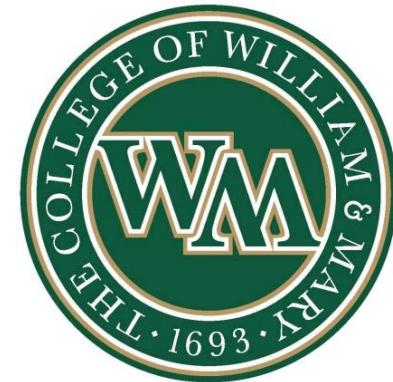


Qweak Ancillary Measurement: Parity-Violating Inelastic $\vec{e}p$ Asymmetry at 3.35 GeV

James Dowd

For the Qweak Collaboration
Jan 15-16, 2015



Ancillary Measurement

- For ~2 weeks, Qweak received polarized electron beam at higher than nominal beam energy
- Opportunity to make an ancillary measurement
- Kinematics tuned to the inelastic region
 - Conducive to probing $F_{1,3}^{\gamma Z}$
- We quickly developed a detector system

Kinematics

$$E_{Beam} = \cancel{1.165 \text{ GeV}}$$

$$Q^2 \approx \cancel{0.025 \text{ GeV}^2}$$

$$\theta \approx 7 - 11^\circ$$

$$\text{Current} = 180 \mu\text{A}$$

$$\text{Polarization} = 89\%$$

$$\text{Target} = 35 \text{ cm LH}_2$$

$$\text{Cryopower} = 3.0 \text{ kW}$$

$$3.35 \text{ GeV}$$

$$0.09 \text{ GeV}^2$$

$$W \approx 2.23 \text{ GeV}$$

$$0.89 * 93\% \text{ Longitudinal}$$

$$0.89 * 37\% \text{ Transverse}$$

Horizontal Drift Chambers

Electron Beam

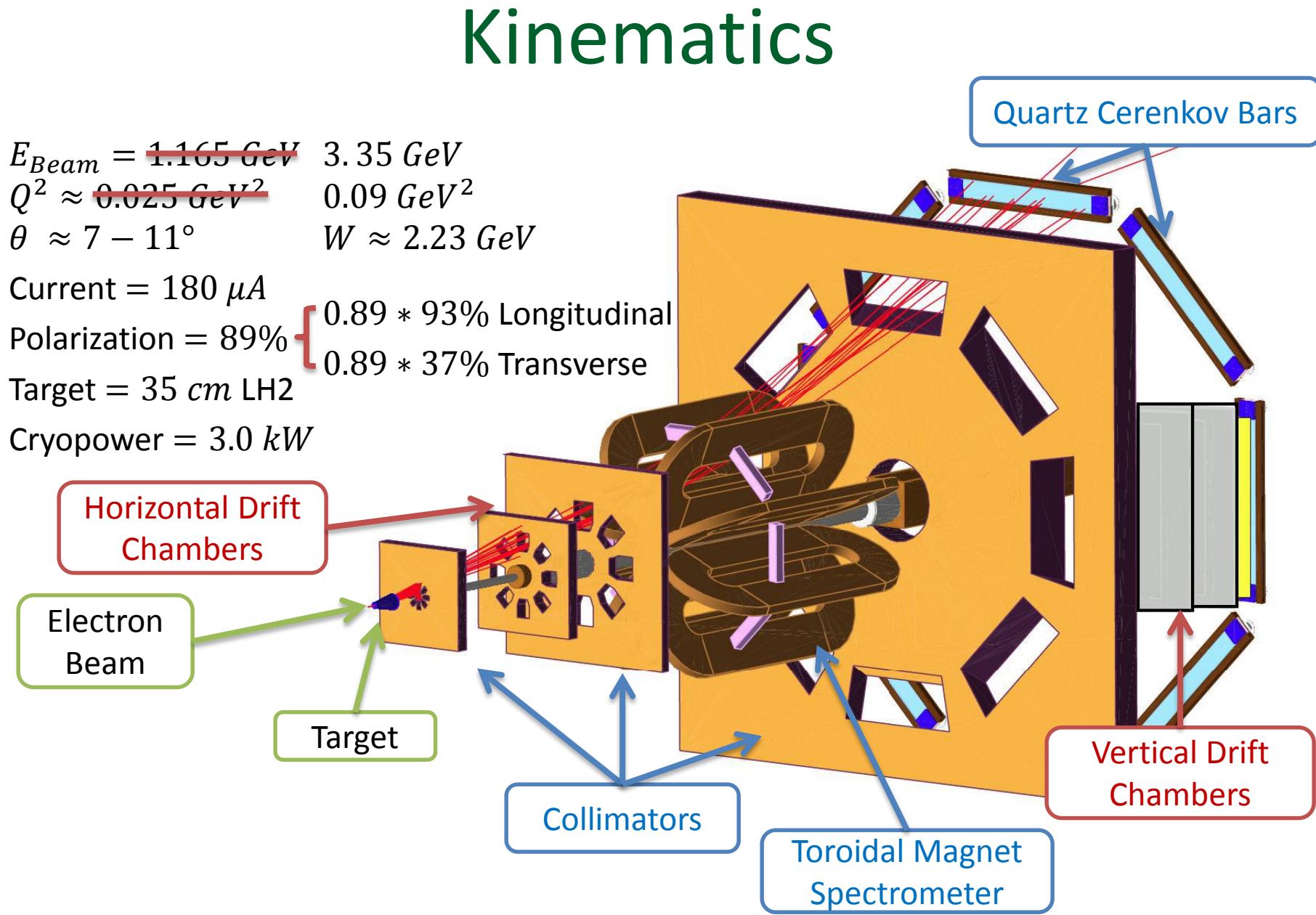
Target

Collimators

Toroidal Magnet Spectrometer

Quartz Cerenkov Bars

Vertical Drift Chambers



Radiative Corrections

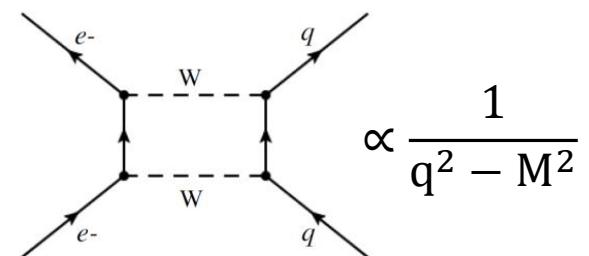
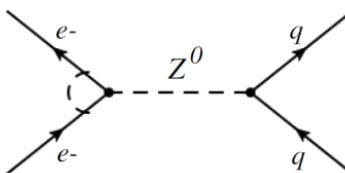
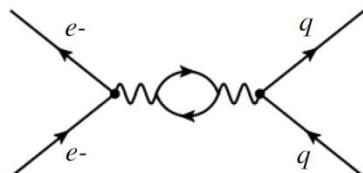
Including 1st order corrections, the proton's weak charge is:

$$Q_W^p = [(\rho_{NC} + \Delta_e)(1 - 4 \sin^2 \theta_W(0)_{\overline{MS}} + \Delta'_e) + \Re e \square_{WW} + \Re e \square_{ZZ} + \Re e \square_{\gamma Z}] \Big|_{E=0, Q^2=0}$$

W and Z mass renormalization

Vertex corrections

Dominated by large momentum transfer & can be calculated by perturbative QCD



γZ -Box

Including 1st order corrections, the proton's weak charge is:

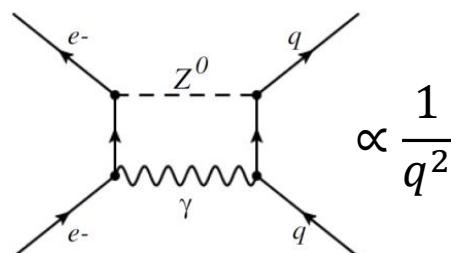
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W and Z mass renormalization

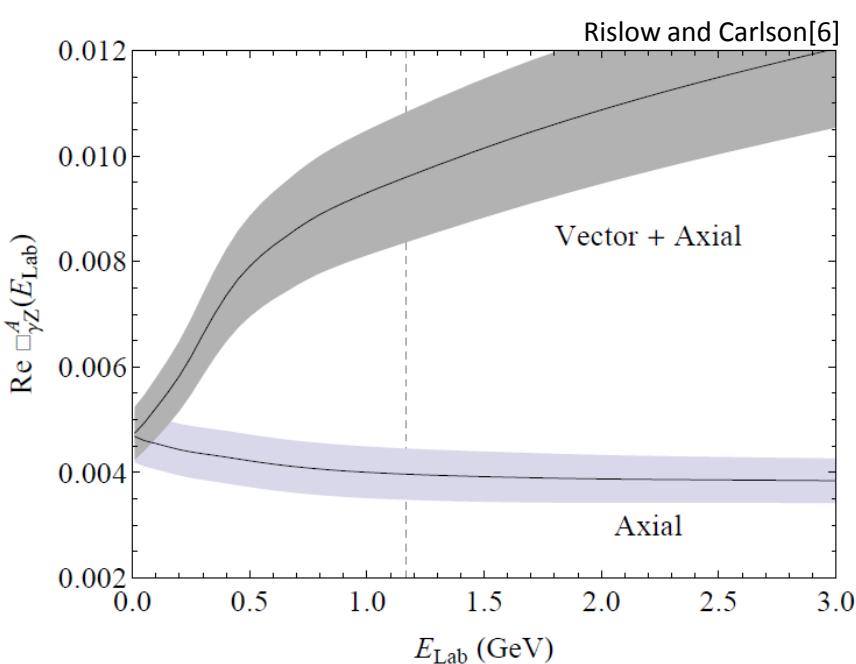
Vertex corrections

Dominated by large momentum transfer & can be calculated by perturbative QCD

Less well known



γZ -Box



$$\Re \square_{\gamma Z}^V \times 10^3 \text{ evaluated at } E = 1.165 \text{ GeV.}$$

Oldest

Sibirtsev *et al.* [1]

$4.7^{+1.1}_{-0.4}$

Rislow and Carlson [2]

5.7 ± 0.9

Gorchtein *et al.* [3]

5.4 ± 2.0

Hall *et al.* [4]

5.60 ± 0.36

Newest

$$\Re \square_{\gamma Z}^A \times 10^3 \text{ evaluated at } E = 1.165 \text{ GeV.}$$

Oldest

Blunden *et al.* [5]

3.7 ± 0.4

Rislow and Carlson [6]

4.0 ± 0.5

Newest

*See reference slide

- Gorchtein and Horowitz showed γZ -box contribution
 - Energy dependent
 - Larger than originally expected ($\sim 8\%$ of Q_W^p)
 - Uncertainty of correction could affect precision aim of Qweak
- Recent thorough analysis of both axial and vector components

$F_{1,2,3}^{\gamma Z}$ Structure Functions

For forward scattering, the dispersion relation is:

$$\Re \square_{\gamma Z}^V(E) = \frac{2E}{\pi} \mathcal{P} \int_0^\infty dE' \frac{1}{E'^2 - E^2} \Im \square_{\gamma Z}^V(E')$$

Using the optical theorem,

$$2\Im \mathcal{M}_{\gamma Z}^{(PV)} = -4\sqrt{2}\pi MG_F \int \frac{d^3 k'}{(2\pi)^3 2E_{k'}} \left(\frac{4\pi\alpha}{Q^2}\right) \frac{1}{1 + Q^2/M_Z^2} L_{\mu\nu}^{\gamma Z} W_{\gamma Z}^{\mu\nu}$$

where

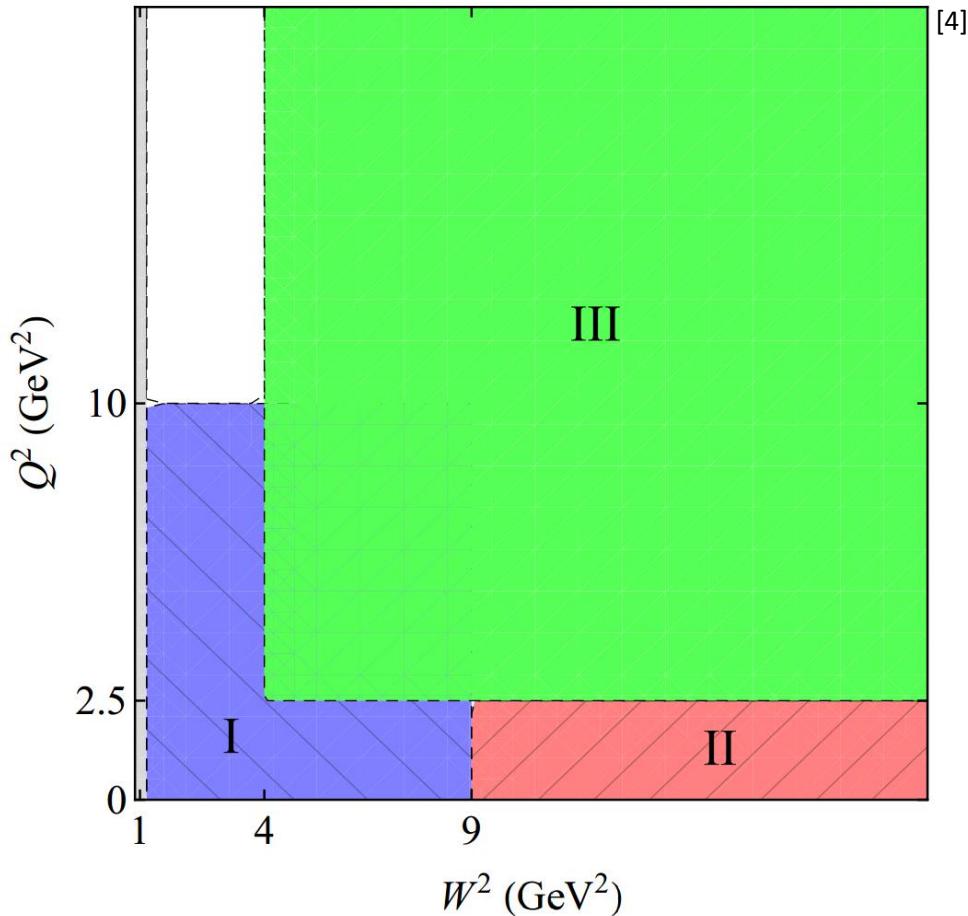
$$\begin{aligned} L_{\mu\nu}^{\gamma Z} &= \bar{u}(k, \lambda)(g_V^e \gamma_\mu - g_A^e \gamma_\mu \gamma_5) \not{k}' \gamma_\nu u(k, \lambda) \\ MW_{\gamma Z}^{\mu\nu} &= -g^{\mu\nu} F_1^{\gamma Z} + \frac{p^\mu p^\nu}{p \cdot q} F_2^{\gamma Z} - i\epsilon^{\mu\nu\lambda\rho} \frac{p_\lambda p_\rho}{2p \cdot q} F_3^{\gamma Z} \end{aligned}$$

Combining everything, the imaginary part of the correction becomes

$$\Im \square_{\gamma Z}^V(E') = \frac{1}{(s - M^2)^2} \int_{W_\pi^2}^s dW^2 \int_0^{Q_{max}^2} dQ^2 \frac{\alpha(Q^2)}{1 + Q^2/M_Z^2} \left[F_1^{\gamma Z} + \frac{s(Q_{max}^2 - Q^2)}{Q^2(W^2 - M^2 + Q^2)} F_2^{\gamma Z} \right]$$

With $s = M^2 + 2ME$, $W_\pi^2 = (M + m_\pi)^2$, and $Q_{max}^2 = 2ME(1 - W^2/s)$

$F_{1,2,3}^{\gamma Z}$ Structure Functions



Integral divided into several regions:

- Region I – Described by the Christy-Bosted $F_{1,2}^{\gamma\gamma}$ fit, transformed to the γZ case
- Region II – GHRM Model II, transformed to γZ
- Region III – Global PDF fits to high-energy data

$$\Im m \square_{\gamma Z}^V(E') = \frac{1}{(s - M^2)^2} \int_{W_\pi^2}^s dW^2 \int_0^{Q_{max}^2} dQ^2 \frac{\alpha(Q^2)}{1 + Q^2/M_Z^2} \left[F_1^{\gamma Z} + \frac{s(Q_{max}^2 - Q^2)}{Q^2(W^2 - M^2 + Q^2)} F_2^{\gamma Z} \right]$$

Non-Resonant inelastic PV measurement

Beam Energy = 3.35 GeV

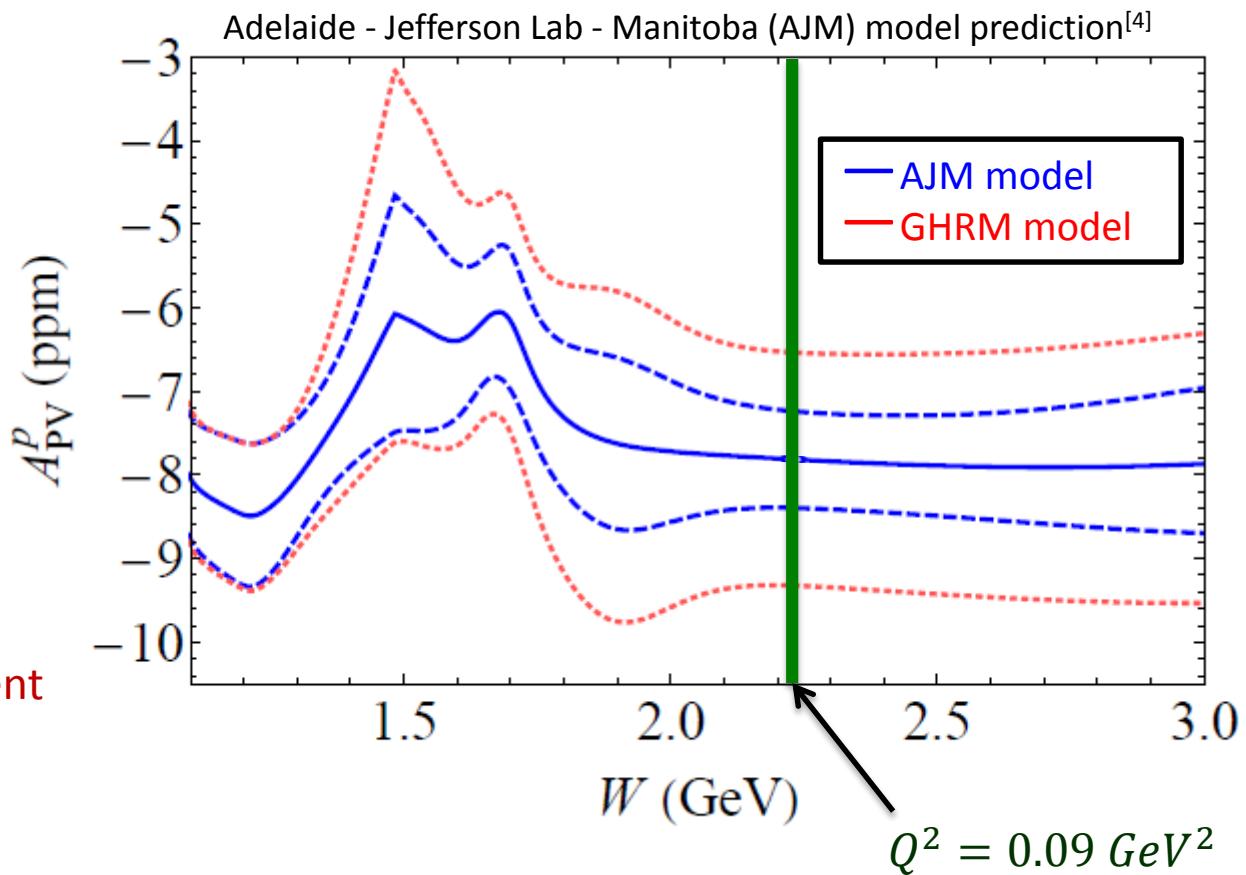
$W = 2.23 \text{ GeV}$

$Q^2 = 0.09 \text{ GeV}^2$

Predicted

$A_{PV}^p \sim (-7.8 \pm 0.6) \text{ ppm}$

An asymmetry measurement
will constrain $F_{1,3}^{\gamma Z}$



- Asymmetry measurement lies in the non-resonant inelastic region
- In a kinematic region with almost no experimental data

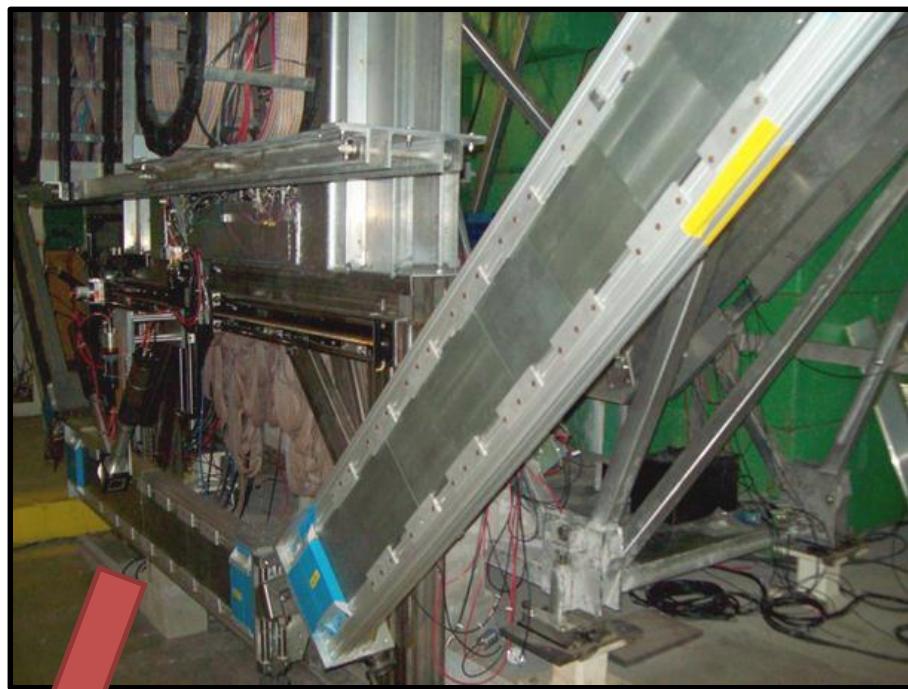
[4] N. L. Hall, P. G. Blunden, W. Melnitchouk, A. W. Thomas, and R. D. Young (2013), arXiv:1304.7877v1 [nucl-th].

Data Analysis

- Large W leads to large pion background
 - Pion and electron signals are integrated together
 - Pion dilution is the largest systematic uncertainty
- Elastic ep radiative tail
 - Well understood
- Partially transversely polarized beam
 - ~37% transverse
 - Pions have a large transverse asymmetry

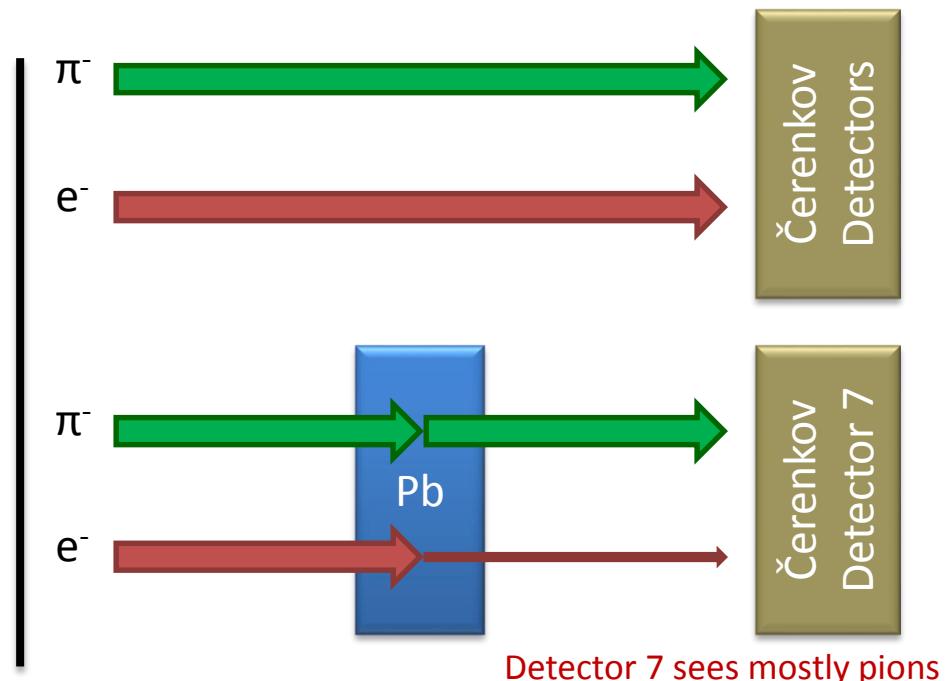
Signal Extraction

- At 3.35 GeV, signals in the Čerenkov detectors will be a combination of π^- & e^-
- Majority of data was taken with 4" lead in front of lowest Čerenkov Detector
- The lead wall ranges out most of the electrons.
 - ~18 radiation lengths



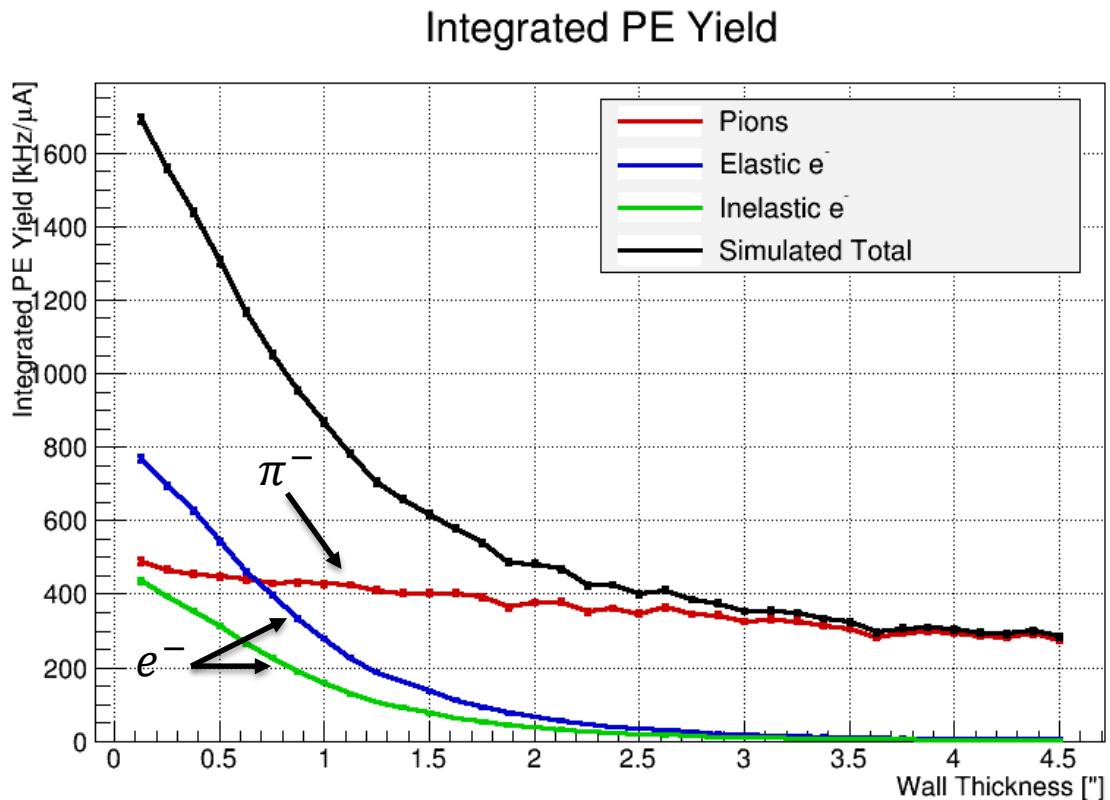
Signal Extraction

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Lead Wall Simulation

- GEANT4 Simulation
- Electron signal is highly suppressed with increasing wall thickness
- Lead wall isolates pions
- Blocked Čerenkov detector becomes a pion detector



ADC Light Spectra

- Some event mode data taken
 - Each event is tracked and counted separately
- Pion rate greater than electron rate by factor of ~ 2
- Pions produce less light than electrons by a factor of ~ 6
 - Electron signal enhanced by lead pre-radiators
- Fitting spectra will help us determine the pion fraction

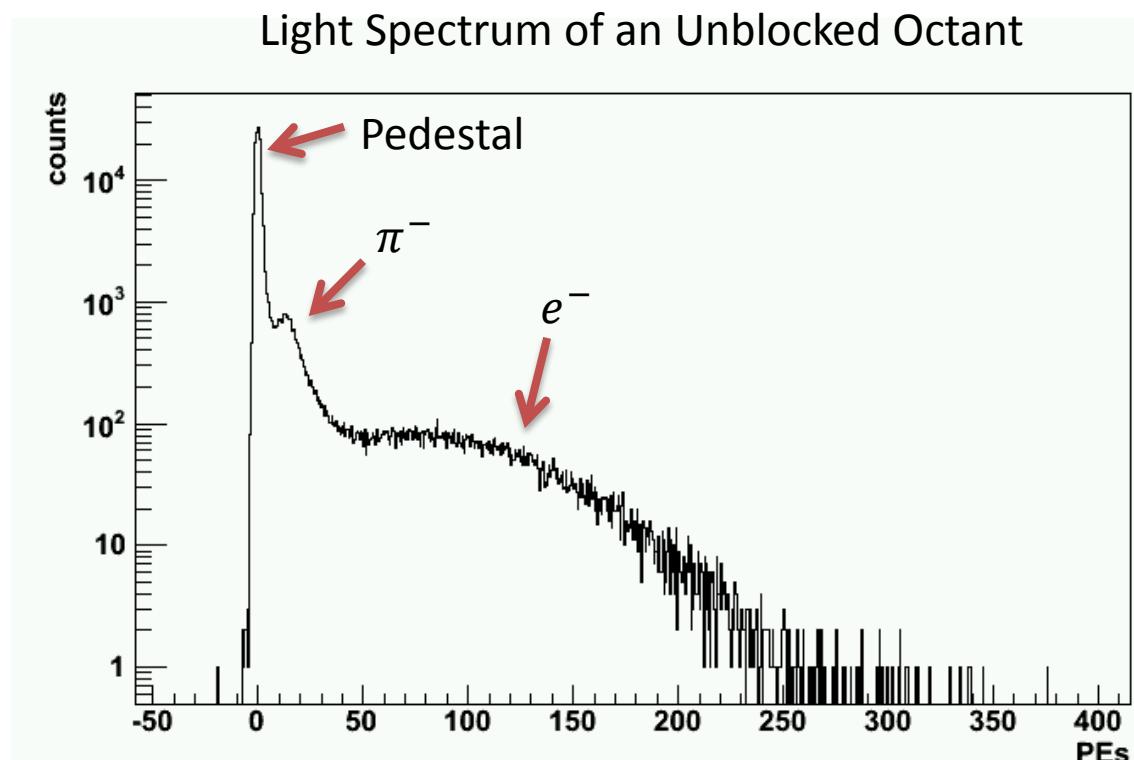
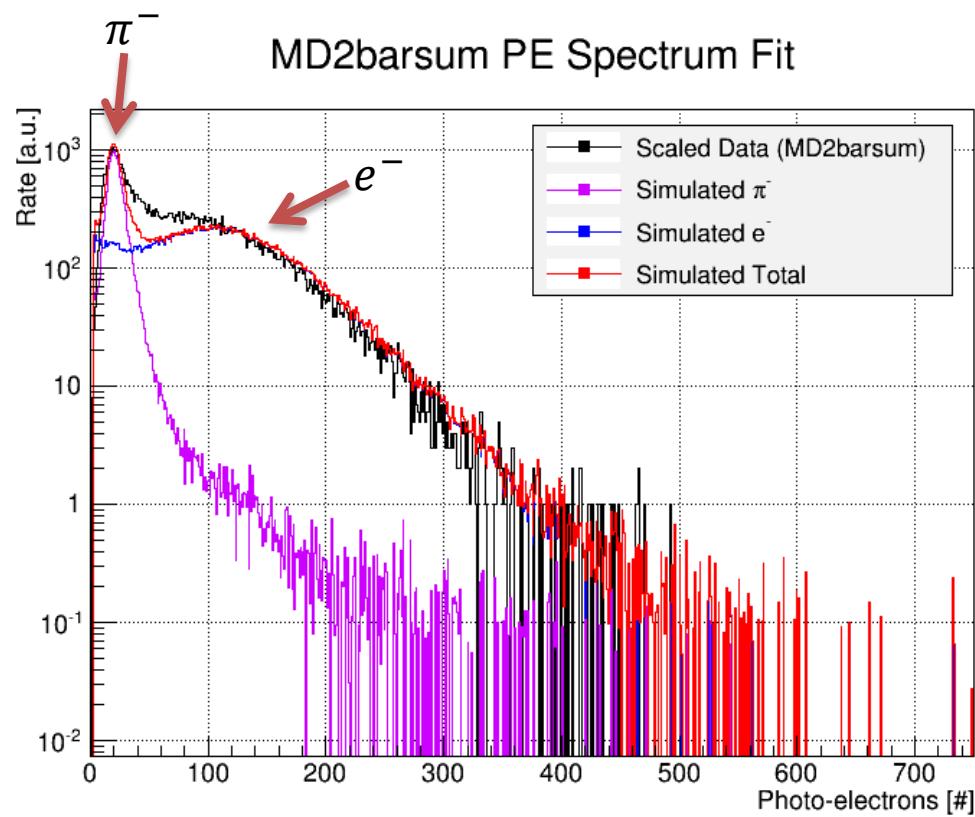


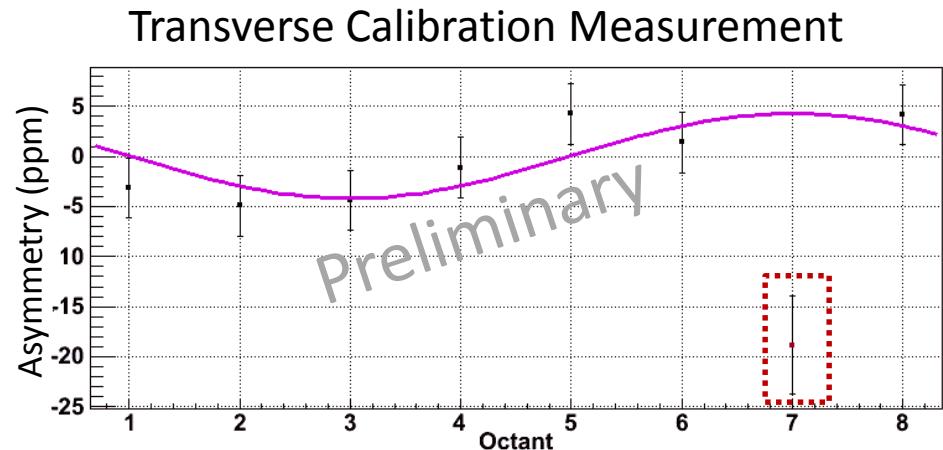
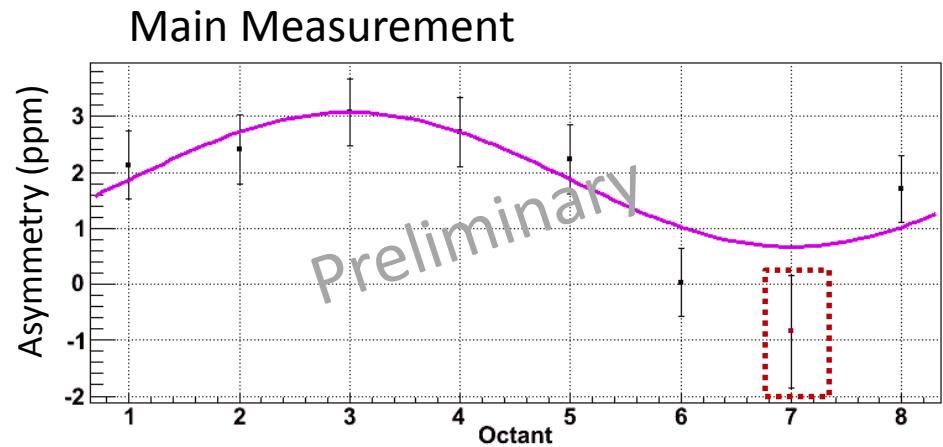
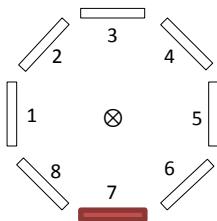
Photo-Electron Spectrum Fit

- Fit data with GEANT4 pion and electron simulations
- Removed electronic pedestal
- Applied a 3 parameter fit
 - Pion Amplitude
 - Electron Amplitude
 - Gain Factor [PEs/ADC channel]
- Simulation under-predicts between electron and pion peaks



Beam Polarization at 3.35 GeV

- Data shown
 - Uncorrected
 - Blinded
- Main measurement beam polarization
 - $\sim 0.93 * P_{Beam}$ Longitudinal
 - $\sim 0.37 * P_{Beam}$ Transverse
- Azimuthally symmetric detectors
- Pure transverse measurement beam polarization
 - $\sim 1.0 * P_{Beam}$ Transverse
- Note the large transverse pion asymmetry with opposite sign



* Octant 7 with lead wall excluded from fit

Extracting A_{PV}^e

Parity-Violating Pion Asymmetry

$$A_{mix}^{MD7} = P_T A_T^{MD7} + P_L A_{PV}^{MD7}$$

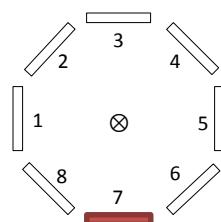
$$\Rightarrow A_{PV}^{MD7} = \frac{1}{P_L} (A_{mix}^{MD7} - P_T A_T^{MD7}) \approx A_{PV}^\pi$$

Parity-Violating Electron Asymmetry

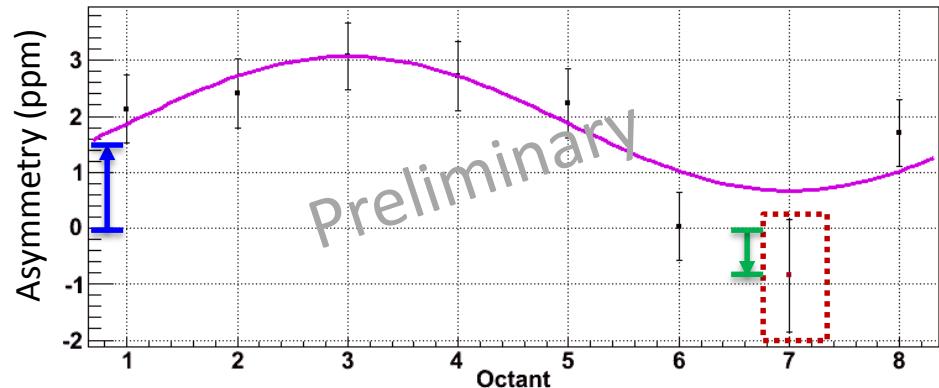
$$A_{PV} = (1 - f_\pi) A_{PV}^e + f_\pi A_{PV}^\pi$$

$$\Rightarrow A_{PV}^e = \frac{A_{PV} - f_\pi A_{PV}^\pi}{1 - f_\pi}$$

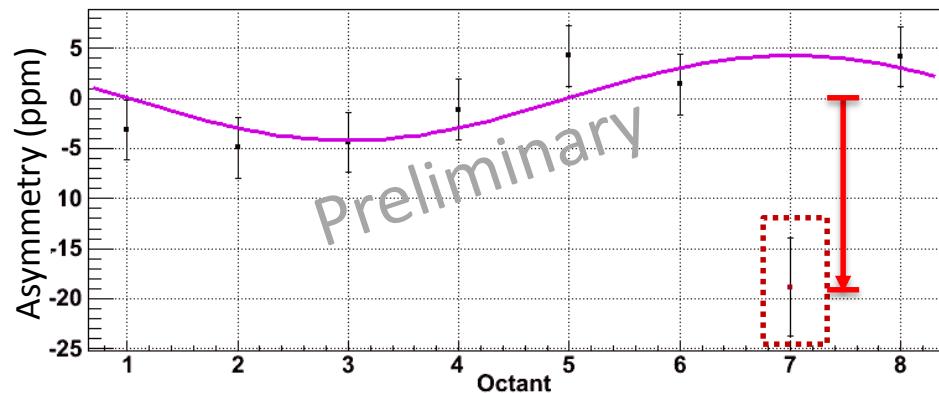
Asymmetry result is highly dependent on f_π



Main Measurement



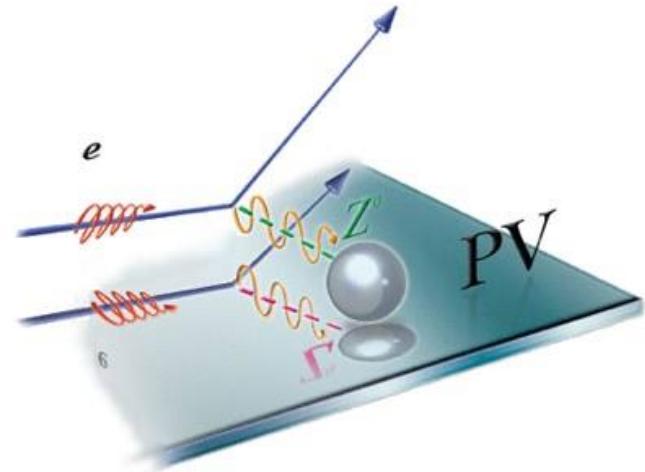
Transverse Calibration Measurement



* Octant 7 with lead wall excluded from fit

Summary

- Pion fraction will be determined by
 - Čerenkov detector with the lead wall
 - Light spectra in the PMTs
 - Monte Carlo simulation
- Asymmetry measurement at 3.35 GeV
 - Constrain γZ structure functions, $F_{1,3}^{\gamma Z}$
 - Part of a small unique dataset
 - Relevant for Qweak and upcoming experiments
- Additional measurements that we get for ‘free’
 - Non-resonant inelastic transverse asymmetries
 - PV asymmetries in pion photoproduction at 3.35 GeV
 - Transverse asymmetries in pion photoproduction at 3.35 GeV



The Qweak Collaboration



95 collaborators 23 grad students
10 post docs 23 institutions

Institutions:

- ¹ University of Zagreb
- ² College of William and Mary
- ³ A. I. Alikhanyan National Science Laboratory
- ⁴ Massachusetts Institute of Technology
- ⁵ Thomas Jefferson National Accelerator Facility
- ⁶ Ohio University
- ⁷ Christopher Newport University
- ⁸ University of Manitoba,
- ⁹ University of Virginia
- ¹⁰ TRIUMF
- ¹¹ Hampton University
- ¹² Mississippi State University
- ¹³ Virginia Polytechnic Institute & State Univ
- ¹⁴ Southern University at New Orleans
- ¹⁵ Idaho State University
- ¹⁶ Louisiana Tech University
- ¹⁷ University of Connecticut
- ¹⁸ University of Northern British Columbia
- ¹⁹ University of Winnipeg
- ²⁰ George Washington University
- ²¹ University of New Hampshire
- ²² Hendrix College, Conway
- ²³ University of Adelaide

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References

- [1] A. Sibirtsev, P. G. Blunden, W. Melnitchouk, and A. W. Thomas, Phys. Rev. **D82**, 013011 (2010), 1002.0740.
- [2] B. C. Rislow and C. E. Carlson, Phys.Rev. **D83**, 113007 (2011), 1011.2397.
- [3] M. Gorchtein, C. Horowitz, and M. J. Ramsey-Musolf, Phys.Rev. **C84**, 015502 (2011), 1102.3910.
- [4] N. L. Hall, P. G. Blunden, W. Melnitchouk, A. W. Thomas, and R. D. Young (2013), arXiv:1304.7877v1 [nucl-th].
- [5] P. Blunden, W. Melnitchouk, and A. Thomas, Phys.Rev.Lett. **107**, 081801 (2011), 1102.5334.
- [6] B. C. Rislow and C. E. Carlson, arXiv:1304.8113v1 [hep-ph].

Backup Slides

Extracting A_{PV}^e

Parity-Violating Pion Asymmetry

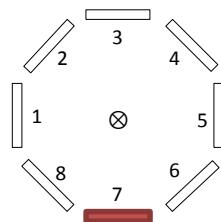
$$A^{MD7} = P_T A_T^{MD7} + P_L A_{PV}^{MD7}$$

$$\Rightarrow A_{PV}^{MD7} = \frac{1}{P_L} (A^{MD7} - P_T A_T^{MD7}) \approx A_{PV}^\pi$$

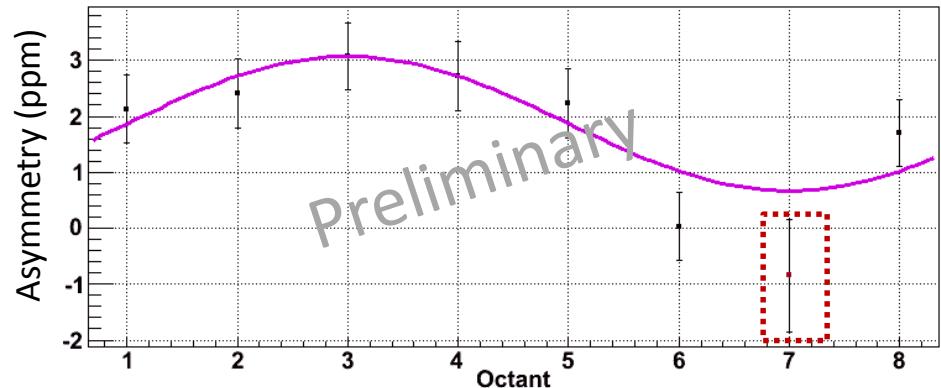
Parity-Violating Electron Asymmetry

$$A_{PV} = (1 - f_\pi) A_{PV}^e + f_\pi A_{PV}^\pi$$

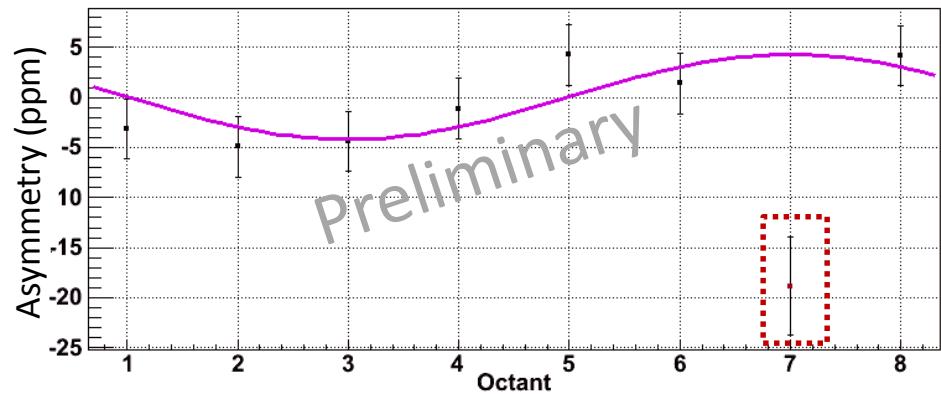
$$\Rightarrow A_{PV}^e = \frac{A_{PV} - f_\pi A_{PV}^\pi}{1 - f_\pi}$$



Main Measurement



Transverse Calibration Measurement



* Octant 7 with lead wall excluded from fit

Non-Resonant inelastic PV measurement

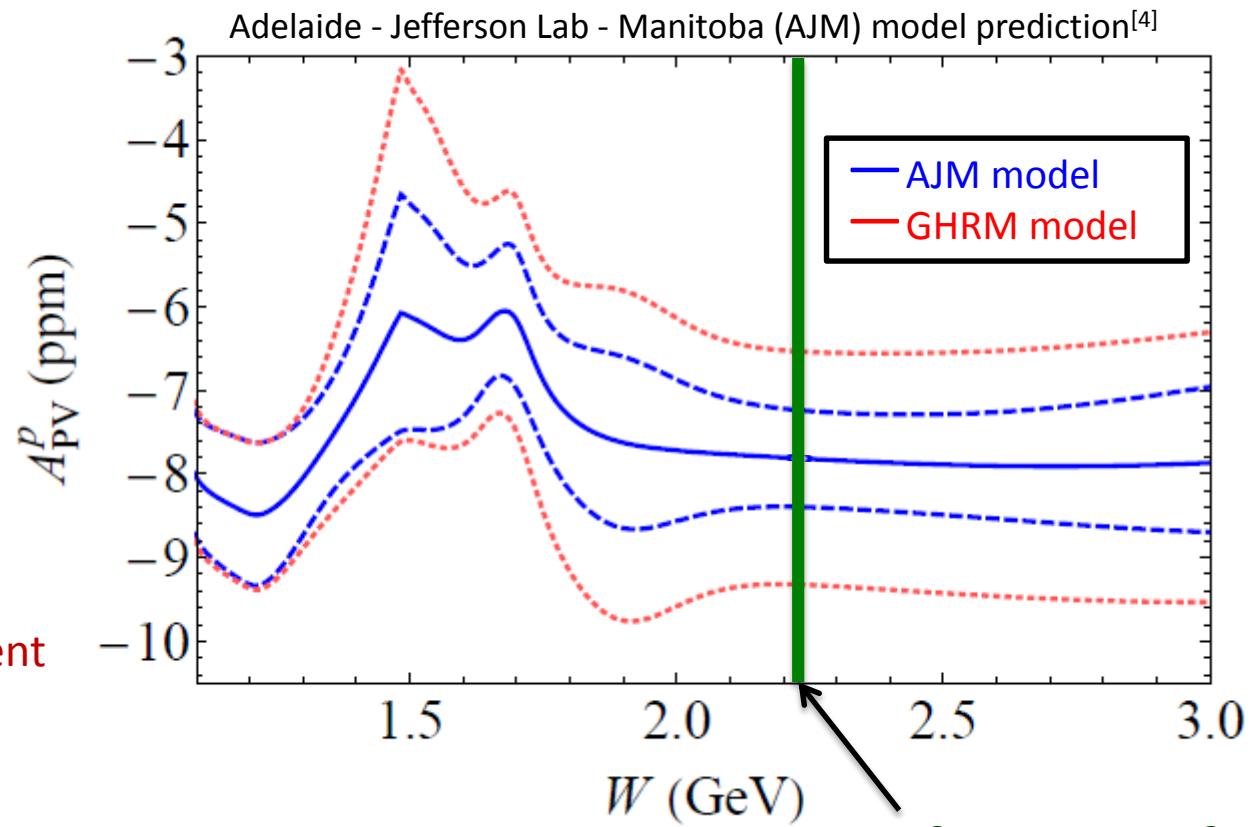
Beam Energy = 3.35 GeV

$W = 2.23 \text{ GeV}$

$Q^2 = 0.09 \text{ GeV}^2$

Predicted

$A_{PV}^p \sim (-7.8 \pm 0.6) \text{ ppm}$



An asymmetry measurement

will constrain $F_{1,3}^{\gamma Z}$

- Asymmetry measurement lies in the non-resonant inelastic region $Q^2 = 0.09 \text{ GeV}^2$

$$A_{PV} = g_A^e \left(\frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \right) \frac{xy^2 F_1^{\gamma Z} + \left(1 - y - \frac{x^2 y^2 M^2}{Q^2} \right) F_2^{\gamma Z} + \frac{g_V^e}{g_A^e} \left(y - \frac{1}{2} y^2 \right) x F_3^{\gamma Z}}{xy^2 F_1^{\gamma\gamma} + \left(1 - y - \frac{x^2 y^2 M^2}{Q^2} \right) F_2^{\gamma\gamma}}$$

[4] N. L. Hall, P. G. Blunden, W. Melnitchouk, A. W. Thomas, and R. D. Young (2013), arXiv:1304.7877v1 [nucl-th].

Noise Correction w/ TDC Cut

Run 17956 – MD triggered

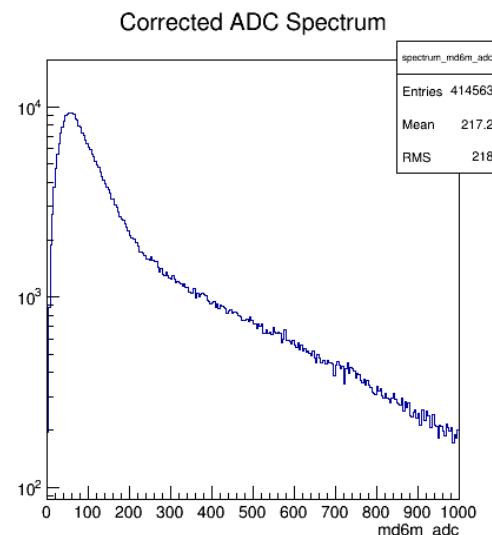
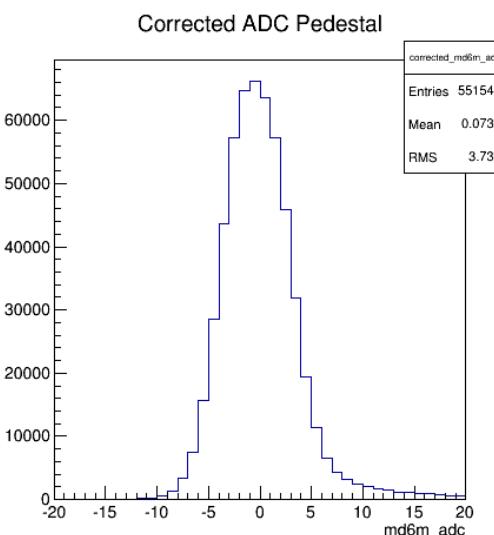
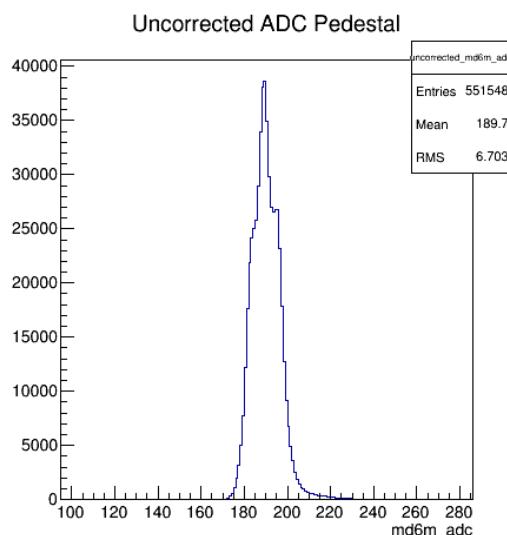
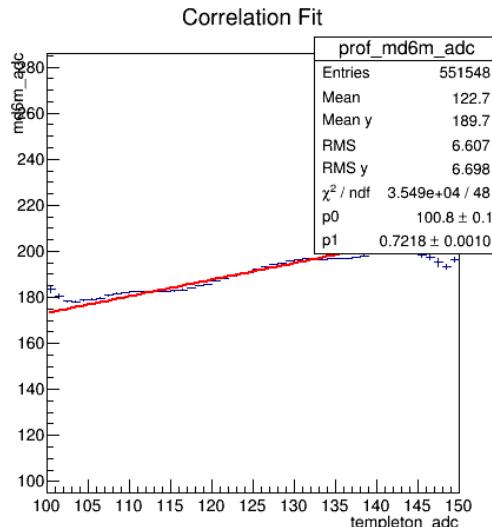
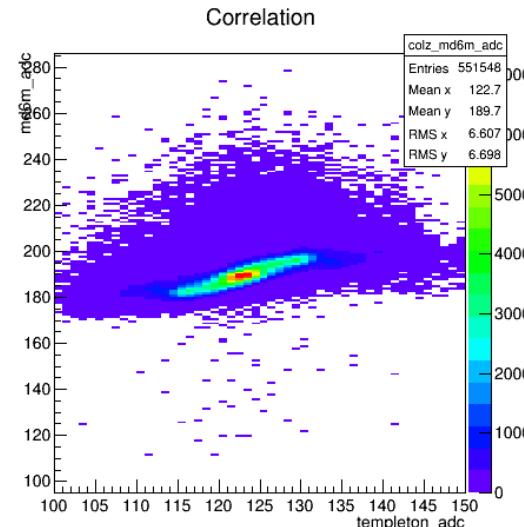
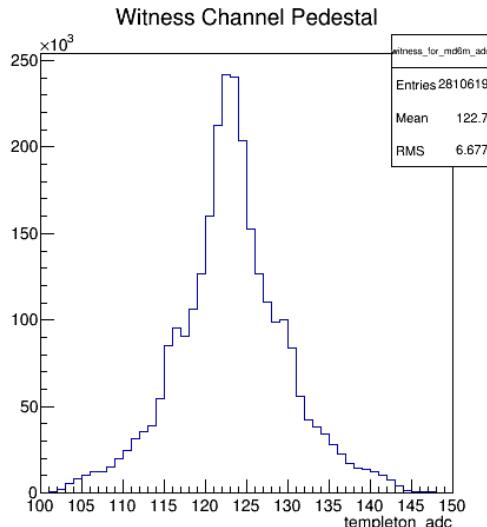
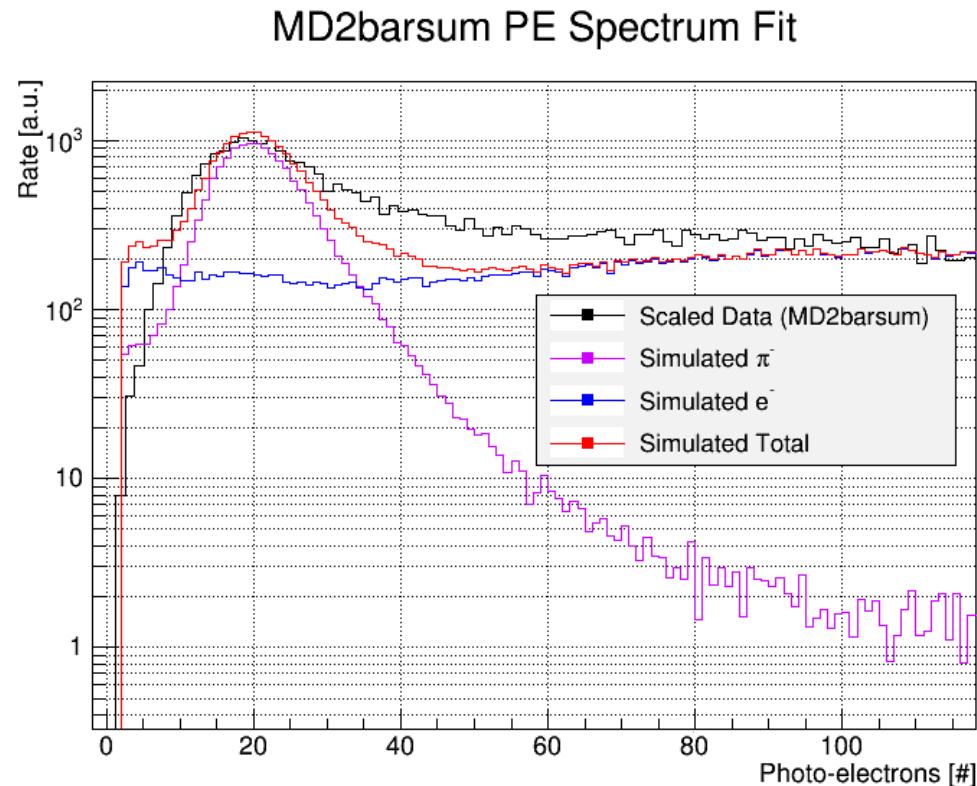


Photo-Electron Spectrum Fit

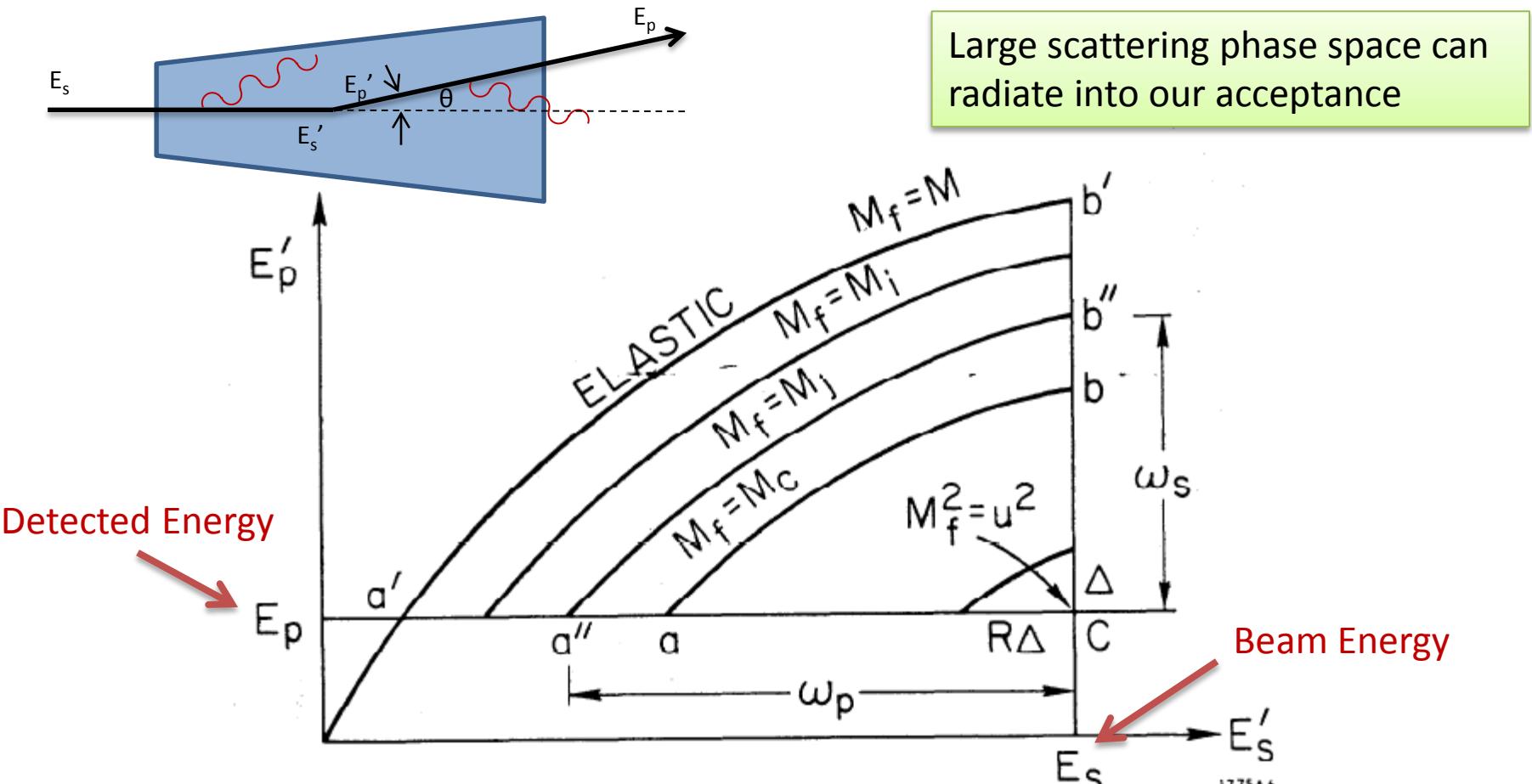
- Fit data with pion and electron simulations
- Removed electronic pedestal
- Applied a 3 parameter fit
 - Pion Amplitude
 - Electron Amplitude
 - Gain Factor [PEs/ADC channel]
- Fit matches well
- Simulation under-predicts between electron and pion peaks



Available Data Set

- Asymmetry and Yields
 - Main Measurement
 - Purely Transverse
 - Aluminum and Carbon Targets
- ADC Spectra
- Tracking Data
- One Čerenkov detector with Lead Wall
 - 4" wall
 - 2" wall
 - Nominal Energy (1.16 GeV)
- Magnet spectrometer current scans
 - LH2
 - Aluminum
 - Reverse Polarity
- Blocked Octant

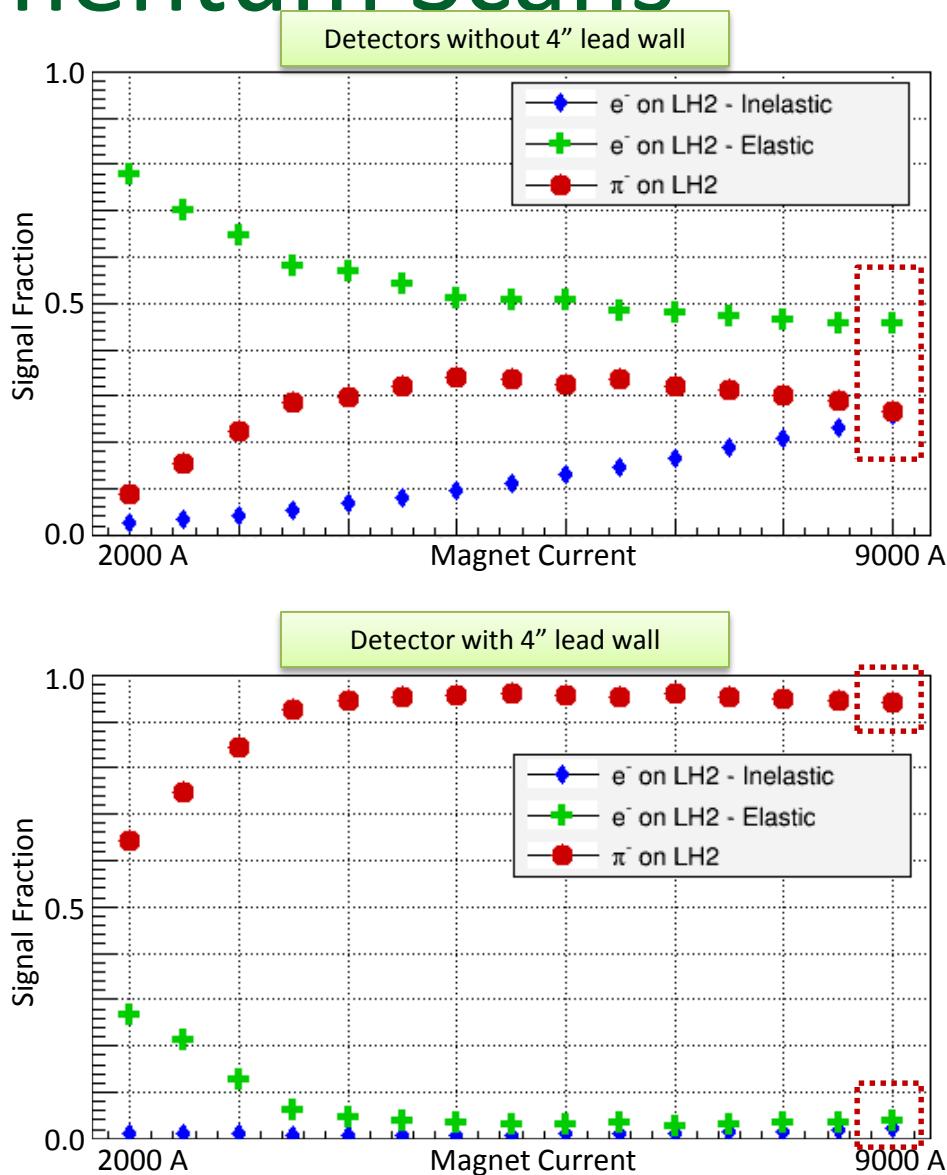
Radiative Corrections



$$\sigma_{exp}(E_s, E_p) = \int_0^T \frac{dt}{T} \int_{E_{s,min}(E_p)}^{E_s} dE'_s \int_{E_p}^{E_{p,max}(E'_s)} dE'_p I(E_s, E'_s, t) \sigma_r(E'_s, E'_p) I(E'_p, E_p, T - t)$$

Simulated Momentum Scans

- Magnet current selects outgoing momentum
- Simulated signal fractions at various magnet currents
 - At 9000 A inelastic signal is maximized at ~25% of total signal
 - At 9000 A elastic signal is minimized
- Most data taken at 9000 A magnet current
- Wall attenuates most electrons
 - ~95% signal comes from pions



Proton's Weak Charge

At tree level, the proton's weak charge is:

$$Q_W^p = 1 - 4 \sin^2 \theta_W = -2(2C_{1u} + C_{1d})$$

The Qweak experiment goal is to measure Q_W^p to $\sim 4\%$

Qweak Apparatus

$E_{Beam} = 1.165 \text{ GeV}$

$Q^2 \approx 0.025 \text{ GeV}^2$

$\theta \approx 7 - 11^\circ$

Current = $180 \mu\text{A}$

Polarization = 89%

Target = 35 cm LH2

Cryopower = 3.0 kW

