Matrix elements from lattice QCD

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reviewing work from many groups

presented at Chiral Dynamics 2012

High-precision calculations of f_{π} ,... in the last decade

Precision calculations of $m_N,\ldots\,$ in the last few years

Well-controlled pion matrix elements are possible now

Intense work on much harder nucleon matrix elements

Lattice QCD can do high-precision calculations



first precision calculations were completed in 2003 there is a notable absence of any baryonic properties

averages of results from arXiv:1110.0016 and http://latticeaverages.org/

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Precision baryon physics is just recently feasible



first calculation in 2009 due to challenge of baryons example of computational thresholds in lattice QCD

BMW collaboration, Science, arXiv:0906.3599, see ETMC and CP-PACS too 3/15

space-like form factor is calculable in Euclidean space

$$\langle \pi, p | V^{\mathsf{em}}_{\mu} | \pi, p'
angle = K_{\mu} F(Q^2) \qquad K_{\mu} = p_{\mu} + p'_{\mu}$$

F(0) = 1, so focus on the slope or the charge radius

$$\langle r^2 \rangle_{\pi} \equiv 6 \left. \frac{dF(Q^2)}{dQ^2} \right|_{Q^2 \to 0}$$

no renormalization is required but $Q^2 \neq 0$ is needed

Precision calculations are essential

good example matrix element to illustrate the need ...



to carefully and precisely control all uncertainties

Pion charge radius



computational thresholds were not met in early results

Pion charge radius



could now calculate pion charge radius to about 5%

Chiral extrapolation of pion charge radius



illustrates the $\chi PT/LQCD$ extrapolation strategy

Pion form factor



lattice calculation is comparable to the measurements

Proton form factors are more challenging

proton leads to two vector form factors, F_1 and F_2

$$\langle N, p | V_{\mu} | N, p' \rangle = K^{1}_{\mu} F_{1}(Q^{2}) + K^{2}_{\mu} F_{2}(Q^{2})$$

now, charge radius and anomalous magnetic moment

$$\langle r^2 \rangle_p \equiv 6 \left. \frac{dF_1(Q^2)}{dQ^2} \right|_{Q^2 \to 0} \qquad \kappa = \lim_{Q^2 \to 0} F_2(Q^2)$$

no renormalization is required but $Q^2 \neq 0$ is needed



no reliable extrapolation, examine raw lattice results



calculations nearly to the physical \mathbf{m}_{π} a major triumph



but well-controlled calculations needed to resolve this



divergence suggests possibly large finite-size effects



finite-size effects at small \mathbf{m}_{π} may resolve this puzzle

Proton axial coupling should be easier

proton requires two axial form factors, $g_{\rm A}$ and $g_{\rm P}$

$$\langle N, p | A^{u-d}_{\mu} | N, p' \rangle = K^A_{\mu} g_A(Q^2) + K^P_{\mu} g_P(Q^2)$$

now both form factors have non-trivial forward limits

$$g_A = g_A(Q^2 = 0)$$
 $g_P = \lim_{Q^2 \to 0} g_P(Q^2)$

 $Q^2 \neq 0$ is not needed for g_A , but a finite Z_A is required



axial coupling has persistently been flat and too low



results approaching the physical point are still too low



but apparent discrepancy is not too large to start with



rule-of-thumb $m_{\pi}L > 4$ is known to be insufficient here



it is known that $m_{\pi}L > 6$ may even be necessary here

Proton parton distributions

GPDs related to generalized form factors $A_{ni}\text{, }B_{ni}\text{, }C_{n}$

$$\langle N, p | O_{\mu\nu} | N, k \rangle = K^A_{\mu\nu} A_{20}(Q^2) + K^B_{\mu\nu} B_{20}(Q^2) + K^C_{\mu\nu} C_2(Q^2)$$

like F_1 or g_A , only A_{n0} has an accessible forward limit

$$\int_{-1}^{1} dx \, x \, q(x) = \langle x \rangle = A_{20}(Q^2 = 0)$$

scale-dep. renormalization is needed but $Q^2 \neq 0$ is not



long standing trend for $\langle x \rangle^{u-d}$ to be quite flat in m_π



will return shortly to this result close to physical limit



again, well-controlled calculations needed here too



corrections by LHPC/ETMC may resolve this puzzle



possible curvature but lightest point has $m_{\pi}L = 2.7$



Well-controlled calculations for the pion are feasible

Intense progress for the nucleon is being made

Apparent conflicts with measurements not justified

Apparent conflicts with χPT not compelling either