Timelike Compton Scattering with CLAS12 at Jefferson Lab

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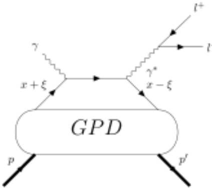
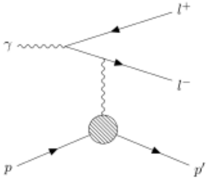
Timelike Compton Scattering

TCS: $\gamma p \to e^+ e^- p'$ **DVCS**: $ep \to e' p' \gamma$



TCS



Bethe-Heitler

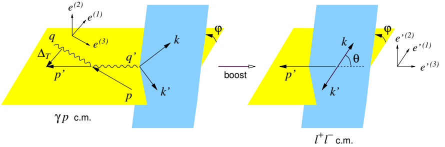


Figure in Berger et al., *EPJ C*, 2002

- BH cross section only depends on electromagnetic FFs
- Unpolarized interference cross section

$$\frac{d^4\sigma_{INT}}{dQ'^2 dt d\Omega} \propto \frac{L_0}{L}\left[\cos(\phi)\frac{1+\cos^2(\theta)}{\sin(\theta)}\mathrm{Re}\tilde{M}^{--} + \ldots\right]$$

- Polarized interference cross section

$$\frac{d^4\sigma_{INT}}{dQ'^2 dt d\Omega} = \frac{d^4\sigma_{INT}\,|_{\mathrm{unpol.}}}{dQ'^2 dt d\Omega} - \nu \cdot A\,\frac{L_0}{L}\left[\sin(\phi)\frac{1+\cos^2(\theta)}{\sin(\theta)}\mathrm{Im}\tilde{M}^{--} + \ldots\right]$$

$$\to \tilde{M}^{--} = \frac{2\sqrt{t_0 - t}}{M}\frac{1-\xi}{1+\xi}\left[F_1\mathcal{H} - \xi(F_1 + F_2)\tilde{\mathcal{H}} - \frac{t}{4M^2}F_2\mathcal{E}\right]$$

Both $\mathrm{Im}\mathcal{H}$ and $\mathrm{Re}\mathcal{H}$ can be accessed in TCS

Motivations to measure TCS

Test of universality of GPDs

- TCS is parametrized by GPDs
- **Comparison between DVCS and TCS** results allows to test the **universality** of GPDs
- TCS does not involve Distribution Amplitudes unlike Deeply Virtual Meson Production
  $\rightarrow$ direct comparison between DVCS and TCS

Real part of CFFs and nucleon D-term

- As for DVCS, TCS unpolarized cross section is sensitive to $\mathrm{Re}\mathcal{H}$, which is still not well constrained by existing data.
- The CFFs dispersion relation at leading order and leading twist :

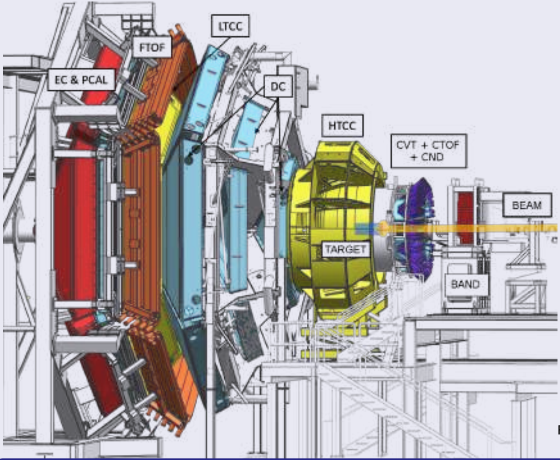
$$\mathrm{Re}\mathcal{H}(\xi, t) = \mathcal{P}\int_{-1}^{1} dx \left(\frac{1}{\xi - x} - \frac{1}{\xi + x}\right) \mathrm{Im}\mathcal{H}(\xi, t) + D(t)$$

- $D(t)$ can be related to the **mechanical properties** of the nucleon.

Review in Polyakov, Schweitzer, *International Journal of Modern Physics A*, 2018

Experimental setup

CLAS12



- **Forward Detector (6 sectors)**
  - Torus magnet
  - Drift Chambers
  - Forward Time-of-Flight
  - Calorimeters (EC and PCAL)
  - Cherenkov counters
- **Central Detector**
  - Solenoid magnet
  - Central Vertex Tracker (Silicon and micromegas)
  - Central Time-of-Flight
  - Central Neutron Detector

Figure in Burkert et al., *NIM A*, 2020

Data set used in this work

- **Fall 2018** run period
- ***LH_2*** target / **10.6 GeV** beam / RG-A
- Inbending torus magnetic field
- Accumulated charge: $\sim$ 150 mC ($\sim$ 200 fb^{-1})

Analysis strategy

Final state selection from PID

$$e\ p \to (e')\gamma p \to (X)\ e^+ e^- p'$$

Exclusivity cuts

$$p_X = p_{beam} + p_{target} - p_{e^+} - p_{e^-} - p_{p'}$$

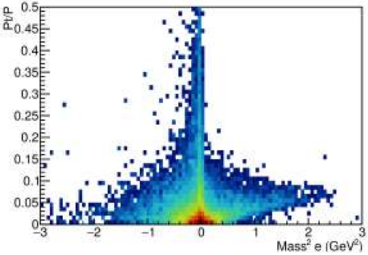
$$|M_X^2| < 0.4\ \mathrm{GeV}^2$$

Quasi-real photoproduction

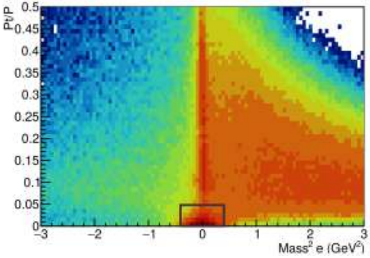
$$\frac{Pt_X}{P_X} < 0.05$$
$$\to Q^2 < 0.1\ \mathrm{GeV}^2$$

after momentum corrections and fiducial cuts

Simulation



Data



Positron identification

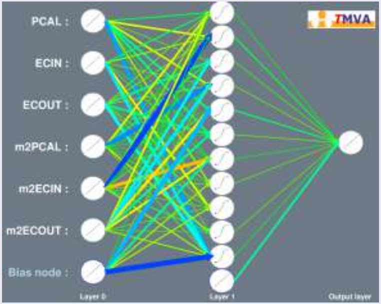
Definitions

Signal: e^+ identified as e^+ Background: π^+ identified as e^+

Strategy and discriminating variables

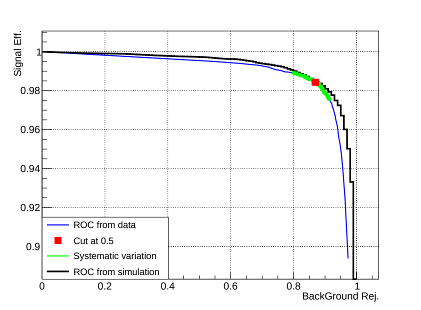
Positron: electromagnetic shower   Pion: Minimum Ionizing Particle (MIP)

$$SF_{\text{EC Layer}} = \frac{E_{dep}(\text{EC Layer})}{P} \qquad M_2 = \frac{1}{3}\sum_{U,V,W}\frac{\sum_{\text{strip}}(x-D)^2\cdot\ln(E)}{\sum_{\text{strip}}\ln(E)} \rightarrow \textbf{6 variables}$$



Output: Signal $\rightarrow$ 1   Background $\rightarrow$ 0

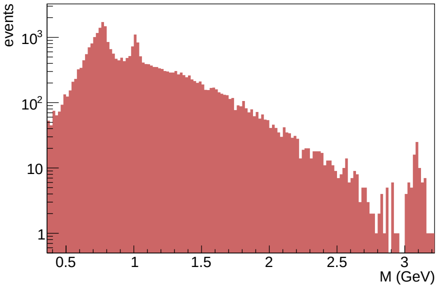
B/S from 50% to 5%



- Signal in data $\Rightarrow$ Outbending electrons
- Background in data $\Rightarrow$ $ep \rightarrow e\pi^+_{PID=e^+}(n)$

Data/Simulation comparison

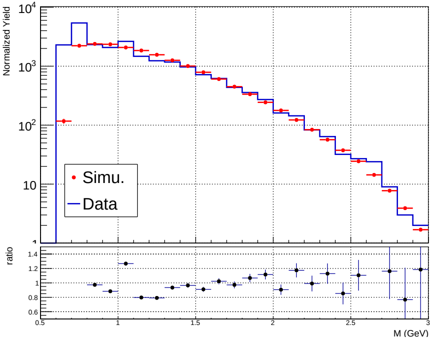
- Vector mesons peaks are visible in data: ω (770 MeV), ρ (782 MeV), Φ (1020 MeV) and J/ψ (3096 MeV)



- Data/BH comparison in the high mass region, no evident high mass vector meson production (ρ (1450 MeV, 1700 MeV))

Phase space of interest

- $1.5\ \text{GeV} < M_{e^+e^-} = \sqrt{Q'^2} < 3\ \text{GeV}$
- $0.15\ \text{GeV}^2 < -t < 0.8\ \text{GeV}^2$
- $4\ \text{GeV} < E_\gamma < 10.6\ \text{GeV}.$



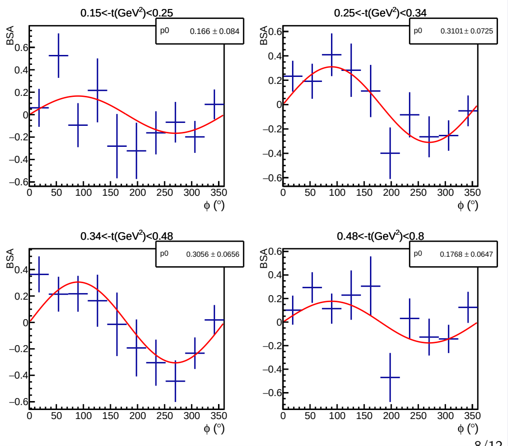
Observable 1: Photon polarization asymmetry (BSA)

Access to the imaginary part of CFFs

$$BSA = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} = \frac{-\frac{\alpha_{em}^3}{4\pi s^2}\frac{1}{-t}\frac{m_p}{Q'}\frac{1}{\tau\sqrt{1-\tau}}\frac{L_0}{L}\sin\phi\frac{(1+\cos^2\theta)}{\sin(\theta)}\mathrm{Im}\tilde{M}^{--}}{d\sigma_{BH}}$$

Experimental measurement

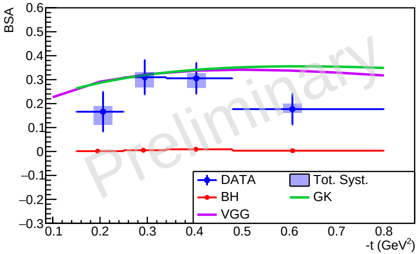
- $BSA(-t, E\gamma, M; \phi) = \frac{1}{Pol_{eff}}\frac{N^+ - N^-}{N^+ + N^-}$ where $N^{\pm} = \sum \frac{1}{Acc} Pol_{transf.}$
- $Pol_{transf.}$ is the **transferred polarization** from the **electron to the photon**
- Pol_{eff} is the **polarization of the CEBAF** electron beam (85%)
- The ϕ-distribution is fitted with a sine function



BSA selected results

- First time measurement
- A **sizeable asymmetry** is measured (above the expected vanishing BSA of BH)
  $\rightarrow$ **signature of TCS**
- Theoretical predictions were provided by M.Vanderhaeghen (using the VGG model) and P.Sznajder (using the GK model)
- Size of the asymmetry is **well reproduced** by VGG and GK models
  $\rightarrow$ **model dependent hints for universality of GPDs**

$< M >= 1.8$ GeV; $< E_\gamma >= 7.29$ GeV; $< \theta >= 92°$



Observable 2: Forward-Backward asymmetry

- Concept explored for J/Ψ production (Gryniuk, Vanderhaeghen, *Phys. Rev. D*, 2016).
- Exploratory studies for TCS performed in my thesis.
- Very first predictions for TCS have been published very recently (Heller, Keil, Vanderhaeghen, *Phys. Rev. D*, 2021).
- Use the different parity of the TCS and BH amplitudes under the inversion of the leptons directions $k \leftrightarrow k' \iff (\theta, \phi) \leftrightarrow (180^\circ - \theta, 180^\circ + \phi)$

BH cross section

$$\frac{d\sigma_{BH}}{dQ^2\,dt\,d\Omega} \propto \frac{1+\cos^2\theta}{\sin^2\theta} \xrightarrow{FB} \frac{d\sigma_{BH}}{dQ^2\,dt\,d\Omega}$$

Int. cross section

$$\frac{d^4\sigma_{INT}}{dQ'^2dtd\Omega} \propto \frac{L_0}{L}cos(\phi)\frac{1+cos^2(\theta)}{sin(\theta)} \xrightarrow{FB} -\frac{d\sigma_{INT}}{dQ^2\,dt\,d\Omega}$$

A_{FB} formula

$$A_{FB}(\theta_0,\phi_0) = \frac{d\sigma(\theta_0,\phi_0) - d\sigma(180^\circ - \theta_0, 180^\circ + \phi_0)}{d\sigma(\theta_0,\phi_0) + d\sigma(180^\circ - \theta_0, 180^\circ + \phi_0)} = \frac{-\frac{\alpha_{em}^3}{4\pi s^2}\frac{1}{-t}\frac{m_p}{Q'}\frac{1}{\tau\sqrt{1-\tau}}\frac{L_0}{L}\cos\phi_0\frac{(1+\cos^2\theta_0)}{\sin(\theta_0)}\mathrm{Re}\tilde{M}^{--}}{d\sigma_{BH}}$$

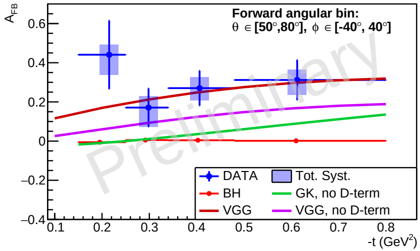
- Access to the real part of the CFFs with no integration over angles
- Removes large dependencies on angular acceptance $\rightarrow$ direct comparison with models
- But smaller phase space $\rightarrow$ lower statistics

A_{FB} selected results

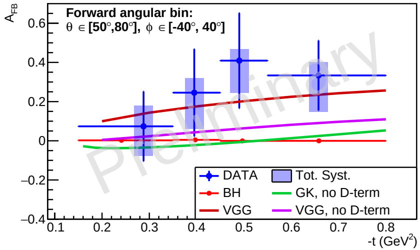
- A_{FB} measured in two mass regions: $M \in [1.5\ \text{GeV}, 3\ \text{GeV}]$ and $M \in [2\ \text{GeV}, 3\ \text{GeV}]$ (known resonance-free region)
- The measured A_{FB} is non-zero: **evidence for signal** beyond pure BH contribution
- Three model predictions
  1. VGG without D-term
  2. VGG with D-term

     D-term in Pasquini et al., *Physics Letters B*, 2014
  3. GK without D-term
- Measured asymmetry is better reproduced by the VGG model **including the D-term** in both mass bins

$$< M >= 1.8\ \text{GeV}; < E_\gamma >= 7.24\ \text{GeV}$$



$$< M >= 2.25\ \text{GeV}; < E_\gamma >= 8.13\ \text{GeV}$$



Conclusions

- TCS observables were measured for the **first time**
- Sizeable BSA (sensitive to Im$\mathcal{H}$) and A_{FB} (sensitive to Re$\mathcal{H}$) are **clear signatures of TCS**
- The results obtained allow to draw physical conclusions:
  - the BSA is well reproduced by models that reproduce existing DVCS data
    $\rightarrow$ hints for **universality of GPDs**
  - the Forward/Backward asymmetry appears to be better reproduced by model with a D-term
    $\rightarrow$ promising path to the measurement of the **mechanical properties of the proton**
- **An article is underway, so stay tuned !**