

Flavor Asymmetry of Proton Sea

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**Exploring Hadrons with Electromagnetic
Probes: Structure, Excitations, Interactions**

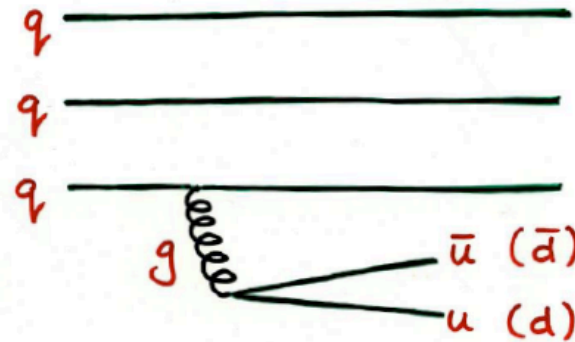
JLab, November 2, 2017

Outline

- Motivation with JLab TDIS Experiment
- Chiral Effective Theory Connection with QCD
- Pion Splitting Functions
- Leading Neutron Electroproduction
- Outlook with Drell-Yan Process

Flavor asymmetry

- Antiquarks in the proton “sea” produced predominantly by gluon radiation into quark-antiquark pairs, $g \rightarrow q\bar{q}$



→ since u and d quark masses are similar, expect flavor-symmetric sea, $\bar{d} \approx \bar{u}$

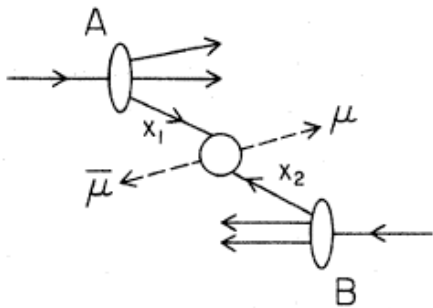
- Experimentally, one finds *large excess of \bar{d} over \bar{u}*

$$\int_0^1 dx (\bar{d}(x) - \bar{u}(x)) = 0.118 \pm 0.012$$

Flavor asymmetry

- Large flavor asymmetry in proton sea suggests important role of chiral symmetry in high-energy reactions

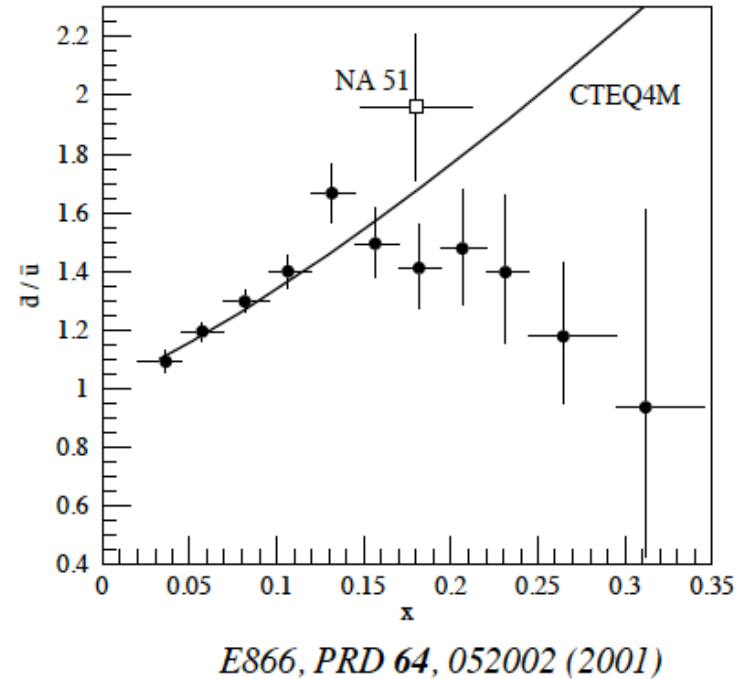
→ Drell-Yan process



$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{9Q^2} \sum_q e_q^2 (q(x_b)\bar{q}(x_t) + \bar{q}(x_b)q(x_t))$$

→ for $x_b \gg x_t$

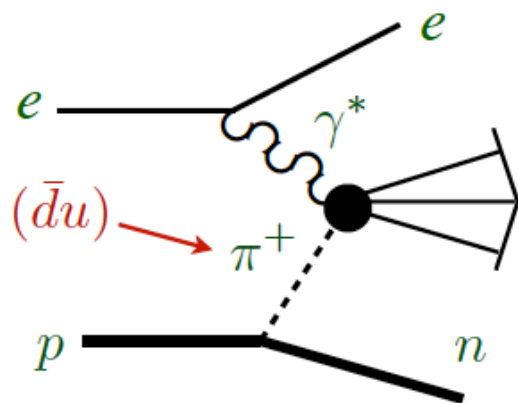
$$\frac{\sigma^{pd}}{2\sigma^{pp}} \approx \frac{1}{2} \left(1 + \frac{\bar{d}(x_t)}{\bar{u}(x_t)} \right) \quad \rightarrow \quad \int_0^1 dx (\bar{d} - \bar{u}) = 0.118 \pm 0.012$$



Flavor asymmetry

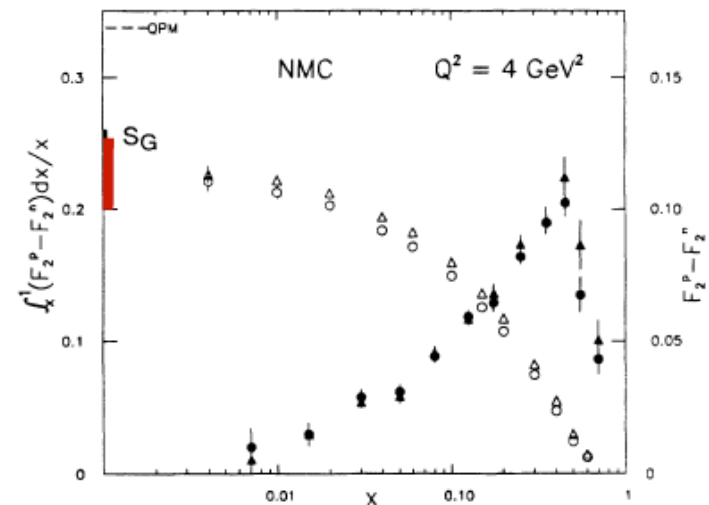
- Large flavor asymmetry in proton sea suggests important role of chiral symmetry in high-energy reactions

→ Sullivan process in DIS



Sullivan, PRD 5, 1732 (1972)

$$\bar{d} > \bar{u}$$



$$\int_0^1 \frac{dx}{x} (F_2^p - F_2^n) = \frac{1}{3} - \frac{2}{3} \int_0^1 dx (\bar{d} - \bar{u})$$

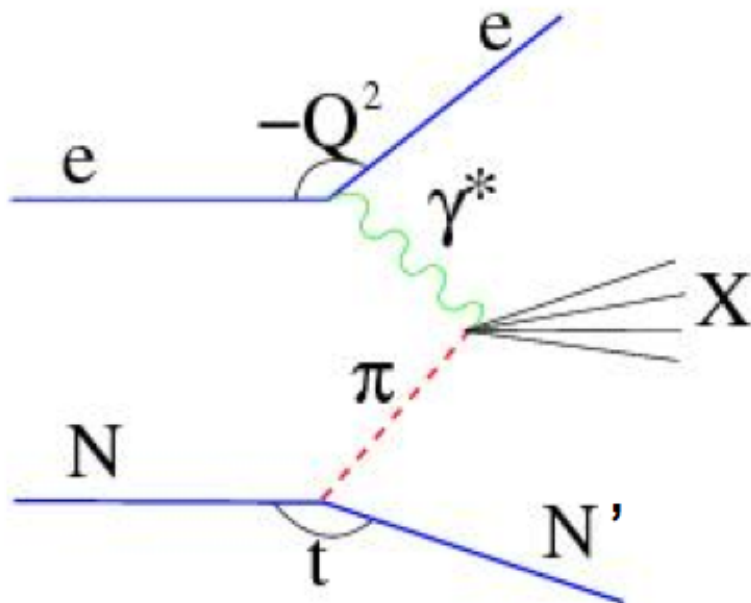
$$= 0.235(26)$$

NMC, PRD 50, 1 (1994)

Measurement of Tagged Deep Inelastic Scattering (TDIS)

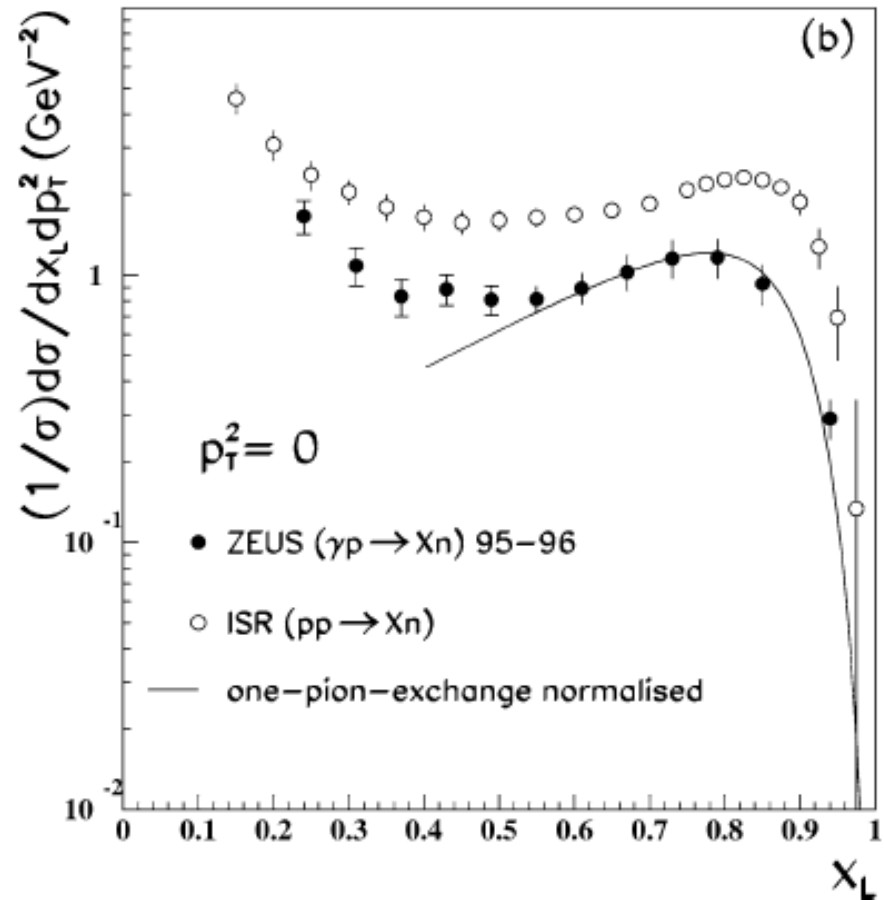
C.Keppel (Contact person)

To probe the elusive mesonic content of the nucleon



$$e + p(\text{or } n) \rightarrow e' + p + X$$

$$e + D \rightarrow e' + p + p + X$$



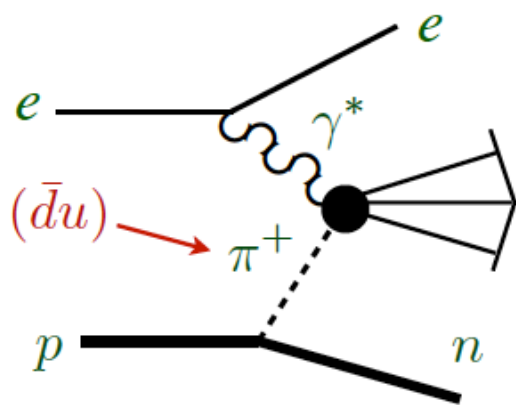
Leading neutron production in e^+p collisions at HERA

ZEUS Collaboration, NPB 637 (2002) 3–56

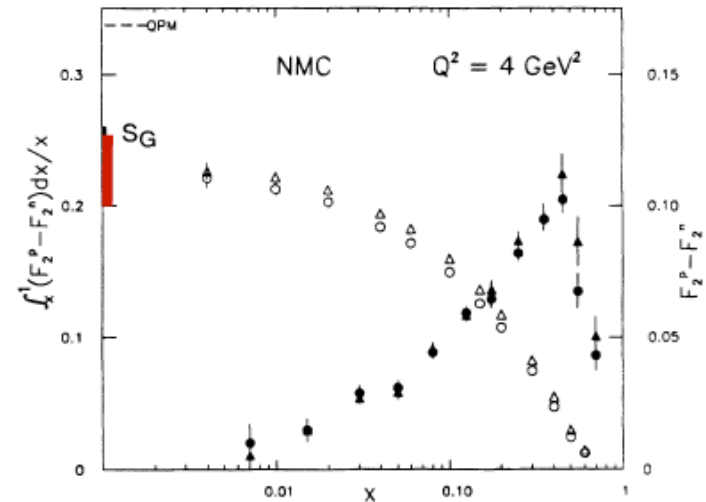
Flavor asymmetry

- Large flavor asymmetry in proton sea suggests important role of chiral symmetry in high-energy reactions

→ Sullivan process in DIS



Sullivan, *PRD* **5**, 1732 (1972)
 Thomas, *PLB* **126**, 97 (1983)
 Miller, Kumano, Strikman, Weiss, ...



$$(\bar{d} - \bar{u})(x) = \frac{2}{3} \int_x^1 \frac{dy}{y} f_\pi(y) \bar{q}^\pi(x/y)$$

pion light-cone momentum distribution in nucleon

$$f_\pi(y) = \frac{3g_{\pi NN}^2}{16\pi^2} y \int dt \frac{-t \mathcal{F}_{\pi NN}^2(t)}{(t - m_\pi^2)^2}$$

connection with QCD?

Connection with QCD

$$\blacksquare (\bar{d} - \bar{u})(x) = \frac{2}{3} \int_x^1 \frac{dy}{y} f_\pi(y) \bar{q}^\pi(x/y) \quad \boxed{f_\pi(y) = \frac{3g_{\pi NN}^2}{16\pi^2} y \int dt \frac{-t \mathcal{F}_{\pi NN}^2(t)}{(t - m_\pi^2)^2}}$$

→ *model-independent* leading nonanalytic (LNA) behavior consistent with Chiral Symmetry of QCD.

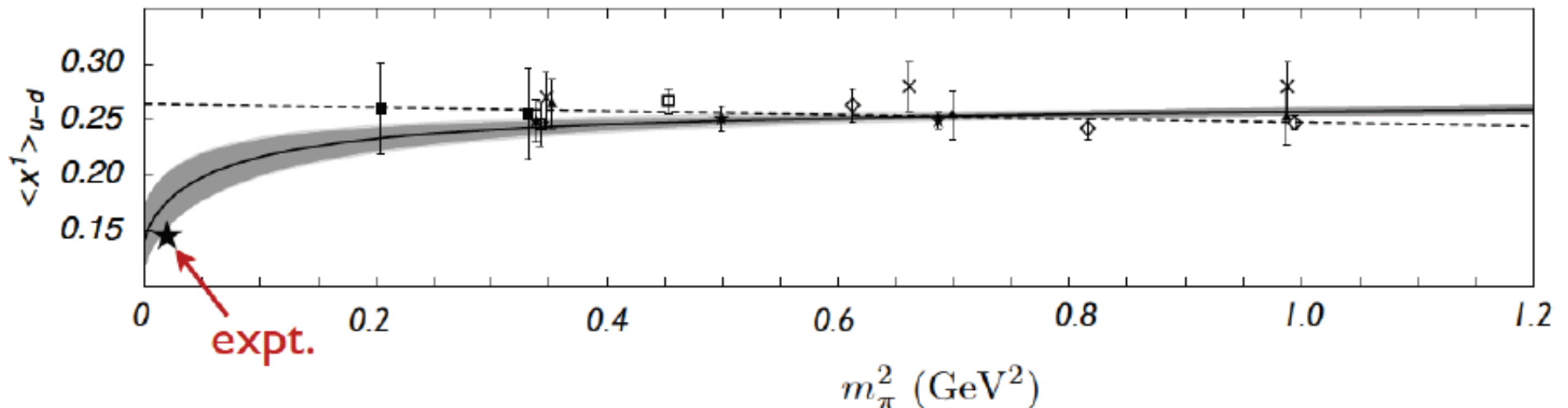
$$\langle x^0 \rangle_{\bar{d}-\bar{u}} \equiv \int_0^1 dx (\bar{d} - \bar{u})$$

$$\boxed{m_\pi^2 f_\pi^2 = -2m_q \langle \bar{q}q \rangle}$$

$$= \frac{2}{3} \int_0^1 dy f_\pi(y) = \frac{2g_A^2}{(4\pi f_\pi)^2} m_\pi^2 \log(m_\pi^2/\mu^2) + \text{analytic terms}$$

- Nonanalytic behavior vital for chiral extrapolation of lattice data

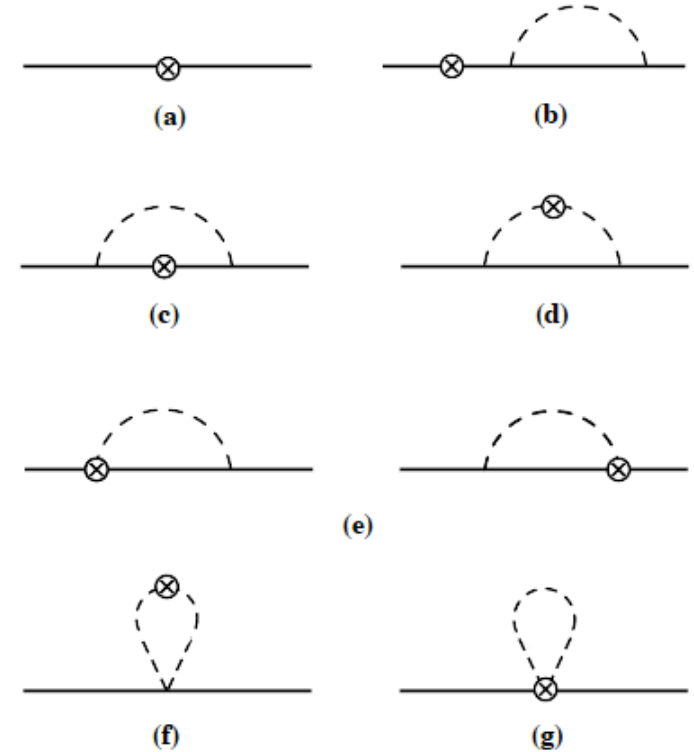
Thomas, Melnitchouk, Steffens PRL 85, 2892 (2000)



Chiral effective theory

■ Pion cloud corrections to electromagnetic N coupling

→ N rainbow (c), π rainbow (d),
Kroll-Ruderman (e),
 π bubble (f), π tadpole (g)



■ Vertex renormalization

$$(Z_1^{-1} - 1) \bar{u}(p) \gamma^\mu u(p) = \bar{u}(p) \Lambda^\mu u(p)$$

→ taking “+” components: $Z_1^{-1} - 1 \approx 1 - Z_1 = \frac{M}{p^+} \bar{u}(p) \Lambda^+ u(p)$

→ e.g. for N rainbow contribution,

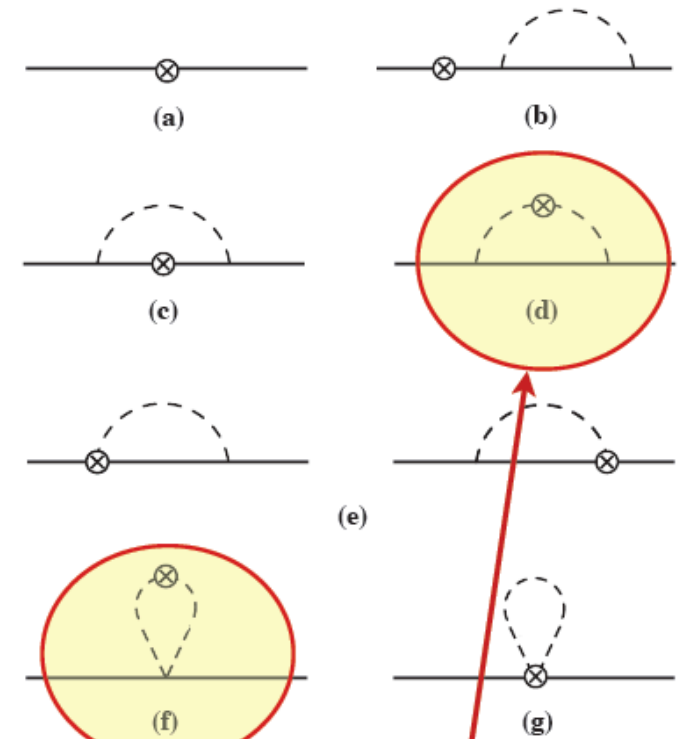
$$\Lambda_\mu^N = -\frac{\partial \hat{\Sigma}}{\partial p^\mu}$$

C.Ji, W. Melnitchouk, A.Thomas,
PRD88,076005(2013)

Flavor asymmetry

■ Pion cloud corrections to electromagnetic N coupling

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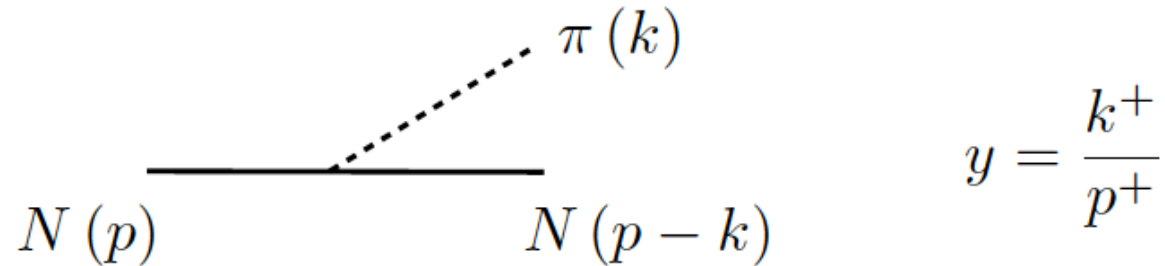
- e.g. for N rainbow contribution,

$$\Lambda_\mu^N = -\frac{\partial \hat{\Sigma}}{\partial p^\mu}$$

contribute to $\bar{d} - \bar{u}$

Pion splitting functions

- Each diagram can be represented by $N \rightarrow N\pi$
“splitting function” $f_i(y)$ (light-cone momentum distribution function)



- Vertex renormalization is k^+ moment of $f_i(y)$

$$1 - Z_1^i = \int dy f_i(y)$$

Pion splitting functions

■ Summary of splitting functions:

$$1 - Z_1^i = \int dy f_i(y)$$

where

$$f_\pi(y) = f^{(\text{on})}(y) + f^{(\delta)}(y)$$
$$f_N(y) = f^{(\text{on})}(y) + f^{(\text{off})}(y) - f^{(\delta)}(y)$$
$$f_{\text{KR}}(y) = f^{(\text{off})}(y) - 2f^{(\delta)}(y)$$
$$f_{\text{tad}}(y) = -f_{\text{bub}}(y) = 2f^{(\text{tad})}(y)$$

with components

$$f^{(\text{on})}(y) = \frac{g_A^2 M^2}{(4\pi f_\pi)^2} \int dk_\perp^2 \frac{y(k_\perp^2 + y^2 M^2)}{[k_\perp^2 + y^2 M^2 + (1-y)m_\pi^2]^2}$$

$$f^{(\text{off})}(y) = \frac{g_A^2 M^2}{(4\pi f_\pi)^2} \int dk_\perp^2 \frac{y}{k_\perp^2 + y^2 M^2 + (1-y)m_\pi^2}$$

$$f^{(\delta)}(y) = \frac{g_A^2}{4(4\pi f_\pi)^2} \int dk_\perp^2 \log\left(\frac{k_\perp^2 + m_\pi^2}{\mu^2}\right) \delta(y)$$

tadpole & bubble
equal & opposite

$$(1 - Z_1^{\text{tad}}) = -(1 - Z_1^{\text{bub}})$$

$$f^{(\text{tad})}(y) = -\frac{4}{g_A^2} f^{(\delta)}(y)$$

UV regularization

- For point-like nucleons and pions, integrals divergent
- Finite size of nucleon provides natural scale to regularize integrals, but does not prescribe form of regularization
 - freedom in choosing regularization prescription (long-distance physics independent of choice!)

$$\mathcal{F} = \Theta(\Lambda^2 - k_{\perp}^2)$$

k_{\perp} cutoff

$$\mathcal{F} = \left(\frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - t} \right)$$

monopole in $t \equiv k^2 = -\frac{k_{\perp}^2 + y^2 M^2}{1 - y}$

$$\mathcal{F} = \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right]$$

exponential in t

$$\mathcal{F} = \exp \left[(M^2 - s) / \Lambda^2 \right]$$

exponential in $s = \frac{k_{\perp}^2 + m_{\pi}^2}{y} + \frac{k_{\perp}^2 + M^2}{1 - y}$

$$\mathcal{F} = \left[1 - \frac{(t - m_{\pi}^2)^2}{(t - \Lambda^2)^2} \right]^{1/2}$$

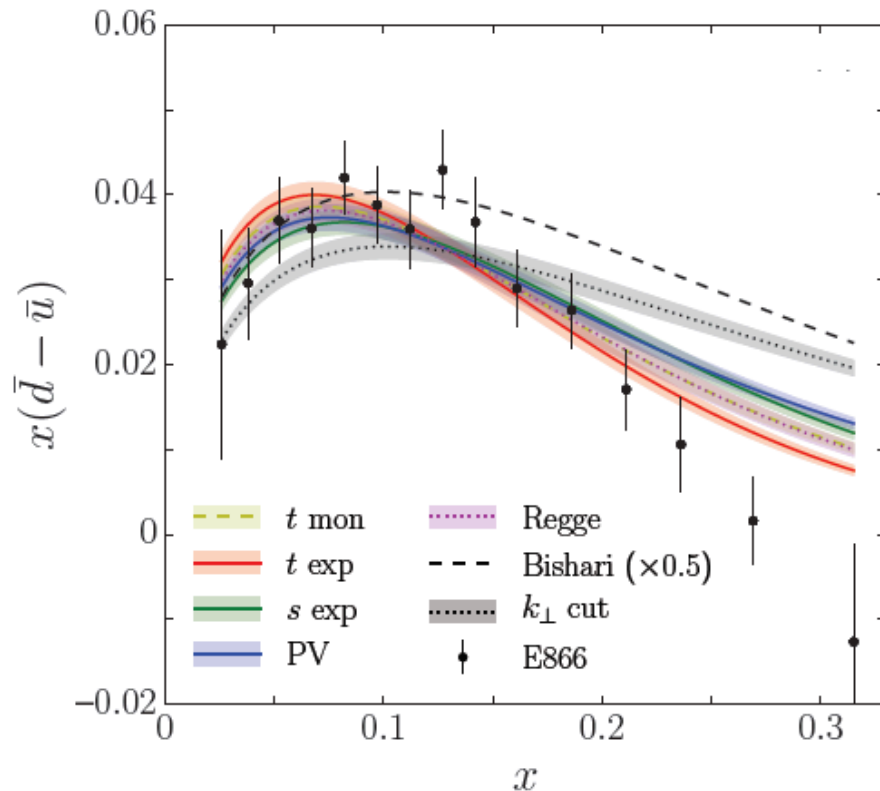
Pauli-Villars

$$\mathcal{F} = y^{-\alpha_{\pi}(t)} \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right]$$

Regge

Flavor asymmetry

- E866 $\bar{d} - \bar{u}$ data can be fitted with range of regulators



average pion “multiplicity”

$$\langle n \rangle_{\pi N} = 3 \int_0^1 dy f_N^{(\text{on})}(y) \\ \sim 0.25 - 0.3$$

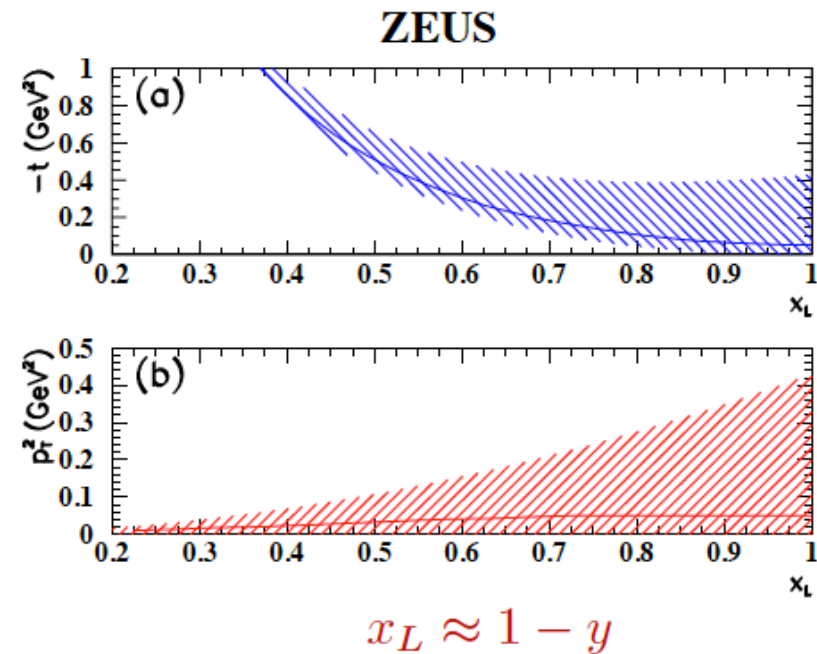
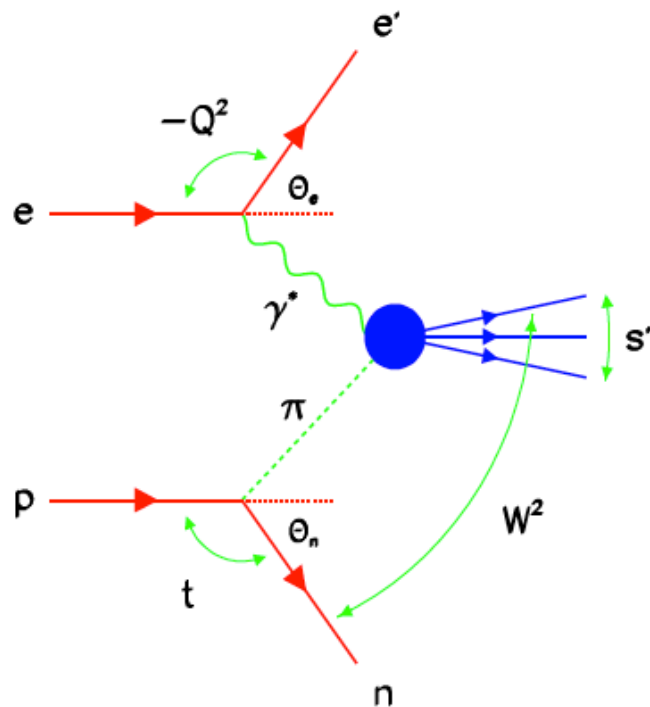
- with exception of k_{\perp} cutoff and Bishari models, all others give reasonable fits, $\chi^2 \lesssim 1.5$
- large- x asymmetry to be probed by FNAL *SeaQuest* expt.

Flavor asymmetry

- E866 $\bar{d} - \bar{u}$ data can be fitted with range of regulators
- Is pion cloud the only explanation for the asymmetry?
 - are there other data that can discriminate between different mechanisms?
 - semi-inclusive production of “leading neutrons” (LN) at HERA!
 - pi-nucleon Drell-Yan process from FNAL (E615) and CERN (NA10) on tungsten target

Leading neutron production at HERA

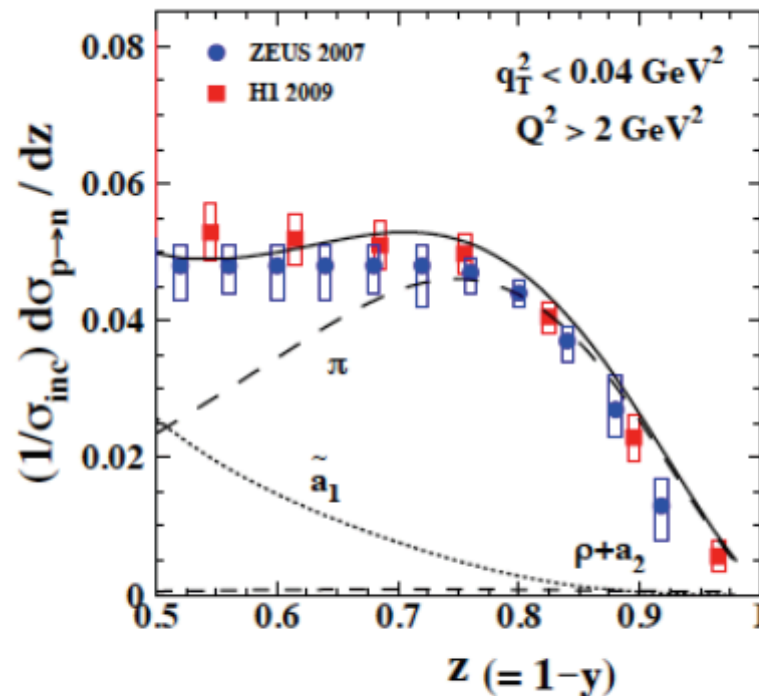
- ZEUS & H1 collaborations measured spectra of neutrons produced at very forward angles, $\theta_n < 0.8$ mrad



- can data be described within same framework as E866 flavor asymmetry?
- simultaneous fit never previously been performed!

Leading neutron production at HERA

- At large y non-pionic mechanisms contribute (e.g. heavier mesons, absorption)

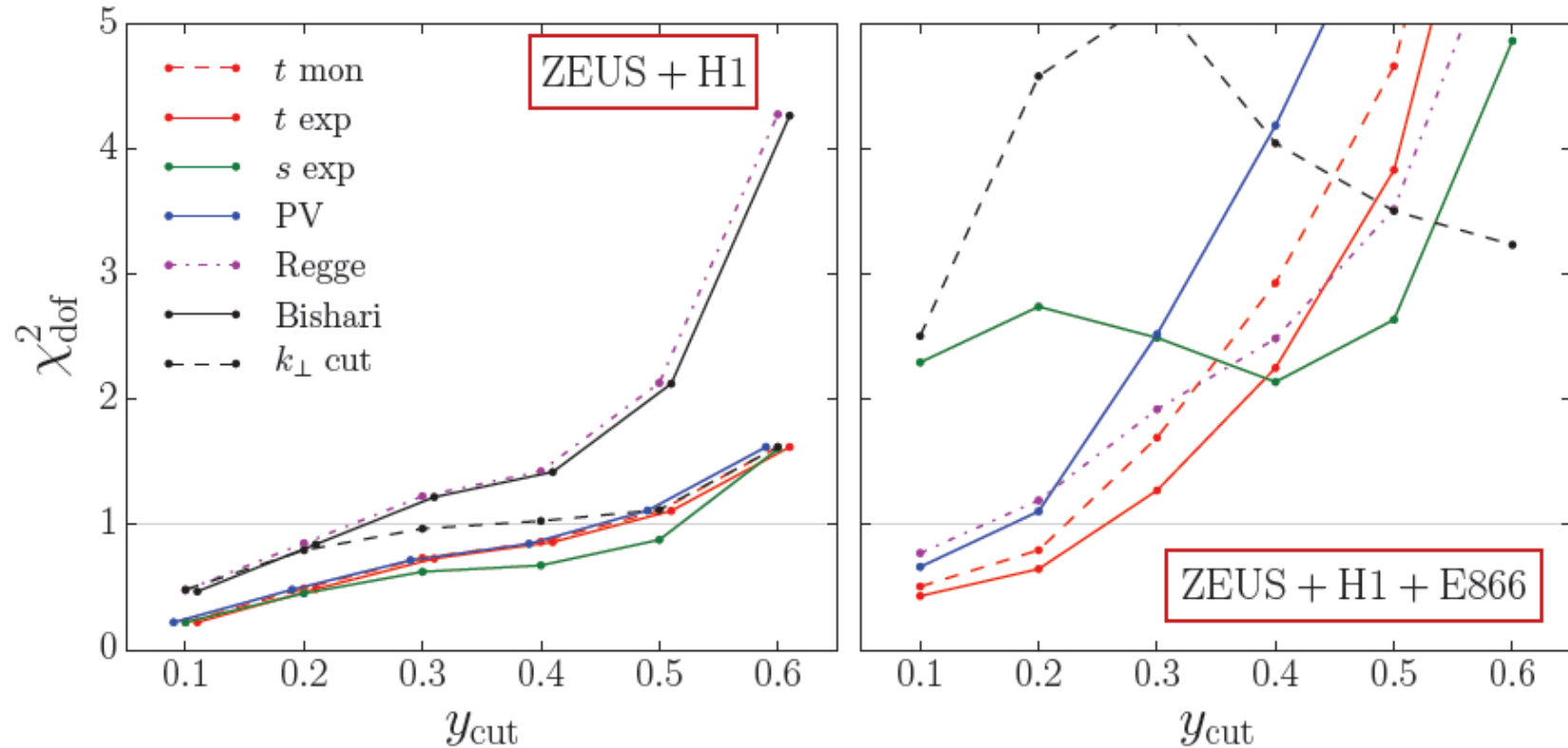


Kopeliovich et al., PRD 85, 114025 (2012)

- To reduce model dependence, fit the value of y_{cut} up to which data can be described in terms of π exchange

Leading neutron production at HERA

■ Combined fit to HERA LN and E866 Drell-Yan data

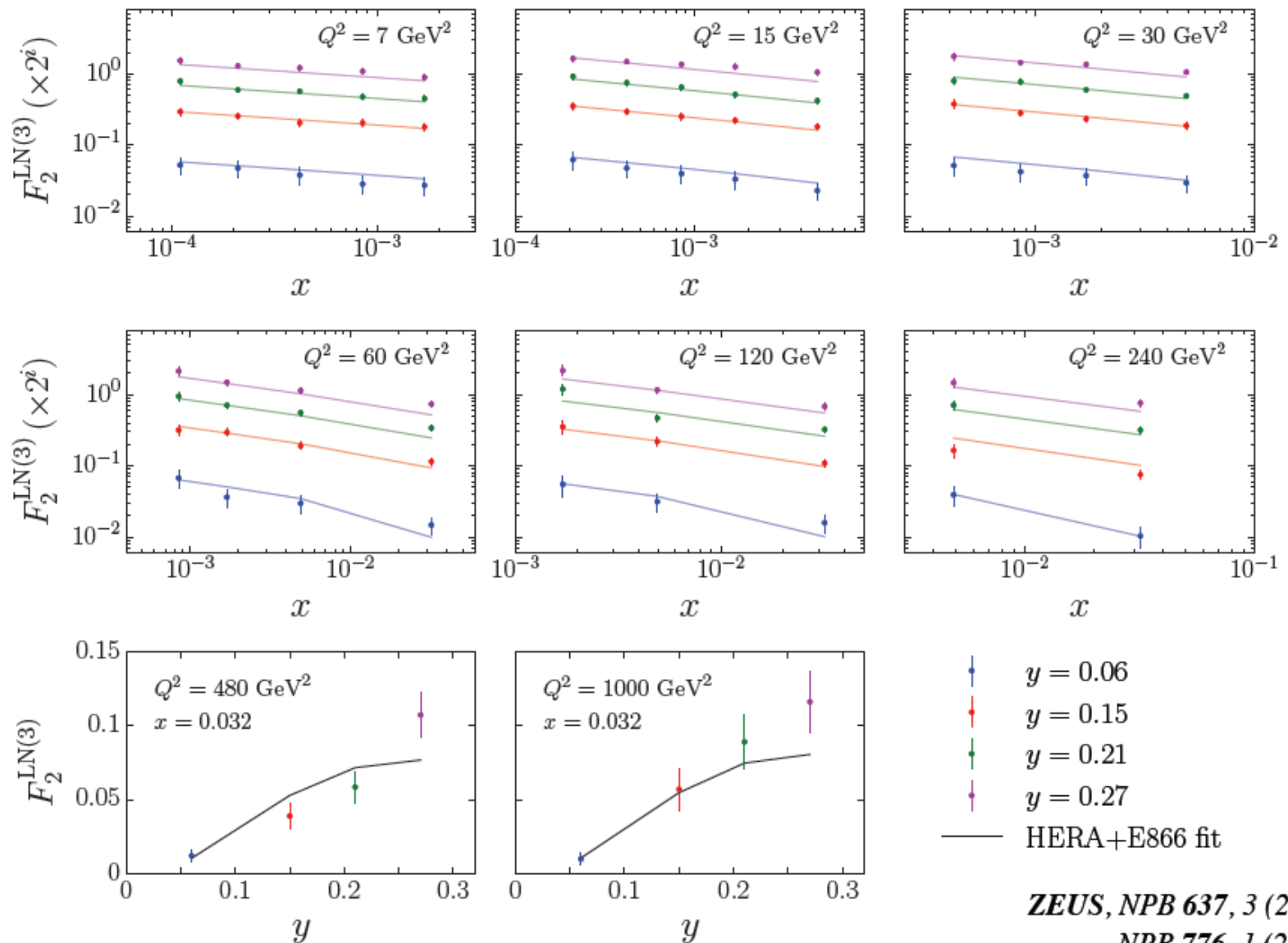


J. McKenney, N. Sato, W. Melnitchouk, C.Ji, PRD93, 054011 (2016)

➔ best fits for largest number of points afforded by t -dependent exponential (and t monopole) regulators

Leading neutron production at HERA

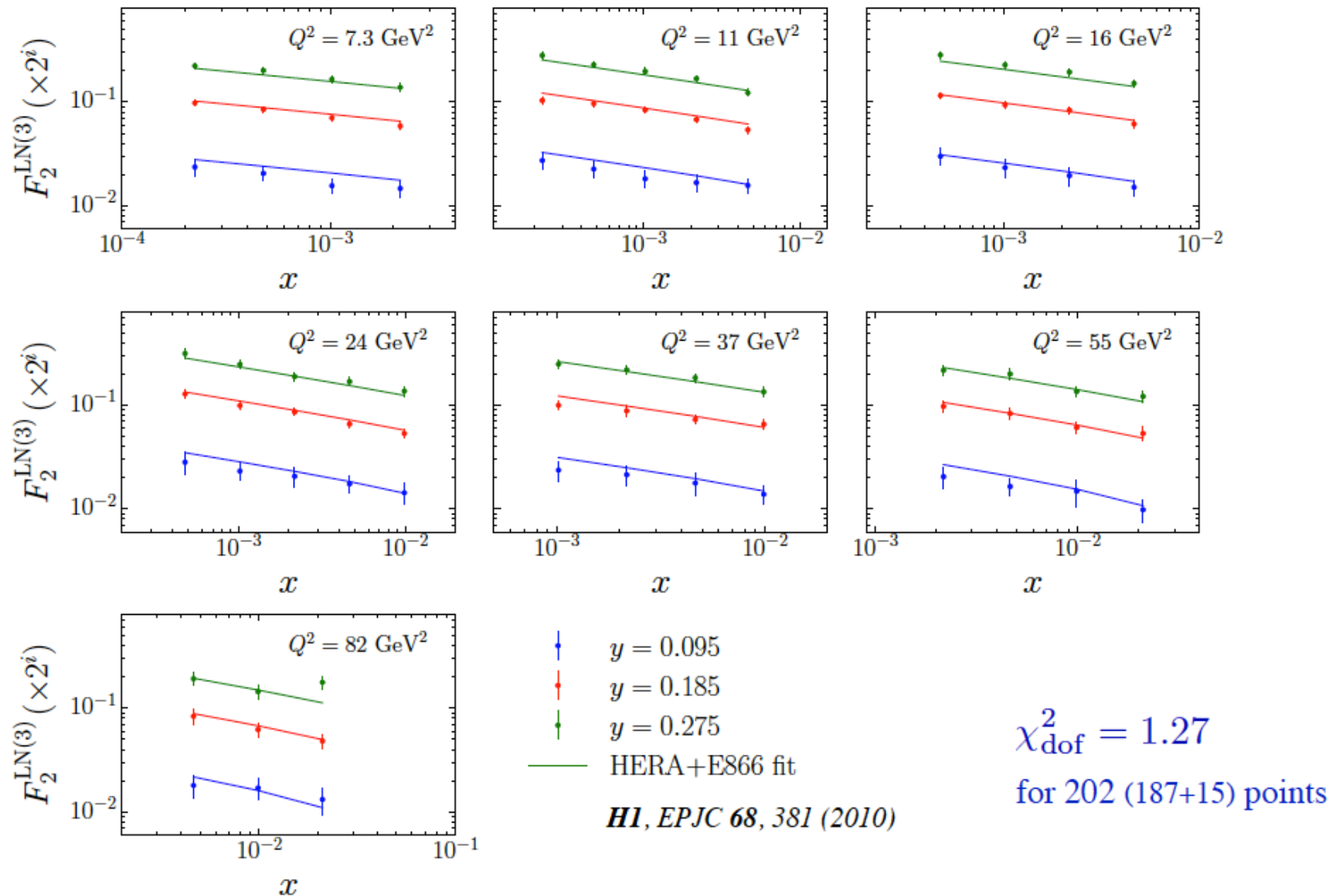
- Fit to ZEUS LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



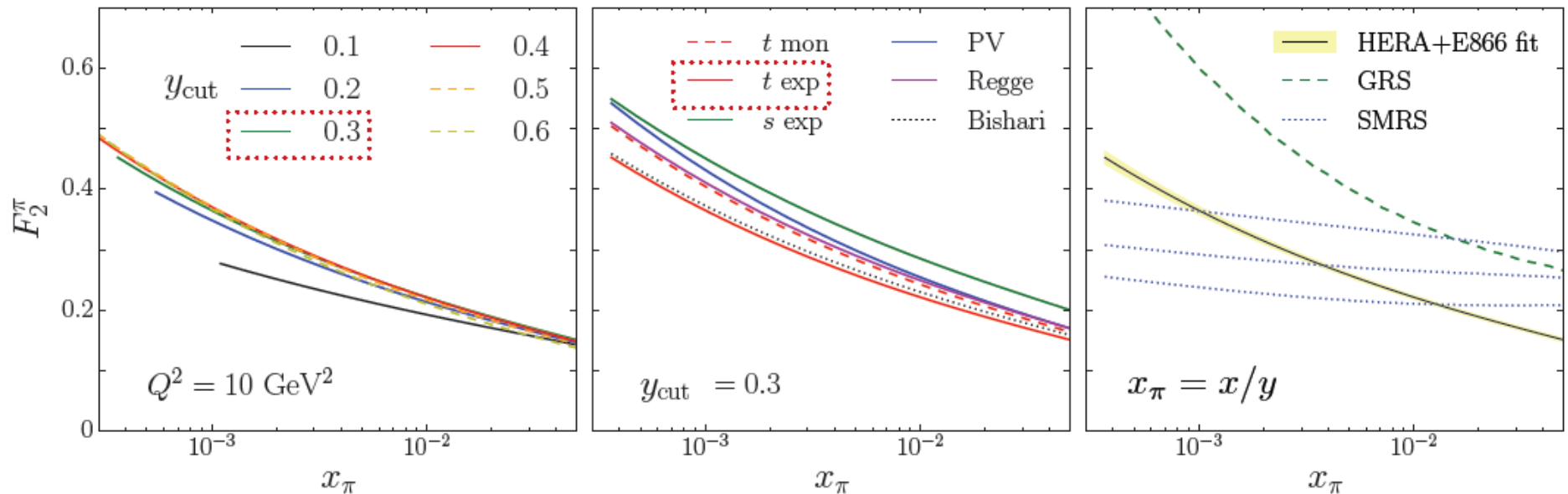
ZEUS, NPB 637, 3 (2002)
NPB 776, 1 (2007)

Leading neutron production at HERA

- Fit to H1 LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



Extracted pion structure function

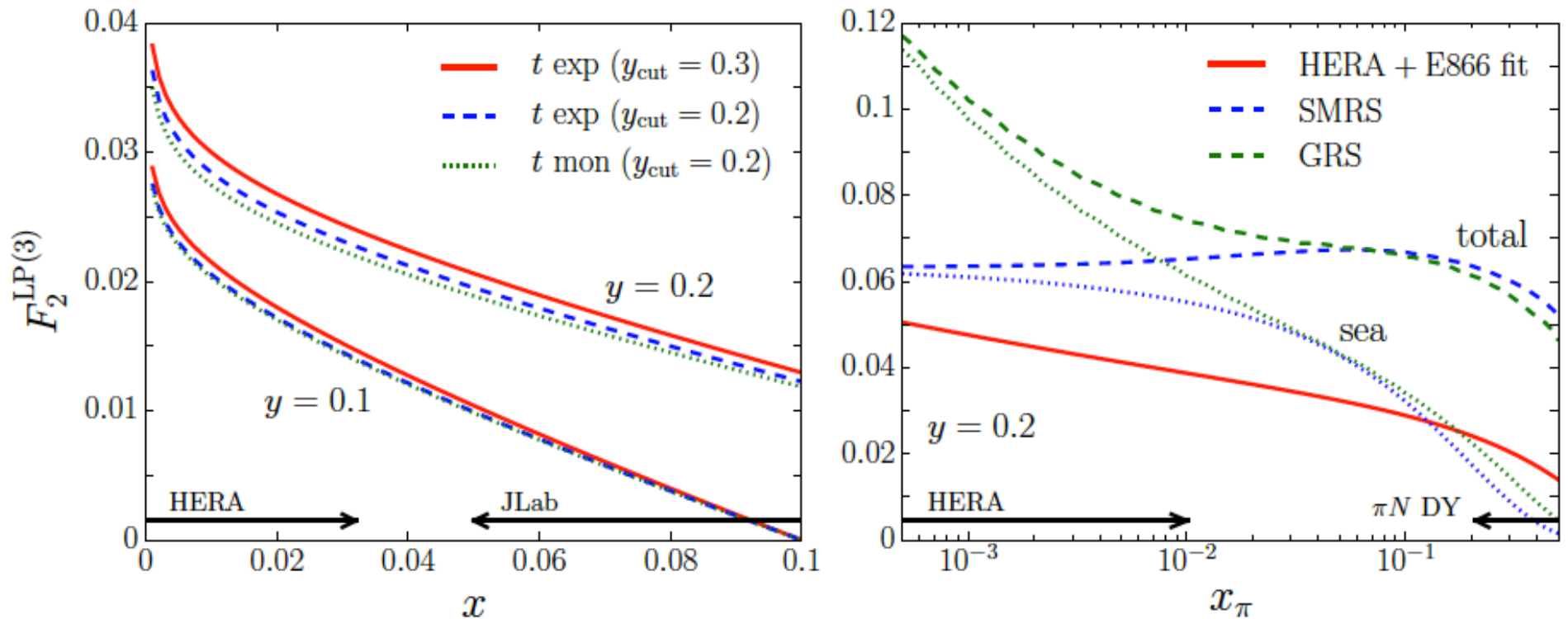


$$F_2^\pi = N x_\pi^a (1 - x_\pi)^b, \quad a = a_0 + a_1 \eta$$

$$\eta \sim \log(\log Q^2)$$

- stable values of F_2^π at $4 \times 10^{-4} \lesssim x_\pi \lesssim 0.03$ from combined fit
- shape similar to GRS fit to πN Drell-Yan data (for $x_\pi \gtrsim 0.2$), but smaller magnitude

Predictions at TDIS kinematics



→ JLab TDIS experiment can fill gap in x_π coverage between HERA and πN Drell-Yan kinematics

J. McKenney, N. Sato, W. Melnitchouk, C.Ji, PRD93, 054011 (2016)

Outlook

- Combined analysis can be extended by including also πN Drell-Yan data
→ constrain large- x_π region ($x_\pi \gtrsim 0.2$)
- Generalize parametrization by fitting individual pion valence and sea quark PDFs, rather than F_2^π
- Ultimate goal will be to use all data sensitive to pion structure (including TDIS, EIC?) to constrain pion PDFs over full range $10^{-4} \lesssim x_\pi \lesssim 1$