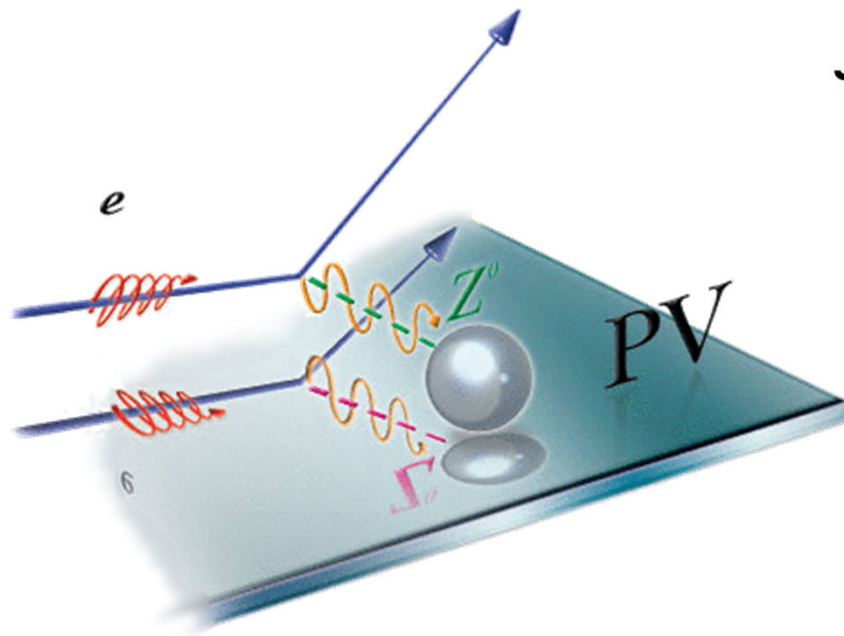


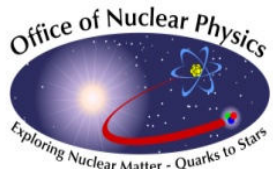
1. Using SBS for new C-REx
2. Strange Quark Program
3. Parity-Violating Deep Inelastic Scattering

Robert Michaels

Jefferson Lab  
Thomas Jefferson National Accelerator Facility



INT Workshop, 2019  
Weak Elastic Scattering



Thomas Jefferson National Accelerator Facility



# Parity Violating Asymmetry

Electron scattering

$$\sigma \approx \left| \begin{array}{c} \text{Diagram 1: } e^- \text{ scattering via } \gamma \text{ (photon) exchange with } ^{208}\text{Pb} \\ \text{Diagram 2: } e^- \text{ scattering via } Z^0 \text{ (Z boson) exchange with } ^{208}\text{Pb} \end{array} \right|^2$$

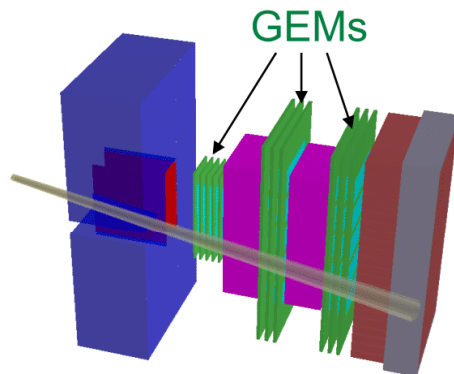
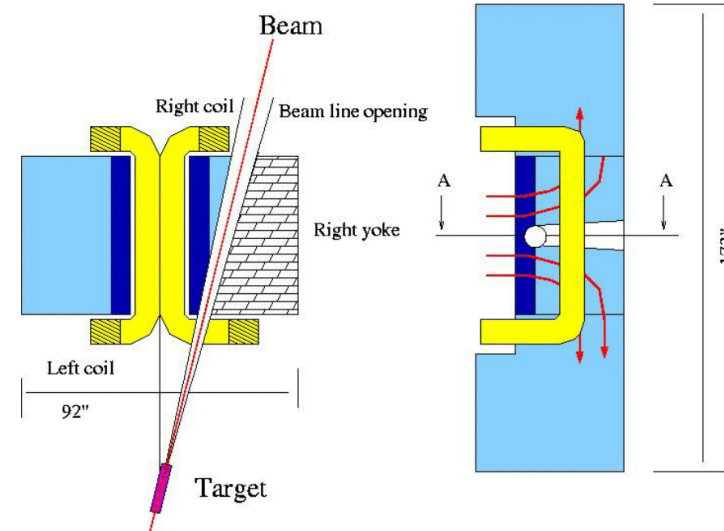
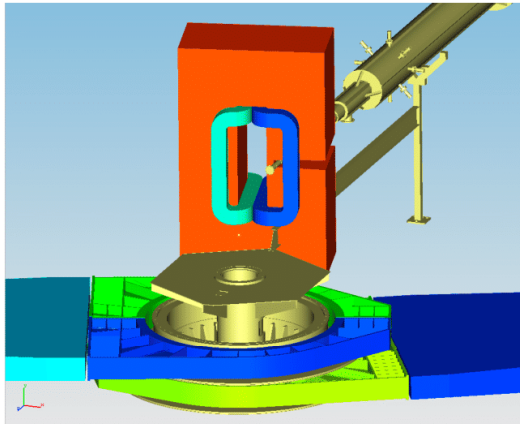
$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim 10^{-4} \times Q^2 \sim 10^{-6}$$

## Applications of $A_{PV}$ at Jefferson Lab

- Nucleon Structure  
Strangeness  $s \bar{s}$  in proton (HAPPEX, G0 expts)
- Test of Standard Model of Electroweak  $\sin^2 \theta_W$   
e-e (MOLLER) or e-q (PVDIS)  
elastic e-p at low  $Q^2$  (QWEAK)
- Nuclear Structure (neutron density) : PREX, CREX

For a possible future C-REX at higher  $Q^2$

## Super Bigbite Spectrometer



- Magnet: 48D48 - 46 cm gap, 2-3 Tesla\*m
- Solid angle is 70 msr at angle 15 deg.
- GEM chambers with 70  $\mu\text{m}$  resolution
- momentum resolution is 0.5% for 5 GeV/c
- angular resolution is 0.5 mr

# Parameters of SBS

$\theta_{central}$ , degree	$\Omega$ , msr	D, meter	Hor. range, degree	Vert. range, degree
3.5	5	9.5	$\pm 1.3$	$\pm 3.3$
5.0	12	5.8	$\pm 1.9$	$\pm 4.9$
7.5	30	3.2	$\pm 3$	$\pm 8$
15	72	1.6	$\pm 4.8$	$\pm 12.2$
30	76	1.5	$\pm 4.9$	$\pm 12.5$

Solid angle

Resolution:

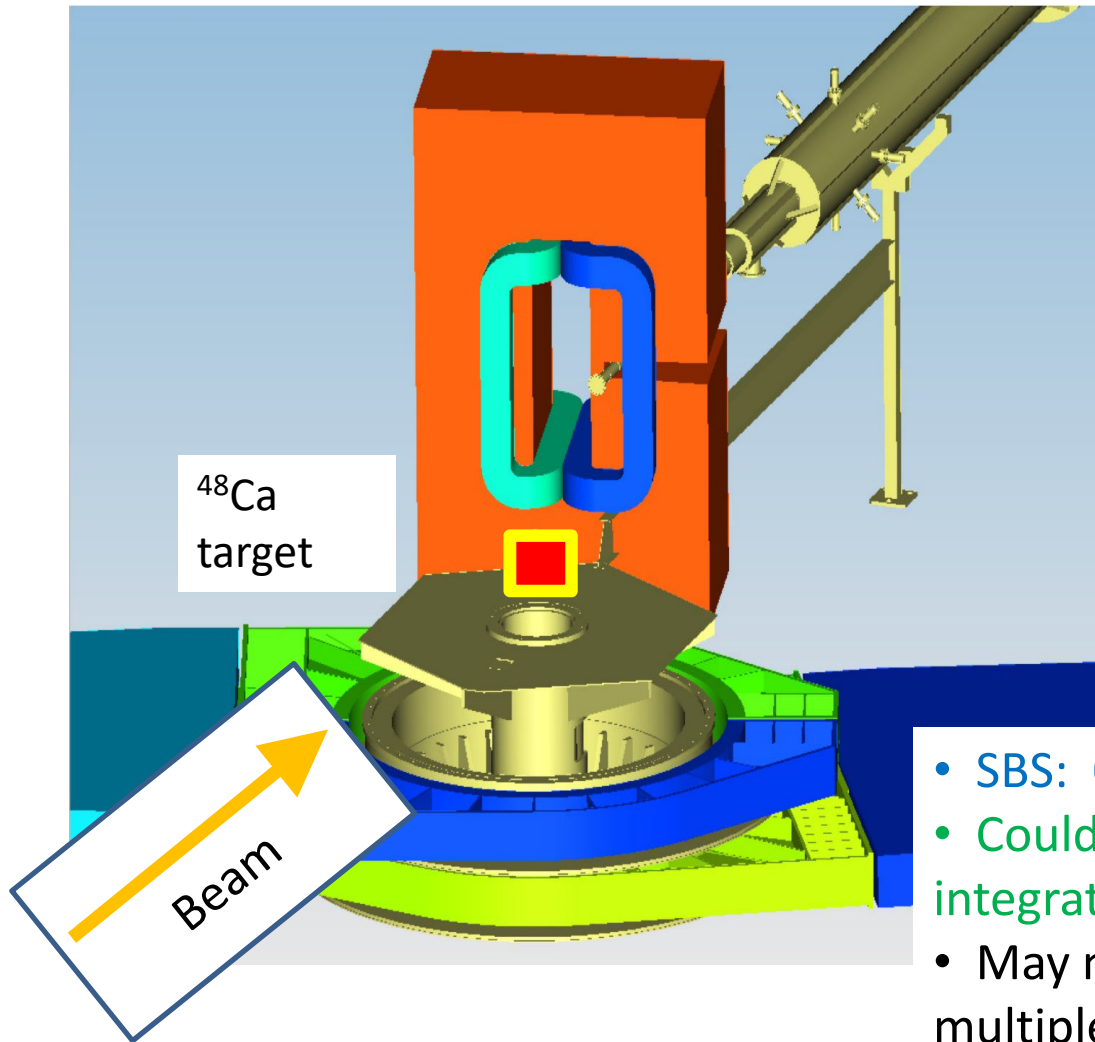
Momentum  $\Rightarrow \frac{\sigma_p}{P} = 0.0029 + 0.0003 \times p[\text{GeV}]$

Angular  $\Rightarrow \sigma_\theta = 0.14 + 1.3/p$  [GeV], mrad

Momentum acceptance  $\Rightarrow P$  range from **2 – 10** , GeV/c

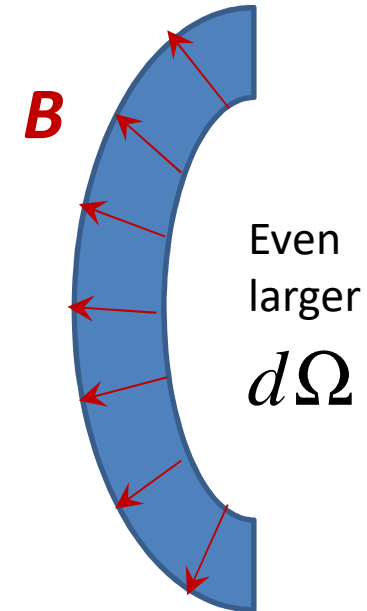
Using SuperBigBite Spectrometer

$d\Omega$  large, but varies with angle setting



or a new Crescent - shaped Spectrometer

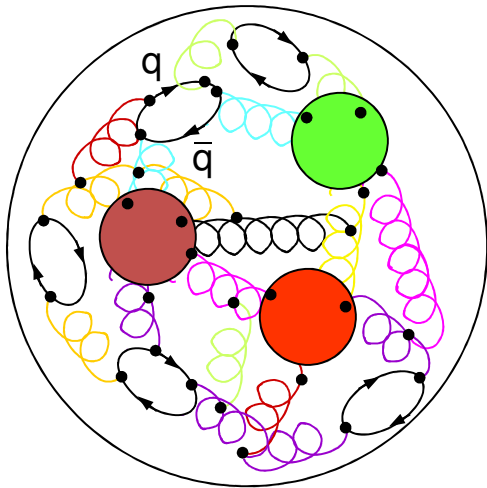
(Bogdan Wojtsekhowski)



- SBS: 0.35% momentum resolution
- Could be improved ( $\sim 1/2$ ) by putting integrating detector near wall of hall A.
- May need helium bag to reduce multiple scattering.

# Historical Review: Strange Quarks in Nucleons

What is a proton made of?



- Parity violating electron scattering provides a third, independent constraint on the vector form factors of the nucleon.

- Early theoretical motivations to look for strangeness effects ( $s\bar{s}$ ), for example:

D.B. Kaplan, A. Manohar Nucl. Phys. **B310** (1988) 527

- Here I briefly summarize the experimental and theoretical conclusions that the strangeness form factors are “small”, i.e. at most a few percent of the electromagnetic FF.

# Quark Currents in the Nucleon

R.D. McKeown, Phys. Lett. B **219** (1989) 140

D.H. Beck, PRD **39** (1989) 3248.

Measure

$$G^{\gamma,p}, G^{Z,p}, G^{\gamma,n} : \quad G \sim \langle N | \sum_i e_i \bar{q}_i \Gamma_\mu q_i | N \rangle$$

e.g.

$$G_{E,M}^{\gamma,p} = \frac{2}{3} G_{E,M}^{u,p} - \frac{1}{3} (G_{E,M}^{d,p} + G_{E,M}^{s,p})$$

Assume  
charge  
symmetry

$$\left. \begin{aligned} G^{u,p} &= G^{d,n} \\ G^{d,p} &= G^{u,n} \\ G^{s,p} &= G^{s,n} \end{aligned} \right\}$$

G. A. Miller PRC **57** (1998) 1492.

R. Lewis, N. Mobed PRD **59** (1999) 073002

B. Kubis, R. Lewis PRC **74** (2006) 015204

therefore

$$\begin{aligned} G_{E,M}^u &= (3 - 4 \sin^2 \theta_W) G_{E,M}^{\gamma,p} - G_{E,M}^{Z,p} \\ G_{E,M}^d &= (2 - 4 \sin^2 \theta_W) G_{E,M}^{\gamma,p} + G_{E,M}^{\gamma,n} - G_{E,M}^{Z,p} \\ G_{E,M}^s &= (1 - 4 \sin^2 \theta_W) G_{E,M}^{\gamma,p} - G_{E,M}^{\gamma,n} - G_{E,M}^{Z,p} \end{aligned}$$

dropping the  $p$  superscripts on the left

# Measuring Strange Vector Form Factors

$$A = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \propto \frac{\text{Diagram with } \gamma \text{ and } Z^0 \text{ exchange}}{2} \sim \frac{10^{-4} Q^2}{\text{GeV}^2}$$

Interference with EM amplitude makes Neutral Current (NC) amplitude accessible

For a proton:

$$A = \left[ \frac{-G_F Q^2}{4\pi\alpha\sqrt{2}} \right] \frac{A_E + A_M + A_A}{\sigma_p} \sim \text{few parts per million}$$

$$A_E = \epsilon G_E^p G_E^Z$$

$$A_M = \tau G_M^p G_M^Z$$

$$A_A = (1 - 4 \sin^2 \theta_W) \epsilon' G_M^p \tilde{G}_A$$



For spin=0, T=0 <sup>4</sup>He:

$G_E^s$  only !

nuclear corrections: forward angle,  
low  $Q^2$  only

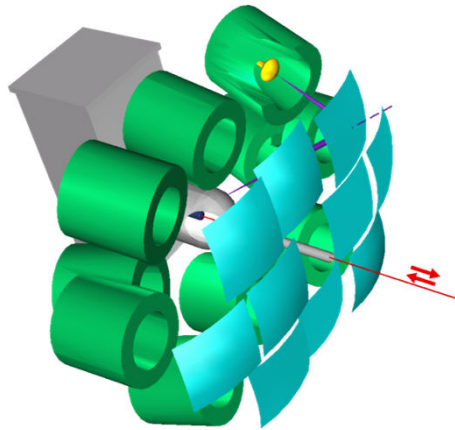
For deuterium:

Enhanced  $G_A$

Back-angle quasi-elastic.



# 20 year World-Wide Effort on Strange Quarks



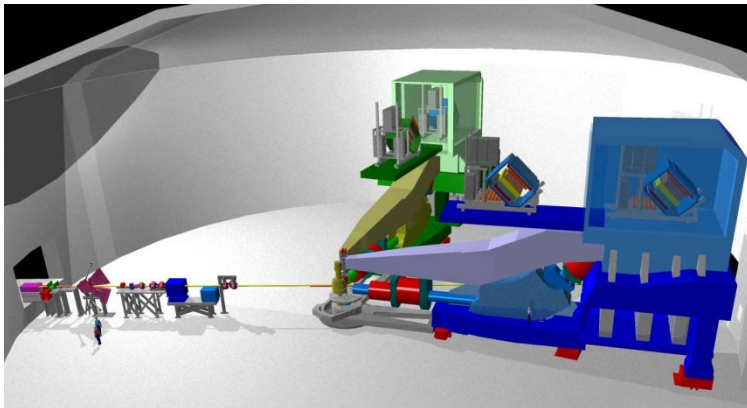
**SAMPLE**

open geometry,  
integrating,  
back-angle only

**HAPPEX**

Precision spectrometer,  
integrating

Forward angle, also  
 $^4\text{He}$  at low  $Q^2$

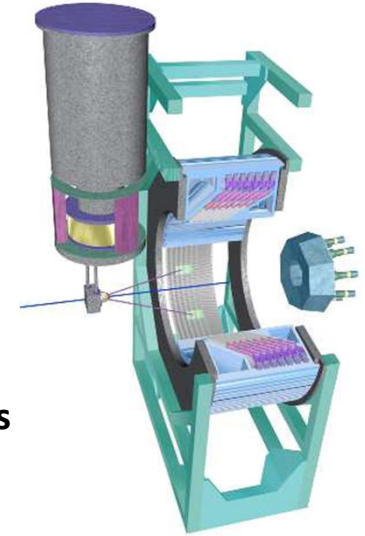


**A4**

Open geometry

Fast counting calorimeter for  
background rejection

Forward and Backward angles

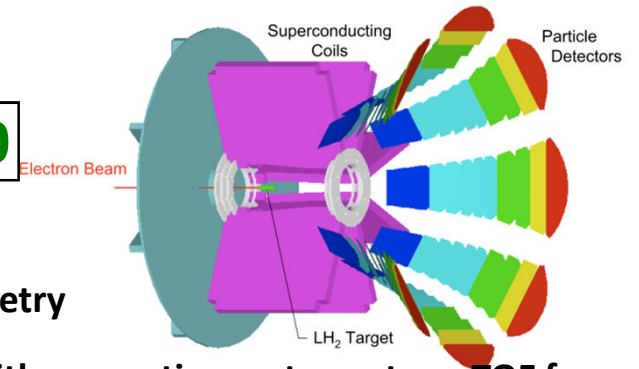


**GO**

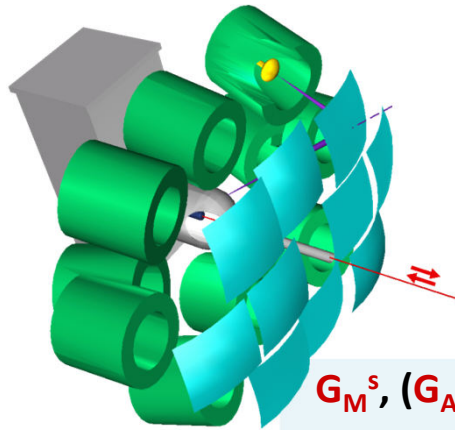
Open geometry

Fast counting with magnetic spectrometer + TOF for  
background rejection

Forward and Backward angles over a range of  $Q^2$



# Kinematics of World Data



**SAMPLE**

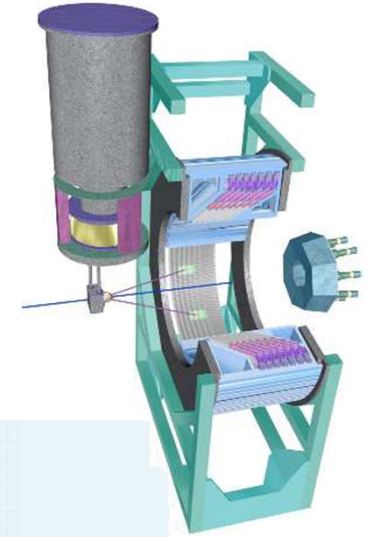
open geometry,  
integrating

$G_M^S, (G_A)$  at  $Q^2 = 0.1 \text{ GeV}^2$

**A4**

Open geometry

Fast counting calorimeter for  
background rejection



$G_E^S + 0.23 G_M^S$  at  $Q^2 = 0.23 \text{ GeV}^2$

$G_E^S + 0.10 G_M^S$  at  $Q^2 = 0.1 \text{ GeV}^2$

$G_M^S, G_A^e$  at  $Q^2 = 0.23 \text{ GeV}^2$

**HAPPEX**

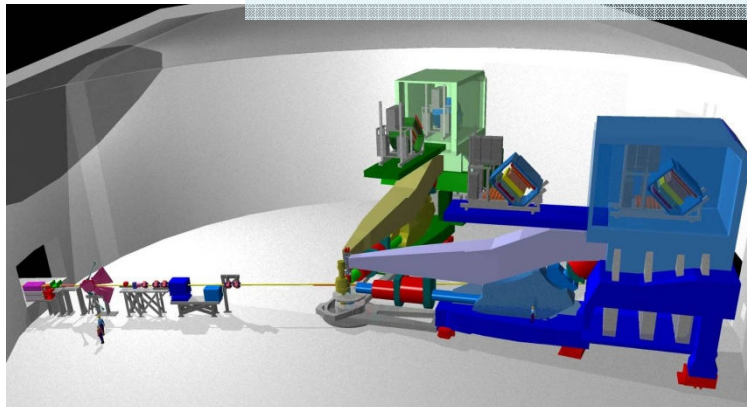
Precision  
spectrometer,  
integrating

$G_E^S + 0.39 G_M^S$  at  $Q^2 = 0.48 \text{ GeV}^2$

$G_E^S + 0.08 G_M^S$  at  $Q^2 = 0.1 \text{ GeV}^2$

$G_E^S$  at  $Q^2 = 0.1 \text{ GeV}^2$  ( $^4\text{He}$ )

$G_E^S + 0.48 G_M^S$  at  $Q^2 = 0.62 \text{ GeV}^2$



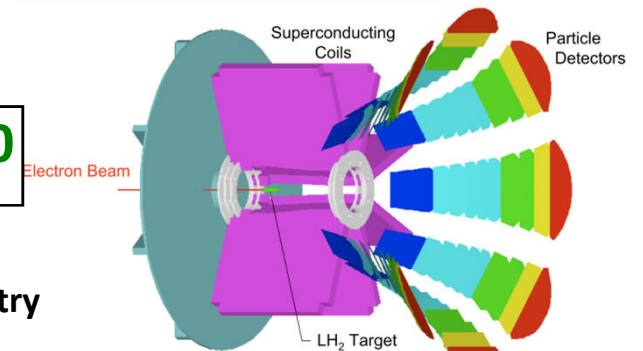
**GO**

Open geometry

Fast counting with magnetic spectrometer + timing for  
background rejection

$G_E^S + \eta G_M^S$  over  $Q^2 = [0.12, 1.0] \text{ GeV}^2$

$G_M^S, G_A^e$  at  $Q^2 = 0.23, 0.62 \text{ GeV}^2$

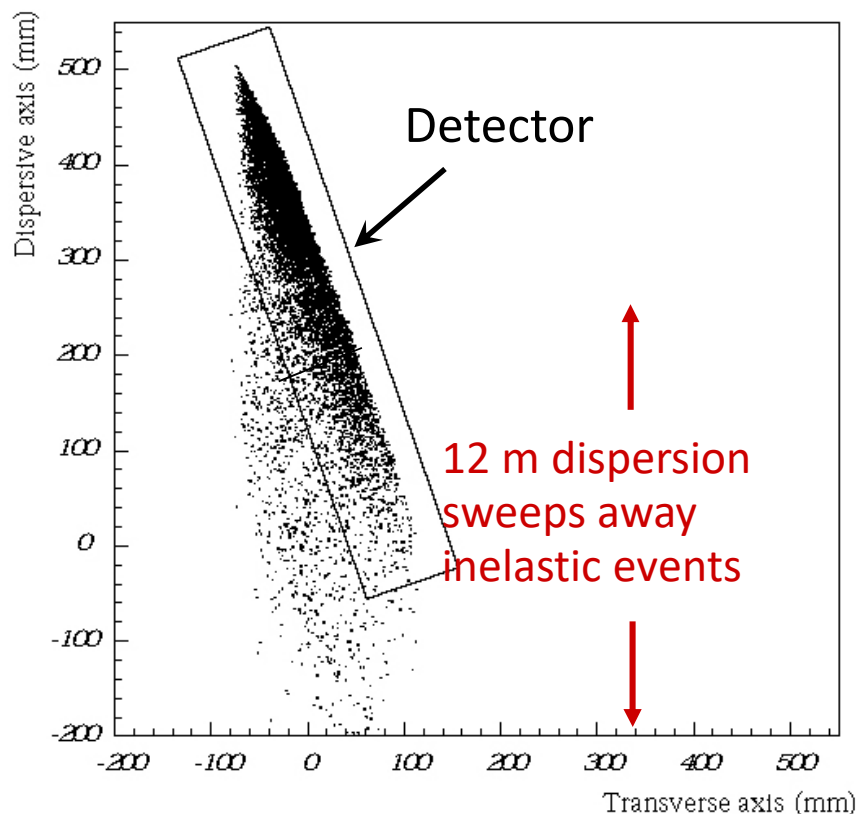


# Example: HAPPEX

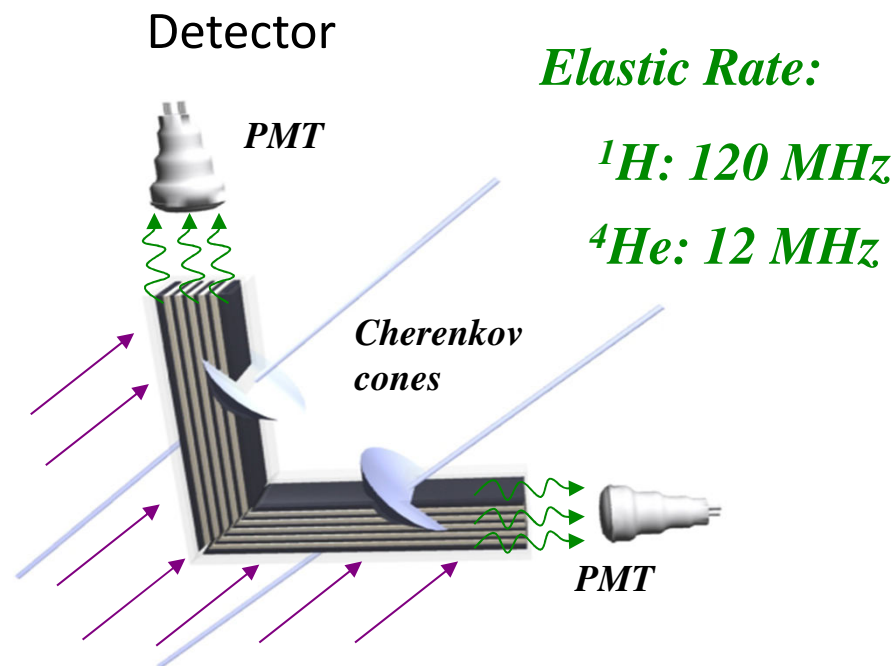
Clean separation of elastic events by HRS optics



Locate detector over elastic line and integrate the flux



Large dispersion & heavy shielding reduce backgrounds at focal plane



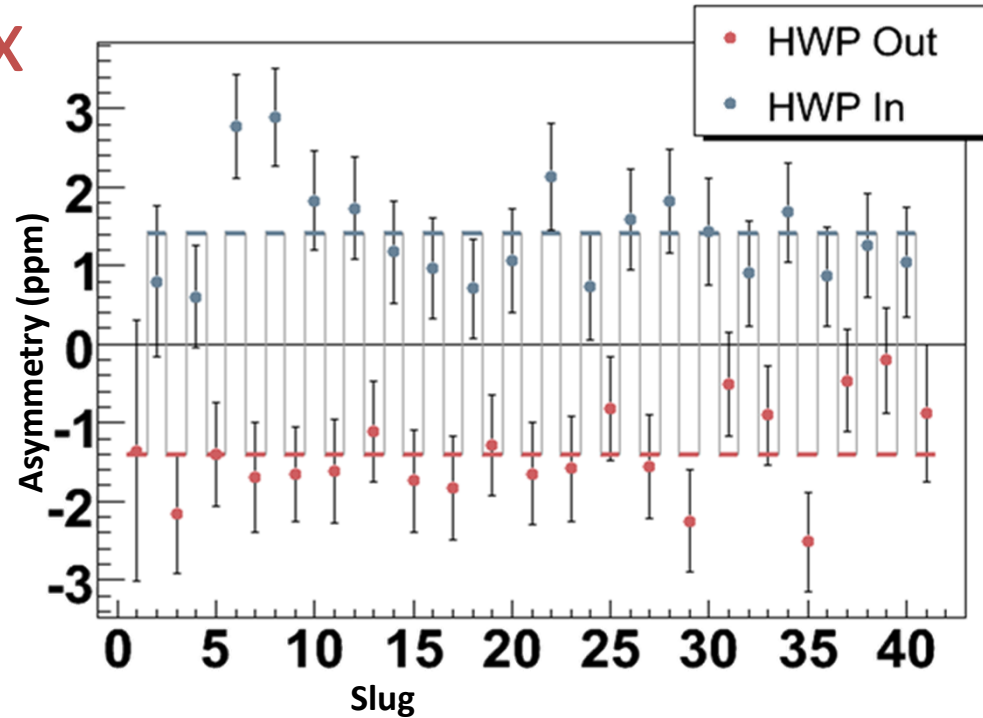
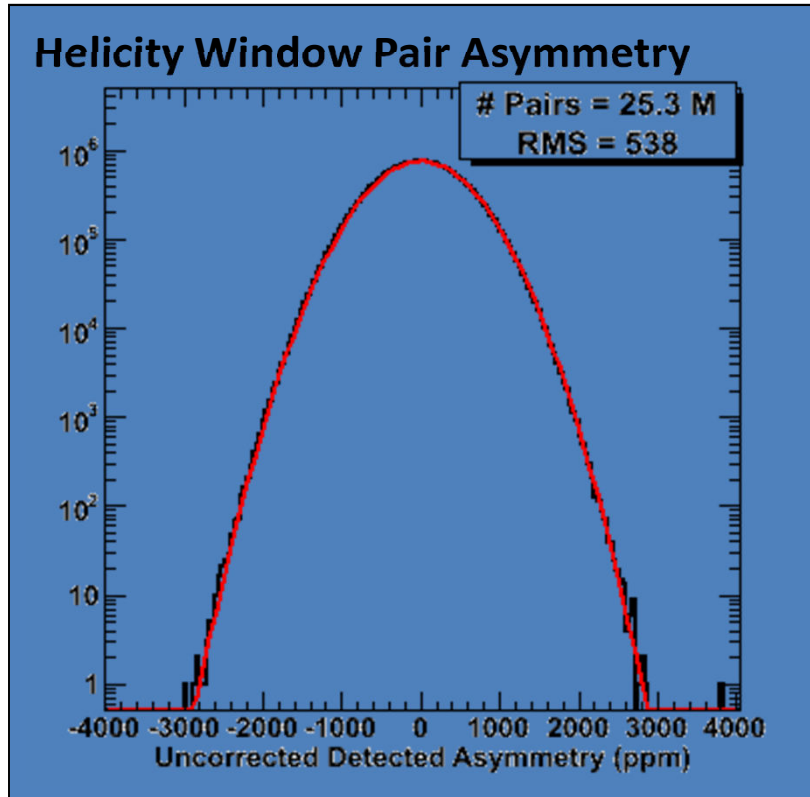
Brass-Quartz Integrating Cerenkov Shower Calorimeter

# Example: HAPPEX

$^1\text{H}$  Results (2005)

Parity Violating Asymmetry

$A_{\text{raw}}$  correction  $\sim 3$  ppb

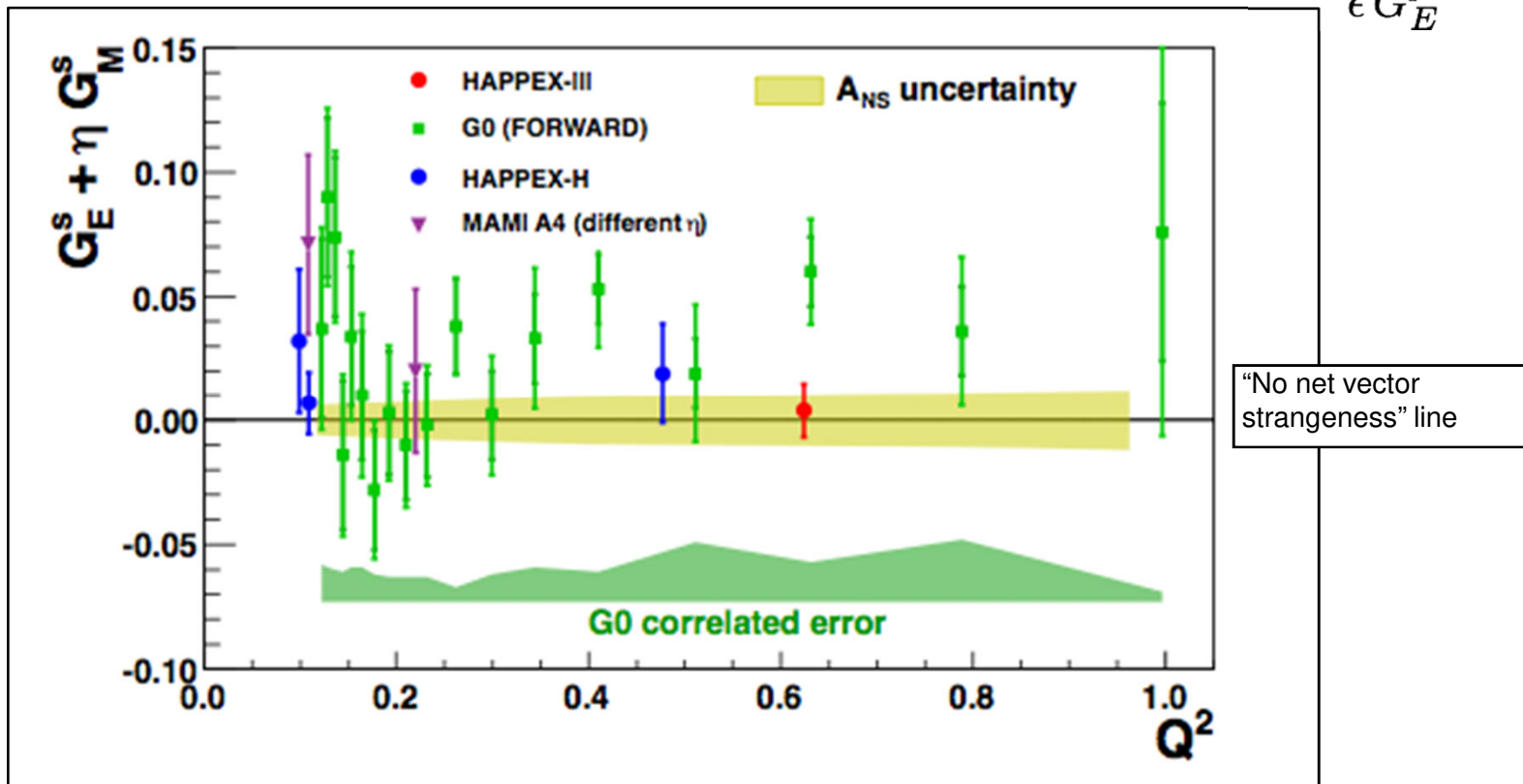


$$Q^2 = 0.1089 \pm 0.0011 \text{ GeV}^2$$

$$A_{\text{raw}} = -1.58 \text{ ppm} \pm 0.12 \text{ (stat)} \\ \pm 0.04 \text{ (syst)}$$

# Forward-angle proton scattering world data prior to Qweak

$$\eta = \frac{\tau G_M^p}{\epsilon G_E^p} \sim Q^2$$

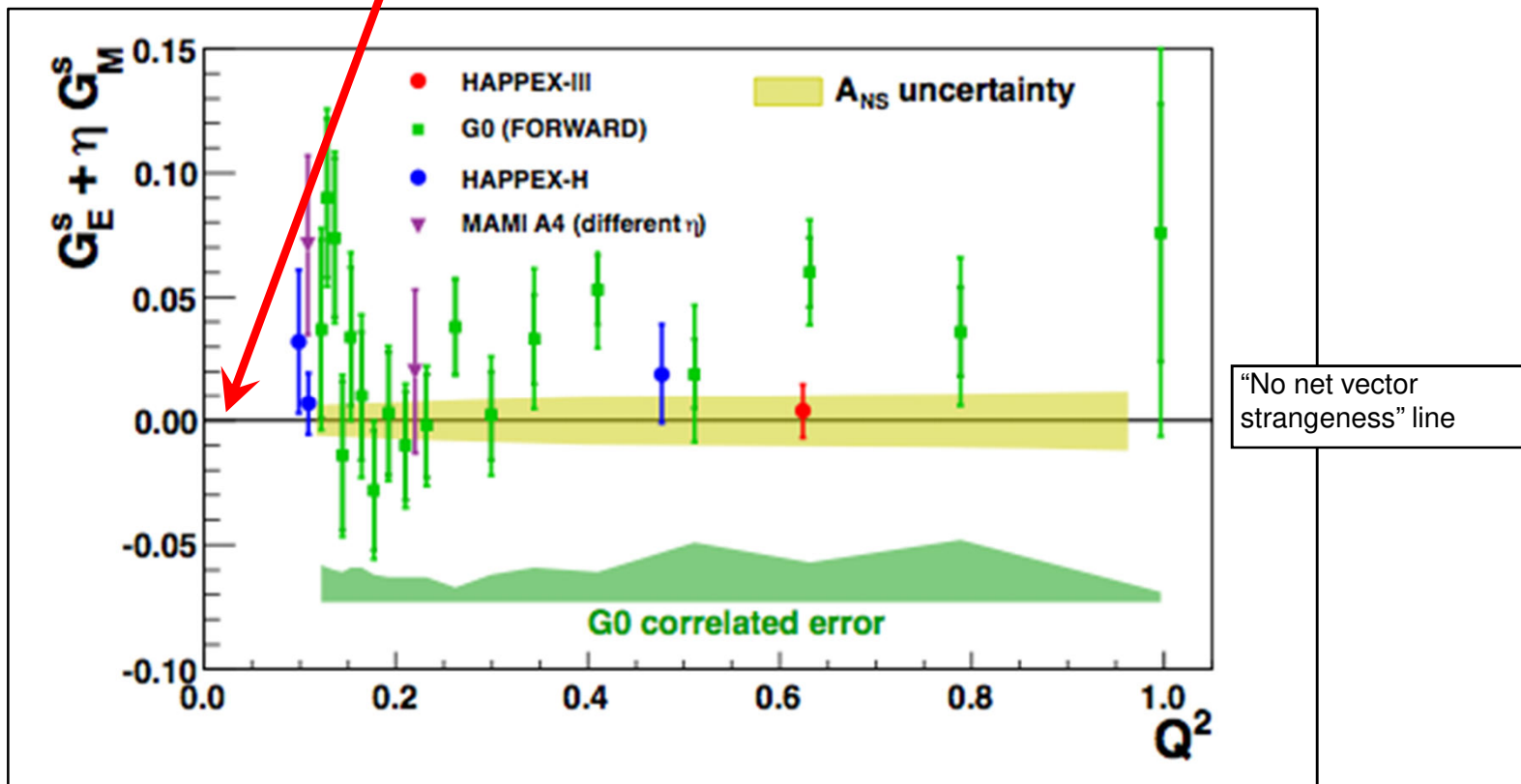


Combining results from forward-angle proton scattering (similar beam energies)

- Additional data at backward angles, and  $^2\text{H}$  and  $^4\text{He}$  target

**Qweak** result consistent with Std. Model w/ no strangeness

$$A = -226.5 \pm 7.3 \pm 5.8 \text{ ppb}, \text{ at } \langle Q^2 \rangle = 0.0248$$

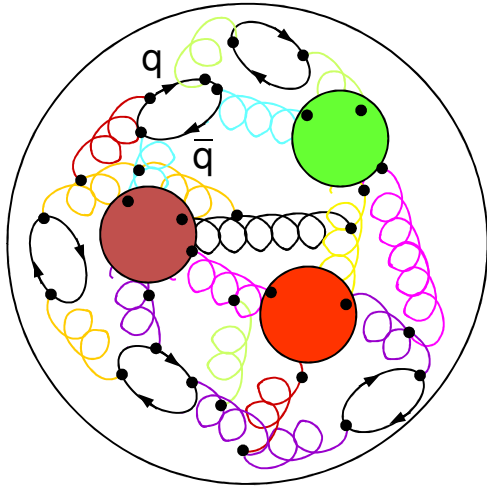


Combining results from forward-angle proton scattering (similar beam energies)

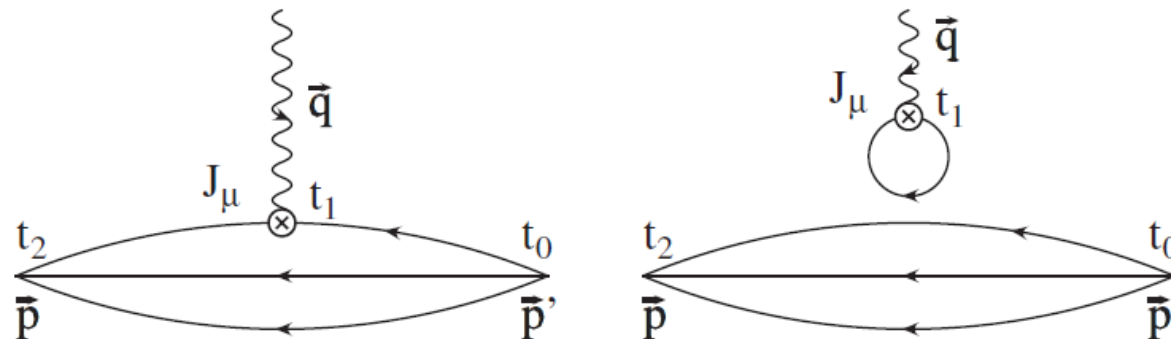
- Additional data at backward angles, and  $^2\text{H}$  and  $^4\text{He}$  target



## Summary: Strange Quarks in Protons



- Strange form factors arise from quantum fluctuations:  $s\bar{s}$  pairs
- Experimentally, these form factors are small, less than a few percent of proton EM FF.
- Precise, direct lattice QCD agree (T. Doi, et al) agree with earlier “lattice + chiral extrapolation” (Leinweber and Thomas) and with experiment.



T. Doi, *et.al.* Phys Rev D **80** (2009) 094503  
(xQCD Collaboration)

# Deep Inelastic Scattering

- Parity violation arises from axial couplings either at the electron or at the quark.
- PVDIS emphasizes the latter: the quark's chirality preference when participating in the in the weak force.

$$2C_{2u} - C_{2d}$$

- This has been measured twice in 40 years.
  - SLAC E122
  - JLAB 2009 -- factor of 5 improvement



# PVDIS Formalism

For deuterium (isoscalar), neglecting sea quarks.

Unique sensitivity to these

$$A_{PV} = \left( \frac{3G_F Q^2}{10\sqrt{2}\pi\alpha} \right) \left[ (2C_{1u} - C_{1d}) + Y_3 \overbrace{(2C_{2u} - C_{2d})} \right]$$

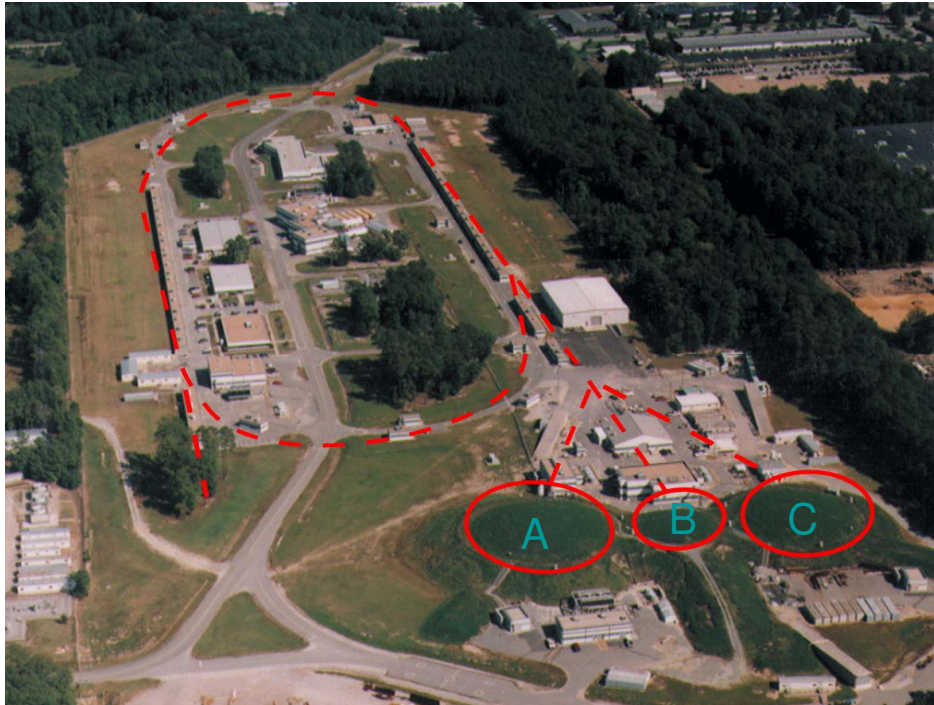
$$C_{1u} = 2 g_A^e g_V^u$$

$$C_{1d} = 2 g_A^e g_V^d$$

$$C_{2u} = 2 g_V^e g_A^u$$

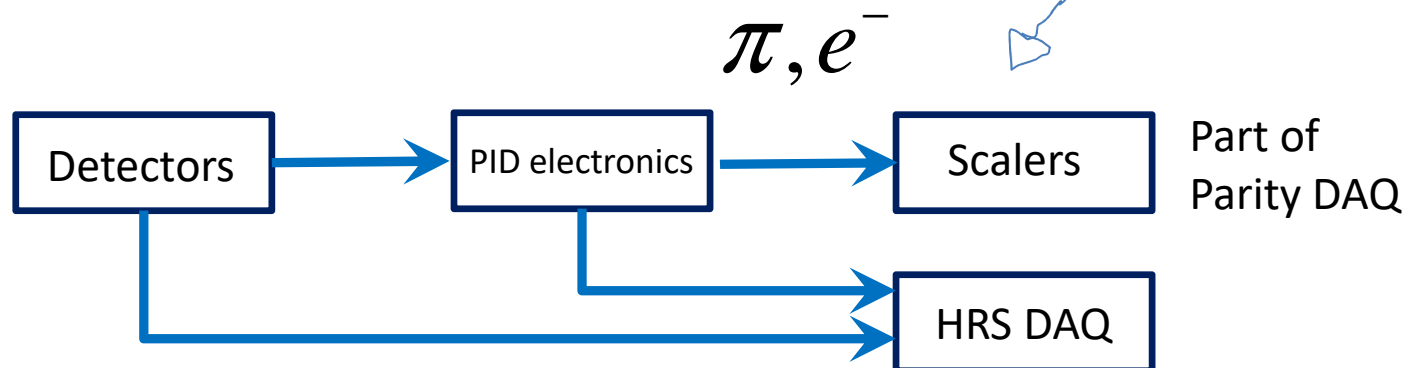
$$C_{2d} = 2 g_V^e g_A^d$$

# PVDIS at 6 GeV (Jefferson Lab)



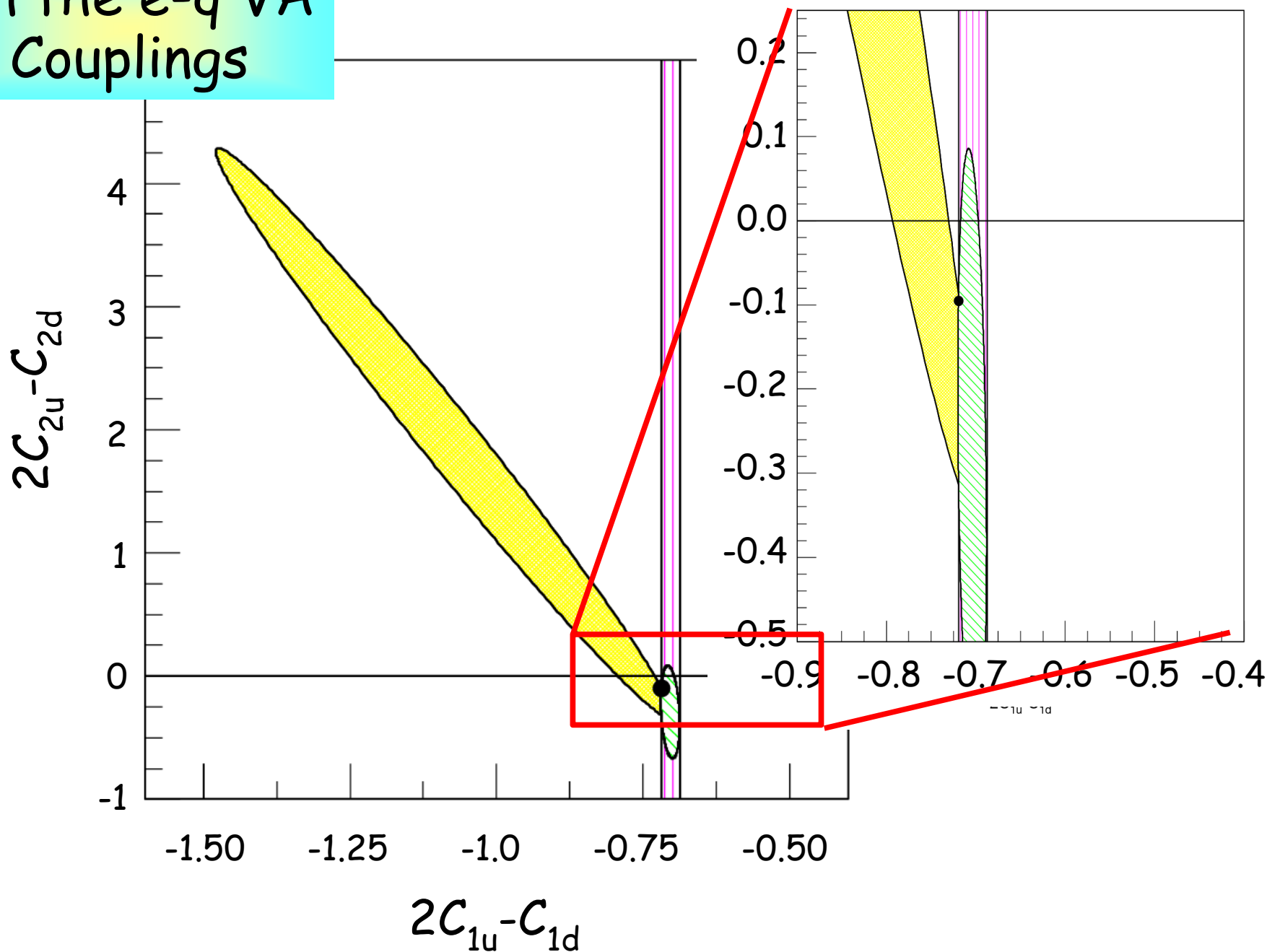
Spokespersons: Xiaochao Zheng, Paul Reimer, and Bob Michaels

- ◆ 100uA, 90% polarized beam on a 20cm liquid deuterium target
- ◆ Measured two DIS points:  $Q^2=1.085$  and  $1.901 \text{ GeV}^2$
- ◆ Ran in Nov-Dec. 2009
- ◆ **four publications**
  - Nature (main result)
  - Archival paper
  - Resonances
  - Specialized DAQ

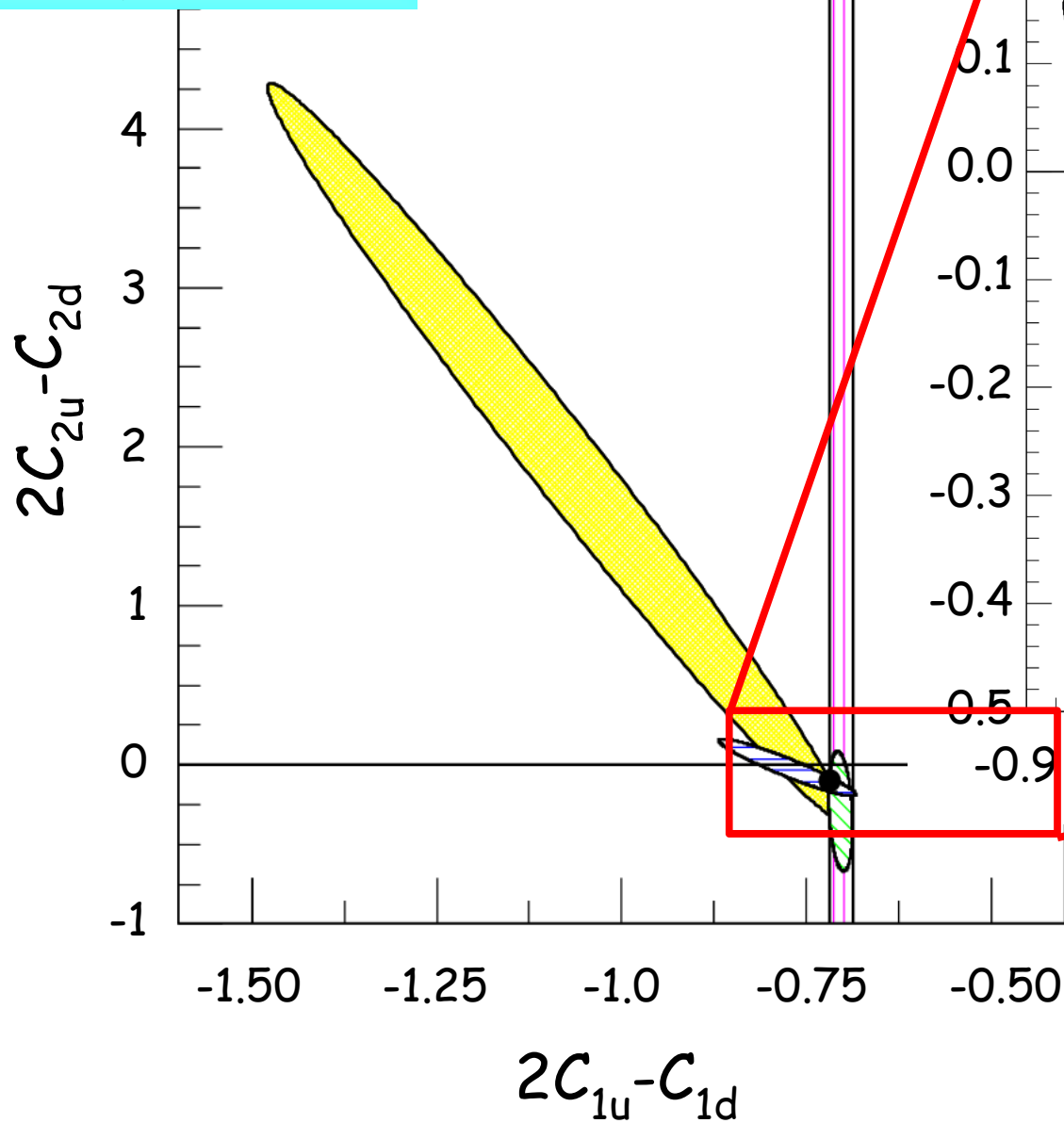


# On the e-q VA Couplings

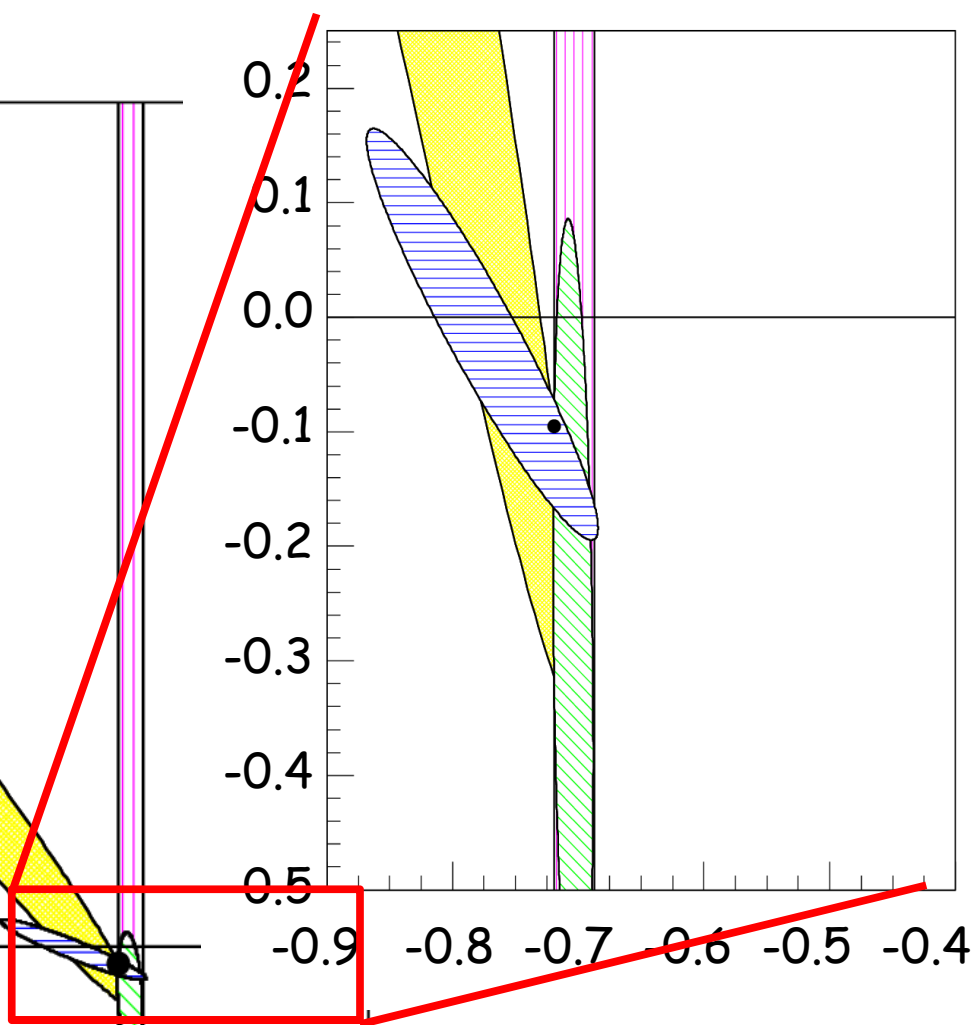
Previous data: E122, Elastic PVES + APV



# On the e-q VA Couplings

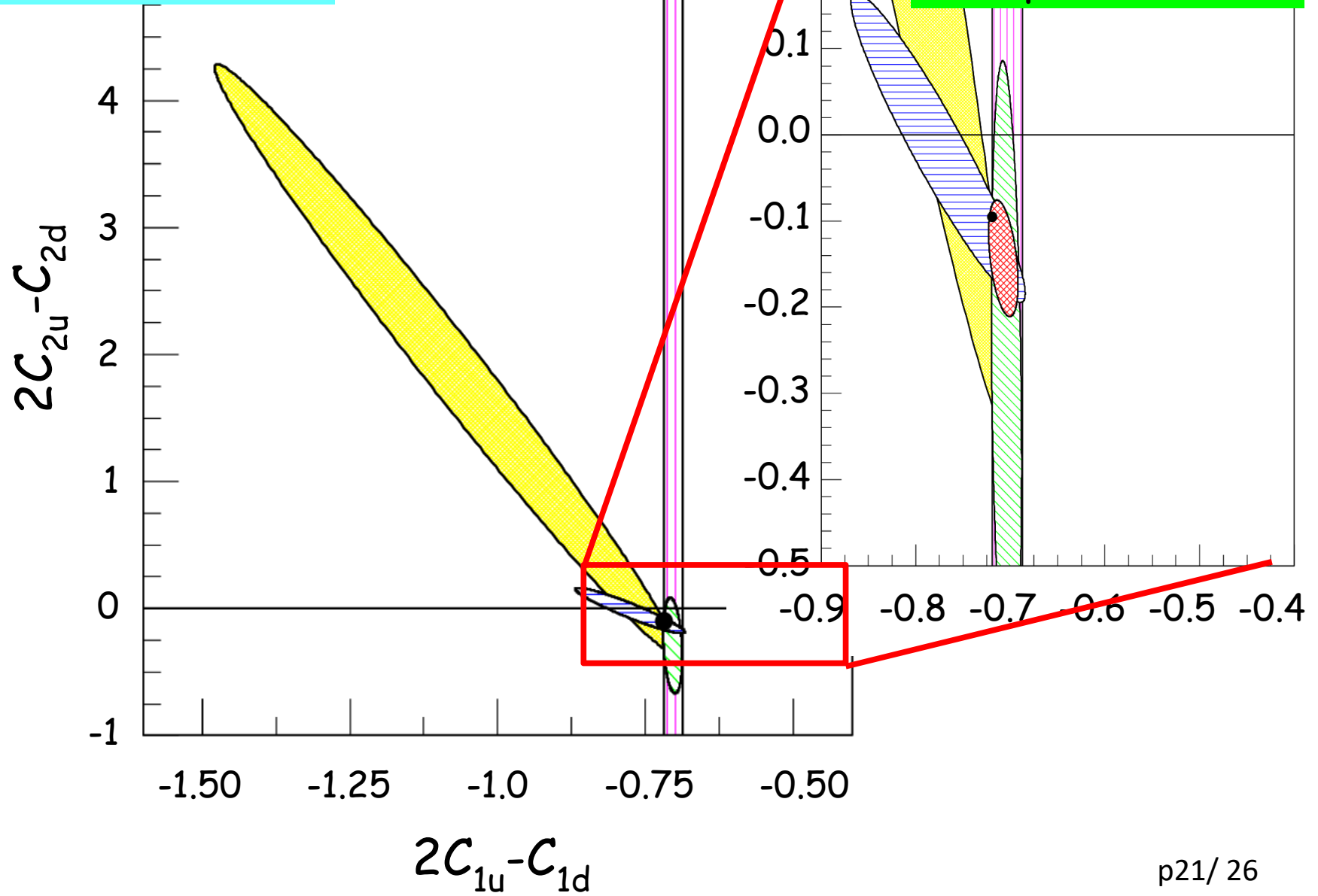


Add JLab 6 GeV PVDIS



# On the e-q VA Couplings

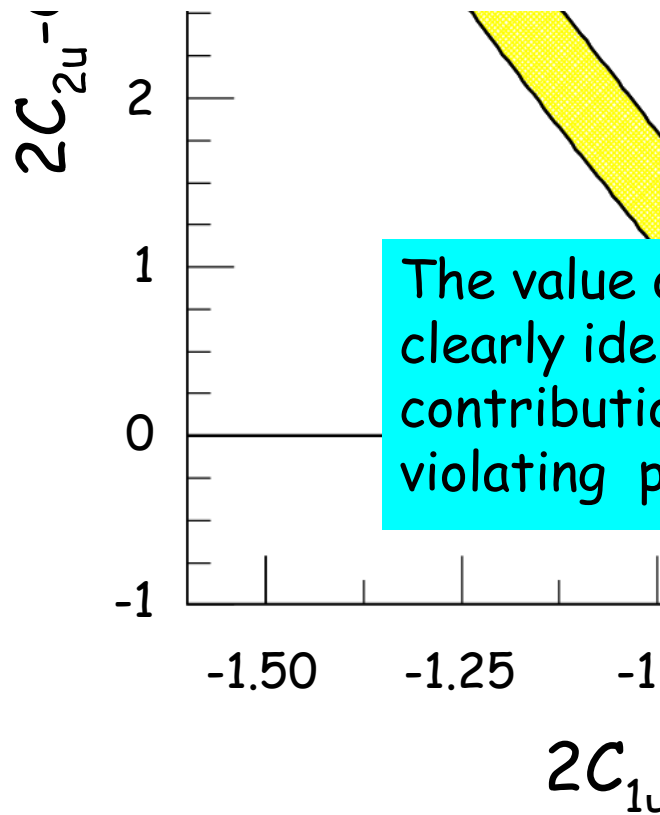
best fit



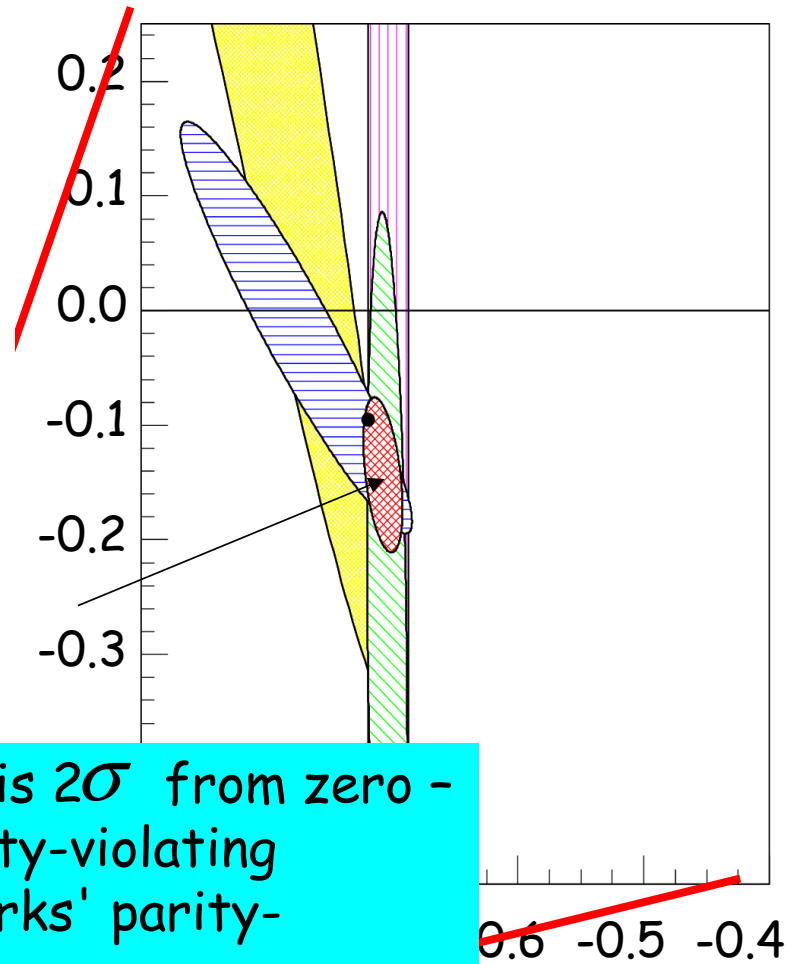
# Quarks are not ambidextrous

By separately scattering right- and left-handed electrons off quarks in a deuterium target, researchers have improved, by about a factor of five, on a classic result of mirror-symmetry breaking from 35 years ago. [SEE LETTER P.67](#)

Marciano., Nature 506, no. 7486, 43 (2014);  
 (Quarks are like people, most prefer to use their right hands, but some prefer left...)



The value of  $2C_{2u} - C_{2d}$  is  $2\sigma$  from zero - clearly identified parity-violating contribution from quarks' parity-violating property



"Measurement of parity violation in electron-quark scattering"  
 Wang et al., Nature 506, no. 7486, 67 (2014);

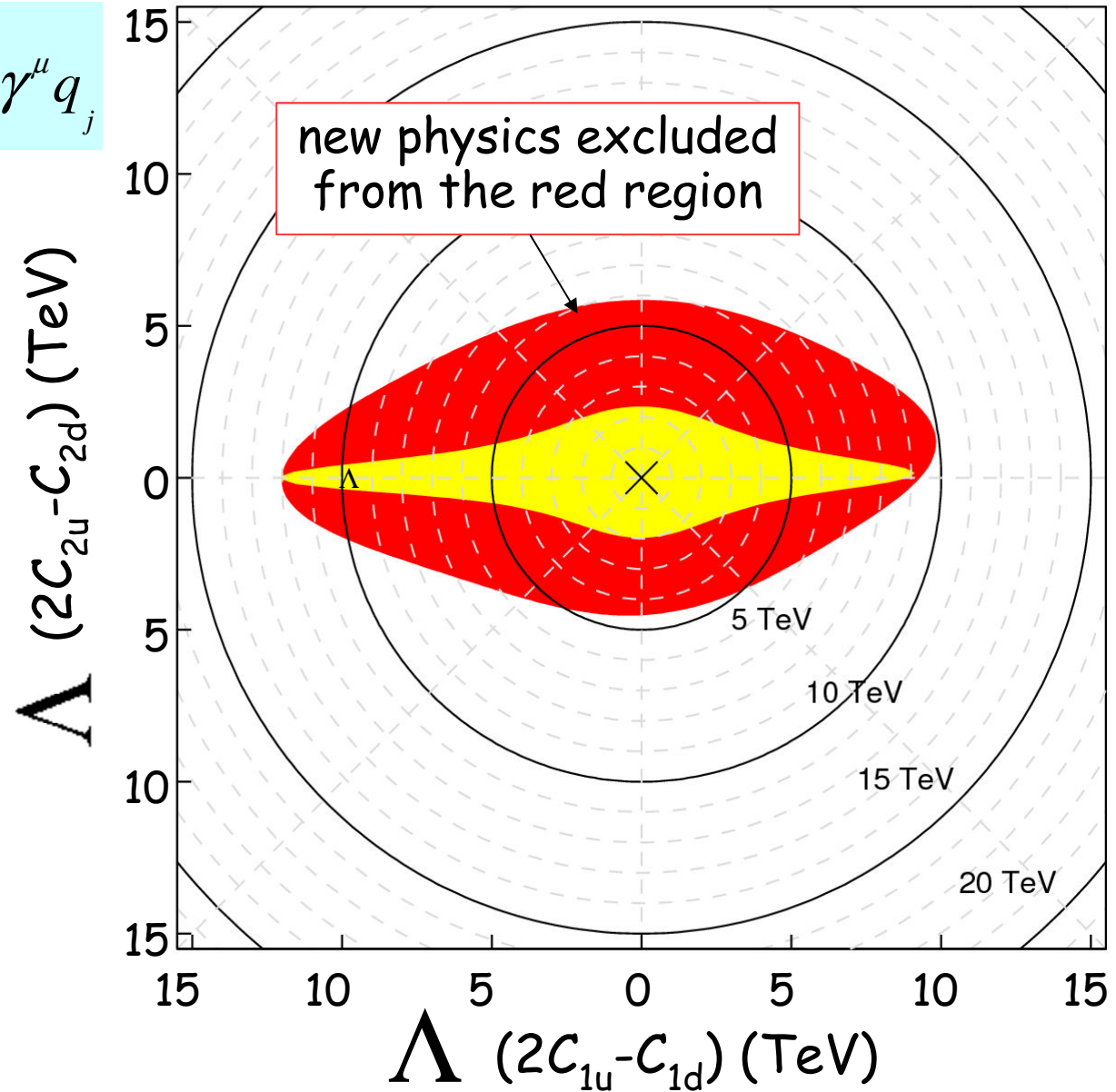
# Limit on new eq VA contact interactions

$$L_{eq} = \frac{g^2}{\Lambda^2} \sum_{i,j=L,R} \eta_{ij} \bar{e} \gamma_\mu e \bar{q} \gamma^\mu q_j$$

Mass Scale

$\Lambda$

$$(g^2 = 4\pi)$$

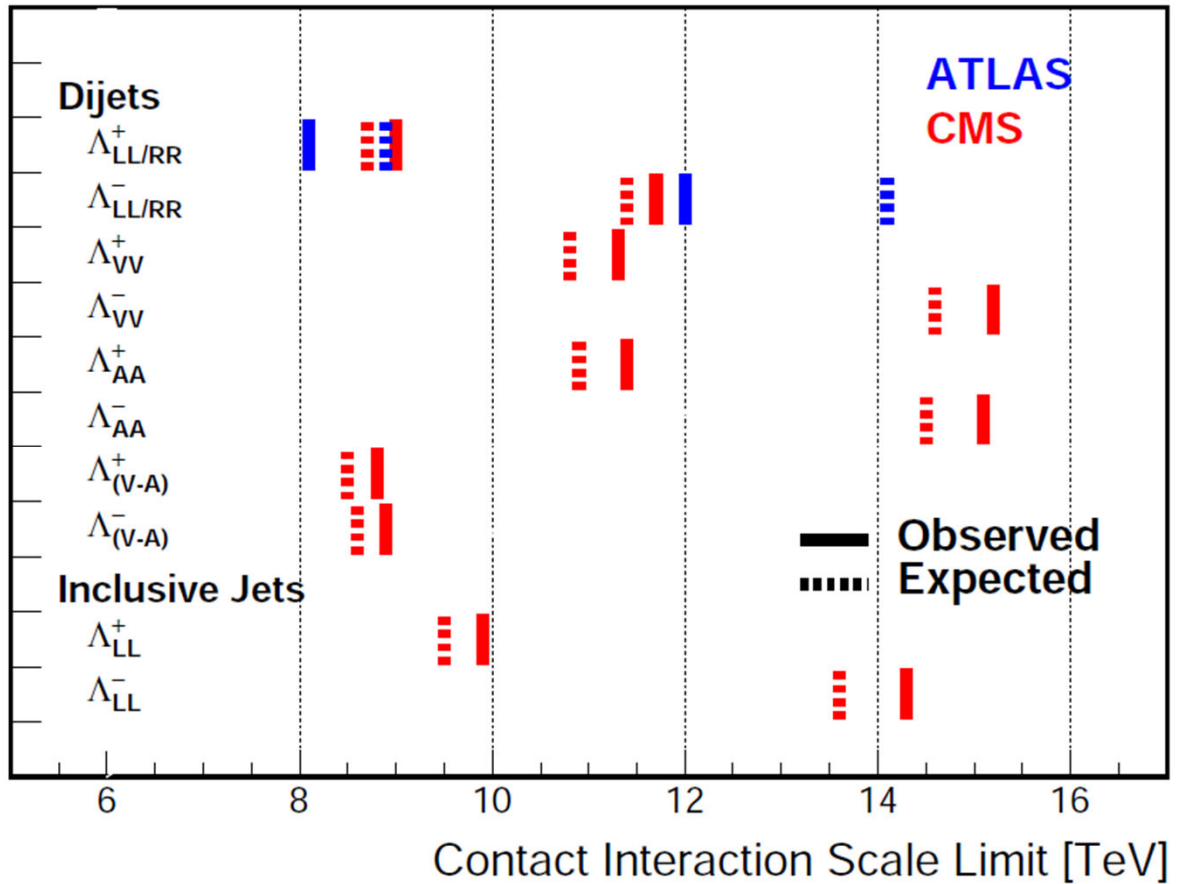




# Contact Interaction Limits from LHC (PDG)

No access to  $AV$  or  $VA$  terms

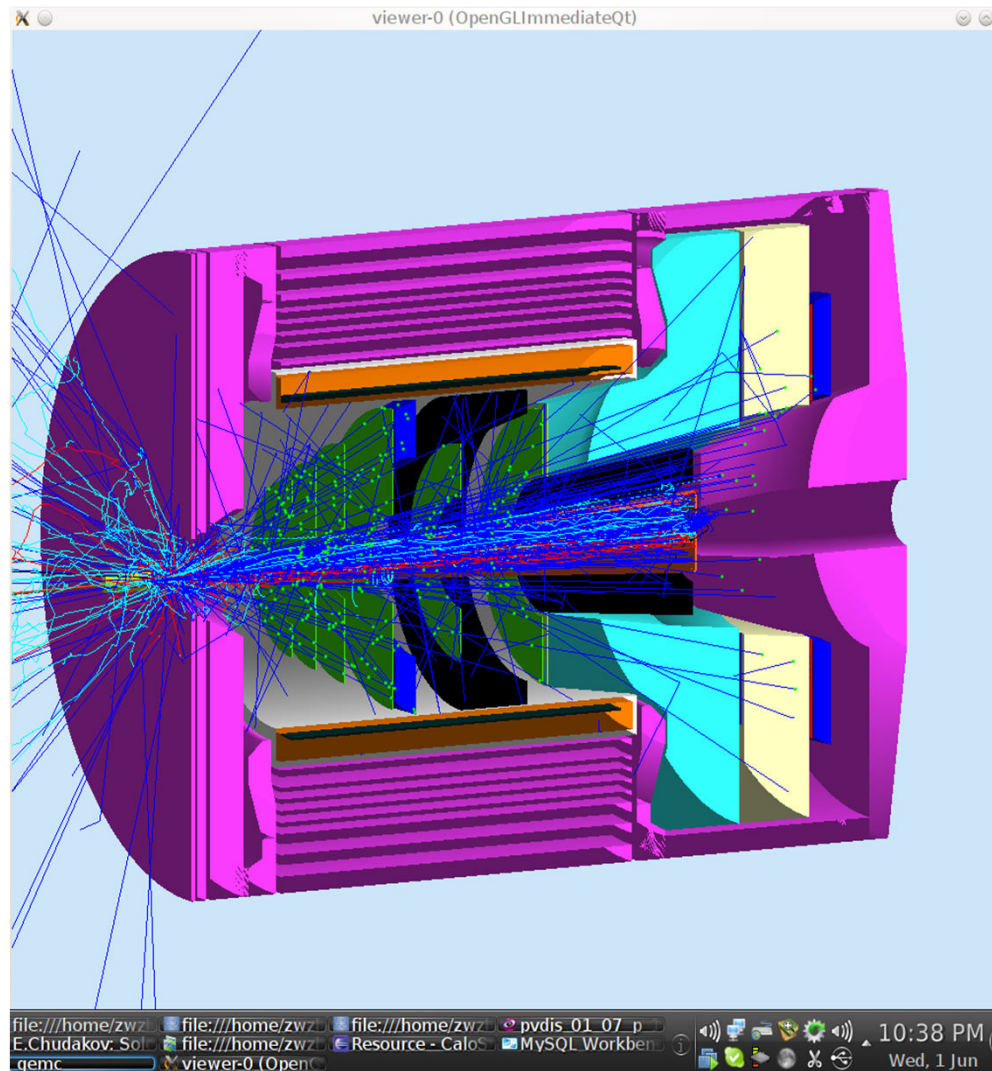
PVES is complementary to collider searches



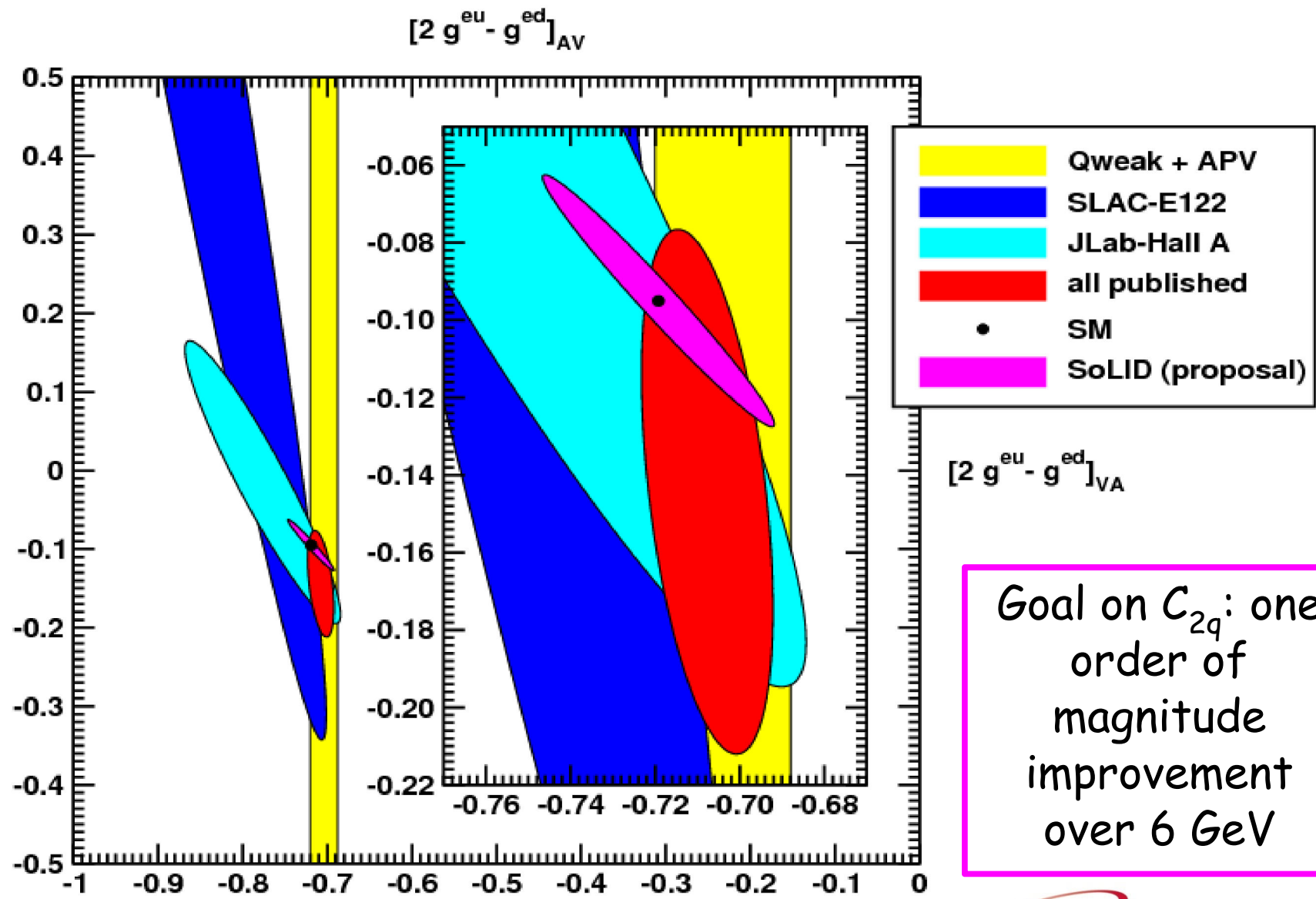


# Future: PVDIS Program with SoLID @ 12 GeV

## "Solenoid Large Intensity Device"



# Future: PVDIS Program with SoLID @ JLab 12 GeV



Robert Michaels, JLAB