

(Dovez 241 Lu's threads)

09/20/12

$$\sigma(l, P_2, P_{22}) \equiv \frac{d\sigma}{d\Omega}(l, P_2, P_{22}) = \Sigma + 4 \cdot \Delta$$

Asymmetry:

$$A = \frac{\sigma(+l, P_2, P_{22}) - \sigma(-l, P_2, P_{22})}{\sigma(+l, P_2, P_{22}) + \sigma(-l, P_2, P_{22})} =$$

$$= \frac{\Sigma + 4\Delta - (\Sigma - 4\Delta)}{\Sigma + 4\Delta + \Sigma - 4\Delta} = \frac{h \cdot \Delta}{\Sigma}$$

$$\Sigma = \frac{d\sigma}{d\Omega} \cdot [1 + \Gamma]; \quad \Delta = \frac{d\sigma}{d\Omega} \cdot \Lambda$$

$$A = \frac{h \cdot \frac{d\sigma}{d\Omega} \cdot \Lambda}{\frac{d\sigma}{d\Omega} \cdot [1 + \Gamma]} = 4 \cdot \frac{\Lambda}{[1 + \Gamma]}$$

$$\Delta = P_2 \cdot \left[\frac{\sqrt{3}}{\sqrt{2}} \cdot \cos \theta^* T_{10}^e - \sqrt{3} \cdot \cos \theta^* \cdot \cos \varphi^* T_{11}^e \right]$$

$$\underline{\Gamma} = P_{zz} \cdot \left[\frac{1}{\sqrt{2}} P_2^0(\cos \theta^*) T_{20} - \frac{1}{\sqrt{3}} P_2^1(\cos \theta^*) \cos \psi^* T_{21} + \frac{1}{2\sqrt{3}} P_2^2(\cos \theta^*) \cos 2\psi^* T_{22} \right] =$$

$$= P_{zz} \left[\frac{1}{\sqrt{2}} \cdot \frac{1}{2} \cdot (3 \cos^2 \theta^* - 1) T_{20} - \frac{1}{\sqrt{3}} \cdot \frac{\sqrt{3}}{2} \cdot \cos 2\theta^* \cos \psi^* T_{21} + \frac{1}{2\sqrt{3}} \cdot \frac{\sqrt{3}}{2} \cdot \cos^2 \theta^* \cdot \cos(2\psi^*) T_{22} \right] =$$

$$= \frac{P_{zz}}{2} \left[\frac{3 \cos^2 \theta^* - 1}{\sqrt{2}} \cdot T_{20} - \sqrt{3} \cdot \cos 2\theta^* \cos \psi^* T_{21} + \sqrt{3} \cos^2 \theta^* \cos(2\psi^*) T_{22} \right]$$

$$\boxed{\begin{aligned} P_2 &= \mu_+ - \mu_- \\ P_{zz} &= 1 - 3\mu_0 \end{aligned}}$$

$(\cos\theta^2)$ according to ZLZ : In ${}^3\text{He}$, d always

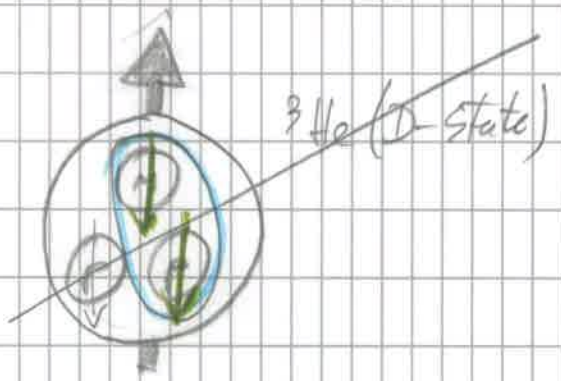
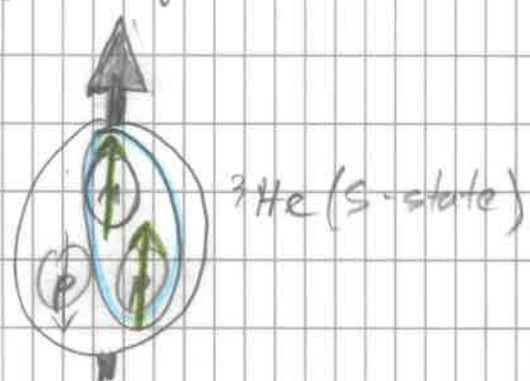
z = only vector polarization, necessary, that can not be found in M_0 state. So:

$M_+, M_- \neq 0, \underline{M_0 = 0}$. Ergo:

$P_{zz} = 1;$

$= P_{{}^3\text{He}} = \frac{2}{3}$

The same reason as $\cos\theta^2$ (for the proton case)



$$1 \text{ GeV}/c = \frac{1}{1.957 \cdot 10^{-16} \text{ m}} = \frac{5.1098 \cdot 10^{15}}{10^{15} \text{ fm}} =$$

$$= \frac{5.1098}{\text{fm}} \Rightarrow \left[\frac{1}{\text{fm}} \right] = \frac{1}{5.1098} \text{ GeV}/c = \underline{\underline{0.1957 \text{ GeV}/c}}$$

Chapter 3

Overview of the Experiment

The experiment described in this thesis is the first experiment performed at the internal-target facility of NIKHEF. The experimental setup allows the measurement of the $(e, e'p)$ and $(e, e'd)$ reactions for tensor-polarized deuterium. In this chapter an overview of the experiment is given, together with its major design considerations. We give a general discussion of the kinematics, the electron storage ring and the principle of the internal target. In order to optimize the detection system, a Monte Carlo code was written which is briefly described in section 3.2. Furthermore, the results of a study of the spatial distribution of the stored electron beam are presented in section 3.3. In section 3.4 the design considerations for the polarized deuterium target will be discussed. A detailed description of the polarized source, internal target, polarimeters and detection system will be given in the next chapters.

3.1 Description of the Experiment

3.1.1 Planned measurements

The goal of our experiment was to measure the spin-dependent observables for both ${}^2\hat{H}(e, e'd)$ elastic and ${}^2\hat{H}(e, e'p)$ quasi-elastic scattering over a large kinematic range. The general expression of the cross section for polarized deuteron electro-disintegration has the following form [6]:

$$\sigma = \sigma_0 \left[1 + \sqrt{3} P_z \sin \theta_d \sin \phi_d i T_{11} + \frac{1}{\sqrt{2}} P_{zz} \left(\frac{3 \cos^2 \theta_d - 1}{2} T_{20} - \sqrt{\frac{3}{2}} \sin 2\theta_d \cos \phi_d T_{21} + \sqrt{\frac{3}{2}} \sin^2 \theta_d \cos 2\phi_d T_{22} \right) \right]. \quad (3.1)$$

Here, σ_0 is the unpolarized cross section, P_z and P_{zz} are the degree of vector and tensor polarization defined as $P_z = n_+ - n_-$ and $P_{zz} = 1 - 3n_0$, where n_+ , n_0 , and n_- are

er 3

View of the Experiment

described in this thesis is the first experiment performed at KHEF. The experimental setup allows the measurement of tensor-polarized deuteron. In this chapter, together with its major design considerations, the electron storage ring and the detection system, a Monte Carlo simulation in section 3.2. Furthermore, the results of the polarized electron beam are presented in section 3.3. The polarized deuteron target will be discussed for the polarized deuteron target in the next chapters.

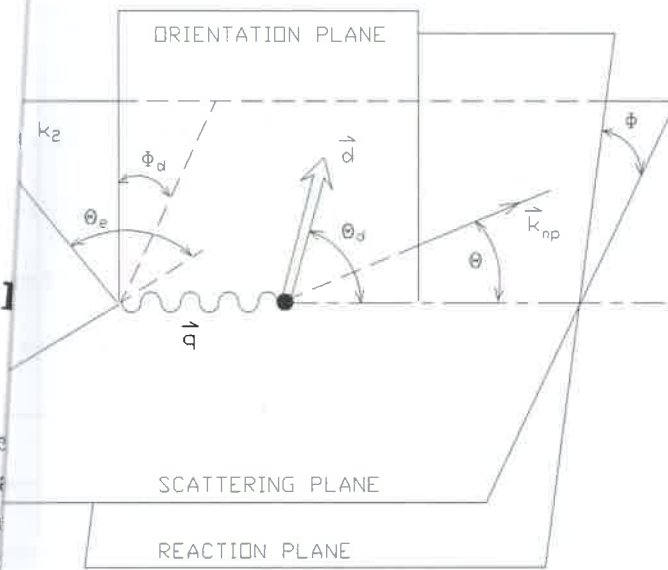
View of the Experiment

Measurements

The experiment was to measure the spin-dependent tensor polarizations of the deuteron in $(e, e'p)$ quasi-elastic scattering over a large range of the cross section for polarized deuteron electron scattering.

$$\sqrt{3} P_z \sin \theta_d \sin \phi_d iT_{11} + \frac{1}{\sqrt{2}} P_{zz} \left(\frac{3 \cos^2 \theta_d}{2} \right) + \sqrt{\frac{3}{2}} \sin 2\theta_d \cos \phi_d T_{21} + \sqrt{\frac{3}{2}} \sin^2 \theta_d \cos 2\phi_d$$

and cross section, P_z and P_{zz} are the degree of polarization, $P_z = n_+ - n_-$ and $P_{zz} = 1 - 3n_0$, where



of exclusive electron-deuteron scattering with polarized electron target. The relative $n - p$ momentum is denoted by \vec{k}_{np} and ϕ and the deuteron orientation axis by \vec{d} characterized by

of the final np relative momentum with respect to the photon helicity and $P_1^d = P_{10}^d = \sqrt{\frac{3}{2}}(n_+ - n_-)$, $P_2^d = P_{20}^d = \frac{1}{\sqrt{2}}(1 - 3n_0)$ tensor polarizations of the deuteron.

merical components of the virtual-photon density matrix. In a coordinate system where the z -axis is parallel to \vec{q} and the y -axis is parallel to $\vec{e} \times \vec{e}'$, they can

$$\rho_L = Q_\mu^2 \frac{\xi^2}{2\eta} \quad (2.70)$$

$$\rho_T = \frac{1}{2} Q_\mu^2 \left(1 + \frac{\xi}{2\eta} \right) \quad (2.71)$$

$$\rho_{LT} = Q_\mu^2 \frac{\xi}{\eta} \sqrt{(\xi + \eta)/8} \quad (2.72)$$

$$\rho_{TT} = Q_\mu^2 \frac{\xi}{4\eta} \quad (2.73)$$

$$\rho'_{LT} = \frac{1}{2} Q_\mu^2 \xi \sqrt{2\eta} \quad (2.74)$$

$$\rho'_T = \frac{1}{2} Q_\mu^2 \sqrt{(\xi + \eta)/\eta}, \quad (2.75)$$

where f is usually called the recoil factor

$$f = 1 + \frac{2E_e}{M_d} \sin^2 \frac{\theta_e}{2}. \quad (2.23)$$

We define Q , related to the space-like four-momentum q transferred from the electron to the target nucleus, by

$$Q^2 = -q_\mu^2 = 4 \frac{E_e^2}{f} \sin^2 \frac{\theta_e}{2}. \quad (2.24)$$

For completeness, the kinetic energy and the angle of the recoil deuteron with respect to the incident electron are given below:

$$T_d = E_e - E_{e'} = \frac{Q^2}{2M_d}, \quad (2.25)$$

$$\sin^2 \theta_d = \frac{\cos^2(\theta_e/2)}{f + (T_d/M_d)^2 \sin^2(\theta_e/2)}. \quad (2.26)$$

The general features of e-d scattering have been investigated by many authors [53, 54]. By using the first Born approximation (one-photon exchange approximation) and imposing Lorentz and gauge invariance, the differential cross section can be written as

$$\frac{d\sigma}{d\Omega} = \frac{1}{f} \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} S, \quad (2.27)$$

where

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} = \frac{\alpha^2 \cos^2(\theta_e/2)}{4E_e^2 \sin^4(\theta_e/2)} \quad (2.28)$$

describes the scattering of an electron from a point-like spinless particle (α is the fine structure constant), and

$$S = A(Q^2) + B(Q^2) \tan^2 \frac{\theta_e}{2}, \quad (2.29)$$

originates from the electromagnetic structure of the deuteron. As a consequence of parity and time-reversal invariance, the structure functions A and B are in turn given by three elementary electromagnetic form factors:

$$A(Q^2) = G_C^2(Q^2) + \frac{8}{9} \eta^2 G_Q^2(Q^2) + \frac{2}{3} \eta G_M^2(Q^2), \quad (2.30)$$

$$B(Q^2) = \frac{4}{3} \eta (1 + \eta) G_M^2(Q^2) \quad (2.31)$$

with $\eta = Q^2/4M_d^2$.

If a target with vector polarization P_z and tensor polarization P_{zz} is used in conjunction with an ultra-relativistic beam of polarized electrons with helicity h , then the cross section takes the following form [3, 55, 56]

$$\begin{aligned} \frac{d\sigma}{d\Omega}(P_z, P_{zz}, h) &= \Sigma + h\Delta \\ \Sigma &= \frac{d\sigma}{d\Omega}[1 + \Gamma] \\ \Gamma &= P_{zz}\left[\frac{1}{\sqrt{2}}P_2^0(\cos\theta_d)T_{20}(Q, \theta_e) - \frac{1}{\sqrt{3}}P_2^1(\cos\theta_d)\cos\phi_d T_{21}(Q, \theta_e)\right. \\ &\quad \left. + \frac{1}{2\sqrt{3}}P_2^2(\cos\theta_d)\cos 2\phi_d T_{22}(Q, \theta_e)\right] \\ \Delta &= \frac{d\sigma}{d\Omega}P_z\left[\frac{\sqrt{3}}{\sqrt{2}}\cos\theta_d T_{10}^e(Q, \theta_e) - \sqrt{3}\sin\theta_d\cos\phi_d T_{11}^e(Q, \theta_e)\right], \end{aligned} \quad (2.32)$$

where the notation has been chosen such that the connection with the Madison Convention [57] is made explicit. The responses T_{ij} can be expressed in terms of the three form factors

$$T_{20} = -\frac{1}{\sqrt{2}S}\left[\frac{8}{3}\eta G_C G_Q + \frac{8}{9}\eta^2 G_Q^2 + \frac{1}{3}\eta[1 + 2(1 + \eta)\tan^2\frac{\theta_e}{2}]G_M^2\right] \quad (2.33)$$

$$T_{21} = +\frac{2}{\sqrt{3}S}\eta[\eta + \eta^2\sin^2\frac{\theta_e}{2}]^{1/2}G_M G_Q \sec\frac{\theta_e}{2} \quad (2.34)$$

$$T_{22} = -\frac{1}{2\sqrt{3}S}\eta G_M^2 \quad (2.35)$$

$$T_{10}^e = \left(\frac{\sqrt{3}}{2}\frac{2}{3}S\right)\eta[(1 + \eta)(1 + \eta\sin^2\frac{\theta_e}{2})]^{1/2}G_M^2 \sec\frac{\theta_e}{2} \tan\frac{\theta_e}{2} \quad (2.36)$$

$$T_{11}^e = \frac{\sqrt{3}A}{2S}\eta^{1/2}[\eta(1 + \eta)]^{1/2}G_M(G_C + \frac{1}{3}\eta G_Q)\tan\frac{\theta_e}{2}, \quad (2.37)$$

where the suffix e in T_{10}^e and T_{11}^e denotes that polarized electron beam is required. This notation is chosen to avoid confusion with the case of the deuteron vector analyzing power using an unpolarized electron beam. In the one-photon exchange approximation of electron-deuteron scattering, and as a result of time-reversal invariance, the deuteron vector analyzing power is identically zero when using an unpolarized electron beam [58].

The angles θ_d and ϕ_d specify the quantization axis of the target with respect to the three-momentum transfer as shown in Fig. 2.7. The associate Legendre polynomials, sometimes expressed as d-functions, multiplying the rank-two responses are

$$P_2^0 = d_{00}^2 = \frac{3}{2}\cos^2\theta_d - \frac{1}{2} \quad (2.38)$$

$$P_2^1 = -\sqrt{6}d_{10}^2 = 3\cos\theta_d\sin\theta_d \quad (2.39)$$

$$P_2^2 = 2\sqrt{6}d_{20}^2 = 3\sin^2\theta_d. \quad (2.40)$$

Subject: RE: Elastic Deuteron Asymmetry
From: Zilu Zhou <zzhou2@slb.com>
Date: 12/19/2011 04:21 PM
To: Miha Mihovilovic <miha.mihovilovic@ijs.si>
CC: Simon Sirca <simon.sirca@fmf.uni-lj.si>

Miha,

Sorry for a late reply. Weekend always busy...

Before we start, I have to warn you that trying to treat $\text{He}_3(e, e'd)p$ as if it is a vector polarized deuteron and then get A_{ed} for elastic asymmetry would be a mistake. A_{ed} elastic is almost near zero as I showed on page 21 in my thesis, which is completely different from e-p elastic.

The vector asymmetry in $\text{He}_3(e, ed)p$ does not come from e-d elastic, but rather with the Faddeev calculation including He_3 ground state properties and isospin reaction process. These interferences cannot be simplified in a intuitive picture like a e-d elastic asymmetry. Whileas $\text{He}_3(e, e'p)d$ or two-body break-up, the isospin interference is negligible, then the e-p elastic picture stands. Then the simplicity prevails, as we have discussed.

Now back to the P_z and P_{zz} calculation:

In Deuteron, this is exact complication. I have both P_z and P_{zz} equations on page 34, or equivalent in some of the publications. For both, there are always dilutions from the n_0 state, as you said. So, strictly speaking, all the P_z measurements from many experiments were wrong without measuring the n_0 state. But everybody did use unpolarized averages in the dominator as a short-cut. This was also one of the arguments that one has to study tensor polarization, in my thesis on page 45.

Other than this, everything is straight forward with the equations on page 34, or page 48. For polarized he_3 , you have + and - states, and then if deuteron is 100% following the he_3 polarization, P_z is 100%.

For P_{zz} , since you do not have n_0 state, then it is fixed at +1, 100%. There is no -2 100% in your case.

Hope this helps. Good luck. Let me know if you have more questions.

Zilu

-----Original Message-----

From: Miha Mihovilovic [<mailto:miha.mihovilovic@ijs.si>]
Sent: Friday, December 16, 2011 11:45 AM
To: Zilu Zhou
Cc: Simon Sirca
Subject: Elastic Deuteron Asymmetry

Dear Zilu,

I have few questions regarding my asymmetry analysis at low missing momenta (that we discusses on Monday), and would like to ask you for your help.

I am trying to calculate double polarized asymmetry for elastic scattering of electrons on deuterium. I am using formalism from your thesis (chapter 2). Unfortunately I do not know, how to properly consider polarization of the deuterium target in my asymmetry expression. It has a following form:

$$A_{\text{elastic}} = P_z * (\text{helicity dependent part of the cross section}) / (\sigma_0 * [1 + P_{zz} * (\text{helicity independent part})])$$

Here P_z and P_{zz} are vector and tensor polarizations of the target. If my target would have 100% polarization, what do I need to consider for P_z and P_{zz} ? Is it possible to take the tensor part out of the equation or do I need to consider it properly?

Simon and I also examined Passchier's paper with hopes to find an answer to this question (although it is not about elastic scattering). In that experiment you were also using vector-polarized deuterium. Unfortunately, we were not able to reproduce formula 3 (which shows how to calculate A_{Ved} from the measured asymmetry) from the expression (2) for the cross-section. In the paper you get very simple relation:

$$A = h * P_{ld} * A_{\text{Ved}}$$

In our calculation we always got an additional tensor term in the denominator. We assumed that in the experiment you were changing vector-polarization but keeping tensor-polarization fixed the whole time. Is this true? The only way that we were able to reproduce formula in Passchier's paper was to consider two additional terms in the denominator with only tensor polarization (vector polarization = 0).

Is it enough to measure asymmetry at only one tensor polarization and just flipping vector-polarization, in order to extract A_{Ved} , or do you need measurements at two different tensor polarizations?

Thank you very much for all your help and suggestions.

Sincerely,
Miha

Subject: RE: Deuteron FF parameterization
From: Ball Jacques <jacques.ball@cea.fr>
Date: 12/16/2011 11:56 AM
To: Miha Mihovilovic <miha.mihovilovic@ijs.si>

Dear Miha,

Thank you for the interest in our past work. The information you asked for can be found at the following web page :

<http://irfu.cea.fr/Sphn/T20/Parametrisations/>

best regards

Jacques Ball

Irfu / SPhN

CEA, Centre de Saclay

F-91191 Gif-sur-Yvette

Tel : 33 1 69-08-87-19

Mail : jacques.ball@cea.fr

Webpage : [http://irfu.cea.fr/Sphn/Phoce/Vie des labos/Ast/ast_groupe.php?id_groupe=495](http://irfu.cea.fr/Sphn/Phoce/Vie_des_labos/Ast/ast_groupe.php?id_groupe=495)

-----Message d'origine-----

De : Miha Mihovilovic [<mailto:miha.mihovilovic@ijs.si>]

Envoyé : vendredi 16 décembre 2011 11:18

À : Ball Jacques

Cc : iball@ilab.org; miham@ilab.org; Simon Sirca

Objet : Deuteron FF parameterization

Dear prof. Ball,

I am a PhD student in Hall A working on the analysis for the E05-102 experiment. In my analysis I am trying to determine asymmetry for the elastic electron-deuteron scattering. To calculate this asymmetry I need to know deuteron form-factors. For the parameterization of the form-factors I would like to use one of the parameterizations from your paper (Eur. Phys. J. A 7, 421 (2000)).

I would like to ask you, if you could send me parameters, that you considered in your fits.

Sincerely,
Miha Mihovilovic

Paramétrisations

The results of the Parameterization I and II, detailed in our article, "Phenomenology of the deuteron electromagnetic form factors" (European Physical Journal A, vol 7 (2000) 421), which were worked out with the World data set are updated on this page.

Parameterization I

$$G_x(Q^2) = G_x(0) [1 - (Q/Q_x(0))^2] \left[1 + \sum_{i=1}^5 a_x(i) Q^{(2i)} \right]^{(-1)}$$

with $x = c, q, m$ and Q in Fm^{-1}

Gc	Gq					Gm						
	Qq(0)	aq(1)	aq(2)	aq(3)	aq(4)	aq(5)	Qm(0)	am(1)	am(2)	am(3)	am(4)	am(5)
Qc(0)	4.21						8.1					7.37
ac(1)	6.740E-01						8.796E-01					5.804E-01
ac(2)	2.246E-02						-5.656E-02					8.701E-02
ac(3)	9.806E-03						1.933E-02					-3.624E-03
ac(4)	-2.709E-04						-6.734E-04					3.448E-04
ac(5)	3.793E-06						9.438E-06					-2.818E-06

$$G_c(0) = 1.$$

$$G_q(0) = 25.83$$

$$G_m(0) = 1.714$$

To see the full results listing of the form factors , click here : [Parameterization I](#)

[Parameterization II](#)

The notations are those used by Kobushkin and Syamtomov in "Deuteron electromagnetic form factors in the transitional region between nucleon-meson and quark gluons pictures" (Phys. At. Nucl 58, 1477, (1995)).

$a(i)$ (fm^{-2})	$b(i)$ (fm^{-1})	$c(i)$	$\alpha^2(i)$ (fm^{-2})	$\beta^2(i)$ (fm^{-2})	$\gamma^2(i)$ (fm^{-2})
1.57057	0.07043	-0.16577	1.52501	43.67795	1.87055
12.23792	0.14443	0.27557	8.75139	30.05435	14.95683
-42.04576	-0.27343	-0.05382	15.97777	16.43075	28.04312
27.92014	0.05856	-0.05598	23.20415	2.80716	41.12940

$$\mu^{(e)} = 240.86 \text{ MeV}/c \quad \mu^{(6)} = 29.138 \text{ MeV}/c \quad \mu^{(v)} = 426.93 \text{ MeV}/c \quad \delta = 898.52 \text{ MeV}/c$$

To see the full results listing of the form factors , click here : Parameterization II

back to t20 home page

Opomba: Dames dne 09/21/20 ma je kutovna prva polovica
 vaze kongre. Hicelw duplez na kega vavchla samovelfw
 na prebravljua! Cepow na, bi hle lildw mewa mafie!

314e

09/21/12

$$|j = \frac{1}{2}, m_j = \frac{1}{2}\rangle_{314e} = \sqrt{\frac{2}{3}} |1, 1\rangle_d | \frac{1}{2}, -\frac{1}{2} \rangle_p - \sqrt{\frac{1}{3}} |1, 0\rangle_d | \frac{1}{2}, +\frac{1}{2} \rangle_p$$

↑
 Tvorba nových, dvoje celostrojových
 spinů. Hodnota není předemná.

$$P_d = \langle \frac{1}{2}, \frac{1}{2} | \frac{1}{\sqrt{2}} \langle 1, 1 | P_{314e} | \frac{1}{2}, \frac{1}{2} \rangle =$$

$$= \left(\sqrt{\frac{2}{3}} \langle 1, 1 | \langle \frac{1}{2}, -\frac{1}{2} | - \sqrt{\frac{1}{3}} \langle 1, 0 | \langle \frac{1}{2}, +\frac{1}{2} | \right) P_{314e} \cdot$$

$$\left(\sqrt{\frac{2}{3}} \cdot 1 \cdot |1, 1\rangle | \frac{1}{2}, -\frac{1}{2} \rangle - \sqrt{\frac{1}{3}} \cdot 0 \right) =$$

$$= \frac{2}{3}$$

TO-DO: + Převést na otázky produktů, ve
 analýze na stránce.

+ Pogleď, kdy je nejvyšší rychlost
 ušetřování elaktů. Kterou identitu
 produkt je nejvyšší vztahem?

ThE: 12.5, Th*: 68.2544, Ph*: 8.53774e-07

Gc(0.27063 GeV**2) = 0.0617677

Gq(0.27063 GeV**2) = 2.3764

Gm(0.27063 GeV**2) = 0.166093

A(0.27063 GeV**2) = 0.00602567

B(0.27063 GeV**2) = 0.000721013

T20(0.27063 GeV**2, 70deg) = -1.07945

T21(0.27063 GeV**2, 70deg) = 0.233359

T22(0.27063 GeV**2, 70deg) = -0.0240091

T10(0.27063 GeV**2, 70deg) = 0.0587881

T11(0.27063 GeV**2, 70deg) = 0.182052

LongAsymmetry: ThE: 10.5, Th*: 71.5778, Ph*: 0, Q2: 0.192864, AL: -0.0212088

TransAsymmetry: ThE: 10.5, Th*: 161.578, Ph*: 0, Q2: 0.192864, AT: -0.0193912

LongAsymmetry: ThE: 11, Th*: 70.7393, Ph*: 0, Q2: 0.211177, AL: -0.0228927

TransAsymmetry: ThE: 11, Th*: 160.739, Ph*: 0, Q2: 0.211177, AT: -0.0233552

LongAsymmetry: ThE: 11.5, Th*: 69.9058, Ph*: 0, Q2: 0.23025, AL: -0.024564

TransAsymmetry: ThE: 11.5, Th*: 159.906, Ph*: 0, Q2: 0.23025, AT: -0.0276787

LongAsymmetry: ThE: 12, Th*: 69.0774, Ph*: 0, Q2: 0.250072, AL: -0.0261893

TransAsymmetry: ThE: 12, Th*: 159.077, Ph*: 0, Q2: 0.250072, AT: -0.0321579

LongAsymmetry: ThE: 12.5, Th*: 68.2544, Ph*: 8.53774e-07, Q2: 0.27063, AL: -0.0277243

TransAsymmetry: ThE: 12.5, Th*: 158.254, Ph*: 0, Q2: 0.27063, AT: -0.0364951

LongAsymmetry: ThE: 13, Th*: 67.4367, Ph*: 0, Q2: 0.291914, AL: -0.0291125

TransAsymmetry: ThE: 13, Th*: 157.437, Ph*: 0, Q2: 0.291914, AT: -0.040332

LongAsymmetry: ThE: 13.5, Th*: 66.6246, Ph*: 0, Q2: 0.313909, AL: -0.0302837

TransAsymmetry: ThE: 13.5, Th*: 156.625, Ph*: 8.53774e-07, Q2: 0.313909, AT: -0.043315

Pal. = 2/3

09/24/12

For 12.5° data

Longitudinal:
$$\frac{-0.03116 + (-0.0252685)}{2} =$$

$$= -0.028214 \pm 0.002946$$

Transverse:
$$\frac{-0.045172 + (-0.0295942)}{2} =$$

$$= -0.037383 \pm 0.00779$$

ThE: 12.5, Th*: 68.2544, Ph*: 8.53774e-07
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 T10(0.27063 GeV**2, 70deg) = 0.0587881
 T11(0.27063 GeV**2, 70deg) = 0.182052
 *****Point ThE: 10, Eprime: 2.37886 is out of acceptance!
 LongAsymmetry: ThE: 10, Eprime: 2.37886, Th*: 72.4212, Ph*: 0, Q2: 0.175324, AL: -0.0195369
 TransAsymmetry: ThE: 10, Eprime: 2.37886, Th*: 162.421, Ph*: 0, Q2: 0.175324, AT: -0.0158946

 *****Point ThE: 10.2857, Eprime: 2.37622 is out of acceptance!
 LongAsymmetry: ThE: 10.2857, Eprime: 2.37622, Th*: 71.9387, Ph*: 0, Q2: 0.185252, AL: -0.0204897
 TransAsymmetry: ThE: 10.2857, Eprime: 2.37622, Th*: 161.939, Ph*: 0, Q2: 0.185252, AT: -0.0178321

 *****Point ThE: 10.5714, Eprime: 2.3735 is out of acceptance!
 LongAsymmetry: ThE: 10.5714, Eprime: 2.3735, Th*: 71.4577, Ph*: 0, Q2: 0.195433, AL: -0.0214491
 TransAsymmetry: ThE: 10.5714, Eprime: 2.3735, Th*: 161.458, Ph*: 8.53774e-07, Q2: 0.195433, AT: -0.0199304

 *****Point ThE: 10.8571, Eprime: 2.37072 is out of acceptance!
 LongAsymmetry: ThE: 10.8571, Eprime: 2.37072, Th*: 70.9784, Ph*: 0, Q2: 0.205866, AL: -0.0224117
 TransAsymmetry: ThE: 10.8571, Eprime: 2.37072, Th*: 160.978, Ph*: 0, Q2: 0.205866, AT: -0.0221796

 *****Point ThE: 11.1429, Eprime: 2.36787 is out of acceptance!
 LongAsymmetry: ThE: 11.1429, Eprime: 2.36787, Th*: 70.5007, Ph*: 0, Q2: 0.216549, AL: -0.0233728
 TransAsymmetry: ThE: 11.1429, Eprime: 2.36787, Th*: 160.501, Ph*: 0, Q2: 0.216549, AT: -0.0245609

 *****Point ThE: 11.4286, Eprime: 2.36496 is out of acceptance!
 LongAsymmetry: ThE: 11.4286, Eprime: 2.36496, Th*: 70.0246, Ph*: 0, Q2: 0.227479, AL: -0.0243272
 TransAsymmetry: ThE: 11.4286, Eprime: 2.36496, Th*: 160.025, Ph*: 0, Q2: 0.227479, AT: -0.0270457

 *****Point ThE: 11.7143, Eprime: 2.36198 is out of acceptance!
 LongAsymmetry: ThE: 11.7143, Eprime: 2.36198, Th*: 69.5502, Ph*: 0, Q2: 0.238654, AL: -0.0252685
 TransAsymmetry: ThE: 11.7143, Eprime: 2.36198, Th*: 159.55, Ph*: 0, Q2: 0.238654, AT: -0.0295947

 LongAsymmetry: ThE: 12, Eprime: 2.35894, Th*: 69.0774, Ph*: 0, Q2: 0.250072, AL: -0.0261893
 TransAsymmetry: ThE: 12, Eprime: 2.35894, Th*: 159.077, Ph*: 0, Q2: 0.250072, AT: -0.0321579

 LongAsymmetry: ThE: 12.2857, Eprime: 2.35583, Th*: 68.6065, Ph*: 0, Q2: 0.26173, AL: -0.0270808
 TransAsymmetry: ThE: 12.2857, Eprime: 2.35583, Th*: 158.606, Ph*: 0, Q2: 0.26173, AT: -0.0346759

 LongAsymmetry: ThE: 12.5714, Eprime: 2.35266, Th*: 68.1372, Ph*: 0, Q2: 0.273627, AL: -0.027933
 TransAsymmetry: ThE: 12.5714, Eprime: 2.35266, Th*: 158.137, Ph*: 0, Q2: 0.273627, AT: -0.0370825

 LongAsymmetry: ThE: 12.8571, Eprime: 2.34942, Th*: 67.6698, Ph*: 0, Q2: 0.28576, AL: -0.0287346
 TransAsymmetry: ThE: 12.8571, Eprime: 2.34942, Th*: 157.67, Ph*: 0, Q2: 0.28576, AT: -0.0393083

 LongAsymmetry: ThE: 13.1429, Eprime: 2.34613, Th*: 67.2041, Ph*: 0, Q2: 0.298126, AL: -0.0294727
 TransAsymmetry: ThE: 13.1429, Eprime: 2.34613, Th*: 157.204, Ph*: 0, Q2: 0.298126, AT: -0.0412855

 LongAsymmetry: ThE: 13.4286, Eprime: 2.34277, Th*: 66.7403, Ph*: 0, Q2: 0.310723, AL: -0.0301327
 TransAsymmetry: ThE: 13.4286, Eprime: 2.34277, Th*: 156.74, Ph*: 0, Q2: 0.310723, AT: -0.0429528

 LongAsymmetry: ThE: 13.7143, Eprime: 2.33935, Th*: 66.2783, Ph*: 0, Q2: 0.32355, AL: -0.0306989
 TransAsymmetry: ThE: 13.7143, Eprime: 2.33935, Th*: 156.278, Ph*: 0, Q2: 0.32355, AT: -0.04426

 LongAsymmetry: ThE: 14, Eprime: 2.33587, Th*: 65.8181, Ph*: 0, Q2: 0.336602, AL: -0.0311538
 TransAsymmetry: ThE: 14, Eprime: 2.33587, Th*: 155.818, Ph*: 8.53774e-07, Q2: 0.336602, AT: -0.0451719

Minimum!
 AL: -0.0252685
 AT: -0.0295947

Maximum!

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 143°

ThE: 12.5, Th*: 68.2544, Ph*: 8.53774e-07
Gc(0.27063 GeV**2) = 0.0617677
Gq(0.27063 GeV**2) = 2.3764
Gm(0.27063 GeV**2) = 0.166093
A(0.27063 GeV**2) = 0.00602567
B(0.27063 GeV**2) = 0.000721013
T20(0.27063 GeV**2, 70deg) = -1.07945
T21(0.27063 GeV**2, 70deg) = 0.233359
T22(0.27063 GeV**2, 70deg) = -0.0240091
T10(0.27063 GeV**2, 70deg) = 0.0587881
T11(0.27063 GeV**2, 70deg) = 0.182052
*****Point ThE: 10, Eprime: 2.37886 is out of acceptance!
LongAsymmetry: ThE: 10, Eprime: 2.37886, Th*: 72.4212, Ph*: 0, Q2: 0.175324, AL: -0.0195369
TransAsymmetry: ThE: 10, Eprime: 2.37886, Th*: 162.421, Ph*: 0, Q2: 0.175324, AT: -0.0158946

*****Point ThE: 10.2941, Eprime: 2.37614 is out of acceptance!
LongAsymmetry: ThE: 10.2941, Eprime: 2.37614, Th*: 71.9245, Ph*: 0, Q2: 0.185547, AL: -0.0205178
TransAsymmetry: ThE: 10.2941, Eprime: 2.37614, Th*: 161.925, Ph*: 0, Q2: 0.185547, AT: -0.0178915

*****Point ThE: 10.5882, Eprime: 2.37334 is out of acceptance!
LongAsymmetry: ThE: 10.5882, Eprime: 2.37334, Th*: 71.4295, Ph*: 0, Q2: 0.19604, AL: -0.0215057
TransAsymmetry: ThE: 10.5882, Eprime: 2.37334, Th*: 161.43, Ph*: 0, Q2: 0.19604, AT: -0.0200587

*****Point ThE: 10.8824, Eprime: 2.37047 is out of acceptance!
LongAsymmetry: ThE: 10.8824, Eprime: 2.37047, Th*: 70.9362, Ph*: 0, Q2: 0.206799, AL: -0.0224966
TransAsymmetry: ThE: 10.8824, Eprime: 2.37047, Th*: 160.936, Ph*: 8.53774e-07, Q2: 0.206799, AT: -0.0223847

*****Point ThE: 11.1765, Eprime: 2.36753 is out of acceptance!
LongAsymmetry: ThE: 11.1765, Eprime: 2.36753, Th*: 70.4446, Ph*: 0, Q2: 0.217822, AL: -0.0234856
TransAsymmetry: ThE: 11.1765, Eprime: 2.36753, Th*: 160.445, Ph*: 0, Q2: 0.217822, AT: -0.0248485

*****Point ThE: 11.4706, Eprime: 2.36452 is out of acceptance!
LongAsymmetry: ThE: 11.4706, Eprime: 2.36452, Th*: 69.9547, Ph*: 0, Q2: 0.229107, AL: -0.0244666
TransAsymmetry: ThE: 11.4706, Eprime: 2.36452, Th*: 159.955, Ph*: 0, Q2: 0.229107, AT: -0.0274176

*****Point ThE: 11.7647, Eprime: 2.36145 is out of acceptance!
LongAsymmetry: ThE: 11.7647, Eprime: 2.36145, Th*: 69.4666, Ph*: 8.53774e-07, Q2: 0.240651, AL: -0.0254327
TransAsymmetry: ThE: 11.7647, Eprime: 2.36145, Th*: 159.467, Ph*: 8.53774e-07, Q2: 0.240651, AT: -0.0300476

LongAsymmetry: ThE: 12.0588, Eprime: 2.3583, Th*: 68.9803, Ph*: 0, Q2: 0.252452, AL: -0.0263755
TransAsymmetry: ThE: 12.0588, Eprime: 2.3583, Th*: 158.98, Ph*: 0, Q2: 0.252452, AT: -0.0326821

LongAsymmetry: ThE: 12.3529, Eprime: 2.35509, Th*: 68.4959, Ph*: 8.53774e-07, Q2: 0.264508, AL: -0.0272852
TransAsymmetry: ThE: 12.3529, Eprime: 2.35509, Th*: 158.496, Ph*: 0, Q2: 0.264508, AT: -0.0352548

LongAsymmetry: ThE: 12.6471, Eprime: 2.35181, Th*: 68.0133, Ph*: 0, Q2: 0.276816, AL: -0.0281507
TransAsymmetry: ThE: 12.6471, Eprime: 2.35181, Th*: 158.013, Ph*: 0, Q2: 0.276816, AT: -0.0376922

LongAsymmetry: ThE: 12.9412, Eprime: 2.34846, Th*: 67.5326, Ph*: 0, Q2: 0.289373, AL: -0.0289589
TransAsymmetry: ThE: 12.9412, Eprime: 2.34846, Th*: 157.533, Ph*: 0, Q2: 0.289373, AT: -0.0399186

LongAsymmetry: ThE: 13.2353, Eprime: 2.34505, Th*: 67.0538, Ph*: 0, Q2: 0.302176, AL: -0.0296955
TransAsymmetry: ThE: 13.2353, Eprime: 2.34505, Th*: 157.054, Ph*: 0, Q2: 0.302176, AT: -0.0418615

LongAsymmetry: ThE: 13.5294, Eprime: 2.34157, Th*: 66.577, Ph*: 8.53774e-07, Q2: 0.315224, AL: -0.0303441
TransAsymmetry: ThE: 13.5294, Eprime: 2.34157, Th*: 156.577, Ph*: 0, Q2: 0.315224, AT: -0.0434574

LongAsymmetry: ThE: 13.8235, Eprime: 2.33803, Th*: 66.1021, Ph*: 0, Q2: 0.328514, AL: -0.0308869
TransAsymmetry: ThE: 13.8235, Eprime: 2.33803, Th*: 156.102, Ph*: 0, Q2: 0.328514, AT: -0.0446568

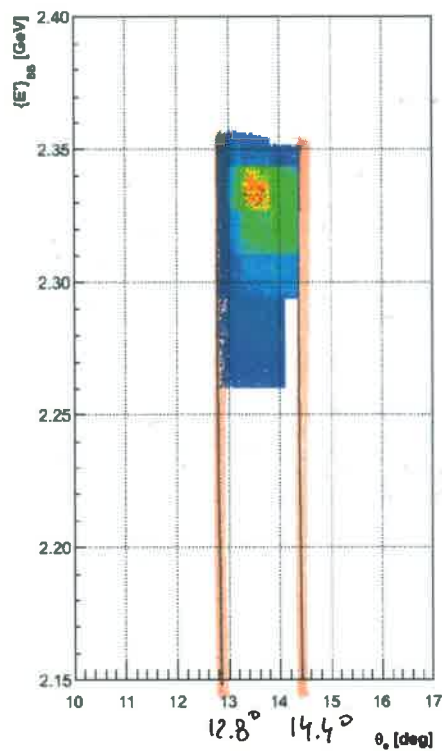
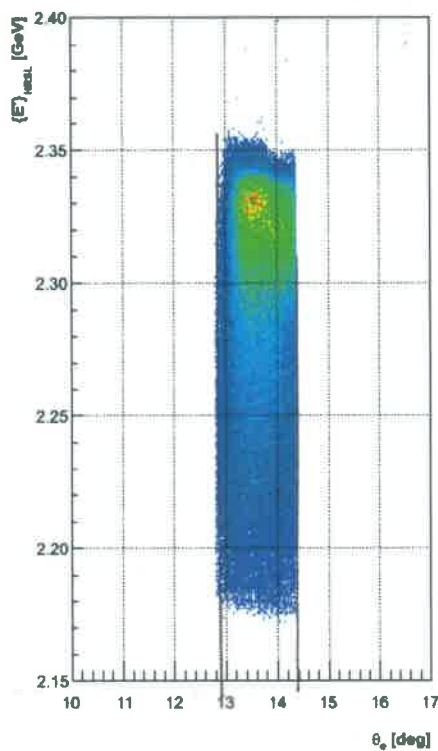
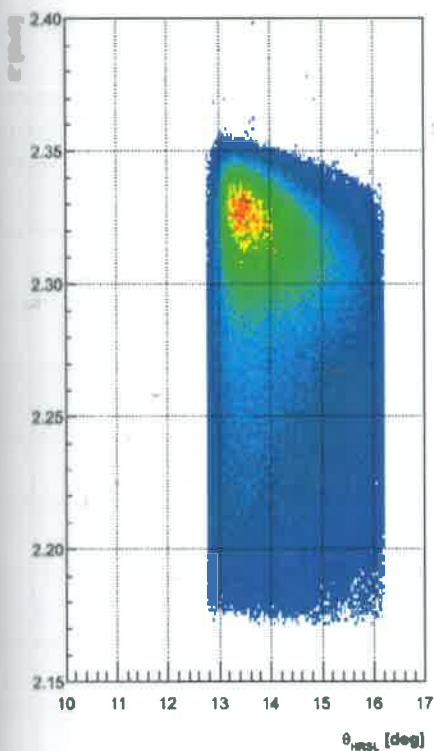
LongAsymmetry: ThE: 14.1176, Eprime: 2.33442, Th*: 65.6292, Ph*: 0, Q2: 0.342042, AL: -0.0313045
TransAsymmetry: ThE: 14.1176, Eprime: 2.33442, Th*: 155.629, Ph*: 0, Q2: 0.342042, AT: -0.0454276

LongAsymmetry: ThE: 14.4118, Eprime: 2.33075, Th*: 65.1582, Ph*: 0, Q2: 0.355806, AL: -0.0315759
 TransAsymmetry: ThE: 14.4118, Eprime: 2.33075, Th*: 155.158, Ph*: 0, Q2: 0.355806, AT: -0.0457568

LongAsymmetry: ThE: 14.7059, Eprime: 2.32702, Th*: 64.6893, Ph*: 0, Q2: 0.369803, AL: -0.0316795
 TransAsymmetry: ThE: 14.7059, Eprime: 2.32702, Th*: 154.689, Ph*: 0, Q2: 0.369803, AT: -0.04565

LongAsymmetry: ThE: 15, Eprime: 2.32323, Th*: 64.2224, Ph*: 0, Q2: 0.384031, AL: -0.0315926
 TransAsymmetry: ThE: 15, Eprime: 2.32323, Th*: 154.222, Ph*: 0, Q2: 0.384031, AT: -0.0451288

Sekej' delov'na se merja toka za 14.5° data:



Longitudinal: $A_L^R = \frac{-0.02896 - 0.031576}{2} = -0.0303 \pm 0.0013$

Transverse: $A_T^R = \frac{-0.039919 - 0.04576}{2} = -0.04284 \pm 0.00292$