

# Single-energy fits

## $\pi$ N Photoproduction

## $\pi$ N scattering

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Have considered **single-energy fits** in the context of amplitude reconstruction.

Here we will describe the single-energy fits displayed on **SAID/MAID** plots of amplitudes.

These have been used in multi-channel fits [ e.g. Geissen , Bonn-Gatchina ] in lieu of actual `data' for  $\pi$ N scattering and  $\pi$ N photoproduction.

Important to understand how they are determined, why they were produced, and what they say about the underlying database.

# Single-energy fits

## $\pi N$ Photoproduction

## $\pi N$ scattering

compare

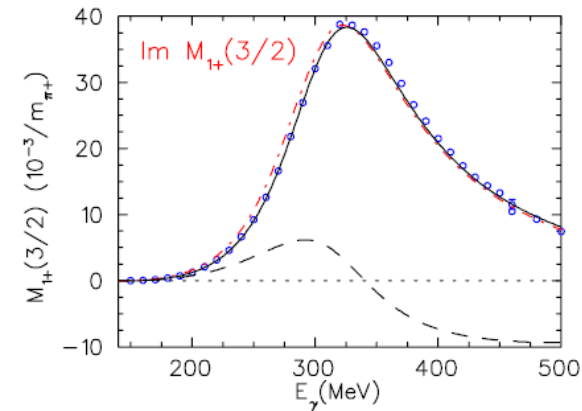
Historical motivation:

Suppose you have a global (energy-dependent) fit to 25K data, over a 2 GeV energy range, varying 0.2K parameters.

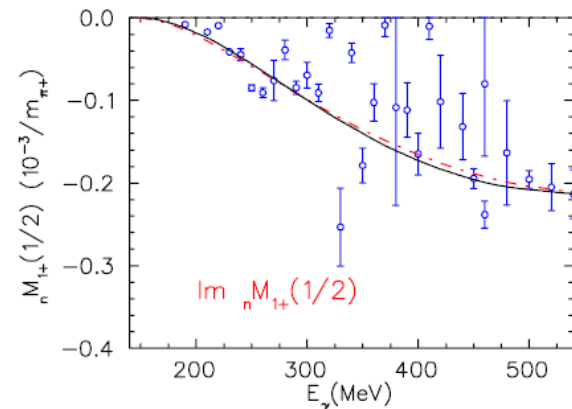
Can the fit be improved?

Bin data over narrow energy ranges. Vary the most significant partial waves – assume the phase found from the global fit is correct. Fit the data along with partial-wave pseudo-data to keep the fit 'close' to global solution – look for systematic deviations.

Consider what may be missing from your model to improve the global fit result.



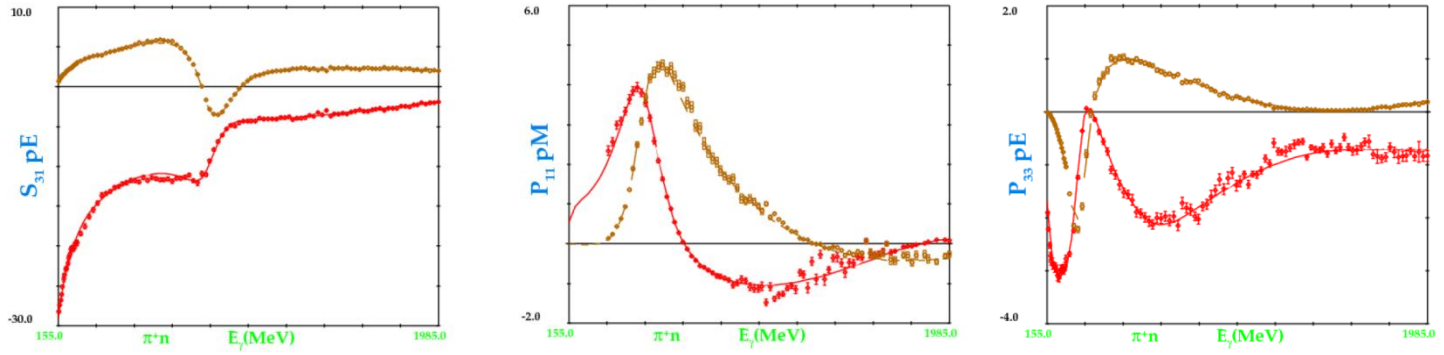
to



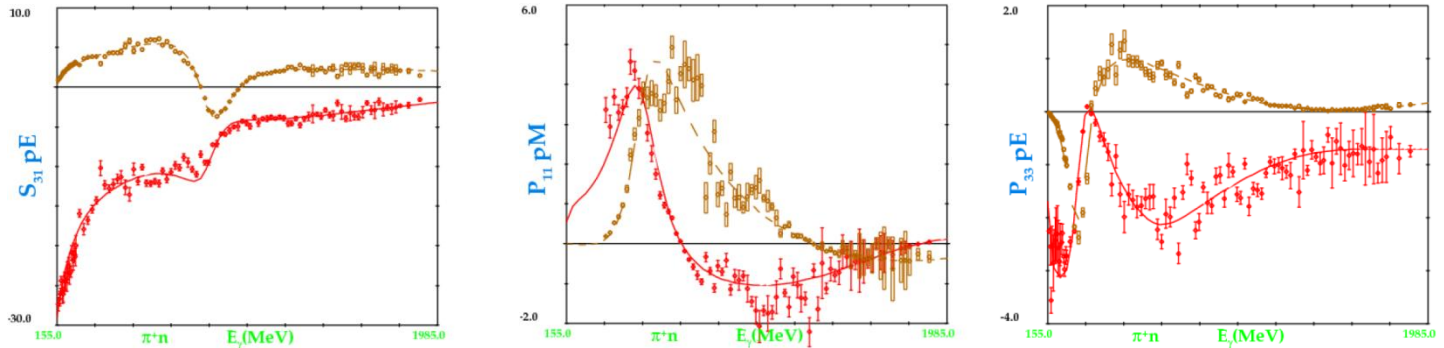
Single-energy fits  
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Single-energy vs global fit

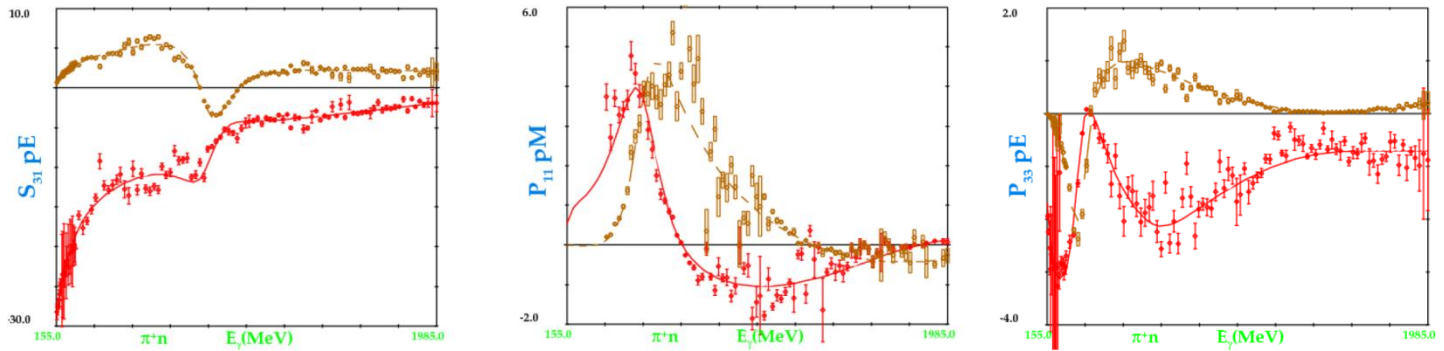
Pseudo-data  
errors/10



Result on SAID



Pseudo-data  
errors \* 10



# Single-energy fits

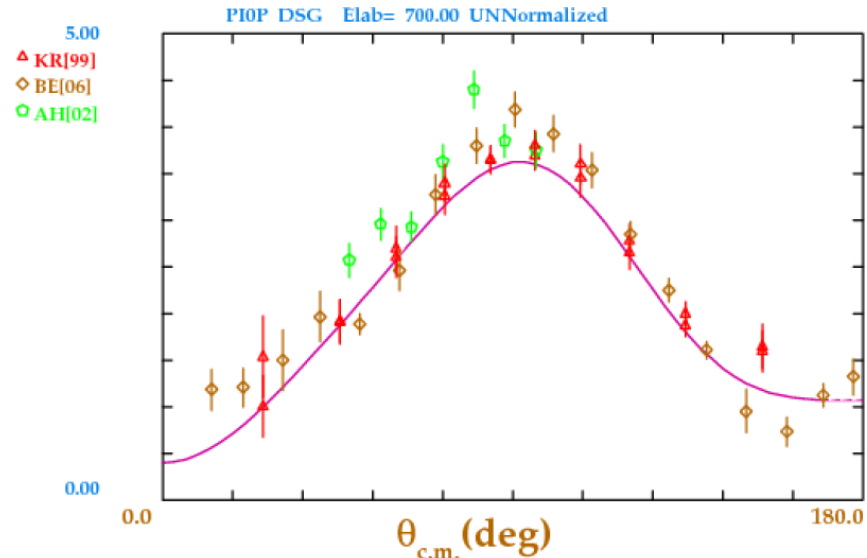
## $\pi$ N Photoproduction

## $\pi$ N scattering

Results depend on (subjective) level of constraint from the global fit. Tighter constraints give back the global result.

Some partial-waves (with little constraint) scatter erratically.

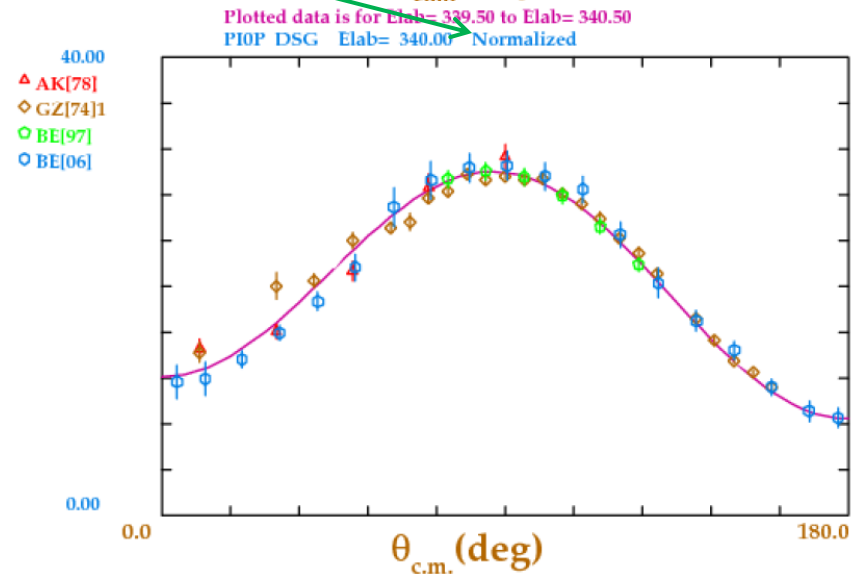
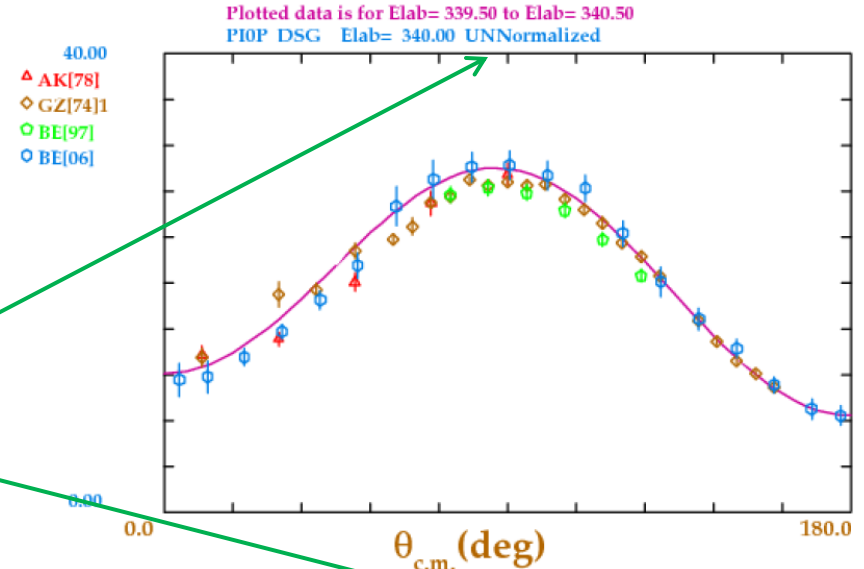
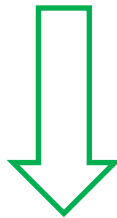
Problem: bins contain old and new data (some contradictory) and may not contain sufficient observables to give a unique result.



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 $\pi$ N scattering

Some discrepancies can be reduced based on the given overall systematic error through a renormalization

This enters into a modified chi-squared definition



# Single-energy fits

## $\pi$ N Photoproduction

## $\pi$ N scattering

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For details see  
 R.A. Arndt, M.H. MacGregor  
*Methods in Computational  
 Physics*, Vol. 6 , 253 (1966).

$$\chi^2 = \sum_{i=1}^{N_D} \left( \frac{\alpha \theta^i(p) - \theta_{exp}^i}{\delta \theta_{exp}^i} \right)^2 + \sum_{j=1}^{N_\alpha} \left( \frac{\alpha_j - 1}{\delta \alpha_j} \right)^2$$

where  $\alpha_j$  factors renormalize ( $N_\alpha$ ) angular distributions, including a total of  $N_D$  data. The  $\alpha_j$  factors do not have to be included explicitly in the parameter ( $p$ ) search.

At the minimum,

$$p \rightarrow p \pm \Delta p, \quad \chi^2 \rightarrow \chi^2 + 1, \quad Prob \rightarrow \frac{1}{e^{1/2}} Prob$$

This was used recently in a  
 kaon photoproduction fit by  
 Sandorfi et al.,  
 J. Phys. G38, 053001 (2011).

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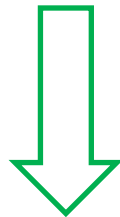
Try to generate single-energy fits based on a database generated from known amplitudes.

Can test effect of more (consistent) measurements on the fit.

Tiator talk  
at NSTAR  
2011.

Exercise was done using MAID Monte Carlo data, then fitted using the SAID global fit.

[ observables generated to be `reasonable' given existing beam/detector quality ]

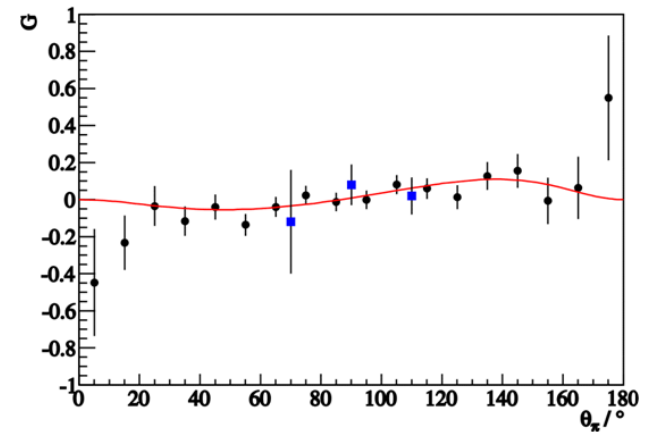
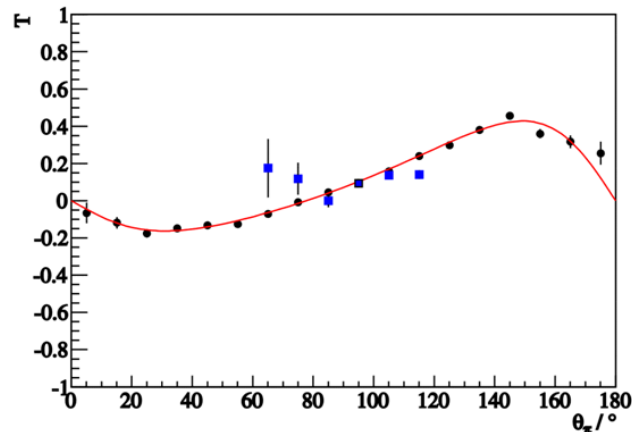
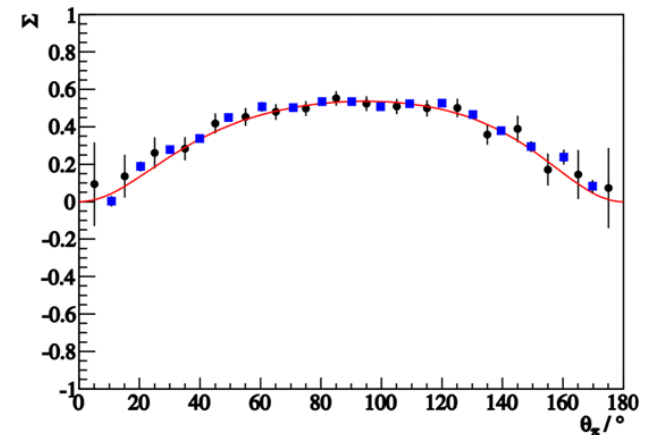
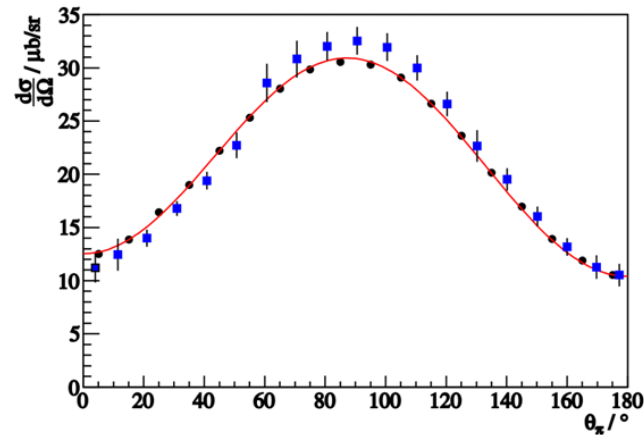


# Single-energy fits

## $\pi N$ Photoproduction

### $\pi N$ scattering

MD07 (red curve) Pseudo-data (**black** points) Real data (**blue** points)  
 $\pi^0 p$  at 320-340 MeV and comparison with real data

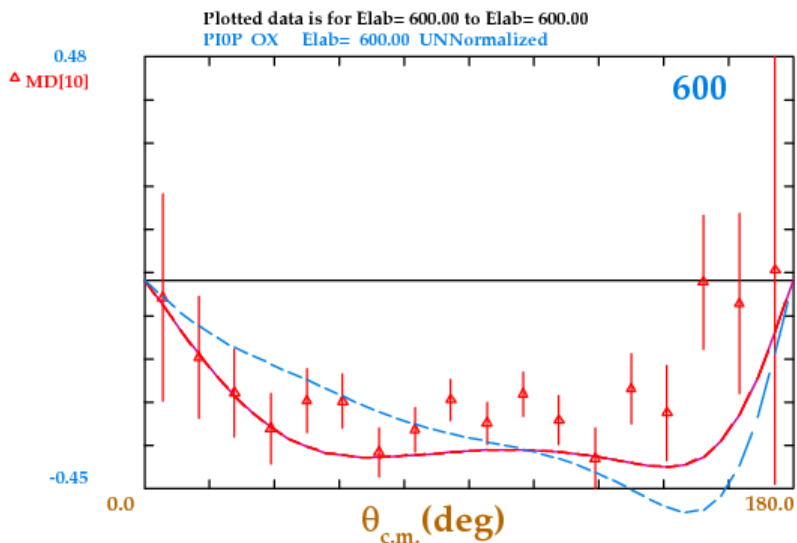




# Single-energy fits

## $\pi$ N Photoproduction

## $\pi$ N scattering

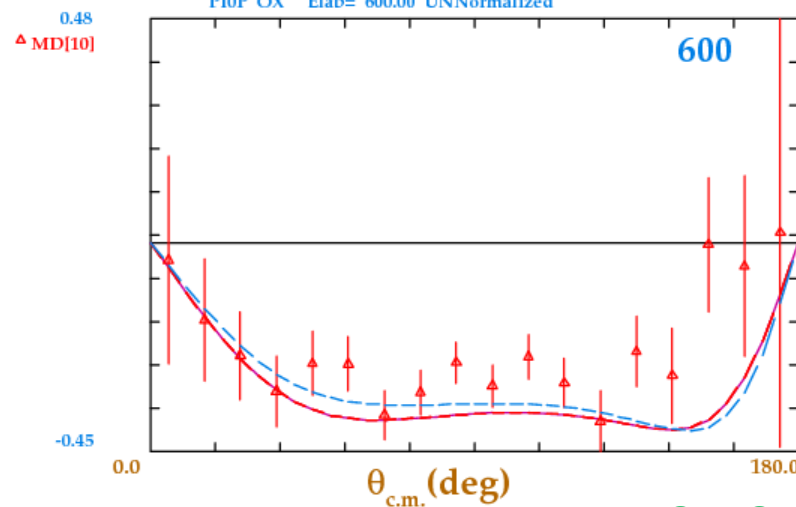
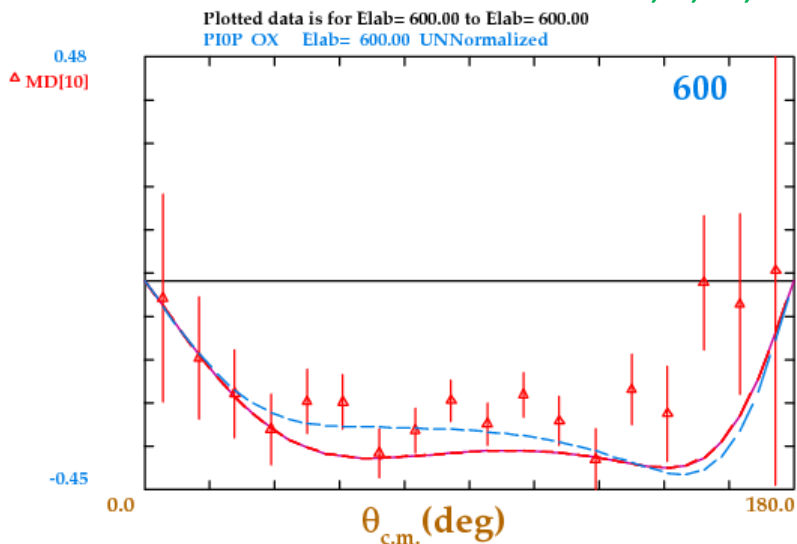


$d\sigma/d\Omega$ , P,  $\Sigma$ , T

Prediction compared to a fit of double-polarization observable

$$O_x \frac{d\sigma}{dt} = -2\text{Re}(b_1 b_3^* - b_2 b_4^*)$$

E, F, G, H

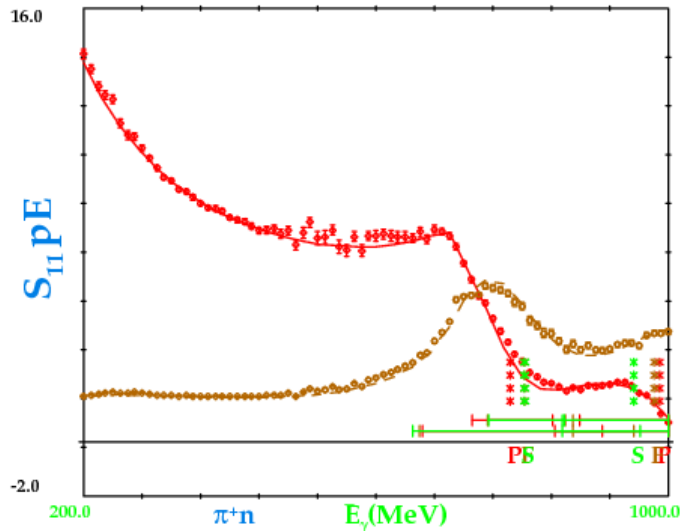


$O_x, O_z, C_x, C_z$

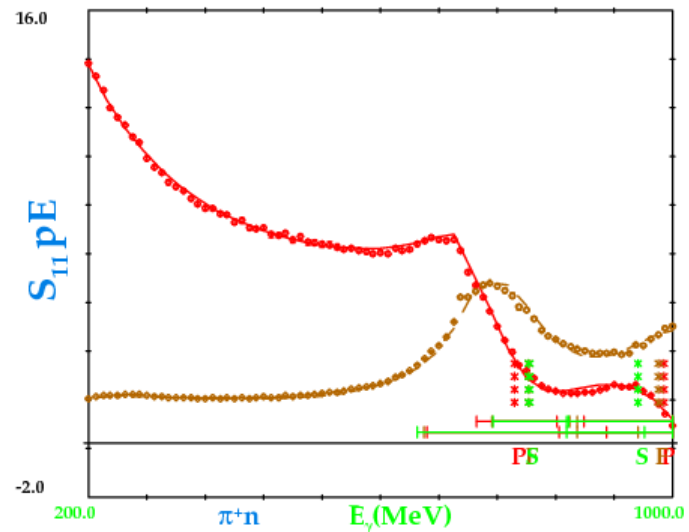
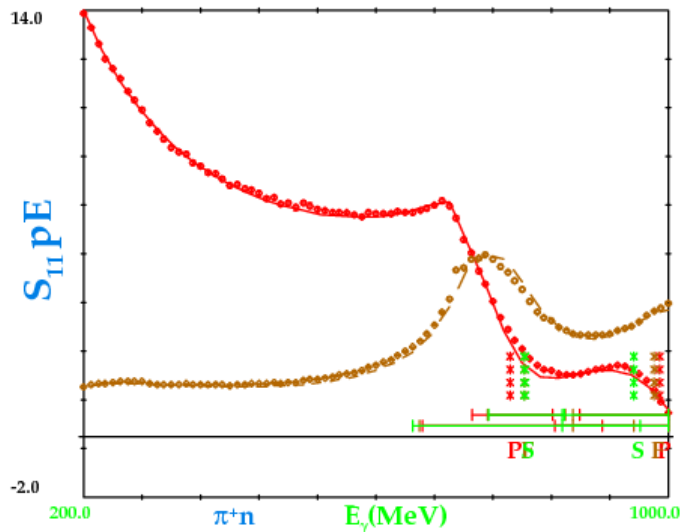
# Single-energy fits

## $\pi$ N Photoproduction

## $\pi$ N scattering



Multipole: predicted vs input



# Single-energy fits

## $\pi$ N Photoproduction

## $\pi$ N scattering

PseudoData - Mozilla Firefox

File Edit View History Bookmarks Tools Help

http://gwdac.phys.gwu.edu/analysis/go56pr.html

Most Visited Getting Started Ohio University : Co...

PseudoData

## Compare Fits to Pseudo Data

[Instructions for comparing PseudoData Fits](#)

Pseudodata generated from MD07 are fitted globally by energy dependent solution LS12. From this starting point, single energy solutions (SES) are generated at several energies. The goal is to study the model-independence of multipole fits without database variations. The fits at 340, 420, and 600 MeV are based on just pseudodata for DSG, P, S, T. Results for other observables are predictions.

Two curves: Red curve: energy-dependent result , Blue curve: SES result  
Single curve: (no energy dependent soln) SES result with error band

[Fits](#) to an 8 observable dataset.  
[Fits](#) to a 12 observable dataset.  
[SES](#) results compared to MD07.

These are preliminary fits for discussion purposes only. Please do not quote the results without permission. The results may change and will be expanded as more fits become available. If you have comments/suggestions, they are welcome: send them to rworkman@gwu.edu

**Choose a Solution:**

[MAID 2007](#)  
 LS12  
 None

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**Choose 2nd Solution:**

SES 340 MeV  
 SES 420 MeV  
 SES 600 MeV

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**Choose a Reaction Type:**

PI0\_P  PI+\_N

**Choose Observables:**

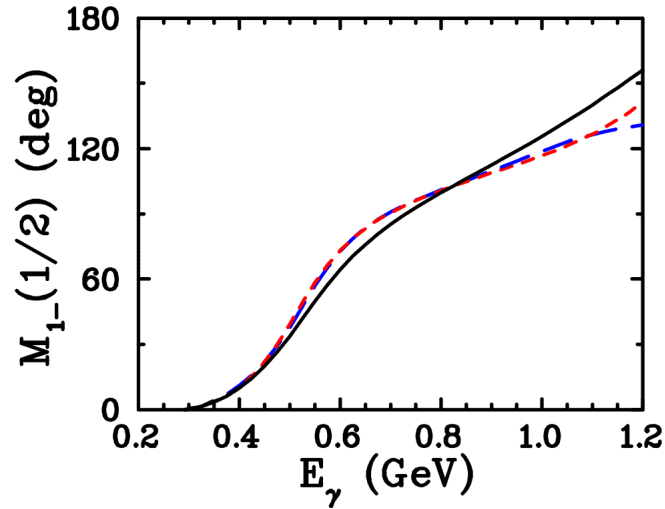
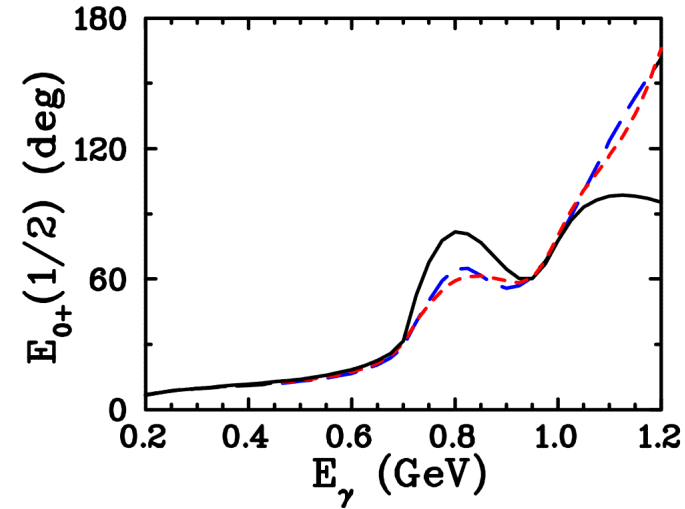
(DSG ) (P ) (S ) (T ) (SGT ) (G ) (H )  
(E ) (F ) (Ox ) (Oz ) (Cx ) (Cz ) (Tx )  
(Tz ) (Lx ) (Lz )  
DSG: differential cross section

**Enter one of the above observable types:**

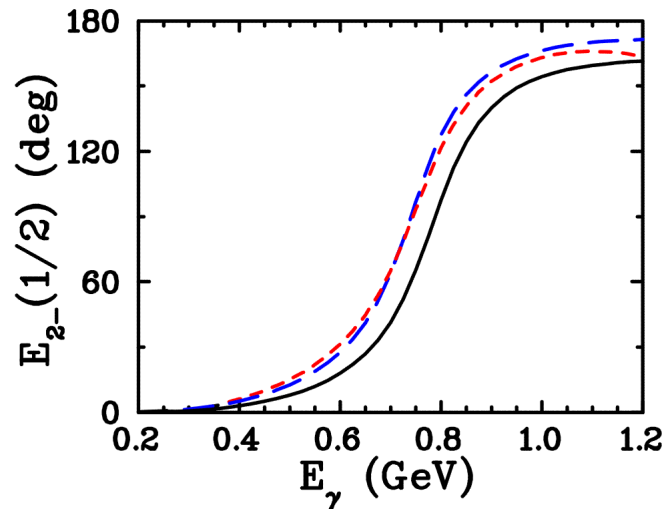
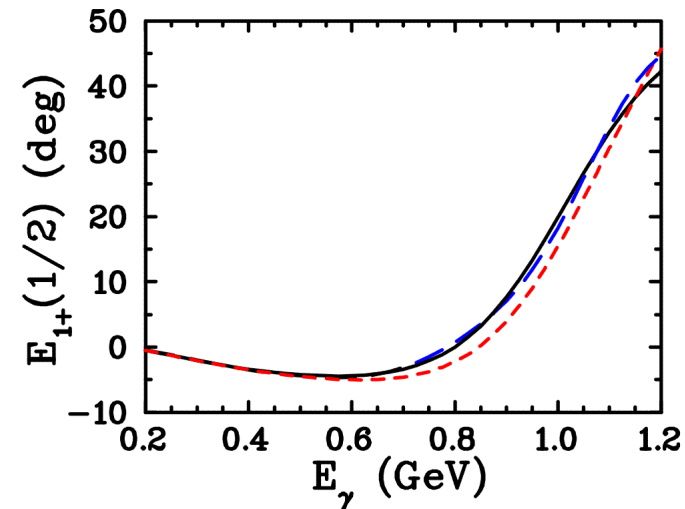
DSG

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SAID vs MAID phase differences



SAID ———  
 MAID - - -  
 ED fit - - -



Effect of searching phases in single-energy fits

Multipole	SP09	LS12	SE4	SE8	MD07
$E_{0+}^{1/2}$	18.4	17.4	16.0(3.0)	16.2(0.9)	16.7
$M_{1-}^{1/2}$	64.4	73.2	68.2(2.4)	73.4(1.6)	72.7
$M_{1-}^{3/2}$	167.1	172.0	176.8(6.5)	167.5(1.8)	163.9
$E_{2-}^{1/2}$	17.9	31.4	27.5(3.1)	26.1(1.0)	27.6
$M_{2-}^{1/2}$	22.2	26.7	25.5(1.8)	26.7(1.0)	26.8

Multipole phases (degrees) from single-energy fits to 4 (SE4) and 8 (SE8) observables at 600 MeV

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 $\pi$ N scattering

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Overall fit quality: single-energy vs global fit

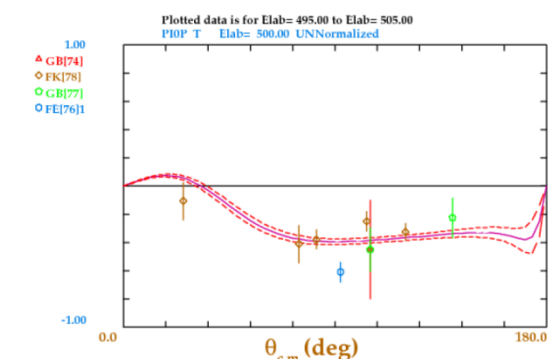
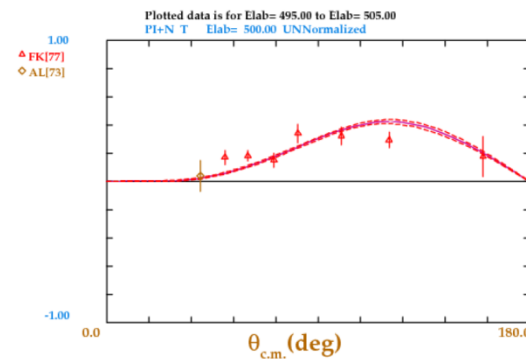
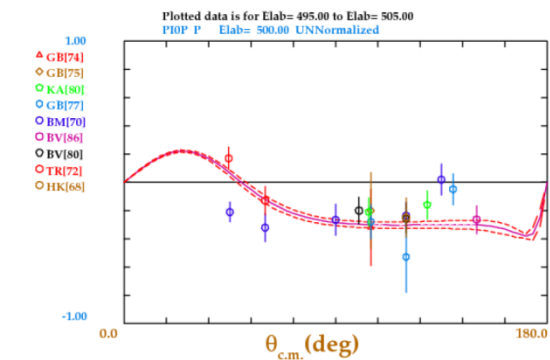
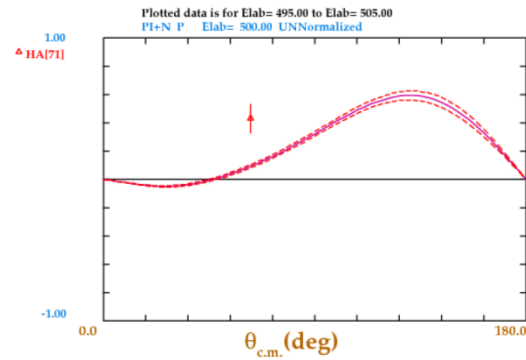
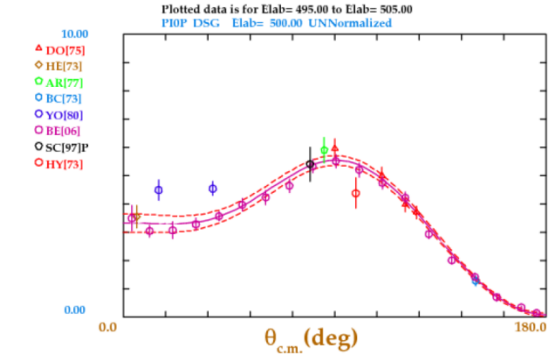
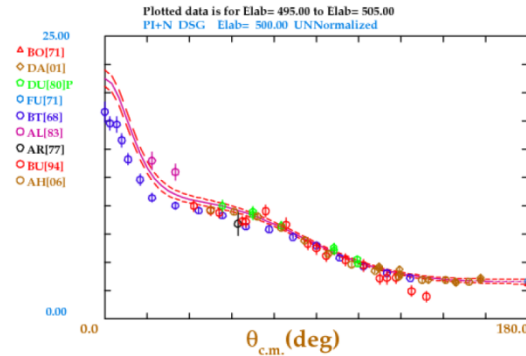
$E_\gamma$ (lab)	global $\chi^2$	single-energy $\chi^2$	data/bin
294-316	1068	957	605
494-516	710	582	377
734-757	1450	1276	778
1494-1515	226	141	113

global  $\chi^2 \sim 2$   
single-energy  $\chi^2 \sim 1.5$

# Single-energy fits

## $\pi N$ Photoproduction

## $\pi N$ scattering



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Some observations:

1. The real-world database is not self-consistent.
2. The single-energy amplitudes are not a representation of the data ( which is not possible – see 1 ).
3. Some data could be pruned without changing the result, but lowering the overall  $\chi^2$  .

Question:

If we drop `really bad' data, can we get a fit  $\chi^2$  / data near unity?



Single-energy fits  
 $\pi$ N Photoproduction  
 $\pi$ N scattering

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Simple exercise:

1. **prune** data if  $\chi^2 > 5$  for point or distribution
2. **Re-fit** remaining data
3. **Repeat** step 1
4. **Stop** if nothing left to prune



Single-energy fits  
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Starting global fit:

$\pi^0$ p	$\pi^+$ n	$\pi^-$ p	$\pi^0$ n
32751/13646	16833/8521	4749/2333	1108/364

Pruned global fit:

$\pi^0$ p	$\pi^+$ n	$\pi^-$ p	$\pi^0$ n
11716/11338	7584/7496	1978/2022	241/251

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Summary :

Pruned for  $\chi^2 > 5$

3757 pruned data contributing 33234 to  $\chi^2$

Some of the pruned data are not old

Multipoles change very little

Need a consensus evaluation of database  
( this was done for  $\pi$ N scattering )

Dangerous to toss data just because it is hard to fit.

Single-energy fits  
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 $\pi$ N scattering

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Single –energy results have been tied to global fits

Two global fit forms, discussed so far, are:

**SAID** ( original form )

$$T_{\gamma\pi} = (\text{Born} + A)(1 + iT_{\pi\pi}) + BT_{\pi\pi}$$


Automatically  
satisfies Watson's  
theorem



**MAID**

$$T_{\gamma\pi} = (\text{Born})(1 + iT_{\pi\pi}) + e^{i\phi}T_{BW}$$

Phase to satisfy  
Watson's theorem



where  $T_{\pi\pi}$  is the associated  $\pi$ N partial-wave amplitude, and  $T_{BW}$  is a Breit-Wigner resonance term.

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 $\pi N$  Photoproduction  
 $\pi N$  scattering

Motivation (simple K-matrix):

Watson's theorem

$$T_{\gamma\pi} = (1 + iT_{\pi\pi})K_{\gamma\pi}$$

Need pole, or is zero at resonance

$$K_{\gamma\pi} = \frac{A}{W - W_R} + B ; K_{\pi\pi} = \frac{C}{W - W_R} + D$$

$$T_{\gamma\pi} = \left( \frac{A}{W - W_R} + B \right) \frac{(W - W_R)}{W - W_R - i[C + (W - W_R)D]}$$

One term  $\rightarrow$  zero at  $W_R$   
 Other term has BW behavior

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Expanding to include another hadronic channel,

$$K = \begin{pmatrix} K_{\gamma\gamma} & K_{\gamma\pi} & K_{\gamma\Delta} \\ K_{\gamma\pi} & K_{\pi\pi} & K_{\pi\Delta} \\ K_{\gamma\Delta} & K_{\pi\Delta} & K_{\Delta\Delta} \end{pmatrix}$$

The relation  $T(1-iK)=K$  gives,

$$T_{\gamma\pi}(1 - iK_{\gamma\gamma}) = (1 + iT_{\pi\pi})K_{\gamma\pi} + iK_{\gamma\Delta}T_{\pi\Delta}$$

Drop  $K_{\gamma\gamma}$   
as small



which can be massaged into the form,

$$T_{\gamma\pi} = (1 + iT_{\pi\pi}) \left( K_{\gamma\pi} - \frac{K_{\gamma\Delta}K_{\pi\pi}}{K_{\pi\Delta}} \right) + \frac{K_{\gamma\Delta}}{K_{\pi\Delta}}T_{\pi\pi}$$

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CGLN  
Schwela & Weizel,  
Z. Phys 221, 71 (1969)

Aside:

The form of the  $M_{1+}^{3/2}$  multipole amplitude from CGLN in the static limit ( neglect  $O(1/M)$  terms ) is also

$$\left[ A ( 1 + i T_{\pi\pi} ) + B T_{\pi\pi} \right]$$

with the factors  $A$  and  $B$  fixed.

See, for example,  
S.N. Yang,  
J. Phys. G 11, L205 (1985).  
L. Resnick,  
Phys. Rev. D2, 1975 (1972).

More elaborate methods, dynamical models, use of the Muskhelishvili-Omnes result will give first term, but the second is generally replaced by a PV integral.

Single-energy fits  
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**Bonn-Gatchina** approach (multi-channel)

This 'P-matrix' approach is similar – but the database includes other reactions.

$$A_a = \hat{P}_b (\hat{I} - i\hat{\rho}\hat{K})_{ba}^{-1}$$

Photoproduction  
amplitude

$$P_b = \sum_{\alpha} \frac{g_{\gamma N}^{(\alpha)} g_b^{(\alpha)}}{M_{\alpha}^2 - s} + \tilde{f}_b$$

P-vector

$$K_{ab} = \sum_{\alpha} \frac{g_a^{(\alpha)} g_b^{(\alpha)}}{M_{\alpha}^2 - s} + f_{ab}$$

Hadronic K-matrix



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A modified SAID approach, which incorporates more info from the  $\pi$ N elastic scattering analysis, is based on a Chew-Mandelstam K-matrix approach

$$T_{\alpha\beta} = \sum_{\sigma} [1 - \overline{KC}]_{\alpha\sigma}^{-1} \overline{K}_{\sigma\beta}$$

← From  $\pi$ N scattering

leading to

$$T_{\alpha\gamma} = \sum_{\sigma} [1 - \overline{KC}]_{\alpha\sigma}^{-1} \overline{K}_{\sigma\gamma}$$

← Fitted to photo data

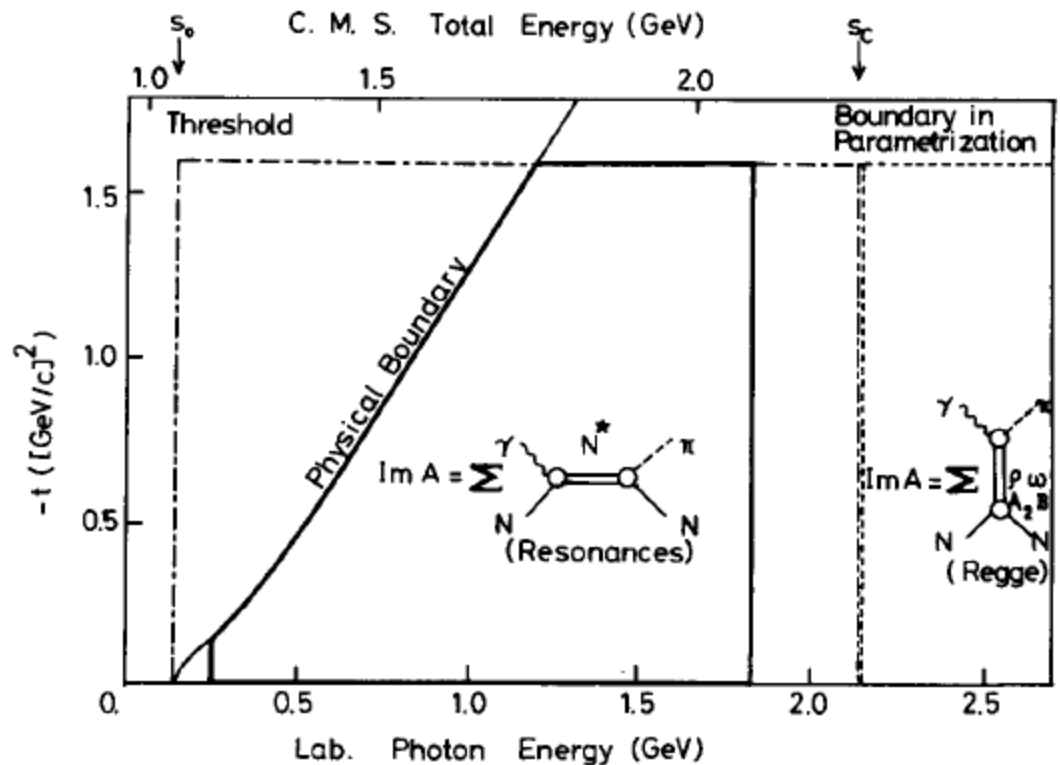
$$C = \text{Re } C + i \rho$$

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Arai/Fujii  
 Crawford  
 Aznauryan

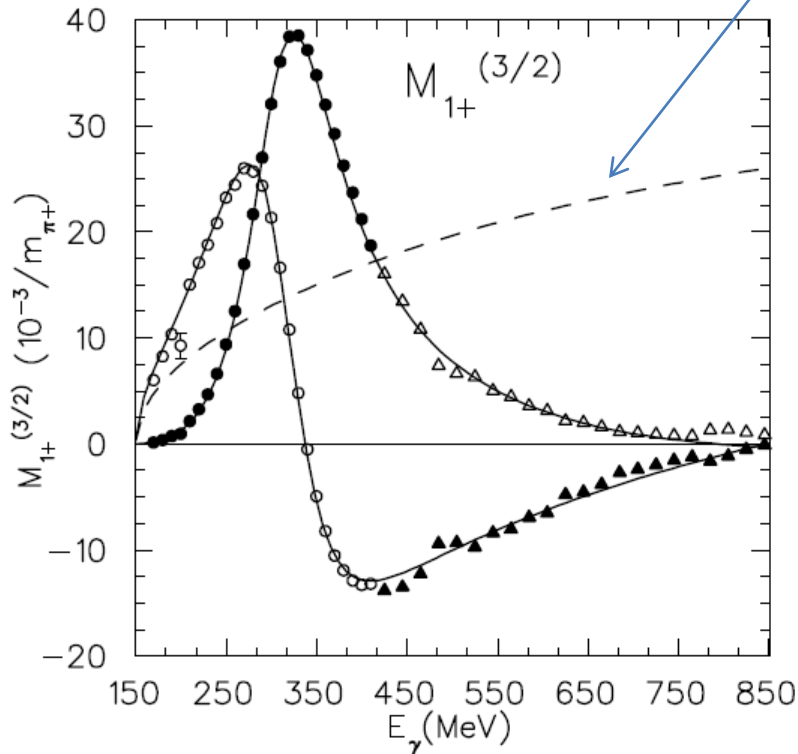
$$\text{Re} A_{i,\text{ch}}(s, t) = R_{i,\text{ch}}(s, t) + \frac{1}{\pi} \times P \int_{(m+M)^2}^{\infty} ds' \left\{ \frac{\text{Im} A_{i,\text{ch}}(s', t)}{(s' - s)} + \xi_i \frac{\text{Im} A_{i,\text{ch}}(s', t)}{(s' - u)} \right\}$$

Parameterize  $\text{Im} A$   
 [A/F: 3-ch K-matrix]  
 calc.  $\text{Re} A$

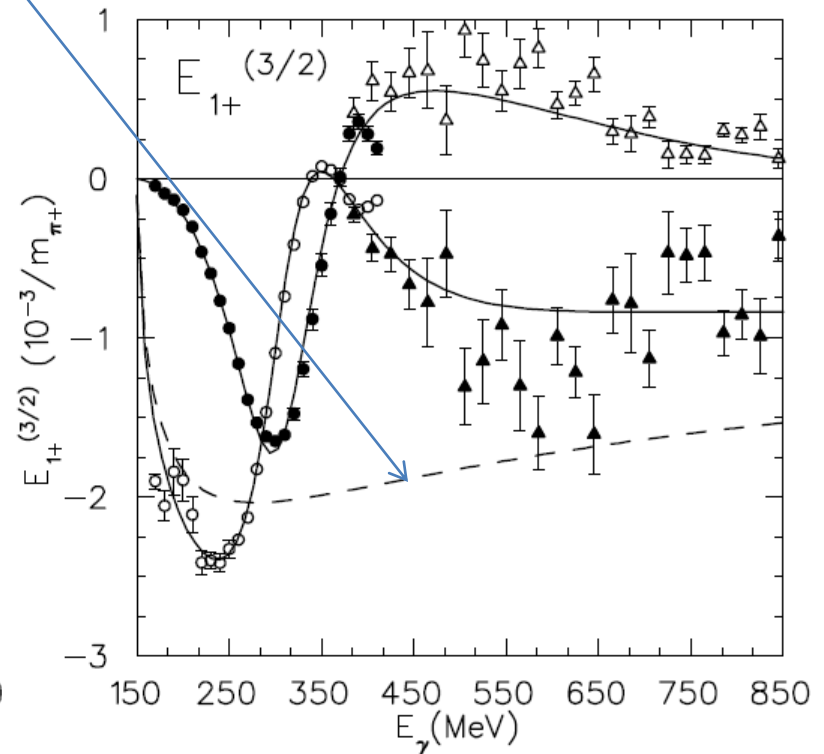


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Born terms



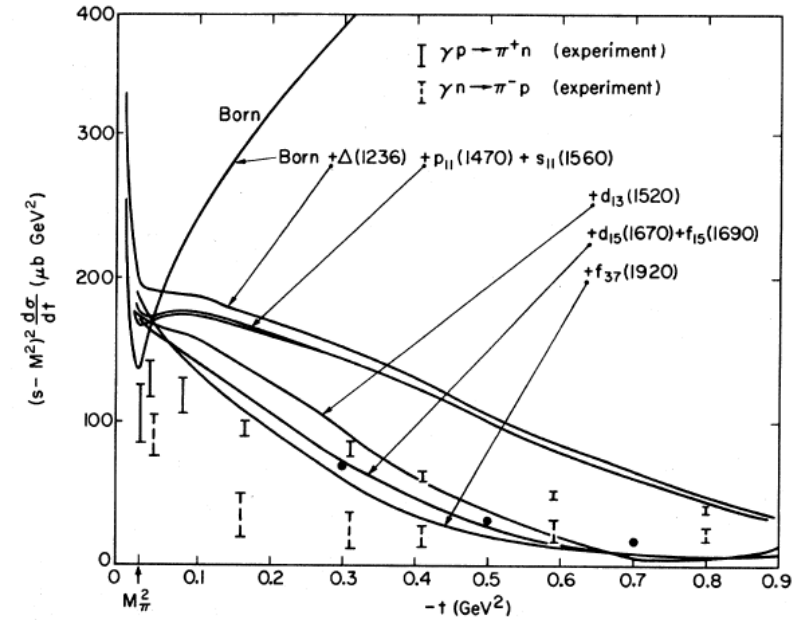
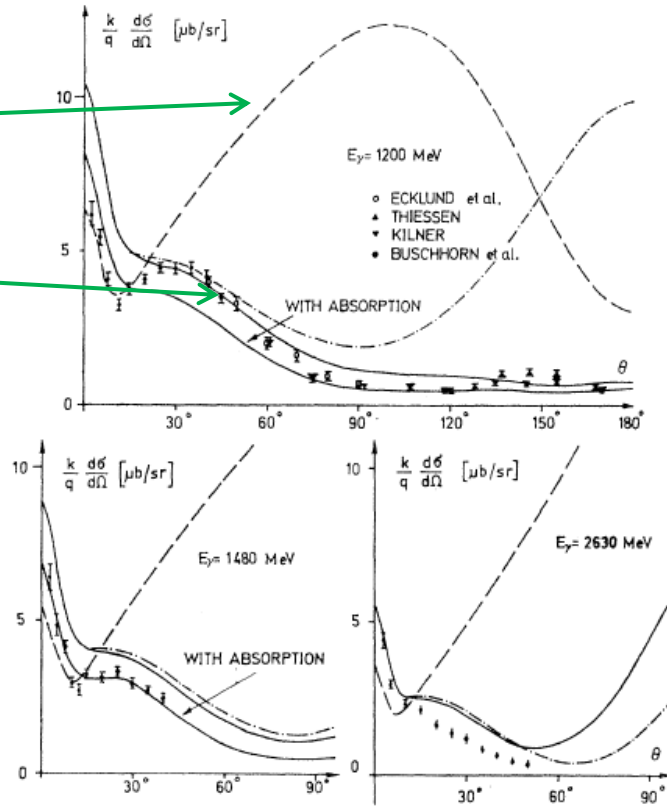
Above 500 MeV  
 Born > largest multipole



Born\*( $1 + iT_{\pi\pi}$ ) gives  
 almost everything here

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Born  $\rightarrow$   
 add  $\Delta(1232)$   
 to RHS of DR



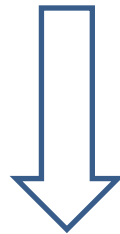
16 GeV  
 Barbour, Malone, Moorhouse,  
 PRD4,1521(1971)

Engels, Schwiderski, Schmidt,  
 PR166,1343(1968).

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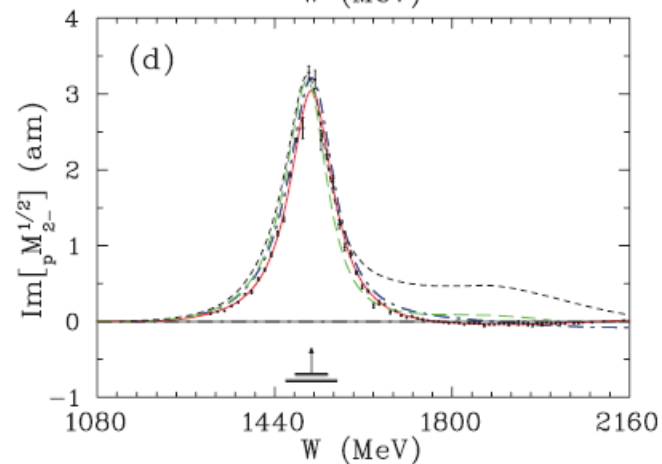
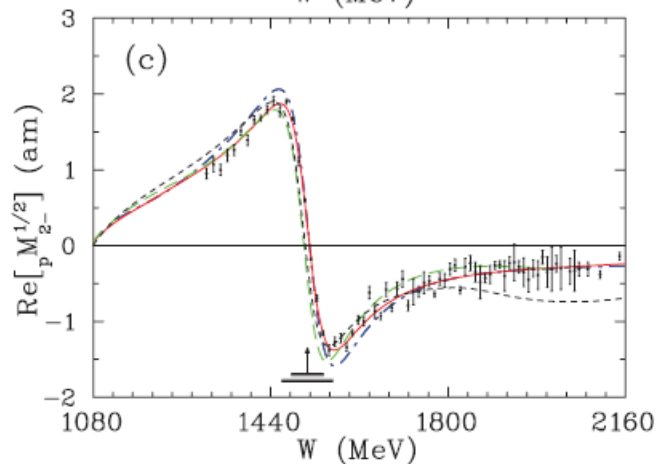
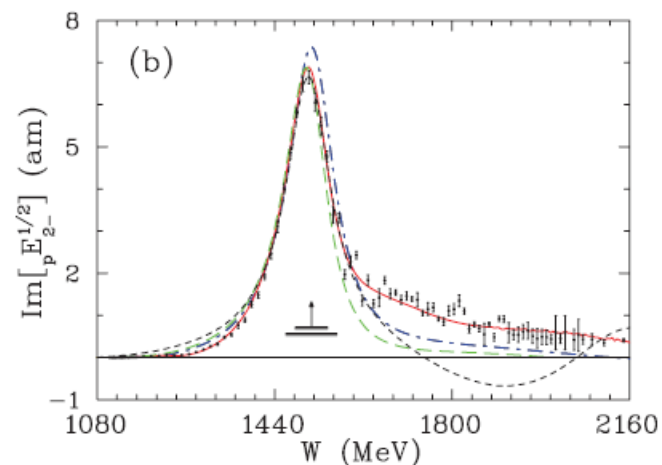
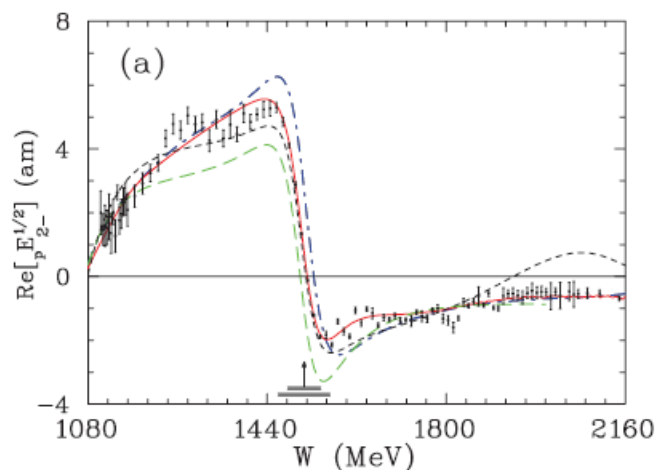
How well do the  
SAID/MAID/BoGa  
analyses agree on the multipoles?



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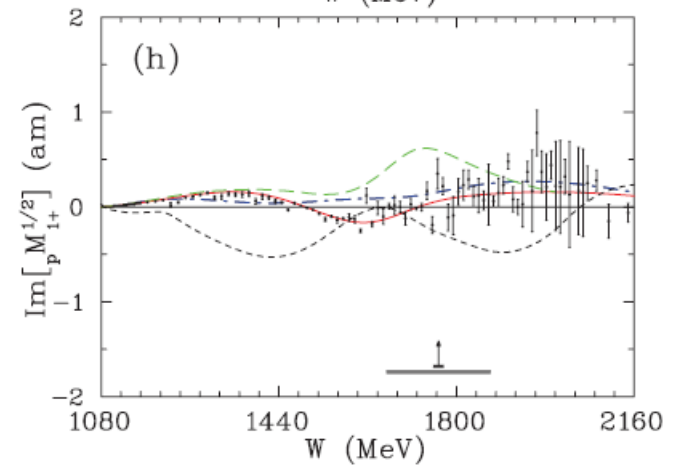
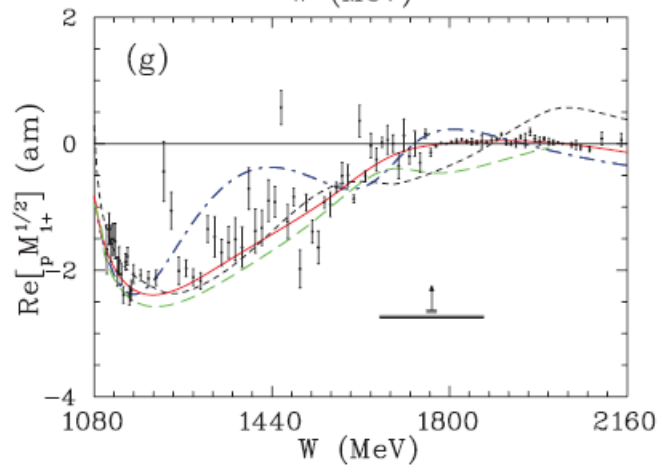
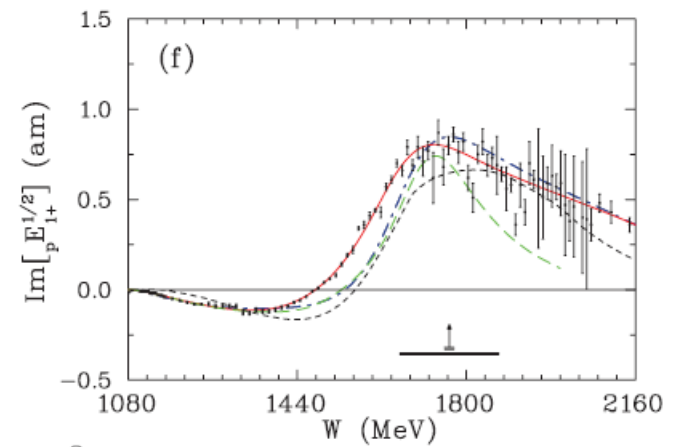
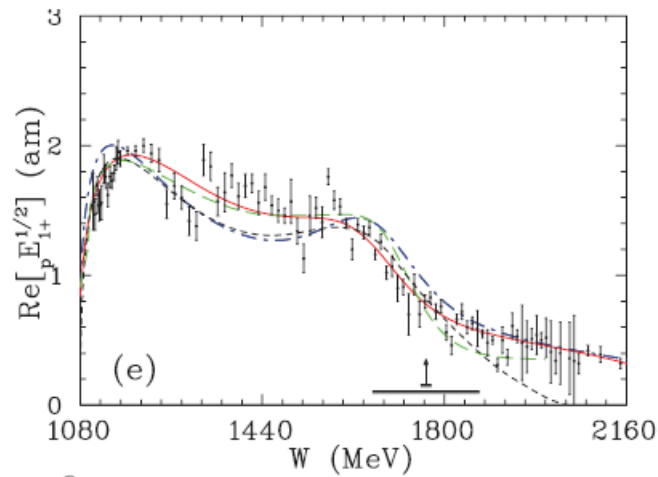
$D_{13}(1520)$  [  $N(1520)3/2^-$  ]

SAID  $CM$  ———  
MAID - - - -  
BoGa - - - -



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 $\pi N$  Photoproduction  
 $\pi N$  scattering

$P_{13}(1720) [N(1720)3/2^+]$



Single-energy fits  
 $\pi N$  Photoproduction  
 $\pi N$  scattering

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The Baryon Summary Table gives  $N\gamma$  branching fractions for those resonances whose couplings are considered to be reasonably well established. The  $N\gamma$  partial width  $\Gamma_\gamma$  is given in terms of the helicity amplitudes  $A_{1/2}$  and  $A_{3/2}$  by

$$\Gamma_\gamma = \frac{k^2}{\pi} \frac{2M_N}{(2J+1)M_R} [ |A_{1/2}|^2 + |A_{3/2}|^2 ] .$$

Here  $M_N$  and  $M_R$  are the nucleon and resonance masses,  $J$  is the resonance spin, and  $k$  is the photon c.m. decay momentum.

Must construct  $A_{1/2}$  ,  $A_{3/2}$   
at the 'resonance' position  
i.e. BW parameterization



Single-energy fits  
 $\pi$ N Photoproduction  
 $\pi$ N scattering

$$A_{l+}^{1/2} = -\frac{1}{2}[(l+2)\bar{E}_{l+} + l\bar{M}_{l+}] ,$$

$$A_{l+}^{3/2} = \frac{1}{2}\sqrt{l(l+2)}[\bar{E}_{l+} - \bar{M}_{l+}] ,$$

$$A_{(l+1)-}^{1/2} = -\frac{1}{2}[l\bar{E}_{(l+1)-} - (l+2)\bar{M}_{(l+1)-}] ,$$

$$A_{(l+1)-}^{3/2} = -\frac{1}{2}\sqrt{l(l+2)}[\bar{E}_{(l+1)-} + \bar{M}_{(l+1)-}] ,$$

Resonance  
 contribution  
 to multipole

Resonance part straightforward in  
 MAID ( BW form)  
 BoGa (K-matrix poles)  
 but not in SAID ( resonances hidden in  $T_{\pi\pi}$  )

In practice, SAID results extracted assuming  
 a MAID-like background-resonance separation