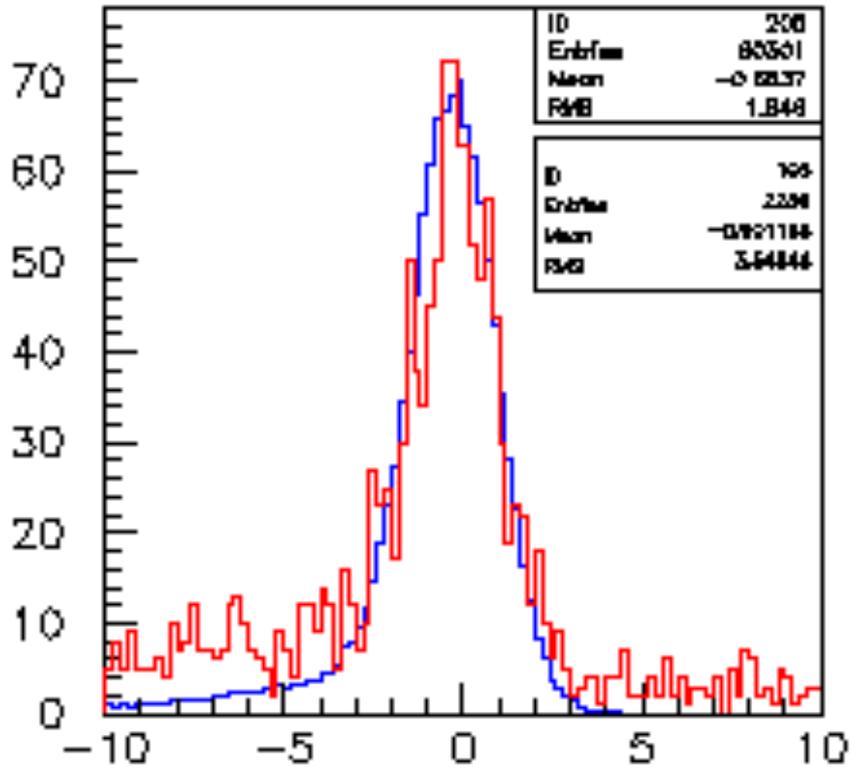
P_{HMS} – Measured proton momentum by HMS

P_{cal} - Calculated proton momentum by knowing the beam energy, E and the proton angle, θ

P_{cent} – HMS central momentum

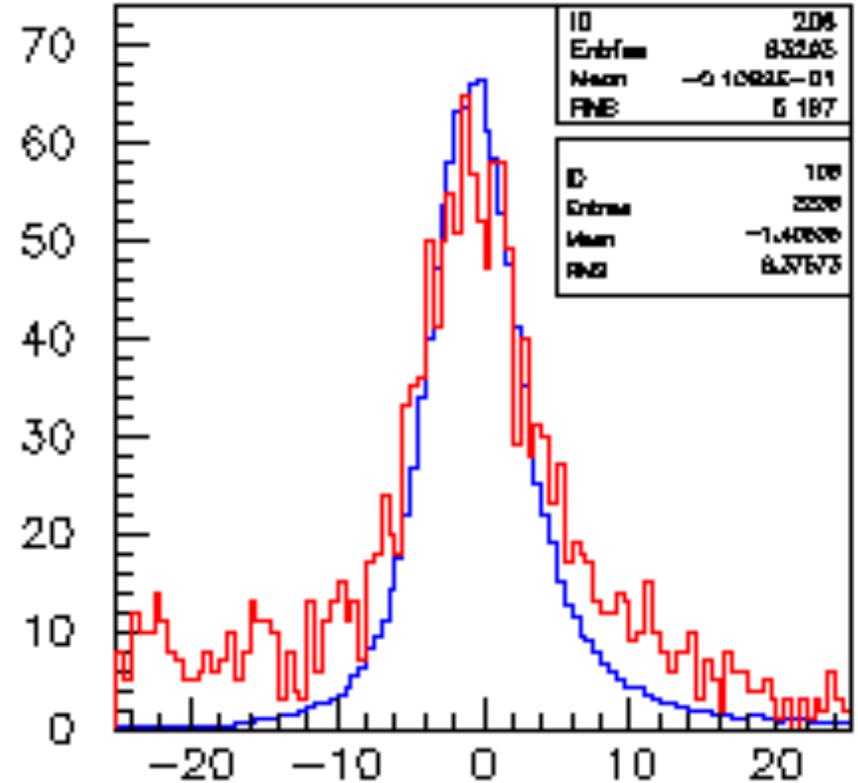
X/Y position difference

X position difference



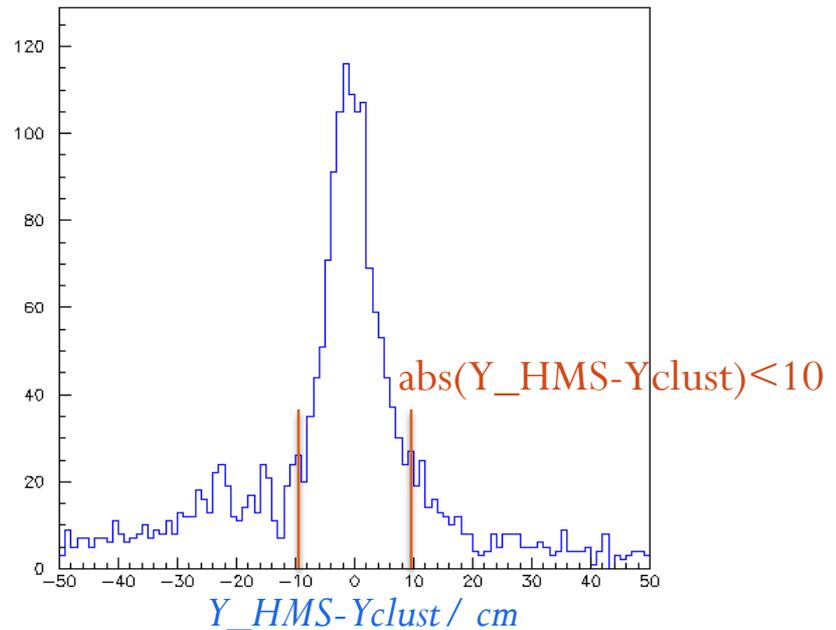
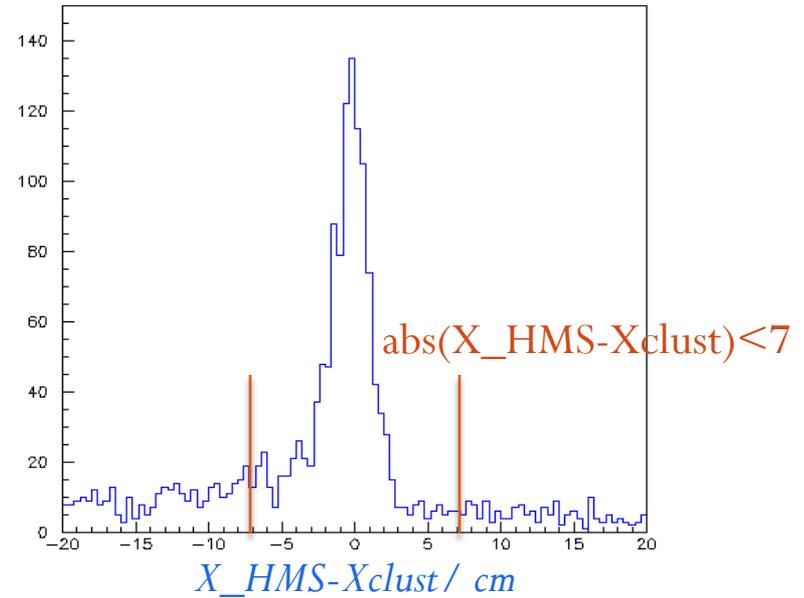
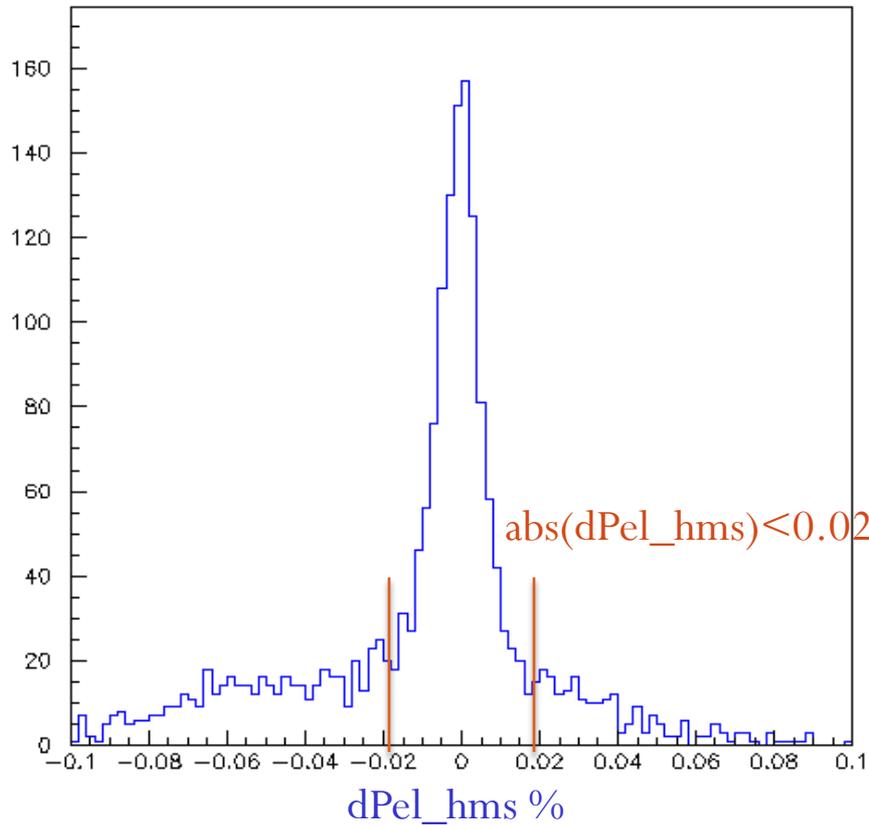
$X_{HMS} - X_{clust} / \text{cm}$

Y position difference

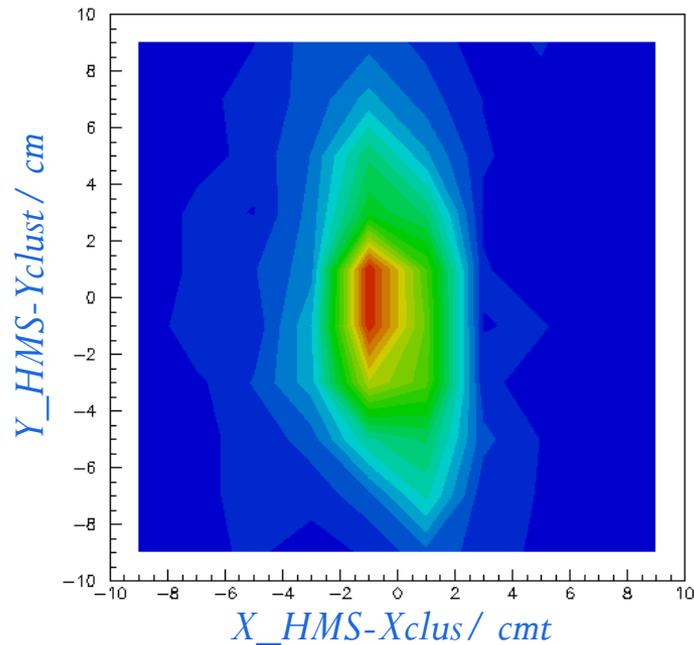


$Y_{HMS} - Y_{clust} / \text{cm}$

Applied the coincidence cuts

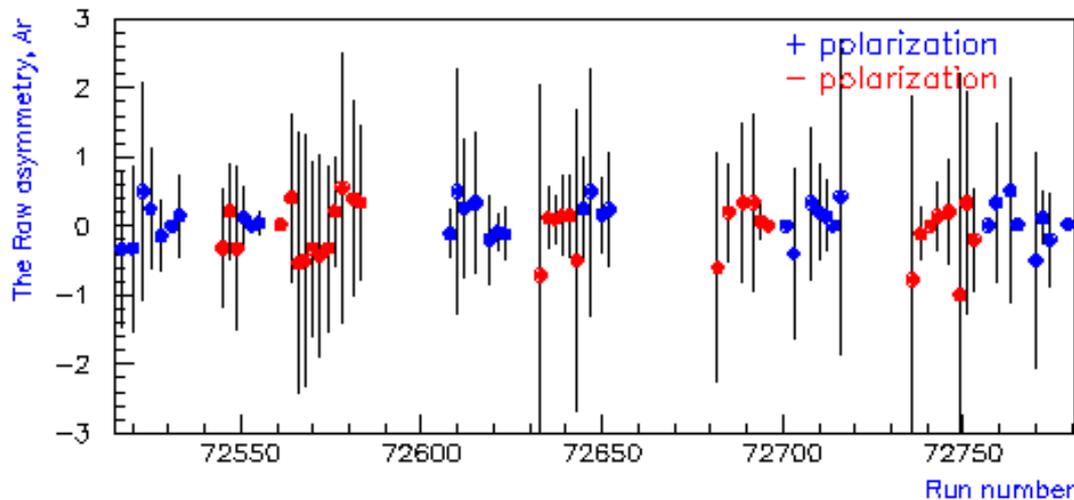


Extract the asymmetries



*Need
dilution factor, f
in order to determine the
physics asymmetry,*

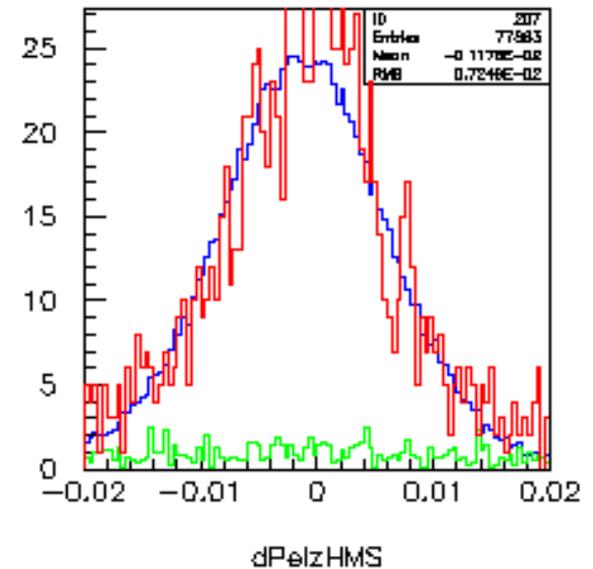
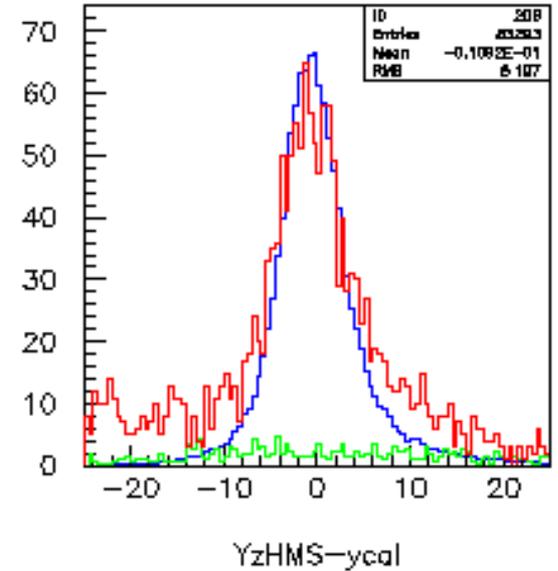
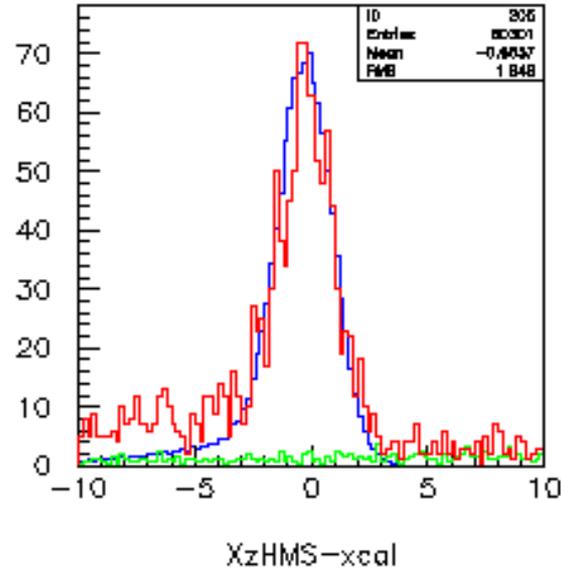
$$A_p = \frac{A_r}{fP_B P_T} + N_C$$



*and G_E^p / G_M^p
(at $Q^2 = 6.26 \text{ (GeV/c)}^2$)*

Estimate the background

- Data
- MC in H
- MC in C



- Use Carbon simulation.
- This is an effort to determine the dilution factor.
- Still working on it.

Conclusion

- Measurement of the beam-target asymmetry in elastic electron-proton scattering offers an independent technique of determining the G_E/G_M ratio.
- This is an 'explorative' measurement, as a by-product of the SANE experiment.
- Extraction of the G_E/G_M ratio from single-arm electron data are shown.
- The preliminary data point at 2.2 (GeV/c)^2 is very consistent with the recoil polarization data (falls even slightly below it)
- The preliminary data points from the coincidence data at $5.17, 6.26 \text{ (GeV/c)}^2$ will become available soon.

SANE Collaborators:

Argonne National Laboratory, Christopher Newport U., Florida International U., [Hampton U.](#), Thomas Jefferson National Accelerator Facility, Mississippi State U., North Carolina A&T State U., Norfolk S. U., Ohio U., Institute for High Energy Physics, U. of Regina, Rensselaer Polytechnic I., Rutgers U., Seoul National U., State University at New Orleans , Temple U., Tohoku U., U. of New Hampshire, U. of Virginia, College of William and Mary, Xavier University of Louisiana, Yerevan Physics Inst.

Spokespersons: S. Choi (Seoul), M. Jones (TJNAF), Z-E. Meziani (Temple),
O. A. Rondon (UVA)

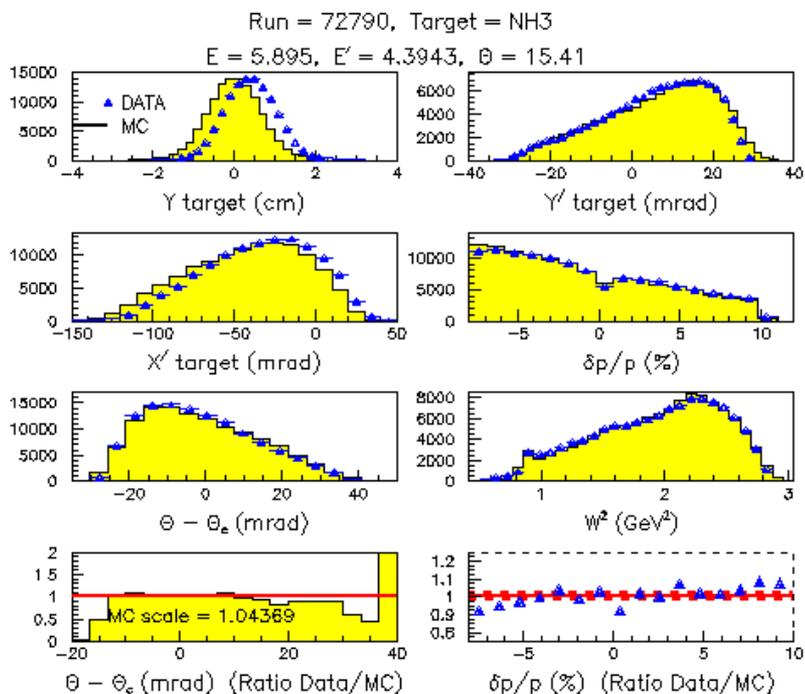
Thank You



Backup Slides

MC with NH3

- Generated N, H and He separately.
- Added Al come from end caps and 4K shields as well.
- Calculated the MC scale factor using the data/MC luminosity ratio for each target type.
- Added all targets together by weighting the above MC scale factors.
- Used 60% packing fraction.



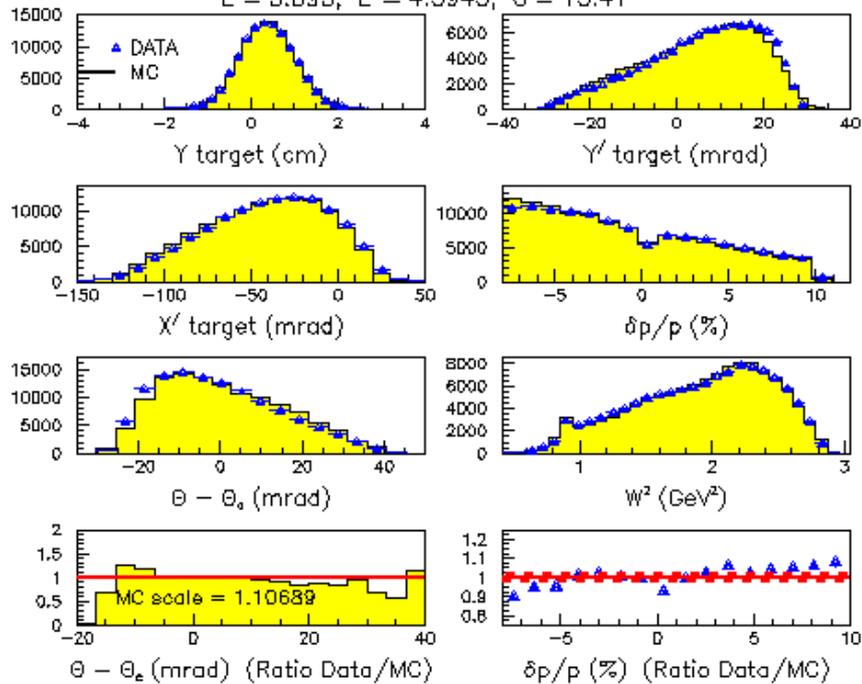
Srast x – Horizontal beam position
(pointing beam left)

Srast y – Vertical beam position
(pointing up)

- Adjust acceptance edges in
Ytar and y' from adjusting srast x offset

Run = 72790, Target = NH3

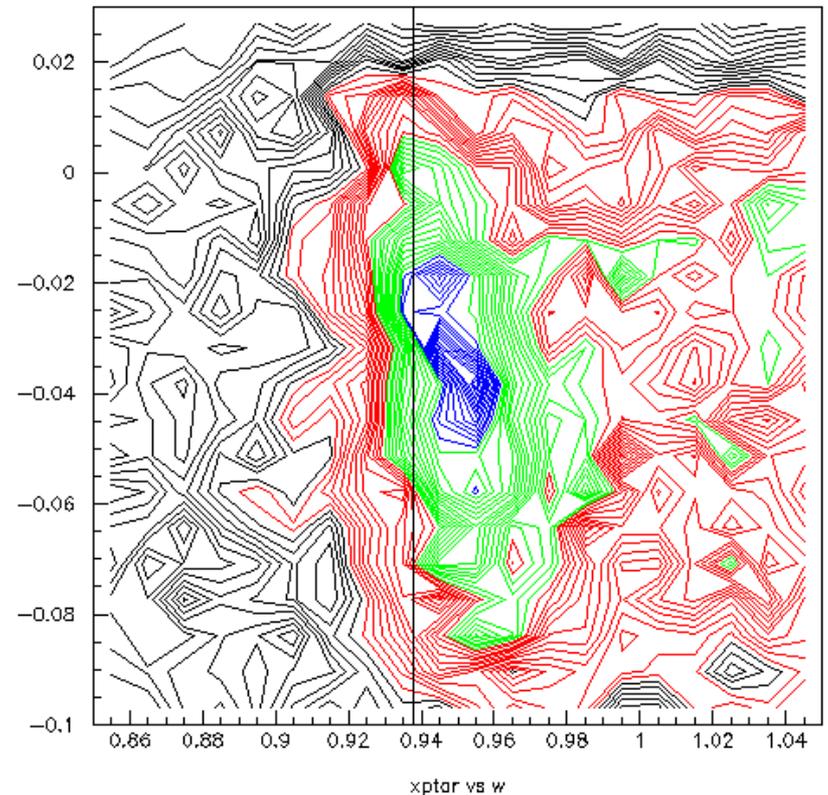
$E = 5.895$, $E' = 4.3943$, $\theta = 15.41$



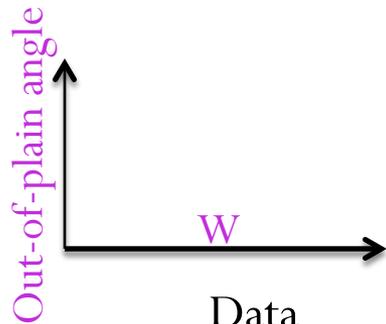
Srast x offset = -0.4 cm

Srast x offset = 0.0 cm

But shows the xptar vs w correlation.

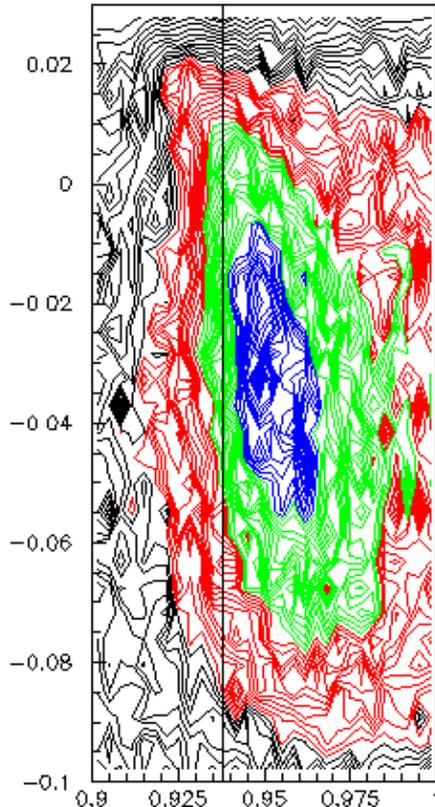


Check the srast x offsets with MC



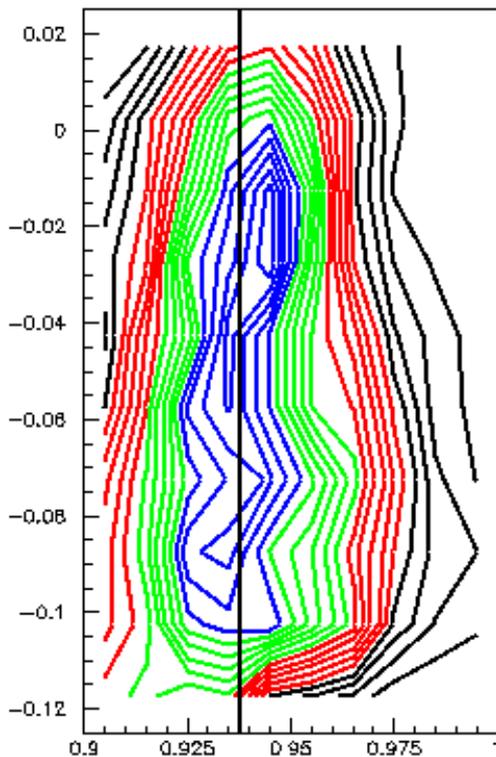
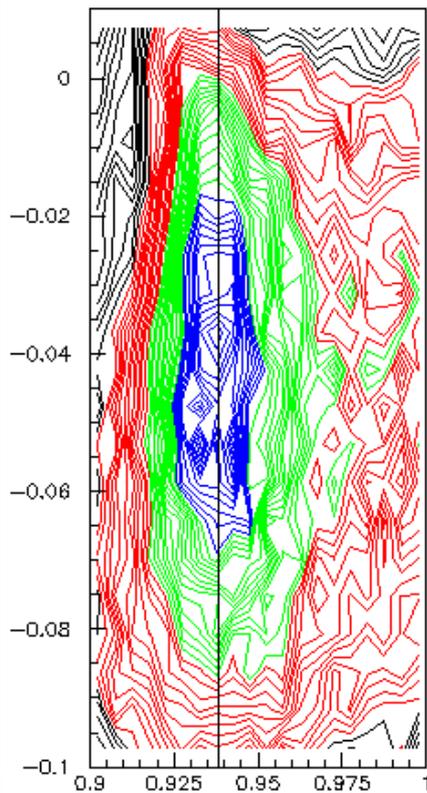
Data

Srast x=0.2 cm



Data

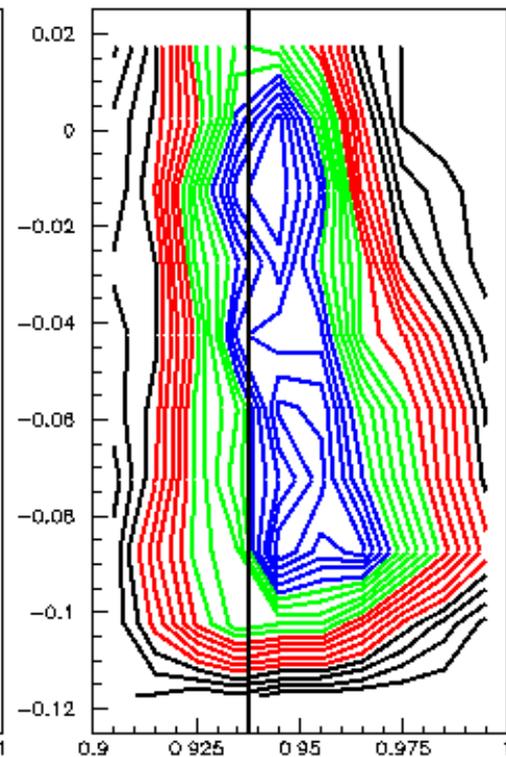
Srast x=-0.9 cm



MC

Gen. Srast x=-0.2 cm

Rec. Srast x=-0.2cm



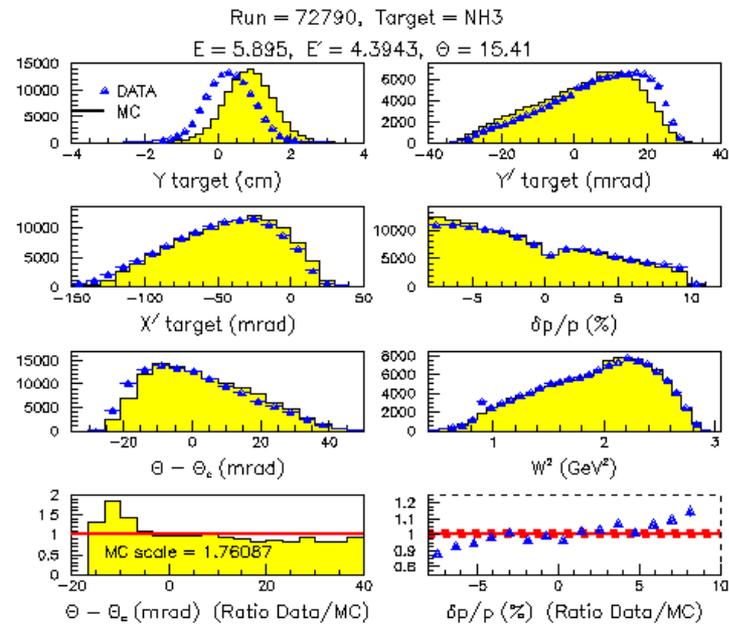
MC

Gen. Srast x=-0.9 cm

Rec. Srast x=0.2cm

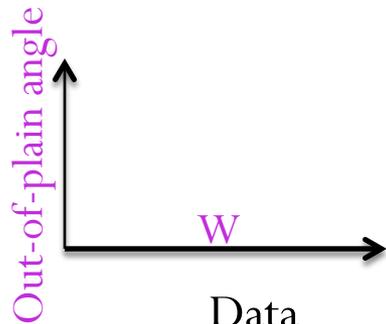
Srast x offset=-0.9 cm ?Too BIG

It does not match with the MC either.



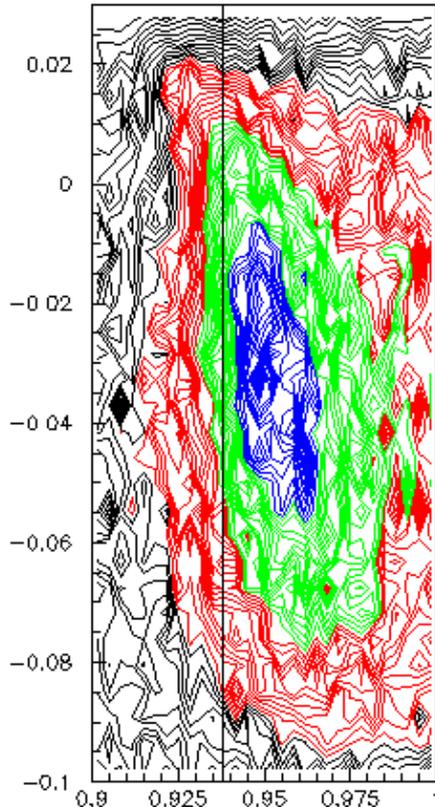
Therefore, this x_{ptar} vs w correlation can be a combination of srast_x and something else.....

Check the srast x offsets with MC



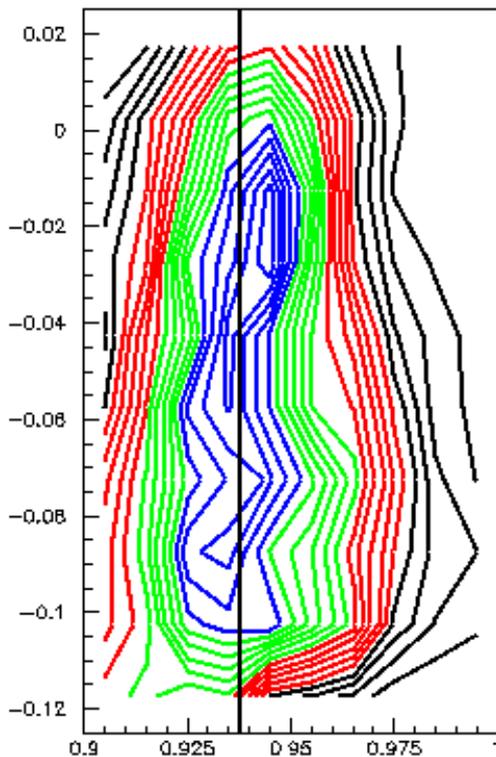
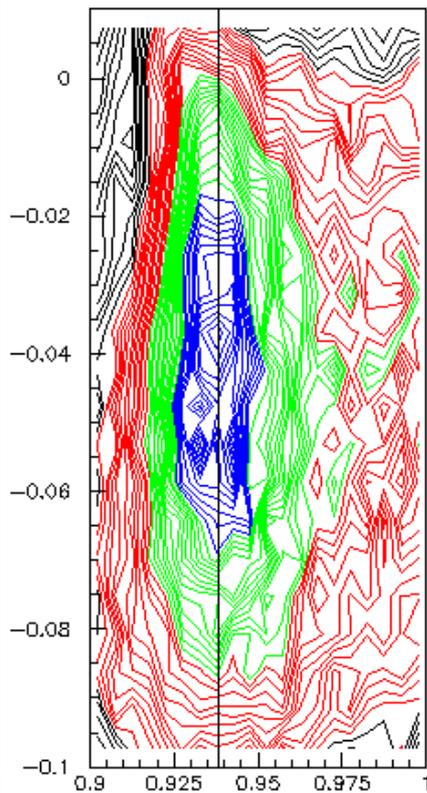
Data

Srast x=0.2 cm



Data

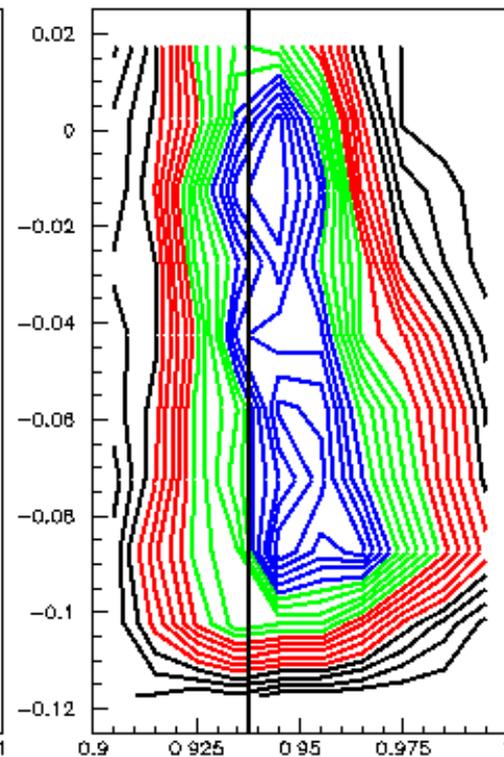
Srast x=-0.9 cm



MC

Gen. Srast x=-0.2 cm

Rec. Srast x=-0.2cm



MC

Gen. Srast x=-0.9 cm

Rec. Srast x=0.2cm

Introduced an azim. Angle correction

We assume that the target magnetic field is symmetric around the target.

In practically, It might not.

So, Introduced the Out-Of-Plane angle (azimuthal angle) dependence field correction.

$$B_corr = (azim-az0)*az_corr$$

Applied the field correction on MC only for the forward direction and changed the parameters “az0” and “az_corr” to make the same xptar vs w correlation as seen on the data.

$$B_scale = 5.003/(5.003+abs(B_corr))$$

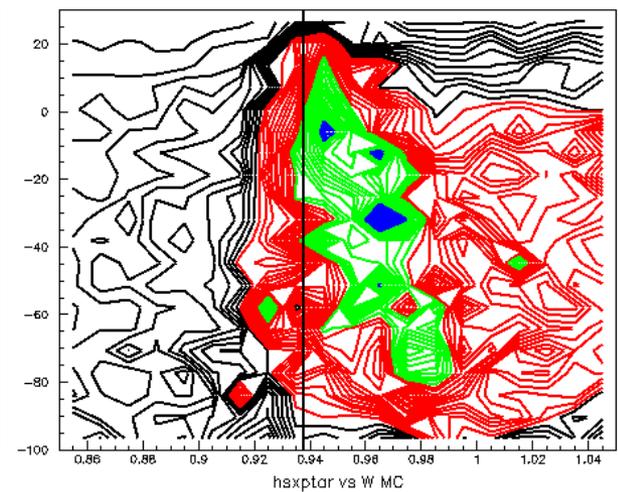
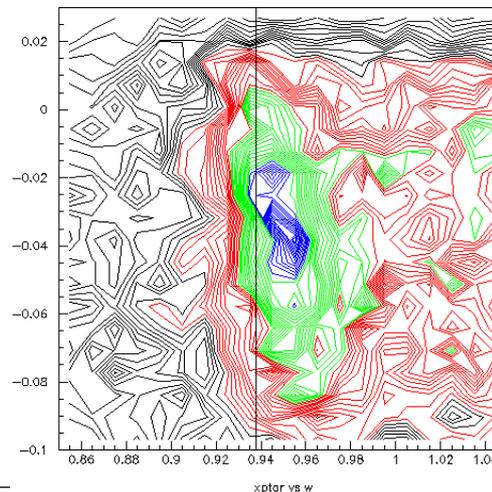
$B(3) = B(3)+B_corr$!B(3) is along Z direction of the field

$$B(3) = B(3)*B_scale$$

$B(1) = B(1)-B_corr$!B(1) is x component pointing down

$$B(1) = B(1)*B_scale$$

Then use that correction on data.



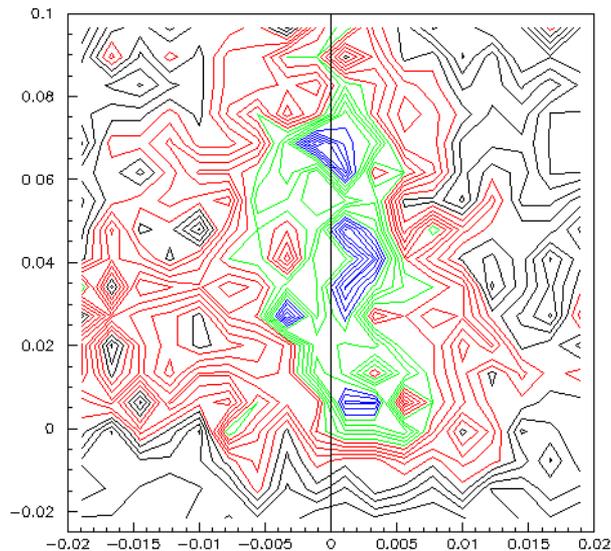
Determine the azimuthal angle correction

Because the azimuthal angle correction depend on the horizontal angle as well,
We need to find the different correction parameters for the HMS and BETA.

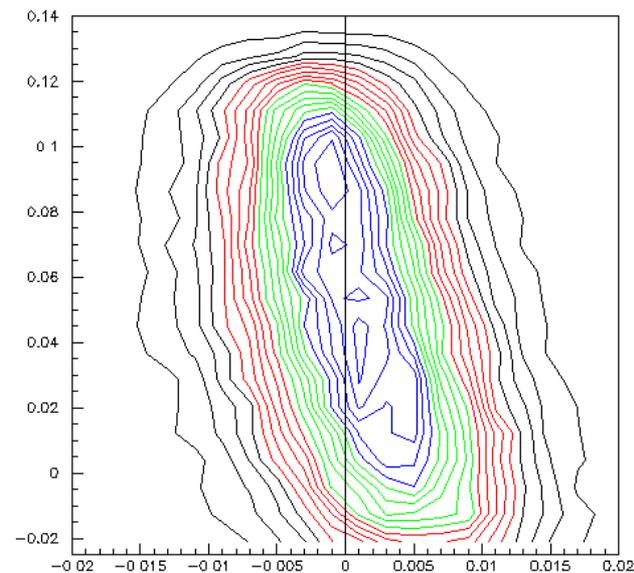
For HMS

By looking at the x_{ptar} vs $d_{\text{pel_hms}}$ correlation on the data, make the same correlation on SIMC by using the azimuthal angle correction only for the forward direction and changing the “azo” and “az_corr” parameters.

Data



SIMC

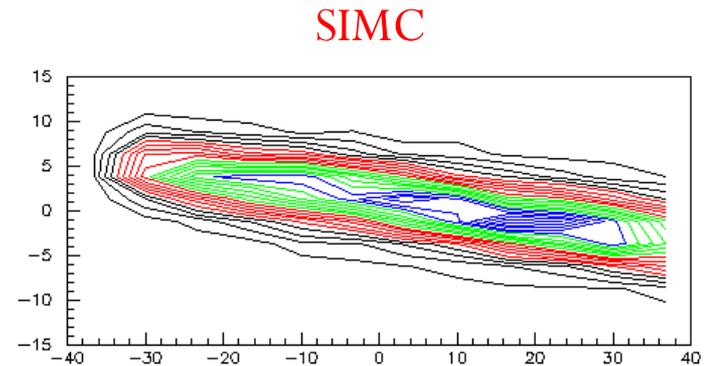
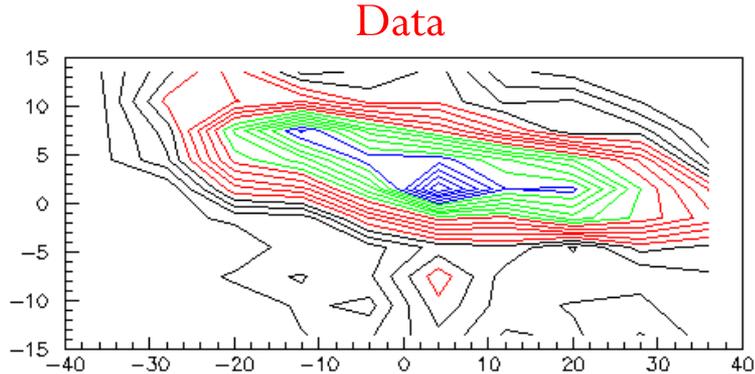


Apply the correction for the HMS side for the data

For BETA

Looked at the y_{diff} vs y_{clust} correlation on the data. Use the azimuthal angle correction on HMS side.

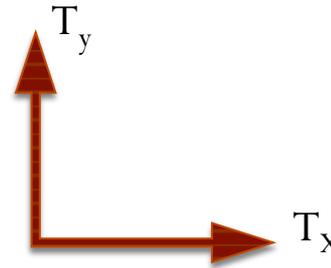
Make the same correlation on SIMC by changing the “az0” and “az_corr” parameters for BETA side.



Apply the correction for the HMS side for the data

Calibrate the EPICS T_x and T_y

- EPICS T_x – BPM horizontal (cm)
- EPICS T_y - BPM vertical (cm)



$$T_x + offset_x = srast_x$$

$$T_y + offset_y = srast_y$$

$$-2.361 + offset_x = -4.00(mm)$$

$$+0.580 + offset_y = +1.00(mm)$$

$$offset_x = -1.64(mm)$$

$$offset_y = +0.42(mm)$$

- Used the above EPICS calibration constants, $offset_{x/y}$ to determine the real beam positions, $srast_{x/y}$ for the other runs for the known T_x and T_y .

$$T_x - 1.64 = srast_x$$

$$srast_x = ?$$

$$T_y + 0.42 = srast_y$$

$$srast_y = ?$$



NOTE:

EPICS calibration constants, $offset_{x/y}$ change with the magnetic field. The girder that holds the BPM is moved when going from perpendicular to parallel.

- In the perpendicular configuration, the electron beam is deflected by different amounts at different beam energies by the target magnetic field and so the girder position has changed. So, it gives different EPICS Tx and Ty at different beam energies and hence we should have different EPICS calibration constants.
- In the parallel configuration, the beam is not deflected by the target field and so the girder position remains the same at different beam energies. Therefore, the EPICS Tx and Ty and hence the EPICS calibration constants are the same at different beam energies.

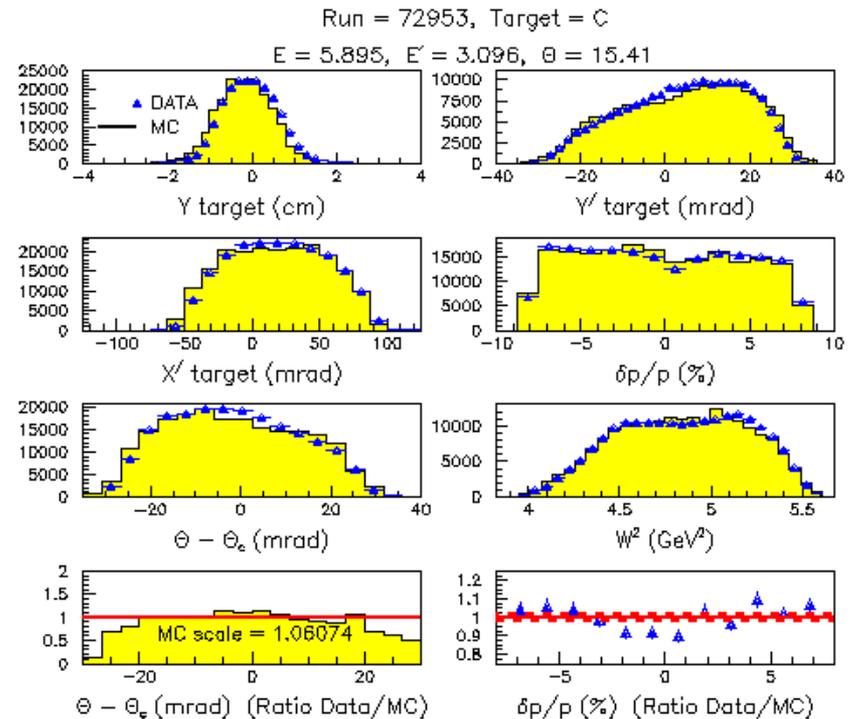
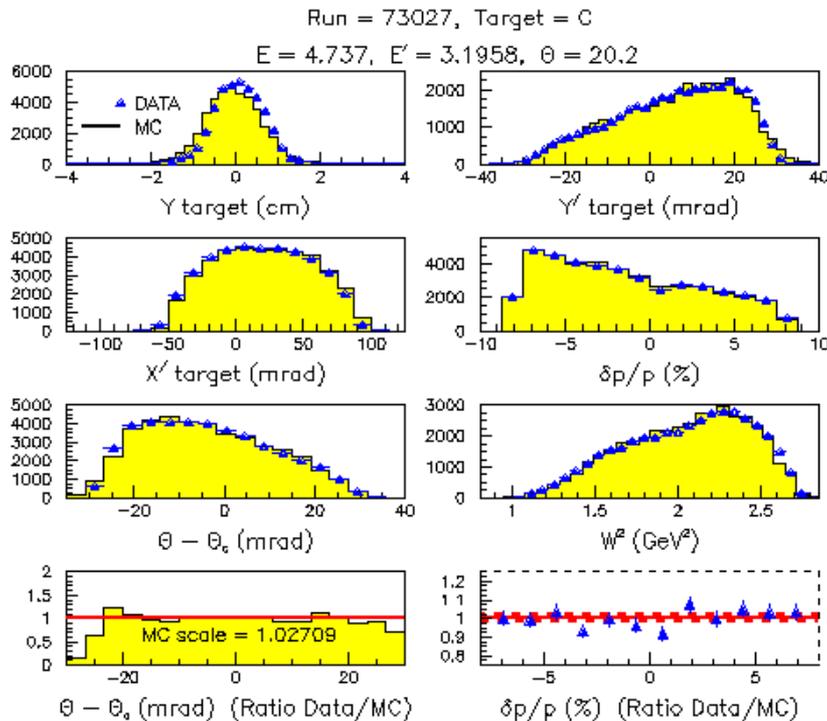
Parallel field Magnetic Configuration

C run 73027 (No He)
 (ebeam = 4.733 GeV, P=3.2 GeV/C,
 $\theta = 20.2^\circ$)

C run 72953 (ebeam = 5.895 GeV,
 P=3.1 GeV/C, $\theta = 15.41^\circ$)

$srast_x = 0.10$ cm
 $srast_y = -0.10$ cm

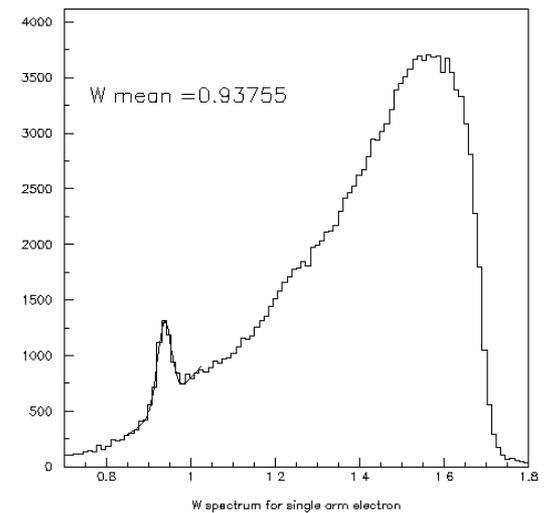
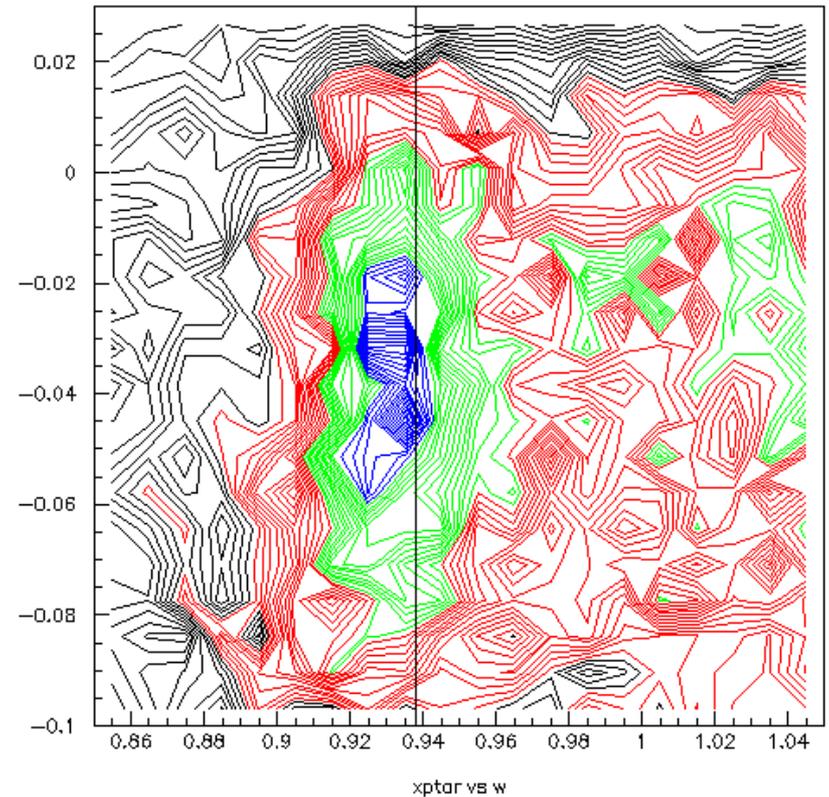
$srast_x = 0.18$ cm
 $srast_y = -0.02$ cm



The azimuthal angle correction fix the x_{ptar} vs w correlation on the data.

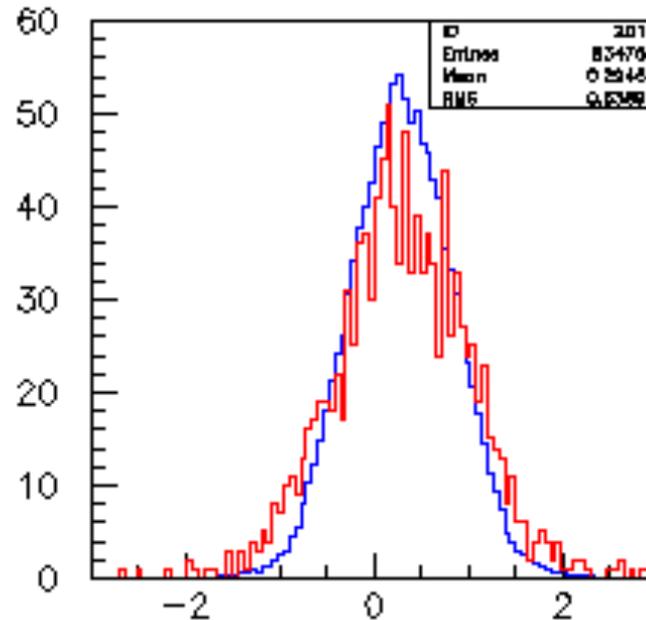
But, w peak is shifted. (~ 0.93 GeV)
Changed $srast_y$ offset to bring the W peak to 0.938 GeV

$srast_y$ offset = 0.1 cm



Find the beam offsets

Changed the `srast_x` offsets to match ‘hsytar’ of data to SIMC each other.



Changed the `srast_y` offsets to match ‘dpel_hms’ of data to SIMC each other.

$$srast_x = -0.25 \text{ cm}$$

$$srast_y = -0.20 \text{ cm}$$

Beam Time

	Energy GeV	Θ_N	Time (Proposal FOM h)		
			Proposal	Actual	Fraction
Calibration	2.4	off, 0, 180	47	25	53%
Production	4.7	180	70	20	29%
	4.7	80	130	98	75%
	5.9	80	200	143	72%
	5.9	180	100	≥ 35	$\geq 35\%$
<hr/>					
Commissioning [calendar days]			14.0	99	
Total [calendar days]			70.0	141	