



Hadron Mass Corrections in semi-inclusive deep-inelastic scattering

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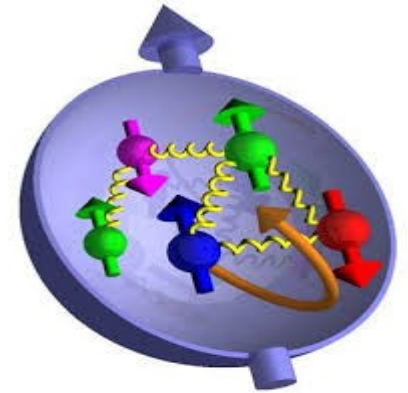
Outline

- Introduction
- Kinematics
- Cross sections at finite Q^2
- Phenomenological implications
- Conclusion and Outlook

Introduction

Hard scattering reactions:

Picture of the nucleon
(partons)



Examples:

Deep inelastic scattering (DIS): $e^-p \longrightarrow e^-X$

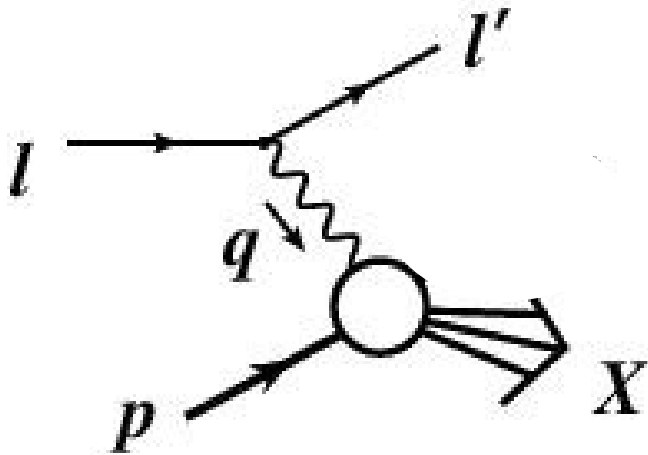
Semi inclusive Deep inelastic scattering (SIDIS): $e^-p \longrightarrow e^-hX$

Drell Yan (DY): $h_A h_B \longrightarrow e^-e^+X$

e^-e^+ annihilation: $e^-e^+ \longrightarrow hX$

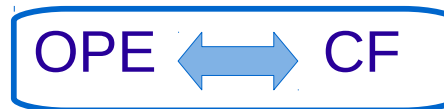
Introduction: DIS

DIS Diagram



Theoretical Framework:

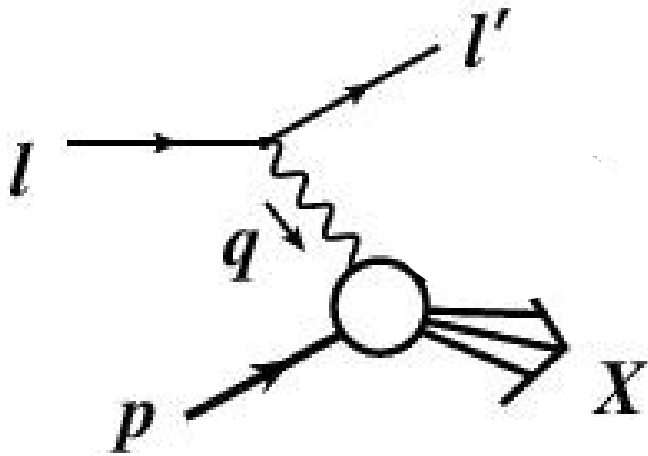
- Operator Product Expansion (OPE).
- Collinear factorization (CF)



To be studied:
massive case

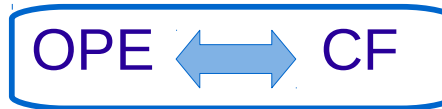
Introduction: DIS

DIS Diagram



Theoretical Framework:

- › Operator Product Expansion (OPE).
- › Collinear factorization (CF)



To be studied:
massive case

DIS Kinematic invariants

Nucleon Mass

Virtuality

$$m_N^2 = (p_\mu)^2$$

$$Q^2 = -(q_\mu)^2$$

Bjorken-x

Final state

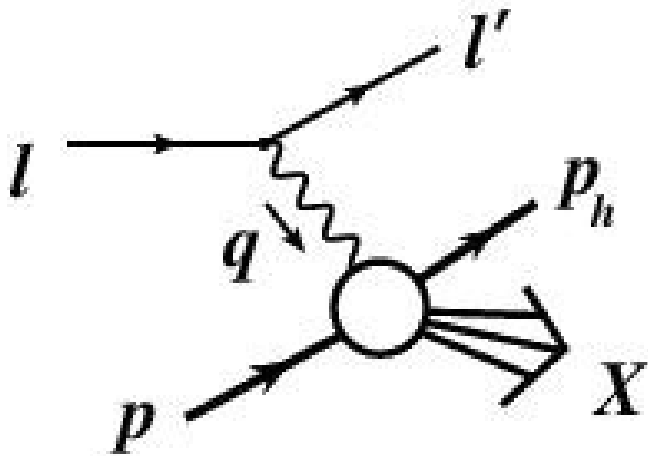
invariant mass

$$x_B = \frac{Q^2}{2p \cdot q}$$

$$W^2 = (p_\mu + q_\mu)^2$$

Introduction: SIDIS

SIDIS Diagram



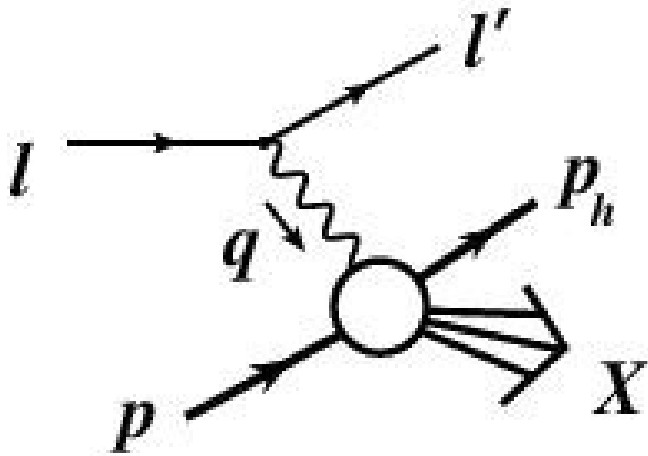
Theoretical Framework:

- Operator Product Expansion (OPE).
- Collinear factorization (CF)

DIS variables

Introduction: SIDIS

SIDIS Diagram



Theoretical Framework:

- Operator Product Expansion (OPE).
- Collinear factorization (CF)

DIS variables + hadronic variables

Hadron Mass Fragmentation
invariant

$$m_h^2 = p_h^2$$

$$z_h = \frac{p_h \cdot p}{q \cdot p}$$

Introduction: Why Hadron Mass Corrections (HMC's)?

Jlab experiments



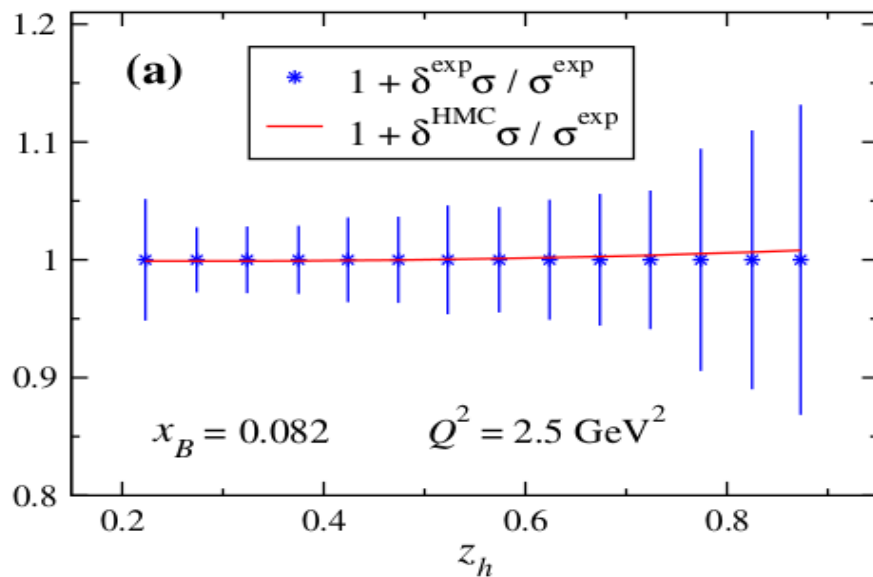
- Usually Low Q^2 ($\sim 1 \text{ GeV}^2$)
- $1/Q^2$ corrections have to be controlled



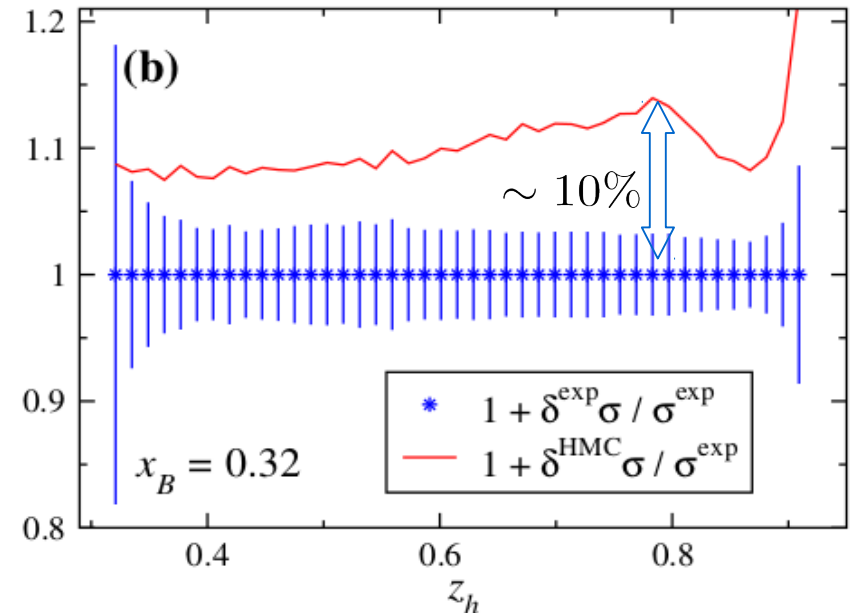
$$O\left(\frac{M^2}{Q^2}\right) \equiv \text{HMC's!!!}$$

Introduction: HMC's for charged pion production (unpolarized SIDIS)

HERMES experiment (DESY)



Jefferson Lab (E00-108)



Accardi, Hobbs, Melnitchouk
JHEP 0911, 084 (2009)

Introduction: Polarized HMC's (Project)

Hadron mass corrections in semi-inclusive deep-inelastic scattering

JHEP 1509 (2015) 169

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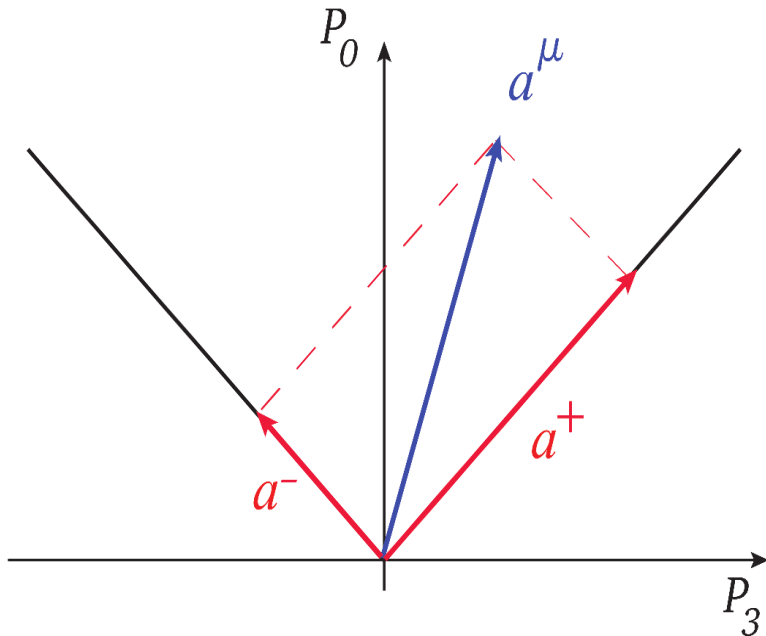
⁴*Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA*

Starting from older JHEP 0911, 084 (2009) we:

- Review all the kinematics.
- Derive the cross section for the unpolarised case and extend to the polarized case.
- Estimate the size of the hadron mass corrections for current and future experiments

SIDIS: Massive scaling variables

Light Cone



$$a^- = \frac{a_0 - a_3}{\sqrt{2}} \quad a^+ = \frac{a_0 + a_3}{\sqrt{2}}$$

Nachtmann Scaling variable

$$\xi = -\frac{q^+}{p^+} = \frac{2x_B}{1 + \sqrt{1 + 4x_B^2 M^2/Q^2}}$$

Bjorken Limit $Q^2 \rightarrow \infty \quad \xi \rightarrow x_B$

Fragmentation Scaling variable

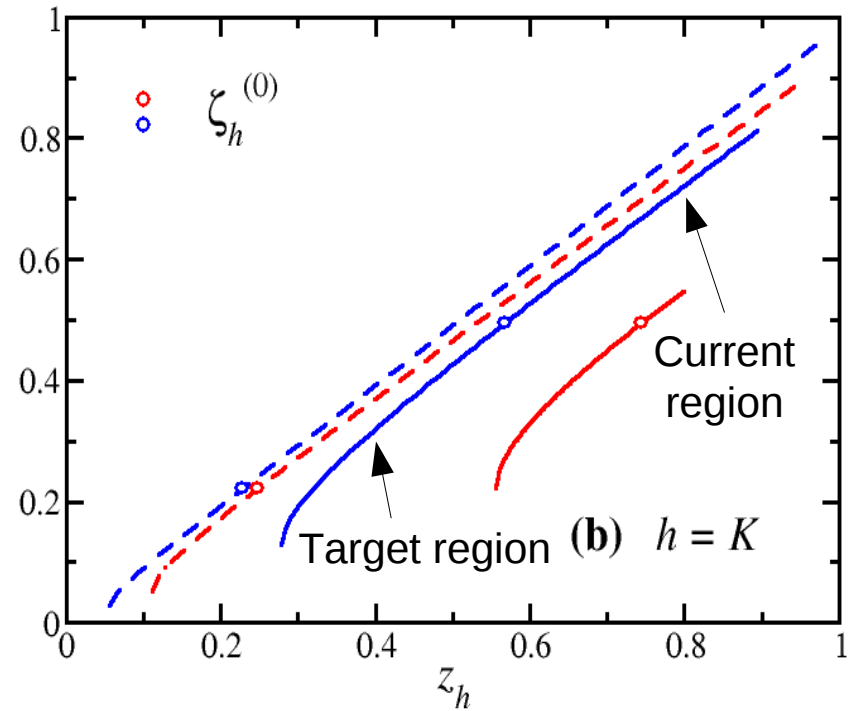
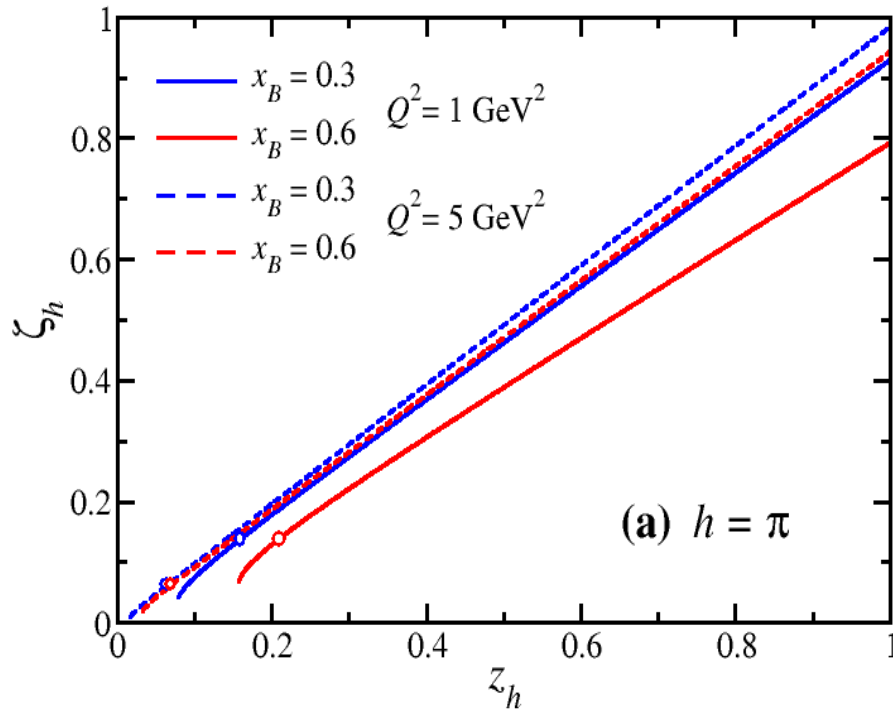
$$\zeta_h = \frac{p_h^-}{q^-} = \frac{z_h}{2} \frac{\xi}{x_B} \left(1 + \sqrt{1 - \frac{4x_B^2 M^2 m_{h\perp}^2}{z_h^2 Q^4}} \right)$$

with

$$m_{h\perp}^2 = m_h^2 + p_{h\perp}^2; \quad \zeta_h \longrightarrow z_h$$

$$Q^2 \rightarrow \infty$$

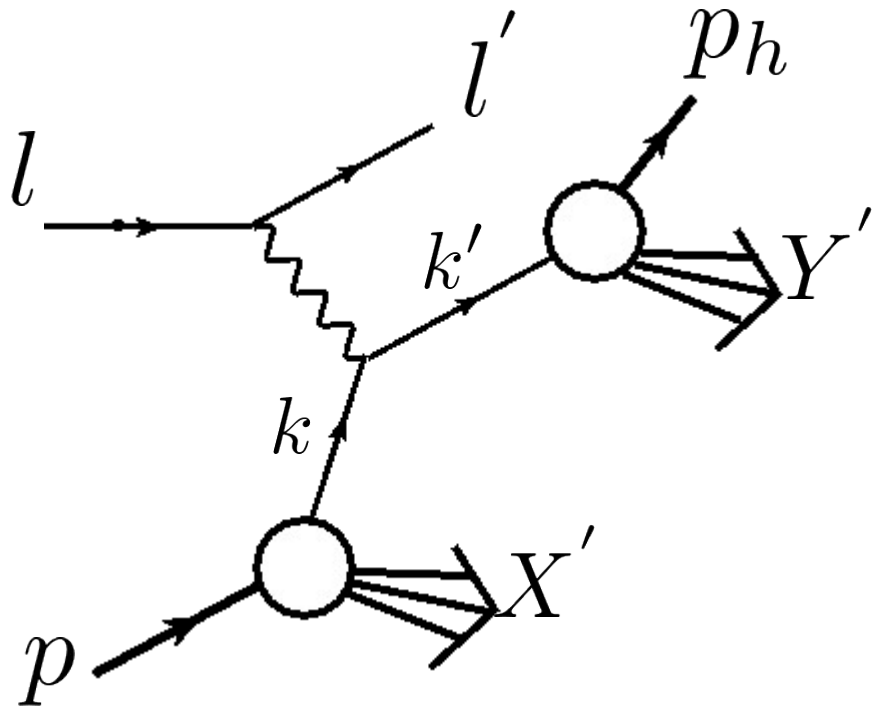
SIDIS: External Kinematics



- Target vs. Current Fragmentation region separated by \circ .
- Limits: 4-momentum + Baryon number conservation

$$\longrightarrow z_h^{\min} < z_h < z_h^{\max}$$

SIDIS: Parton Kinematics



Internal (parton) Kinematics

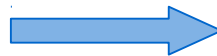
$$x = \frac{k^+}{p^+} \quad (\text{vs. } \xi = -\frac{q^+}{p^+})$$

$$z = \frac{p_h^-}{k'^-} \quad (\text{vs. } \zeta_h = -\frac{p_h^-}{q^-})$$

“Collinear” Factorization:

Expand around “collinear” momenta \tilde{k}_μ and \tilde{k}'_μ

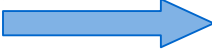
Initial and scattered quarks are off-shell



Virtualities \tilde{k}^2 and \tilde{k}'^2 will be fixed later

SIDIS: Internal Kinematics

Virtuality choice:

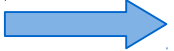
- Usual Choice: $\tilde{k}^2 = 0 \quad \tilde{k}'^2 = 0$

 at LO

$$x = \xi \xrightarrow{Q^2 \rightarrow \infty} x_B$$

$$z = \zeta_h \xrightarrow{Q^2 \rightarrow \infty} z_h$$

BUT kinematic analysis shows that

$$\tilde{k}^2 \geq x(\zeta_h - 1)Q^2/\xi$$

- Alternative Choice $\tilde{k}^2 = 0 \quad \tilde{k}'^2 = \frac{m_h^2}{\zeta_h}$

 at LO

$$x = \xi_h = \xi \left(1 + \frac{m_h^2}{\zeta_h Q^2}\right)$$

$$z = \zeta_h$$

(respects kinematic bound)



NOT UNIQUE: Needs more analysis at NLO

LO Cross section at finite Q^2

Spin averaged Cross section

$$\sigma_h \equiv \frac{1}{2} \frac{d\sigma_h^{\uparrow\uparrow+\downarrow\uparrow}}{dx_B dQ^2 dz_h} = \frac{2\pi\alpha^2}{Q^4} \frac{y^2}{1-\varepsilon} J_h \sum_q e_q^2 q(\xi_h, Q^2) D_q^h(\zeta_h, Q^2)$$

Jacobian $J_h = \frac{d\zeta_h}{dz_h}$

Unpolarised PDF


Fragmentation Function

LO Cross section at finite Q^2


Spin averaged Cross section

$$\sigma_h \equiv \frac{1}{2} \frac{d\sigma_h^{\uparrow\uparrow+\downarrow\uparrow}}{dx_B dQ^2 dz_h} = \frac{2\pi\alpha^2}{Q^4} \frac{y^2}{1-\varepsilon} J_h \sum_q e_q^2 q(\xi_h, Q^2) D_q^h(\zeta_h, Q^2)$$

Jacobian $J_h = \frac{d\zeta_h}{dz_h}$




Unpolarised
PDF



Fragmentation
Function

Spin dependent Cross section

$$\Delta\sigma_h \equiv \frac{d\sigma_h^{\uparrow\uparrow-\downarrow\uparrow}}{dx_B dQ^2 dz_h} = \frac{4\pi\alpha^2}{Q^4} \frac{y^2 \sqrt{1-\varepsilon^2}}{1-\varepsilon} J_h \sum_q e_q^2 \Delta q(\xi_h, Q^2) D_q^h(\zeta_h, Q^2)$$




Polarised PDF

LO Cross section at finite Q^2


Spin averaged Cross section

$$\sigma_h \equiv \frac{1}{2} \frac{d\sigma_h^{\uparrow\uparrow+\downarrow\uparrow}}{dx_B dQ^2 dz_h} = \frac{2\pi\alpha^2}{Q^4} \frac{y^2}{1-\varepsilon} J_h \sum_q e_q^2 q(\xi_h, Q^2) D_q^h(\zeta_h, Q^2)$$

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Unpolarised
PDF



Fragmentation
Function

Spin dependent Cross section

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Polarised PDF

Bjorken Limit: $\left(\frac{M^2}{Q^2}, \frac{m_h^2}{Q^2} \rightarrow 0\right)$

$$\sigma_h^{(0)} \equiv \sigma_h(\xi_h \rightarrow x_B, \zeta_h \rightarrow z_h)$$

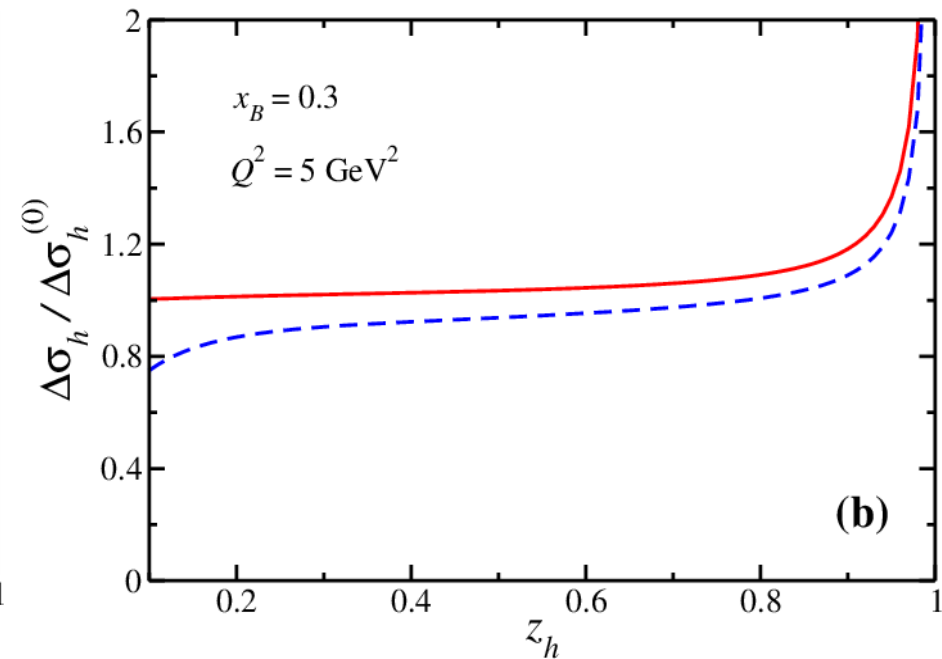
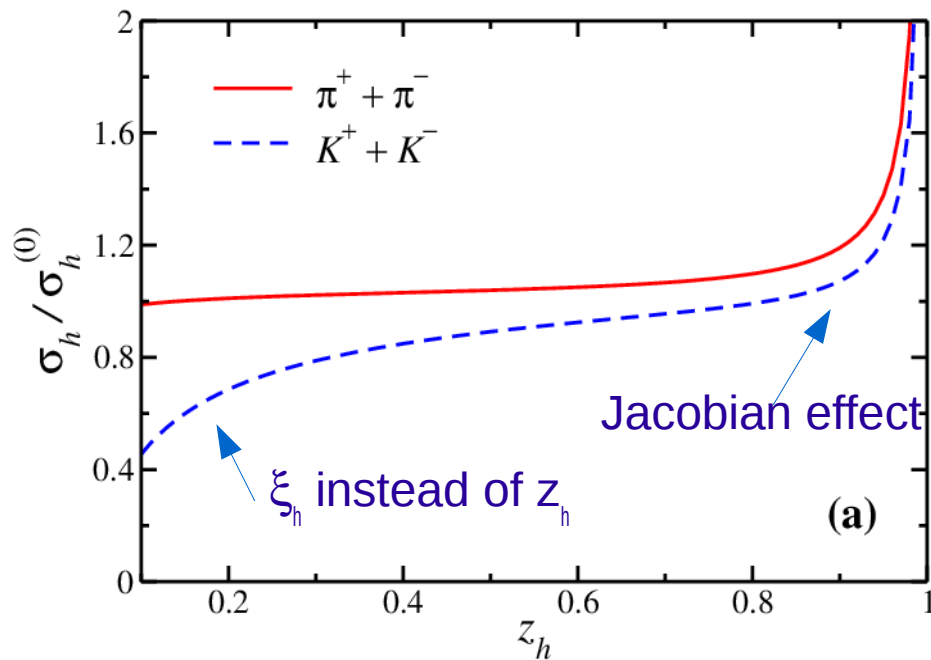
$$\Delta\sigma_h^{(0)} \equiv \Delta\sigma_h(\xi_h \rightarrow x_B, \zeta_h \rightarrow z_h)$$

Phenomenological implications

Pions vs. Kaons

Unpolarised

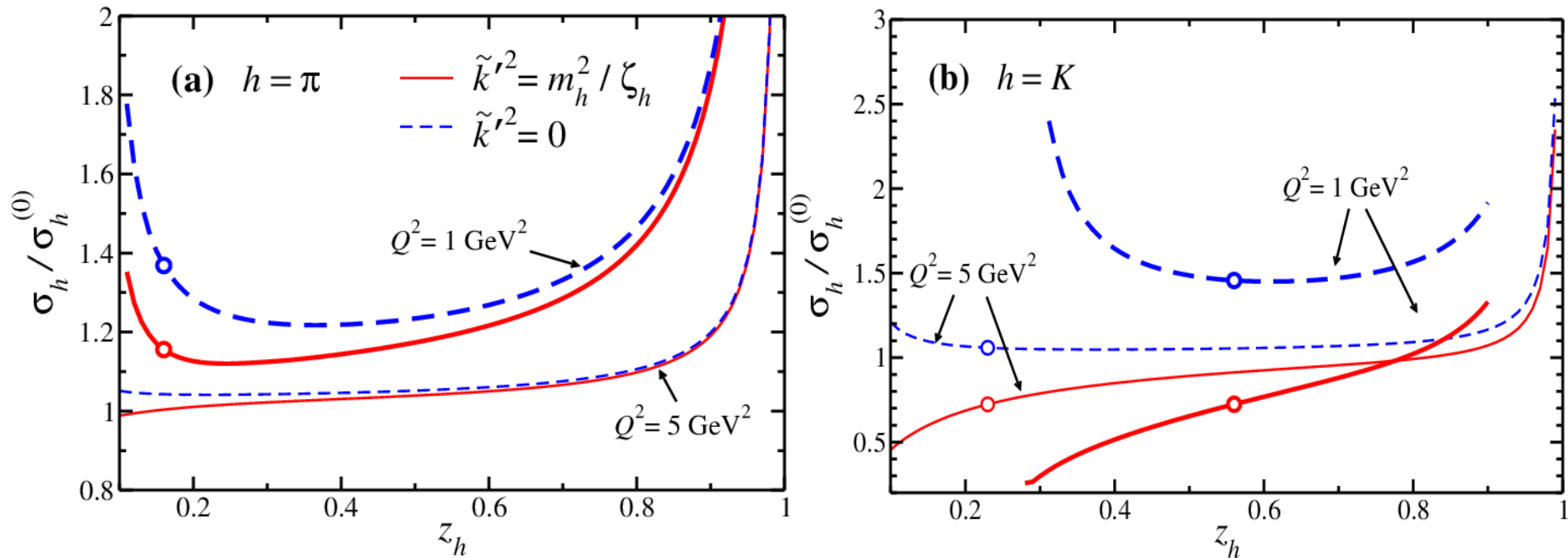
Polarised



Kaons mass effects much larger

Phenomenological implications

Dependence on the parton virtuality



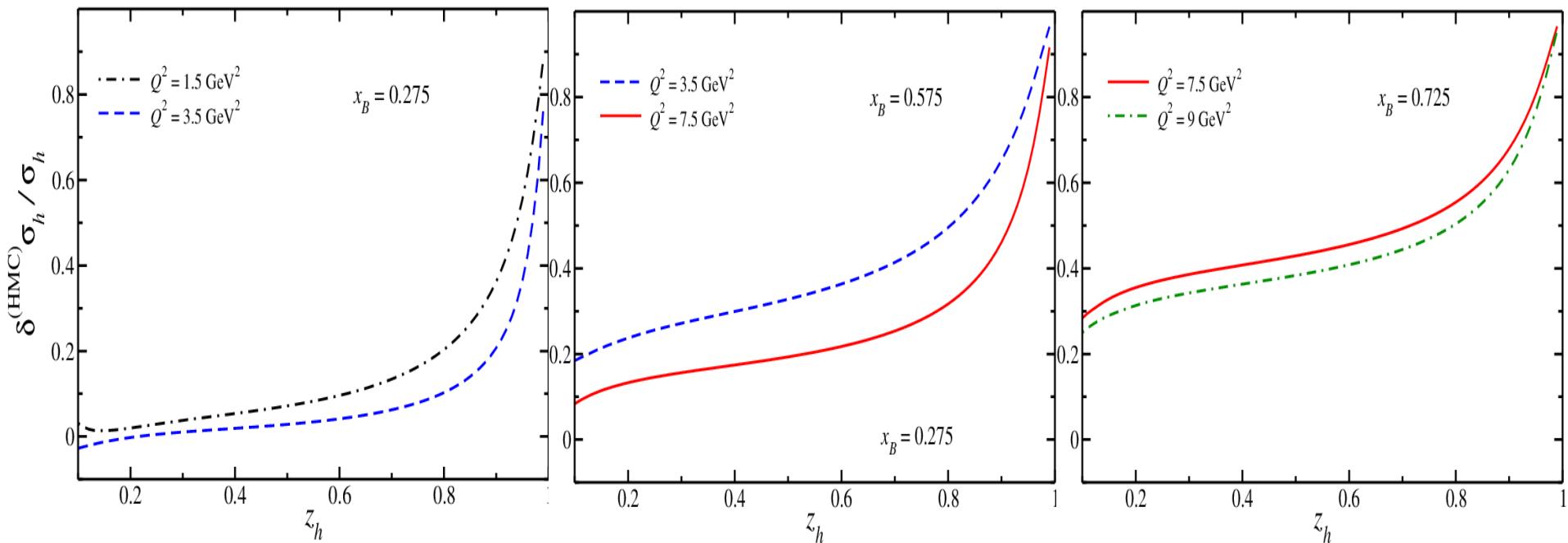
Red lines: Our choice $\tilde{k}'^2 = m_h^2 / \zeta_h$

Thin lines: Albino et al. Nucl.Phys B803, 42 (2008) $\tilde{k}'^2 = 0$

Mass corrections for specific experiments

Jlab 11 GeV (Unpolarized Pions)

(Experiments: E12-09-007, E12-13-007, E12-06-109)



$$\frac{\delta^{(\text{HMC})} \sigma_h}{\sigma_h} = \frac{\sigma_h - \sigma_h^{(0)}}{\sigma_h}$$

Relative magnitude of the Mass Correction

Strange quark PDF

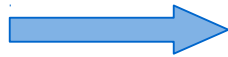
- **HERMES**



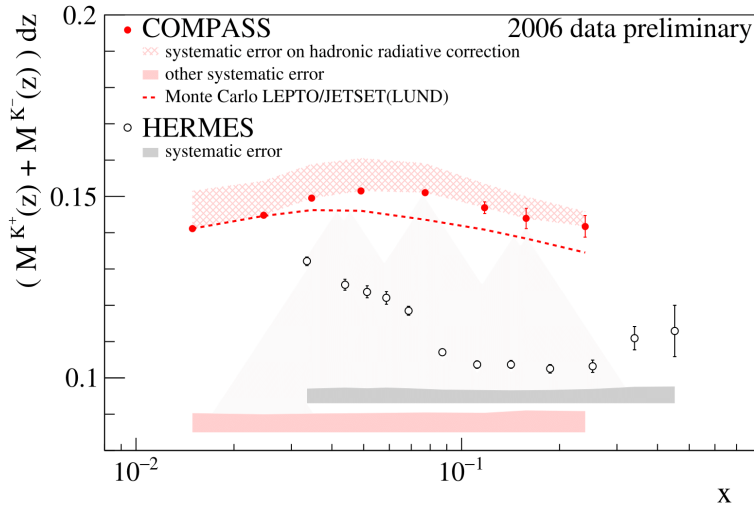
- Claim very different s-quark shape compared to CTEQ6L
- Measurements from ATLAS/CMS at LHC also show different s-PDF
- Strange pdf may not be what we think!

Strange quark PDF

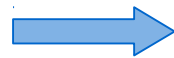
• HERMES



- Claim very different s-quark shape compared to CTEQ6L
- Measurements from ATLAS/CMS at LHC also show different s-PDF
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• But COMPASS



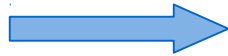
- Different x_B dependence

$$\langle Q^2 \rangle_{\text{COMPASS}} \gtrsim \langle Q^2 \rangle_{\text{HERMES}}$$

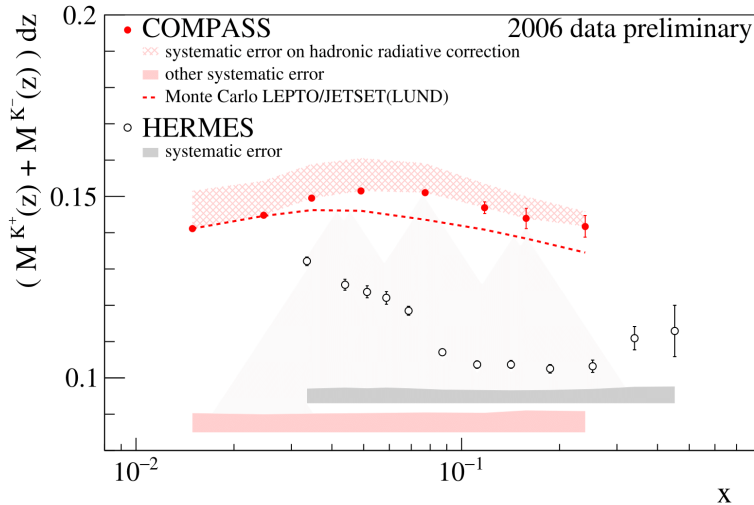
- Overall value higher

Strange quark PDF

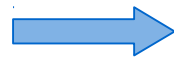
• HERMES



- Claim very different s-quark shape compared to CTEQ6L
- Measurements from ATLAS/CMS at LHC also show different s-PDF
- Strange pdf may not be what we think!



• But COMPASS



- Different x_B dependence

$$\langle Q^2 \rangle_{\text{COMPASS}} \gtrsim \langle Q^2 \rangle_{\text{HERMES}}$$

- Overall value higher

Can HMCs reduce the discrepancy HERMES vs. COMPASS?

Strange quark PDF

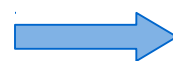
Hadron Multiplicity(Kaon) \longrightarrow Strange quark “Tagging”

Experimental definition:
$$M^k(x_B, z_h, Q^2) = \frac{\frac{dN^K}{dx_B dQ^2 dz_h}}{\frac{dN^{DIS}}{dx_B dQ^2}}$$

Parton model:
$$M^K(x_B, z_h, Q^2) = \frac{\sum_q e_q^2 q(x_B, Q^2) D_q^h(z_h, Q^2)}{\sum_q e_q^2 q(x_B, Q^2)}$$

Integrated Multiplicity:
$$M^k(x_B, Q^2) = \int_{0.2}^{0.8(0.85)} dz_h M^k(x_B, Q^2, z_h)$$

HERMES (COMPASS)



Sensitive to s-quark

Strange quark PDF

SIDIS on Deuteron (HERMES and COMPASS)

Parton model:
$$M^K(x_B, Q^2) = \frac{Q(x) \int \mathcal{D}_Q^K(z) dz + S(x) \int \mathcal{D}_S^K(z) dz}{5Q(x) + 2S(x)}$$

$$Q(x) \equiv u(x) + \bar{u}(x) + d(x) + \bar{d}(x)$$

$$S(x) \equiv s(x) + \bar{s}(x)$$

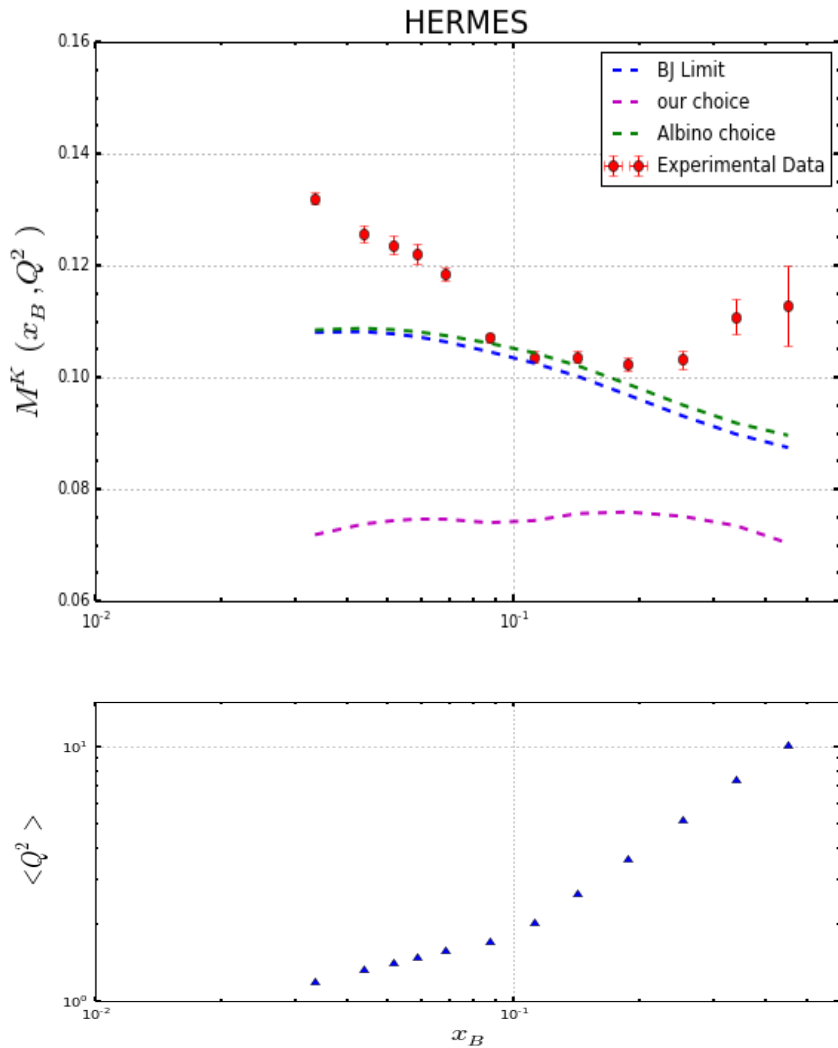
$$\mathcal{D}_Q^K(z) \equiv 4D_u^K(z) + D_d^K(z)$$

$$\mathcal{D}_s^K(z) \equiv 2D_s^K(z)$$

$$K = K^+ + K^-$$

Strange quark PDF

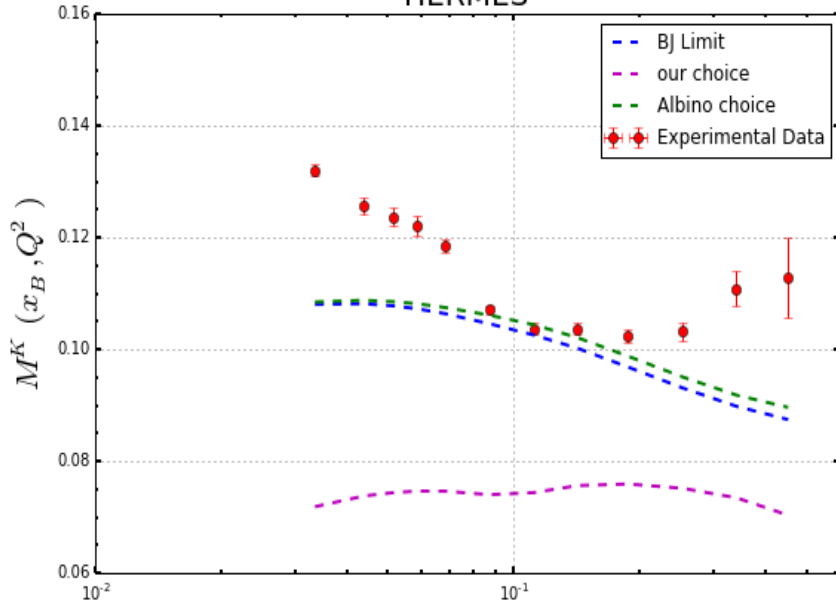
PDFs: CTEQ6L, FFs: DSSV07



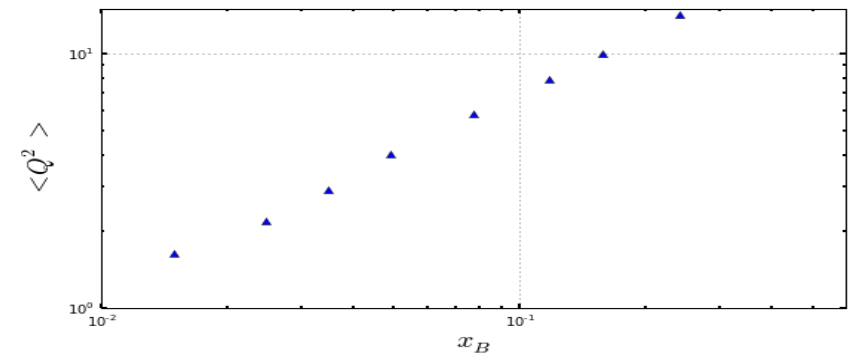
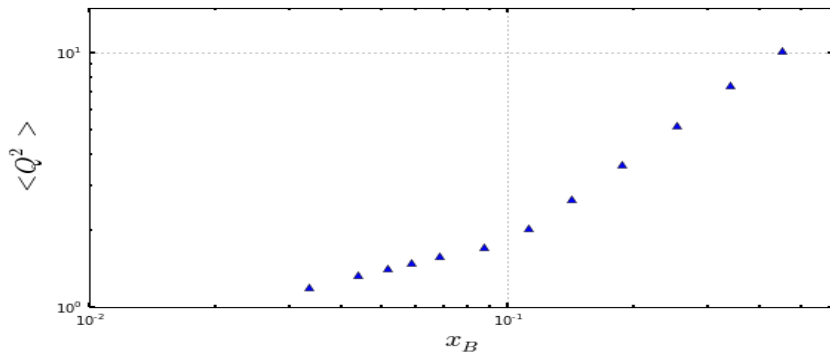
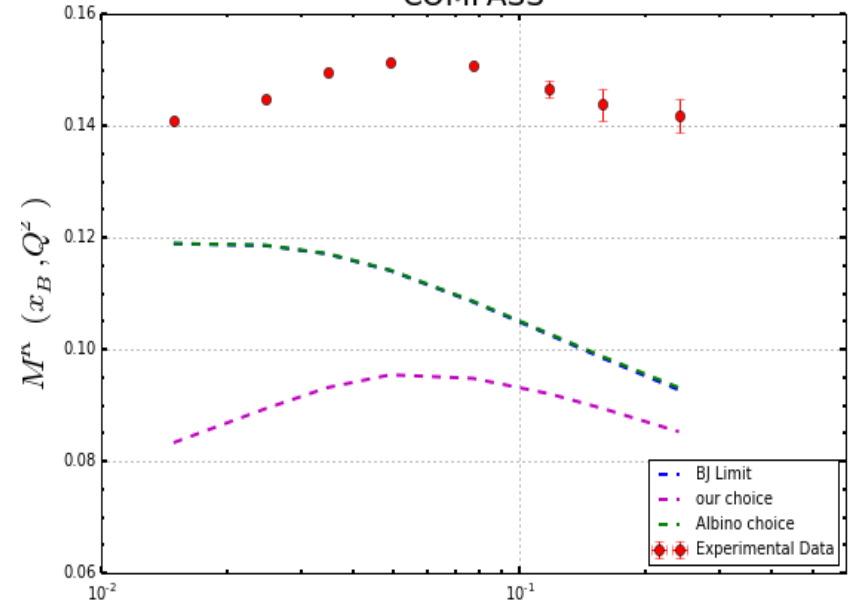
Strange quark PDF

PDFs: CTEQ6L, FFs: DSSV07

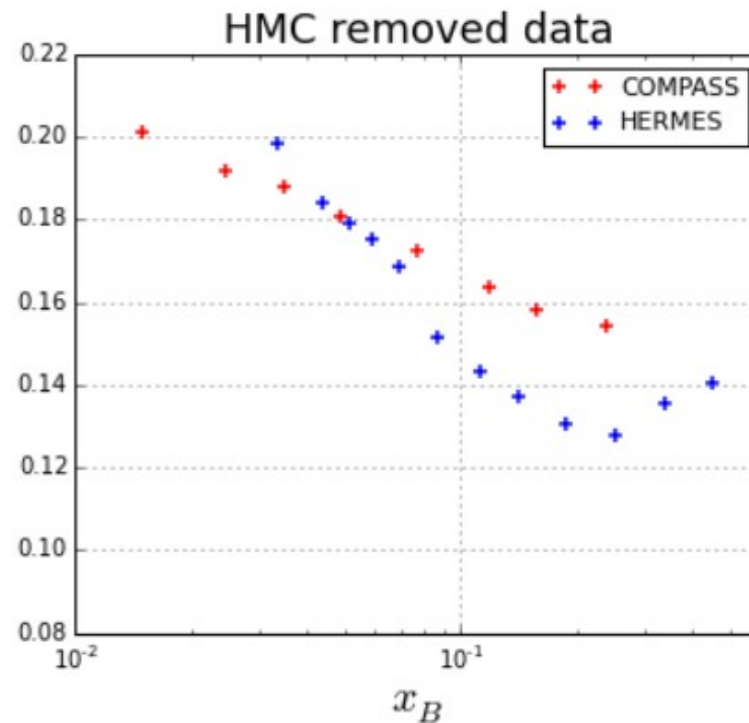
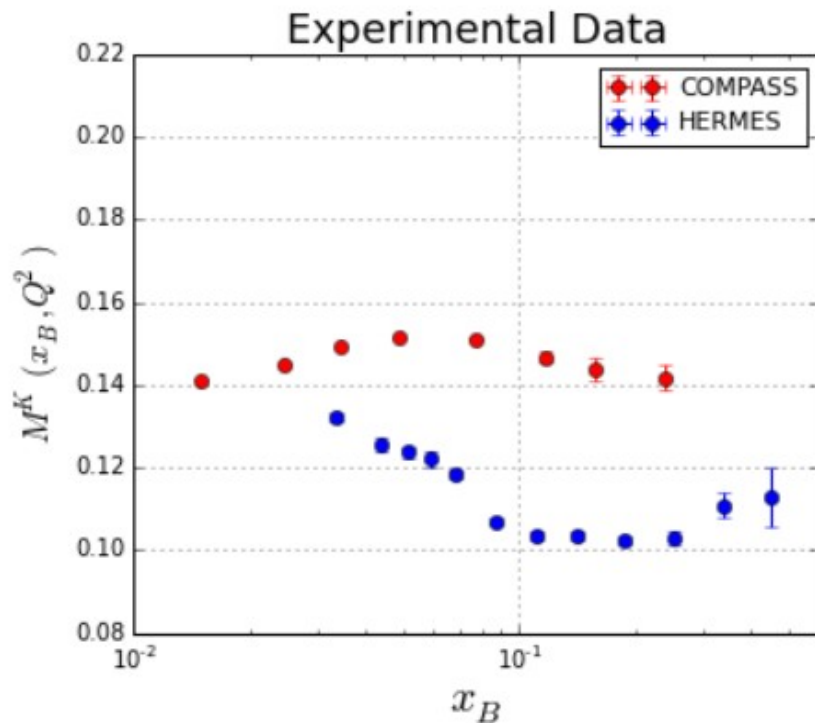
HERMES



COMPASS



COMPASS vs. HERMES



Data without HMCs: $D_{w/oHMCs} = D_{exp} \frac{M^{K(0)}}{M^K} \Big|_{Theory}$

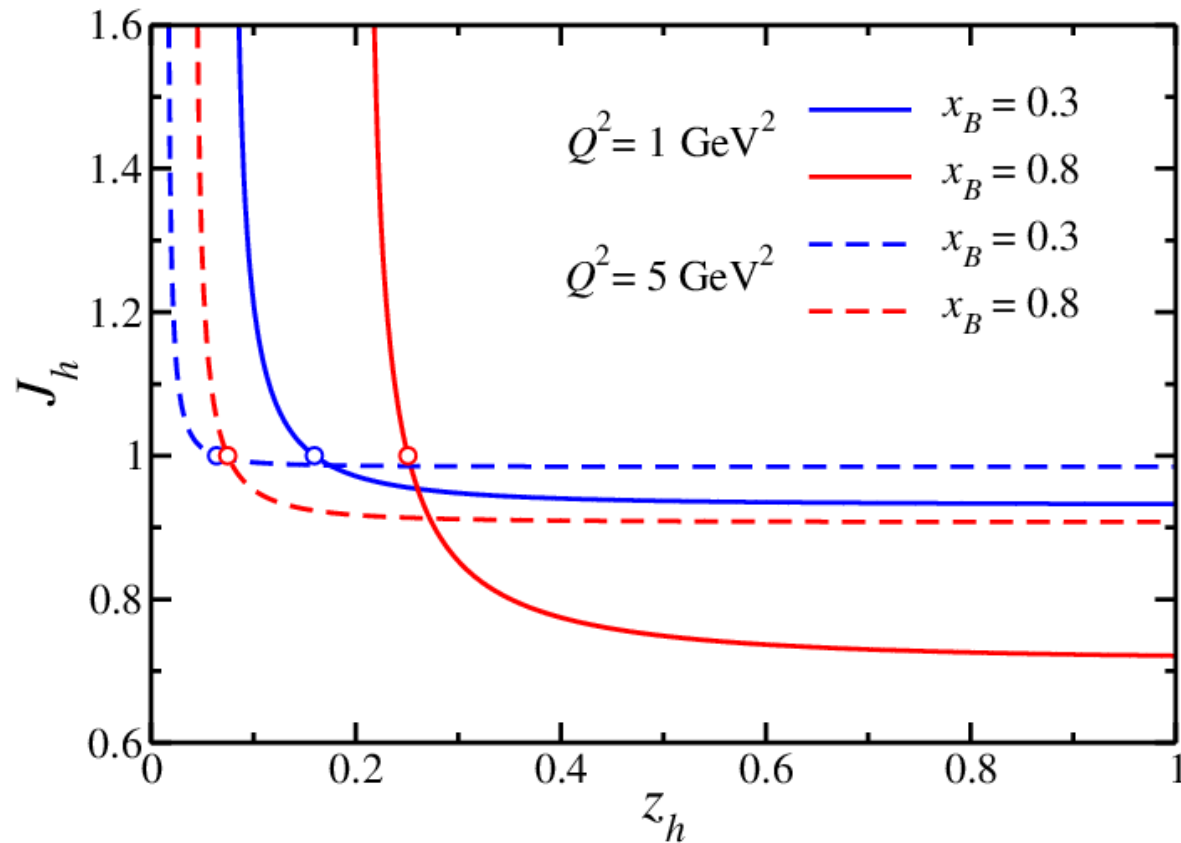
Conclusion and outlook

- HMC's at LO are captured by new scaling variables ξ_h and ζ_h .
 - Matching “internal” and external kinematics requires a massive fragmenting parton with $\tilde{k}'^2 \geq m_h^2/\zeta_h$.
 - x_B and z_h mixed in the new scaling variables ξ_h and ζ_h
- We have quantified HMC's effects numerically: stronger for large as well very small values of z_h , stronger effects at large x_B and low Q^2 (as DIS).
- More dramatic effects for Kaons compared to pions.
- Jefferson Lab Pions: corrections up to ~40 %- 50 % even at moderately large Q^2 (small effects for HERMES and COMPASS)
- HMCs reduce the discrepancy between HERMES and COMPASS data for averaged Kaon integrated multiplicities
- Future development: prove factorization at NLO in the presence of massive fragmenting partons with non-zero virtuality, and “correct” choice of \tilde{k}'^2 .

Thank you!

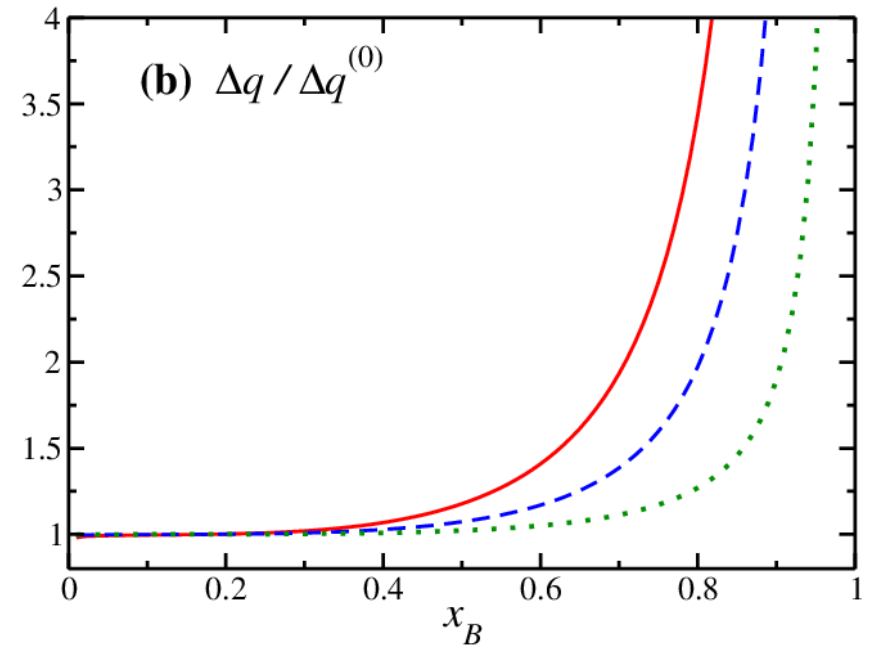
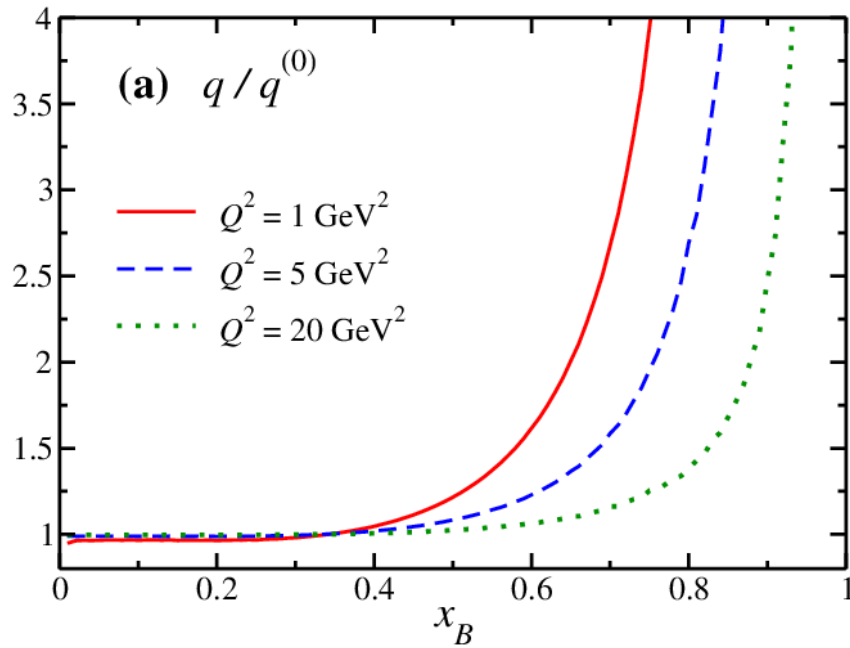
Backup slides

Jacobian



Backup slides

PDF's



Backup slides

Fragmentation function

