



QCD Frontier 2013

Tensor Polarized Deuteron

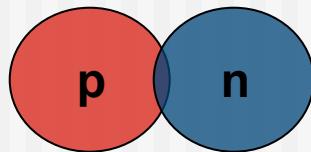
**QCD Frontier 2013 Meeting
October 21-22, 2013**

Narbe Kalantarians
Hampton University



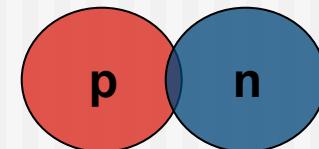
Outline

- Background/Motivation
- Spin-1/Tensor-Polarization Concept
- Previous and planned measurements
- Possible (Spin-1) Physics with an EIC

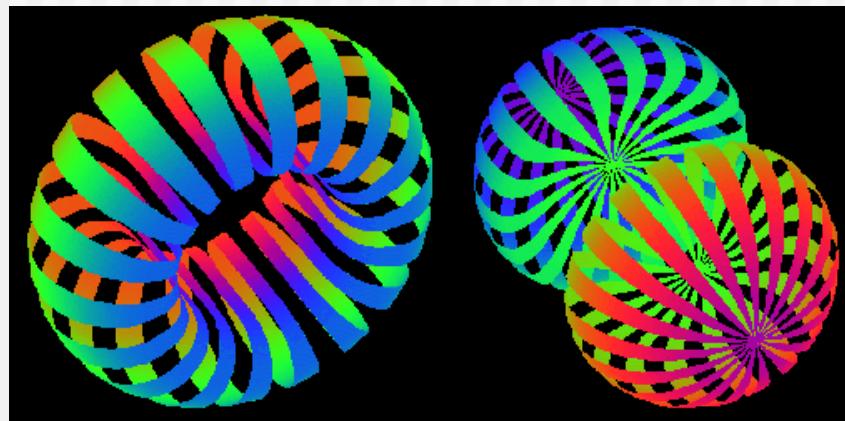


Why Deuteron?

- Spin-1 system
- Simple lab for nuclear physics
- Reasonably “easy” to polarize.



Spatial distribution depends on the spin state

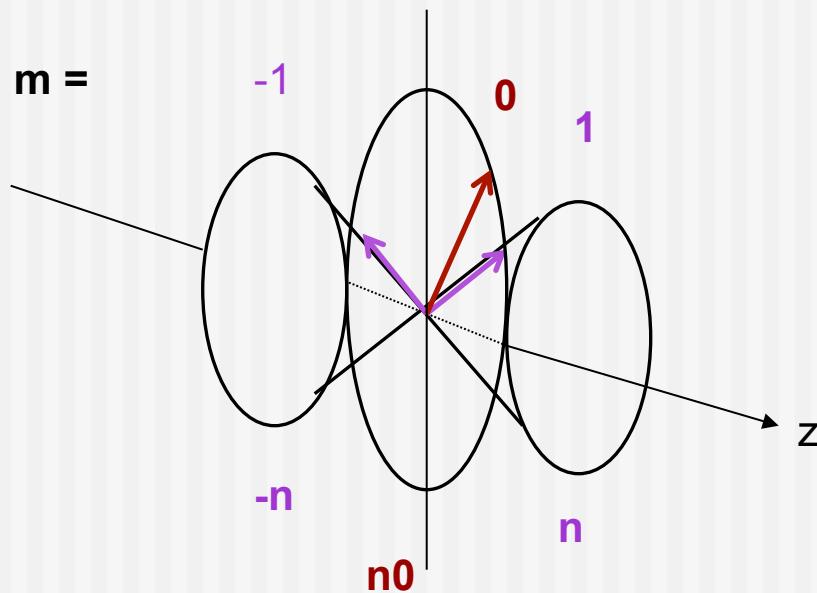


$M = 0$

$M = \pm 1$

Spin-1

Spin-1 system in a B-field leads to 3 sublevels via Zeeman interaction.

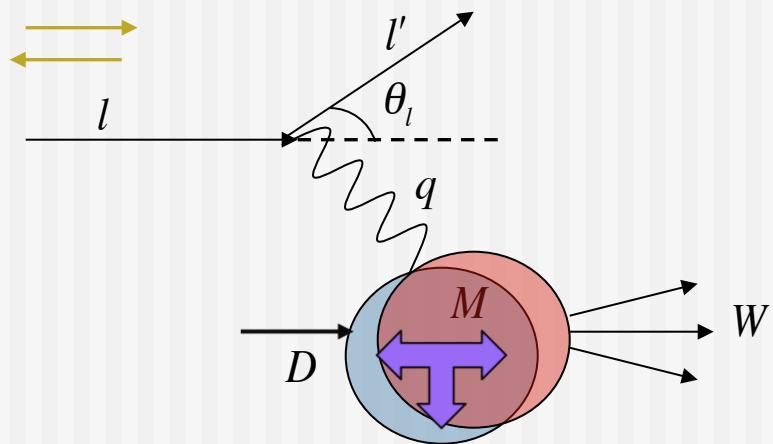


Vector polarization: $(n^+ - n^-)$; $-1 < P_z < +1$

Tensor polarization: $(n^+ - n^0) - (n^0 - n^-)$; $-2 < P_{zz} < +1$ Normalization: $(n^+ + n^- + n^0) = 1$

Some research has been done with deuteron beams (Thesis: V. Morozov)

Inclusive Scattering with Spin-1



Frankfurt & Strikman (1983)

Hoodbhoy, Jaffe, Manohar (1989)

$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{Mott} \left[\frac{1}{v} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2(\theta/2) \right] \\ + \gamma g_1(x, Q^2) + \delta g_2(x, Q^2)$$

$$+ \zeta b_1(x, Q^2) + \varepsilon b_2(x, Q^2) + \xi b_3(x, Q^2) + \eta b_4(x, Q^2)$$

$$b_1, b_2 \sim \frac{1}{P_{zz}^{eff}}$$

Spin-1 => 4 more structure-functions: b_1, b_2, b_3, b_4

Spin-1 Structure Functions

Leading Twist: F_1, g_1, b_1

	Nucleon	Deuteron
F_1	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^{1/2} + q_{\downarrow\uparrow}^{-1/2}]$	$\frac{1}{3} \sum_q e_q^2 [q_{\uparrow}^1 + q_{\uparrow}^{-1} + q_{\uparrow}^0]$
g_1	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^{1/2} - q_{\downarrow\uparrow}^{-1/2}]$	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^1 - q_{\downarrow}^{-1}]$
b_1	...	$\frac{1}{2} \sum_q e_q^2 [q^0 - q^1]$ $q^0 = (q_{\uparrow}^0 + q_{\downarrow}^0) = 2q_{\uparrow}^0$ $q^1 = (q_{\uparrow}^1 + q_{\downarrow}^1) = (q_{\uparrow}^1 + q_{\uparrow}^{-1})$

F_1 : quark distributions averaged over spin states

g_1 : difference of distributions of quarks aligned/anti-aligned with nucleon

b_1 : difference of helicity-0/helicity non-zero states of the deuteron

Spin-1 Structure Functions

Leading Twist: F_1, g_1, b_1

	Nucleon	Deuteron
F_1	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^{1/2} + q_{\downarrow\uparrow}^{-1/2}]$	$\frac{1}{3} \sum_q e_q^2 [q_{\uparrow}^1 + q_{\uparrow}^{-1} + q_{\uparrow}^0]$
g_1	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^{1/2} - q_{\downarrow\uparrow}^{-1/2}]$	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^1 - q_{\downarrow}^{-1}]$
b_1	...	$\frac{1}{2} \sum_q e_q^2 [q_{\uparrow}^0 - q_{\uparrow}^1]$

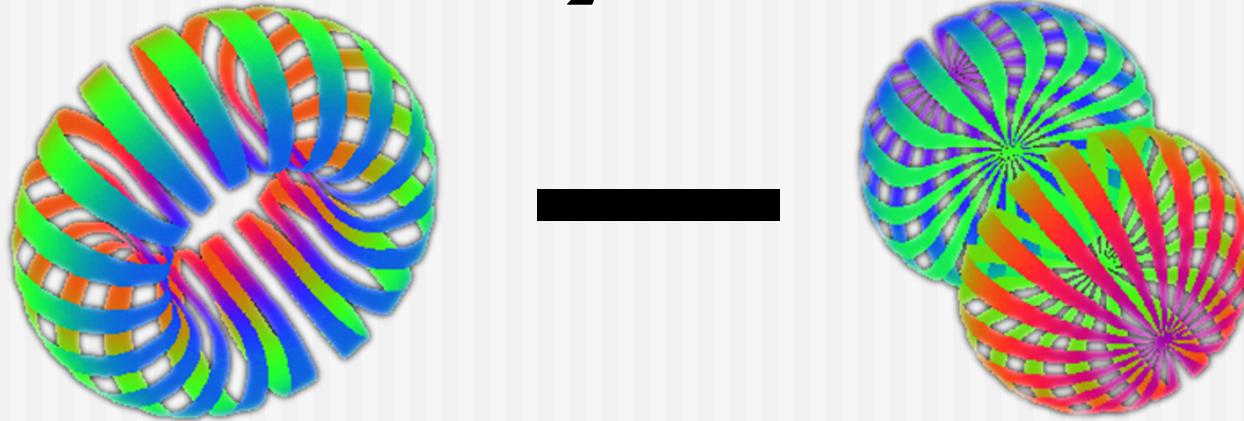
b_2 : related to b_1 , by relation similar to Callan-Gross.

b_4 : kinematically suppressed at longitudinal polarization. Also, leading twist.

b_3 : higher twist, similar to g_2 .

b_1^d

$$b_1^d \approx \frac{1}{2}(q^0 - q^1)$$



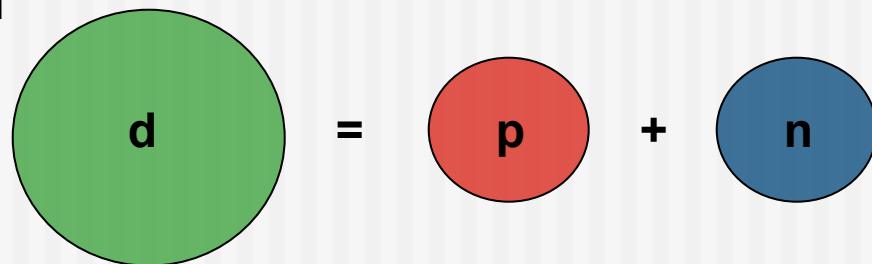
- Deuteron essentially combination of nuclear and quark physics.
- Measured via DIS, but dependent on deuteron spin-state.
- Allows for investigation of nuclear effects at parton level.

b_1^d

Hoodbhoy, Jaffe, Manohar (1989)

b_1 vanishes in the absence of nuclear effects.

i.e., if



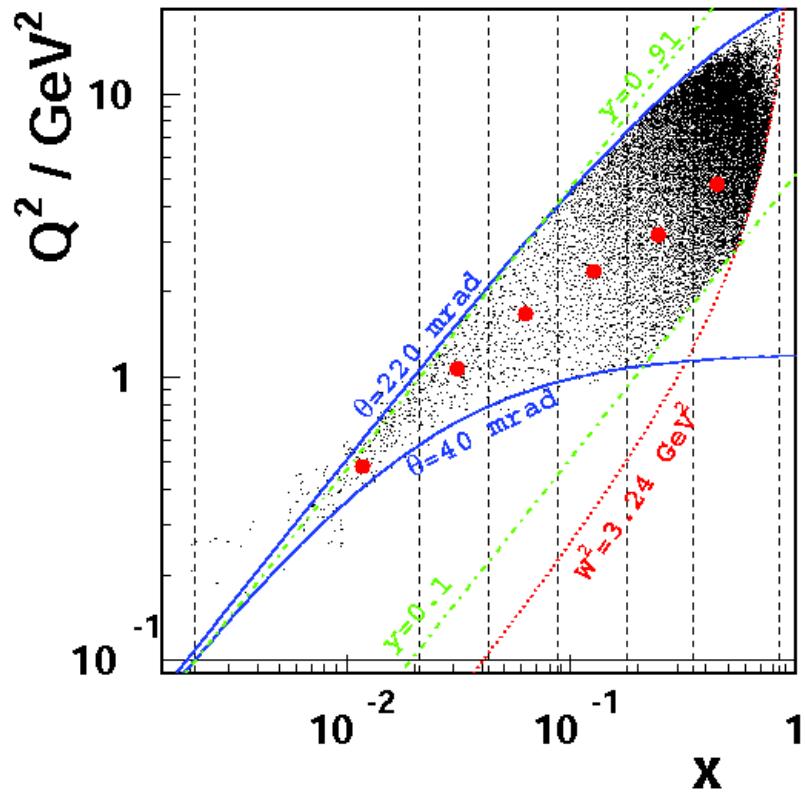
p,n in relative S-state

D-state has component with nuclear spin 1 from orbital angular momentum 1, but total nucleon spin 0. But even accounting for this, expected to be very small.

Khan & Hoodbhoy, PRC 44 , 1219 (1991) : $b_1 \approx O(10^{-4})$
Relativistic convolution model with binding

Umnikov, PLB 391, 177 (1997) : $b_1 \approx O(10^{-3})$
Relativistic convolution with Bethe-Salpeter formalism

HERMES Measurement: kinematics

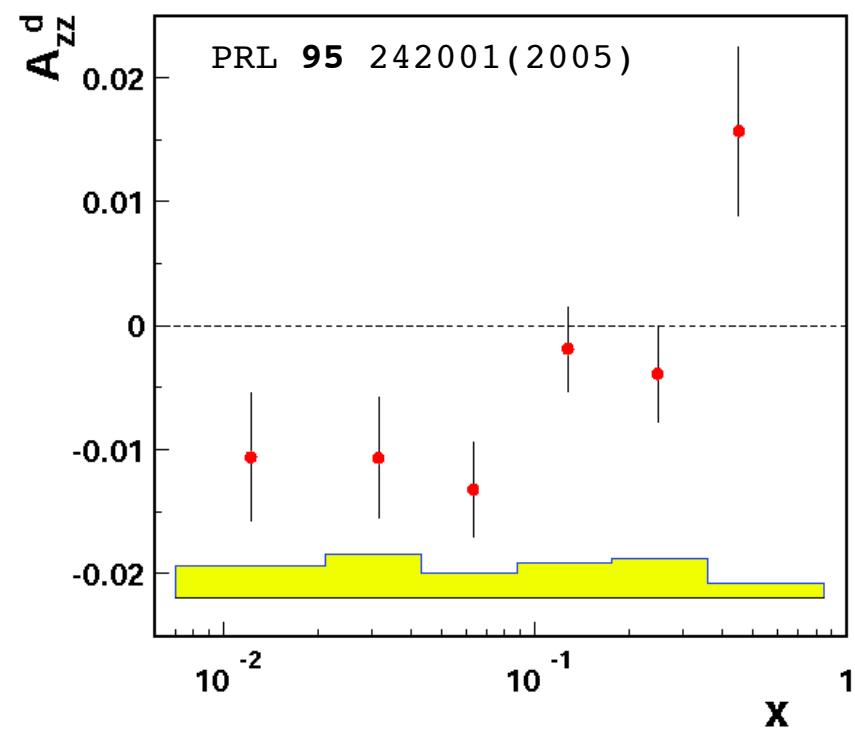


$0.01 < x < 0.45$

$0.5 < Q^2 < 5 \text{ GeV}^2$

- 27.6 GeV longitudinally polarized positron beam
- Internal tensor polarized d_2 gas target; $P_{zz} \sim 0.8$ (negligible P_z), dilution ~ 0.9 .
- 1 month of data taking.

HERMES Measurement: A_{zz}^d

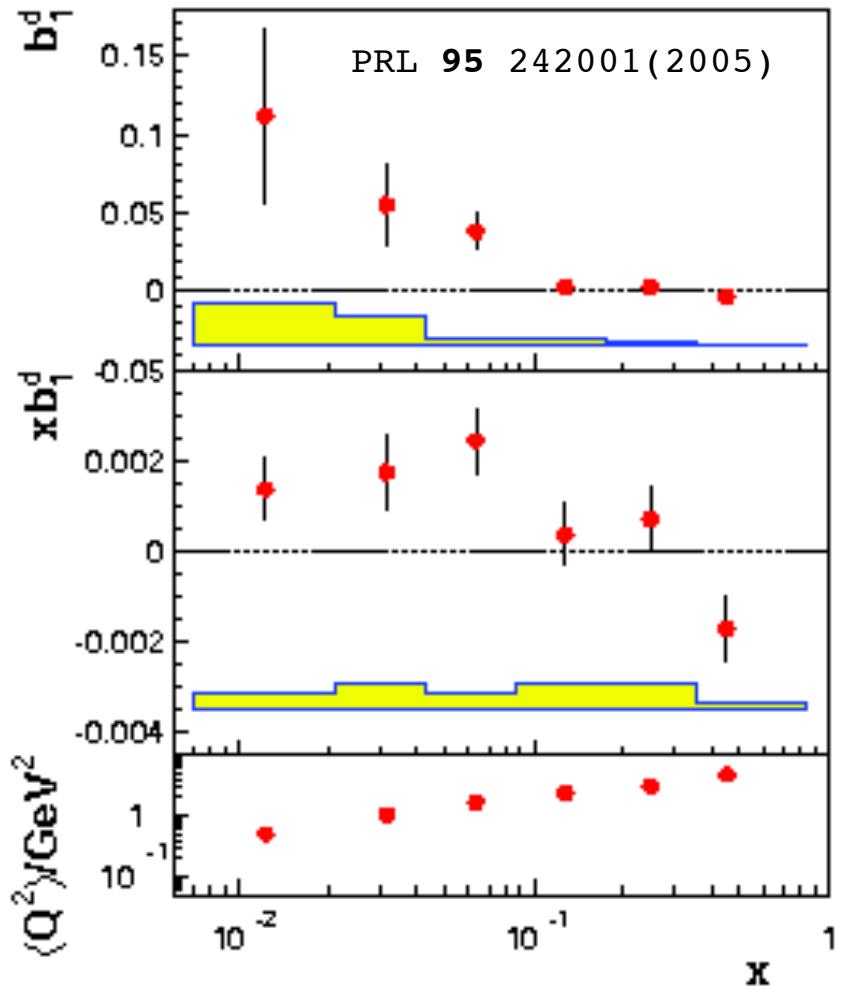


Tensor spin asymmetry

$$A_{zz} = \frac{1}{P_{zz}} \frac{2\sigma^1 - 2\sigma^0}{3\sigma^U}$$

HERMES result was about 2σ from 0.

HERMES Measurement: b_1^d



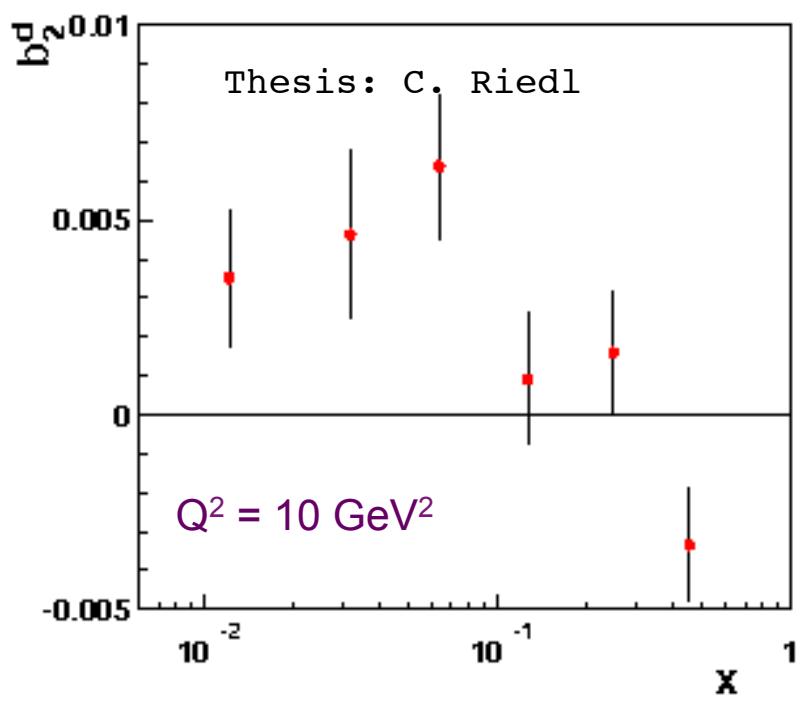
$$b_1 = -\frac{3}{2} F_1 A_{zz}$$

Rising of b_1 as $x \rightarrow 0$ can be related to same mechanism responsible for nuclear shadowing.

Ashman, et al. PLB 206 364(1988)

Can also be described in models involving double-scattering of leptons (first from proton, then neutron).

HERMES Measurement: b_2^d



b_2 related to b_1 via Callan-Gross-type relation.

$$b_2 = 2xb_1 \left(\frac{1+R}{1+\gamma^2} \right)$$

$$R = (1+\gamma^2) \frac{F_2}{2xF_1} - 1$$

HERMES Close-Kumano Sum Rule

F.E.Close, S.Kumano, PRD**42** 2377 (1990)

If sea quark and antiquark tensor polarization vanishes
i.e.

$$\int b_1(x)dx = 0$$

HERMES measurement:

$$\int_{0.02}^{0.85} b_1(x)dx = 0.0105 \pm 0.0034 \pm 0.0035$$

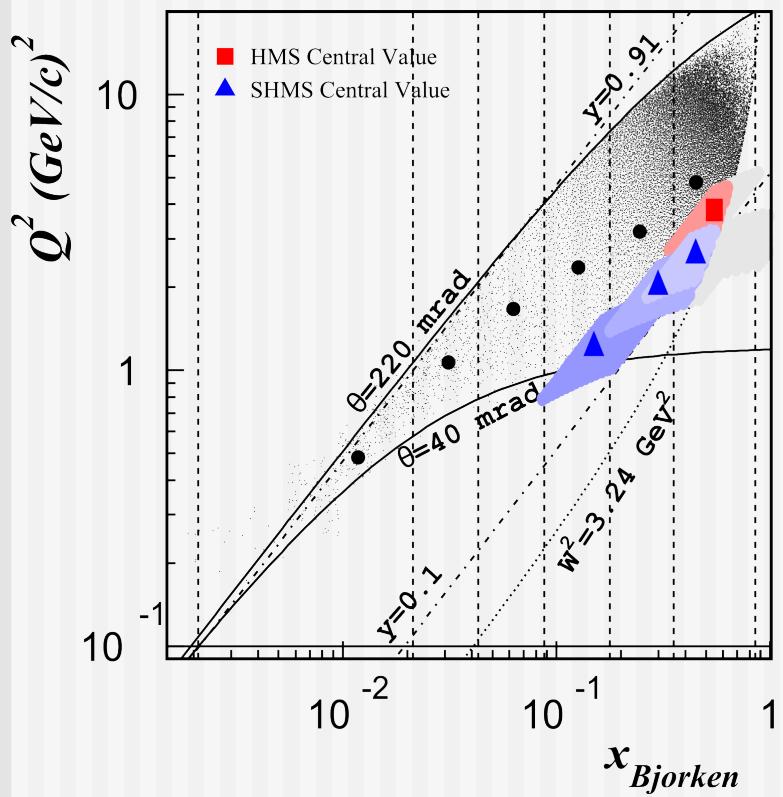
2σ result, over measured range

$$\int_{0.02}^{0.85} b_1(x)dx = 0.0035 \pm 0.0010 \pm 0.0018$$

1.7σ result, with $Q^2 > 1 \text{ GeV}^2$

PRL **95** 242001 (2005)

Proposal To Determine b_1^d at JLab



- Measurement at Jlab 12GeV could be complementary to HERMES.
- Advantage would be higher luminosity: $\sim 10^{35} \text{ cm}^{-2} \text{s}^{-1}$ compared to $\sim 10^{31} \text{ cm}^{-2} \text{s}^{-1}$.
- Some research has been done tensor polarizing solid deuterium (ND_3) target via NMR*: $P_{zz} \sim 0.2$, dilution $\sim 0.24, 0.36$.
- Submitted at PAC 40; Conditionally approved.

Experimental Method

$$A_{zz} = \frac{2}{fP_{zz}} \frac{\sigma_{\dagger} - \sigma_0}{\sigma_0}$$

$$= \frac{2}{fP_{zz}} \left(\frac{N_{\dagger}}{N_0} - 1 \right)$$

$$b_1 = -\frac{3}{2} F_1^d A_{zz}$$

Observable is the Normalized XS Difference

B-Field, density, temp, etc. held same in both states

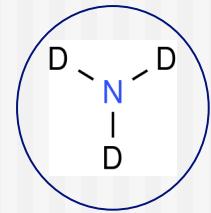
σ_{\dagger} : Tensor Polarized cross-section

σ_0 : Unpolarized cross-section

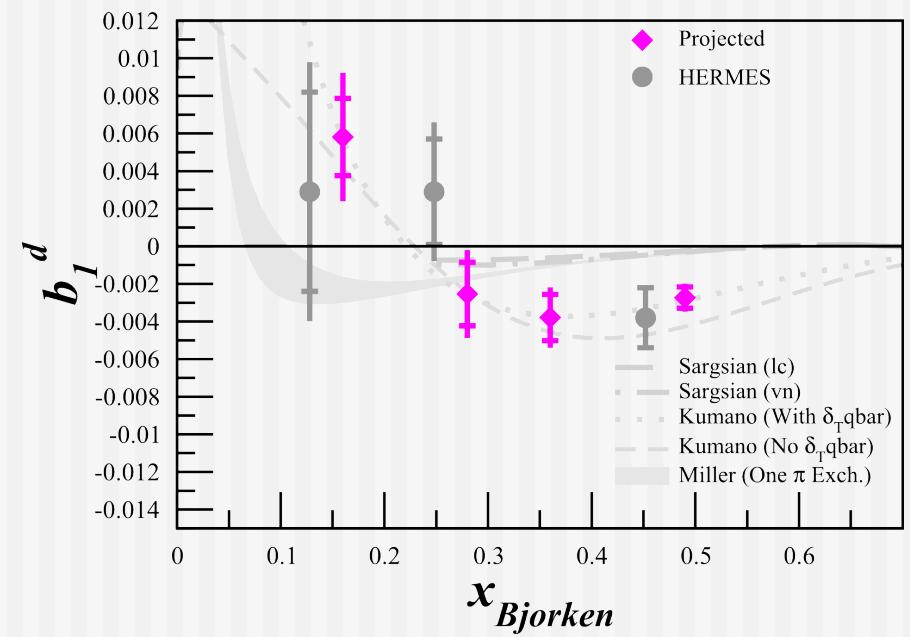
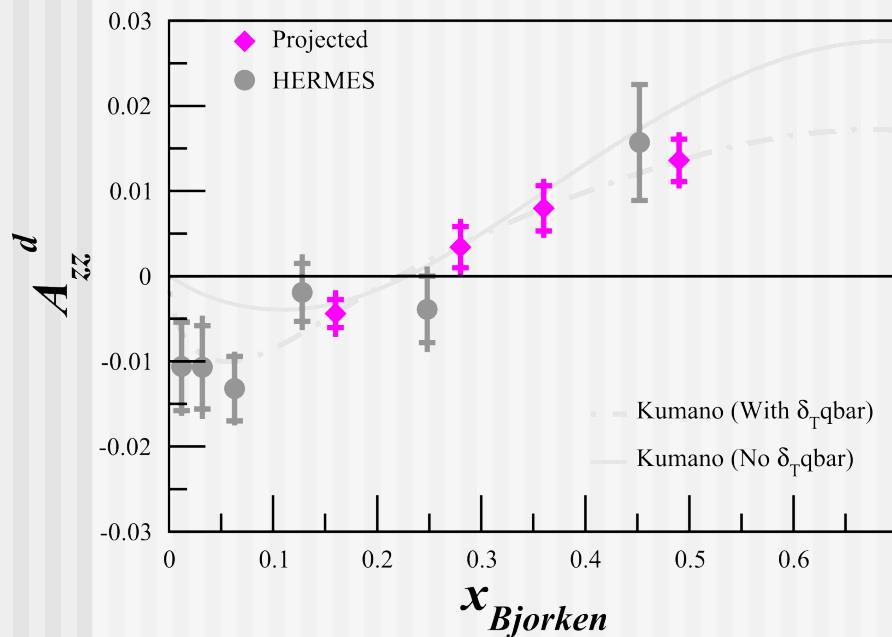
P_{zz} : Tensor Polarization

dilution factor

$$f \approx \frac{6}{20}$$



Projected Results: P_{zz} near 30%



RF saturation has demonstrated P_{zz} enhancement to ~30%

$$*Crabb \text{ et al. } P_{zz} = 2 - (4 - 3P_z^2)^{1/2}$$

false asymmetries suppressed by $1/P_{zz}$

$$\delta A_{zz} = \pm \frac{2}{f P_{zz} \sqrt{N_{cycles}}} \delta \xi$$

Tens. Pol. Scattering at low x

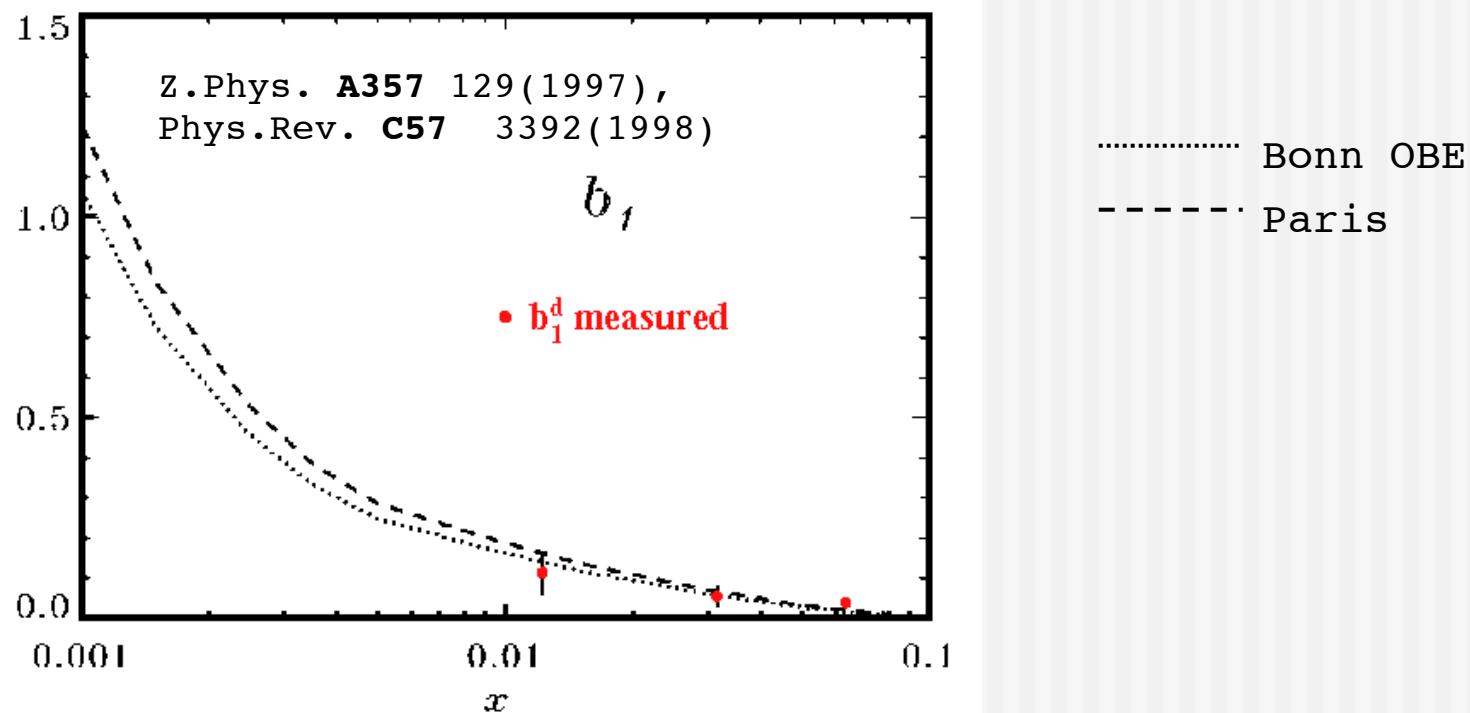
Possibilities

- Small x aspect of tensor pol. (deuteron) could access anti-shadowing and 2 nucleon scattering.
- Could also provide complementary information for OAM.
- A good starting point would be to extract b_1^d with an EIC.

Issues to Address:

- How well can polarization, beam stability be understood and controlled?
- Need to do simulation studies.

b_1^d Predictions

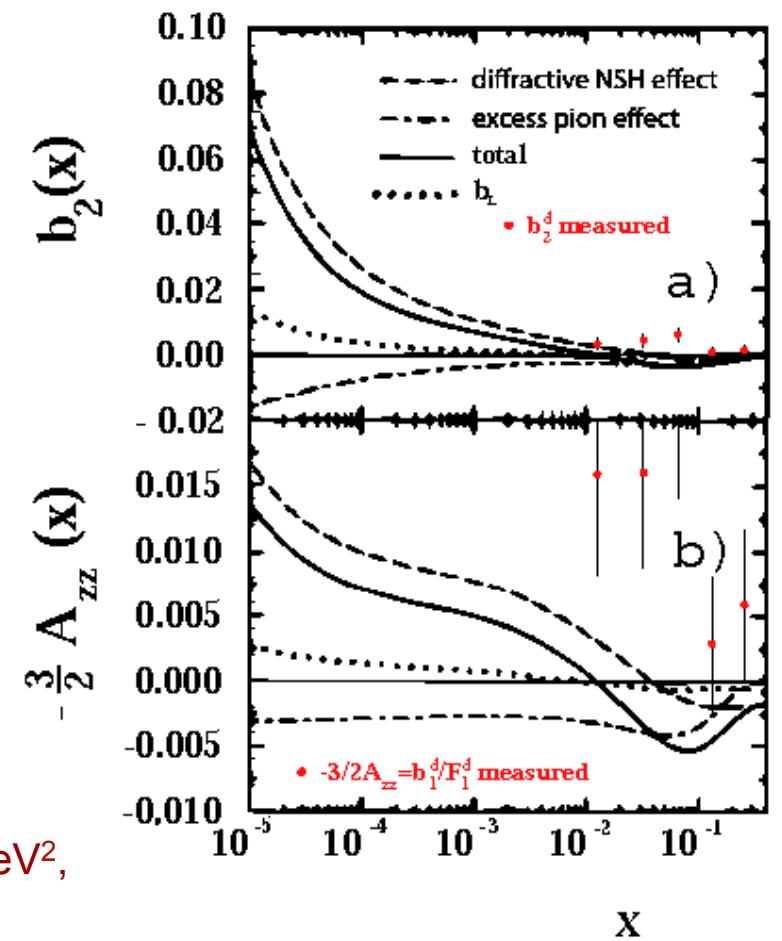
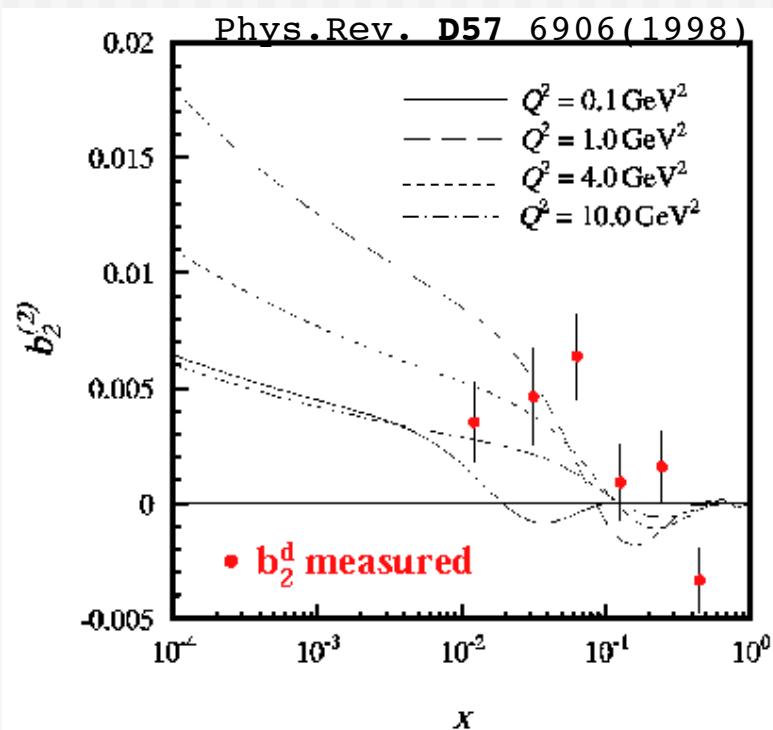


- Both models predict b_1 (rapidly) increasing as $x \rightarrow 0$.
- Errors for (HERMES) data shown are statistical only.

L. Frankfurt, V. Guzey, M. Strikman
Mod. Phys.Lett. **A21**(2006) 23-40

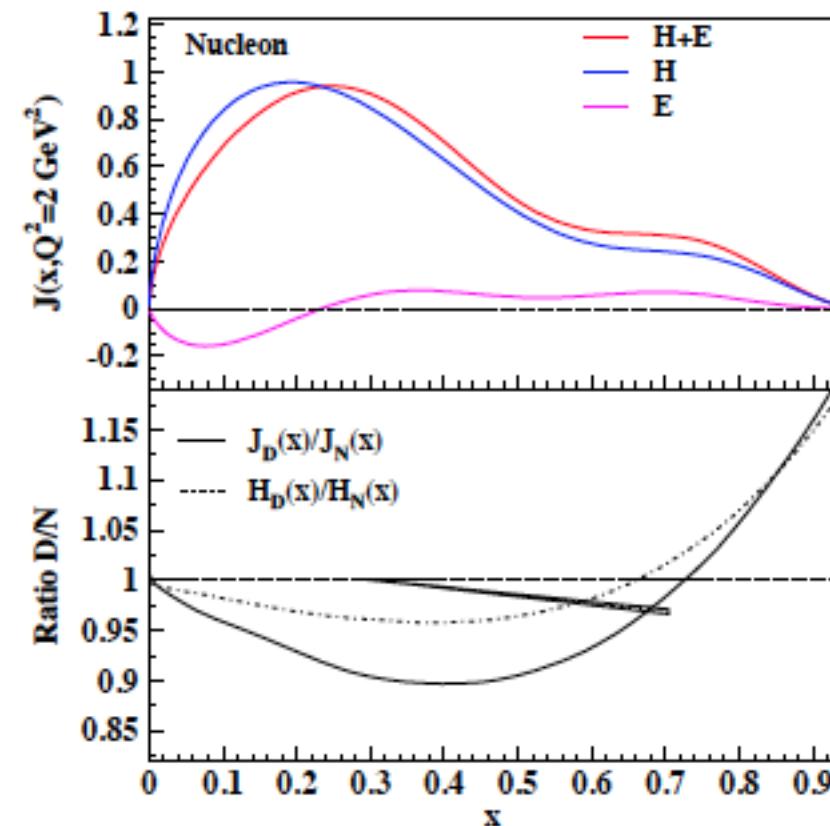
Predictions for b_2^d , A_{zz}^d

Phys.Lett. **B398** 245 (1997)



- Disentangling possible at lower x .
- (HERMES)Results shown are for $Q^2=10 \text{ GeV}^2$, errors are statistical.

OAM Sum Rule



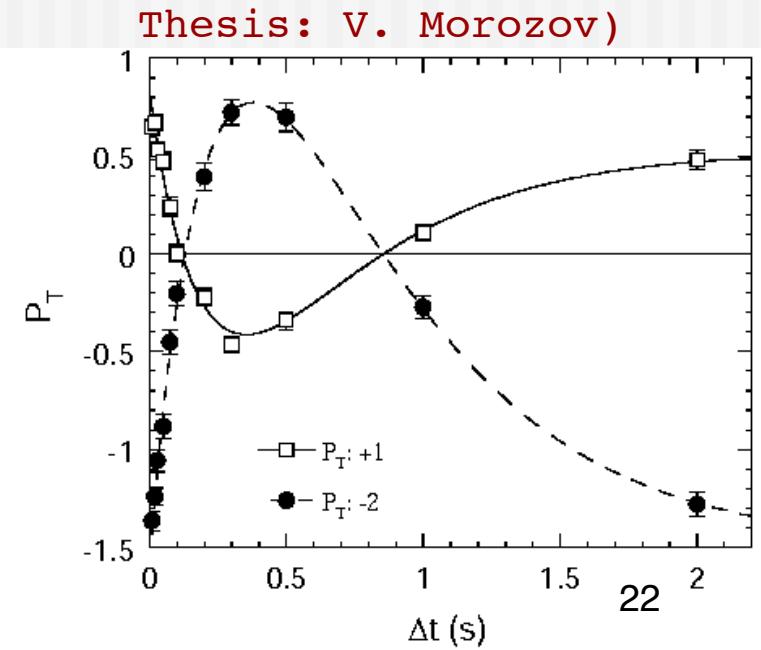
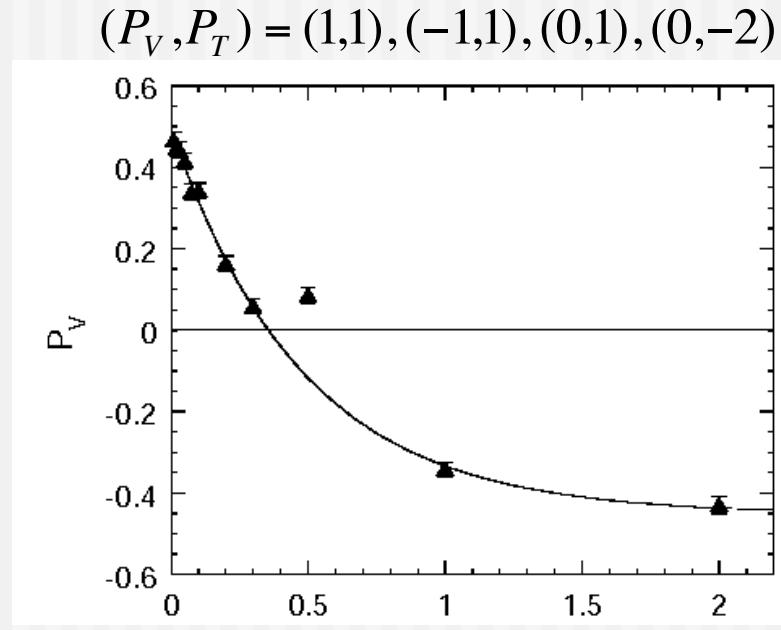
- OAM obtained from A_{UT} (vector pol.)
- Small, hatched area, for ratio, experimental (1109.6197 [hep-ph])
- b_1^d adds complementary information.
- Further development in progress.

$$b_1^d \equiv H_5(x, 0, 0)$$

S. K. Taneja, K. Kathuria, S.
Liuti, G. R. Goldstein Phys.Rev. D.
86 036008

Deuteron Beam Polarization Studies

- Studied deuteron spin manipulation with a 270 MeV vertically polarized beam stored in IUCF storage ring. Similar study done at COSY.
- Beam Fast RF cycled through 4 vertical polarization states (to reduce systematic errors).
- Spin-1 linear combination: Flip by bunches or extract at experiment.
- Simulation in progress for MEIC (figure-8) concept.

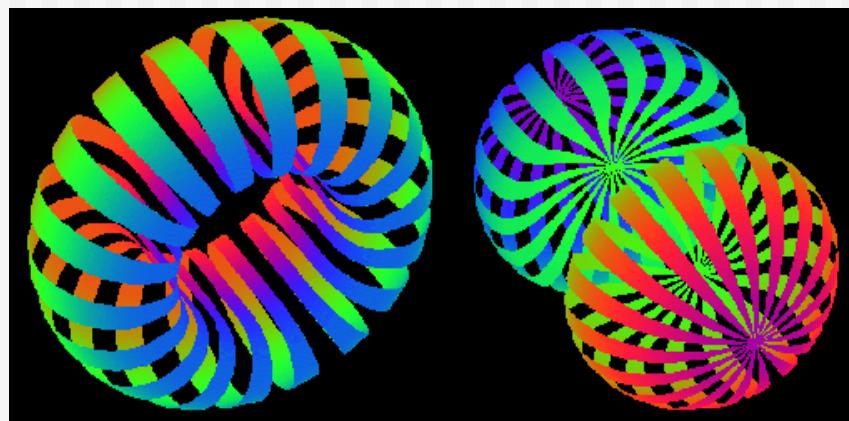


Summary

- Tensor Polarized deuteron provides Spin-1 quark/nuclear system.
- Spin-1 produces 4 new SSFs.
- HERMES measurement, complementary proposal at Jlab.
- Access to lower x , with tensor polarized deuteron, could open new physics capabilities.
- Study underway for polarized deuteron beam for MEIC.
- Physics needs some considerable development.

*Many thanks to C. Weiss, V. Morozov, S. Liuti

Support Slides



Spin-1 Structure Functions

Nucleon	Deuteron
b_1 ...	$\frac{1}{2} \sum_q e_q^2 [2q_\uparrow^0 - (q_\uparrow^1 + q_\downarrow^{-1})]$

From reflection-symmetry

$$q_\uparrow^m = q_\downarrow^{-m}$$

b_1 , d.n.e for spin-1/2 and vanishes in absence of nuclear effects.

In relative S-state b_1 describes difference between helicity-0 and averaged nonzero.

$$q^0 = (q_\uparrow^0 + q_\downarrow^0) = 2q_\uparrow^0$$

$$q^1 = (q_\uparrow^1 + q_\downarrow^1) = (q_\uparrow^1 + q_\uparrow^{-1})$$

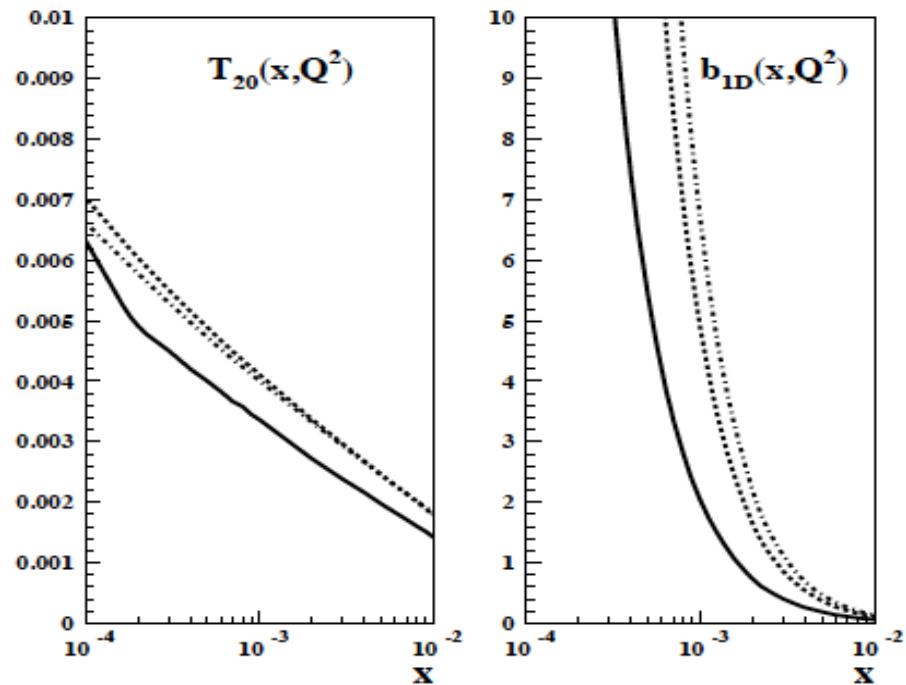
b_1 depends only spin-averaged distributions

$$\frac{1}{2} \sum_q e_q^2 [q^0 - q^1]$$

Hoodbhoy, Jaffe, Manohar (1989)

Tens. Pol. Scattering at low x

L. Frankfurt, V. Guzey, M. Strikman
Mod. Phys. Lett. **A21**(2006) 23–40



Solid curve: $Q^2 = 2 \text{ GeV}^2$
Dashed: 5 GeV^2
Dotted: 10 GeV^2

$$T_{20} = 2 \left(\frac{\sigma^+ - \sigma^0}{\sigma^+ + \sigma^0} \right)$$

$$b_1^d(x, Q^2) = -\frac{F_2^d(x, Q^2)}{2x} T_{20}(x, Q^2)$$

b_1 Collaboration

K. Allada, A. Camsonne, *J.-P. Chen*, A. Deur,
D. Gaskell, M. Jones, C. Keith, C. Keppel, D. Mack,
J. Piece, *P. Solvignon*, S. Wood, J. Zhang

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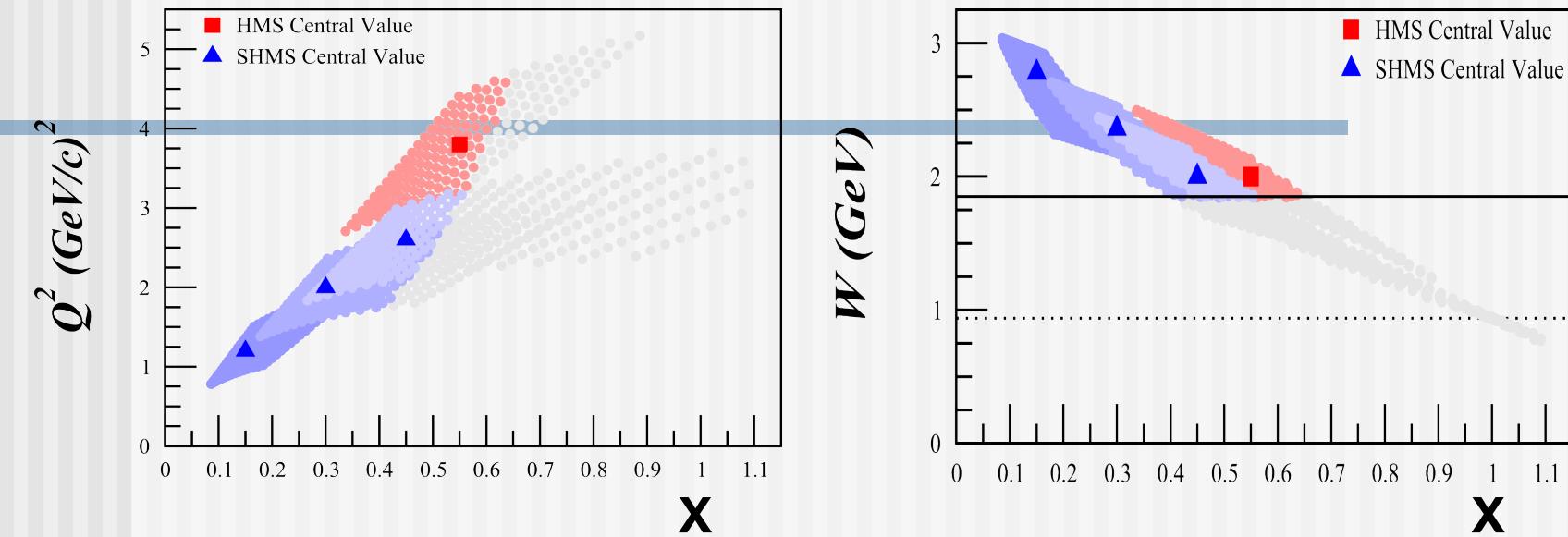
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Hyekoo Kang, Yoomin Oh
Seoul National University

Y. X. Ye, P. J. Zhu
University of Science and Technology of China

Abdellah Ahmidouch
North Carolina A&T State University

Kinematics



Detector	x	Q^2 (GeV^2)	W (GeV)	$E_{e'}$ (GeV)	$\theta_{e'}$ (deg.)	θ_q (deg.)	Rates (kHz)	Time (Days)
SHMS	0.15	1.21	2.78	6.70	7.35	11.13	1.66	6
SHMS	0.30	2.00	2.36	7.45	8.96	17.66	0.79	9
SHMS	0.45	2.58	2.00	7.96	9.85	23.31	0.38	15
HMS	0.55	3.81	2.00	7.31	12.50	22.26	0.11	30

Technically Challenging Experiment

I) Systematics

TAC : Important to control measured false asymmetries to better than 6×10^{-4} .

TAC : “We believe this is possible with a combination of upgrades to Hall C infrastructure and sufficient commitment by the collaboration to control the unusual systematic issues of this experiment.”

II) Development of Large Tensor Polarizations

- 1) Incremental : Higher B field (7.55T, 212 GHz), better fridge, pumps, tempering, FM'ing.
- 2) RF Saturation : Has been demonstrated to produce large P_{zz} (30%).
For full saturation $P_{zz} \approx P_z$, so range of expectation is about 20-50%.
- 3) Additional Microwave Source: No theoretical limit to P_{zz} , but expensive and unproven.
- 4) Adiabatic Fast Passage :

Systematics

$\delta\xi$

Charge Determination

$< 2 \times 10^{-4}$, mitigated by thermal isolation of BCMs and addition of 1 kW Faraday cup

Luminosity

$< 1 \times 10^{-4}$, monitored by Hall C lumi

Target dilution and length step like changes observable in polarimetry

$< 1 \times 10^{-4}$

Beam Position Drift effect on Acceptance

$< 1 \times 10^{-4}$ (we can control the beam to 0.1 mm, raster over 2cm diameter)

Effect of using polarized beam

$< 2.2 \times 10^{-5}$, using parity feedback

Impact on the observable

$$\delta A_{zz} = \pm \frac{2}{f P_{zz} \sqrt{N_{cycles}}} \delta\xi$$

Dedicated team to systematics/false asymms

similar manpower requirement to g2p exp.
where we had several teams completely
separate from the polarized target effort.

Systematics

False Asymmetries from Time Dependent Drifts

False Asymmetries

Spec. $\langle x \rangle$	Hours	Stat. Err ($\times 10^{-3}$)	Cycles	$\delta A_{zz} (\times 10^{-3})$
0.15	144	2.6	12	4.3
0.30	216	3.0	9	4.9
0.45	360	3.7	15	3.8
0.55	720	4.1	36	2.4

Normalization Factors

Source	Relative Uncertainty
Polarimetry	8.0%
Dilution/Packing Fraction	4.0%
Radiative Corrections	1.5%
Charge Determination	1.0%
Detector Resolution and Efficiency	1.0%
Total	9.2%

Tensor Target opens new possibilities

Few Examples

Tensor Structure function b_2, b_3, b_4

Azimuthal Asymmetries b_4

Elastic e-D scattering

$$\begin{matrix} T_{20} \\ T_{11} \end{matrix}$$

D(e,e'p) Cross Section on Tensor Polarized Deuterium.
H. Anklin, W. Boeglin et al., PR97-102, PAC13 rated A-

X>1 Scattering, connection to SRCs : M. Sargsian et al.

D-Wave Components of Deuteron Wave function : S. Luiti et al.

Polarized Target

Dynamic Polarization of ND₃

5 Tesla or 7.5 Tesla

3 cm target length, 2 cm diameter

Longitudinally polarized

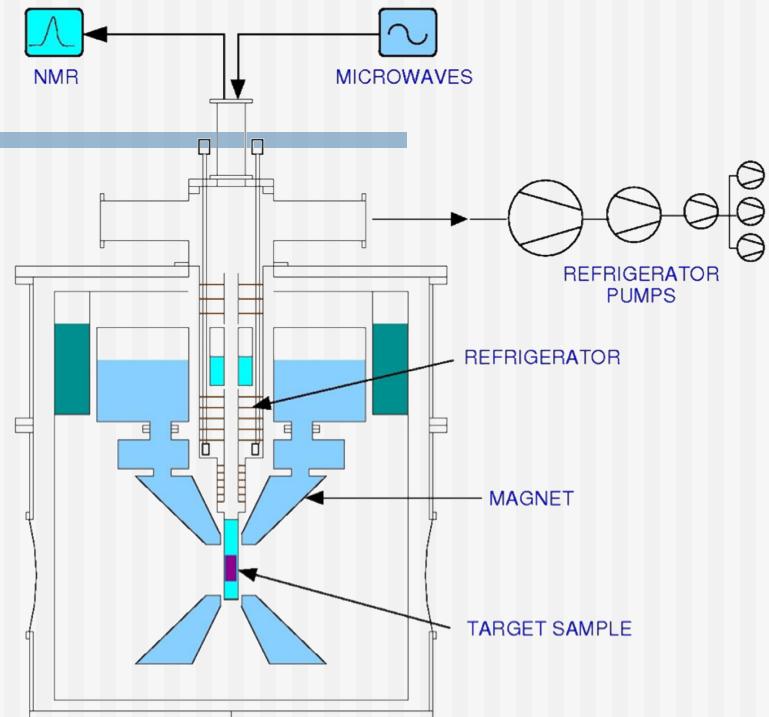


fig. courtesy of C. Keith

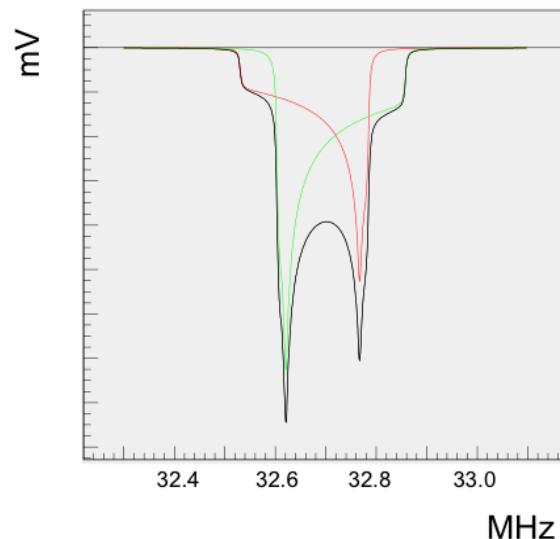
Run in Polarized and Unpolarized Mode.

B Field held at const value for both states

LHe level, temp. etc. held const for both states

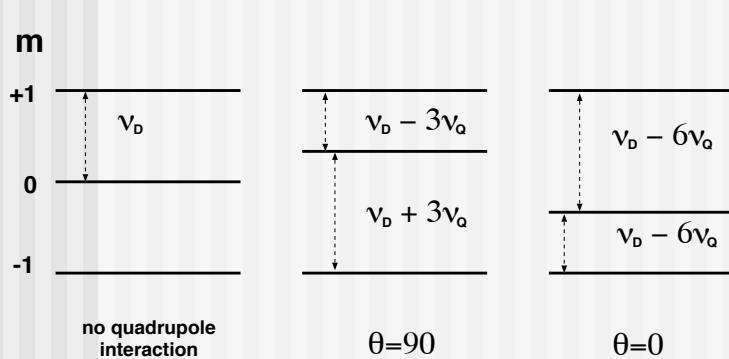
RF Saturation to Enhance P_{zz}

ND_3 Vector polarized



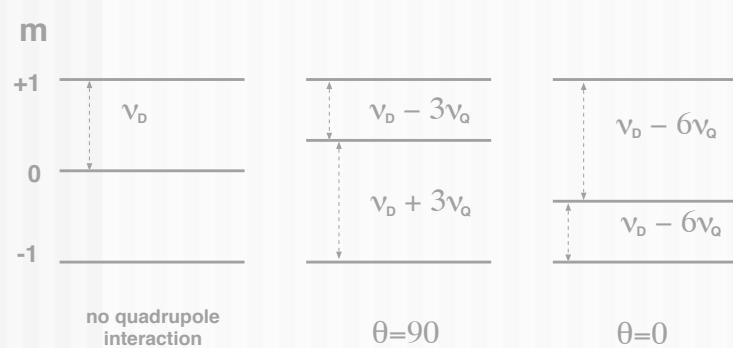
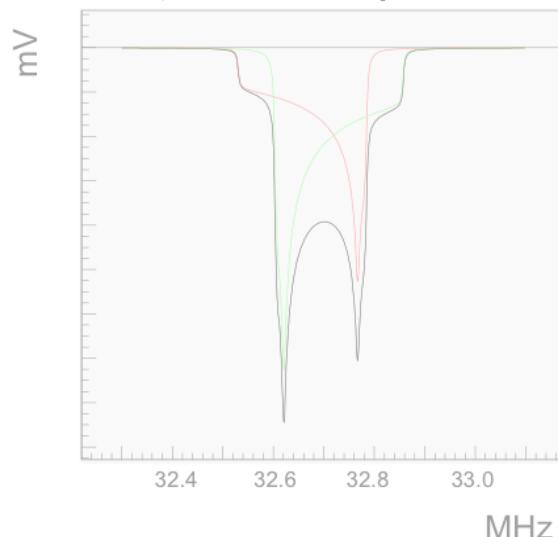
Vector Polarization \propto Sum of Peak Areas

Tensor Polarization \propto Diff of Peak Areas



RF Saturation to Enhance P_{zz}

ND_3 Vector polarized

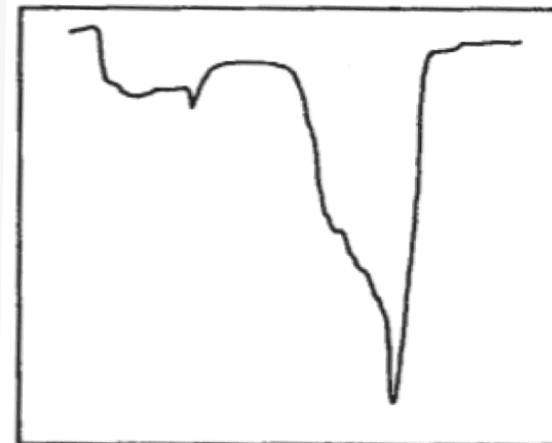


RF Saturate one of the peaks

kill the $m=0 \leftrightarrow m=-1$ transition, which enhances the $m=1 \leftrightarrow m=0$ transition

$P_{zz} = 20\%$ for 2.5T at 1K

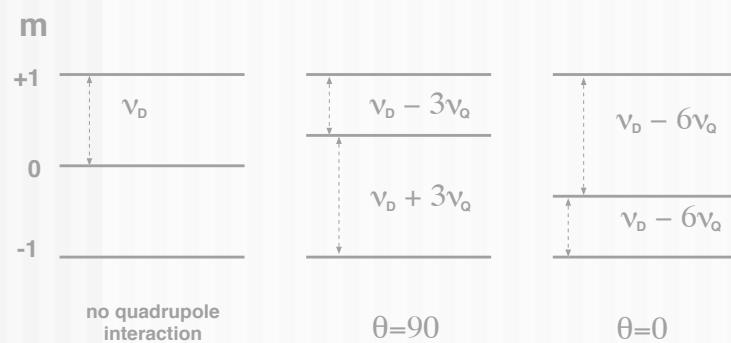
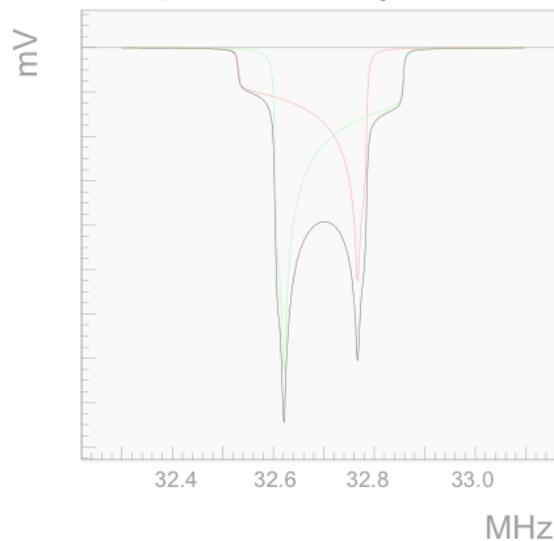
good results even with only 2.5 T field



Meyer and Schilling , 1984 Proceedings of the 4th
Int.
Workshop on Polarized Target Materials &
Techniques

RF Saturation to Enhance P_{zz}

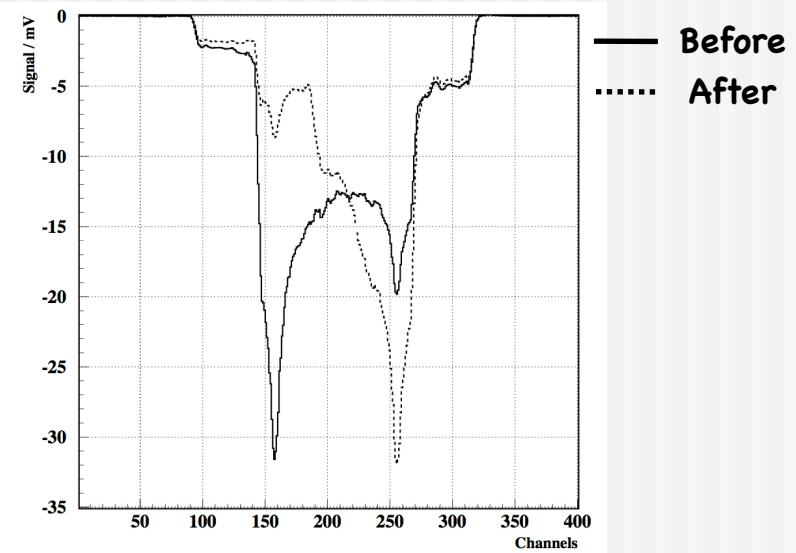
ND₃ Vector polarized



RF Saturate one of the peaks

kill the $m=0 \leftrightarrow m=-1$ transition, which enhances the $m=1 \leftrightarrow m=0$ transition

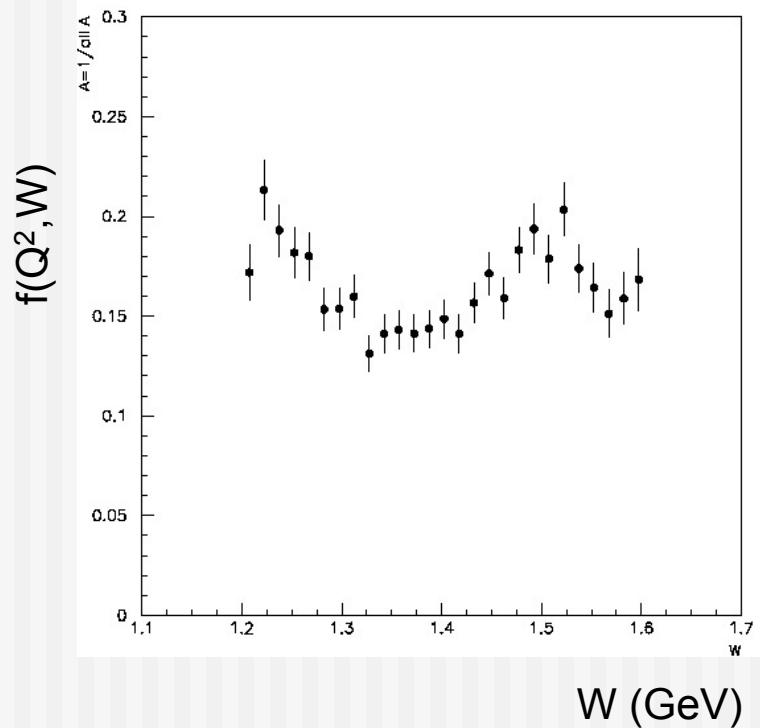
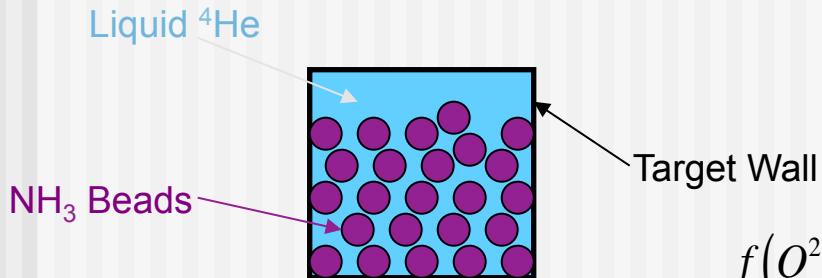
P_{zz} = 30% for 5.0 T at 1K



S. Bueltmann, et al (D. Crabb Lab) 1999.

Packing-Fractions & Dilution-Factors

- Packing Fraction essentially amount of material in target cup. This is a number.
- Dilution Factor (f) ratio of rates of free polarizable nucleons (proton) to all nucleons composing the target sample (nitrogen, NMR coils, ...). This is kinematics dependant.
- Need Packing Fraction and Dilution Factor for each target load used during running of experiment.



$$f(Q^2, W) = \frac{N_1 \sigma_1(Q^2, W)}{N_{14} \sigma_{14}(Q^2, W) + N_1 \sigma_1(Q^2, W) + \sum N_A \sigma_A(Q^2, W)}$$