

Structure of Antikaon-Nucleon Scattering

Maxim Mai

(The George Washington University)

M. Mai and U. G. Meißner, Nucl. Phys. A **900** (2013)

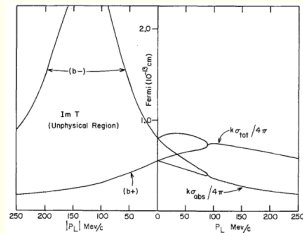
M. Mai and U.-G. Meißner, Eur. Phys. J. A **51** (2015)

L. Roca, M. Mai, E. Oset and U.-G. Meißner, Eur. Phys. J. C **75** (2015)

A. Cieplý, M. Mai, U.-G. Meißner and J. Smejkal, Nucl. Phys. A **954** (2016)

- “Four sets of scattering amplitudes are obtained consistent with all the present data on K -proton interactions and the possibilities for discrimination between them are discussed. Two of these amplitudes are found to correspond to a **resonance-like** behavior just within the unphysical region.”

Dalitz, Tuan (1959)



⇒ *What is this resonance?*

$(S = -1; I = 0; J^P = 1/2^{-(?)}; M \sim 1410 - 18i \text{ MeV})$

A. Dynamically generated from coupled-channel effects

- K-matrix
- unitarized coupled-channel amplitude from ChPT
- ⇒ **two-pole solution**

Dalitz, Tuan (1960!)

Kaiser, Siegel, Weise (1995)

Oller, Meißner (2001)

B. Quark model → genuine qqq state

Capstick, Isgur (1986)

PDG favors A

PDG MiniReview (2015)

Lattice QCD favors A

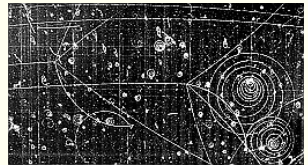
Hall et al. (2014)

Experimental situation (no photons)

✓ Total cross sections on $K^-p \rightarrow K^-p, \bar{K}^0n, \dots$

LNL (1960s), Rutherford Lab(1980s), ...

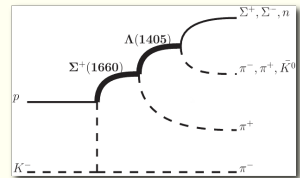
- bubble chamber experiments
- huge error bars
- large deviations btw. experiments



✓ $\pi\Sigma$ mass distribution

- 2m bubble chamber @ CERN
- low energy resolution
- multistep production \Rightarrow new parameter

Hemingway (1985)



✓ Strong energy shift and width in $\bar{K}H$

$\Rightarrow a_{K^-p}$ from modified Deser-type relation

Bazzi et al.(2011)

Meißner, Raha, Rusetsky (2004)

✗ Plans for an upgrade to $\bar{K}D$

$\Rightarrow a_1, a_0$ from *Faddeev equations/EFT*

Shevchenko (2014).../MM et al.(2014)...

✗ pp collisions

COSY (2008) HADES (2013)

\Rightarrow high quality data, **but** theoretical analysis very intricate

Scattering amplitude - framework

- ChPT is an appropriate tool to study low-energy hadronic interactions.

Weinberg (1979) Gasser, Leutwyler (1981)

Here it has to fail! Because:

1. Kaon mass is large → convergence
 2. Relevant thresholds are widely separated → convergence
 3. Resonance just below $\bar{K}N$ threshold → non-perturbative effect
- Non-perturbative methods:
 - Dispersion relations, N/D , Roy-Steiner equations
 - K-Matrix, JÜLICH-BONN model, ...
 - IAM, *Chiral Unitary Models*, ...
 - Chiral Unitary Models - driving term*

$$V(\not{q}_2, \not{q}_1; p) = A_{WT}(q_1 + q_2) + Born(s) + Born(u) \\ + A_{14}(q_1 \cdot q_2) + A_{57}[\not{q}_1, \not{q}_2] + A_M + A_{811} \left(q_2(q_1 \cdot p) + q_1(q_2 \cdot p) \right)$$

⇒ $A_{..}$ depend on low energy constants ⇒ free parameters

Scattering amplitude - re-summation

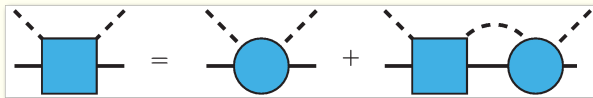
- Bethe-Salpeter equation

Salpeter et al. (1951)

$$T(q_2, q_1; p) = V(q_2, q_1; p) + i \int \frac{d^d l}{(2\pi)^d} \frac{V(q_2, l; p) T(l, q_1; p)}{((\not{p} - \not{l}) - m + i\epsilon)(l^2 - M^2 + i\epsilon)}$$

→ Intermediate particles are *off-shell*

⇒ exactly corresponding to a series of Feynman loop diagrams



⇒ BSE can be solved analytically, if $(\mathbf{f}) V \sim$ local terms Bruns, MM, Meißner (2011)

⇒ drop the Born graphs

→ *Bubble chain* in s direction ⇒ topologies are missing

⇒ scale dependence **does not** cancel out ⇒ additional model parameters

→ Off-shell effects are moderate

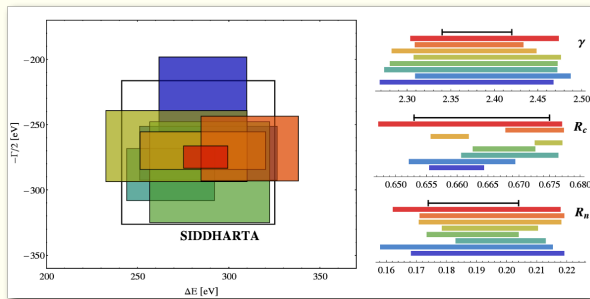
MM, Meißner (2013)

⇒ use on-shell approximation first (scan of 20-dim. parameter space)

Fits and results

- Randomly chosen sets of starting values ($\# \approx 10000$)
- Solutions having poles on I. RS sorted out
- **Results:** 8 best fits obtained with similar $\chi^2_{\text{d.o.f.}}$.

Fit #	1	2	3	4	5	6	7	8
$\chi^2_{\text{d.o.f.}}$	1.35	1.14	0.99	0.96	1.06	1.02	1.15	0.90

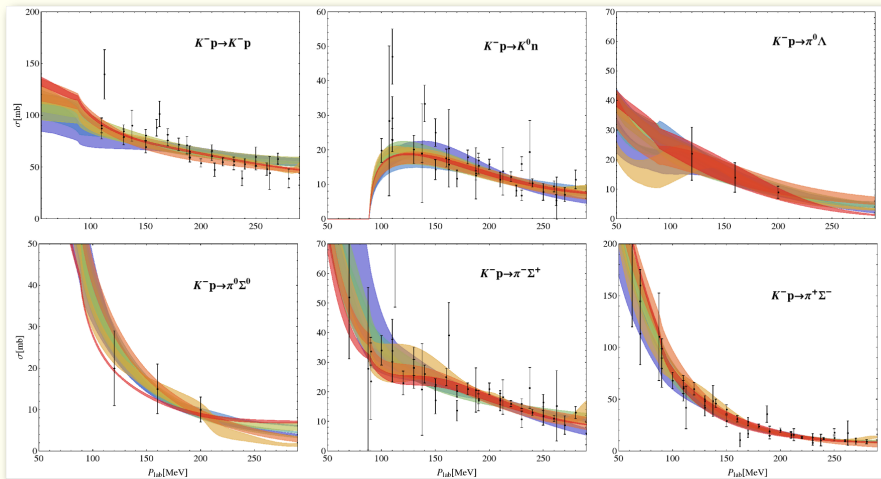


→ error bars are twofold

1. parameter: variation of best fit parameters, such that $\Delta\chi^2_{\text{d.o.f.}} < 1.15$
2. systematic: spread of solutions

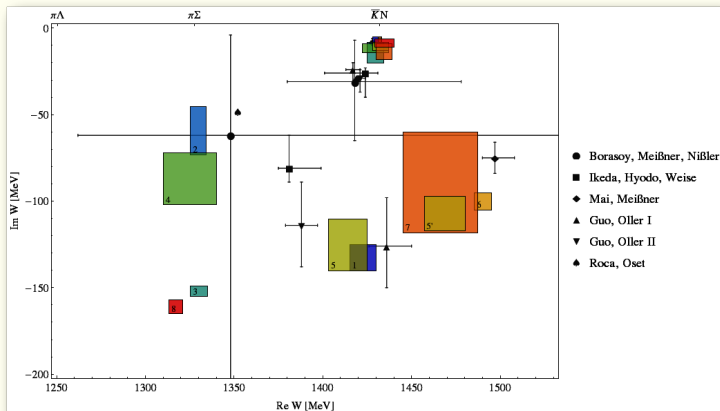
Results

→ similar cross sections



Results - complex plane

- Analytic continuation to the complex energy plane



- two poles in all solutions on II. RS
- stable position of the narrow pole
- position of the second pole is rather unstable

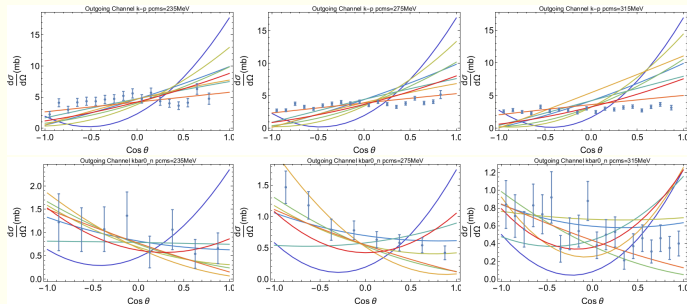
Can we reduce this ambiguity? PART 1

1) Differential cross sections

Mast et al. (1976)

- Chiral Unitary Meson-Baryon amplitude has a genuine P -wave part
- Does the prediction agree with data? [analyzed by D. Sadasivan]

$$K^- p \rightarrow K^- p$$



$$K^- p \rightarrow \bar{K}^0 n$$

- solution #1 seems to deviate strongly from data
- further constraints require a refit

Can we reduce this ambiguity? PART 2

2) New data on total cross sections - synthetic data from fit #4:

$$p_{lab} = 100 \dots 300 \text{ MeV}, \Delta E = 5, 10 \text{ MeV}, \Delta\sigma = 2.5, 5, 10$$

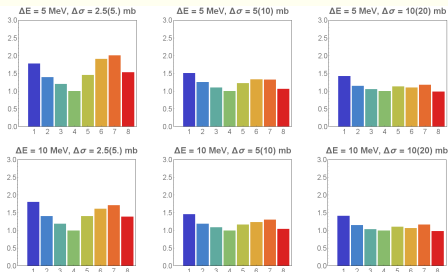
Compare $\chi^2_{d.o.f.}/\chi^2_{d.o.f.} (\#4)$

→ threshold ratios, SIDDHARTA

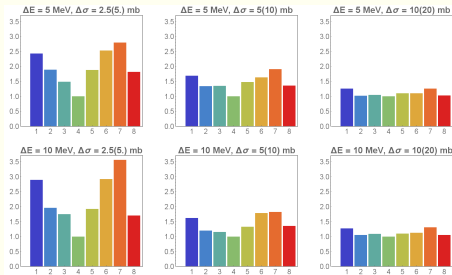
→ pseudo and real scattering data

→ threshold ratios, SIDDHARTA

→ pseudo scattering data



⇒ $\Delta\sigma < 5(10) \text{ mb}$ and $\Delta E < 10 \text{ MeV}$ desired



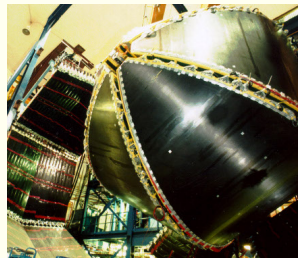
⇒ much larger values of $\Delta\sigma$ and ΔE are sufficient

Can we reduce this ambiguity? PART 3

3) CLAS data on $\gamma p \rightarrow K^+ \pi \Sigma$

- $\pi \Sigma$ mass distribution
- electro- and photoproduction:
 $\gamma p \rightarrow (K^+) \Lambda(1405) \rightarrow \pi \Sigma$
- $J^P = \frac{1}{2}^-$ “confirmed” experimentally
- high statistics and good angular resolution

⇒ theoretical analysis requires a **photoproduction amplitude**



Photoproduction amplitude

I Gauge invariant approaches

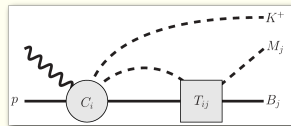
1. “Turtle” approximation Gross, Riska (1987), Kvinikhidze et al.(1999) Borasoy et al.(2005)
 - attach photon everywhere to *off-shell* hadronic amplitude
 - single meson case is done for the NLO-kernel MM et al.(2012)
2. Gauged vertices Nakamura, Jido (2014)
 - photon attached to meson production amplitude at the tree level
 - ⇒ at the LO driving term no good fit to CLAS data
 - ⇒ good fit with additional vector meson d.o.f. - **15 per energy bin!**

II Test model

- most simple ansatz to test the hadronic solution:

$$\mathcal{M}_j(W, M_{\pi\Sigma}) = C_i(W) \cdot G_i(M_{\pi\Sigma}) \cdot T_{i \rightarrow j}^{on}(M_{\pi\Sigma})$$

- flexible enough for the CLAS data Oset, Roca (2013)
 - ⇒ less free parameters (15 ↦ 10)
- no gauge invariance, parameters are not physical
 - ⇒ global fit is meaningless
 - ⇒ **conservative** test of the hadronic solutions

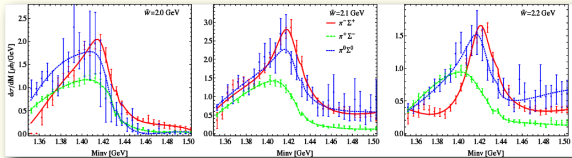


Results

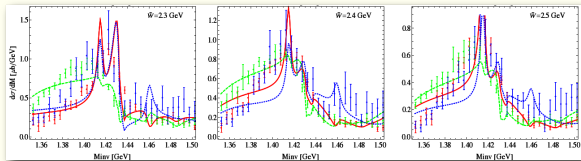
- Test of hadronic solutions

Fit #	1	2	3	4	5	6	7	8
$\chi^2_{\text{d.o.f.}} (\text{hadr.})$	1.35	1.14	0.99	0.96	1.06	1.02	1.15	0.90
$\chi^2_{\text{p.p.}} (\text{CLAS})$	3.18	1.94	2.56	1.77	1.90	6.11	2.93	3.14

- Hadronic fits #2, #4 and #5 lead to *good* fits

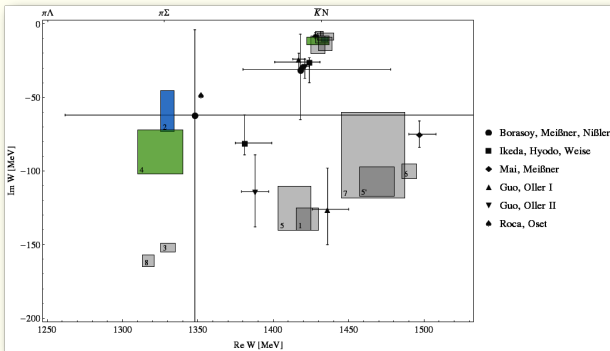


- Hadronic fits #1, #3, #6, #7 and #8 do not!!!



Results - comparison

⇒ after comparison with Hemingway data ($K^- p \rightarrow \Sigma^+ \pi^- \pi^+ \pi^-$) two solutions remain: #2 and #4



⇒ both solutions have similar pole positions

... also similar to the estimation by Oset and Roca (2013)

⇒ **universal feature** demanded by CLAS data!

Summary

- The NLO chiral unitary $\bar{K}N$ amplitude used to analyze hadronic data
- 8 solutions are found in the *on-shell* approximation
 - *the position of the narrow pole is quite certain*
 - *broad pole has large systematic uncertainty*
- Fit to differential cross section is possible ... **work in progress**
- Photoproduction amplitude constructed from the hadronic part
 - *simple, but very flexible ansatz ... conservative test*
 - *5 solutions disagree with the CLAS data, 2 remain after all tests*
- New data can actually reduce the ambiguity of the $\bar{K}N$ amplitude
 - *desired accuracy is **not a part of science-fiction***