# The CLAS12 RICH Reconstruction

M. Contalbrigo – INFN Ferrara – on behalf of the CLAS12 RICH group

KAON2022 Workshop on Kaons with CLAS12

14<sup>th</sup> December 2022, Laboratori INFN di Frascati





# The CLAS12 RICH Reconstruction

Marco Contalbrigo – INFN Ferrara

KAON 2022: Workshop on kaons at CLAS12 – 14<sup>th</sup> December 2022



# RICH @ CLAS12







Completed in June 2022 with the symmetric configuration dedicated to the runs with polarized targets (now ongoing)







Features: large volume, > 50k readout channels, complicated topology

- ightarrow need large statistics for calibration
  - $\rightarrow$  calibration possible from DSTs (not only calibration runs)

single photon reconstruction for high-level pattern recognition and PID

- ightarrow require DC track impact point and angle, use EB PID
- $\rightarrow$  interplay (merge) with EB only possible for two sectors and relevant for multi-particle events
  - $\rightarrow$  post-process with specific hipo banks









#### Read CSG volumes from CAD stl files

**Convert volumes** into tracking surfaces (Shape3D) and spheres (Sphere3D) with given orientation

Each Sphere3D has an associated Shape3D to define its solid angle of acceptance

Align the tracking surfaces (as per mounting points)

Global RICH Layer (aerogel, MaPMTs, spherical mirror assembling)

Components (each single mirror, aerogel tile)

**Detail MaPMT pixel geometry** (on the misaligned plane)







# **Ray Tracing**



Complex geometry with various photon paths (reflections) off the same particle

## From CLAS12:

particle momentum photon emission point

## From RICH:

hit time and position

## **Direct ray-tracing:**

assume an ID hypothesis (e,  $\pi$ , k, p)

ray-trace a limited sample of photon trials (selection of  $\phi$  's for given  $\theta)$ 

adjust the angles to match the hit starting from the closest trial (convergence in 2-3 iterations)

validate photon reconstructed Cherenkov angle and transit time







# **Ray Tracing**



Trial position and refraction at boundaries  $[\theta, n(\beta)]$  depend on particle hypothesis Stop when closer than a given fraction of the expected (angular) resolution





# **Ray Tracing**



Photon path reconstruction allow to assign the photon to the most likely hypothesis:

- be robust and easy to control (easy to handle multi-reflections, up to e.g. 5)
- discriminate background (hit far from trials, no solution foreseeable)
- provide full information (photon path, time, position and component of each reflection)
- allow relation with nominal optical components, resolution and efficiency





Run Group C



Example of 3 particle event into two RICHes (no calibration)







Observe *n* events  $x_i$  and expect a total of  $\mu$  based on model *f* and parameters  $\theta$ 

In the *logL* any normalization constant is irrelevant, i.e. any term that does not depend on the parameters  $\theta$ .

In case f( $x_i$ ;  $\theta$ ) is a binned PDF with poisson probability the above reduces (except for an irrelevant normalization constant) to the binned maximum likelihood

$$L f_{\mathrm{P}}(\boldsymbol{n};\boldsymbol{\theta}) = \prod_{i=1}^{N} \frac{\mu_{i}^{n_{i}}}{n_{i}!} e^{-\mu_{i}}$$



# **Binned Log Likelihood**



$$-2\ln\lambda(\boldsymbol{\theta}) = 2\sum_{i=1}^{N} \left[ \mu_i(\boldsymbol{\theta}) - n_i + n_i \ln \frac{n_i}{\mu_i(\boldsymbol{\theta})} \right]$$

Taking a likelihood ratio, vs an ideal model corresponding to the "observed pattern" ( $\mu_i = n_i$ ), provides a chi2-like estimator (goodness of fit) if  $\mu_i$  are not too small

- \* pdf normalization to 1 is given by Poisson
- \* the  $\mu_i n_i$  term is optional (less stringent except close to threshold)
- \* one can have bins with zero counts (last term is taken to be zero)
- \* there is some arbitrary choice in the bin selection (N)
  (i.e. total PMT surface or just the area potentially illuminated by the photons)
- \*  $\mu_i$  is the expected yield in bin i (signal + background)

$$f_{PIXEL}(i;\theta) = 1 - e^{-\mu(i,\theta)} \sim \mu(i,\theta)$$

Poisson probability to have zero (no hit)







All these quantities are defined at pixel level

 $d\theta$  and  $d\phi$  define the pixel solid angle and are known by RICH reconstruction

 $N_0$ ,  $\varepsilon$ ,  $\theta$ , t,  $\sigma$ ,  $\sigma_t$ , B can be extracted from data (control samples) and enter the CCDB database

 $\epsilon$ (i) can reflect dead (0), hot (1), or the quantum/reflection efficiency ([0:1])

B(i) can be derived from random triggers or electron control sample

This definition is effectively pretty close to the simpler ones used in pass1 However it is more general, seems better defined and could accounts for second order effects (photon path and pattern change among various mass hypotheses).





## **RICH Calibration Suites:**

- **1. Dark count measurement** with scaler readout and random triggers
  - 1. extended to 2 modules
  - 2. dedicated data taking
  - 3. extract hot channel list, estimate dark count rate and pixel efficiency

#### 2. Time calibration from CLAS data

- 1. extended to two modules
- 2. input is calibration data (full runs)
- 3. extract time offsets and time walk corrections

### 3. Cherenkov angle calibration from CLAS data

- 1. new software
- 2. input is DST data (high statistics)
- 3. extract measured Cherenkov angle mean and sigma per photon
  - 4. detection topology and particle charge
- 4. Alignment from CLAS data and MC
  - Check response for specific photon paths
  - Account for correlations at once with AI

ε, Β

 $\boldsymbol{t}$  ,  $\boldsymbol{\sigma}_{t}$ 

Ν<sub>0</sub>, θ, σ



# Likelihood Ratio



$$\mu(i,\theta) = \varepsilon(i) \frac{d\phi}{2\pi} e^{-\frac{(\theta-\theta_i)^2}{2\sigma^2}} \frac{d\theta}{\sqrt{2\pi\sigma}} e^{-\frac{(t(\theta)-t_i)^2}{2\sigma_t^2}} \frac{dt}{\sqrt{2\pi\sigma_t}}$$

With negligible B(i), at first order all the terms not depending on  $\theta$  change the likelihood value but not the minimum location (best  $\theta$ , or mass, estimator).

One might take the likelihood ratio with a model corresponding to the "observed pattern" in which all the hit are at the right (expected) angle:

$$\mu^{\prime(i,\theta)} = \varepsilon(i) \frac{d\phi}{2\pi} \qquad 1 \qquad \frac{d\theta}{\sqrt{2\pi}\sigma} \qquad 1 \qquad \frac{dt}{\sqrt{2\pi}\sigma_t}$$
$$\theta = \theta_i \qquad t = t_i$$

$$-2ln\lambda(\theta) = 2\sum_{i=1}^{N} \left[ \mu_i(\theta) - n_i + n_i ln \frac{\mu_i'(\theta)}{\mu_i(\theta)} \right]$$





As expected, the log ratio should reduce to a sort of chi2.

$$-2ln\lambda(\theta) = 2\left(N_{exp} - N\right) + \sum_{i=1}^{N} \left[\frac{(\theta - \theta_i)^2}{\sigma^2}\right] + \left[\frac{(t(\theta) - t_i)^2}{\sigma_t^2}\right]$$

Except for second order effects, the real difference is in the background that defines a sort of cutoff: an accepted hit that is background for all the hypotheses does not count in the likelihood, whereas the ordinary chi2 weights anyway its distance from the expected value (provides a preference)







## Providing best particle hypothesis (PID) with quality estimators

"nam	e": "RICH::Particle",	
"gro	oup": 21800,	
"ite	m" : 37 <i>,</i>	
"inf	o": "Reconstructed Che	renov information per track",
"ent	tries": [	
	{"name":"id", "t	ype":"B", "info":"id"},
	{"name":"hindex",	"type":"S", "info":"related row in the RICH::clusters bank (if any)"},
	{"name":"pindex",	"type":"B", "info":"related row in the REC::Particle bank"},
	{"name":"emilay",	"type":"B", "info":"aerogel layer of photon emission"},
	{"name":"emico",	"type":"B", "info":"aerogel component of photon emission"},
	{"name":"t", "typ	e":"B", "info":"aerogel component of particle entrance point"},
	{"name":"emqua",	"type":"S", "info":"aerogel quadrant of photon emission"},
	{"name":"mchi2",	"type":"F", "info":"track-cluster matching chi2 (if any)"},
	{"name":"best_PID",	"type":"S", "info":"most probable PID choice"},
	{"name":"RQ_prob",	"type":"F", "info":"goodness of hadron choice parameter (1=anambiguos, 0=random)"},
	{"name":"ReQ_prob",	"type":"F", "info":"goodness of elecgtron choice parameter (1=anambiguos, 0=random)"},
	{"name":"el_prob",	"type":"F", "info":"probability to be an electron"},
	{"name":"pi_prob",	"type":"F", "info":"probability to be an pion"},
	{"name":"k_prob",	"type":"F", "info":"probability to be an kaon"},
	{"name":"pr_prob",	"type":"F", "info":"probability to be an proton"},
	{"name":"best_etaC",	"type":"F", "info":"Average etaC for best hypothesis"},
	{"name":"best_c2",	"type":"F", "info":"chi2 for best hypothesis"},
	{"name":"best_RL",	"type":"F", "info":"Likelihood ratio for best hypothesis"},
	{"name":"best_ntot",	"type":"F", "info":"Number of photon used for likelihood"},
	{"name":"best_mass",	"type":"F", "info":"Reconstructed mass for best hypothesis"}
]		



**RICH PID** 



- Pass 1 inbending data
- Particles selection based on the EB
- One good trigger electron
- Only 1 charged track in the RICH
- RICH with direct photons only
  - alignment to be completed
- > Average Cherenkov angle of the track
  - 4 p.e. minimum



## NOTE:

- No fiducial cuts
- No PID refinements (chi2pid, calo SF, etc.)
- No kinematical cuts (SIDIS)

"Natural" binning for the RICH analysis:

- tile number (variation in the refractive index)
- particle momentum (beta dependence of Cherenkov angle)

### M. Mirazita pass1 analysis



## **Pion vs Electron**





Too low momenta for SIDIS kinematics

#### M. Mirazita pass1 analysis



0.4

0.38

0.36

0.34

0.32

0.3

0.28

0.26

0.24

0.22

0.2<sup>C</sup>

1. Select positive hadrons with EB

## **Positive Hadrons**



taC\_I1\_t12\_p2

Std Dev 0.00999

taC\_11\_t12\_p2

td Dev 0.0134

4349 0.303

0.3043

Layer 1

Tile 12

2. Look at the RICH Cherenkov angle row 1 Track  $\theta_{\rm C}$ , Layer 1, tile 12 , P=3.70 Track  $\theta_{C}$ , Layer 1, tile 12 , P=3.30 Track  $\theta_{C}$ , Layer 1, tile 12 , P=3.50 hEtaC I1 t12 1200 1000 800 p=3.30 100565 Entries Mean x 4.009 Mean y 0.304 0.9666 Std Dev x Std Dev y 0.009572 10<sup>2</sup> Track  $\theta_{\rm C}^{}$ , Layer 1, tile 12 , P=3.90 Track  $\theta_{\rm C}$ , Layer 1, tile 12 , P=4.10 Track θ<sub>c</sub>, Layer 1, tile 12 , P=4.30 taC\_I1\_t12\_p aC\_I1\_t12\_p2 6626 ntries vlean 0.3048 Std Dev 0.00725 700 0.3045 0.00772 10 Track  $\boldsymbol{\theta}_{C}^{},$  Layer 1, tile 12 , P=4.70 Track  $\theta_{C}$ , Layer 1, tile 12 , P=4.90 Track  $\boldsymbol{\theta}_{C},$  Layer 1, tile 12 , P=4.50 aC\_I1\_t12\_p23 taC\_11\_t12\_pt 700 600 500 4922 0.3027 700 600 500 400 0.303 400 td Dev 0.01327 td Dev 0.01417 p=4.90 2 3 8 9 10 5 6 300 200 100  $\succ$ **Extension to higher momenta** requires better alignment (i.e.

M. Mirazita pass1 analysis

more photons)



## Kaon Yield





M. Mirazita pass1 analysis





# pi+/k+ misidentification







#### M. Mirazita pass1 analysis





Alignment requires large statistics and full reconstruction to deal with the various photon paths Angular resolution is comparable for direct and reflected photons

#### Direct photons: electrons vs pi+

Photon 0, (B=1) vs Momentum, Layer 0, Tile 13, direct photons

#### **Electrons: direct vs planar reflection**

Layer 0 Tile 5

**Electrons: direct vs spherical reflection** 









Contalbrigo M.



## Pass2 Alignment





Contalbrigo M.

## **RICH Pass2 Reconstruction**

Pass1: direct photons and single reflection

Pass2: full acceptance (multiple reflections), RICH PID and 2<sup>nd</sup> module

multi-thread safe, all particle ID hypothesis, multi-particles likelihood PID for single photon and single particle complete photon path tagging for proper calibration and alignment



# Secondary ID



Clear pion (electron). Equally bad kaon and proton, but LH with N term prefers smaller expected yield.

Contalbrigo M.

## Pion vs Kaon







M. Mirazita analysis