

Figure 1: Radiative corrections to the measured asymmetries.

0.1 Proton Radiative Corrections

External radiative correction were applied following the formalism of Mo and Tsai [1], using the peaking approximation procedure of Ref. [2]. The internal radiative corrections utilized POLRAD [3] which had been updated to include a model of the RSS polarized data. This model was radiated, compared to the experimental data and adjusted accordingly in an iterative procedure until the chi-squared of the fit was minimized. The details of this fit have been previously detailed in the appendix of [4]. The size of the radiative corrections and the related systematic uncertainy can be seen in Fig. 1.

The radiative correction consisted of an additive and a multiplicative contribution:

$$A = \frac{A_{raw}}{f_{RC}} + A_{RC} \tag{1}$$

 f_{RC} contains contributions from the elastic tail subtraction only, while A_{RC} has both elastic and inelastic contributions. In the following sections we detail how these quantities are determined.

0.1.1 Elastic Tail Subtraction

The measured asymmetry contains a contribution from elastic tail as follows:

$$A_T^{\mathbf{r}} = \frac{\Delta_T^{\mathbf{r}}}{\Sigma_T^{\mathbf{r}}} = \frac{\Delta_{\mathrm{in}}^{\mathbf{r}} + \Delta_{\mathrm{el}}^{\mathbf{r}}}{\Sigma_{\mathrm{in}}^{\mathbf{r}} + \Sigma_{\mathrm{el}}^{\mathbf{r}}}$$
(2)

where the subscripts T,el and in refer to the total measured quantity, elastic and inelastic contributions respectively. The superscript r implies the the quantity is radiated. For example, $A_T^{\rm r}$ represents the measured asymmetry.

To remove the radiated elastic tail, we must deal with the polarized cross section Δ :

$$\Delta_{\rm in}^{\rm r} = \Delta_T^{\rm r} - \Delta_{\rm el}^{\rm r} \tag{3}$$

The measured asymmetry is related to the polarized cross section as:

$$A_{\rm in}^{\rm r} = \frac{\Delta_{\rm in}^{\rm r}}{\Sigma_{\rm in}^{\rm r}} \tag{4}$$

$$= \frac{\Delta_T^{\rm r} - \Delta_{\rm el}^{\rm r}}{\Sigma_{\rm in}^{\rm r}} \tag{5}$$

$$= \frac{\Sigma_T^{\mathbf{r}} A_T^{\mathbf{r}} - \Delta_{\mathrm{el}}^{\mathbf{r}}}{\Sigma_{\mathrm{in}}^{\mathbf{r}}} \tag{6}$$

$$= \left(\frac{\Sigma_T^{\rm r}}{\Sigma_{\rm in}^{\rm r}}\right) A_T^{\rm r} - \frac{\Delta_{\rm el}^{\rm r}}{\Sigma_{\rm in}^{\rm r}} \tag{7}$$

$$= \frac{1}{f}A_T^{\mathbf{r}} - A_{rc} \tag{8}$$

where we have made the following definitions:

$$f = \frac{\Sigma_{\text{in}}^{\text{r}}}{\Sigma_T^{\text{r}}} \tag{9}$$

$$A_{rc} = \frac{\Delta_{el}^{r}}{\Sigma_{in}^{r}} \tag{10}$$

So f_{RC} is completely determined by the model [5] of the unpolarized cross section f_{RC} is the only part of the radiative correction to affect the statistical errors. A_{RC} requires both the unpolarized model and knowledge of the elastic asymmetry [6] to form Δ_{el}^{r} .

0.1.2 Inelastic Radiative Correction

To be written...

References

- [1] L. W. Mo and Y.-S. Tsai, Rev. Mod. Phys. 41, 205 (1969).
- [2] S. Stein et al., Phys. Rev. D 12, 1884 (1975).
- [3] I. Akushevich, A. Ilichev, N. Shumeiko, A. Sorok, and A. Tolkachev, Comput. Phys. Commun. 104, 201 (1997).
- [4] F.R. Wesselmann et al. Phys.Rev. Lett 98, 132003 (2007).
- [5] E. Christy model, to be published.
- [6] Elastic Form Factors citation needed.