

*Proton Form Factor Ratio,  $G_E^P/G_M^P$   
From  
Double Spin Asymmetries*

Spin  
Asymmetries of the  
Nucleon  
Experiment  
( E07-003)

**Analysis Updates**

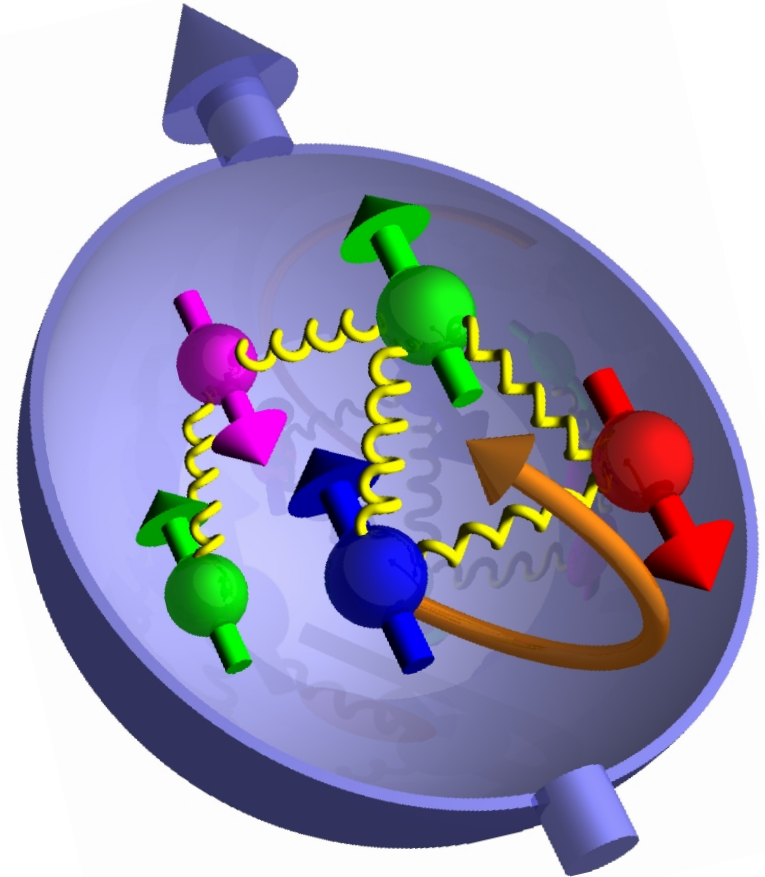


**Jefferson Lab**  
● Thomas Jefferson National Accelerator Facility

Anusha Liyanage  
HU Group Meeting  
(December 04, 2012)

# Outline

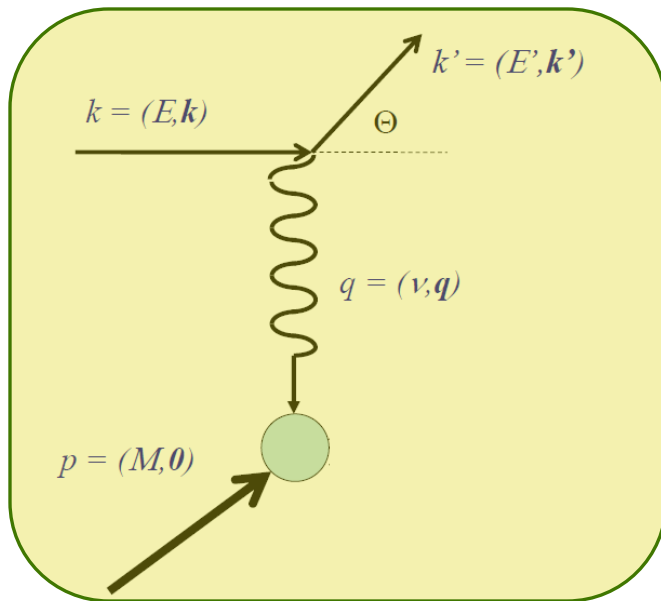
- Introduction
- Physics Motivation
- Experiment Setup
  - BETA Detector
  - HMS Detector
  - Polarized Target
- Elastic Kinematic
- Data Analysis &  
MC/SIMC Simulation
- Conclusion



# Introduction

## Nucleon Elastic Form Factors

- Defined in context of single-photon exchange.
- Describe how much the nucleus deviates from a point like particle.
- Describe the internal structure of the nucleons.
- Provide the information on the spatial distribution of electric charge (by electric form factor,  $G_E$ ) and magnetic moment (by magnetic form factor,  $G_M$ ) within the proton.
- Can be determined from elastic electron-proton scattering.
- They are functions of the four-momentum transfer squared,  $Q^2$



The four-momentum transfer squared,

$$Q^2 = -q^2 = 4EE' \sin^2\left(\frac{\Theta}{2}\right)$$

$$E - E' = \frac{Q^2}{2M}$$

# General definition of the nucleon form factor is

$$\langle N(P') | J_{EM}^\mu(0) | N(P) \rangle = \bar{u}(P') \left[ \gamma^\mu F_1^N(Q^2) + i\sigma^{\mu\nu} \frac{q_\nu}{2M} F_2^N(Q^2) \right] u(P)$$

Sachs Form Factors  $G_E = F_1 - \tau F_2$  ;  $G_M = F_1 + F_2$  ;  $\tau = \frac{Q^2}{4M^2}$

$F_1$  – non-spin flip (Dirac Form Factor) describe the charge distribution

$F_2$  – spin flip (Pauli form factor) describe the magnetic moment distribution

At low  $|q^2|$

$$G_E(q^2) \approx G_E(\vec{q}^2) = \int e^{i\vec{q}\cdot\vec{r}} \rho(\vec{r}) d^3\vec{r}$$

$$G_M(q^2) \approx G_M(\vec{q}^2) = \int e^{i\vec{q}\cdot\vec{r}} \mu(\vec{r}) d^3\vec{r}$$

Fourier transforms of the charge,  $\rho(r)$   
and magnetic moment,  $\mu(r)$  distributions  
in Breit Frame

At  $q^2 = 0$

$$G_E(0) = \int \rho(\vec{r}) d^3\vec{r} = 1$$

$$G_M(0) = \int \mu(\vec{r}) d^3\vec{r} = \mu_p = +2.79$$

$$\mu \frac{G_E^p}{G_M^p} = 1$$

# Form Factor Ratio Measurements

## 1. Rosenbluth separation method.

- Measured the electron - unpolarized proton elastic scattering cross section at fixed  $Q^2$  by varying the scattering angle,  $\theta_e$ .
- Strongly sensitive to the radiative corrections.

$$\frac{d\sigma}{d\Omega} = \underbrace{\frac{\alpha^2 E' \cos^2 \frac{\theta_e}{2}}{4(1+\tau)E^3 \sin^4 \frac{\theta_e}{2}}}_{\sigma_{Mott} / (1+\tau)} \left[ G_E^2 + \frac{\tau}{\epsilon} G_M^2 \right]$$

$$Q^2 = 2EE'(1 - \cos\theta_e)$$

$$\tau = \frac{Q^2}{4M^2}$$

$$\epsilon = \left[ 1 + 2(1+\tau) \tan^2 \frac{\theta_e}{2} \right]^{-1}$$

$$\frac{d\sigma}{d\Omega} \cdot \frac{\epsilon(1+\tau)}{\sigma_{Mott}} = G_E^2 \epsilon + \tau G_M^2$$

$$Y = m X + C$$

The gradient =  $G_E^2$  , The Intercept =  $\tau G_M^2$  ,

E - Incoming electron energy

E' - Outgoing electron energy

$\theta_e$  - Outgoing electron's scattering angle

M - Proton mass

## 2. Polarization Transfer Technique.

- Measured the recoil proton polarization from the elastic scattering of polarized electron-unpolarized proton.
- Insensitive to absolute polarization, analyzing power.
- Less sensitive to radiative correction.

$$\frac{G_E}{G_M} = -\frac{P_T}{P_L} \frac{(E + E') \tan\left(\frac{\theta_e}{2}\right)}{2M_p}$$

$E$  - Incoming going electron energy

$E'$  - Out going electron energy

$\theta_e$  - Outgoing electron's scattering angle

$M_p$  - Proton mass

$$P_L = M_p^{-1} (E + E') \sqrt{\tau(1 + \tau)} G_M^2 \tan^2(\theta_e / 2) \longrightarrow \text{Polarization along } q$$

$$P_T = 2\sqrt{\tau(1 + \tau)} G_E G_M \tan(\theta_e / 2) \longrightarrow \text{Polarization perpendicular to } q$$

(in the scattering plane)

$$P_N = 0 \longrightarrow \text{Polarization normal to scattering plane.}$$

### 3. Double-Spin Asymmetry.

- Measured the cross section asymmetry between + and – electron helicity states in elastic scattering of a polarized electron on a polarized proton.
- The systematic errors are different when compared to either the Rosenbluth technique or the polarization transfer technique.
- The sensitivity to the form factor ratio is the same as the Polarization Transfer Technique.

$$A_p = \frac{-br \sin \theta^* \cos \phi^* - a \cos \theta^*}{r^2 + c}$$

$$\frac{G_E}{G_M} = -\frac{b}{2A_p} \sin \theta^* \cos \phi^* + \sqrt{\frac{b^2}{4A_p^2} \sin^2 \theta^* \cos^2 \phi^* - \frac{a}{A_p} \cos \theta^* - c}$$

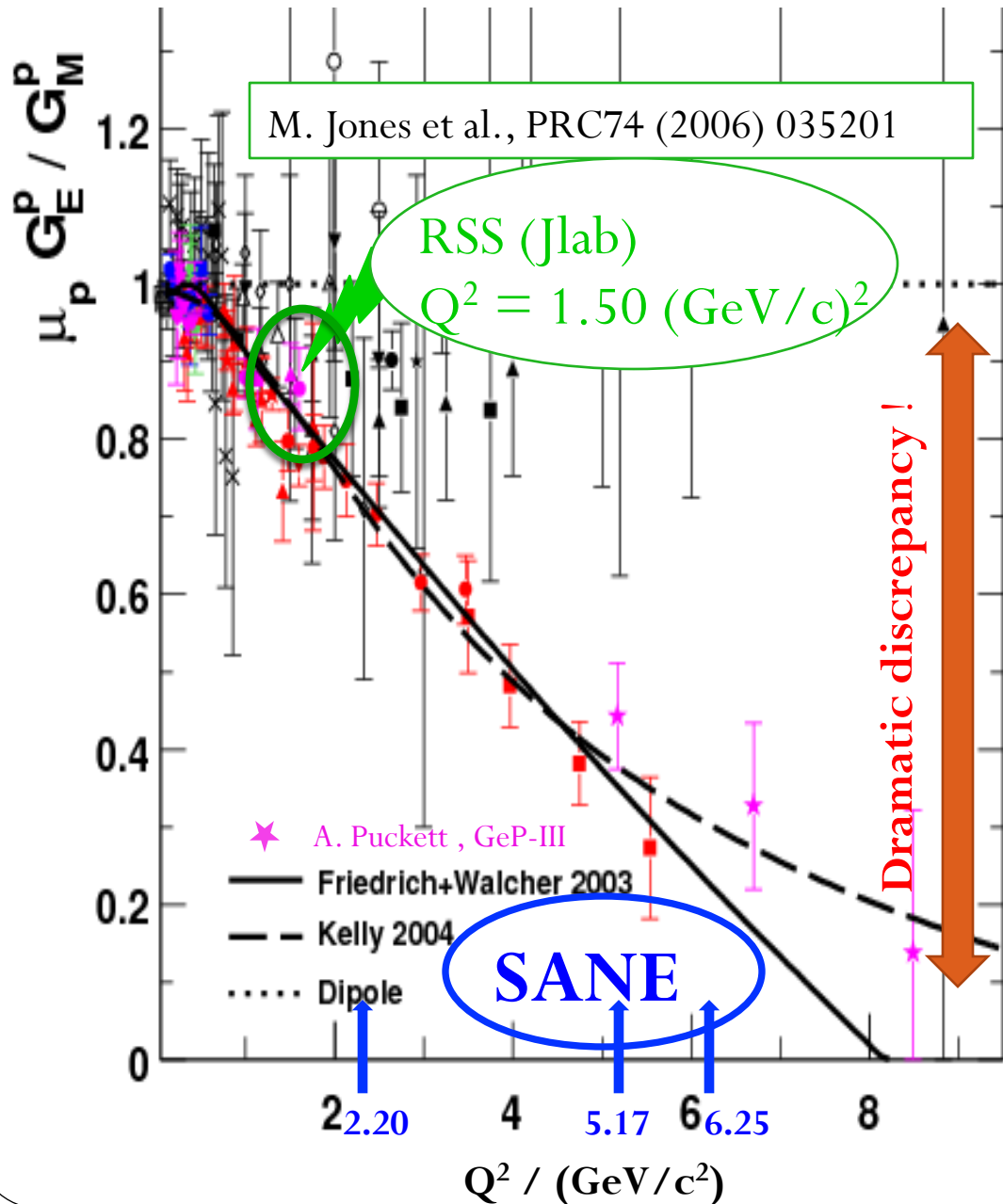
Here,  $r = G_E / G_M$

$a, b, c =$  kinematic factors

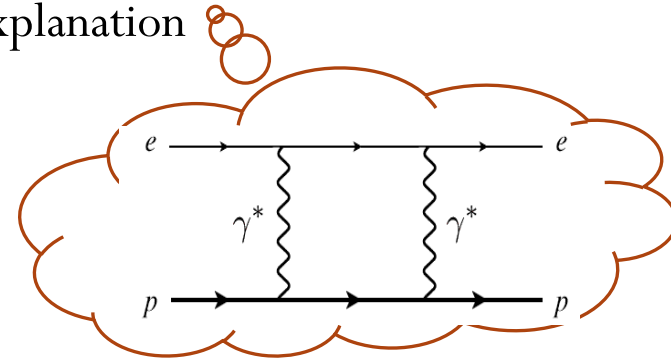
$\theta^*, \phi^* =$  pol. and azi. Angles between  $\vec{Q}$  and  $\vec{S}$

$A_p =$  The beam - target asymmetry

# Physics Motivation



- Dramatic discrepancy between Rosenbluth and recoil polarization technique.
- Multi-photon exchange considered the best candidate for the explanation

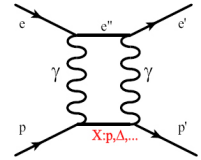
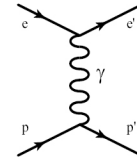


- **Double-Spin Asymmetry** is an Independent Technique to verify the discrepancy



# Two-Photon Exchange

- Both Rosenbluth method and the polarization transfer technique account for radiative correction, but neither consider two photon exchange.



- Contribution of the TPE amplitude has calculated theoretically and, has an  $\epsilon$  dependence that has the same sign as the  $G_E$  contribution to the cross section and is large enough to effect the extracted value of  $G_E$ .

Therefore, the extracted  $G_E/G_M$  for the Rosenbluth technique is reduced.

- The effect of TPE amplitude on the polarization components is small, though the size of the contribution change with  $\epsilon$
- The size of the TPE would measure by taking the  $\epsilon$  dependence of the ratio of cross sections, R for elastic electron-proton scattering to positron-proton scattering at a fixed  $Q^2$  and measuring the deviation from 1.

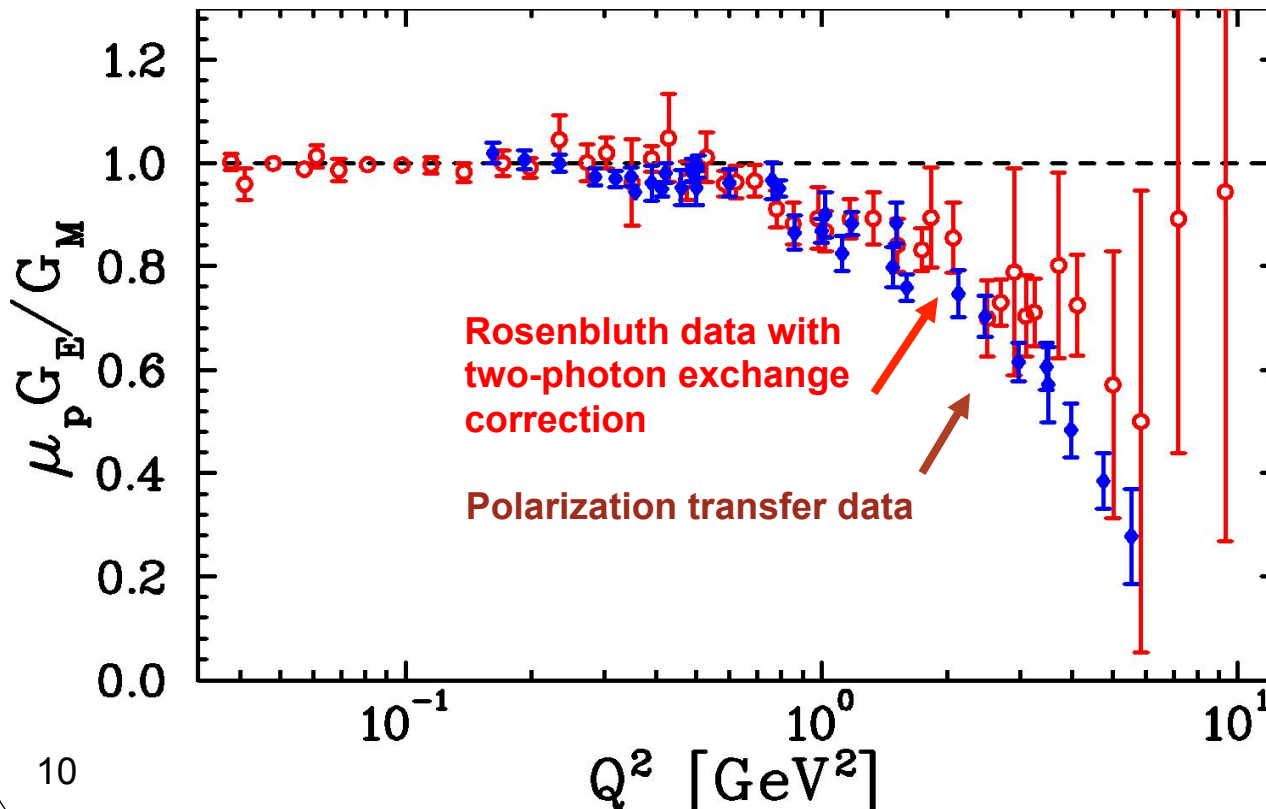
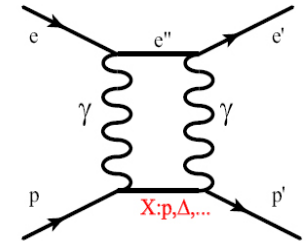
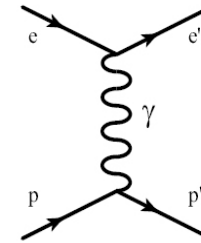
$$R = \frac{\sigma_{e^+}}{\sigma_{e^-}} = \frac{(A_{1\gamma} + A_{2\gamma})^2}{(A_{1\gamma} - A_{2\gamma})^2} \approx 1 + 4 \operatorname{Re} \left( \frac{A_{2\gamma}}{A_{1\gamma}} \right)$$

# Two-Photon Exchange: Exp. Evidence

Two-photon exchange theoretically suggested

TPE can explain form factor discrepancy

J. Arrington, W. Melnitchouk, J.A. Tjon,  
Phys. Rev. C 76 (2007) 035205



# Asymmetry measurements

$$\sigma = \sigma_0 + P_E P_T \Delta\sigma$$



$$\sigma_{++} = \sigma_0 + P_E P_T \Delta\sigma$$

$$\sigma_{+-} = \sigma_0 - P_E P_T \Delta\sigma$$

$$\frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = P_E P_T \cdot \frac{\Delta\sigma}{\sigma_0} = \frac{N_+ - N_-}{N_+ + N_-} = A_r$$

$$\frac{A_r}{P_E P_T} = \frac{\Delta\sigma}{\sigma_0} = A_p$$

Hence,

$A_p$ , known as the physics asymmetry is the relative scattering cross section correction due to the spin.

$A_r$  is the raw asymmetry

$\sigma$  - Scattering cross section

$\sigma_0$  - Scattering cross section at unpolarized target

$\sigma_B$  - Scattering cross section from background

$\Delta\sigma$  -  $\sigma$  correction due to the spin

$P_E$  - Beam polarization

$P_T$  - Target polarization

$f$  - Dilution factor

With background....

$$\sigma_{++} = \sigma_0 + P_E P_T \Delta\sigma + \sigma_B$$

$$\sigma_{+-} = \sigma_0 - P_E P_T \Delta\sigma + \sigma_B$$

$$A_r = P_E P_T \cdot \frac{\Delta\sigma}{(\sigma_0 + \sigma_B)}$$

$$A_r = P_E P_T \cdot \frac{\Delta\sigma}{\sigma_0} \cdot \frac{\sigma_0}{(\sigma_0 + \sigma_B)} f$$

$$A_p = \frac{A_r}{f P_E P_T}$$

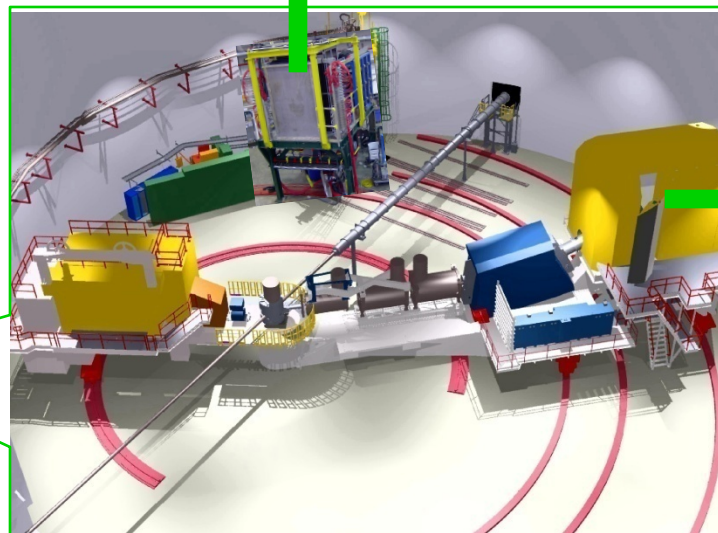
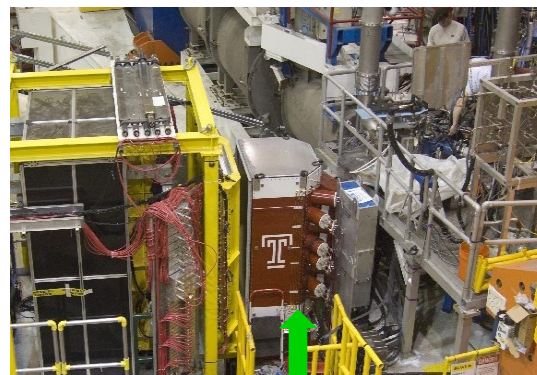
# Experiment Setup



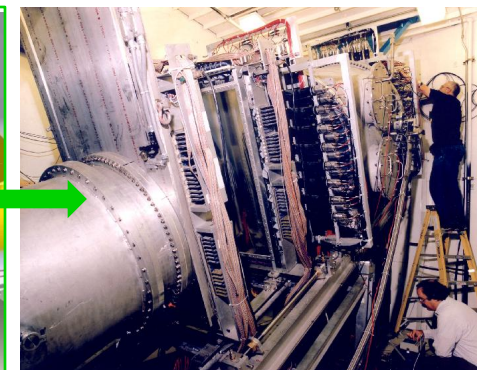
Hall C at  
Jefferson Lab

Elastic ( $e, e'p$ ) scattering from the polarized  $\text{NH}_3$  target using a longitudinally polarized electron beam

(Data collected from Jan – March, 2009)



- BETA for coincidence electron detection
- Central scattering angle :  $40^\circ$
- Over 200 msr solid angle coverage

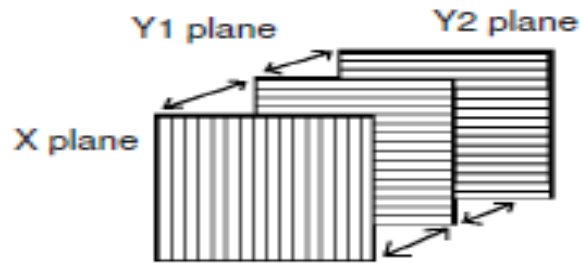


- HMS for the scattered proton detection
- Central angles are  $22.3^\circ$  and  $22.0^\circ$
- Solid angle  $\sim 10$  msr

# Big Electron Telescope Array – BETA

## Forward Tracker

- 3 planes of Bicron Scintillator provide early particle tracking

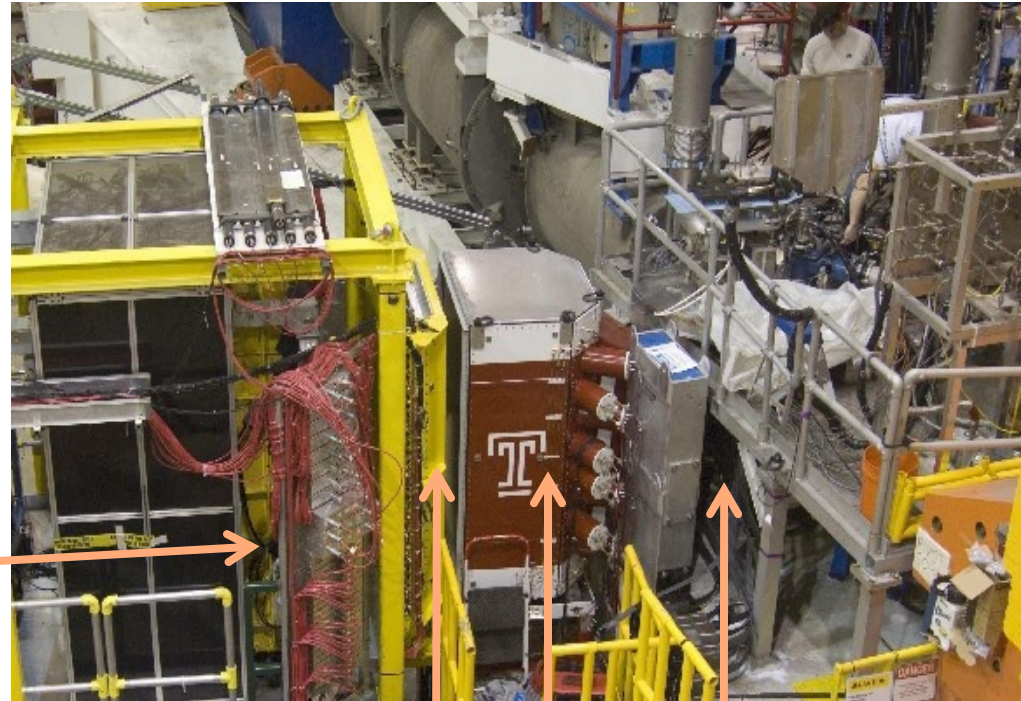


## Cerenkov

- N<sub>2</sub> gas cerenkov
- Provides particle ID
- 8 mirrors and 8 PMTs

## Lucite Hodoscope

- 28 bars of 6cm wide Lucite
- Bars oriented horizontally for Y tracking
- PMTs on either side of bar provides X resolution



BigCal

Lucite Hodoscope

Tracker

Cerenkov

Big Cal (GEP III Collaboration)

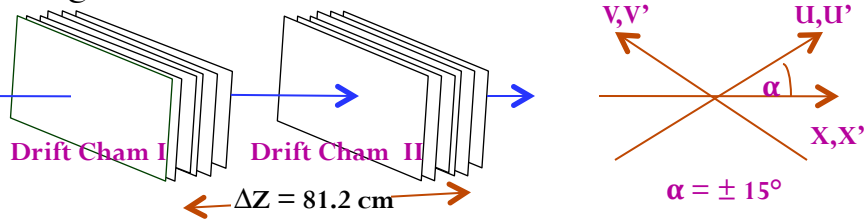
Lead glass calorimeter

- 1744 blocks aprx. 4cm x 4cm
- energy and position measurement

# High Momentum Spectrometer – HMS

## Drift Chambers

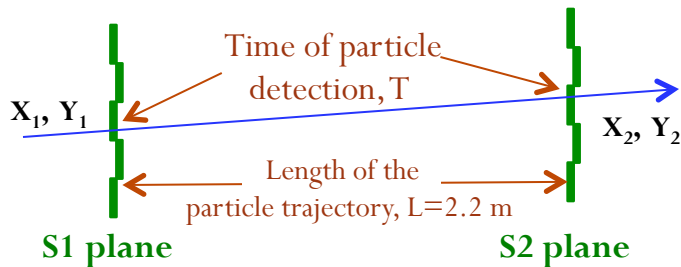
- Each plane has a set of alternating field and sense wires Filled with an equal parts Argon-Methane mixture



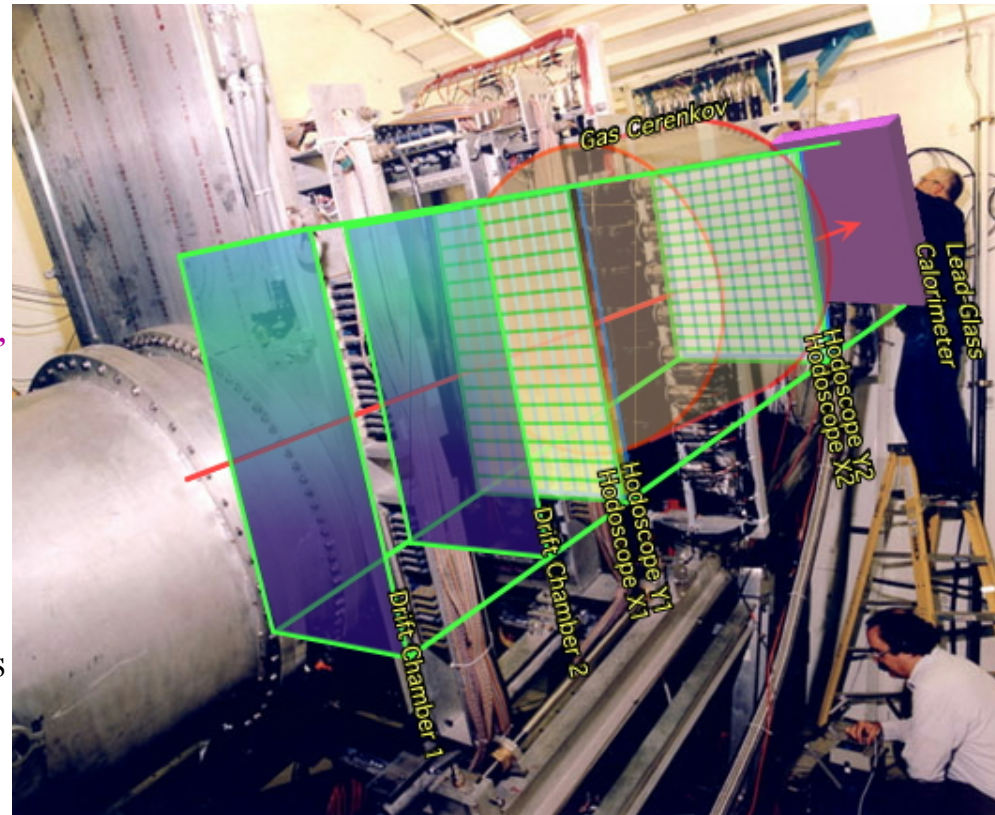
- Track particle trajectory by multiple planes.
- $\chi^2$  fitting to determine a straight trajectory.

## Hodoscopes

- Each plane contains 10 to 16 Scintillator paddles with PMTs on both ends
- Each Paddle is 1.0 cm thick and 8.0 cm wide



- Fast position determination & triggering
- Time of Flight (TOF) =  $T_2 - T_1$  determines  $\beta$   
( $\beta = L/c \times \text{TOF}$ )



## Gas Cerenkov

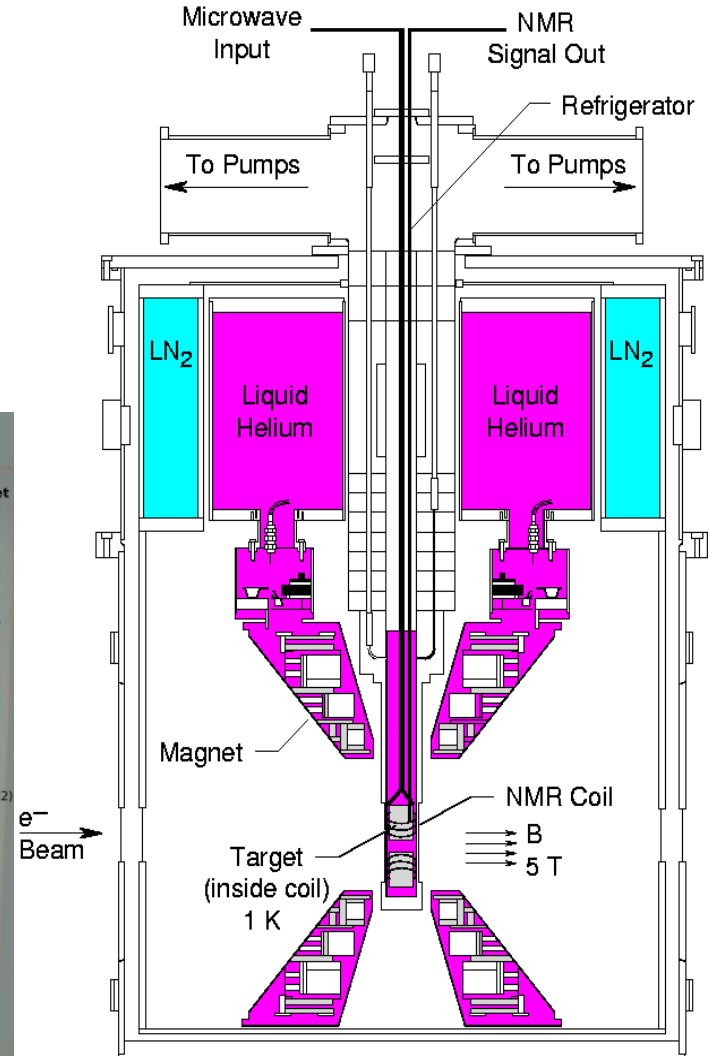
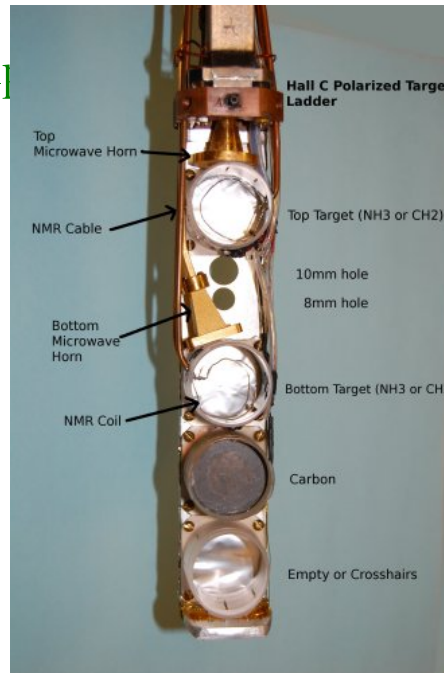
- Two mirrors (top & bottom) connected to two PMTs
- Used as a Particle ID

## Lead Glass Calorimeter

- 4 layers of 10 cm x 10cm x 70cm blocks stacked 13 high.
- Used as a Particle ID

# Polarized Target

- C, CH<sub>2</sub> and NH<sub>3</sub>
- Dynamic Nuclear Polarization (DNP) polarized the protons in the NH<sub>3</sub> target up to 90% at  
1 K Temperature  
5 T Magnetic Field
- Temperature is maintained by immersing the entire target in the liquid He bath
- Used microwaves to excite spin flip transitions (55 GHz - 165 GHz)
- Polarization measured using NMR coils
- To maintain reasonable target polarization, the beam current,
  - limited to 100 nA
  - Was uniformly rastered.

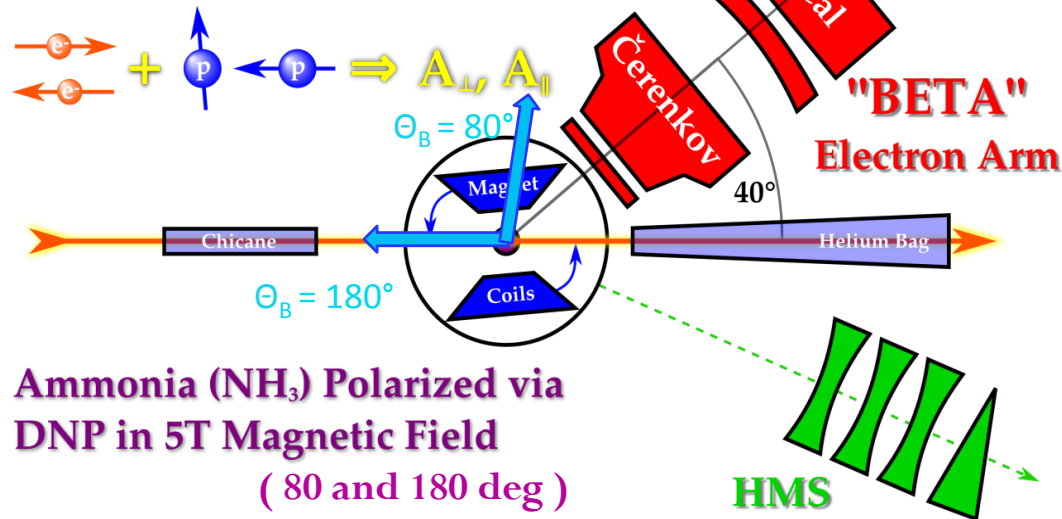


The Polarized Target Assembly

# Polarized Target Magnetic Field

**Polarized Electron Beam: 4.7, 5.9 GeV**

**Polarized Proton Target:  $\sim \perp, \parallel$**



- Used only perpendicular magnetic field configuration for the elastic data
  - Average target polarization is  $\sim 70\%$
  - Average beam polarization is  $\sim 73\%$



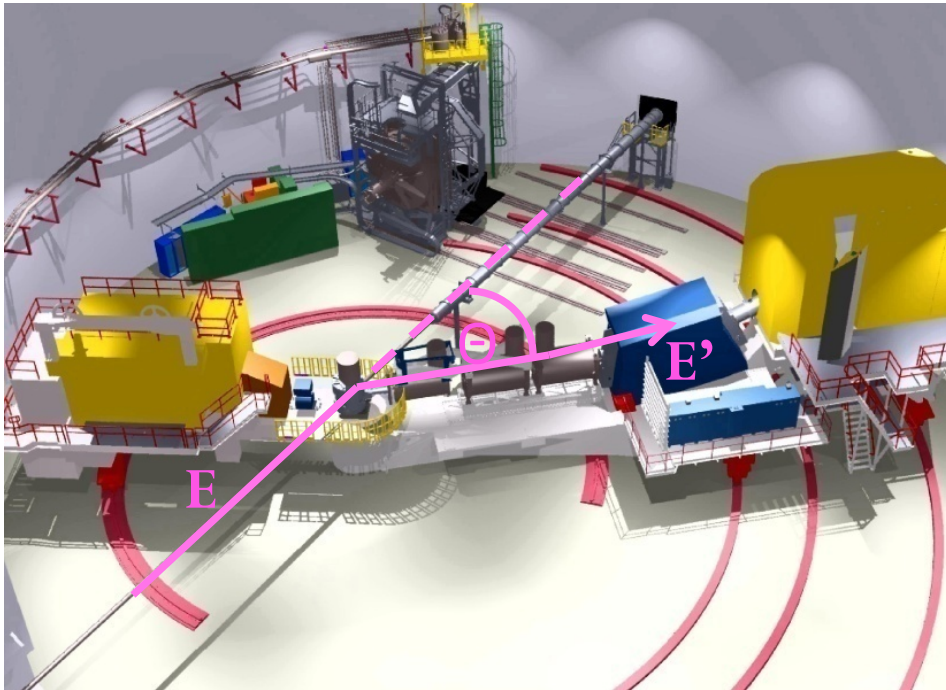
# Elastic Kinematics

( From HMS Spectrometer )

Spectrometer mode	Coincidence	Coincidence	Single Arm
HMS Detects	Proton	Proton	Electron
E Beam GeV	4.72	5.89	5.89
$P_{\text{HMS}}$ GeV/c	3.58	4.17	4.40
$\Theta_{\text{HMS}}$ (Deg)	22.30	22.00	15.40
$Q^2$ (GeV/c) <sup>2</sup>	5.17	6.26	2.20
Total Hours (h)	~40 (~44 runs)	~155 (~135 runs)	~12 (~15 runs)
Elastic Events	~113	~1200	-

# Data Analysis

## Electrons in HMS



$$\vec{e}^- \vec{p} \longrightarrow e^- p$$

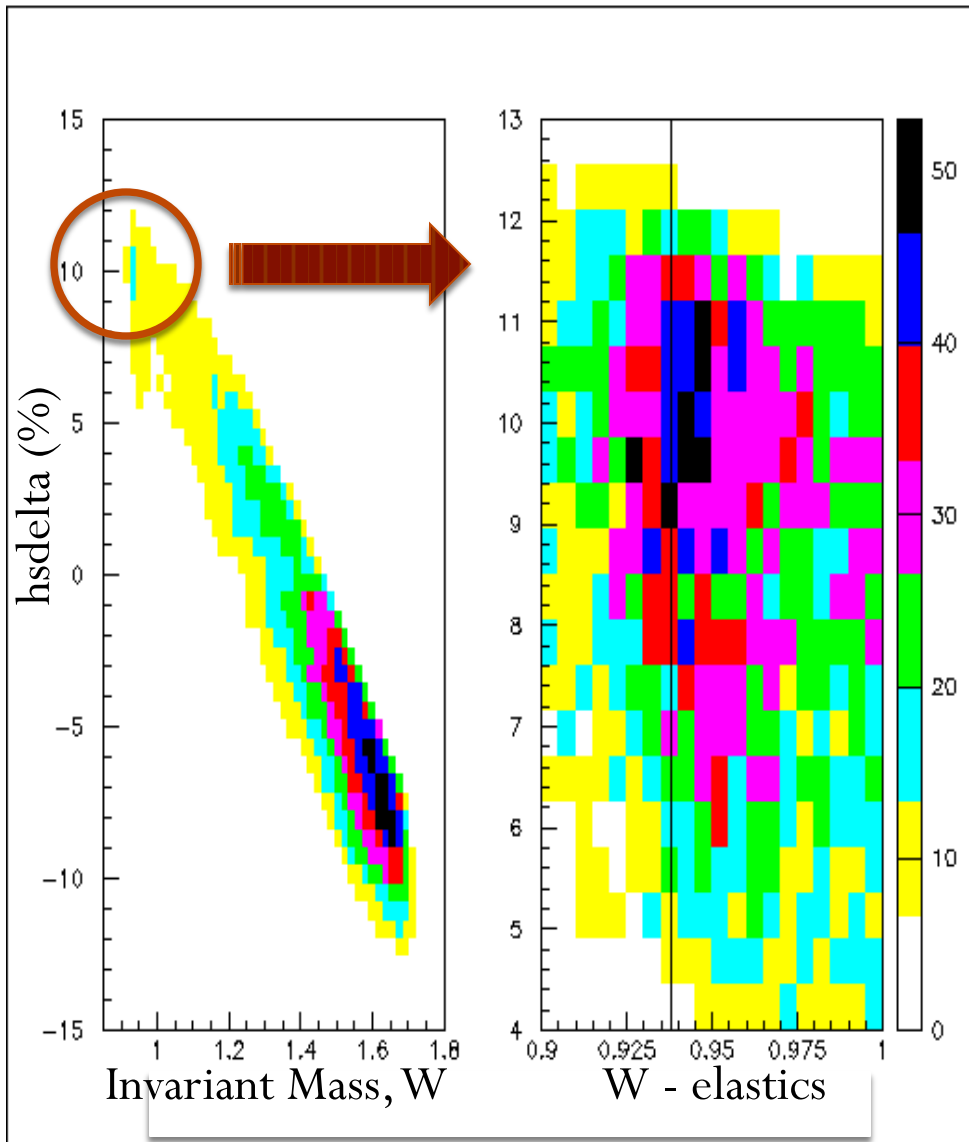
By knowing,  
the incoming beam energy,  $E$ ,  
scattered electron energy,  $E'$   
and  
the scattered electron angle,  $\theta$

$$Q^2 = 4EE' \sin^2\left(\frac{\theta}{2}\right)$$



$$W^2 = M^2 - Q^2 + 2M(E - E')$$

- Momentum Acceptance



$$hsdelta = \left( \frac{P - P_c}{P_c} \right) = \frac{\delta p}{p}$$

$P$  - Measured momentum in HMS

$P_c$  - HMS central momentum

The elastic data are outside of the usual delta cut  $\pm 8\%$

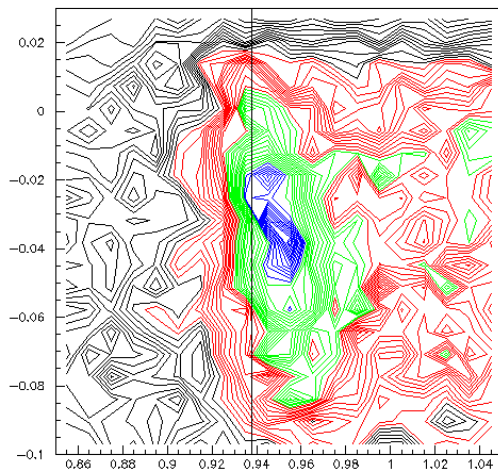
Because HMS reconstruction matrix elements work fine up to 10



Use  $-8\% < hsdelta < 10\%$

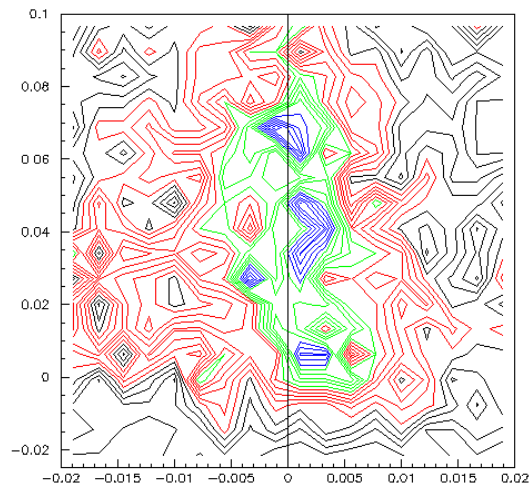
# Perp. target magnetic field make some correlations....

In Single Arm electron data



Xptar vs W

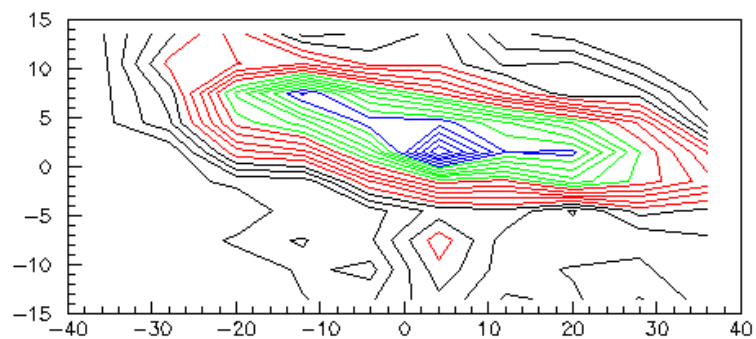
In COIN HMS data



Xptar vs dpel\_hms

- Introduced an 'azimuthal angle correction' which correct the target magnetic field in vertical direction in terms of the azimuthal angle. (First make the same correlations on MC/SIMC by applying the correction only for the forward direction and then use the correction on data)
- Different corrections for different detector angles.

In COIN BETA data



Y\_HMS-y\_clust vs y\_clust

# Extract the electrons

- Used only Electron selection cuts.

# of Cerenkov photoelectrons  $> 2$

$$E_{sh}/E' > 0.7$$

$$\left( \frac{P - P_c}{P_c} \right) < 10 \text{ and } \left( \frac{P - P_c}{P_c} \right) > -8$$

- Cerenkov cut

- Calorimeter cut

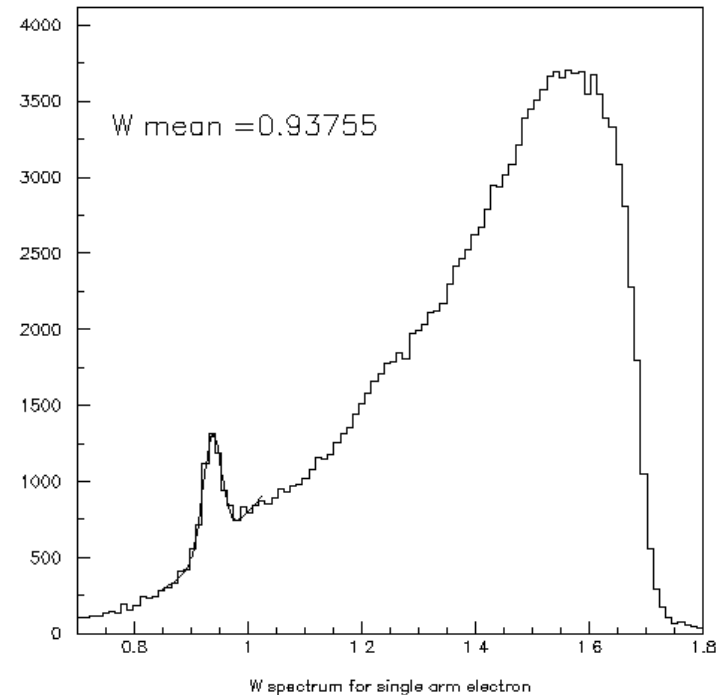
- HMS Momentum Acceptance cut

Here,

$P/E'$  - Detected electron momentum/  
energy at HMS

$P_c$  - Central momentum of HMS

$E_{sh}$  - Total measured shower energy  
of a chosen electron track by  
HMS Calorimeter



# Extracted the Asymmetries .....

The raw asymmetry,  $A_r$

$$A_r = \frac{N^+ - N^-}{N^+ + N^-}$$

$$\Delta A_r = \frac{2\sqrt{N^+} \sqrt{N^-}}{(N^+ + N^-)\sqrt{(N^+ + N^-)}}$$

$N^+ / N^- =$  Charge and life time normalized counts for the +/- helicities

$\Delta A_r =$  Error on the raw asymmetry

$P_B P_T =$  Beam and Target polarization

$N_c =$  A correction term to eliminates the contribution from quasi-elastic  $^{15}\text{N}$  scattering under the elastic peak

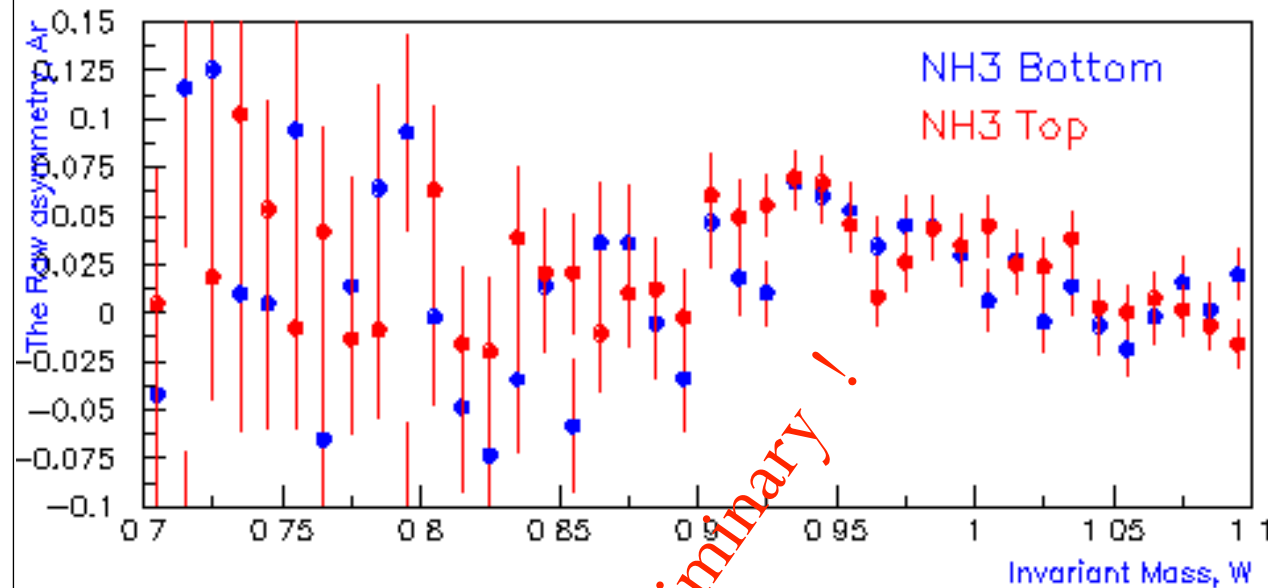
## The Asymmetries

## Need

*dilution factor,  $f$*   
in order to determine the  
*physics asymmetry,*

$$A_p = \frac{A_r}{f P_B P_T} + N_c$$

and  $G_E^p / G_M^p$   
(at  $Q^2 = 2.2 \text{ (GeV/c)}^2$ )

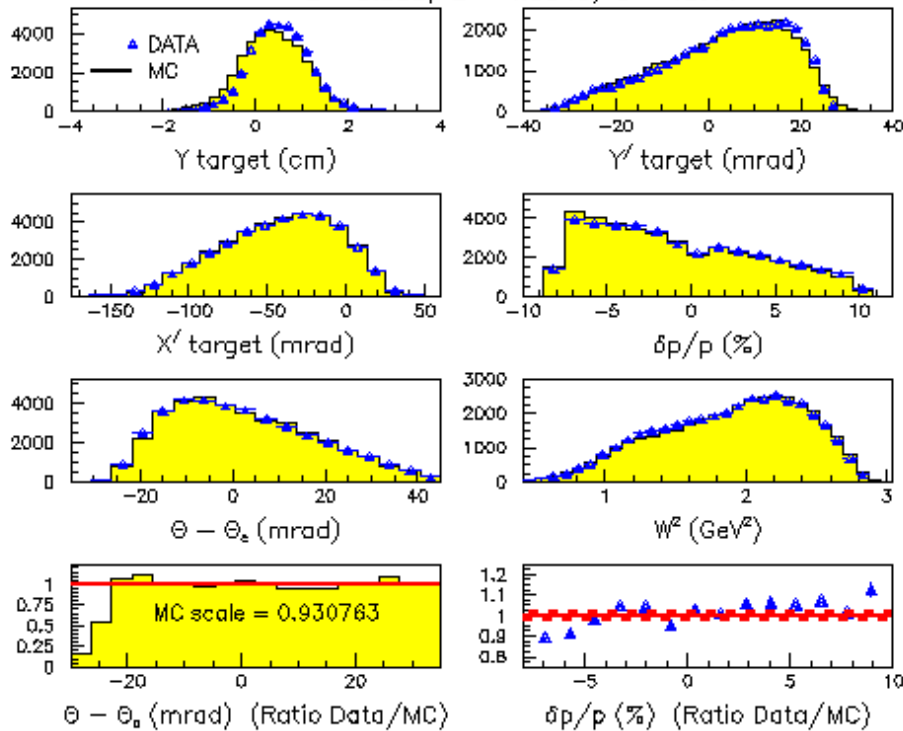


**Preliminary!**

# MC for C run

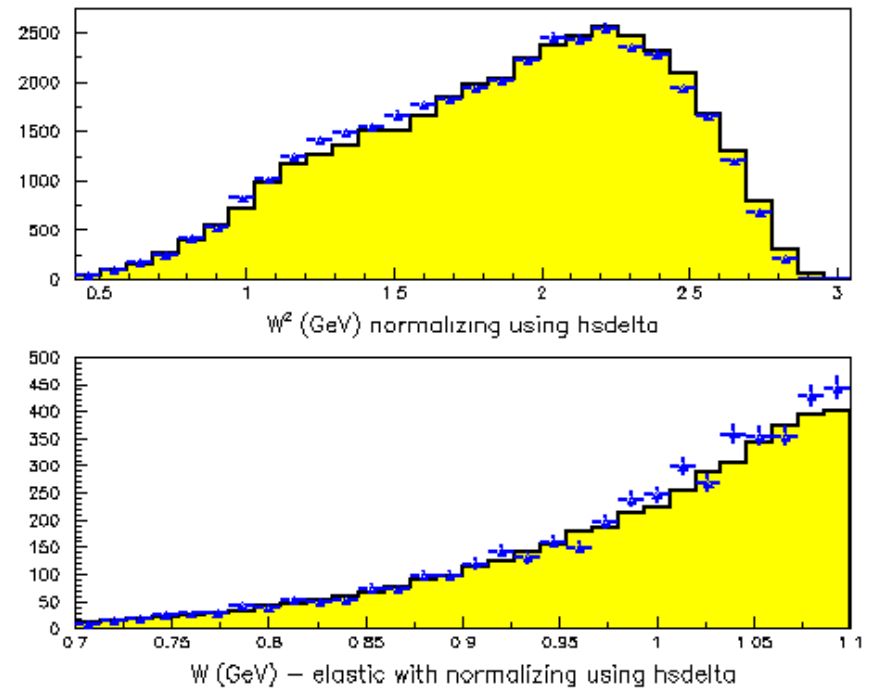
Run = 72782, Target = C

$E = 5.895$ ,  $E' = 4.3943$ ,  $\Theta = 15.41$



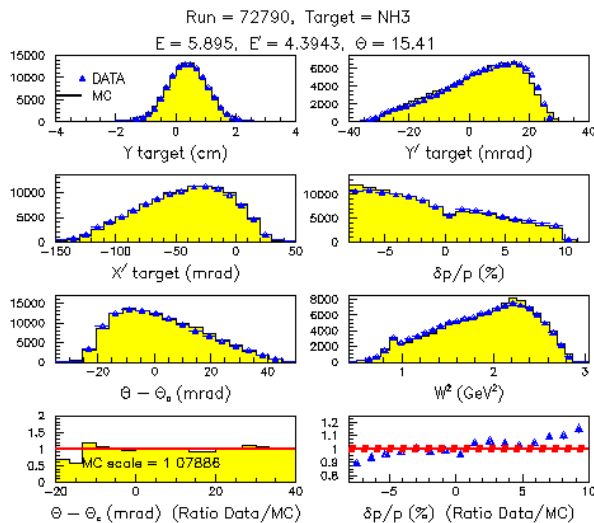
Srast x offset=-0.4 cm

Srast y offset=0.1 cm



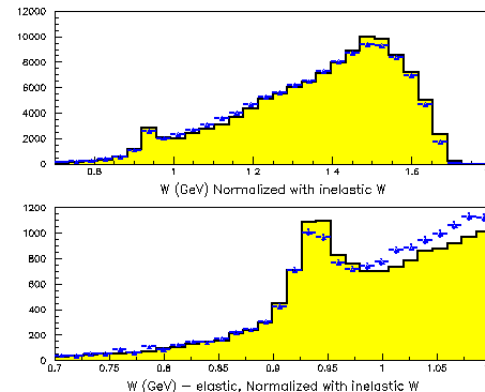
# MC with NH3

- Generated N, H and He separately.
- Added Al come from target end caps and 4K shields as well.
- Calculated the MC scale factor using the data/MC luminosity ratio for each target type.
- Added all targets together by weighting the above MC scale factors.
- Used 60% packing fraction.
- Adjust acceptance edges in Ytar and yptar from adjusting the horizontal beam position.
- Adjust the vertical beam position to bring the W peak to 0.938 GeV



$$srast_x = -0.40 \text{ cm}$$

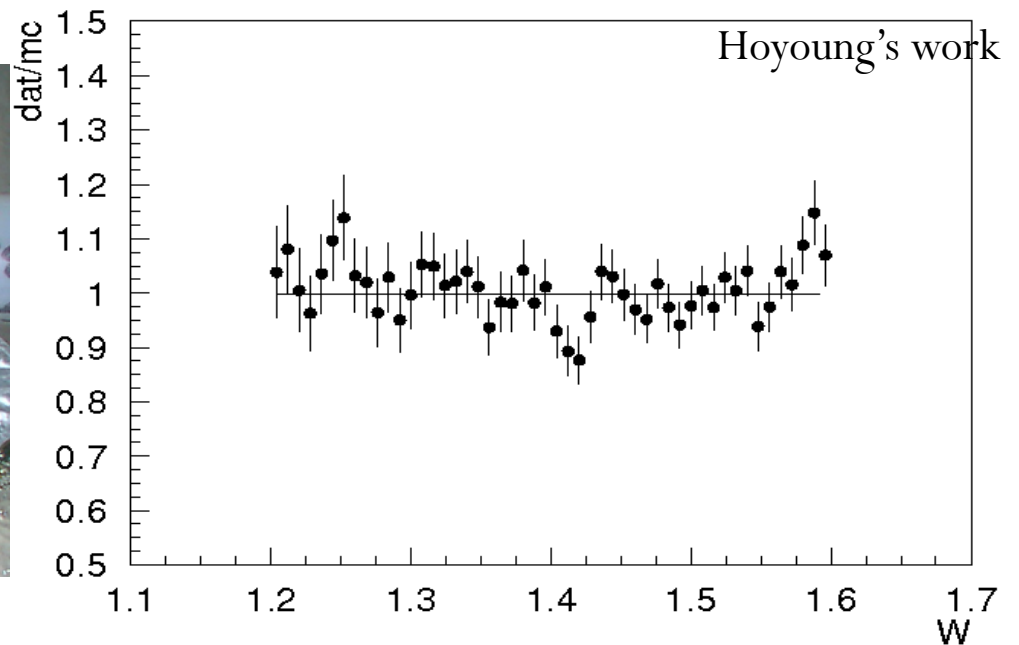
$$srast_y = 0.10 \text{ cm}$$





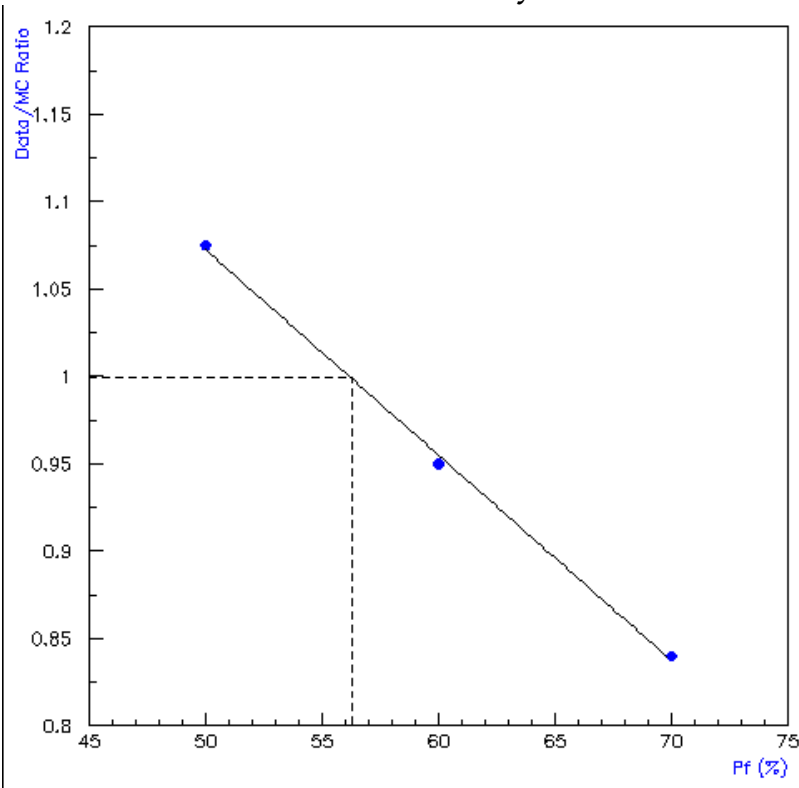
# Packing Fraction.

- Packing Fraction is the actual amount of target material used.
- Determined by taking the ratio of data to MC as a function of  $W$ .
- Need to determine the packing fractions for each of the  $\text{NH}_3$  loads used during the data taking.



# Determine the Packing Fraction

- Looked data to SIMC comparison for the NH<sub>3</sub> target for 3 different Packing Fractions.
- Normalized MC\_NH<sub>3</sub> by 0.93 which is the factor that brings C data/MC ratio to 1.



- Determined the packing fraction which brings Data/MC ratio to 1 from the plot.
- Packing Fraction=56.3 %

Pf (%)	50	60	70
Data/MC Ratio	1.00	0.88	0.78
Data/MC Ratio/0.93	1.075	0.95	0.84

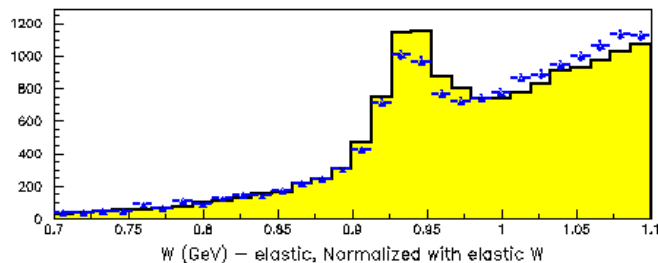
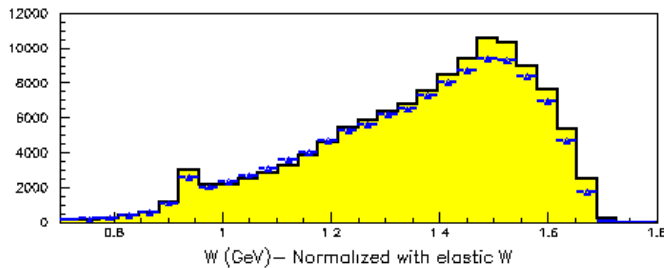
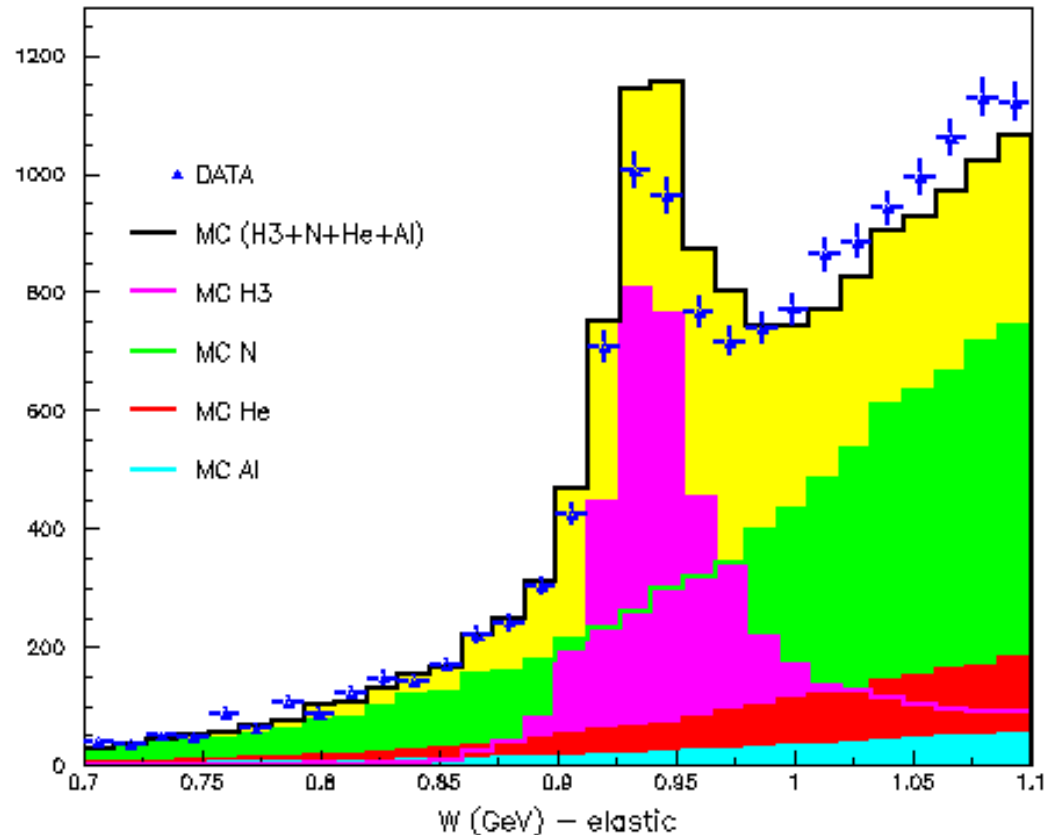
# Determination of the Dilution Factor

## What is the Dilution Factor ?

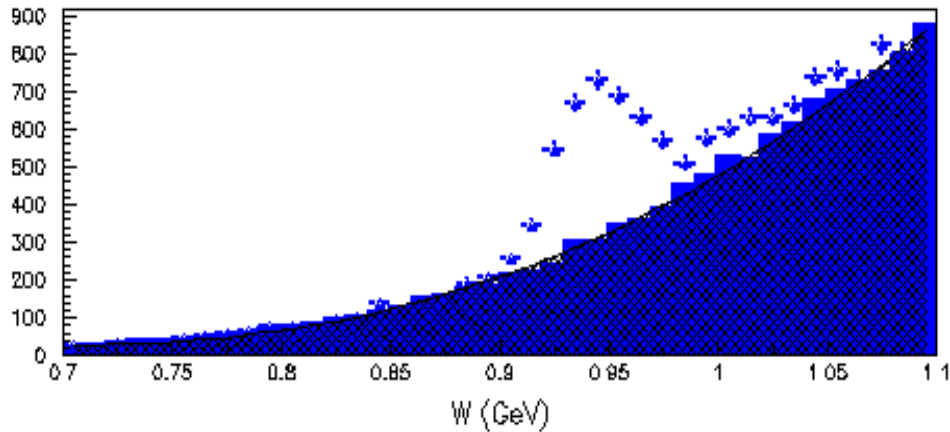
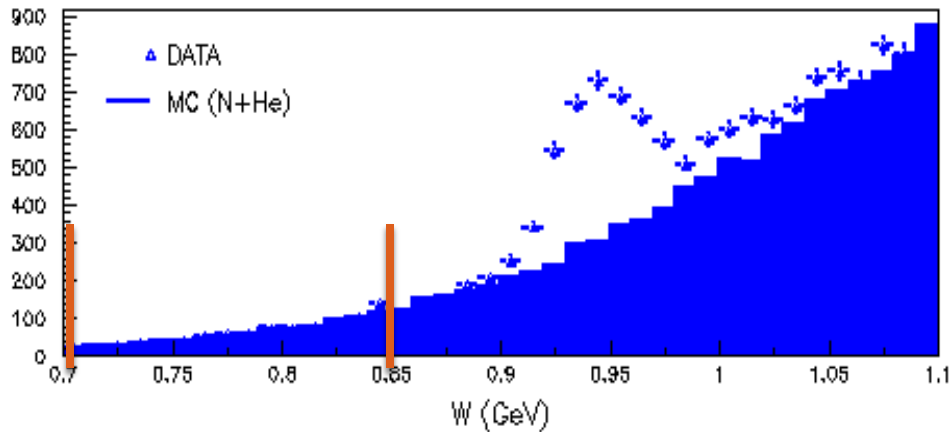
The dilution factor is the ratio of the yield from scattering off free protons (protons from H in NH<sub>3</sub>) to that from the entire target (protons from N, H, He and Al)

Dilution Factor,

$$F = \frac{Yield_{Data} - Yield_{MC(N+He)}}{Yield_{Data}}$$



- MC Background contributions (Only He+N+Al)



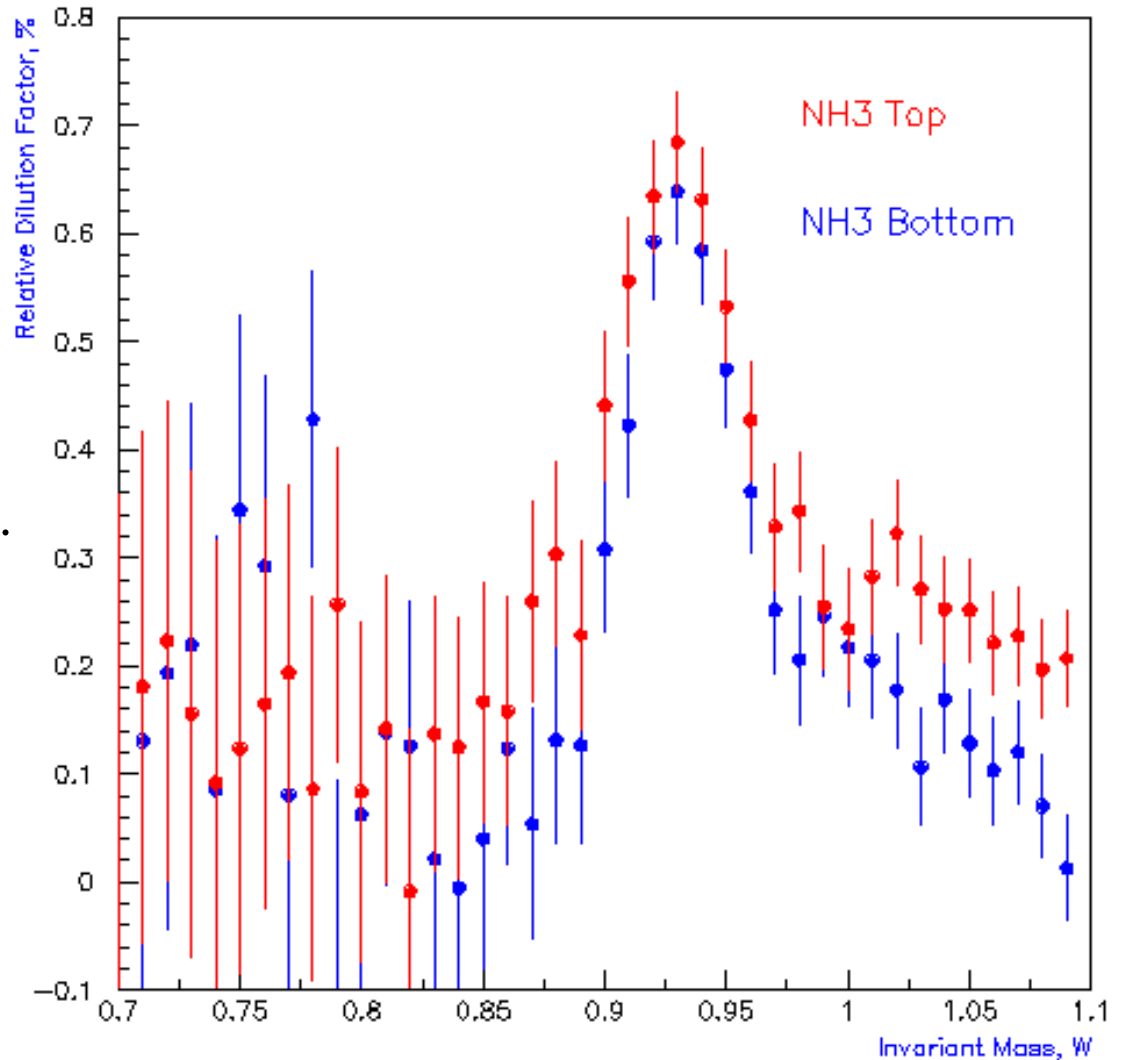
- Calculate the ratio of  $\text{Yield}_{\text{Data}} / \text{Yield}_{\text{MC}}$  for the W region  $0.7 < W < 0.85$  and MC is normalized with this new scaling factor.
- Used the polynomial fit to N+ He+Al in MC and
- Subtract the fit function from data

# ▪ The relative Dilution Factor (**Preliminary**)

*Dilution Factor,*

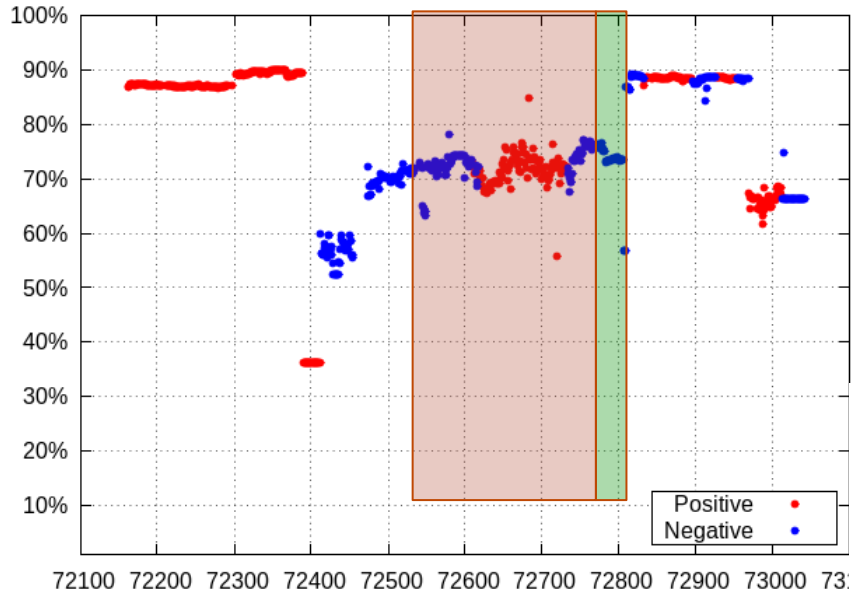
$$F = \frac{Yield_{Data} - Yield_{MC(N+He)}}{Yield_{Data}}$$

- We have taken data using both NH3 targets, called NH3 top and NH3 bottom.
- NH3 crystals are not uniformly filled in each targets which arise two different packing fractions and hence two different dilution factors.



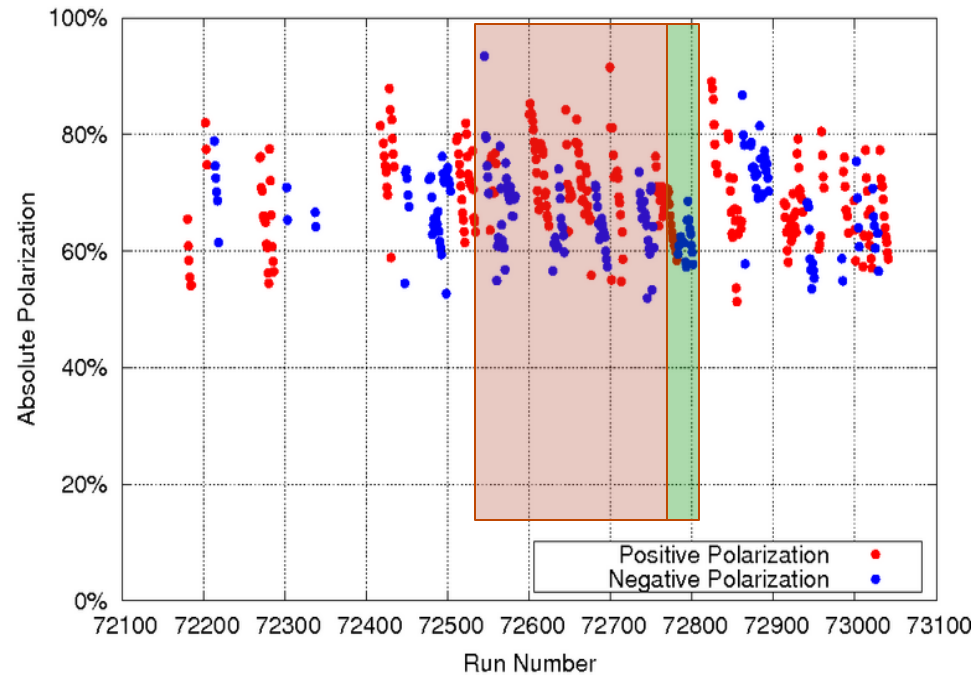
# Beam / Target Polarizations

SANE Beam Polarization Per Run

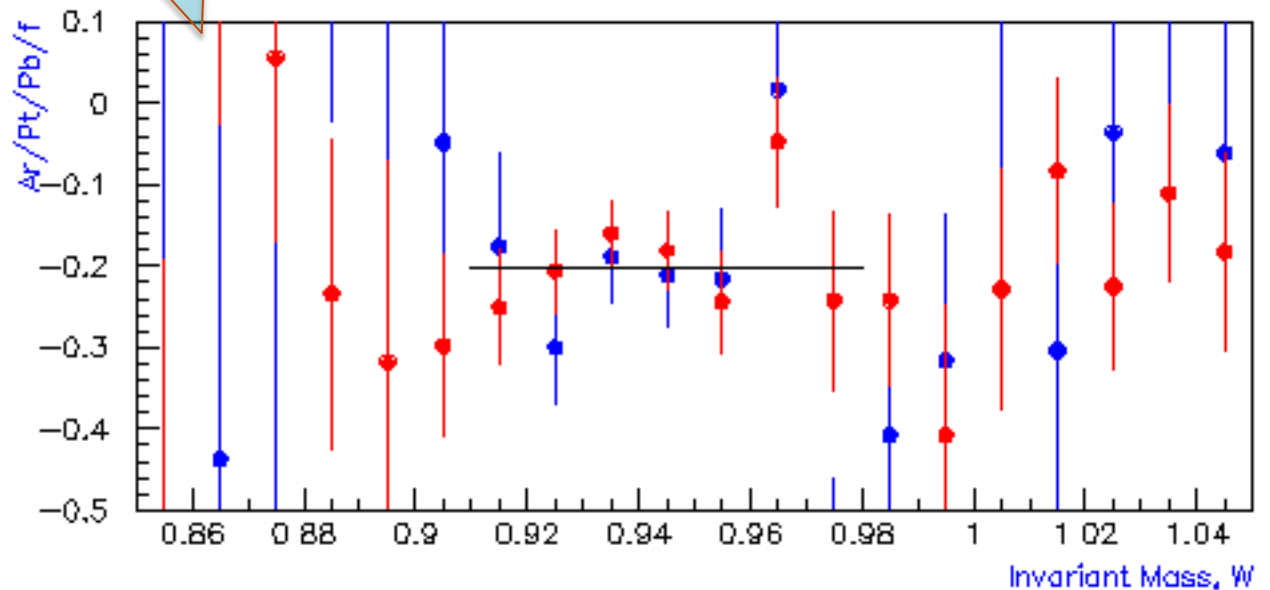
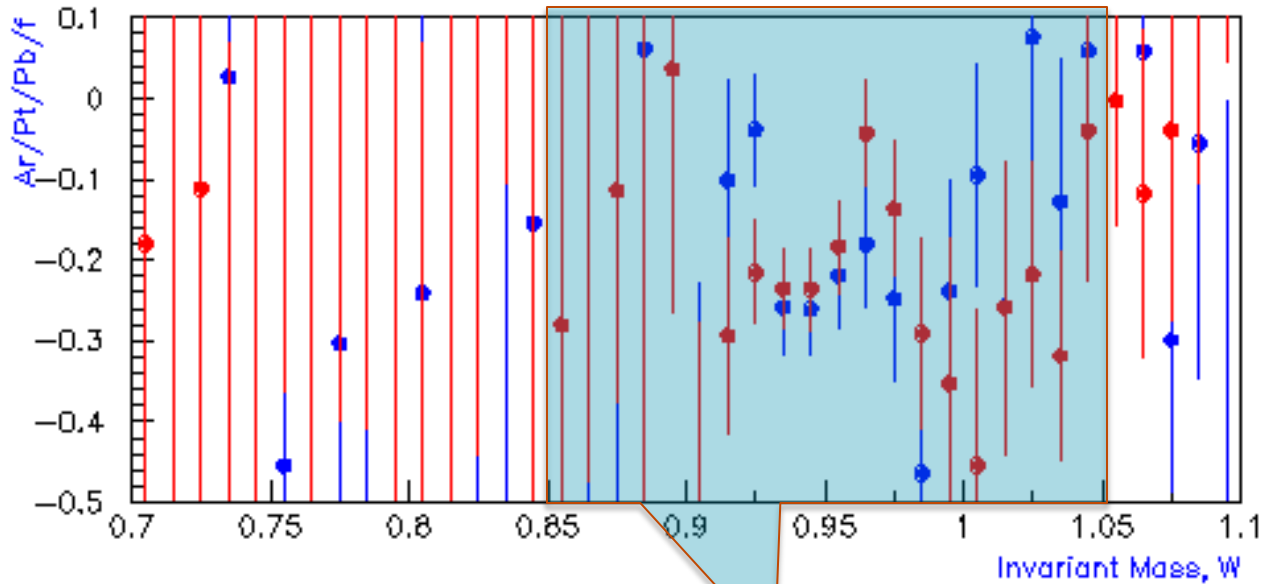


COIN data  
Single arm electron data

Absolute Target Polarization for All SANE Runs



# The Physics Asymmetry (Preliminary)



$A_{\text{phy}}$	Error $A_{\text{phy}}$
-0.201	0.0174

- The beam - target asymmetry,  $A_p$

$$A_p = \frac{-br \sin \theta^* \cos \phi^* - a \cos \theta^*}{r^2 + c}$$

$$\frac{G_E}{G_M} = -\frac{b}{2A_p} \sin \theta^* \cos \phi^* + \sqrt{\frac{b^2}{4A_p^2} \sin^2 \theta^* \cos^2 \phi^* - \frac{a}{A_p} \cos \theta^* - c}$$

Using the experiment data at

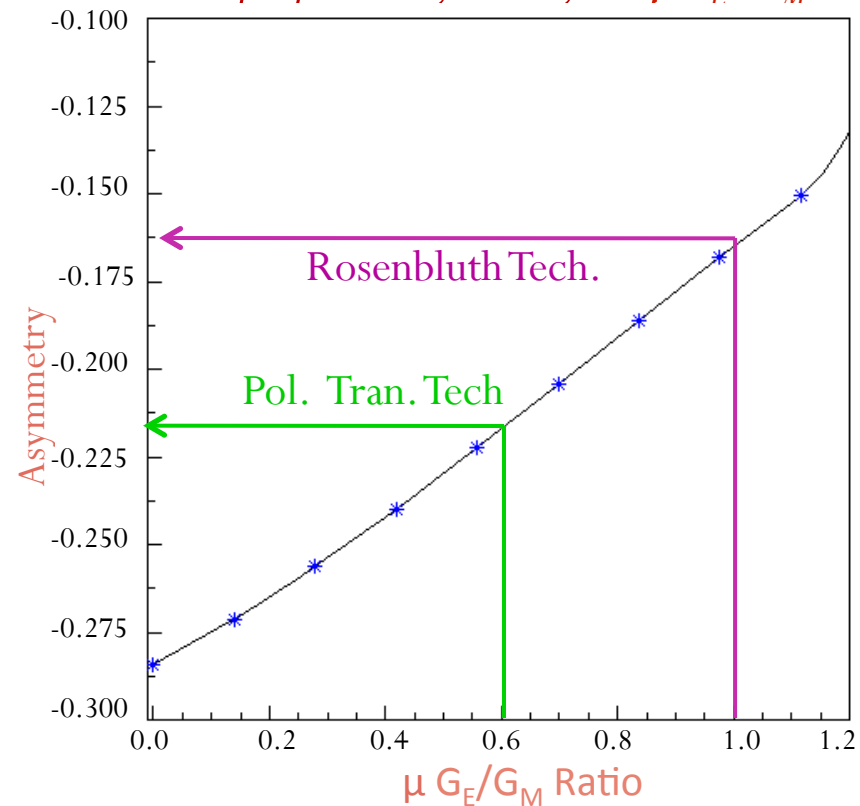
$$Q^2 = 2.2 \text{ (GeV/c)}^2$$

$$\theta^* \approx 34.55^\circ \text{ and } \phi^* = 180^\circ$$

From the HMS kinematics,  $r^2 \ll c$

$$A_p = \frac{-b \sin \theta^* \cos \phi^* r}{c} - \frac{a \cos \theta^*}{c}$$

The projected asymmetry vs  $\mu G_E/G_M$





Using the experiment data at  $Q^2=2.2 \text{ (GeV/c)}^2$  and by knowing the  $A_p=-0.201$ ,

$$r = \left( \frac{G_E}{G_M} \right) = 0.2416$$

$$\mu r = \mu \left( \frac{G_E}{G_M} \right) = 0.674$$

Where,  $\mu$  – Magnetic Moment of the Proton=2.79

## ■ Error propagation from the experiment

$$A_p = \frac{-b \sin \theta^* \cos \phi^* r}{c} - \frac{a \cos \theta^*}{c}$$

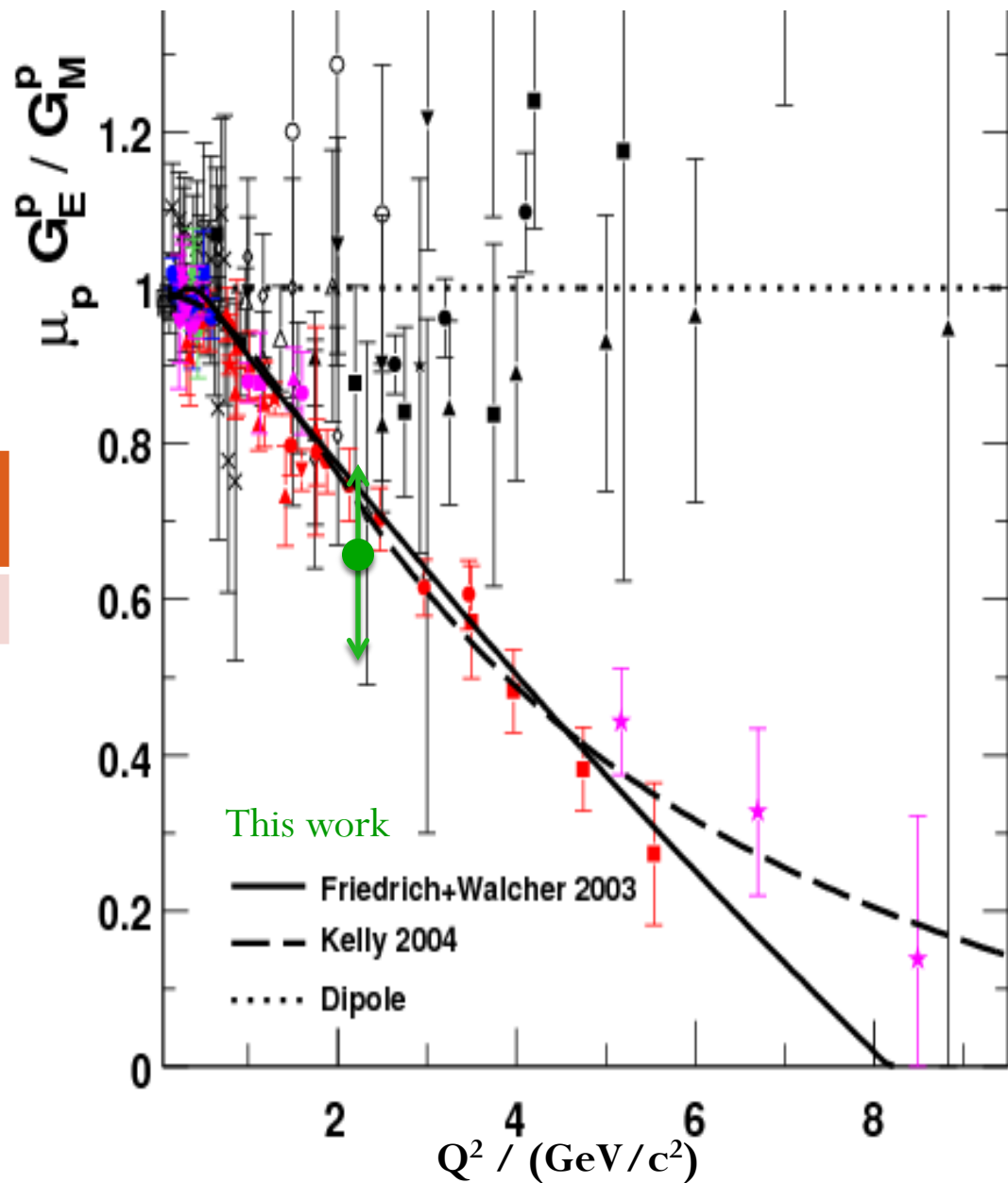
$$\Delta r = \Delta \left( \frac{G_E}{G_M} \right) = \left| \frac{c}{b \sin \theta^* \cos \phi^*} \right| \Delta A_p$$

By knowing the  $\Delta A_p=0.017$ ,

$$\Delta(\mu r) = \Delta \left( \mu \frac{G_E}{G_M} \right) = 0.13$$

# Preliminary .....

$\mu G_E/G_M$	$\Delta(\mu G_E/G_M)$
0.674	0.13



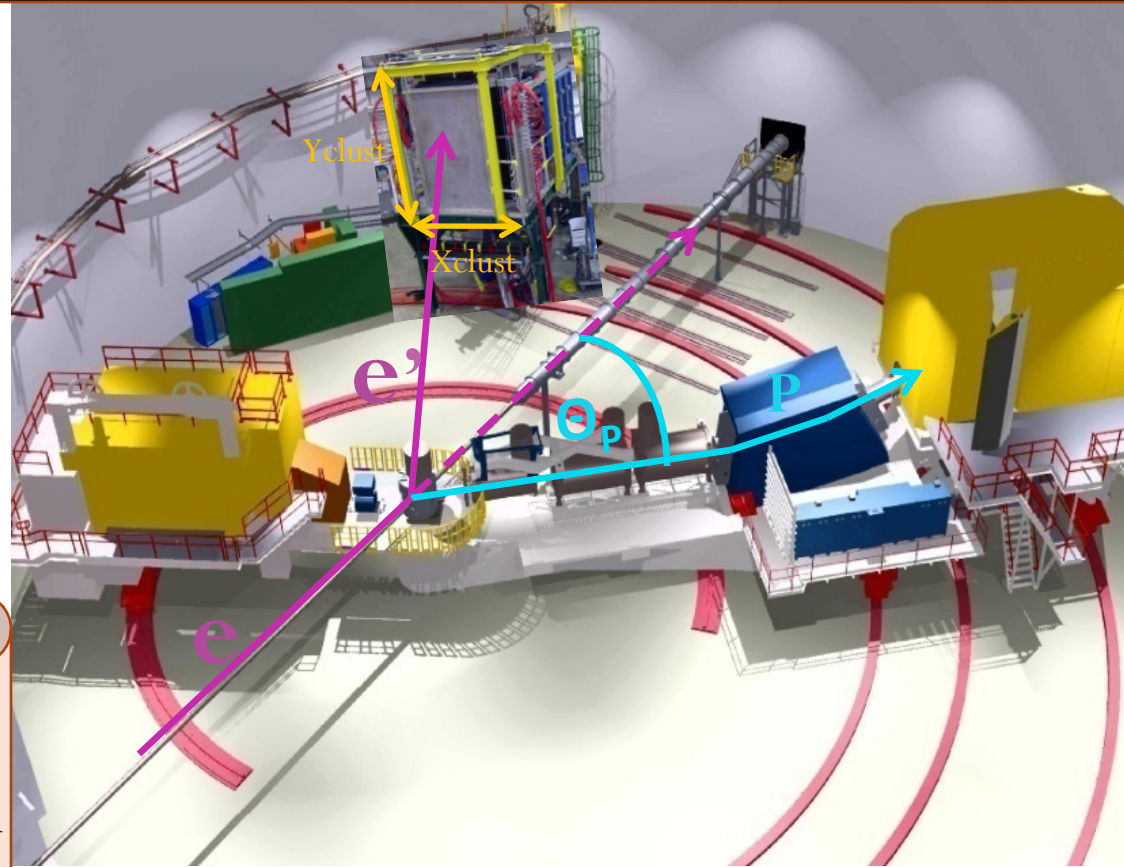
# Coincidence Data (Electrons in BETA and Protons in HMS)

## Definitions :

$X/Y_{clust}$  - Measured  $X/Y$  positions on the BigCal

- $X$  = horizontal / in-plane coordinate
- $Y$  = vertical / out-of-plane coordinate

$E_{clust}$  - Measured electron energy at the BigCal



By knowing  
the energy of the polarized electron  
beam,  $E_B$   
and  
the scattered proton angle,  $\Theta_p$



We can predict the

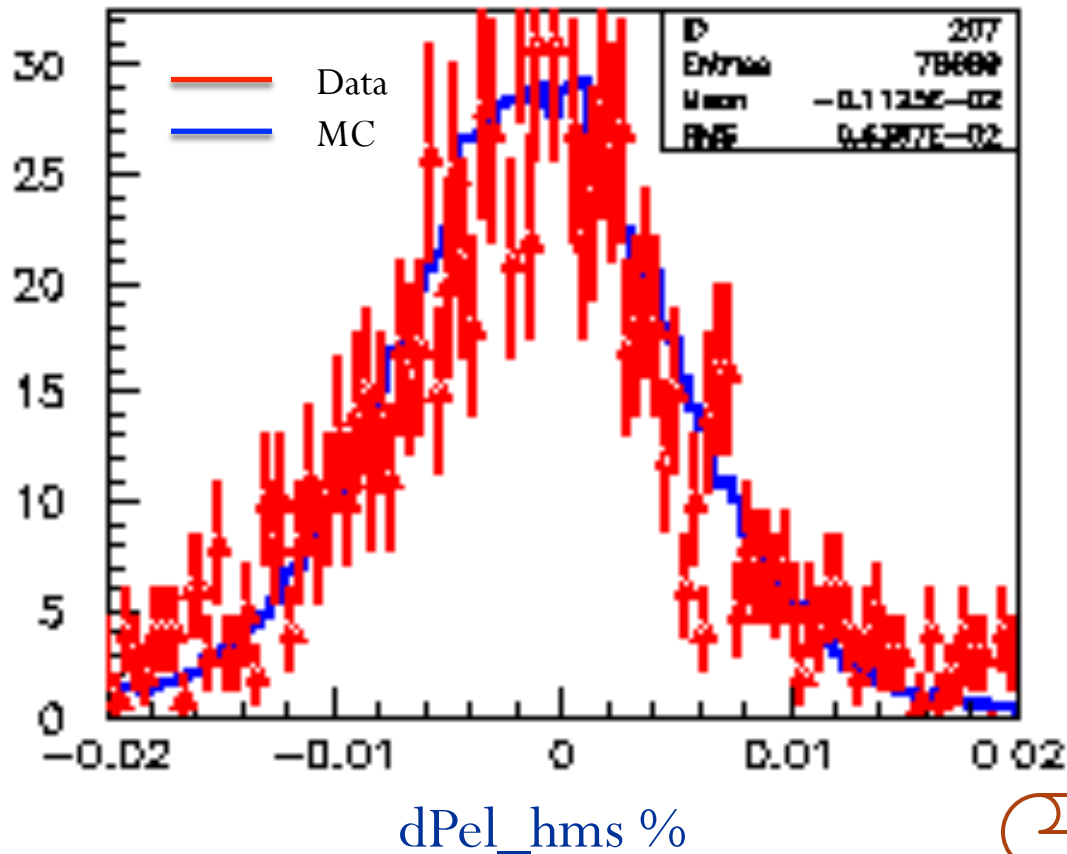
- $X/Y$  coordinates -  $X_{HMS}$ ,  $Y_{HMS}$  and  
(Target Magnetic Field Corrected)
- The Energy -  $E_{HMS}$   
of the coincidence electron on the BigCal

# Elastic Kinematics

( From HMS Spectrometer )

Spectrometer mode	Coincidence	Coincidence	Single Arm
HMS Detects	Proton	Proton	Electron
E Beam GeV	4.72	5.89	5.89
$P_{\text{HMS}}$ GeV/c	3.58	4.17	4.40
$\Theta_{\text{HMS}}$ (Deg)	22.30	22.00	15.40
$Q^2$ (GeV/c) <sup>2</sup>	5.17	6.26	2.20
Total Hours (h)	~40 (~44 runs)	~155 (~135 runs)	~12 (~15 runs)
e-p Events	~113	~1200	-

# Fractional momentum difference



$$dPel\_hms = \frac{P_{HMS} - P_{Cal}}{P_{cent}}$$

$$P_{Cal} = \sqrt{v^2 + 2Mv}$$

$$v = \frac{Q^2}{2M}$$

$$Q^2 = \frac{4M^2 E^2 \cos^2 \theta}{M^2 + 2ME + E^2 \sin^2 \theta}$$

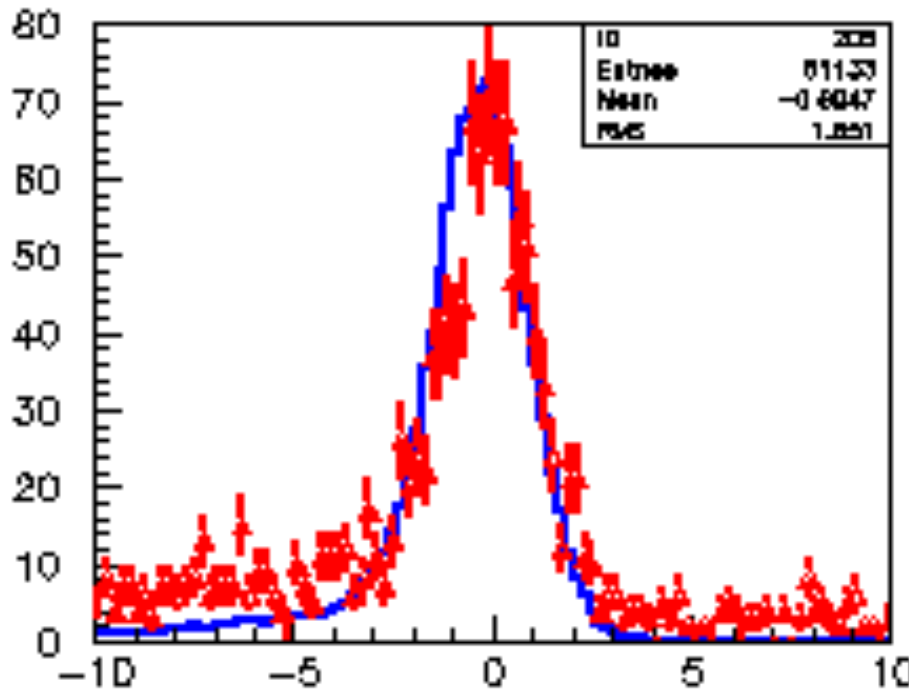
$P_{HMS}$  – Measured proton momentum by HMS

$P_{cal}$  - Calculated proton momentum by knowing the beam energy,  $E$  and the proton angle,  $\theta$

$P_{cent}$  – HMS central momentum

# X/Y position difference

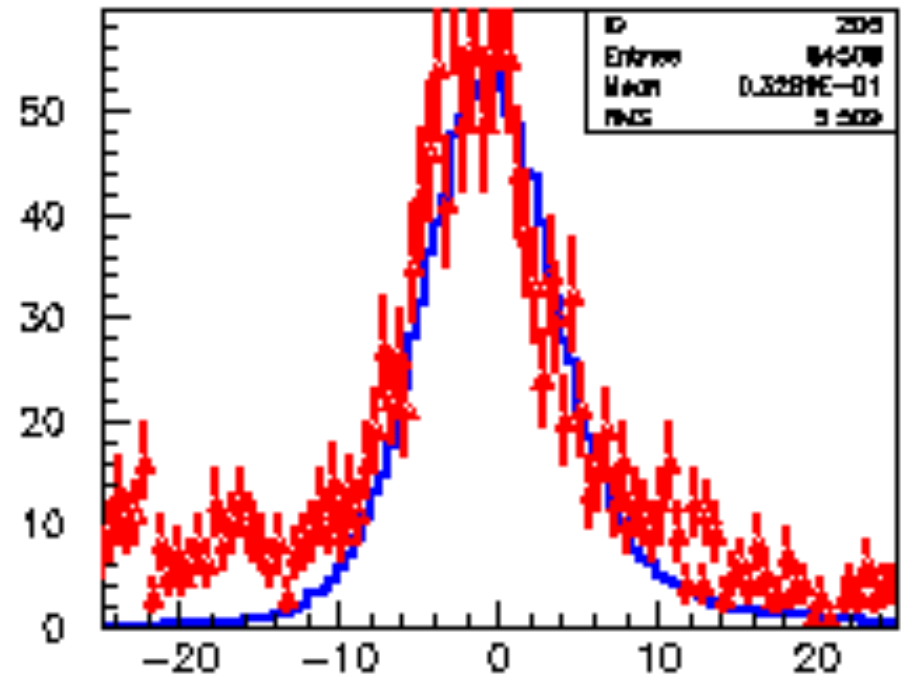
*X position difference*



$X_{HMS} - X_{clust} / \text{cm}$

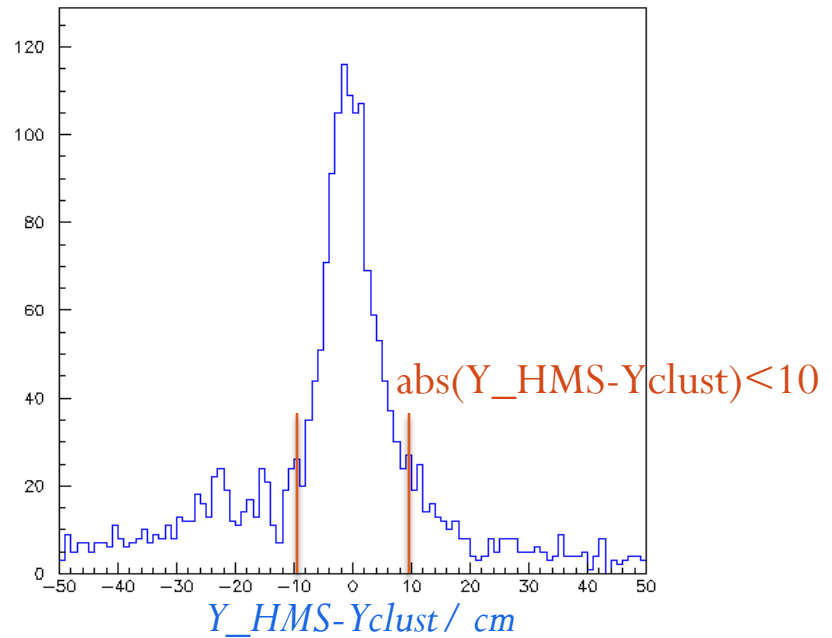
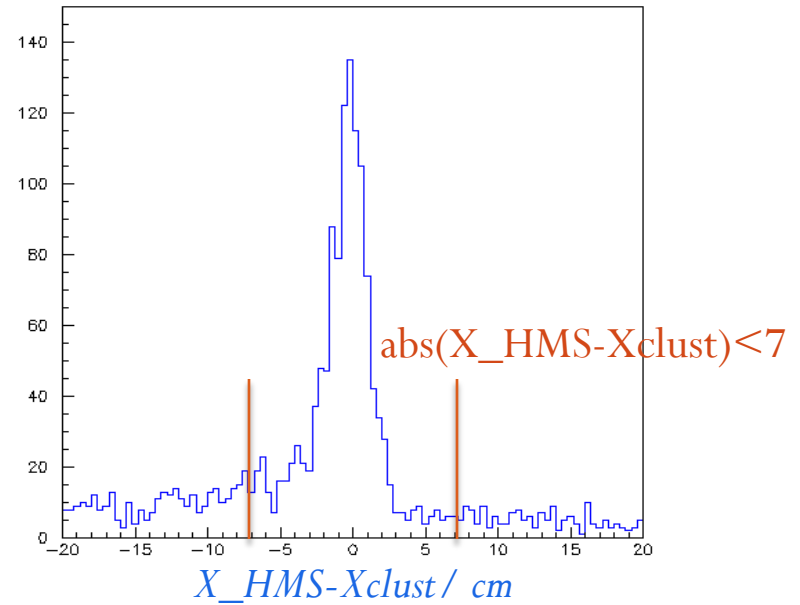
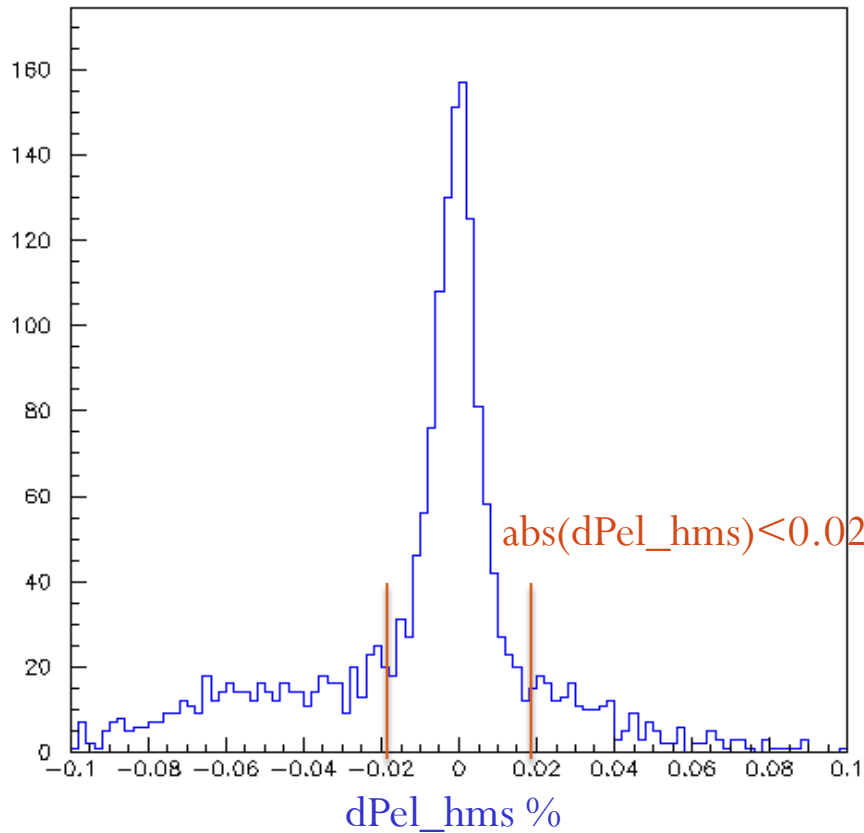
— Data  
— MC

*Y position difference*



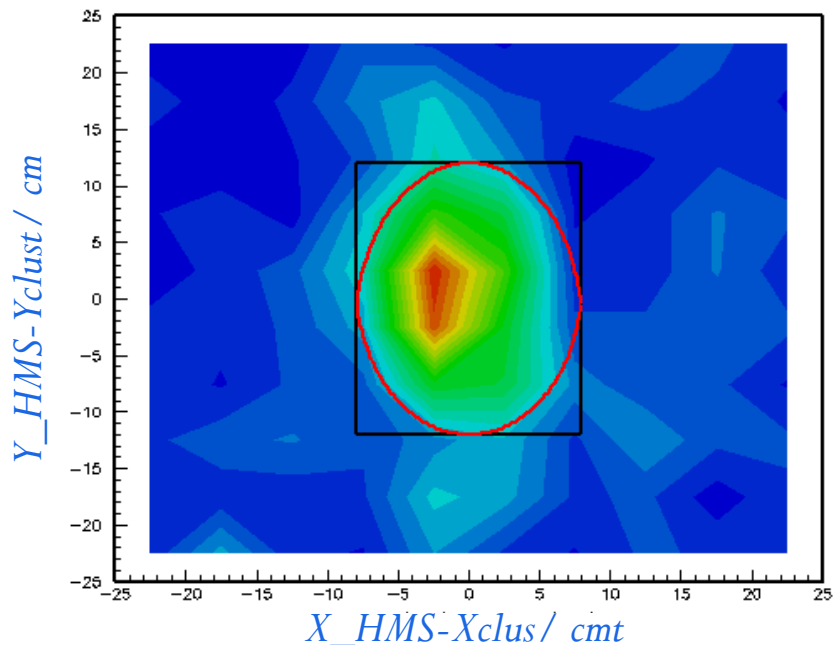
$Y_{HMS} - Y_{clust} / \text{cm}$

# Applied the coincidence cuts

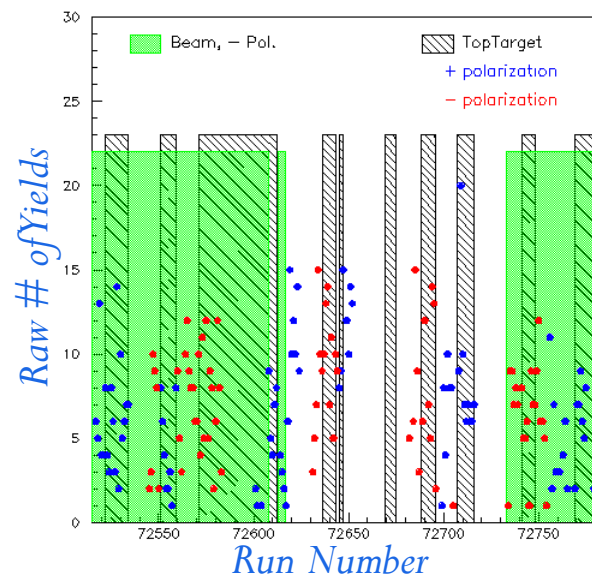
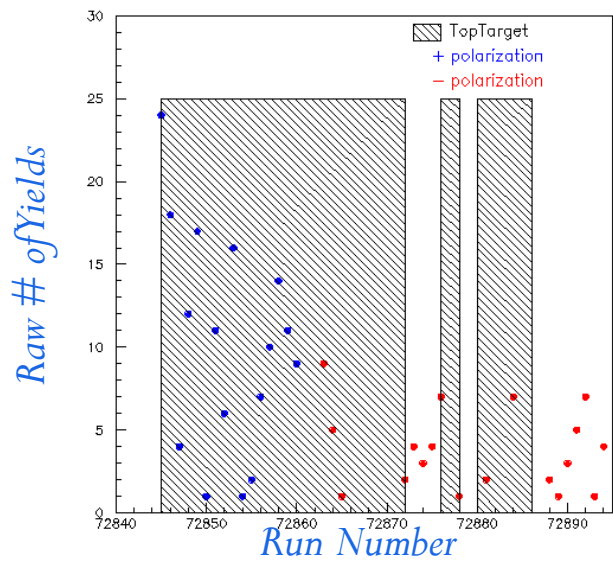
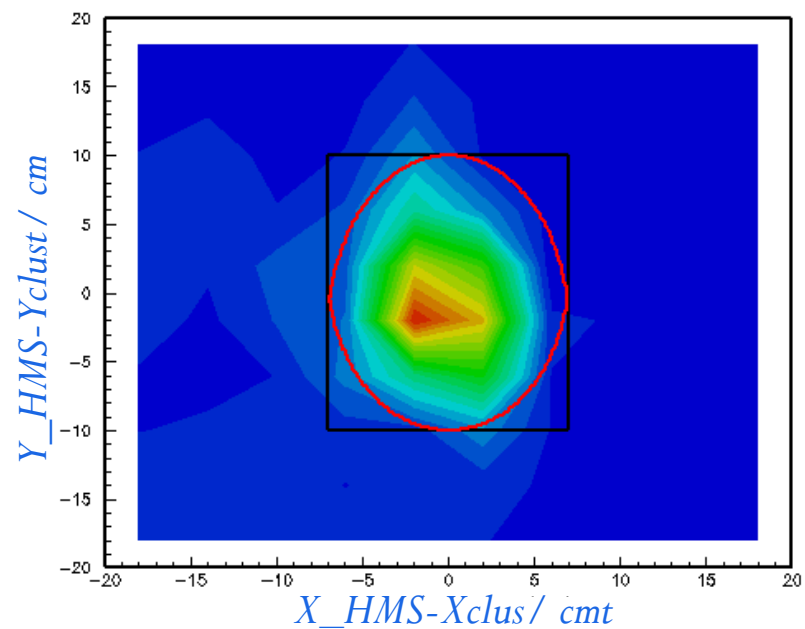


# Elastic Events

4.72 GeV data

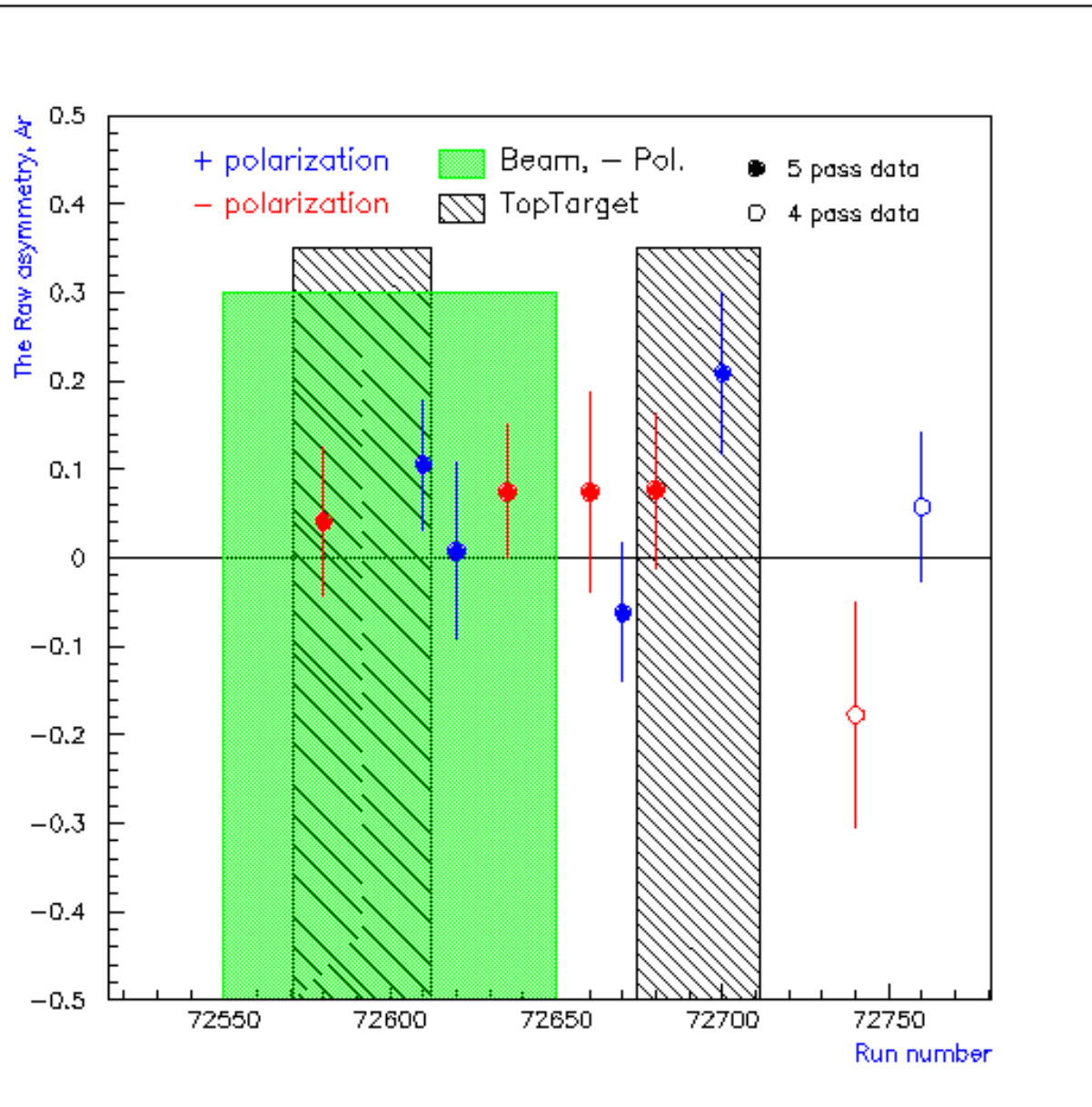


5.89 GeV data





# Extract the Raw Asymmetries



- Raw yields are normalized with
- Total Charge
  - charge average +/- life times

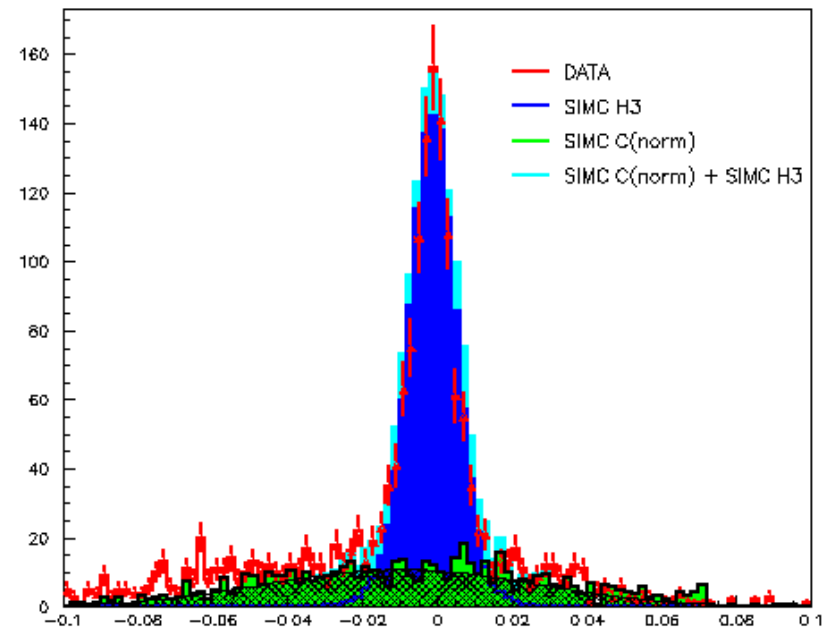
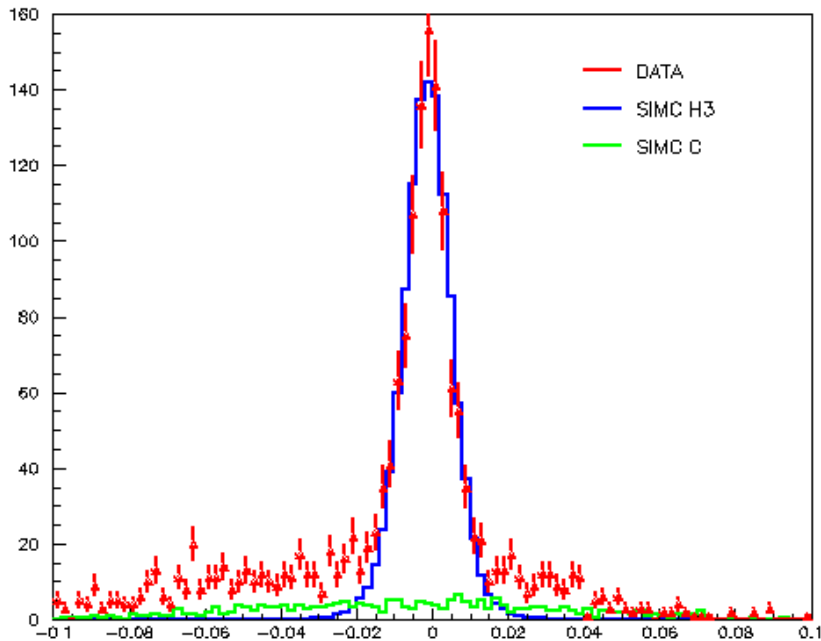
*Need  
dilution factor,  $f$   
in order to determine the  
physics asymmetry,*

$$A_p = \frac{A_r}{fP_B P_T} + N_C$$

*and  $G^P_E / G^P_M$*

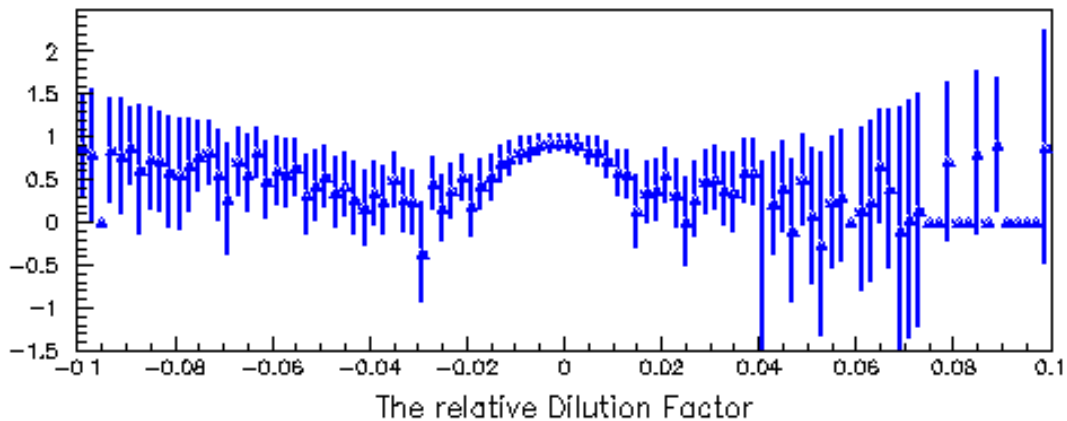
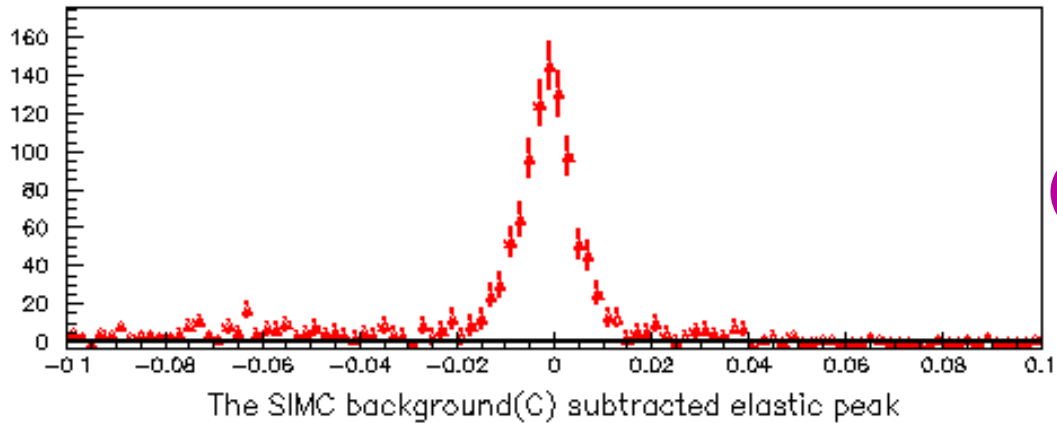
# Determine The Dilution Factor

- Estimate The Background



- Get the ratio of data/SIMC\_C for the region of  $0.03 < \text{dpel\_hms} < 0.08$ . (ratio=2.73893)
- Normalized the SIMC\_C with that ratio (2.73893) for the region of  $-0.1 < \text{dpel\_hms} < 0.1$  and added SIMC\_H3 to it. Compare with the data.  
$$\text{Data}/\text{SIMC}(\text{H3}+2.73893*\text{C}) = 0.991536$$
- Used the Gaussian fit for the SIMC\_C (normalized with 2.73893) and subtract it from the data
- Get the relative dilution factor by taking the ratio of SIMC\_C subtracted data to data.  
the relative df. =  $(\text{data}-\text{SIMC}_C)/\text{data}$

- Get The Relative Dilution Factor



*Two different target cups  
(NH<sub>3</sub> Top and NH<sub>3</sub> Bottom)*



*Two different packing  
fractions*

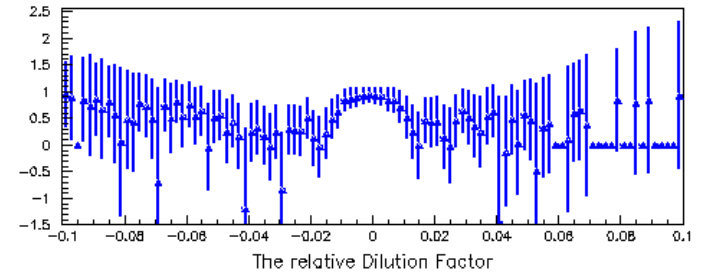
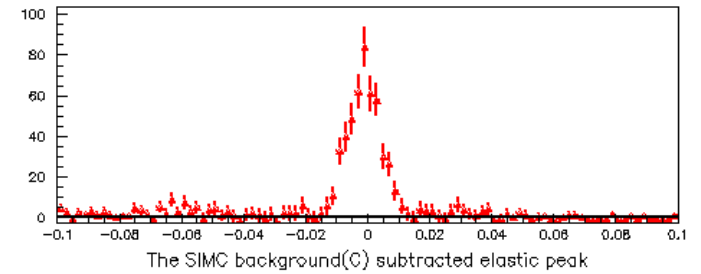
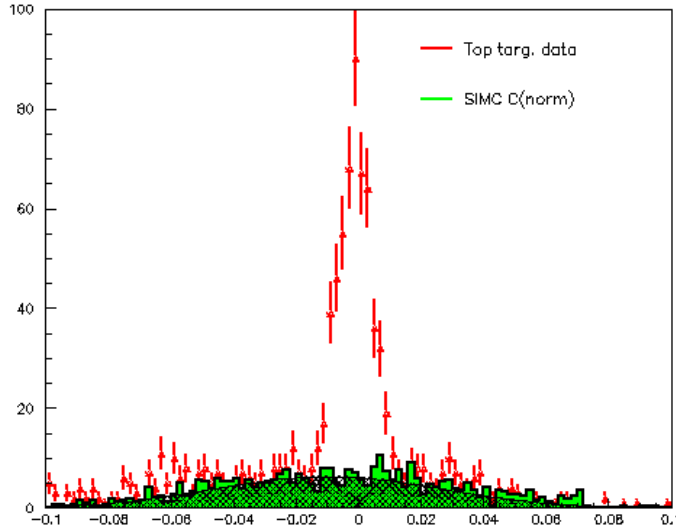


*Need*

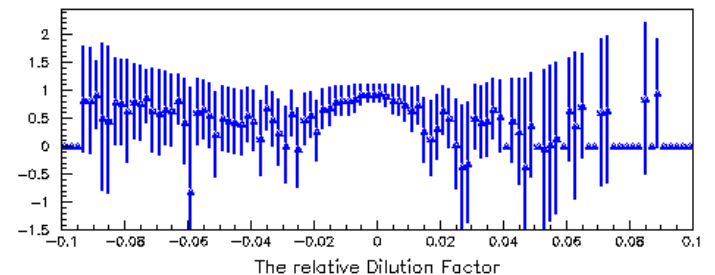
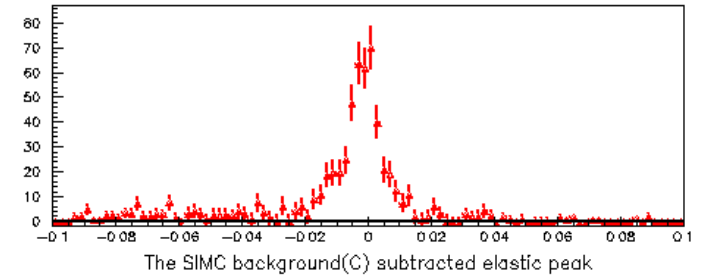
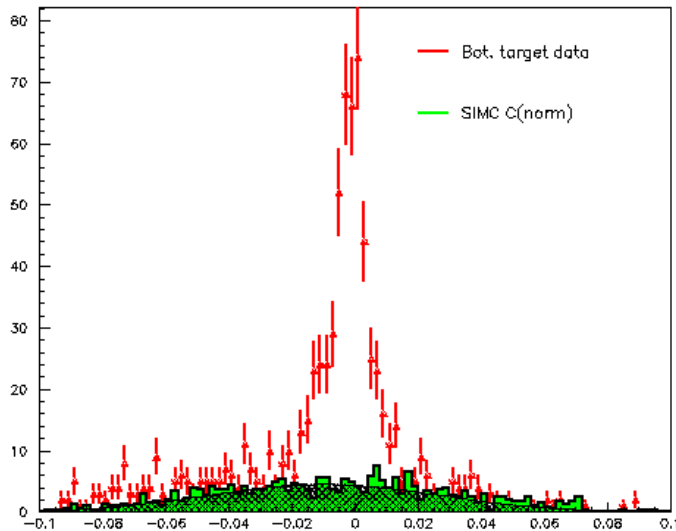
*Two different dilution  
factors*

- The Relative Dilution Factors For

Top Target



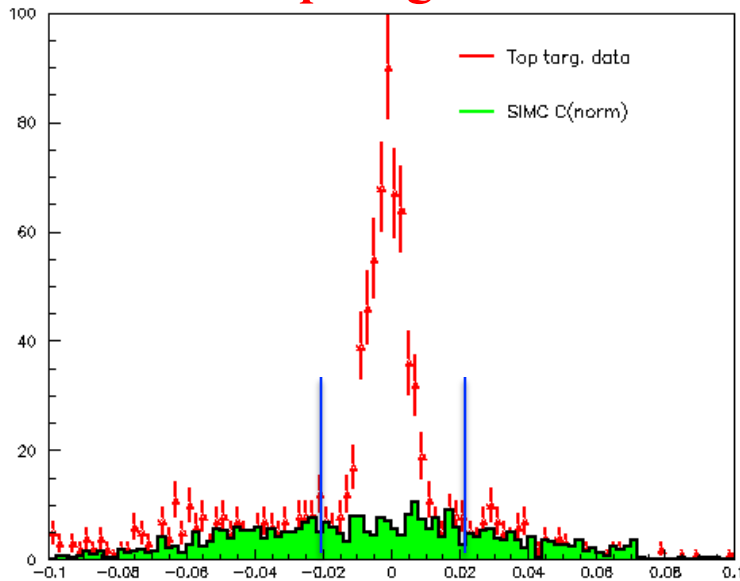
Bottom Target



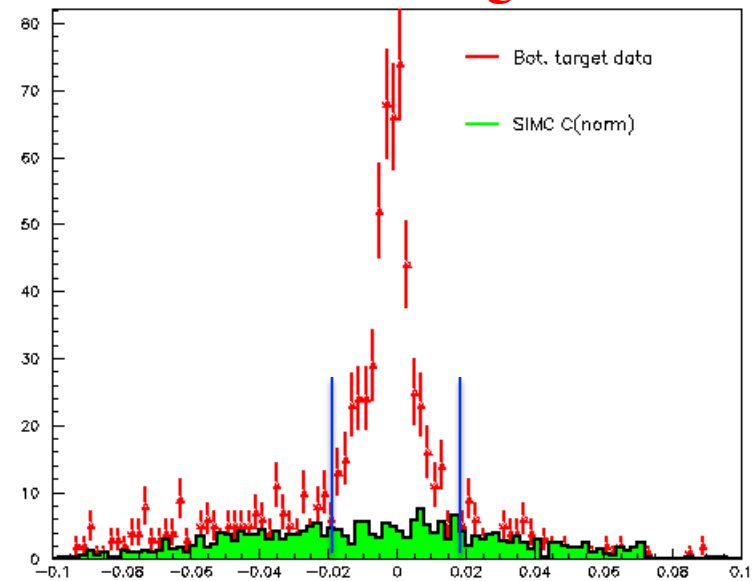
- The Relative Dilution Factor  
(Used the Integration Method)

- Because of the law statistics, It is hard to correct the raw asymmetry for the df as a function of dpel\_hms
- Just integrate over the dpel\_hms region of +/- 0.02 for the top and bottom.

**Top Target**



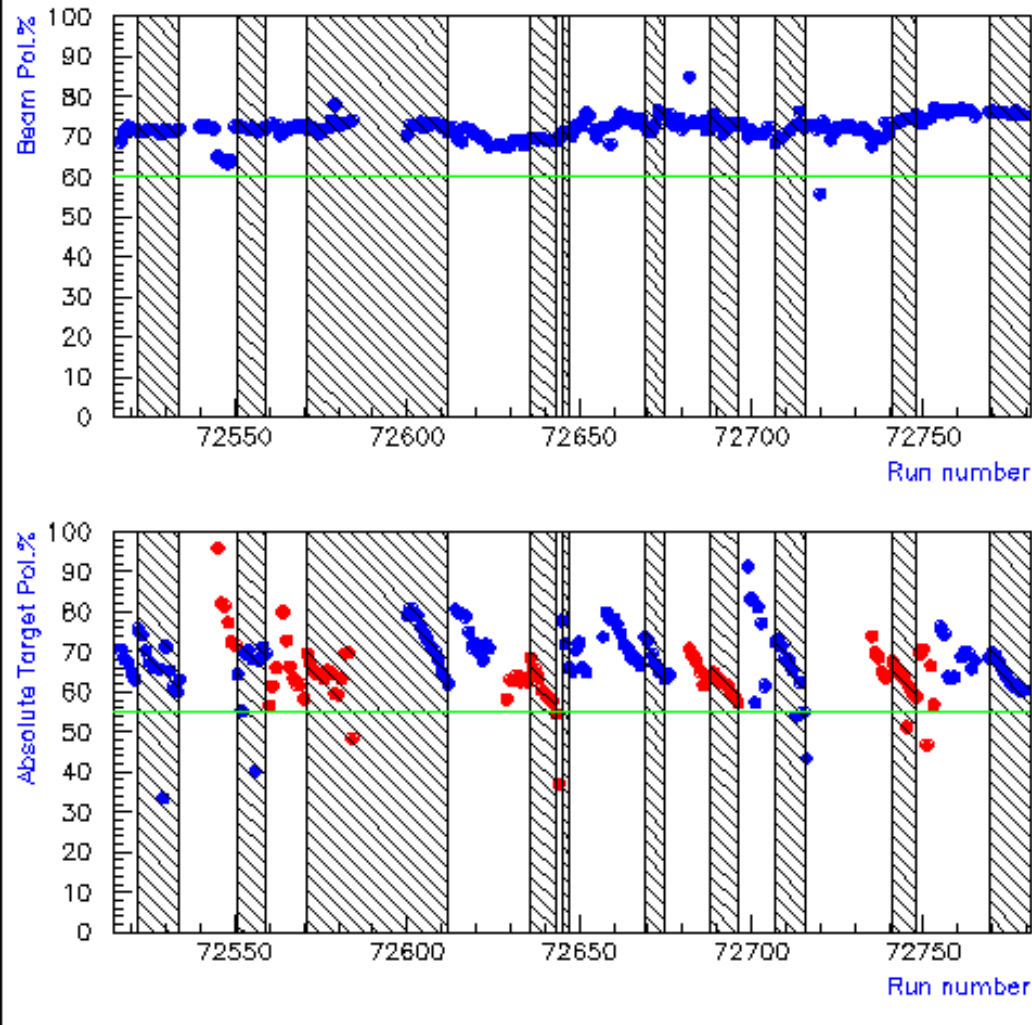
**Bottom Target**



$$\begin{aligned} \text{The relative D.F} &= (\text{data-SIMC\_C})_{\text{top}}/\text{data}_{\text{top}} \\ &= 606-130/606 \\ &= 0.785 \end{aligned}$$

$$\begin{aligned} &= (\text{data-SIMC\_C})_{\text{bot}}/\text{data}_{\text{bot}} \\ &= 606-130/606 \\ &= 0.785 \end{aligned}$$

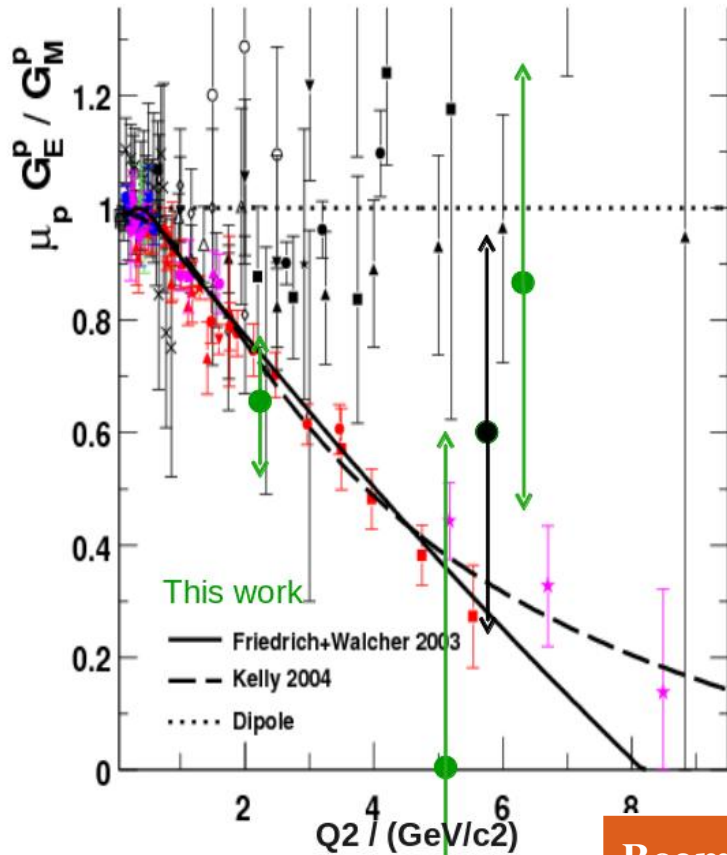
# Beam and Target Polarizations



- Used the runs of beam polarization  $> 60\%$  and  $\text{abs}(\text{target polarization}) > 55\%$
- Used the charge average target and beam polarizations to calculate the physics asymmetries



# Extract the Form Factor Ratio, $G_E/G_M$



Beam(GeV)	4.72	5.89	Weighted Avg.
$Q^2$ (GeV/c) <sup>2</sup>	5.17	6.26	5.72
$\mu G_E/G_M$	-0.032	0.875	0.6145
$\Delta(\mu G_E/G_M)$	0.668	0.424	0.358



## To Do

- Determine the new dilution factor, raw / physics asymmetries and hence the form factor ratio,  $G_E / G_M$  using the new packing fraction of 56.3% for the single arm electron data.
- Estimate the systematic errors for both single arm electron and coincidence data

# Conclusion

- Measurement of the beam-target asymmetry in elastic electron-proton scattering offers an independent technique of determining the  $G_E/G_M$  ratio.
- This is an ‘explorative’ measurement, as a by-product of the SANE experiment.
- Extraction of the  $G_E/G_M$  ratio from single-arm electron and Coincidence data are shown.
- The preliminary data point at  $2.2 \text{ (GeV/c)}^2$  is very consistent with the recoil polarization data (falls even slightly below it)
- The preliminary weighted average data point of the coincidence data at  $5.72 \text{ (GeV/c)}^2$  consistent with the recoil polarization data within its  $3 \sigma$  error.

## **SANE Collaborators:**

Argonne National Laboratory, Christopher Newport U., Florida International U., [Hampton U.](#), Thomas Jefferson National Accelerator Facility, Mississippi State U., North Carolina A&T State U., Norfolk S. U., Ohio U., Institute for High Energy Physics, U. of Regina, Rensselaer Polytechnic I., Rutgers U., Seoul National U., State University at New Orleans , Temple U., Tohoku U., U. of New Hampshire, U. of Virginia, College of William and Mary, Xavier University of Louisiana, Yerevan Physics Inst.

**Spokespersons:** S. Choi (Seoul), M. Jones (TJNAF), Z-E. Meziani (Temple),  
O. A. Rondon (UVA)

Thank You

