

# Electromagnetic Calorimeters

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JLab Summer Detector/Computer Lectures

`http:`

`//www.jlab.org/~gen/talks/calor_lect_3.pdf`

# Outline

- 1 Introduction
- 2 Calorimeters
  - Generic calorimeter
  - Light collecting calorimeters
- 3 Front-End Electronics
- 4 Procedures
- 5 Summary

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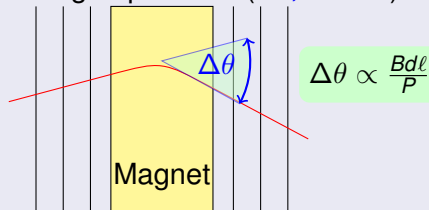
# What is a calorimeter?

Particle detection main goal: measure 3-momenta  $\vec{P}$

## Magnetic spectrometers

- Coordinate detectors
- Magnetic field

Charged particles ( $e^\pm$ ,  $\pi^\pm$  etc)



Momentum resolution:

$$\sigma(P)/P \propto P \quad (\text{for large } P)$$

## Calorimeters

Detectors thick enough to absorb nearly all of the particle's energy released via cascades (showers)

Neutral ( $\gamma$ ,  $n$ ) and charged particles

The energy goes mainly into heat.

- "True" C. -  $E_0$  (heat)
- "Pseudo" C. -  $\mathcal{O}(E_0)$ :  
ionization, Cherenkov light

Poisson process:  $N_e \propto E_0$ ,

$$\sigma N_e = \sqrt{N_e} \text{ and } \frac{\sigma E}{E} \propto \frac{1}{\sqrt{E}}$$

# "True" Calorimeters

"True" calorimeters measure the temperature change of the absorber:  $\Delta T = \frac{E_0}{c \cdot M} \sim \frac{1 \cdot 10^{10} \text{ eV} \cdot 1.6 \cdot 10^{-19} \text{ J/eV}}{10^3 \text{ J/kg} \cdot 1 \text{ kg}} \approx 10^{-12} \text{ K}$  too low!

- High particle flux
  - History: W. Orthmann -  $1 \mu\text{W}$  sensitivity; 1930, with L. Meitner they measured the mean energy of  $\beta$  from  $^{210}\text{Bi}$  (6% accuracy)  $\Rightarrow$  W.Pauli's neutrino hypothesis.
  - Precise beam current measurements (SLAC-1970s, JLab-2003)
- Ultra-cold temperatures (low C), superconductivity - new detectors for exotic particle search, like "dark matter" candidates.



# “Pseudo” Calorimeters

“Pseudo” calorimeters detect  $\mathcal{O}(E_0)$ : ionization, Cherenkov light

- History: N.L. Grigorov 1954 - idea, 1957 - implementation in cosmic ray studies (Pamir, 3900 m). Layers of an absorber and layers of proportional counters - counting the number of particles in the shower (calibration needed).
- Starting in 1960s - revolution in compact electronics  $\Rightarrow$  affordable ADC (Analog-to-Digital Converters). New accelerators - various types of calorimeters with  $\sim 10 \rightarrow 10^5$  ADC channels.

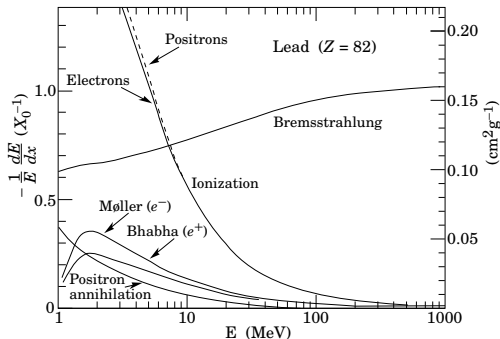
## Applications

- detecting neutrals
- good energy resolution at high energies
- fast signals for trigger
- particle identification ( $e^\pm/h$ )

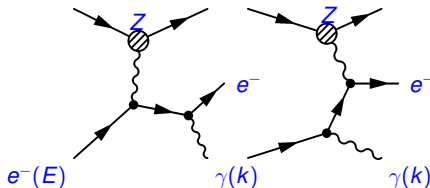
# $e^\pm$ interactions

## Energy loss in medium

- Bremsstrahlung  
 $e^\pm Z \rightarrow e^\pm \gamma Z$
- Ionization
- Bhabha/Møller scattering  
 $e^\pm e^- \rightarrow e^\pm e^-$
- $e^+$  annihilation



## Bremsstrahlung



$$\sigma \propto \frac{Z^2}{m^2} \Rightarrow \frac{\sigma_\mu}{\sigma_e} \approx 2 \cdot 10^{-5}$$

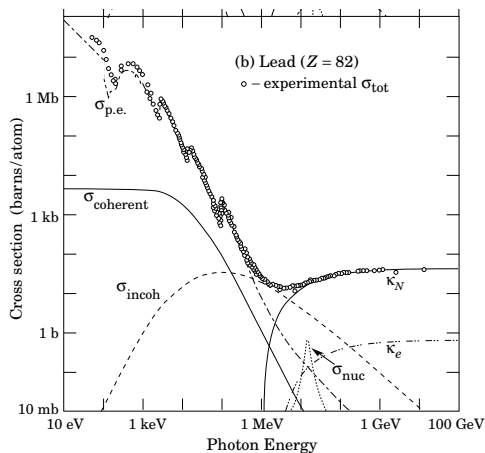
$$\frac{dN_\gamma}{dk} \propto \frac{1}{k}$$

$$\frac{dE_\gamma}{dk} = c(k)$$

# $\gamma$ interactions

## Interaction in medium

- Pair production  
 $\gamma Z \rightarrow e^+ e^- Z$  ( $K_N$ )
- Pair production  
 $\gamma e^- \rightarrow e^+ e^- e^-$  ( $K_e$ )
- Compton scattering  
 $\gamma e^- \rightarrow \gamma e^-$  ( $\sigma_{incoherent}$ )
- Rayleigh scattering  
( $\sigma_{coherent}$ )
- Photonuclear absorption  
( $\sigma_{nuc}$ )
- Atomic photoeffect ( $\sigma_{p.e.}$ )



# Scaling of Material Properties

## Radiation length

$X_0$  - the material thickness for a certain rate of EM:

$$e^\pm: \frac{dE_{loss}}{dx} \simeq \frac{E}{X_0}$$

$$\gamma: \lambda_{e^+e^-} \simeq \frac{9}{7} \cdot X_0$$

Derived from EM calculations:

$$X_0 \simeq \frac{716 \text{ g}\cdot\text{cm}^{-2}\cdot\text{A}}{Z(Z+1)\cdot\ln(287/\sqrt{Z})}$$

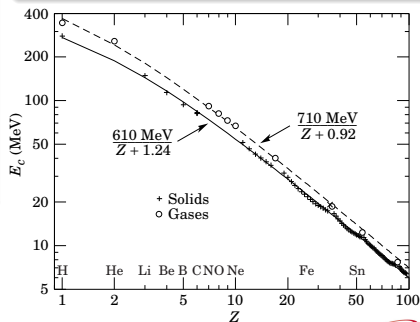
## Critical Energy

$E_c$ : cascade stops

Losses: Ionization = Radiation

$$\text{B.Rossi: } \left. \frac{dE_{ioniz}}{dx} \right|_{E_c} \simeq \frac{E}{X_0}$$

$$E_c \simeq \frac{610(710) \text{ MeV}}{Z+1.24(0.92)} \text{ solids(gasses)}$$



# Electromagnetic Showers

Photons and light charged particles ( $e^\pm$ ) interact with matter:

- electrons radiate  $e^\pm \rightarrow e^\pm \gamma$
- photons convert  $\gamma \rightarrow e^+ e^-$

A cascade develops till the energy of the particles go below a certain limit.

The charged particles of the cascade ( $e^\pm$ ) leave detectable signals.

# Electromagnetic Shower: longitudinal development

Scaling variables:

$$t = \frac{x}{X_0} \quad y = \frac{E}{E_c}$$

## Simple model

A simple example of a cascade:

$\times 2$  at  $\Delta t = 1$ .

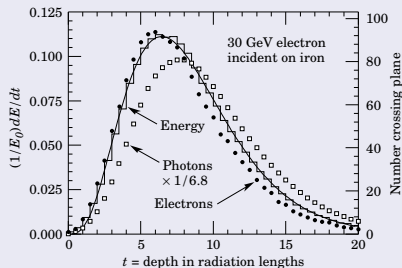
$$E(t) = \frac{E_0}{2^t} \Rightarrow t_{max} = \ln \frac{E_0}{E_c} / \ln 2$$

$$t_{max} \propto \ln \left( \frac{E_0}{E_c} \right)$$

Detectable signal:

$$L_{charged} \propto E_0 / E_c$$

## Simulation: EGS4, GEANT



$$t_{max} \simeq \ln(y) + \begin{cases} -0.5 & e^- \\ +0.5 & \gamma \end{cases}$$

$$t(> 95\%) \simeq t_{max} + 0.08Z + 9.6$$

Fluctuations: mid of cascade

$$\sigma N \simeq N \Rightarrow t_{calor} \sim t(> 95\%)$$

# Electromagnetic Shower: transverse size

Molière radius:  $R_M = \frac{X_0 \cdot 21 \text{ MeV}}{E_c}$

$R < 2 \cdot R_M$  contains 95% of the shower

# Properties of Materials

Material	Density $g/cm^3$	$X_0$ $g/cm^2$	$X_0$ $cm$	$\lambda_I$ $g/cm^2$	Molière $R_M cm$	$E_{crit}$ $MeV$	Refr. index
W	19.3	6.5	0.35	185.	0.69	10.6	1.58
Pb	11.3	6.4	0.56	194.	1.22	9.6	
Cu	8.96	13.	1.45	134.	1.15	26.	
Al	2.70	24.	8.9	106.	3.3	56.	
C	2.25	42.	18.8	86.	3.5	111.	
Plastic	1.0	44.	42.	82.	6.1		
H <sub>2</sub>	0.07	61.	860.	50.	50.	360.	

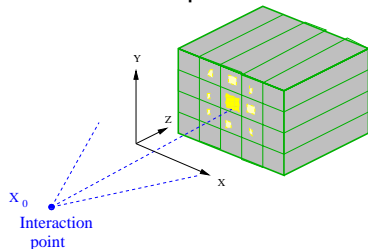


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# Generic Calorimeter

A matrix of separate elements:



Measured:

- $A_i$  - measured amplitudes
- $\alpha_i$  - calibration factors  
(slow variation)
- $x_i|y_i$  - module coordinates

$$E = \sum_{i \in k \times k} \mathcal{E}_i$$

Typically  $k = 3, 5$

$$\mathcal{E}_i = \alpha_i \cdot A_i$$

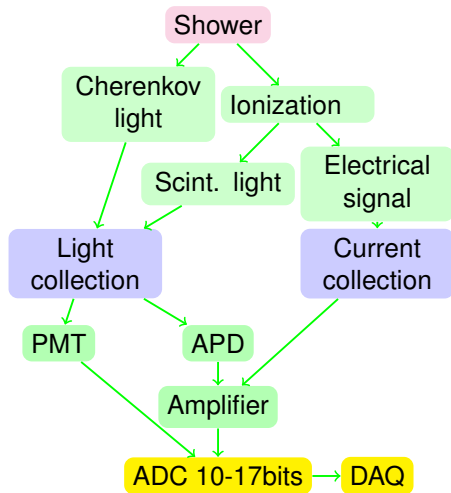
$$x|y = f(., x_i|y_i, E_i, ..)$$

$\vec{X}_0 \Rightarrow$  direction

## Important parameters

- Energy resolution  $\frac{\sigma E}{E}$
- Linearity
- Coordinate resolution  $\sigma x$
- Time resolution
- Stability
- Specific requirements:  
radiation hardness. mag. field
- Cost

# Generic Calorimeter



## Important procedures

- Calibration:  $A_i$  - measured  
 $\rightarrow E_i = \alpha_i \cdot A_i$ .  
 $\alpha_i$  have to be **measured** using particles of known energies.
- Monitoring of the calibration factors  $\alpha_i$  using detector response to a simple excitation (ex: light from a stable source).

# Homogeneous and Sampling Calorimeters

Consider: EM shower in plastic scintillator

Needed length  $\sim 15 \cdot X_0 = 600 \text{ cm}$  - not practical!

## Homogeneous calorimeters (EM)

Heavy active material, no passive absorber

- Best energy resolution
- Higher cost

## Sampling calorimeters

Heavy material absorber and the active material are interleaved.

Features:

- Compact
- Relatively cheap
- Sampling fluctuations  $\Rightarrow$  impact on  $\frac{\sigma E}{E}$

# Resolutions

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## Energy resolution

$$\frac{\sigma E}{E} = \alpha \oplus \frac{\beta}{\sqrt{E}} \oplus \frac{\gamma}{E}$$

- $\alpha$  - constant term (calibration)
- $\beta$  - stochastic term (signal/shower fluctuations)
- $\gamma$  - noise

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## Spatial resolution

$$\sigma X = \alpha_1 \oplus \frac{\beta_1}{\sqrt{E}}$$

# Energy resolution

- Fluctuations of the track length (EM):  $\frac{\sigma E}{E} \simeq \frac{0.005}{\sqrt{E}}$
- Statistics of the observed signal (EM):  $\frac{\sigma E}{E} > \frac{0.01}{\sqrt{E}}$
- Sampling fluctuations (EM):  $\frac{\sigma E}{E} \simeq \frac{\sqrt{E_c \cdot t}}{\sqrt{E}}$ , where  $t$  is the layer thickness in  $X_0$  (B.Rossi),  
 $\sim \frac{0.1 \cdot \sqrt{t}}{\sqrt{E}}$  for lead absorber ( $t > 0.2$ )
- Noise, pedestal fluctuations  $\frac{\sigma E}{E} < \frac{0.01}{E}$
- Calibration drifts  $\frac{\sigma E}{E} \sim 0.01$  for a large detector
- Other ...

# Spatial resolution

- Module lateral size < shower size
- Calculating the shower centroid
- EM:  $\sigma_X > 0.05 \cdot R_M$

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# Light Collecting Homogeneous EM Calorimeters

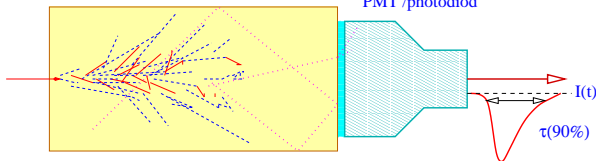
Heavy transparent materials (small  $X_0$ ) are preferable  $\Rightarrow$  compact, larger signal

- Heavy crystal scintillators: NaI, CsI, BGO, PbW etc: high light yield  $\Rightarrow$  good resolution, expensive
- Heavy crystal Cherenkov detectors: PbF, etc: compact, radiation hard
- Lead glass ( $\text{SiO} \rightarrow \text{PbO}$ ) Cherenkov detectors: medium performance, affordable

Glass / Crystal

Optical coupling

PMT / photodiode



Time resolution:

- Scintillation time
- Light bouncing
- Photodetector

Typically:

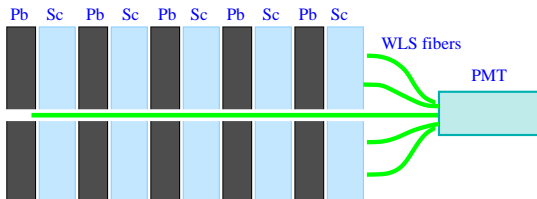
$\tau(90\%) \sim 100 \text{ ns}$  for Cherenkov detectors

Light collection 20 - 50%

# Light Collecting Sampling EM Calorimeters

Heavy absorber (Pb,Cu,W...) and a scintillator (plastic) or Cherenkov radiator (quartz fibers ...). Problem: how to collect the light? The most popular solutions for this moment:

- SPACAL (Pb, sc. fibers). The fibers can be bundled to the PM. Very good resolution. Difficult to manufacture.
- Sandwich with WLS fibers crossing through (“shashlik”). The fibers are bundled to the PM. Good resolution. Easy to build.



Time resolution:

- Scintillation time
- Photodetector time

Typically

$$\tau(90\%) \sim 50 \text{ ns}$$

# Light Detectors

## Photomultiplier Tubes (PMT)

A vacuum vessel with a photocathode and a set of electrodes (dynodes) for electron multiplication.

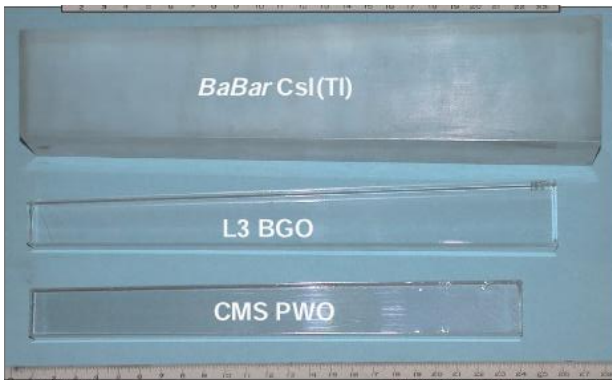
- Very high gain  $\sim 10^5 - 10^7$
- Very low electronic noise
- Size: diameter 2-40 cm
- • Slow drift of the gain
- • Sensitive to the magnetic field
- • Relatively low QE  $\sim 20\%$
- Radiation hard

## Avalanche Photodiodes (APD)

A silicon diode in avalanche mode and an electronic amplifier

- Gain  $\sim 50 - 300$
- • High electronic noise
- • Size:  $1 \times 2 \text{ cm}^2$
- • Very sensitive to the bias voltage
- Not sensitive to the magnetic field
- High QE  $\sim 75\%$  at 430 nm
- • Temperature sensitive  $-2\%/K$
- • Radiation hardness may be a problem

# Crystals in big experiments



BaBar CsI(Tl)  $\sim 10000$

L3 BGO -  $\sim 11000$

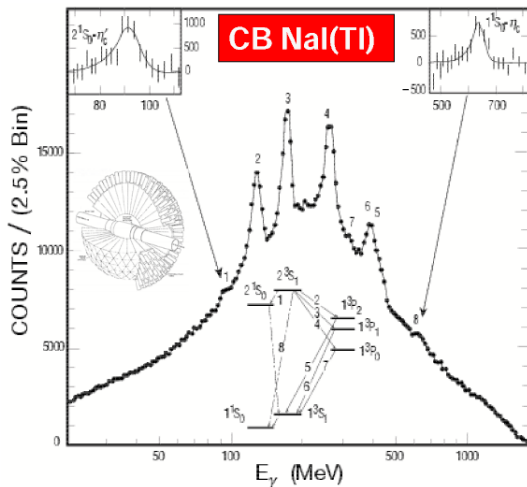
CMS PbWO -  $\sim 80000$

# EM calorimeters with optical readout

Material	Density $g/cm^3$	$X_0$ $cm$	$R_M$ $cm$	$\lambda_I$ $cm$	Refr. index	$\tau$ $ns$	Peak $\lambda$ $nm$	Light yield	$\frac{N_{p.e.}}{GeV}$	rad	$\frac{\sigma E}{E}$
Crystals											
Nal(Tl)**	3.67	2.59	4.5	41.4	1.85	250	410	1.00	$10^6$	$10^2$	$1.5\%/E^{1/4}$
CsI *	4.53	1.85	3.8	36.5	1.80	30	420	0.05	$10^4$	$10^4$	$2.0\%/E^{1/2}$
CsI(Tl)*	4.53	1.85	3.8	36.5	1.80	1200	550	0.40	$10^6$	$10^3$	$1.5\%/E^{1/2}$
BGO	7.13	1.12	2.4	22.0	2.20	300	480	0.15	$10^5$	$10^3$	$2\%/E^{1/2}$
PbWO <sub>4</sub>	8.28	0.89	2.2	22.4	2.30	5/39% 15/60% 100/01%	420 440	0.013	$10^4$	$10^6$	$2.0\%/E^{1/2}$
LSO	7.40	1.14	2.3		1.81	40	440	0.7	$10^6$	$10^6$	$1.5\%/E^{1/2}$
PbF <sub>2</sub>	7.77	0.93	2.2		1.82	Cher	Cher	0.001	$10^3$	$10^6$	$3.5\%/E^{1/2}$
Lead glass											
TF1	3.86	2.74	4.7		1.647	Cher	Cher	0.001	$10^3$	$10^3$	$5.0\%/E^{1/2}$
SF-5	4.08	2.54	4.3	21.4	1.673	Cher	Cher	0.001	$10^3$	$10^3$	$5.0\%/E^{1/2}$
SF57	5.51	1.54	2.6		1.89	Cher	Cher	0.001	$10^3$	$10^3$	$5.0\%/E^{1/2}$
Sampling: lead/scintillator											
SPACAL	5.0	1.6				5	425	0.3	$2 \cdot 10^4$	$10^6$	$6.0\%/E^{1/2}$
Shashlyk	5.0	1.6				5	425	0.3	$10^3$	$10^6$	$10\%/E^{1/2}$
Shashlyk(K)	2.8	3.5	6.0			5	425	0.3	$4 \cdot 10^5$	$10^5$	$3.5\%/E^{1/2}$

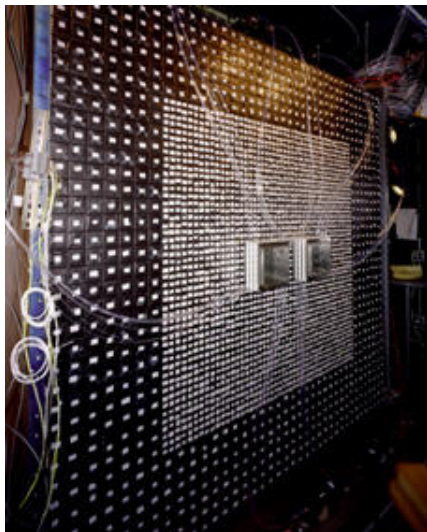
\* - hygroscopic

# Crystal Ball (SLAC, DESY)



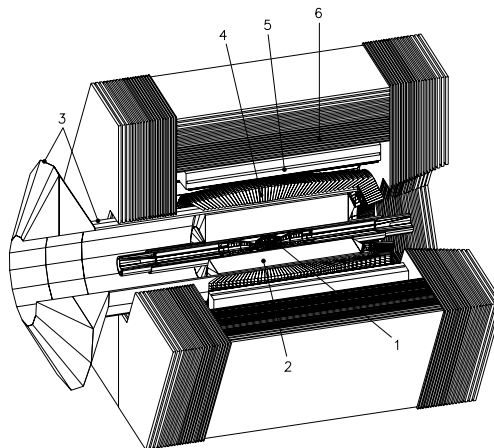
- $\sim 600$  NaI crystals
- $\gamma$  detection
- Charmonia spectra  
 $\Rightarrow$  QCD tune!

# KTeV (FNAL)

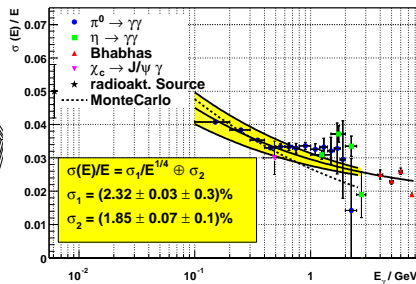


- 3256 CsI crystals
- $\pi^0 \rightarrow \gamma\gamma$  detection
- $\sigma E/E \approx 2.0\% \sqrt{E} + 0.5\%$

# BaBar (SLAC)



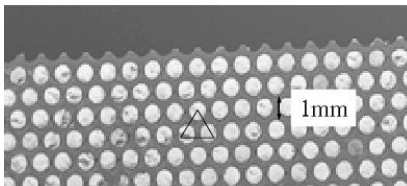
- $\sim 10000$  CsI(Tl) crystals
- $\sigma E/E \approx 2.3\%/E^{1/4} + 1.9\%$





# SpaCal (CERN, Frascati, JLab)

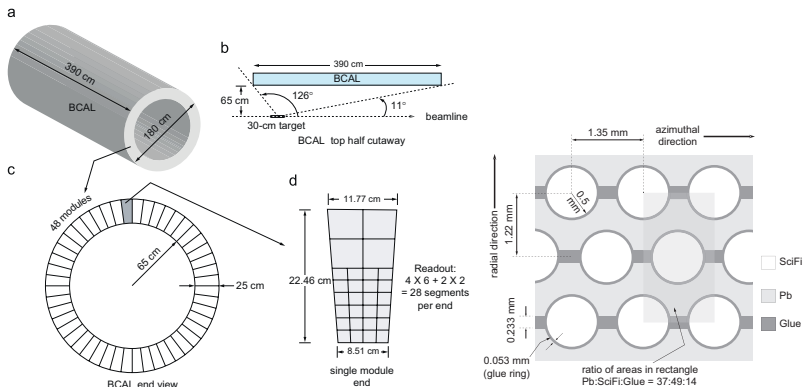
scintillating fibers / lead matrix



- Fibers/lead 50% / 50% in volume
- $X_o = 1.2 \text{ cm}$
- $5 \text{ g/cm}^3$

- CERN - original R&D
- KLOE (DAFNE) - 5000 PMTs
- KLOE  $\sigma E/E \approx 5.7\%/E^{1/2}$
- KLOE  $\sigma\tau \approx 50/E^{1/2} + 50 \text{ ps}$

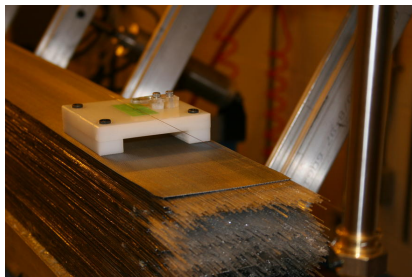
# SpaCal: Barrel Calorimeter in Hall D



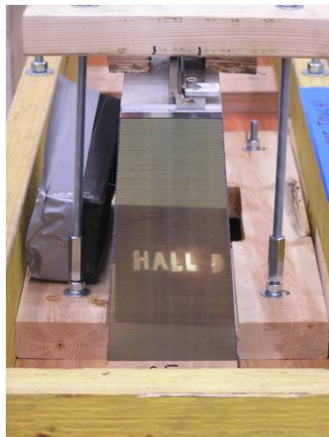
# SpaCal: Barrel Calorimeter in Hall D

## Built at Regina

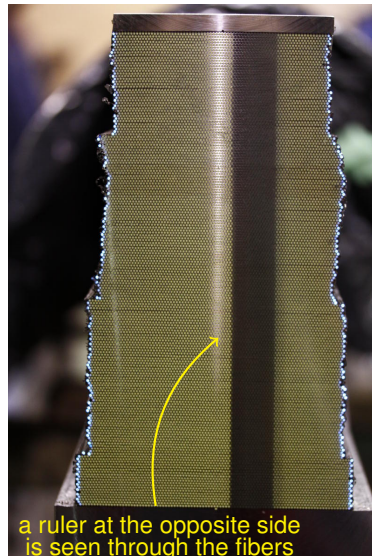
- Lead swaging (grooves)
- Glue lead and fibers layer by layer
- Cut and polish



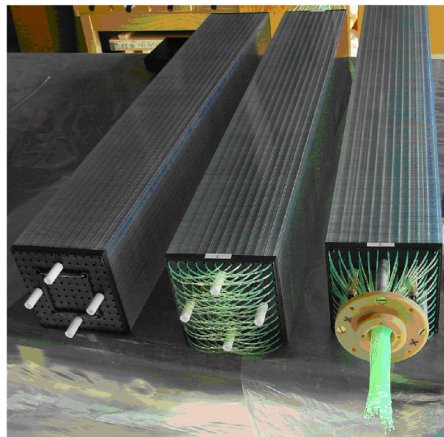
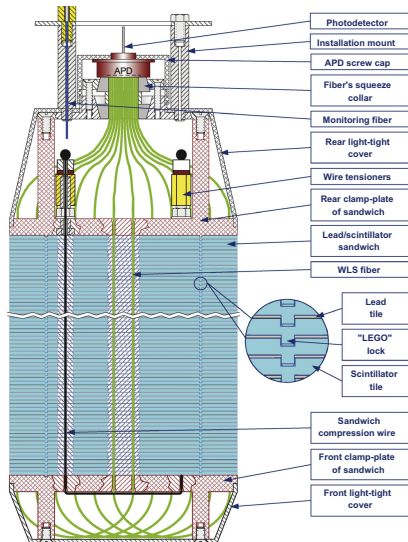
# Barrel Calorimeter - Construction of Modules



- 40 modules built
- very regular matrix



# Shashlyk: Experiment KOPIO



$$\bullet \quad \sigma E/E \approx 2.0 \oplus 3.0\%/E^{1/2}$$

$$\bullet \quad \sigma\tau \approx 70/E^{1/2} \oplus 14/E \text{ ps}$$

# Front-End Electronics

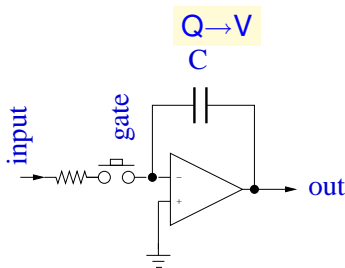
## Requirements

- Resolution  $\sim 10^{-3}$
- Dynamic range  $> 10^2$ :  
needed to measure the  
shower profile and the  
coordinates
- Differential linearity  $< 1\%$
- Digitization speed ( $> 1$  MHz)
- Readout speed ( $> 100$  kHz)
- Cost

## Existing generic solutions

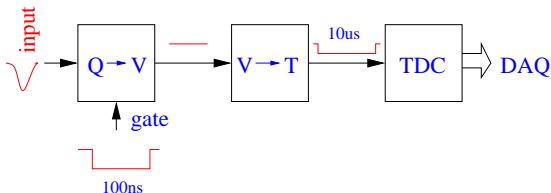
- Charge integrating ADC
- Flash ADC
- Combinations (pipeline ADC)

# Charge Integrating ADC

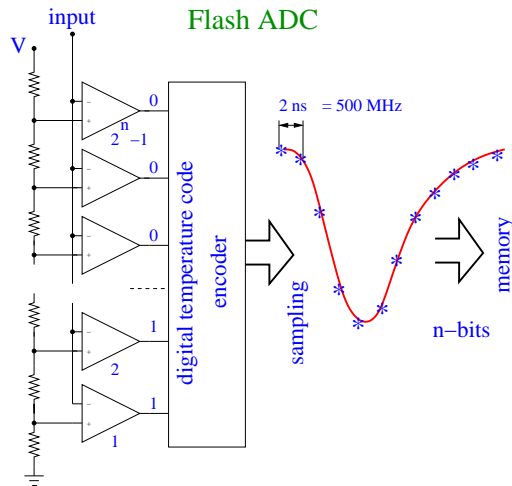


- Many products on the market
- Precise: 12-15 bits
- Gate must come in time  $\Rightarrow$  long ( $>300\text{-}500\text{ ns}$ ) delay for each channel is needed (cables)
- Slow conversion time  $> 10\text{ }\mu\text{s} \Rightarrow$  not suitable for trigger logic
- Problems at very high rate: pileup, deadtime
- Pedestal

## Integrating ADC



# Flash ADC



- Cost  $\times 10$  of the QDC (250 MHz, 12 bits)
- Huge memory buffers needed
- Resolution  $n$  bits  $\Rightarrow 2^n$  comparators
- Pipeline readout - no dead time
- No delay cables needed
- Pileup can be partially resolved
- Time resolution without extra discr.& TDCs
- FPGA computing - trigger logic
- Became the mainstream



# Calibration

The detector has to be calibrated at least once.

- Test beam
- Better: in-situ, using an appropriate process:
  - $e^+e^-$  collider: Bhabha scattering  $e^+e^- \rightarrow e^+e^-$ ,  
 $e^+e^- \rightarrow e^+e^-\gamma$
  - LHC:  $Z \rightarrow e^+e^-$  (1 Hz at low luminosity)
  - $h+h \rightarrow \pi^0+X$ ,  $\pi^0 \rightarrow \gamma\gamma$
  - RCS (JLab):  $e^-p \rightarrow e^-p$

Procedure: for event  $n$ :

$$\mathcal{E}^{(n)} = \sum_{i \in k \times k} \alpha_i \cdot A_i^{(n)}$$

$$\chi^2 = \sum_n (E^{(n)} - \sum_{i \in k \times k} \alpha_i \cdot A_i^{(n)}) / \sigma_n$$

- System of linear equations
- $\Rightarrow N \times N$  matrix - nearly diagonal
- Easy to solve

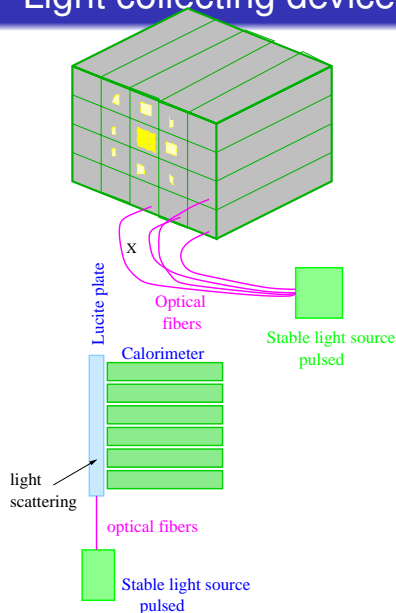
# Monitoring

## Instabilities:

- All avalanche-type devices tend to drift (PMT, gas amplification ...)
- Optical components may lose transparency
- Temperature dependence
- Many other sources of instability ...

Calibration is typically done once per many days of running  $\Rightarrow$  signal monitoring in between is needed.

# Light collecting devices



- Stable pulsed light source:
  - Xe flash lamp: 1% stability,  $>100$  ns pulse
  - Laser: 2-5% stability,  $\ll 1$  ns pulse
  - LED: 1-3% stability in thermostate,  $>30$  ns pulse
- Usually the light source has to be monitored
- Light distribution
- Material transparency: not easy to monitor ( $\lambda$ -dependence)
- Scintillation yield - no monitoring this way

# Summary

## Calorimeters are used for:

- Detecting neutrals
- Energy and coordinate measurements
- Trigger
- Separation of hadrons against  $e^{\pm}, \gamma$  and muons

The calorimeters are of increasing importance with higher energies. They became the most important/expensive/large detectors in the current big projects (LHC etc).

## Summary (continued)

There are various techniques to build calorimeters for different resolution, price, radiation hardness and other requirements.

The typical energy resolutions are:

- EM: from  $\frac{\sigma E}{E} \sim \frac{2\%}{\sqrt{E}} \oplus 0.3\%$  for scintillating crystals to about  $\frac{\sigma E}{E} \sim \frac{10\%}{\sqrt{E}} \oplus 0.8\%$  for sampling calorimeters.
- HD calorimeters:  $\frac{\sigma E}{E} \sim \frac{30-50\%}{\sqrt{E}} \oplus 3\%$

The coordinate resolutions could be about 1-3 mm for EM calorimeters and 20-30 mm for HD ones.