# Spin Asymmetries on the Nucleon Experiment HMS results 

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## Outline

- Introduction to SANE-HMS
- Experimental Setup and SANE-HMS
- Dilution Factor
- Radiative Corrections
- HMS Asymmetries
- Spin Structure Functions
- $\mathrm{d}_{2}$, Twist-3 Matrix Element
- Summary


## Purpose of SANE-HMS

SANE-HMS resonance data explores high Bjorken $x$ region at intermediate $Q^{2}$ :

- Resonances and $Q^{2}$ dependence of $A_{1}$ and $A_{2}$
- SSF $g_{2}\left(x, Q^{2}\right)$
- Higher twist effects
- Twist-3 $d_{2}$ matrix element


## Experimental Setup

Polarized Electron Beam: 4.7, 5.9 GeV
Polarized Proton Target: ~ $\perp$, ||


Electron Arm

Ammonia ( $\mathbf{N H}_{3}$ ) Polarized via DNP in 5T Magnetic Field

## HMS Coverage for SANE



## from Packing Fraction

The target and beam are not completely polarized. It contains also un-polarizable materials.

$$
A=\frac{1}{P_{b} P_{t} f} \frac{d \sigma^{\downarrow \uparrow}-d \sigma^{\uparrow \uparrow}}{d \sigma^{\downarrow \uparrow}+d \sigma^{\uparrow \uparrow}}
$$

Dilution factor $f$ is the ratio of free polarizable nucleons to the total amount of nucleons in the sample.
$\mathrm{A}=\mathrm{A}_{180}$ or $\mathrm{A}_{80}$ is the measured asymmetry without radiative corrections. ammonia to the target cell, or the fraction of the cell's length that would be filled with ammonia by cylindrical symmetry.

$$
f=\frac{3 H}{(3 H+N) p f+H e(1-p f)+\text { Others }}
$$



## from Packing Fraction

Total yield has linear relation with packing fraction:
$Y_{T}=m p f+b$
Using MC (P. E. Bosted and M. E. Christy, PRC81 (2010) 055213) assuming two different $p f$, the slope( $m$ ) and intercept(b) can be calculated and then the yield of real data produces pf of real target.

SANE packing fractions are $56 \%-62 \%$ with $\sim 4.5 \%$ error.

## Packing Fraction



Data and MC comparison (Red is MC)

## Packing fraction



■ Unique PF(\%) * Run average

## Dilution Factor

Dilution factor is calculated using MC , comparing cross sections of each materials in target cell. And packing fraction is the only necessary input for each target cell.


Dilution factor for resonance with PF of 61.9\%


## Dilution Factor

Dilution factor is calculated using MC , comparing cross sections of each materials in target cell. And packing fraction is the only necessary input for each target cell.


Dilution factor for DIS with PF of 58.8\%


## Radiative Corrections

Following S. Stein et al., PRD12 (1975) 1884, corrections were mainly done by POLRAD 2.0
Initial fit parameters came from RSS,
basically Breit-Wigner resonance
and polynomial (with some correction) deep inelastic tail.

Newly corrected data was refitted to iterate.


## HMS Asymmetries



## HMS Asymmetries



## HMS Asymmetries

| Setting | Beam energy (GeV) | HMS central momentum (GeV) | HMS angle from beamline (degree) | $\begin{aligned} & <Q^{2}> \\ & \left(G e V^{2}\right) \end{aligned}$ | $\begin{aligned} & <W> \\ & (\mathrm{GeV}) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| (1) | $\begin{gathered} 4.7 \text { (par) / } \\ 5.9 \text { (per) } \end{gathered}$ | $\begin{aligned} & 3.2 \text { (par) / } \\ & 4.4 \text { (per) } \end{aligned}$ | $\begin{gathered} 20.2 \text { (par) / } \\ 15.4 \text { (per) } \end{gathered}$ | 1.863 | 1.353 |
| (2) | 5.9 | 3.1 | 15.4 | 1.313 | 2.196 |
| (3) | 4.7 | 2.2 | 16 | 0.806 | 2.196 |

## SANE-HMS Region 1

$Q^{2}=1.86 \mathrm{GeV}^{2}$
Resonance region

## Parallel and Perpendicular Asymmetries

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



## Asymmetries $\boldsymbol{A}_{1}$ and $\boldsymbol{A}_{2}$

$A_{1}$ and $A_{2}$ are virtual photoabsorption asymmetries.

$$
\begin{aligned}
& A_{1}=\frac{\sigma_{1 / 2}^{T}-\sigma_{3 / 2}^{T}}{\sigma_{1 / 2}^{T}+\sigma_{3 / 2}^{T}}=\frac{\sigma_{T T}}{\sigma_{T}}=\frac{g_{1}-\gamma^{2} g_{2}}{F_{1}} \\
& A_{2}=\frac{2 \sigma_{L T}}{\sigma_{1 / 2}^{T}+\sigma_{3 / 2}^{T}}=\frac{\sigma_{L T}}{\sigma_{T}}=\frac{\gamma\left(g_{1}+g_{2}\right)}{F_{1}}
\end{aligned}
$$

$\sigma_{1 / 2}^{T}$ and $\sigma_{3 / 2}^{T}$ are the virtual photon absorption transverse cross sections when total helicity of photon and nucleon is $1 / 2$ and $3 / 2$ respectively. $\sigma_{L T}$ is the interference term between the transverse and longitudinal photon-nucleon amplitude.
Radiative correction done by iterating the fits of $A_{1}$ and $A_{2}$ until they converged.

## Asymmetry $A_{1}=\frac{\sigma_{T T}}{\sigma_{T}}$

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



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## Asymmetry $A_{1}=\frac{\sigma_{T T}}{\sigma_{T}}$

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



Asymmetry $A_{2}=\frac{\sigma_{L T}}{\sigma_{T}}$

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



## Fitting Function

$$
\text { fit }=\underbrace{\sum_{i=1}^{4} B W_{i}}_{\text {Resonances }}+\underbrace{x^{\alpha} \sum_{n=0}^{3} \beta_{n} x^{n}}_{\text {DIS }} \times \underbrace{\times \frac{1}{\sqrt{Q^{2}}}}_{A_{2} \text { only }}
$$

where

$$
B W_{i}=\frac{a_{i} \kappa_{i}^{2} w_{i}^{2} \Gamma_{i} \Gamma_{i}^{\gamma}}{\kappa_{c m}^{2}\left[\left(w_{i}^{2}-W^{2}\right)^{2}+w_{i}^{2} \Gamma_{i}^{2}\right]}
$$

$$
\begin{aligned}
& =g_{i}\left(\frac{q_{m}}{q_{i}}\right)^{\left(q_{i}+1\right)}\left(\frac{q_{i}^{2}+X_{i}^{2}}{g_{m}^{2}+X_{i}^{2}}\right)^{h} \\
& =g_{i}\left(\frac{k_{m}}{k_{i}}\right)^{(2 i j)}\left(\frac{k_{i}^{2}+X_{i}^{2}}{k_{m}^{2}+X_{i}^{2}}\right)^{j i}
\end{aligned}
$$

$$
\begin{aligned}
\kappa_{i} & =\sqrt{\frac{\left(w_{i}^{2}+M^{2}+Q^{2}\right)^{2}}{4 w_{i}^{2}}-M^{2}} \\
q_{i} & =\sqrt{\frac{\left(w_{i}^{2}+M^{2}-m_{\pi}^{2}\right)^{2}}{4 w_{i}^{2}}-M^{2}} \\
\kappa_{c m} & =\sqrt{\frac{\left(W^{2}+M^{2}+Q^{2}\right)^{2}}{4 W^{2}}-M^{2}} \\
q_{c m} & =\sqrt{\frac{\left(W^{2}+M^{2}-m_{\pi}^{2}\right)^{2}}{4 W^{2}}-M^{2}}
\end{aligned}
$$

Table 3.2: The fitting parameters of $A_{1}$ and $A_{2} . a_{i}$ is the amplitude, $\omega_{i}$ is the centroid, and $g_{i}$ is the width of the i-th BW peak.

| Parameter | $A_{1}$ Fit | $A_{2}$ Fit |
| :---: | :---: | :---: |
| $a_{1}$ | $-0.553 \pm 0.204$ | $-0.306 \pm 0.152$ |
| $a_{2}$ | $0.724 \pm 0.267$ | $-0.474 \pm 0.210$ |
| $a_{3}$ | $0.615 \pm 0.071$ | - |
| $\omega_{1}$ | $1.186 \pm 0.016$ | 1.232 (fixed) |
| $\omega_{2}$ | $1.381 \pm 0.006$ | $1.323 \pm 0.010$ |
| $\omega_{3}$ | $1.547 \pm 0.012$ | - |
| $g_{1}$ | $0.031 \pm 0.025$ | $0.070 \pm 0.057$ |
| $g_{2}$ | $0.053 \pm 0.036$ | $0.058 \pm 0.035$ |
| $g_{3}$ | $0.197 \pm 0.068$ | - |

## Spin Structure Function $g_{1}$

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



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$$
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$$



## Spin Structure Function $g_{1}$

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



## Spin Structure Function $\boldsymbol{g}_{2}$

$$
Q^{2}=1.86 \mathrm{GeV}^{2}
$$



## Preliminary Twist-3 $\mathrm{d}_{2}$ for the Region 3

$$
d_{2}=3 \int_{0}^{1} x^{2}\left(g_{2}-g_{2}^{W W}\right) d x=\int_{0}^{1} x^{2}\left(2 g_{1}+3 g_{2}\right) d x
$$

## OPE valid to

$$
N=2<Q^{2} / M_{0}^{2} \sim 1.8 / 0.5^{2}
$$

per DIS - resonances duality Ji \& Unrau, PRD52 (1995) 72


## $\mathrm{d}_{2}$ Matrix Element



## $\mathrm{d}_{2}$ Matrix Element

$$
\begin{gathered}
\overline{d_{2}}=-0.0087 \pm 0.0014 \\
\overline{d_{2}}=\int_{0.47}^{0.87} x^{2}\left(2 g_{1}+3 g_{2}\right) d x
\end{gathered}
$$

## $\mathrm{d}_{2}$ Matrix Element



## SANE-HMS Region 2

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$

Extension of RSS data into DIS

Parallel and Perpendicular Asymmetries

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$



## Asymmetry $A_{1}=\frac{\sigma_{T T}}{\sigma_{T}}$

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$



## Asymmetry $A_{2}=\frac{\sigma_{L T}}{\sigma_{T}}$

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$



## Spin Structure Function $g_{1}$

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$



## Spin Structure Function $g_{1}$



## Spin Structure Function $g_{1}$

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$





## Spin Structure Function $g_{2}$

$$
Q^{2}=1.31 \mathrm{GeV}^{2}
$$



## SANE-HMS Region 3

## $Q^{2}=0.81 \mathrm{GeV}^{2}$ DIS region

## Parallel and Perpendicular Asymmetries <br> $$
Q^{2}=0.81 \mathrm{GeV}^{2}
$$



Asymmetry $A_{1}=\frac{\sigma_{T T}}{\sigma_{T}}$

$$
Q^{2}=0.81 \mathrm{GeV}^{2}
$$



Asymmetry $A_{2}=\frac{\sigma_{L T}}{\sigma_{T}}$

$$
Q^{2}=0.81 \mathrm{GeV}^{2}
$$



## Spin Structure Function $g_{1}$

$$
Q^{2}=0.81 \mathrm{GeV}^{2}
$$



## Spin Structure Function $\boldsymbol{g}_{2}$

$$
Q^{2}=0.81 \mathrm{GeV}^{2}
$$



## SANE collaboration

U. Basel, Florida International U., Hampton U., Norfolk S. U., North Carolina A\&T S. U., IHEPProtvino, U. of Regina, Rensselaer Polytechnic I., Rutgers U., Seoul National U., Temple U., TJNAF, U. of Virginia, College of William \& Mary, Yerevan Physics I.

Spokespersons: S. Choi (Seoul), M. Jones(Jlab), Z-E. Meziani (Temple), O. A. Rondon (U. of Virginia)

## Summary

SANE-HMS is a measurement of spin structure functions in high Bjorken $x$ and intermediate $\mathbf{Q}^{2}$.

Parallel and perpendicular asymmetries and structure functions show good agreement with previous experiments.
$A_{2}$ and $g_{2}$ show significant $Q^{2}$ evolution. Negative $A_{2}$ at $\mathrm{W}=1.3 \mathrm{GeV}$ is shown. And negative $\mathrm{A}_{2}$ at DIS region can affect $g_{1}$ deduced from parallel asymmetry only (e.g. Hall B results).
, $\overline{d_{2}}=-0.0087 \pm 0.0014$ is the first negative result of $d_{2}$ matrix element, although its integration range is limited.

## Backup Slides




## Systematic Errors

| Error Source | Average |
| :--- | ---: |
| Target Polarization | $4.0 \%$ |
| Beam polarization | $1.5 \%$ |
| Dilution Factor | $3.3 \%$ |
| Nitrogen Correction | $0.4 \%$ |
| Radiative Corrections | $10 \%\left(A_{1}\right.$ and $\left.A_{2}\right)$ |
| Kinematic Reconstruction |  |






## Twist-3 d 2



## Useful relations

$$
\begin{aligned}
g_{1} & =\frac{F_{1}}{1+\gamma^{2}}\left(A_{1}+\gamma A_{2}\right), & A_{1}=\frac{g_{1}-\gamma^{2} g_{2}}{F_{1}} \\
g_{2} & =\frac{F_{1}}{1+\gamma^{2}}\left(-A_{1}+\frac{A_{2}}{\gamma}\right) & A_{2}=\gamma \frac{g_{1}+g_{2}}{F_{1}}, \\
A_{\|} & =D\left(A_{1}+\eta A_{2}\right), &
\end{aligned}
$$

## Spin Structure Functions

Inclusive DIS cross section depends on four structure functions, two unpolarized ( $\mathrm{F}_{1}, \mathrm{~F}_{2}$ ) and two polarized $\left(g_{1}, g_{2}\right)$. The spin structure functions $g_{1}$ and $g_{2}$ can be experimentally determined by measuring spin asymmetries:

$$
\begin{aligned}
& A_{\|}=\frac{\sigma^{\downarrow \uparrow}-\sigma^{\uparrow \uparrow}}{\sigma^{\downarrow \uparrow}+\sigma^{\uparrow \uparrow}}, \quad A_{\perp}=\frac{\sigma^{\downarrow \rightarrow}-\sigma^{\uparrow \rightarrow}}{\sigma^{\downarrow \rightarrow}+\sigma^{\uparrow \rightarrow}} \\
& g_{1}\left(x, Q^{2}\right)=\frac{F_{1}\left(x, Q^{2}\right)}{d^{\prime}}\left[A_{\|}+\tan (\theta / 2) A_{\perp}\right] \\
& g_{2}\left(x, Q^{2}\right)=\frac{y F_{1}\left(x, Q^{2}\right)}{2 d^{\prime}}\left[\frac{E+E^{\prime} \cos (\theta)}{E^{\prime} \sin (\theta)} A_{\perp}-A_{\|}\right]
\end{aligned}
$$

## ep deep inelastic scattering

High-energy electron-nucleon scattering(Deep Inelastic Scattering) ep $\rightarrow$ e'X
k and $\mathrm{k}^{\prime}$ are the four-momenta of the incoming and outgoing electrons, $P$ is the four-moemntum of a proton with mass $M$, and W is the mass of the recoiling system X .
q is the four-momentum of the virtual photon(the exchanged particle). ( $Q^{2}=-q^{2}$ )


## Spin structure functions

When the spins of electron and nucleon are all polarized, we can see the dependence of scattering cross section on the spin structure functions $g_{1}\left(x, Q^{2}\right)$ and $g_{2}\left(x, Q^{2}\right)$.

$$
\begin{aligned}
g_{1} & =\frac{1}{2} \sum_{i} e_{i}^{2}\left[q_{i}^{+}-q_{i}^{-}\right] \\
g_{2} & =g_{2}^{W W}+\overline{g_{2}}
\end{aligned}
$$

$g_{2}^{W W}\left(x, Q^{2}\right)=-g_{1}\left(x, Q^{2}\right)+\int_{x}^{1} \frac{g_{1}\left(x^{\prime}, Q^{2}\right)}{x^{\prime}} d x^{\prime}$

## High Momentum Spectrometer

## 2/25/2009 9:29:15 AM

Hall A 4K Supply Flow: $35.4 \mathrm{~g} / \mathrm{s}$
Hall B 4K Supply Flow:
Hall C 4K Supply Flow:
H.
gis

High Momentum Spectrometer

| $\begin{aligned} & \text { Momentum } \\ & \text { Input }(\mathrm{GeV}) \\ & \hline \end{aligned}$ | Rumning | Stop | $\square \square \square$ Protron $\square \square \square \square$ Electron | Ramp All to Zero Amps |
| :---: | :---: | :---: | :---: | :---: |

$\square$ Input OK
Q1 at Current $Q 2$ at Current $Q 3$ at Current $\quad$ Dipole at Field


## World Data and SANE Region

World data lacks big region, especially in the perpendicular asymmetry. SANE covers broad region of


## Packing Fraction

Comparing data with Monte Carlo results assuming 50\% and 60\% packing fraction of target, 60.9\% packing fraction is determined for the target material \#9 6-28-07 14NH3.


## Unpolarized Strucuture Functions



## ep deep inelastic scattering

Invariant quantities:
$\nu=\frac{q \cdot P}{M}=E-E^{\prime}$ is the lepton's energy loss in the nucleon rest frame (in earlier literature sometimes $\nu=q \cdot P)$. Here, $E$ and $E^{\prime}$ are the initial and final lepton energies in the nucleon rest frame.
$Q^{2}=-q^{2}=2\left(E E^{\prime}-\vec{k} \cdot \vec{k}^{\prime}\right)-m_{\ell}^{2}-m_{\ell^{\prime}}^{2}$ where $m_{\ell}\left(m_{\ell^{\prime}}\right)$ is the initial (final) lepton mass. If $E E^{\prime} \sin ^{2}(\theta / 2) \gg m_{\ell}^{2}, m_{\ell^{\prime}}^{2}$, then
$\approx 4 E E^{\prime} \sin ^{2}(\theta / 2)$, where $\theta$ is the lepton's scattering angle with respect to the lepton beam direction.
$x=\frac{Q^{2}}{2 M \nu}$ where, in the parton model, $x$ is the fraction of the nucleon's momentum carried by the struck quark.
$y=\frac{q \cdot P}{k \cdot P}=\frac{\nu}{E}$ is the fraction of the lepton's energy lost in the nucleon rest frame.
$W^{2}=(P+q)^{2}=M^{2}+2 M \nu-Q^{2}$ is the mass squared of the system $X$ recoiling against the scattered lepton.

## Radiative Correction

$$
A=\frac{1}{f C_{N} P_{b} P_{t} f_{R C}} \frac{d \sigma^{\downarrow \uparrow}-d \sigma^{\uparrow \uparrow}}{d \sigma^{\downarrow \uparrow}+d \sigma^{\uparrow \uparrow}}+A_{R C}
$$

1. Incoming and
outgoing electron lose energy before and after scattering.



CLAS eg1b data


CLAS eg1b data

