Nucleons in the Medium

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Theme

- Gain insight into nuclear structure from a QCD viewpoint
- Highlight opportunities provided by nuclear systems to study QCD
- Present complementary approach to traditional nuclear physics
 - formulated as a covariant quark theory
 - grounded in good description of mesons and baryons
 - at finite density self-consistent mean-field approach
- Fundamental difference
 - bound nucleons differ from free nucleons
 - medium modification effects typically ~10–15%
- Possible answers to many long standing questions: we address
 - the EMC effect [European Muon Collaboration]
 - the NuTeV anomaly [Neutrinos at the Tevatron]

EMC Effect





- J. J. Aubert *et al.* [European Muon Collaboration], Phys. Lett. B **123**, 275 (1983).
- Immediate parton model interpretation:
 - valence quarks in nucleus carry less momentum than in nucleon
- Nuclear effects seem to influence the quarks bound in the nucleons
- What is the mechanism? After 25 years no consensus
- EMC \implies medium modification

Medium Modification

- 50 years of traditional nuclear physics tells us that the nucleus is composed of nucleon-like objects
- However if a nucleon property is not protected by a symmetry its value may change in medium – for example:
 - mass, magnetic moment, size
 - quark distributions, form factors, GPDs, etc
- There must be medium modification:
 - nucleon propagator is changed in medium
 - off-shell effects ($p^2 \neq M^2$)
 - in-medium nucleon has 12 form factors instead of 2

$$\langle J^{\mu} \rangle = \sum_{\alpha,\beta=+,-} \Lambda^{\alpha}(p') \left[\gamma^{\mu} f_{1}^{\alpha\beta} + \frac{1}{2M} i \sigma^{\mu\nu} q_{\nu} f_{2}^{\alpha\beta} + q^{\mu} f_{3}^{\alpha\beta} \right] \Lambda^{\beta}(p)$$

Need to understand these effects as first step toward QCD based understanding of nuclei

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 $\langle J^{\mu} \rangle = \bar{u}(p') \left[\gamma^{\mu} F_1(Q^2) + \frac{1}{2M} i \sigma^{\mu\nu} q_{\nu} F_2(Q^2) \right] u(p)$

 Need to understand these effects as first step toward QCD based understanding of nuclei

Finite nuclei quark distributions

• Definition of finite nuclei quark distributions

$$q_A(x_A) = \frac{P^+}{A} \int \frac{d\xi^-}{2\pi} e^{\frac{iP^+ x_A \xi^-}{A}} \langle A, P | \overline{\psi}_q(0) \gamma^+ \psi_q(\xi^-) | A, P \rangle$$

Approximate using a modified convolution formalism

$$q_A(x_A) = \sum_{\alpha,\kappa,m} \int dy_A \int dx \,\,\delta(x_A - y_A \, x) \,f_{\alpha,\kappa,m}(y_A) \,\,q_{\alpha,\kappa}(x)$$



Nambu–Jona-Lasinio Model

Interpreted as low energy chiral effective theory of QCD



 Can be motivated by infrared enhancement of gluon propagator e.g. DSEs and Lattice QCD



- Investigate the role of quark degrees of freedom.
- NJL has same symmetries as QCD
 - Lagrangian: $\mathcal{L}_{NJL} = \overline{\psi} \left(i \partial m \right) \psi + G \left(\overline{\psi} \Gamma \psi \right)^2$

Nucleon in the NJL model

- Nucleon approximated as quark-diquark bound state
- Use relativistic Faddeev approach:



Nucleon quark distributions

$$q(x) = p^{+} \int \frac{d\xi^{-}}{2\pi} e^{i x p^{+} \xi^{-}} \langle p, s | \overline{\psi}_{q}(0) \gamma^{+} \psi_{q}(\xi^{-}) | p, s \rangle_{c}, \quad \Delta q(x) = \langle \gamma^{+} \gamma_{5} \rangle$$

Associated with a Feynman diagram calculation



Results: proton quark distributions



- Empirical distributions:
 - Martin, Roberts, Stirling and Thorne, Phys. Lett. B **531**, 216 (2002).
 - M. Hirai, S. Kumano and N. Saito, Phys. Rev. D 69, 054021 (2004).
- NJL model gives good description of free nucleon quark distributions
- Approach is covariant, satisfies all sum rules & positivity constraints
- DGLAP equations [Dokshitzer (1977), Gribov-Lipatov (1972), Altarelli-Parisi (1977)]

$$\frac{\partial}{\partial \ln Q^2} q_v(x, Q^2) = \alpha_s(Q^2) P(z) \otimes q_v(y, Q^2)$$

Asymmetric Nuclear Matter

• Fundamental physics: mean fields couple to the quarks in nucleons





Finite density mean-field Lagrangian: σ , ω , ρ fields

 $\mathcal{L} = \overline{\psi} \left(i \, \partial \!\!\!/ - M^* \!\!- V \!\!\!/ \right) \psi + \mathcal{L}'_I$

- σ: isoscalar-scalar attractive
 ω: isoscalar-vector repulsive
 ρ: isovector-vector attractive/repulsive
- Finite density quark propagator

$$S(k)^{-1} = k - M - i\varepsilon \quad \Rightarrow \quad S_q(k)^{-1} = k - M^* - V_q - i\varepsilon$$

Effective Potential

$$\mathcal{E} = \mathcal{E}_V - \frac{\omega_0^2}{4 G_\omega} - \frac{\rho_0^2}{4 G_\rho} + \mathcal{E}_p + \mathcal{E}_n$$

- \mathcal{E}_V : vacuum energy $\mathcal{E}_{p(n)}$: energy of nucleons moving in σ , ω , ρ fields
- Effective potential provides

$$\omega_0 = 6 G_\omega \left(\rho_p + \rho_n \right), \quad \rho_0 = 2 G_\rho \left(\rho_p - \rho_n \right), \quad \frac{\partial \mathcal{E}}{\partial M^*} = 0$$

• $G_{\omega} \Leftrightarrow Z = N$ saturation & $G_{\rho} \Leftrightarrow$ symmetry energy

Quark vector fields:

$$V_{u(d)} = \omega_0 \pm \rho_0$$

• Recall: quark propagator: $S_q(k) = [k - M^* - V_q]^{-1}$

Isovector EMC effect



• EMC ratio:
$$R = \frac{F_{2A}}{F_{2A}^{\text{naive}}} = \frac{F_{2A}}{Z F_{2p} + N F_{2n}} \sim \frac{4 u_A(x) + d_A(x)}{4 u_0(x) + d_0(x)}$$

- Density is fixed only Z/N ratio is changing
 - non-trivial isospin dependence
- proton excess: *u*-quarks feel more repulsion than *d*-quarks
- **neutron excess**: *d*-quarks feel more repulsion than *u*-quarks
- Isovector interaction \implies isovector EMC Effect

Weak mixing angle and the NuTeV anomaly



- NuTeV: $\sin^2 \theta_W = 0.2277 \pm 0.0013 (\text{stat}) \pm 0.0009 (\text{syst})$
 - G. P. Zeller et al. Phys. Rev. Lett. 88, 091802 (2002)
- World average $\sin^2 \theta_W = 0.2227 \pm 0.0004$: 3 $\sigma \implies$ "NuTeV anomaly"
- Huge amount of experimental & theoretical interest [over 400 citations]
- No universally accepted <u>complete</u> explanation

Paschos-Wolfenstein ratio motivated the NuTeV study:

$$R_{PW} = \frac{\sigma_{NC}^{\nu A} - \sigma_{NC}^{\bar{\nu} A}}{\sigma_{CC}^{\nu A} - \sigma_{CC}^{\bar{\nu} A}}, \qquad NC \Longrightarrow Z^0, \quad CC \Longrightarrow W^{\pm}$$

• For an isoscalar target $u_A \simeq d_A$ and if $s_A \ll u_A + d_A$

$$R_{PW} = \frac{1}{2} - \sin^2 \theta_W + \left(1 - \frac{7}{3}\sin^2 \theta_W\right) \frac{\langle x \, u_A^- - x \, d_A^- - x \, s_A^- \rangle}{\langle x \, u_A^- + x \, d_A^- \rangle}$$

- NuTeV "measured" R_{PW} on an Fe target ($Z/N \simeq 26/30$)
- $Z/N \neq 1 \implies$ need neutron excess corrections: $\Delta R_{PW} \neq 0$
- New realization concerning isovector EMC effect
 - isovector forces effect all u and d quarks in the nucleus
 - for N > Z there is a shift in momentum from u to d quarks
- Must correct for this isovector EMC effect in NuTeV analysis

Summary of corrections to NuTeV Analysis

- Isovector EMC effect: $\Delta R^{\rho^0} = -0.0019 \pm 0.0006$
 - sign of correction is fixed by nature of vector fields

$$q(x) = \frac{p^+}{p^+ - V^+} q_0 \left(\frac{p^+}{p^+ - V^+} x - \frac{V_q^+}{p^+ - V^+} \right), \qquad N > Z \implies V_u < V_d$$

- ρ^0 -field shifts momentum from u to d quarks
- size of correction is constrained by Nucl. Matt. symmetry energy
- Charge symmetry violation: $\Delta R^{CSV} = -0.0026 \pm 0.0011$
 - $m_u < m_d$ shifts momentum from u to d quarks
- Strange quarks: $\Delta R^s = 0.0 \pm 0.0018$
- All results include correction for NuTeV experimental acceptances, etc
- A negative ΔR_{PW} means $\sin^2 \theta_W$ decreases
 - Final result: $\sin^2 \theta_W = 0.2232 \pm 0.0013 (\text{stat}) \pm 0.0024 (\text{syst})$

Total NuTeV correction



- No evidence for physics beyond the Standard Model
- Instead "NuTeV anomaly" is evidence for medium modification
 - Equally interesting
 - EMC effect has generated over 1000 papers

Observable elsewhere ... Parity Violating DIS



• Isovector EMC effect can be observed in PV DIS [γZ^0 interference]

$$A^{PV} = \frac{d\sigma_R - d\sigma_L}{d\sigma_R + d\sigma_L} \propto [a_2(x) + \ldots]; \quad a_2(x) = g_A^e \frac{F_2^{\gamma Z}}{F_2^{\gamma}}$$

• For $N \simeq Z$ target

$$a_2(x) \simeq \left(\frac{9}{5} - 4\sin^2\theta_W\right) - \frac{12}{25} \frac{u_A^+(x) - d_A^+(x)}{u_A^+(x) + d_A^+(x)}$$

Large x dependence of $a_2(x)$ even after naive correction

evidence for medium modification

Finite nuclei EMC effects

EMC ratio

$$R = \frac{F_{2A}}{F_{2A}^{\text{naive}}} = \frac{F_{2A}}{Z F_{2p} + N F_{2n}}$$

• Polarized EMC ratio

$$R_s^H = \frac{g_{1A}^H}{g_{1A}^{H,\text{naive}}} = \frac{g_{1A}^H}{P_p^H \, g_{1p} + P_n^H \, g_{1n}}$$

- Spin-dependent cross-section is suppressed by 1/A
 - Must choose nuclei with $A \lesssim 27$
 - protons should carry most of the spin e.g. \implies ⁷Li, ¹¹B, ...
- Ideal nucleus is probably ⁷Li
 - From Quantum Monte–Carlo: $P_p^J = 0.86$ & $P_n^J = 0.04$
- Ratios equal 1 in non-relativistic and no-medium modification limit

EMC ratio ⁷Li, ¹¹B, ¹⁵N and ²⁷Al



Is there medium modification



Is there medium modification



- Medium modification of nucleon has been switched off
- Relativistic effects remain
- Large splitting would be strong evidence for medium modification

Nuclear Spin Sum

Proton spin states	Δu	Δd	Σ	g_A
p	0.97	-0.30	0.67	1.267
7 Li	0.91	-0.29	0.62	1.19
^{11}B	0.88	-0.28	0.60	1.16
15 N	0.87	-0.28	0.59	1.15
27 Al	0.87	-0.28	0.59	1.15
Nuclear Matter	0.79	-0.26	0.53	1.05

- Angular momentum of nucleon: $J = \frac{1}{2} = \frac{1}{2} \Delta \Sigma + L_q + J_g$
 - in medium $M^* < M$ and therefore quarks are more relativistic
 - Iower components of quark wavefunctions are enhanced
 - quark lower components usually have larger angular momentum
 - $\Delta q(x)$ very sensitive to lower components
- Conclusion: quark spin → orbital angular momentum in-medium

Conclusion

- Effective quark theories can be used to incorporate quarks into a traditional description of nuclei
 - complementary approach to traditional nuclear physics
- Major outstanding discrepancy with Standard Model predictions for Z⁰ is NuTeV anomaly
 - may be resolved by CSV and isovector EMC effect corrections
- EMC effect and NuTeV anomaly are interpreted as evidence for medium modification of the bound nucleon wavefunction
- This result can be tested using PV DIS measurements
 - predict large medium modification in PV DIS
 - predict flavour dependence of EMC effect can be large
- In nuclei quark spin converted to orbital angular momentum
 - Polarized EMC effect