### Vector Meson Photoproduction at Threshold: from Omega to Upsilon

Igor Strakovsky The George Washington University





IS, S. Prakhov, Ya. Azimov et al, Phys Rev C 91, 045207 (2015)

IS, D. Epifanov, & L. Pentchev, Phys Rev C **101**, 042201 (2020)

IS, L. Pentchev, & A.I. Titov, Phys Rev C **101**, 045201 (2020)

Meng-Lin Du, V. Baru, Feng-Kun Guo, Ch. Hanhart, U.-G. Meissner,

A. Nefediev, & IS, Eur Phys J C **80**, 1053 (2020)

L. Pentchev & IS, Eur Phys J A **57**, 56 (2021)

IS, W.J. Briscoe, L. Pentchev, & A. Schmidt, Phys Rev C 104, 074028 (2021) Xiao-Yun Wang, Fancong Zeng, & IS, Phys Rev C 106, 015202 (2022)

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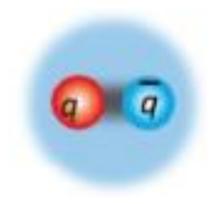


### Outline



- Vector Meson Zoo
  - VM Nucleon SL
  - Thr Kinematics & EM Properties of VM
  - VMD Phenomenology
  - Brief tour through Photoproduction Experiments
  - VM Nucleon SLs
- Outlook: A View of Future

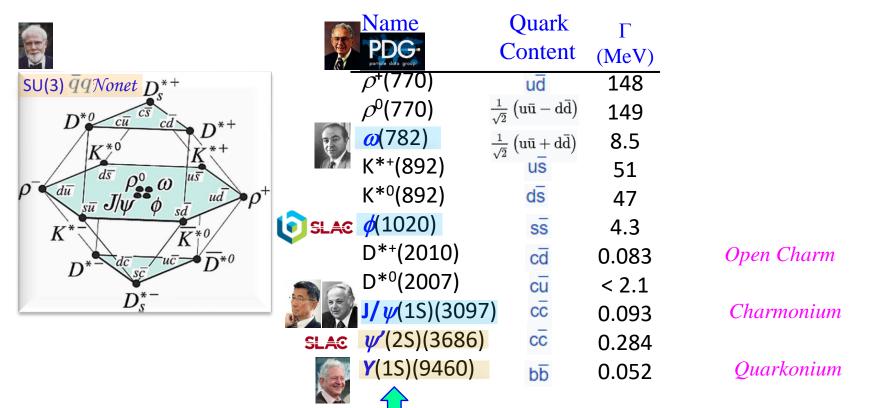
# Vector Meson goo



#### Vector Meson Domestic Zoo

- Some *vector mesons* can, compared to other mesons, be measured to very high precision.
- This stems from fact that *vector mesons* have *same* quantum numbers as *photon*.

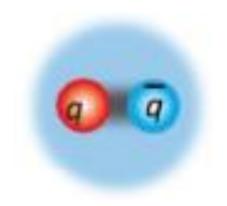
$$I^{G}(J^{PC}) = O^{-}(1^{--})$$

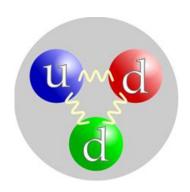


• We will focus on 5 *vector mesons* from  $\overline{q}q$  *Nonet* which widths are **narrow** enough to study *meson photoproduction* @ threshold & where data are available.



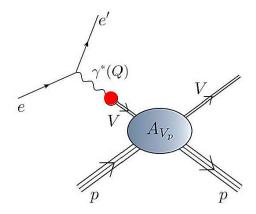
### Vector Meson — Mucleon SI







• Don't have *vector meson* beams, so experiments @ modern EM-accelerators attempt to access such interactions via EM production reactions  $ep \rightarrow e'Vp$ .



• Interaction of *heavy vector mesons*,  $J/\psi$  or Y, with *proton* offers prospects for access to *QCD van der Waals* interaction, generated by multiple gluon exchange may relate to observation, *e.g.*, of *hidden-charm* 5q states.



• *Gluonic van der Waals* interaction between color singlet hadrons can be described in *QCD* & is equivalent to *EM* interaction between two neutrally charged atoms.

#### Vector Meson - Nucleon Scattering Length Determination

IS, D. Epifanov, & L. Pentchev, Phys Rev C 101, 042201 (2020) IS, L. Pentchev, & A.I. Titov, Phys Rev C 101, 045201 (2020)

• Small *positive* or *negative VN SL* may indicate weakly *repulsive* or *attractive VN* interaction if there is no VN bound state below experimental  $q_{min}$ .



• For evaluation of absolute value of VN SL, we apply VMD approach that links near-threshold photoproduction Xsections of  $\gamma p \rightarrow Vp$  & elastic  $Vp \rightarrow Vp$ 

$$\frac{d\sigma^{\gamma p \to Vp}}{d\Omega}|_{\text{thr}} = \frac{q}{k} \frac{1}{64\pi} |T^{\gamma p \to Vp}|^2 = \frac{q}{k} \cdot \frac{\pi \alpha}{g_V^2} \frac{d\sigma^{Vp \to Vp}}{d\Omega}|_{\text{thr}} = \frac{q}{k} \frac{\pi \alpha}{g_V^2} |\alpha_{Vp}|^2$$

k is photon CM momentum  $k = (s - M^2) / 2 s^{1/2}$ 



q is vector-meson CM momentum

 $T^{p \rightarrow Vp}$  is the invariant amplitude of vector-meson photoproduction





 $g_V$  is VMD coupling constant, related to vector-meson EM decay width  $\Gamma(V \to e^+e^-)$ 

$$g_V^2 = \frac{\pi \cdot \alpha^2 \cdot m_V}{3 \cdot \Gamma(V \to e^+ e^-)}$$

• Finally, one can express absolute value of VN SL as product of pure **EM VMD**-motivated kinematic factor

$$B_V^2 = \frac{\alpha \cdot m_V \cdot k}{12\pi \cdot \Gamma(V \to e^+e^-)}$$





& hadronic factor  $h_{Vp} = \sqrt{b_1}$ 

$$h_{Vp} = \sqrt{b_1}$$

where  $b_1$  came from **best fit** 

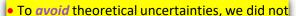
$$\sigma_t(q) = 0 \cdot q + b_3 \cdot q^3 + b_5 \cdot q^5$$





that is determined by interplay of strong (hadronic) & EM dynamics as

$$|\alpha_{Vp}| = B_V \cdot h_{Vp}$$



- determine sign of SL,
- separate Re & Im parts of SL,
- extract spin 1/2 & 3/2 contributions.



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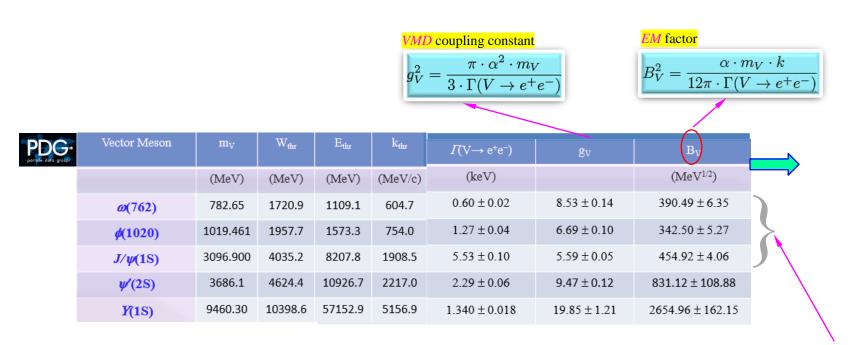
### Threshold Kinemalics



### Vector Meson Properties

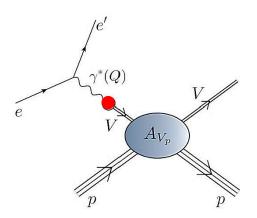


### Kinematics for VM Photoproduction off Proton @ Thresholds & VM EM Properties



• EM factor  $B_V$  for low-lying vector meson is close to each other.

# VIII Phenomenology



#### VMD Model

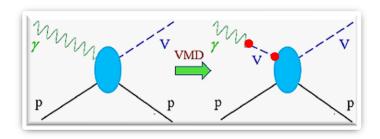
• Vector Meson Dominance model
relying on transparent current-field identities

M. Gell-Mann & F. Zachariasen, Phys Rev 124, 953 (1961)

J.J. Sakurai, Currents and Mesons (The University of Chicago Press, 1969)

N.M. Kroll, T.D. Lee, & B. Zumino, Phys. Rev. 157, 1376 (1967)

• In *VMD*, *real photon* can *fluctuate* into virtual *vector meson*, which subsequently scatters off target proton.



• *VMD* does not contain *free parameters* & can be used for variety of qualitative estimates of observables in *vector meson* photoproductions @ least as first step towards their more extended theoretical studies.

#### VMD for VN Interaction

• There is no alternative VMD to get  $J/\psi p$  SL from meson photoproduction. [Possible alternative is to develop sophisticated, nonperturbative Courtesy of Arkady Vainshtein & Misha Ryskin, July 2020 reaction theory that can explain *quark+anti-quark* scattering from hadron targets into vector meson final-states.]

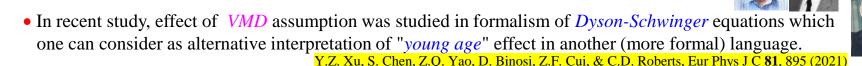
• To estimate theoretical uncertainty related to VMD model, one refer to estimation of cross section of  $J/\psi$ photoproduction in *peripheral model* & found strong energy dependence close to threshold because non-diagonal  $p \rightarrow Vp$  & elastic  $Vp \rightarrow Vp$  must have larger transfer momenta vs elastic scattering. This result in violation of *VMD* by factor of **5**.

> K.G. Boreskov & B.L. Ioffe, Sov J Nucl Phys 25, 331 (1977) B.Z. Kopeliovich, I. Schmidt, & M. Siddikov, Phys Rev C 95, 065203 (2017)

• Color factor for *charmonium* is 1/9 while for *open charm* is 8/9.



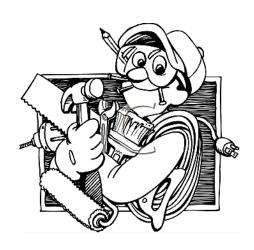
• Strong suppression in  $\overline{VN}$  interaction close to threshold is observed because of  $\overline{qq}$  pair in point-like configuration lacks sufficient time to form complete wave function of vector meson; that is, proton interacts with "young" (undressed) vector meson whose size is smaller than that of "old" one participating in elastic  $Vp \rightarrow Vp$  scattering. E.L. Feinberg, Sov Phys Usp, 23, 629 (1980); Courtesy of Misha Ryskin, July 2020





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# Brief Sour through Photoproduction Experiments





### P2 for Omega





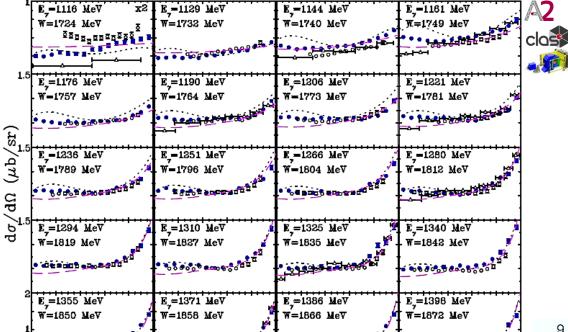


### $\gamma p \rightarrow \omega p \rightarrow \pi^0 \gamma p \rightarrow 3\gamma p$ Measurements

 $BR(\omega p \rightarrow \pi^0 \gamma) = 8.4\%$ 

IS, S. Prakhov, Ya. Azimov et al, Phys Rev C 91, 045207 (2015)





-0.6 -0.2 0.2

0.6

 $\cos\theta$ 

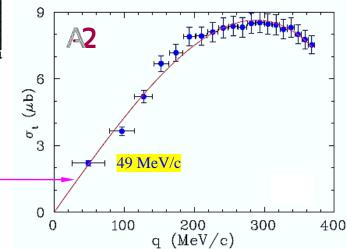
1 -0.6 -0.2 0.2 0.6

- Full production-angle coverage allows to determine σ.
- Legendre polynomial extension

$$d\sigma/d\Omega(E_{\gamma},\cos\theta) = \sum_{j=0} A_{j}(E_{\gamma})P_{j}(\cos\theta)$$

confirms  $\sigma_t$  determination

$$\sigma_t = 4\pi A_0(E_\gamma)$$





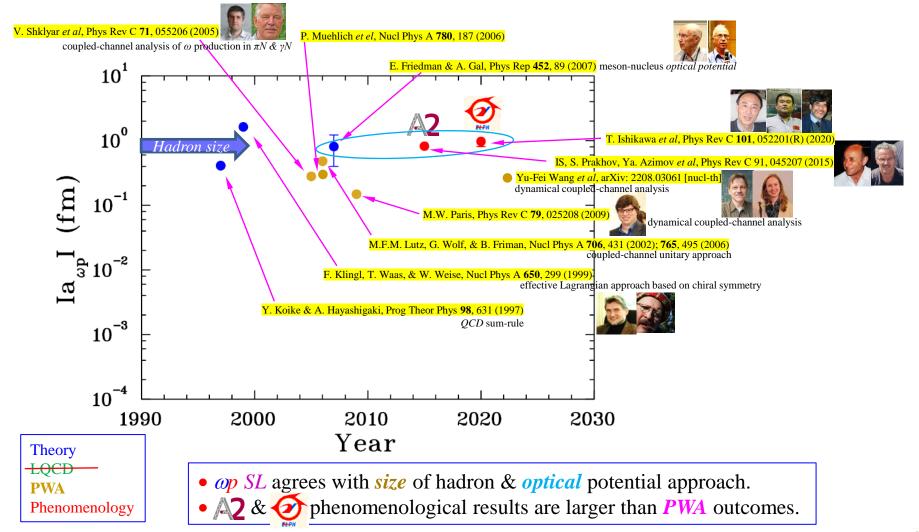
 $q + b_3 \cdot q^3 + b_5 \cdot q^5$ 

-0.6 -0.2 0.2 0.6

0.0 2.0 2.0 0.8

#### What is Known for $\omega N$ Scattering Length

- To avoid theoretical uncertainties, we did not
  - determine sign of \$1.
  - separate Re & Im parts of SL,
  - extract spin 1/2 & 3/2 contributions.



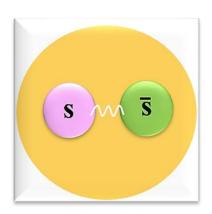
# Flab for Phi & \$/Psi







# Fab for Phi



BARYON2022, Seville, Spain, November 2022



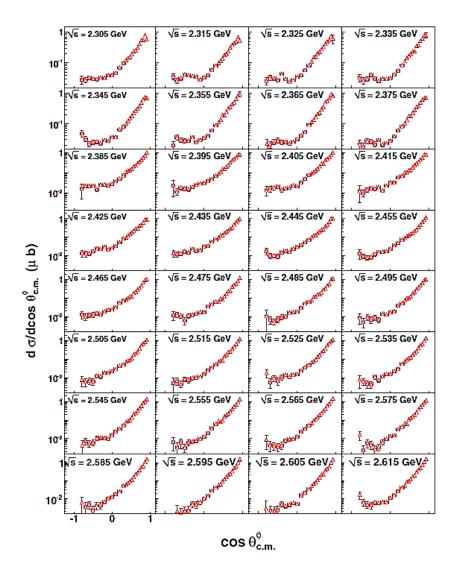




 $BR(\phi p \rightarrow K^+K^-) = 49.2\%$ 

B. Dey et al, Phys Rev C 89, 055208 (2014)







•  $\cos \theta$  of cos spans from -0.80 to 0.93.

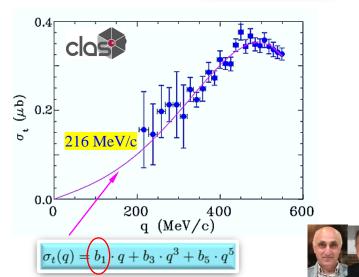


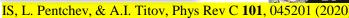
• Legendre polynomial extension

$$d\sigma/d\Omega(E_{\gamma}, \cos\theta) = \sum_{j=0} A_j(E_{\gamma})P_j(\cos\theta)$$

is way to determine  $\sigma_t$ 



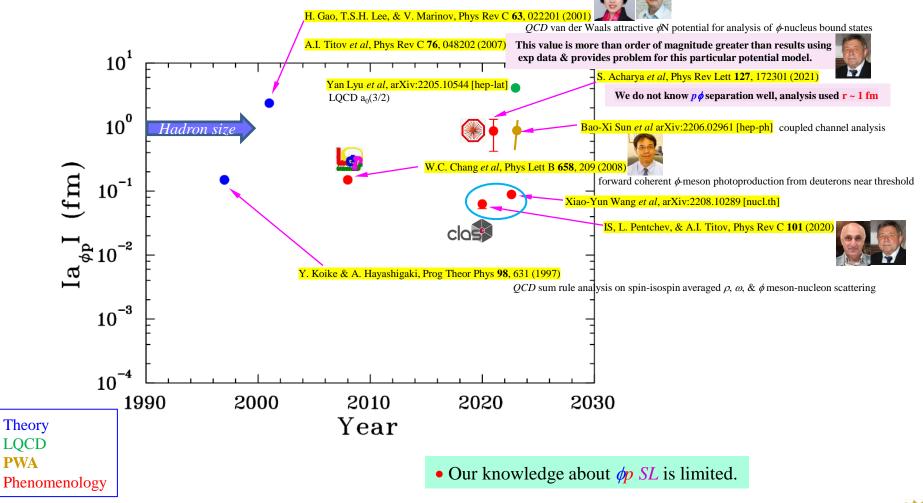






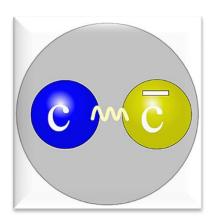
### What is Known for $\phi N$ Scattering Length

- To avoid theoretical uncertainties, we did not
  - determine sign of \$1.
  - separate Re & Im parts of \$1,
  - extract spin 1/2 & 3/2 contributions.





# Flab for 5/Psi

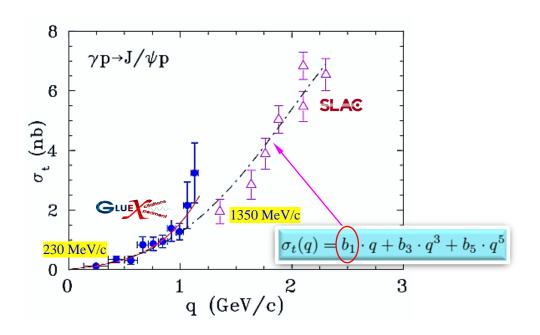


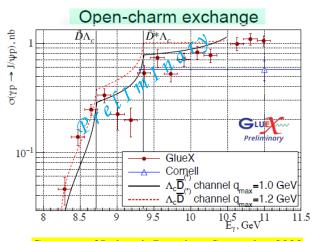




#### New J/\psi-p Scattering Length

• All previous theoretical results (including potential approaches & LQCD calculations) gave much-much larger SL.





Courtesy of Lubomir Pentchev, September 2022

$a_i$	GLUE	Guil & SLA
$a_1 [nb/(GeV/c)]$	$0.46 \pm 0.16$	$0.53 \pm 0.12$
$a_3 \left[ \text{nb}/(\text{GeV/c})^3 \right]$	$0.83 \pm 0.91$	$0.78 \pm 0.16$
$a_5 \left[ \text{nb}/(\text{GeV}/c)^5 \right]$	$0.28 \pm 0.87$	$-0.06 \pm 0.03$
χ²/dof	0.67	0.98



• There is *no discrepancy* between

A. Ali *et al*, Phys Rev Lett **123**, 072001 (2019)

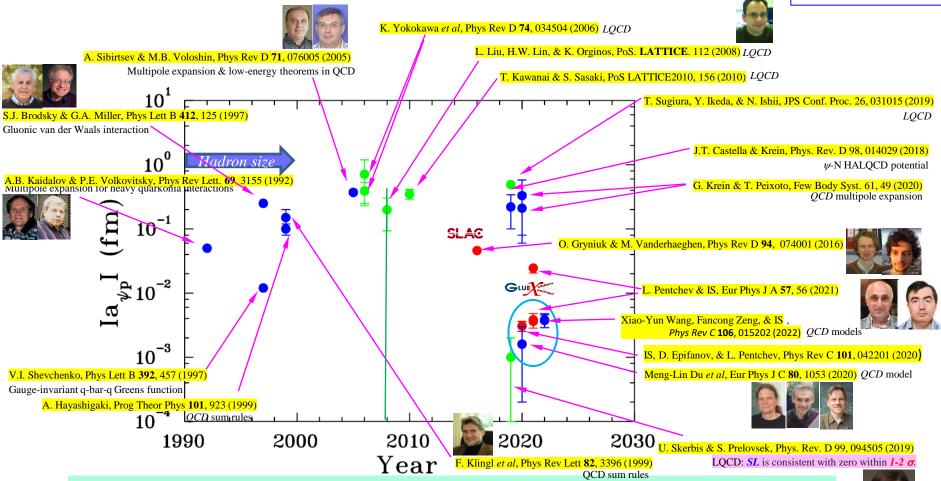
U. Camerini et al, Phys Rev Lett 35, 483 (1975)

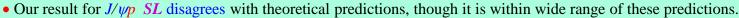


#### What is known for J/wp Scattering Length

- To *avoid* theoretical uncertainties, we did not
  - determine sign of \$1.
  - separate Re & Im parts of \$1.
  - extract *spin* 1/2 & 3/2 contributions.

Theory
LQCD
PWA
Phenomenology

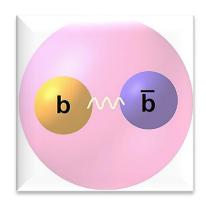


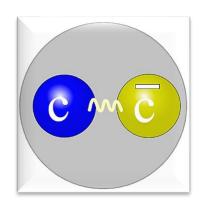


- They do not consider "young" meson effect.
- LQCD results are still uncertain.



### ESC for Upsilon & Psi'







### Expectation from

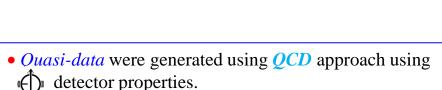
IS, W.J. Briscoe, L. Pentchev, & A. Schmidt, Phys Rev C 104, 074028 (2)



• *OCD* production amplitude can be factorized in terms of gluonic generalized parton distributions (GPD) & *quarkonium* distribution amplitude on one side & hard *quark-gluon* interaction on other side.

Y. Guo, X. Ji, & Y. Liu, Phys Rev D 103, 096010 (2021)

- - Just *theoretical* uncertainties.
  - Experimental uncertainties depend on luminosity, detector acceptance, & efficiency.
  - One can expect enormous Y rate, & uncertainties will be comparatively small.



Theoretical fit of Guilage data @ 95% C.L.

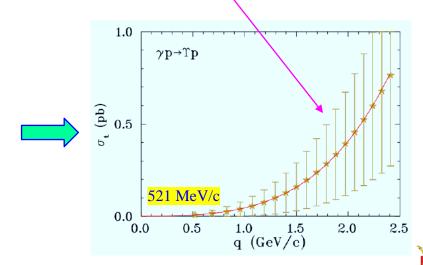
q (GeV/c)

CENTER for NUCLEAR FEMTOGRAPHY

γp→Yp x2000

2.0

- Further optimization of the low- $Q^2$  taggers may allow even smaller  $q_{min}$  to be achieved.
- It was assumed total integrated luminosity of 100 fb<sup>-1</sup> for photoproduction @ ( ), which corresponds to 116 days of beam with  $10^{34}$  cm<sup>-2</sup> s<sup>-1</sup>, for MC calculations. O. Gryniuk *et al*, Phys Rev D **102**, 014016 (2020)







6

5 -

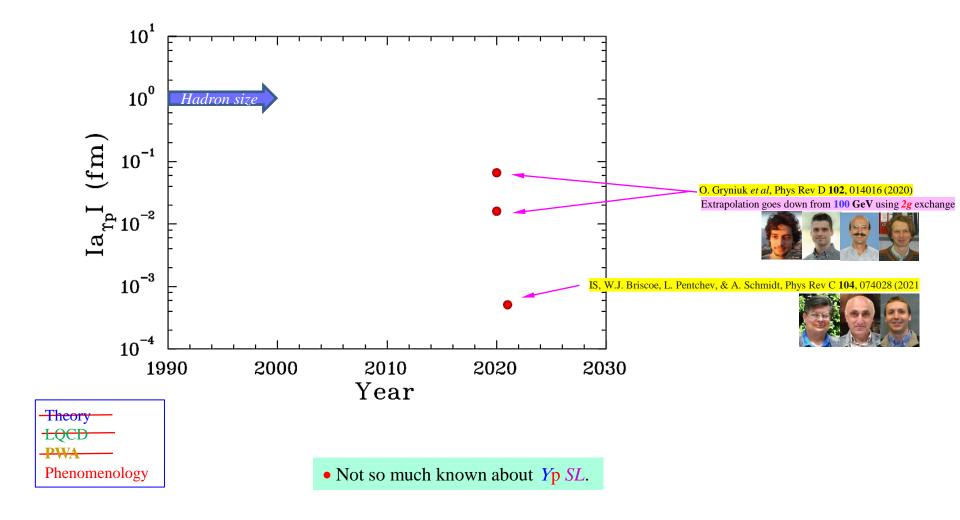
4

0.0

 $\gamma p \rightarrow J/\psi p$ 

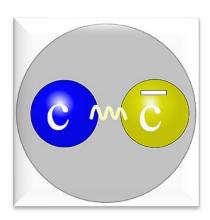
GLUE

### What may be known for Yp Scattering Length





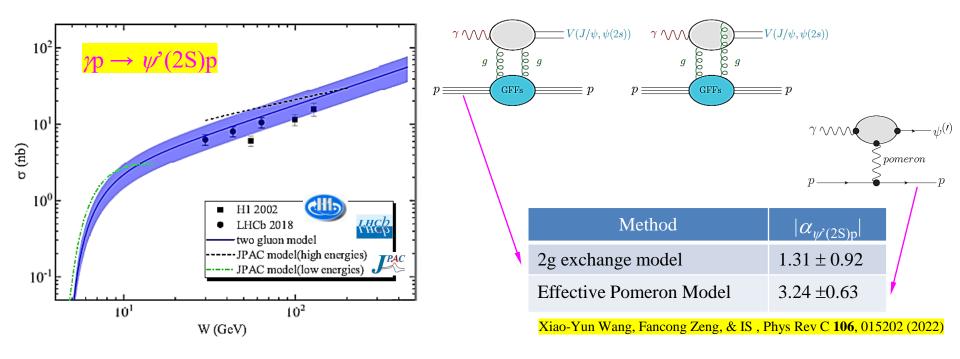
# ESC for Psi'







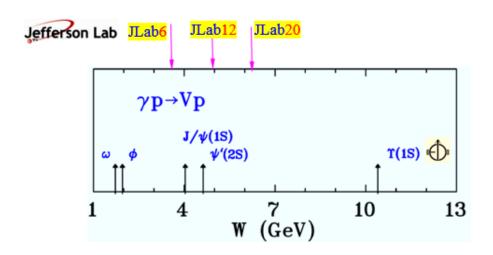
Xiao-Yun Wang, Fancong Zeng, & Quanjin Wang, Phys Rev D 105, 096033 (2022)



- Masses of  $\psi'(2S) \& J/\psi(1S)$  are close to each other but due to another *radial wave function*, *SL*s will be different.
- Phenomenological  $J/\psi(1S)$  SL agrees with theoretical  $\psi'(2S)$  SL.

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### Vector Meson - Mucleon SI



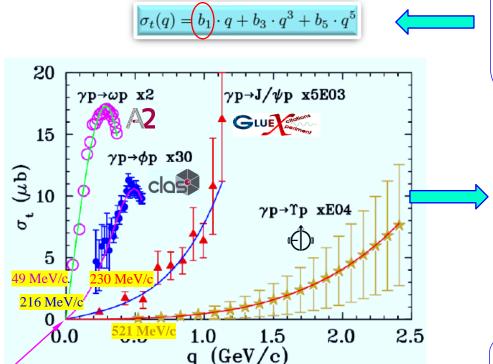


#### Total Cross Sections for Vector Meson Photoproduction off Proton

• Traditionally,  $\sigma_t$  behavior of near-threshold binary *inelastic* reaction

$$m_a + M_b < m_c + M_d$$

is described as series of *odd* powers in *q* (*even* powers in case of *elastic*).



- *Linear* term is determined by two independent *S*-waves only with total spin *1/2* &/or *3/2*.
- Contributions to *cubic* term come from both
   P-wave amplitudes &
   W dependence of S-wave amplitudes,
- *Fifth*-order term arises from *D*-waves & *W* dependencies of *S* & *P*-waves.

$$b_1 = (4.42 \pm 0.14) \times 10^{-2} \, \mu b / (MeV/c)$$
IS, S. Prakhov, Ya. Azimov *et al*, Phys Rev C **91**, 045207 (2015)
$$b_1 = (3.40 \pm 1.15) \times 10^{-4} \, \mu b / (MeV/c)$$
IS, L. Pentchev, & A.I. Titov, Phys Rev C **101**, 045201 (2020)

GLUE  $b_1 = (0.46 \pm 0.16) \times 10^{-6} \, \mu b/(MeV/c)$ 

IS, D. Epifanov, & L. Pentchev, Phys Rev C 101, 042201 (2020)

 $b_1 = (0.37 \pm 0.04) \times 10^{-9} \, \mu b/(MeV/c)$ 

S, W.J. Briscoe, L. Pentchev, & A. Schmidt, Phys Rev C **104**, 074028 (2021)

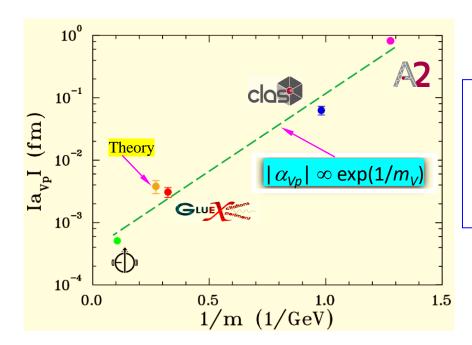
- Dramatic differences in hadronic factors  $\frac{h_{V_p} = \sqrt{b_1}}{\text{as slopes } (b_1) \text{ of } \sigma_t \text{ @ threshold as function of } q}$ varies significantly from  $\omega$  to  $\phi$  to  $J/\psi$ .
- Therefore, such big difference in SL is determined mainly by hadronic factor  $h_{Vp}$ .



• Our assumption is that

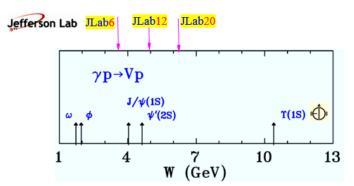
there is no VN bound state

below experimental  $q_{min}$ 



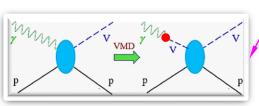
#### Vector Meson – Nucleon SL

- Such big difference in SLs of Vp systems is determined mainly by hadronic factor  $h_{V_p}$ , & reflects strong weakening of interaction in  $bb-p \& \bar{c}c-p$  systems compared to that of *light*  $\bar{q}q$ -p (q = u, d) configurations.
- Interaction in  $\bar{s}s-p$  has intermediate strength that is manifested in intermediate value of *p SL*.



• Such small value of  $\phi p$  SL compared to typical *hadron* size of 1 fm, indicates that proton is more transparent for *ϕ*-meson compared to *ω*-meson, & is much less transparent than for  $J/\psi$ -meson.

$$|\alpha_{\gamma_p}| \ll |\alpha_{\psi'p}| \leq |\alpha_{J/\psi p}| \ll |\alpha_{\phi p}| \ll |\alpha_{\omega p}|$$



- $p \rightarrow V$  coupling  $\bar{q}q$  is proportional to  $\alpha_s$  & separation of corresponding quarks.
- This *separation* (in *zero approximation*) is proportional to  $\frac{1}{m_V}$ .

# Outlook: P View of Future







- $|\alpha_{\mathsf{Vp}}| \ll |\alpha_{\mathsf{Vp}}| \leq |\alpha_{\mathsf{J/Vp}}| \ll |\alpha_{\mathsf{pp}}| \ll |\alpha_{\mathsf{pp}}|$ • It is remarkable that proton is quite so *transparent* to  $J/\psi$ , though general progression from  $\omega$  to  $\phi$  to  $J/\psi$  to probably  $Y & \psi'$
- Due to *small size* of "young" V vs "old" V, measured & predicted SL is very small. V crated by photon @ threshold then most probably V is not formed completely & its radius is smaller than that for normal ("old") V. Therefore, one observe stronger suppression for  $V_p$  interaction.
- *Light Vs* can be "young" as well. This depends on kinematics. Another point is that for slow *heavy* quark, one need more time to reach *equilibrium*, *i.e.*, to form final (long-living/static) V.
- Our phenomenology determined q-bar-q p SL which is smaller than V-p SLQuantitatively, there will be some difference between V-p SL & that for g-bar-q pair & p. Or our results are low level of **V-p** SL determination.
- Most theoretical calculations using gluonic van der Waals interaction disagree with our *phenomenological* results. Specifically, they do not consider *V young* effect.
- This should be calculated within some *model*. In general, result depends on energy, quark mass, & overlap integral between q-bar-q pair WF & V WF (this put some constrain on size of *q-bar-q* pair).
- We found strong exponential increase of V-p SL with inverse mass of Vs.  $|\alpha_{Vn}| \propto \exp(1/m_V)$



Igor Strakovsky



- Obviously, & facilities will open new *window* in solving the *VN SL* puzzle. It will allow to make deal with "*young*" *Y*-meson as well.
- Jefferson Lab upgrade will help to solve puzzle as well.
- It was observed that J/ψ-N cross section measured via J/ψ re-scattering/absorption inside nucleus is anomaly small in case of low energy photoproduction.
   This can be explained by fact that we dealt with "young" J/ψ of too small radius.
   Y-photoproduction on both proton & nucleus will extend our J/ψ study.
- In case of  $J/\psi$  (even Y) *electroproduction*, we deal with the "young"  $J/\psi(Y)$  for larger  $Q^2$  & we will have smaller formation time & correspondingly smaller radius of heavy *Charmonium* & *Quarkonium*.



Gracias por la invitación y su atención.



10/27/2022



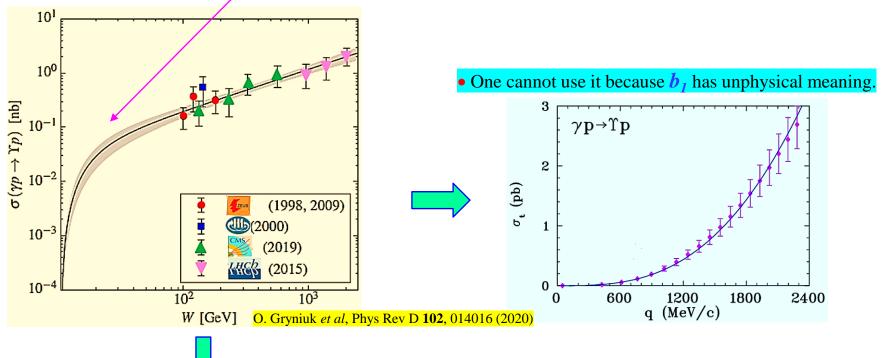




### Alternative Expectation from (1)



• Extrapolation goes down from 100 GeV using 2g exchange which disagreed with  $\bigcirc 1/\psi$  threshold data.





$$|\alpha_{\gamma_p}| = 0.066 \pm 0.001 \text{ fm}$$
 or  $|\alpha_{\gamma_p}| = 0.016 \pm 0.001 \text{ fm}$ 

Factor of difference with our phenomenology is 600 or 150